

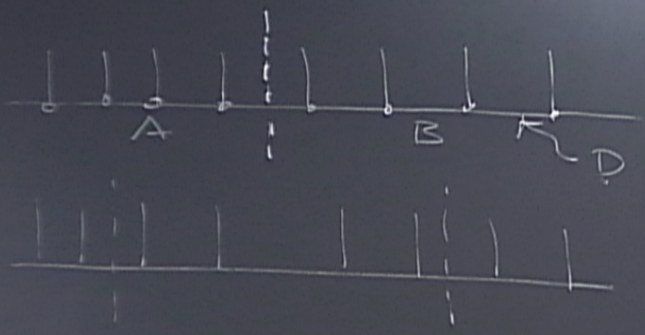
Title: Explorations in Condensed Matter-14

Date: Apr 02, 2015 10:15 AM

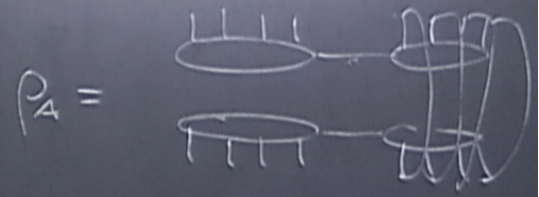
URL: <http://pirsa.org/15040003>

Abstract:

Entanglement - boundary law how MPS scale



$$S \leq \log D$$

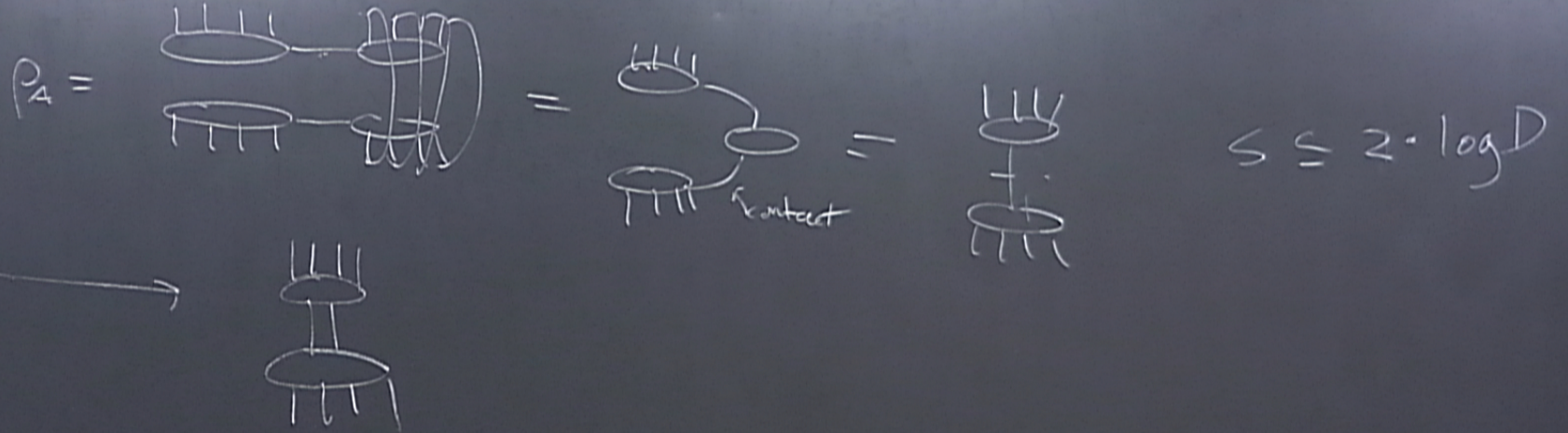


$$S \leq 2 \log D$$

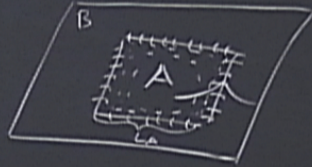


MPS scale

$\log D$



Crapped phases obey boundary law (we think)



$$S_A = a L_A + \dots$$

↑ non-universal ↑ universal numbers

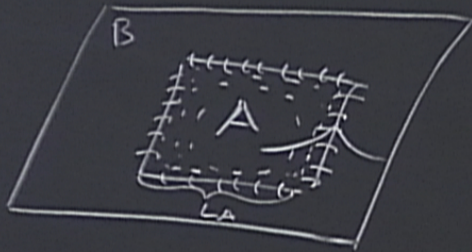
universal - independent of (most) microscopic details

Concrete example (valence bond solid)





Gapped phases obey boundary law (we think)



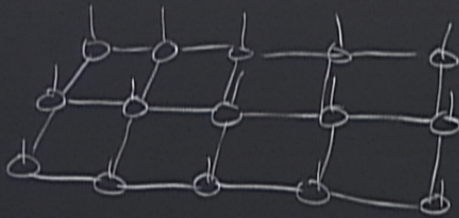
$$S_A = \underbrace{a L_A}_{\text{non-universal}} + \dots \underbrace{\dots}_{\text{universal numbers}}$$

universal = independent of (most) microscopic details

Concrete example (valence bond solid)



Research into extensions of MPS ideas \rightarrow Z^d
tensor product states / projected entangled pair states (PEPS)

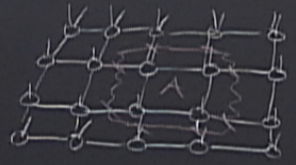


law (we think)

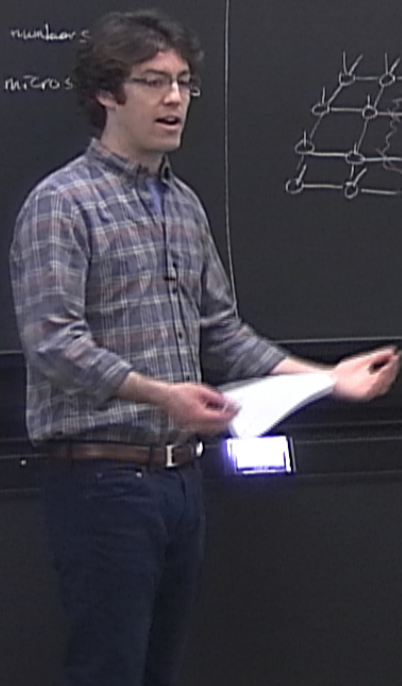
$L_A + \dots$
non-universal universal numbers
independent of (most) micros
solid)

Research into extensions of MPS ideas \rightarrow 2d
tensor product states / projected entangled pair states (PEPS)

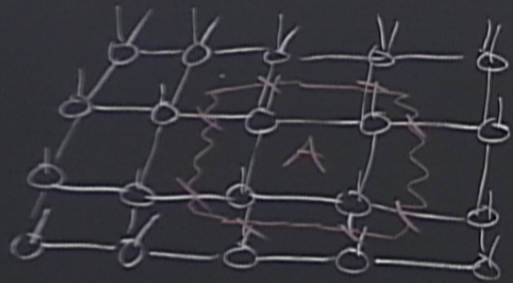
- can have both exponential and power-law decaying correlators (finite D)
- have boundary law entanglement



$$S \leq L_A \log D$$



Research into extensions of MPS ideas \rightarrow 2d
tensor product states / projected entangled pair states (PEPS)



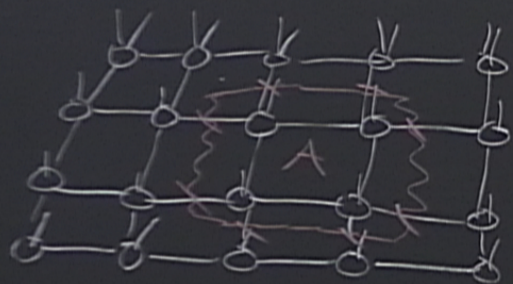
- can have both exponential and power-law decaying correlators (finite D)

- have boundary law entanglement

- not efficient to contract exactly

$$S \leq L_A \log D \quad (\text{recall MPS } D^3)$$

Research into extensions of MPS ideas $\rightarrow 2d$
tensor product states / projected entangled pair states (PEPS)



- can have both exponential and power-law decaying correlators (finite D)

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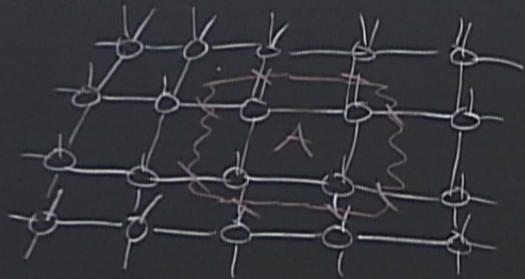
- not efficient to contract exactly

$$S \leq L_A \log D \quad (\text{recall MPS } \underline{\underline{D^3}})$$

Research into extensions of MPS ideas \rightarrow 2d
tensor product states / projected entangled pair states (PEPS)

o can have both exponential
and power-law decaying correlators (finite D)

o have boundary law entanglement
o not efficient to contract exactly



$$|S| \leq P_A \log D \quad (\text{recall MPS } \underline{\underline{D^3}})$$

P_A = perimeter of A



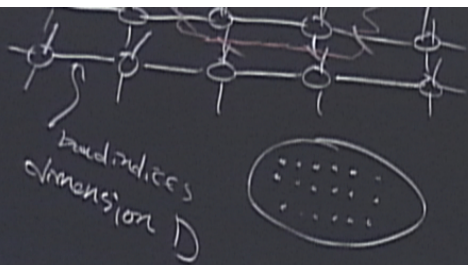
boundaries
dimension D

Is boundary law the end? No —

1d critical (gapless, disordered) systems violate boundary law

Boundary: $S_{\text{gapped, 1d}} \sim \text{const.} \sim a L^{d-1} \sim a L^0$

Critical (1d): $S_{\text{crit, 1d}} \sim \text{const.} \log L \rightarrow$ numerics (Vidal 2003)
field theory (Cardy 2004)



$$|S| \leq P_A \log D$$

(recall MPS D^3)

P_A = perimeter of A

$D \ll 10$

Thinking about units $\log L \rightarrow \log(L/\epsilon)$

Critical system $\xi \rightarrow \infty$, $\epsilon \sim$ lattice spacing, microscopic

$$\epsilon \rightarrow 2\epsilon = \epsilon'$$

$$S_{\text{crit}} = \kappa \log(L/\epsilon) + b$$

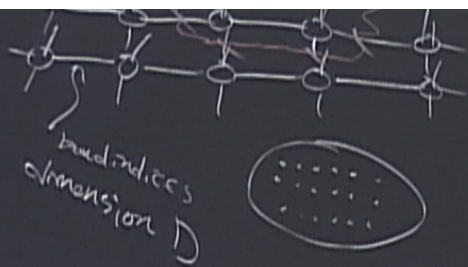
$$= \kappa \log\left(\frac{2L}{\epsilon'}\right) + b = \kappa \log(L/\epsilon') + \underbrace{[\kappa \log 2 + b]}_{b'}$$

$$= \kappa \log(L/\epsilon') + b'$$

ϵ boundary law

(Ida 2003)

(Cardy 2004)



$$|S| \leq P_A \log D$$

(recall MPS D^3)

P_A = perimeter of A

$$D \ll 10$$

Thinking about units $\log L \rightarrow \log(L/\epsilon)$

Critical system $\xi \rightarrow \infty$, $\epsilon \sim$ lattice spacing, microscopic

$$\epsilon \rightarrow 2\epsilon = \epsilon'$$

$$S_{\text{crit}} = k \log(L/\epsilon) + b$$

$$= k \log\left(\frac{2L}{\epsilon'}\right) + b = k \log(L/\epsilon') + \underbrace{[k \log 2 + b]}_{b'}$$

$$= \underbrace{k}_{\text{universal number}} \log(L/\epsilon') + b'$$

universal number

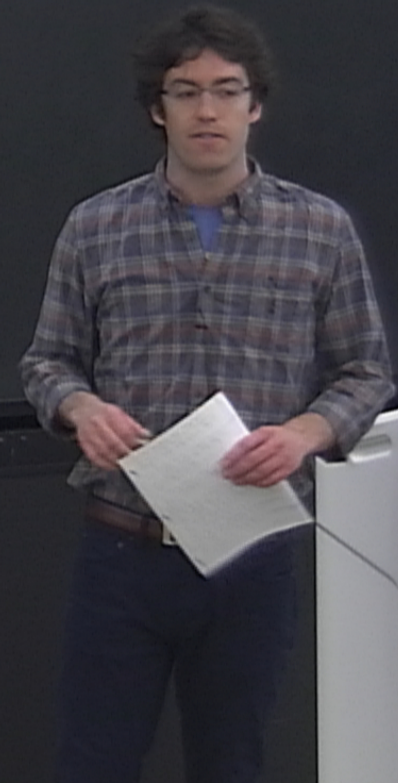
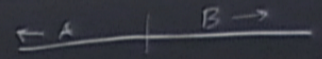
ϵ boundary law

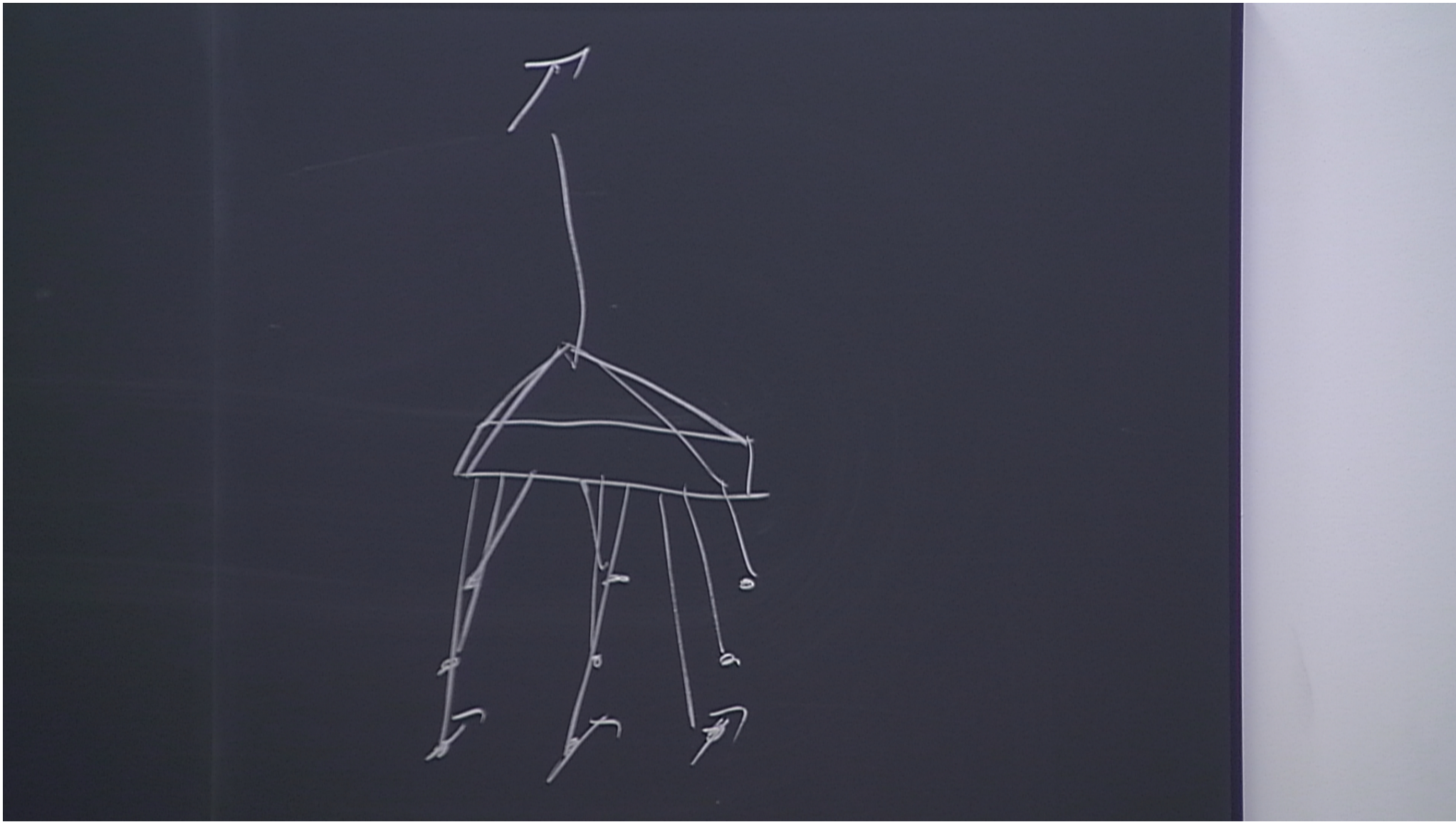
(Vidal 2003)

(Cardy 2004)

For 1d systems described by conformal field theory $K = \frac{c}{6}$
Similar to uncovering scaling exponents (ν, α, \dots)

c = central charge

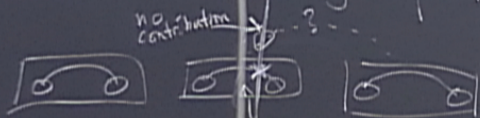




For 1d systems described by conformal field theory $\boxed{k = \frac{c}{6}}$ $c = \text{central charge}$ $\leftarrow A \quad | \quad B \rightarrow$
 Similar to uncovering scaling exponents (ν, α, \dots)

Why $\log L$ for critical 1d?

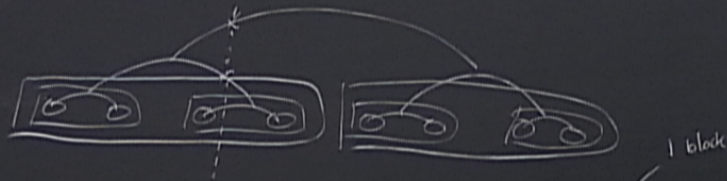
renormalization group thinking (block spin, real space RG)



contribution to entanglement at lattice scale

$\rightarrow S \leq \text{const}$ independent of global system size

To see $\log L$, assume scale invariance: do one step of RG, properly rescale, physics same



$$S = \underbrace{1+1+1+\dots}_{\log L} \approx \log L$$

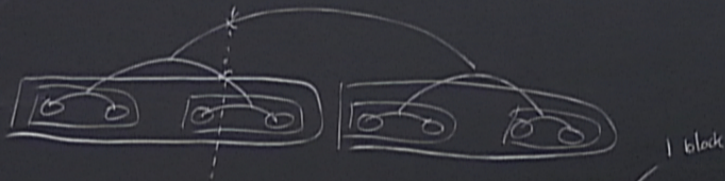
Finite, length L system

$$\frac{L}{2^n} = 1$$

(n = number of RG steps)

$$\Rightarrow \boxed{n = \log_2 L}$$

To see $\log L$, assume scale invariance: do one step of RG, properly rescale, physics same



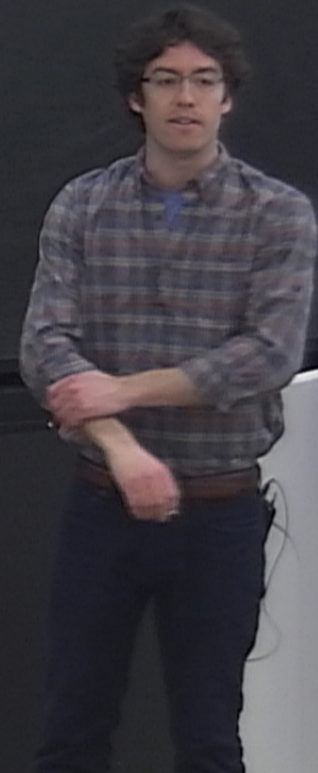
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Finite, length L system

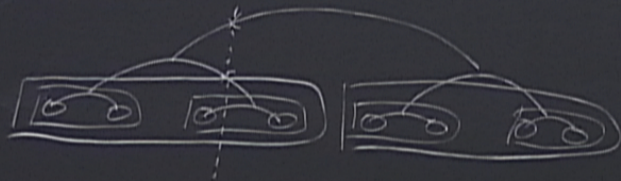
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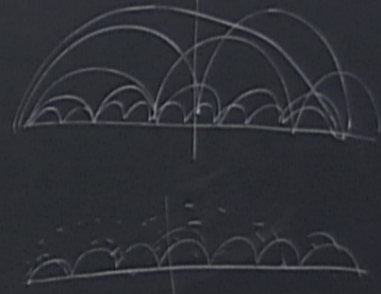


1 block

$S = 1 + 1 + 1 + \dots$ in $\log L$
 Finite, length h L system

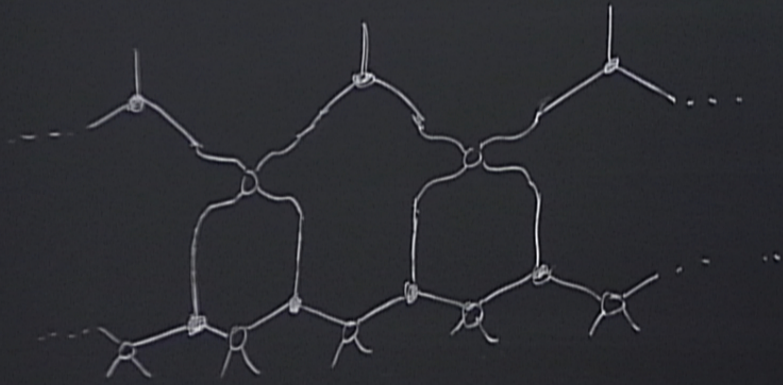
$$\frac{L}{2^n} = 1 \Rightarrow n = \log_2 L$$

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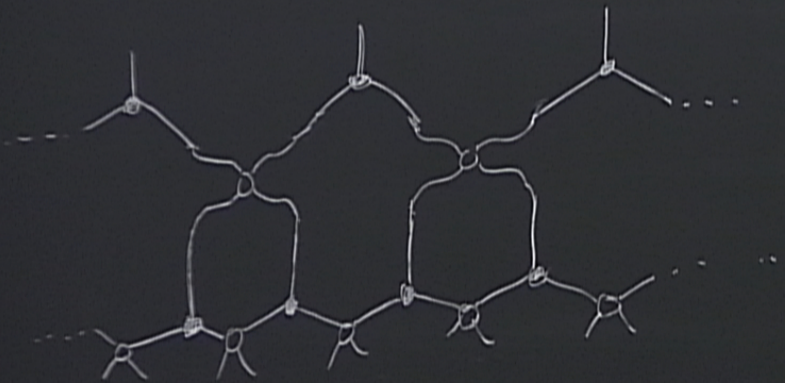
Tensor network wavefunction capturing critical $1d$ scaling?
Clue: invariant under rescaling, used by Vidal 2007



multi-scale
entanglement

(n = number of RG steps)

Tensor network wavefunction capturing critical $1d$ scaling?
Clue: invariant under rescaling, proposed by Vidal 2007



multi-scale
entanglement
renormalization
ansatz

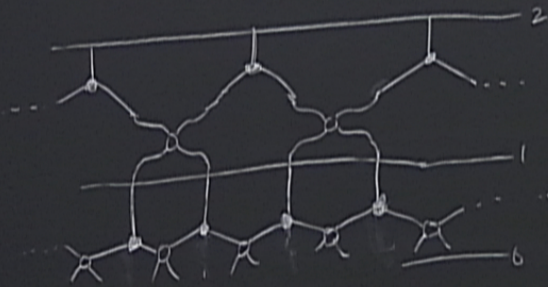
(MERA)

$S = \frac{1}{2} \log \frac{L}{h}$
Finite, length L system

$$\frac{L}{2^n} = 1 \Rightarrow n = \log_2 L$$

(n = number of RG steps)

Tensor network wavefunction capturing critical Id scaling?
Clue: invariant under rescaling, proposed by Vidal 2007



multi-scale
entanglement
renormalization
ansatz

(MERA)

Properties:

- layered structure
- applying each layer, half # of

tion capturing critical ld scaling?

rescaling, proposed by Vidal 2007

-2
multi-scale
entanglement
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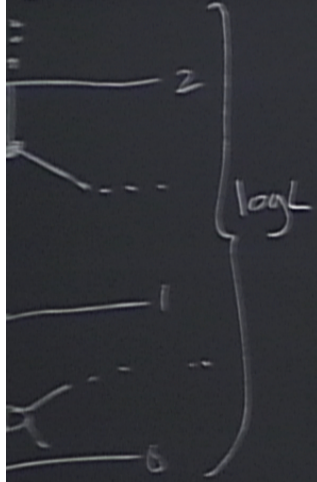
(MERA)

Properties:

- layered structure
- applying each layer, half # of sites
- system size L , $\log_2 L$ layers

function capturing critical Id scaling?

tensor rescaling, proposed by Vidal 2007

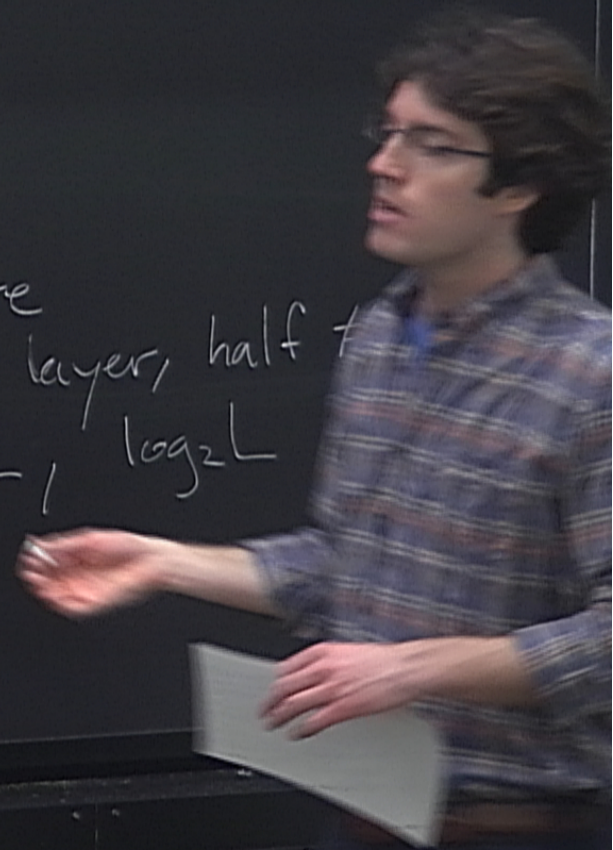


multi-scale
entanglement
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ansatz

(MERA)

Properties:

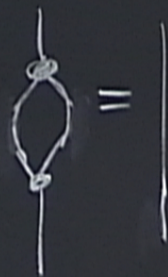
- layered structure
- applying each layer, half of system size
- system size L , $\log_2 L$



o two types of tensors



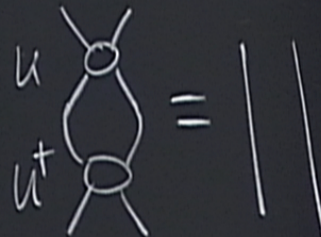
"isometry"



Isometric



"disentangler"



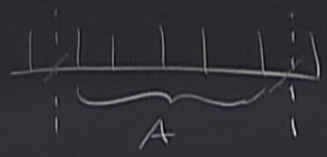
of sites
layers

and power-law decaying correlators (finite D)

o have boundary law entanglement
o not efficient to contract exactly

$$|S \leq P_A \log D \quad (\text{recall MPS } D^3)$$

P_A = perimeter of A $D \sim 10$

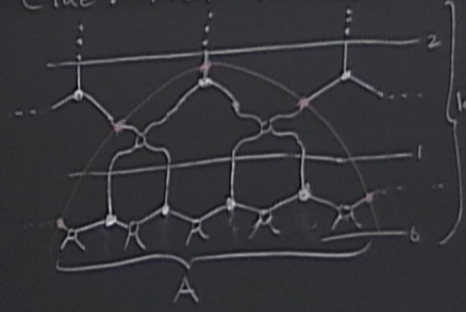


$$S \leq 2 \log D$$

$S = \frac{1}{\log 2} \log L$
 Finite, length L system $\frac{L}{2^n} = 1 \Rightarrow n = \log_2 L$
 (n = number of RG steps)

1 block

Tensor network wavefunction capturing critical Id scaling
Clue: invariant under rescaling, proposed by Vidal 2



multi-scale
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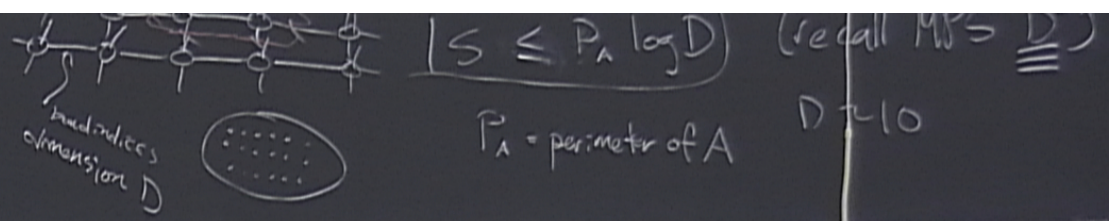
Properties:
o layered structure
o applying to system state

bonds cut (MERA) region A $\sim \log L_A$

$$S \leq (\log L_A) \underbrace{(\log D)}_{\text{const.}} \leq \text{const.} \log L_A$$

standard RG argument (Cardy)

$$\rightarrow \text{MERA } \langle \mathcal{O}_i \mathcal{O}_j \rangle_c \sim \frac{1}{|j-i|^p} \quad p > 0$$



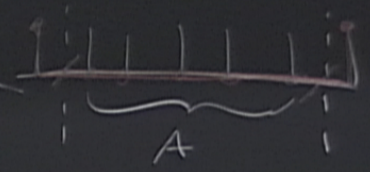
(MERA) region $A \sim \log L_A$ ($S \leq$)

$\underbrace{(\log D)}_{\text{const.}} \leq \text{const.} \log L_A$

argument (Cardy)

$(Q_i - Q_j)_c \sim \frac{1}{|j-i|^p}$ $p > 0$

	gapped	Critical	$d=1$ Fermi surface
1d	const	$\log L$	$\log L$
2d	L	L (dimensional Fermi surface)	$L \log L$



$S \leq 2 \log D$

$L \sim$ linear dimension of A