

Title: Explorations in Condensed Matter-13

Date: Apr 01, 2015 10:15 AM

URL: <http://pirsa.org/15040002>

Abstract:

→ MPS  $\text{---} \text{---} \text{---} \text{---} \text{---}$  → exponentially decaying (connected)  $\langle \hat{O}_i \hat{O}_j \rangle - \langle \hat{O}_i \rangle \langle \hat{O}_j \rangle \sim e^{-\xi |i-j|}$

→ MPS appropriate ground states gapped 1d systems

correlation length  $\xi = \frac{1}{\Delta}$

Theorem (Hastings 2007)

Local 1d Hamiltonian, finite energy gap, unique ground state  
 $\exists$  MPS bond dimension  $D$ ,  $p > 0$  (independent of  $D$ )  
such that  $|\langle \psi_{\text{MPS}} | \hat{O} | \psi_{\text{MPS}} \rangle - \langle \psi | \hat{O} | \psi \rangle| \leq C_D D^{-p}$

$\Delta$

nt

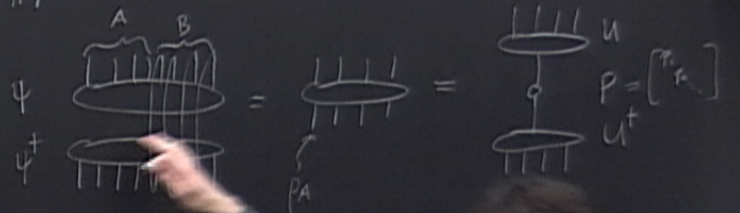
picture of why MPS are "right" form of g.s. of gapped 1d H's

amount/scaling of entanglement

: entanglement entropy (bipartite, von Neumann)  
state, divide regions A, B (A+B = whole system)

$$\text{Tr}[p_A \log p_A] = -\text{Tr}[p_B \log p_B]$$

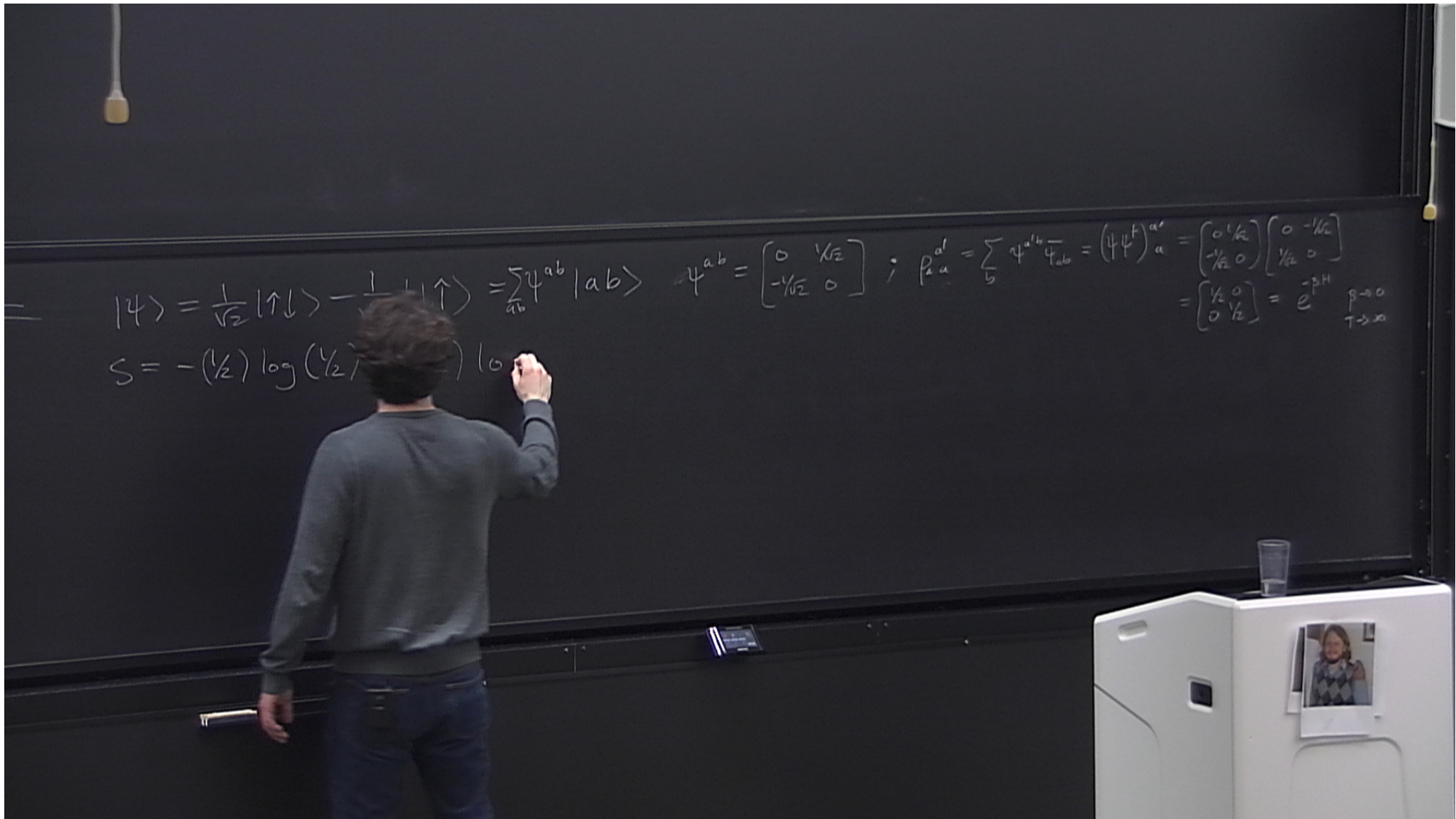
$$S_B[\rho] = \text{Tr}_B[|\psi\rangle\langle\psi|]$$



$$S = \sum_n -p_n \log p_n$$







$$|\psi\rangle = \frac{1}{\sqrt{2}}|\uparrow\downarrow\rangle - \frac{1}{\sqrt{2}}|\downarrow\uparrow\rangle = \sum_{ab} \psi^{ab} |ab\rangle \quad \psi^{ab} = \begin{bmatrix} 0 & 1/\sqrt{2} \\ -1/\sqrt{2} & 0 \end{bmatrix} ; \quad P_a^d = \sum_b \psi^{ab} \bar{\psi}_{ab} = (\psi \psi^\dagger)_a^a = \begin{bmatrix} 0 & 1/\sqrt{2} \\ -1/\sqrt{2} & 0 \end{bmatrix} \begin{bmatrix} 0 & -1/\sqrt{2} \\ 1/\sqrt{2} & 0 \end{bmatrix} \\ = \begin{bmatrix} 1/2 & 0 \\ 0 & 1/2 \end{bmatrix} = e^{-\beta H} \quad \beta \rightarrow 0 \\ T \rightarrow \infty$$
$$S = -(\frac{1}{2}) \log(\frac{1}{2}) - (\frac{1}{2}) \log(\frac{1}{2}) = \log 2$$



Ex 1

$$|\psi\rangle = \frac{1}{\sqrt{2}}|\uparrow\downarrow\rangle - \frac{1}{\sqrt{2}}|\downarrow\uparrow\rangle = \sum_{ab} \psi^{ab} |ab\rangle \quad \psi^{ab} = \begin{bmatrix} 0 & 1/\sqrt{2} \\ -1/\sqrt{2} & 0 \end{bmatrix}; \quad \rho_a^{ab} = \sum_b \psi^{ab} \overline{\psi^{ab}} = (\psi \psi^\dagger)_a^a = \begin{bmatrix} 0 & 0 \\ 0 & 1/2 \end{bmatrix}$$

$$S = -\left(\frac{1}{2}\right) \log\left(\frac{1}{2}\right) - \left(\frac{1}{2}\right) \log\left(\frac{1}{2}\right) \rightarrow \text{max entanglement } p \text{ equal eigenvalues}$$
$$= -\log\left(\frac{1}{2}\right) = \log 2$$

Ex 2

$$|\psi\rangle = \psi^{ab} |ab\rangle \quad \psi^{ab} = \frac{1}{2} \begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix} \quad |\psi\rangle = \frac{1}{2}(|\uparrow\uparrow\rangle + |\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle + |\downarrow\downarrow\rangle)$$

$$\rho_a^{ab} = (\psi \psi^\dagger)_a^a = \frac{1}{2} \begin{bmatrix} 1 & 1 \\ 1 & 1 \end{bmatrix} = \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ 0 & 0 \end{bmatrix} \begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \frac{1}{\sqrt{2}}$$
$$S = -(1) \log(1) = 0$$

→ MPS  $\bullet-\bullet-\bullet-\bullet-\bullet$  → exponentially decaying (connected)  $\langle \mathcal{O}_i \mathcal{O}_j \rangle \sim \langle \mathcal{O}_i \rangle \langle \mathcal{O}_j \rangle$

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Local 1d Hamiltonian, finite energy gap, unique ground state

$\exists$  MPS bond dimension  $D$ ,  $\mu > 0$  (independent of  $D$ )

such that  $|\langle \psi_{\text{MPS}} | \hat{\mathcal{O}} | \psi_{\text{MPS}} \rangle - \langle \psi | \hat{\mathcal{O}} | \psi \rangle| \leq C_{\mathcal{O}} D^{-\mu}$

$\hat{\mathcal{O}}$  local operator





$= -\log(1/2) = \log 2$   
Ex 2  $|\psi\rangle = \psi^{ab} |a,b\rangle$   $\psi^{ab} = \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix}$   $|\psi\rangle = \frac{1}{2} (|\uparrow\uparrow\rangle + |\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle + |\downarrow\downarrow\rangle)$   $D \begin{bmatrix} 1 & 1 & 1 & 1 \\ 1 & 1 & 1 & 1 \\ \dots & \dots & \dots & \dots \end{bmatrix}$   $|\psi\rangle = \left(\frac{1}{\sqrt{2}} |\uparrow\downarrow\rangle\right)^{\otimes 2}$   
 $\rho_a^{a'} = (\psi \psi^\dagger)_{a'}^a = \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix} \frac{1}{\sqrt{2}}$   $S = -(1) \log(1) = 0$   
 $= | \rightarrow \rangle | \rightarrow \rangle$

How much entanglement can an MPS carry?

"structured" density matrices

$\rho_a^{a'}$   $a = (a_1, a_2, a_3, \dots, a_{N_a})$   
 $a_j = 1, \dots, d$   
 $\hat{a} = 1, 2, \dots, d^{N_a}$   $N_a = \dots$



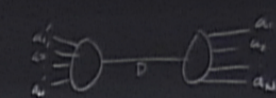


$= -\log(1/2) = \log 2$   
Ex 2  $|4\rangle = \psi^{ab} |ab\rangle$   $\psi^{ab} = \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix}$   $|4\rangle = \frac{1}{2} (|11\rangle + |1\bar{1}\rangle + |\bar{1}1\rangle + |\bar{1}\bar{1}\rangle)$   $D = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$   
 $\rho_{A|A} = (\psi\psi^\dagger)_{aa} = \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \frac{1}{\sqrt{2}}$   $S = -(1)\log(1) = 0$   
 $|4\rangle = \left(\frac{1}{\sqrt{2}} |1\bar{1}\bar{1}\bar{1}\rangle\right)^{\otimes 2} = |1\rangle \rightarrow |1\rangle \rightarrow$

How much entanglement can an MPS carry?

"structured" density matrices

$\rho_{A|A}$   
 $a = (a_1, a_2, a_3, \dots, a_{N_A})$   
 $a_j = 1, \dots, d$   
 $\hat{a} = 1, 2, \dots, \frac{d^{N_A}}{A}$   $(N_A = \# \text{ sites in } A)$

assume  $\rho_{A|A} = \sum_{i=1}^D t_i^{a_i} v_{i a_i} =$    
 $v_{i a} = (v_i)_a = i$  vectors each dimension  $A$   
 generally  $\text{span}\{\vec{v}_i\} = D$  dimensional space  
 Kernel/null space  $v_{i a}$   $(A-D)$  dimensional  $\Rightarrow v_{i a}$  rank  $D$   
 $\rho_A$  also has rank  $D$

$D < A$   
 $\sum_i v_{i a} w_i^a = 0$   
 $\uparrow$   $A-D$



$= -\log(1/2) = \log 2$   
Ex 2  $|\psi\rangle = \psi^{ab} |ab\rangle$   $\psi^{ab} = \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix}$   $|\psi\rangle = \frac{1}{2} (|\uparrow\uparrow\rangle + |\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle + |\downarrow\downarrow\rangle)$   $D \begin{pmatrix} 1/2 & 1/2 \\ 1/2 & 1/2 \end{pmatrix}$   
 $\rho_{a,a} = (\psi\psi^\dagger)_{a,a} = \frac{1}{2} \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix} = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & 0 \end{pmatrix} \begin{pmatrix} 1 & 1 \\ 1 & -1 \end{pmatrix} \frac{1}{\sqrt{2}}$   $S = -(1)\log(1) = 0$   
 $|\psi\rangle = \left(\frac{1}{\sqrt{2}} |\uparrow\downarrow\rangle\right)^{\otimes 2}$   
 $= |\rightarrow\rangle |\rightarrow\rangle$

How much entanglement can an MPS carry?  
 "structured" density matrices  $\rightarrow$  assume  $\rho_{A,a} = \sum_{i=1}^D t_i^{a'} v_{i,a}$

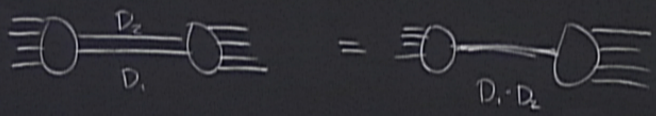
$\rho_{A,a} = (a_1, a_2, a_3, \dots, a_{N_A})$   
 $a_j = 1, \dots, d$   
 $\hat{a} = 1, 2, \dots, \frac{d^{N_A}}{A}$   $(N_A = \# \text{ sites in } A)$

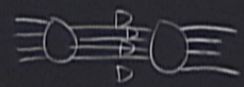
$v_{i,a} = (v_i)_a = i$  vectors each dimension  $A$   
 generally  $\text{span}\{\vec{v}_i\} = D$  dimensional space  
 Kernel/null space  $v_{i,a}$   $(A-D)$  dimensional  $\Rightarrow v_{i,a}$  rank  $D$   
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$D < A$   
 $\sum_i v_{i,a} w_i^a = 0$   $\uparrow$  basis of  $A-D$  of these



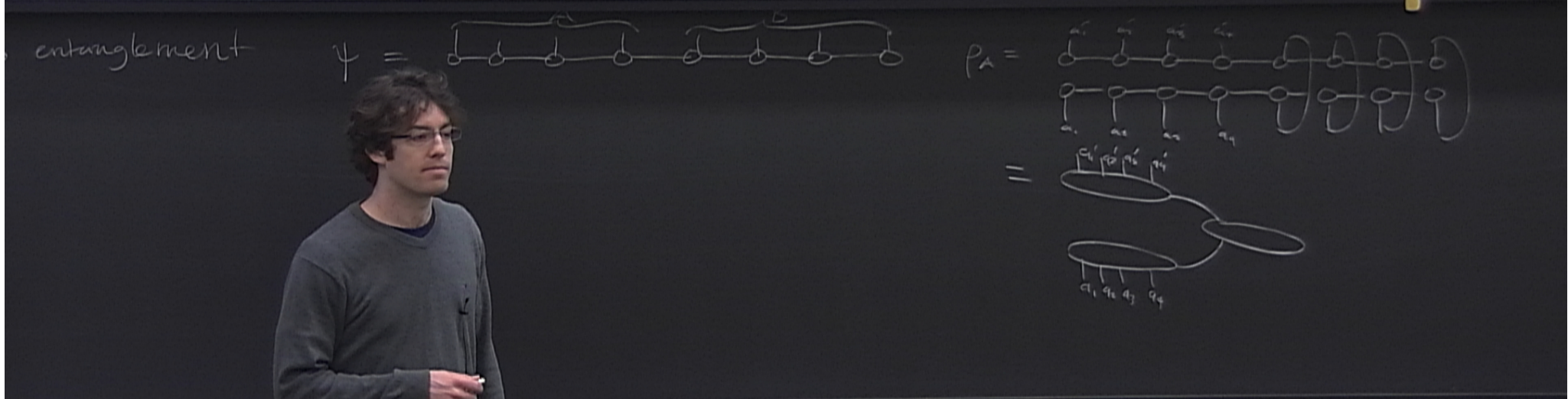
• most entangled case  $\rho = \frac{1}{D_1 D_2} \sum_{i,j} |i\rangle\langle i| \otimes |j\rangle\langle j|$  rank  $D_1 D_2$   $S[\rho] \leq \log D$

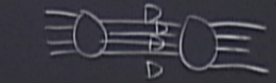
•   $S[\rho] \leq \log(D_1 \cdot D_2) = \log D_1 + \log D_2$


•   $S \leq N(\text{lines cut}) \cdot \log D$

•   $S \leq 2 \log D$

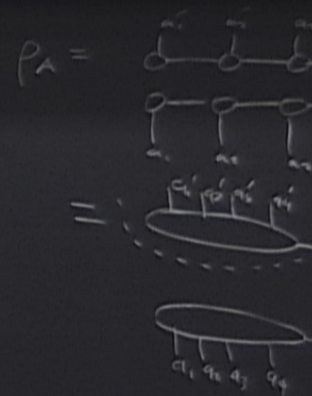
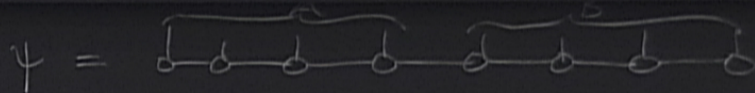





 $S \leq N_{\text{lines cut}} \cdot \log D$


 $S \leq 2 \log D$

MPS entanglement

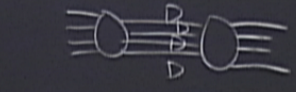


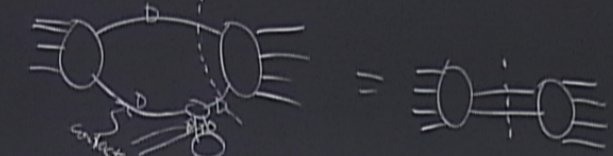
Theorem (Hastings 2007)

Consider 1d gapped Hamiltonian (unique g.s.)

Entanglement




 $S \leq N_{\text{lines cut}} \cdot \log D$


 $S \leq 2 \log D$

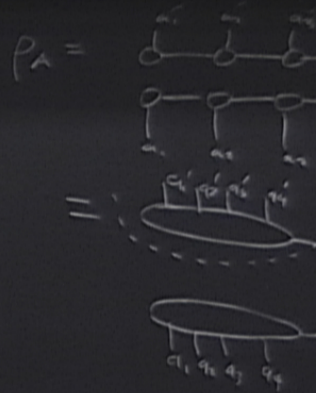
MPS entanglement  $\psi =$  

Theorem (Hastings 2007)

Consider 1d gapped Hamiltonian (unique g.s.)

for  $A | B$  bipartition

$S_A \leq \text{const.}$  independent of system size

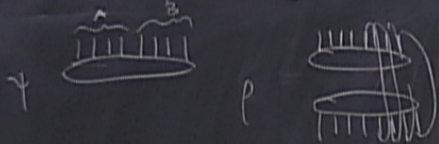


$S \leq \log D$  (indep of  $N$ ) "boundary law" scaling ("area law")





Maximally entangled  $\psi$



$\rho = D_A \times D_A$  matrix  $D_A = d^{N_A}$

$\rho$  full rank, all eigenvalues same

$$S_{\max} = \log D_A = \log(d^{N_A}) = \underline{\underline{N_A \log d}} \Rightarrow S_{\max} \sim \text{size (volume) of } A$$

"volume law"

random  $\psi$  will be volume law  
Hilbert space

