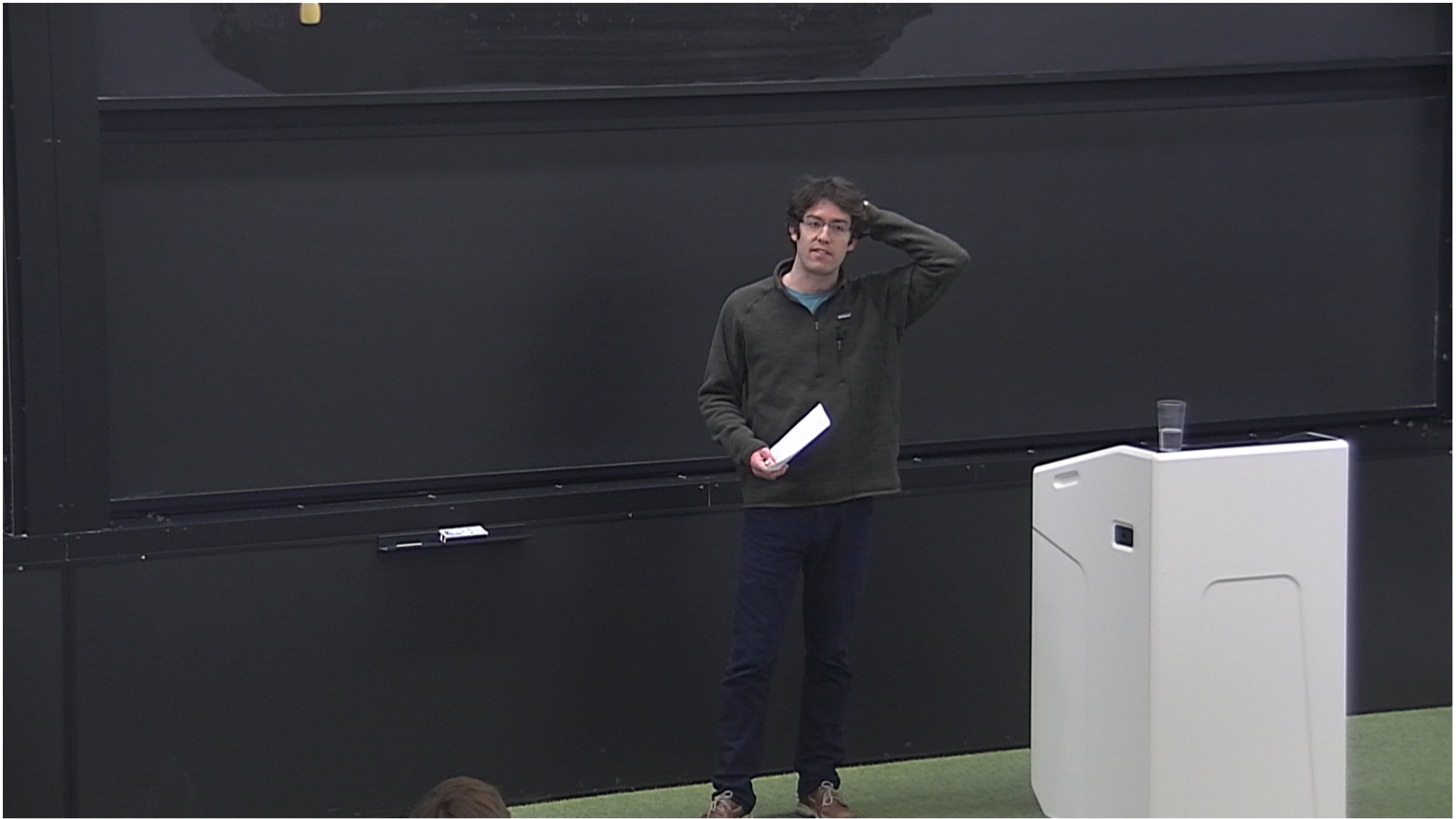


Title: Explorations in Condensed Matter-6

Date: Mar 23, 2015 10:15 AM

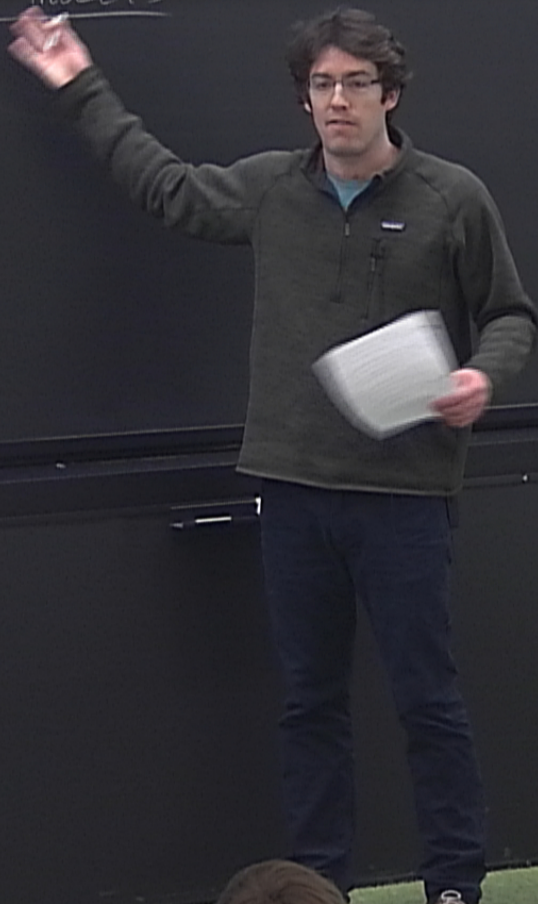
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Abstract:



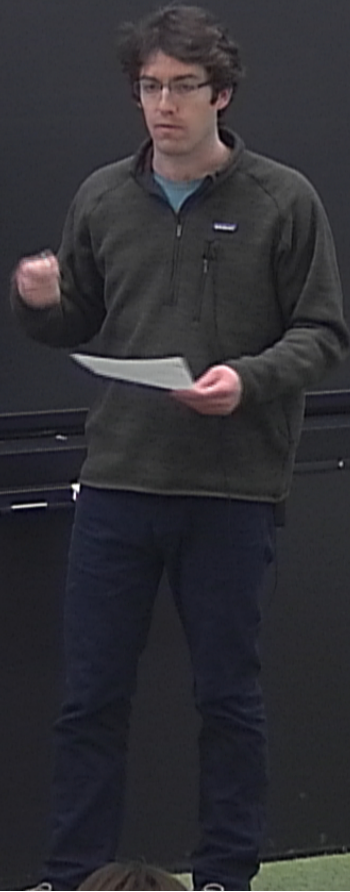


Lattice models





Lattice models





Lattice models

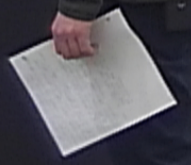
correlation (strong)  $\rightarrow$  interactions dominate kinetic + potential



Lattice models

correlation (strong)  $\rightarrow$  interactions dominate kinetic + potential

band structure  
 $\rightarrow$   
 $\rightarrow$





Lattice models

correlation (strong)  $\rightarrow$  interactions dominate kinetic

$\rightarrow$  band structure

$\rightarrow$  KS-DFT

build in interactions into bands



Lattice models

correlation (strong)  $\rightarrow$  mt is dominate kinetic + potential

$\hookrightarrow$  band structure

$\hookrightarrow$  KS-DFT

build in

into bands

tight-binding





Lattice models

correlation (strong)  $\rightarrow$  interactions dominate kinetic energy

$\hookrightarrow$  band structure

$\hookrightarrow$  KS-DFT

tight-binding



build in interactions into bands

$$t \sim \int \phi_i^*(r) \left[ -\frac{1}{2} \nabla^2 + U(\vec{r}) \right]$$



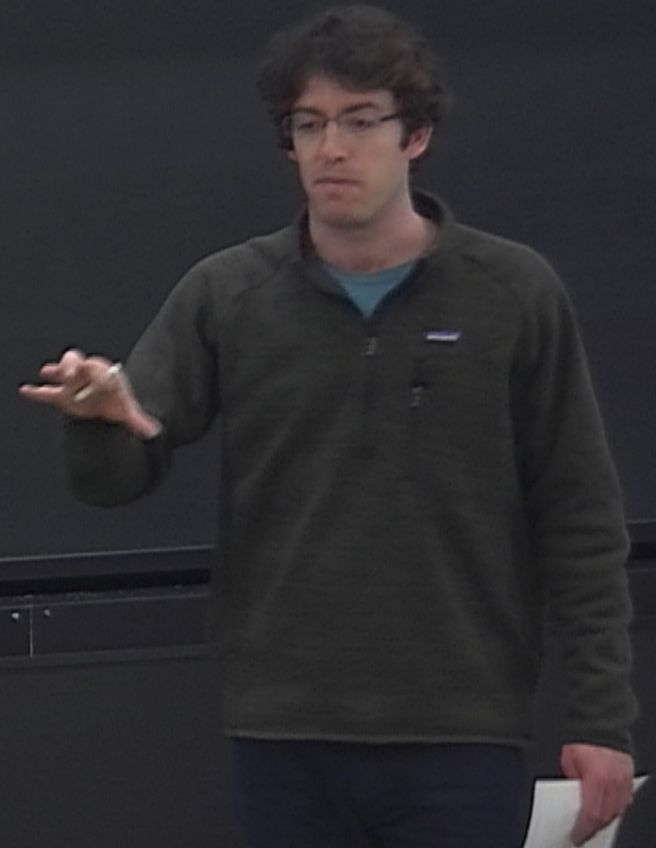
interactions dominate kinetic + potential

in interactions into bands  
 $T \sim \int \phi_1^*(r) \left[ -\frac{1}{2} \nabla^2 + U(\vec{r}) \right] \phi_2(r) \rightarrow$   
 $U \sim \int V_{\text{ext}}(r-r') |\phi_1(r)|^2 |\phi_1(r')|^2$



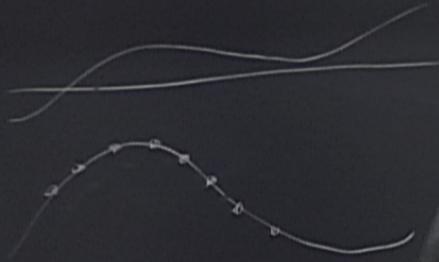


strong correlations  $\rightarrow$  d & f bands



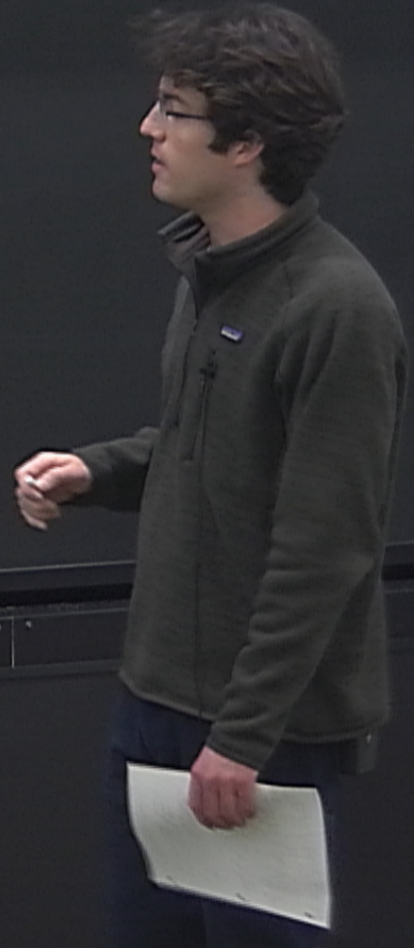
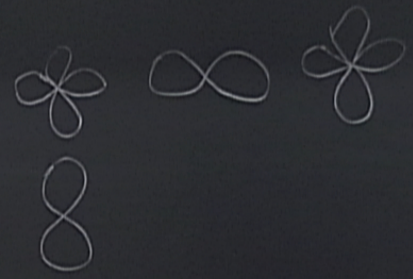
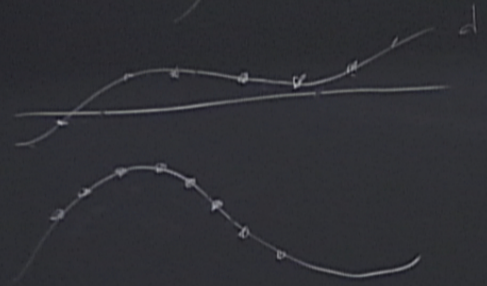


strong correlations  $\rightarrow$  d & f bands

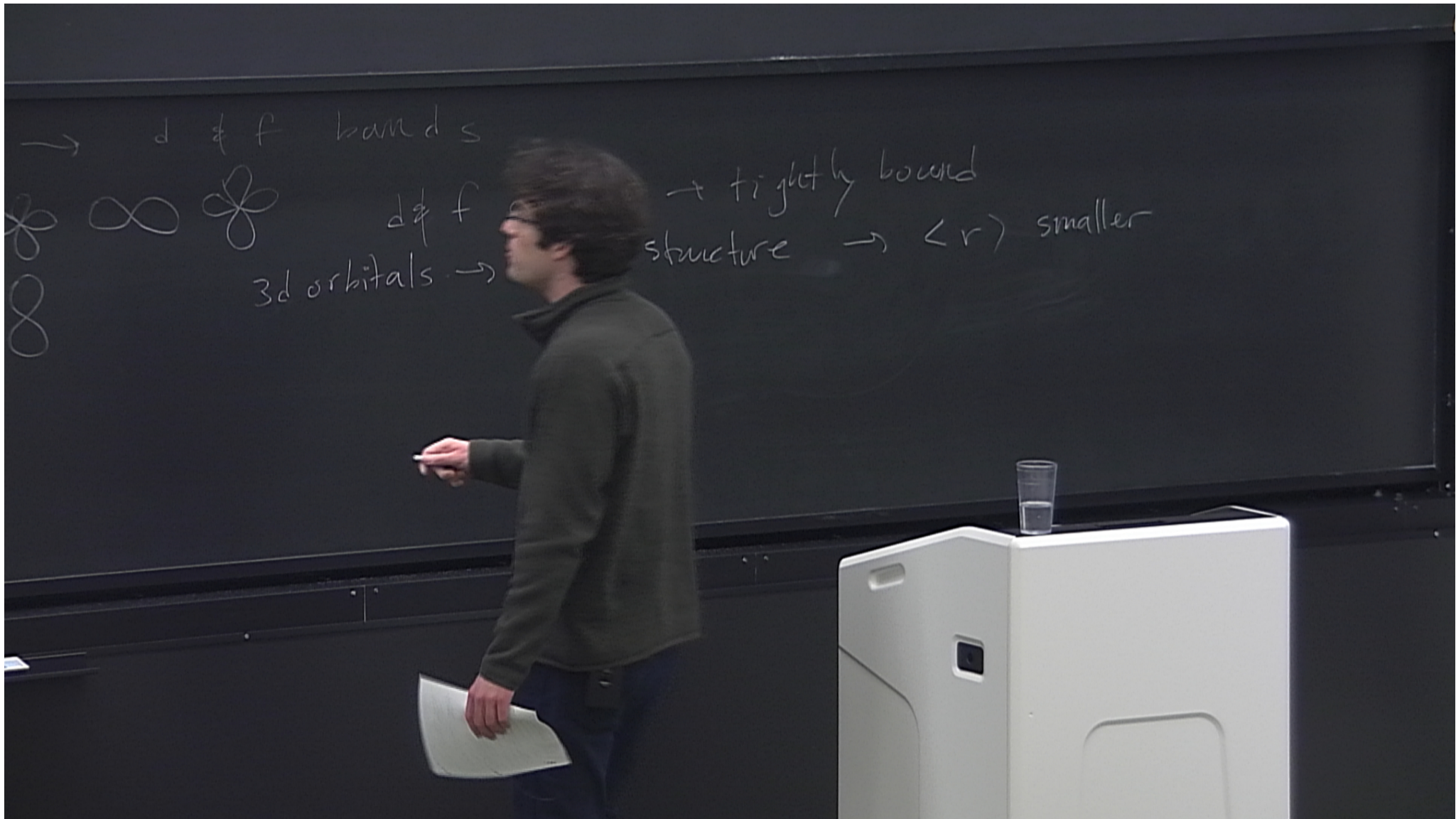




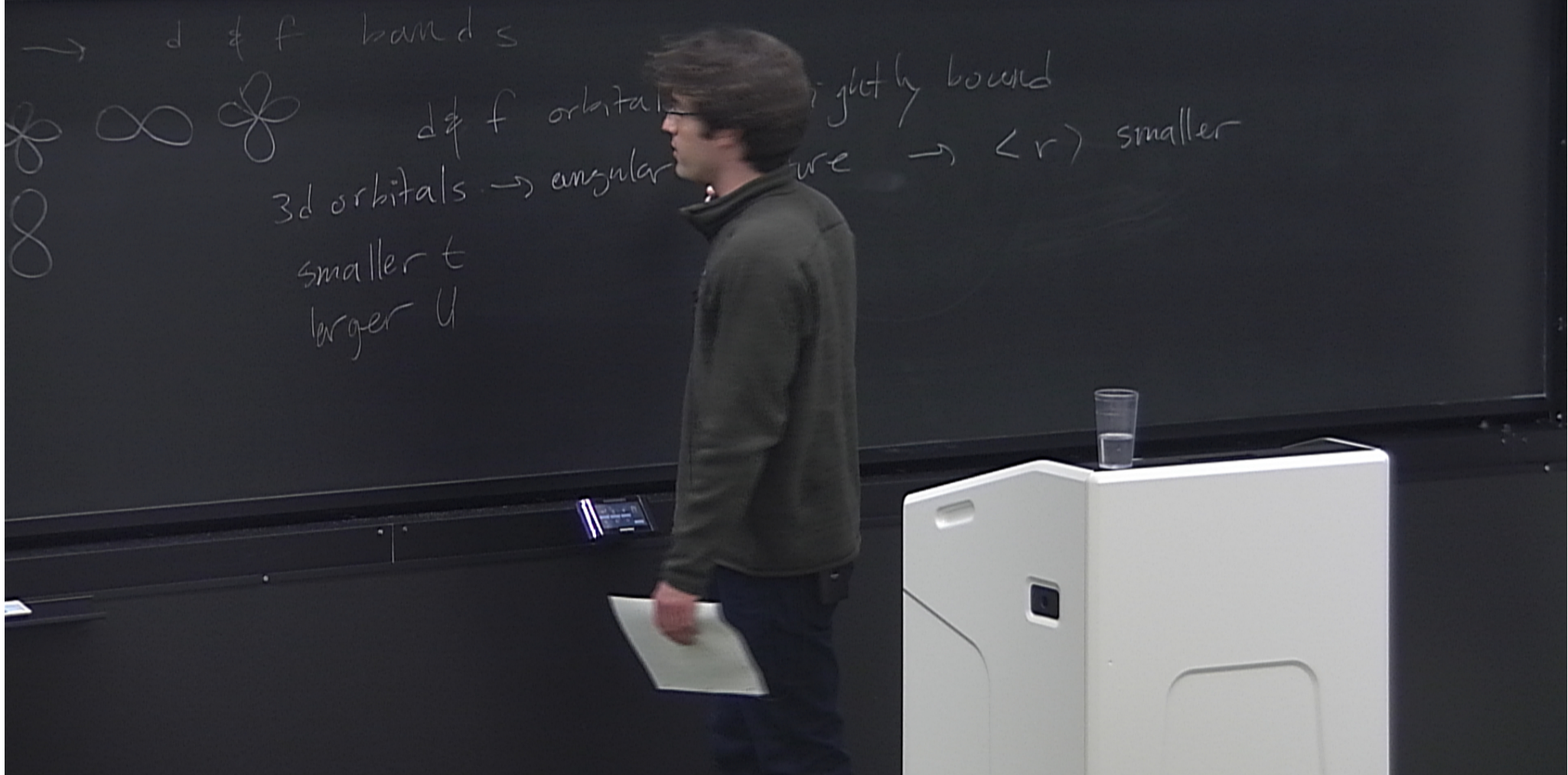
strong correlations  $\rightarrow$  d & f bands





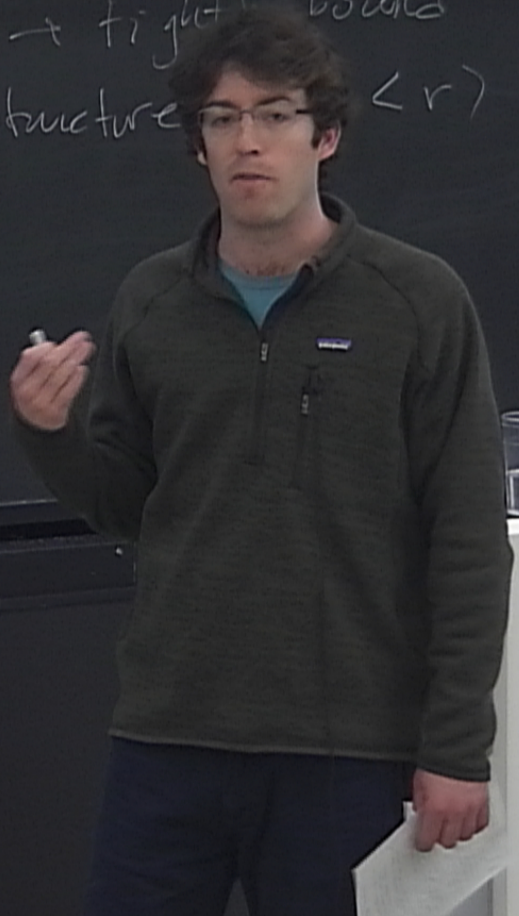
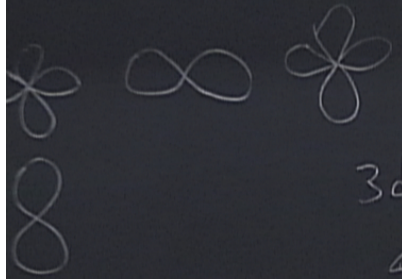




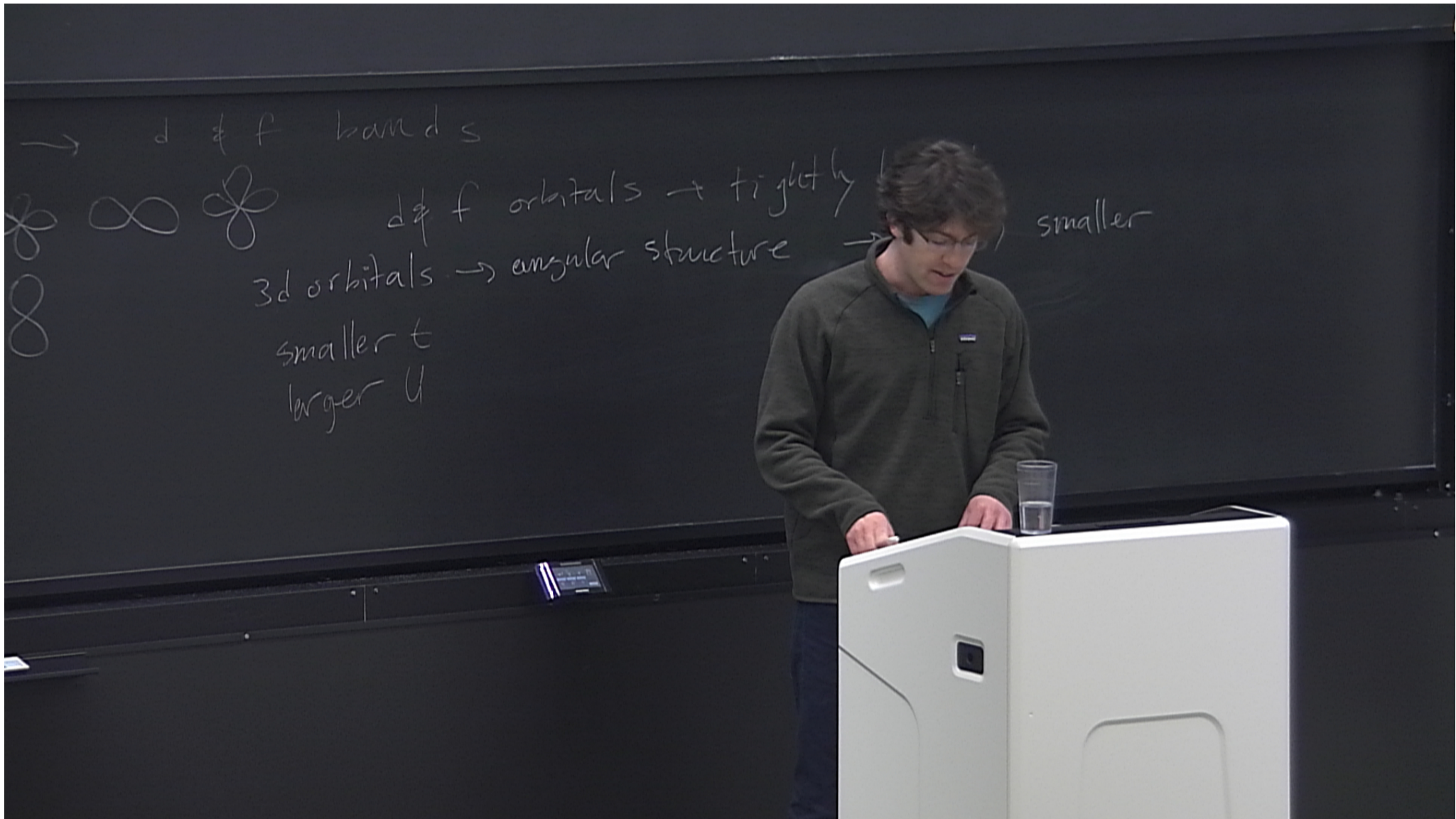




→ d & f bands  
d & f orbitals → tightly bound  
3d orbitals → angular structure <math>\langle r \rangle</math> smaller  
smaller  $t$   
larger  $U$





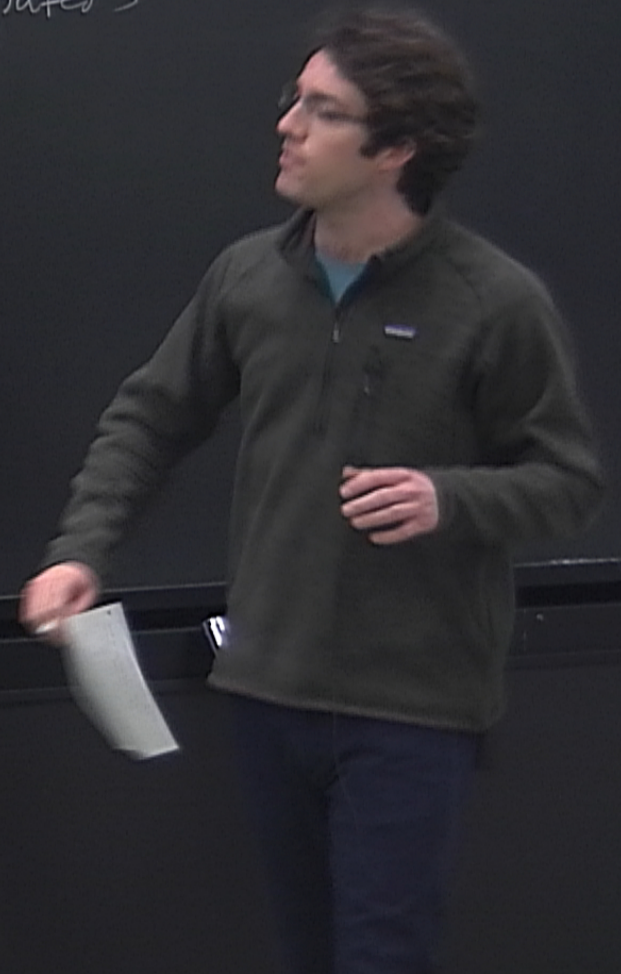




lattice models  $\begin{cases} \rightarrow \\ \rightarrow \end{cases}$  simple for

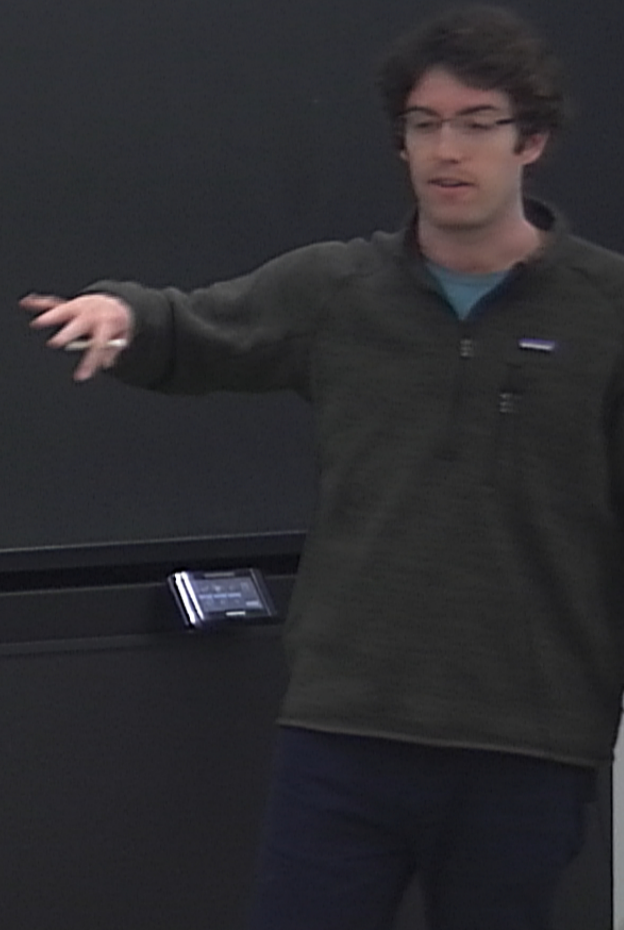


lattice models  $\begin{cases} \rightarrow \text{simple for people} \\ \rightarrow \text{simple for computers} \end{cases}$



lattice models  $\begin{cases} \rightarrow \text{simple for people } \checkmark \\ \rightarrow \text{simple for computers } \checkmark \end{cases}$

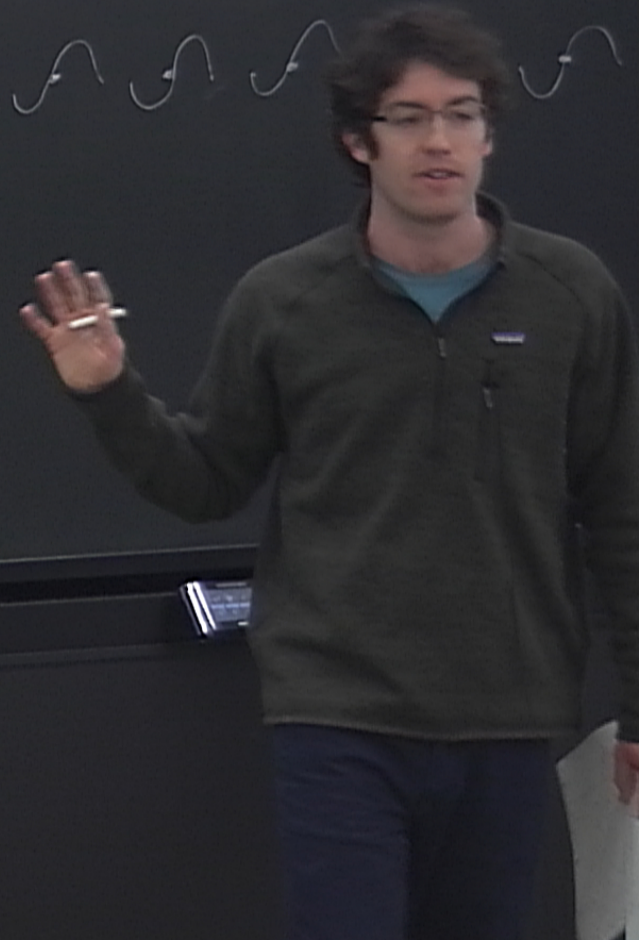
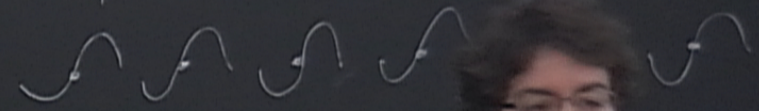
bands Bloch states  $\phi_{\alpha\vec{k}}(\vec{r})$





lattice models  $\begin{cases} \rightarrow \text{simple for people } \checkmark \\ \rightarrow \text{simple for computers } \checkmark \end{cases}$

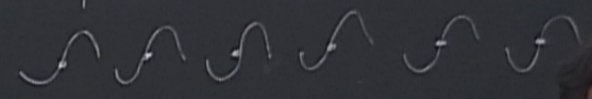
bands Bloch states  $\phi_{\alpha\mathbf{k}}(\vec{r})$





lattice models  $\begin{cases} \rightarrow \text{simple for people } \checkmark \\ \rightarrow \text{simple for computers } \checkmark \end{cases}$

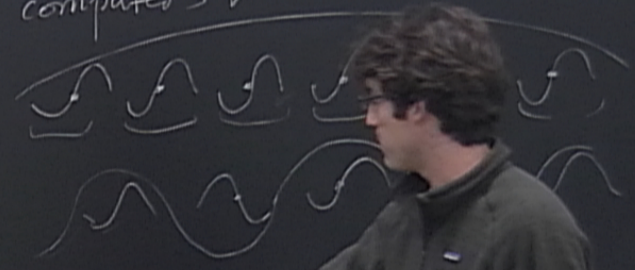
bands Bloch states  $\phi_{\vec{k}}(\vec{r})$





lattice models  $\begin{cases} \rightarrow \text{simple for people } \checkmark \\ \rightarrow \text{simple for computers } \checkmark \end{cases}$

bands Bloch states  $\phi_{\vec{k}}(\vec{r})$



$k \approx 0$

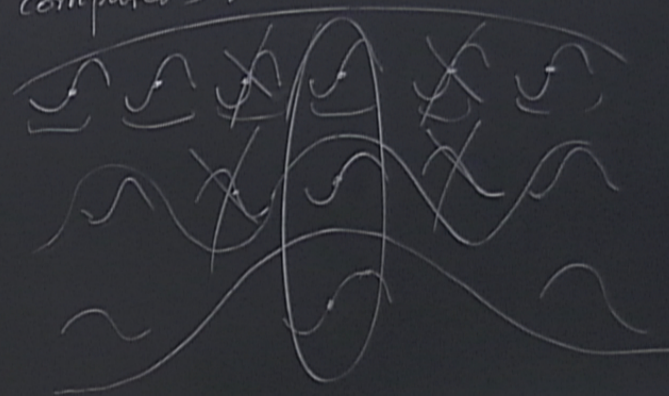
$k \approx \pi$



lattice models  $\begin{cases} \rightarrow \text{simple for people } \checkmark \\ \rightarrow \text{simple for computers } \checkmark \end{cases}$

bands Bloch states  $\psi_{\vec{k}}(\vec{r})$

$$\phi_{\alpha j}(\vec{r}) = \frac{1}{\sqrt{N}}$$



$k=0$

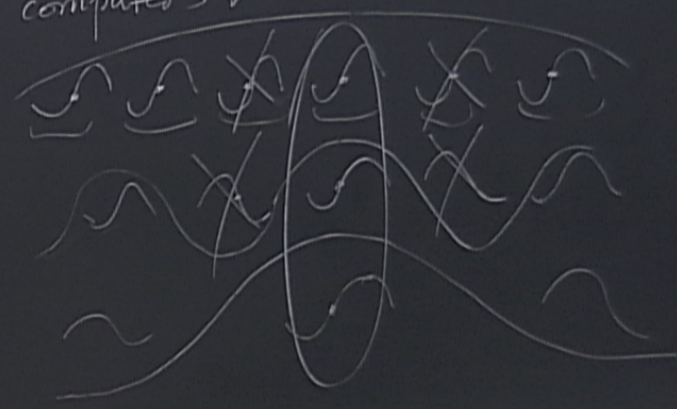
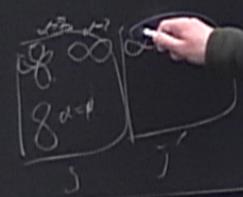
$k=\pi$



lattice models  $\begin{cases} \rightarrow \text{simple for people } \checkmark \\ \rightarrow \text{simple for computers } \checkmark \end{cases}$

bands Bloch  $\phi_{\alpha\vec{k}}(\vec{r})$

$$\phi_{\alpha j}(\vec{r}) = \frac{1}{\sqrt{N}} \sum_{\vec{k} \in \text{BZ}} e^{i\vec{k}\cdot\vec{r}} \phi_{\alpha j}(\vec{k})$$



$k \approx 0$

$k \approx \pi$



$$\hat{c}_{i\sigma}^+ = \int_{\vec{r}} \phi_i(\vec{r}) \hat{\psi}_{\vec{r}\sigma}^+$$

$(\alpha i) \rightarrow j$



$$\hat{c}_{i\sigma}^{\dagger} = \int_{\vec{r}} \phi_i(\vec{r}) \hat{\psi}_{\vec{r}\sigma}^{\dagger} \quad (\Rightarrow \hat{\psi}_{\vec{r}\sigma}^{\dagger} = \sum_i \phi_i(\vec{r}) \hat{c}_{i\sigma}^{\dagger})$$

$(\alpha i) \rightarrow j$  (orbitals)

$\hat{H}$

$$\hat{c}_{i\sigma}^\dagger \stackrel{\text{def}}{=} \int_{\vec{r}} \phi_i(\vec{r}) \hat{\psi}_{\vec{r}\sigma}^\dagger \quad (\Rightarrow \hat{\psi}_{\vec{r}\sigma}^\dagger = \sum_i \phi_i(\vec{r}) \hat{c}_{i\sigma}^\dagger = \sum_i \phi_i(\vec{r}) \phi_i^\dagger)$$

$(\alpha i) \rightarrow j$  (orbitals)

$$\hat{H} = \sum_{\sigma} \int_{\vec{r}} \hat{\psi}_{\vec{r}\sigma}^\dagger \left[ -\frac{1}{2} \nabla^2 + U(\vec{r}) \right] \hat{\psi}_{\vec{r}\sigma} + \frac{1}{2} \int_{\vec{r}, \vec{r}'} U_E(\vec{r} - \vec{r}')$$



$$\hat{C}_{i\sigma}^{\dagger} \stackrel{\text{def}}{=} \int_{\vec{r}} \phi_i(\vec{r}) \hat{\psi}_{\vec{r}\sigma}^{\dagger} \quad (\Rightarrow \hat{\psi}_{\vec{r}\sigma}^{\dagger} = \sum_i \phi_i(\vec{r}) \hat{C}_{i\sigma}^{\dagger} = \sum_i \int_{\vec{r}'} \phi_i(\vec{r}) \phi_i^*(\vec{r}') \hat{\psi}_{\vec{r}'\sigma}^{\dagger})$$

$(\propto i) \rightarrow j$  (orbitals)

$$\hat{H} = \sum_{\sigma} \int_{\vec{r}} \hat{\psi}_{\vec{r}\sigma}^{\dagger} \left[ -\frac{1}{2} \nabla^2 + U(\vec{r}) \right] \hat{\psi}_{\vec{r}\sigma} + \frac{1}{2} \int_{\vec{r}, \vec{r}'} U_E(\vec{r}, \vec{r}') \hat{\psi}_{\vec{r}\sigma}^{\dagger} \hat{\psi}_{\vec{r}'\sigma}$$

$$\hat{C}_{i\sigma}^{\dagger} \stackrel{\text{def}}{=} \int_{\vec{r}} \phi_i(\vec{r}) \hat{\psi}_{\vec{r}\sigma}^{\dagger} \quad \left( \Rightarrow \hat{\psi}_{\vec{r}\sigma} = \sum_i \phi_i(\vec{r}) \hat{C}_{i\sigma} = \sum_i \int_{\vec{r}'} \underbrace{\phi_i(\vec{r}) \phi_i^*(\vec{r}')}_{\delta(\vec{r}-\vec{r}')} \hat{\psi}_{\vec{r}'\sigma} = \hat{\psi}_{\vec{r}\sigma} \right)$$

( $\alpha i$ )  $\rightarrow$   $j$  (orbitals)

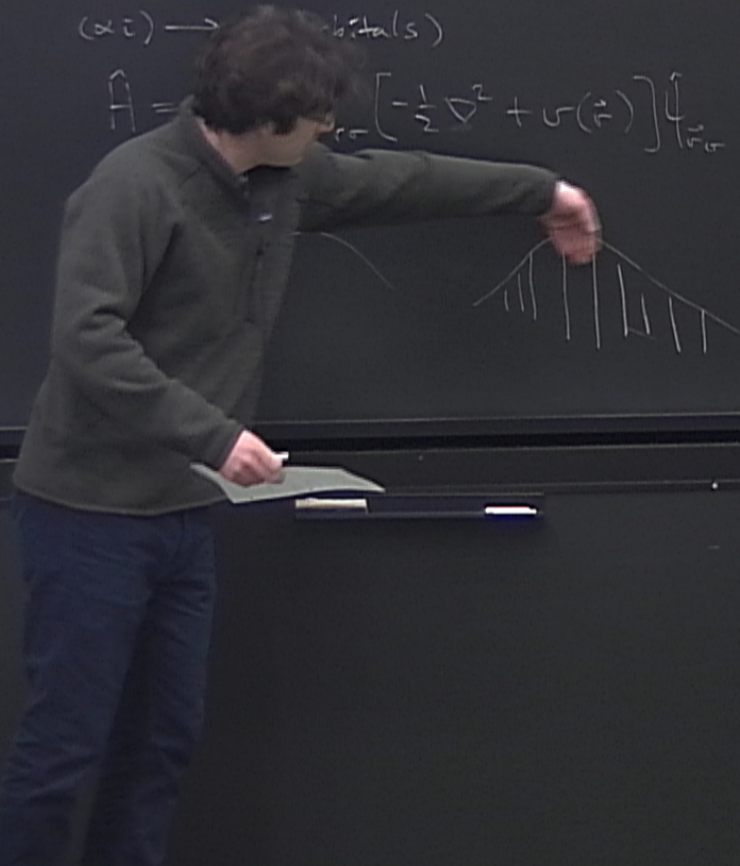
$$\hat{H} = \sum_{\sigma} \int_{\vec{r}} \hat{\psi}_{\vec{r}\sigma}^{\dagger} \left[ -\frac{1}{2} \nabla^2 + U(\vec{r}) \right] \hat{\psi}_{\vec{r}\sigma} + \frac{1}{2} \int_{\vec{r}, \vec{r}'} U_E(\vec{r}-\vec{r}')$$



$$\hat{c}_{i\sigma}^\dagger \stackrel{\text{def}}{=} \int_{\vec{r}} \phi_i(\vec{r}) \hat{\psi}_{\vec{r}\sigma}^\dagger \quad \left( \Rightarrow \hat{\psi}_{\vec{r}\sigma} = \sum_i \phi_i(\vec{r}) \hat{c}_{i\sigma} = \sum_{i'} \underbrace{\int_{\vec{r}'} \phi_i(\vec{r}) \phi_{i'}^*(\vec{r}')}_{\delta(\vec{r}-\vec{r}')} \hat{\psi}_{\vec{r}'\sigma} = \hat{\psi}_{\vec{r}\sigma} \right)$$

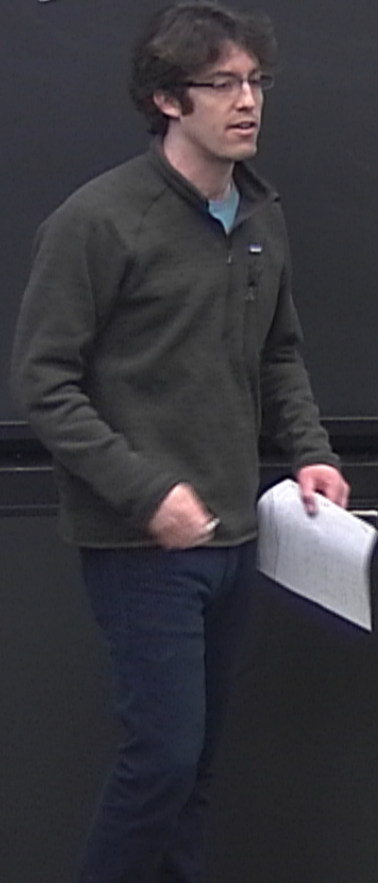
$(\propto i) \rightarrow$  orbitals

$$\hat{H} = \sum_{\vec{r}\sigma} \left[ -\frac{1}{2} \nabla^2 + U(\vec{r}) \right] \hat{\psi}_{\vec{r}\sigma} + \frac{1}{2} \int_{\vec{r}, \vec{r}'} U_E(\vec{r}-\vec{r}')$$



(tunneling  
hopping)

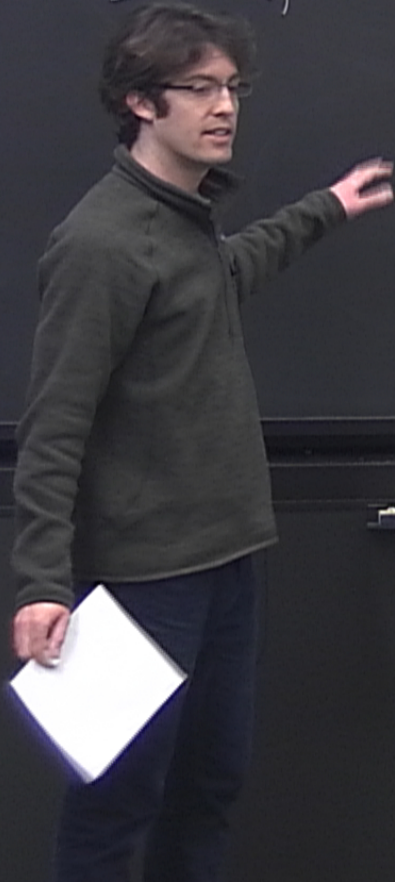
$$t_{ij} = -\langle \phi_i | \hat{H}_0 | \phi_j \rangle = -\int_{\mathbb{R}} \phi_i^*(\mathbb{R}) \left[ -\frac{1}{2} \nabla^2 + V(\mathbb{R}) \right] \phi_j(\mathbb{R})$$





(tunneling  
-  
hopping)

$$t_{ij} = -\langle \phi_i | \hat{H}_0 | \phi_j \rangle = -\int_{\mathbb{R}^d} \phi_i^*(\mathbf{r}) \left[ -\frac{1}{2} \nabla^2 + V(\mathbf{r}) \right] \phi_j(\mathbf{r})$$





(tunneling  
hopping)

$$t_{ij} = -\langle \phi_i | \hat{H}_0 | \phi_j \rangle = -\int_{\vec{r}} \phi_i^*(\vec{r}) \left[ -\frac{1}{2} \nabla^2 + v(\vec{r}) \right] \phi_j(\vec{r})$$

$$V_{ijkl} = \frac{1}{2} \int v_{ee}(\vec{r}-\vec{r}') \phi_i^*(\vec{r}) \phi_j(\vec{r}) \phi_k^*(\vec{r}') \phi_l(\vec{r}')$$

→ exact transformation / change of basis

→ hope that  $t_{ij}, V_{ijkl}$  short ranged



→ hope that  $t_{ij}, V_{ij}$  short ranged

What  $t$ 's and  $V$ 's are there?

- crystal f:  $t_{ij}$   $E_{\text{exp}} = t_{ij}(\beta, \epsilon)$



→ hope that  $t_{ij}, V_{ijke}$  short ranged

What  $t$ 's and  $V$ 's are there?

crystal field matrix  $E_{\alpha\beta} = t_{\alpha i}(\beta i)$

→ symmetries, group theory

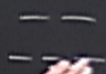


→ transformation / change of basis  
→ hope that  $t_{ij}, V_{ijke}$  short ranged

What  $t$ 's and  $V$ 's are there?

- crystal field matrix  $E_{\alpha\beta} = t_{(\alpha i)}$   
→ symmetries, group theory

$3d e^-$ , octahedral environment, phenomenological potential (x-tal field)



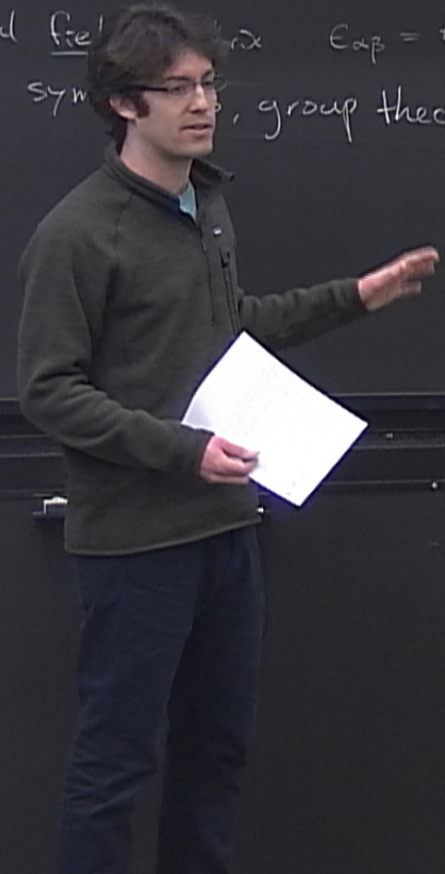


→ transformation / change of basis  
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What  $t$ 's and  $V$ 's are there?

- crystal field matrix  $E_{\alpha\beta} = t_{\alpha i}(\beta \epsilon)$   
→ sym, group theory

$3d e^-$ , octahedral environment, phenomenological potential (x-tal field)  
→  $3e_g$   
→  $3t_{2g}$





→ hope that  $t_{ij}, V_{ij}$  short ranged

What  $t$ 's and  $V$ 's are there?

- crystal field matrix  $E_{\alpha\beta} = t_{\alpha\beta}$   
→ symmetries, group theory

$3d e^-$ , octahedral environment, phenomenological potential (crystal field)

→  $-\text{---} 3g$   
 $-\text{---} 3t_{2g}$  → Jahn-Teller distortions

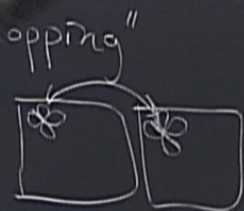


→ hope that  $t_{ij}$ ,  $V_{ij}$  short ranged

What  $t$ 's and  $V$ 's are there?

- crystal field matrix  $E_{sp} = t_{ij}(p_i)$   
 → symmetry theory

- next-neighbour hopping  
 $t \stackrel{\text{def}}{=} t$



$3d^5$ , octahedral environment, phenomenological parameters

→  $3g$   
 $3t_{2g}$

Jahn-Teller distortions

An energy level diagram showing the splitting of five degenerate d-orbitals into a higher energy  $3g$  level and a lower energy  $3t_{2g}$  level. The  $3t_{2g}$  level is further split into two levels due to Jahn-Teller distortion, with the upper level circled.



• Most important interaction

$$\frac{U}{Z} \stackrel{\text{def}}{=} V_{\text{coul}} = \frac{1}{2} \int_{\vec{r}, \vec{r}'} V_{\text{ec}}(\vec{r} - \vec{r}') |\phi_c(\vec{r})|^2 |\phi_c(\vec{r}')|^2 \quad (\sim \text{cl})$$



• Most important interaction

$$\frac{U}{2} \stackrel{\text{def}}{=} V_{cccc} = \frac{1}{2} \int_{\vec{r}, \vec{r}'} v_{cc}(\vec{r}-\vec{r}') |\phi_c(\vec{r})|^2 |\phi_c(\vec{r}')|^2 \quad \left( \sim \overbrace{c_{cc}^{\dagger} c_{cc}^{\dagger} c_{cc} c_{cc}} \right)$$

further-neighbor  $V$ 's

$$V_{ccjj} = \frac{1}{2} \int v_{cc}(\vec{r}-\vec{r}') |\phi_c(\vec{r})|^2 |\phi_j(\vec{r}')|^2$$

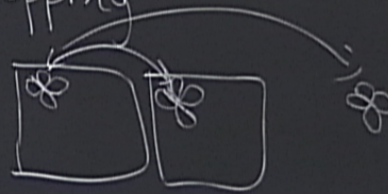


→ hope that  $t_{ij}, V_{ij}$  short ranged

What  $t$ 's and  $V$ 's are there?

- crystal field matrix  $E_{\alpha\beta} = t_{\alpha\beta}(\beta\bar{\alpha})$   
 → symmetries, group theory

- next-neighbor "hopping"  
 $t \stackrel{\text{def}}{=} t_{(i\bar{i})(k\bar{k})}$   
 further-neighbor  
 $t$



$3d^5$ , octahedral environment phenomena  
 →  $3g$   
 $3t_{2g}$



## "Exchange" mechanisms

- Coulomb exchange → short range, ferromagnetic  $\leftrightarrow$  int.
- Kinetic exchange → neighboring sites, antiferromagnetic spin



## "Exchange" mechanisms

◦ Coulomb exchange → short range, ferromagnetic spin int.

◦ Kinetic exchange → neighboring sites, magnet spin

$$\sum_{ij} V_{ij} n_i n_j$$

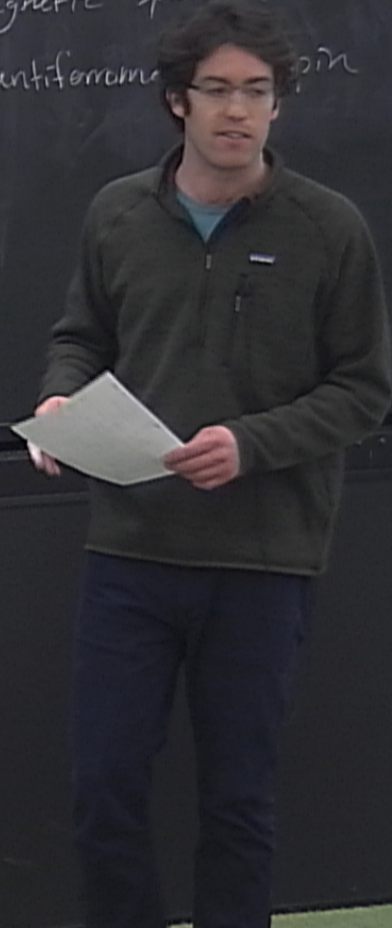


"Exchange" mechanisms (spin interactions)

◦ Coulomb exchange → short range, ferromagnetic spin int.

◦ Kinetic exchange → neighboring sites, antiferromagnetic spin

$$\sum_{\sigma\sigma'} V_{ijj'j''} c_{i\sigma}^\dagger c_{j\sigma} c_{j'\sigma'}^\dagger c_{j''\sigma'}$$





"Exchange" mechanisms (spin interactions)

◦ Coulomb exchange → short range, ferromagnetic spin int.

◦ Kinetic exchange → neighboring sites, antiferromagnet spin

$$\sum_{\sigma\sigma'} V_{ijji} \underbrace{c_{i\sigma}^\dagger c_{j\sigma} c_{j\sigma'}^\dagger c_{i\sigma'}}_{\text{}} = -V_{ijji} \sum_{\sigma\sigma'} (c_{i\sigma}^\dagger c_{i\sigma'}) (c_{j\sigma'}^\dagger c_{j\sigma}) + (\sim \hat{S}_i \cdot \hat{S}_j)$$



# "Exchange" mechanisms (spin interactions)

o Coulomb exchange → short range, ferromagnetic spin int.

o Kinetic exchange → neighboring sites, ferromagnetic spin

$$\sum_{\sigma\sigma'} V_{ijji} c_{i\sigma}^\dagger c_{j\sigma} c_{j\sigma'}^\dagger c_{i\sigma'} = -V_{ijji} \sum_{\sigma\sigma'} (c_{j\sigma'}^\dagger c_{j\sigma}) + (\sim \hat{n})$$

$$S_i^z \stackrel{\text{def}}{=} \frac{1}{2} \sum_{\alpha\beta} \hat{\sigma}_{\alpha\beta} c_{i\alpha}^\dagger c_{i\beta}$$