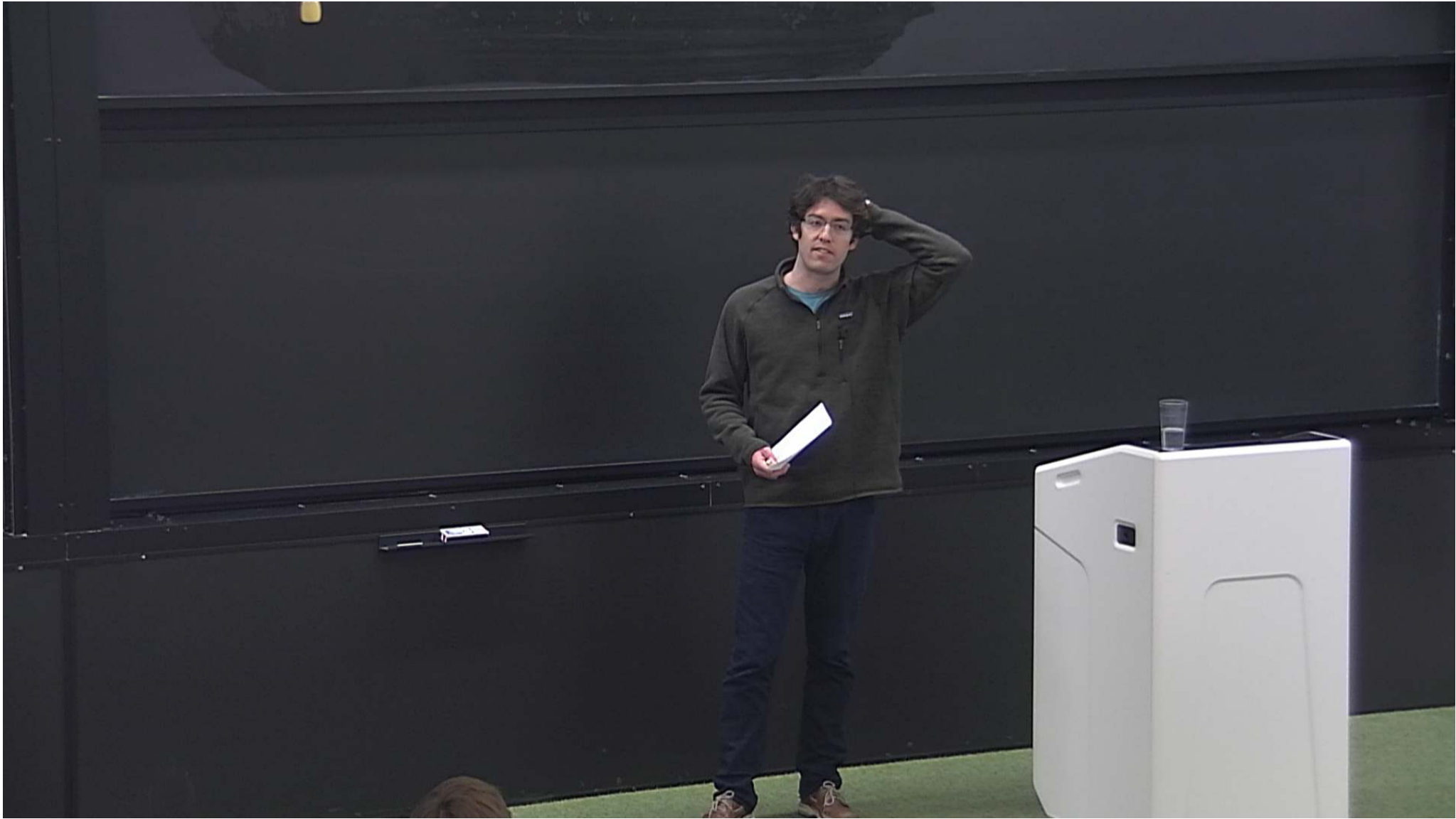


Title: Explorations in Condensed Matter-6

Date: Mar 23, 2015 10:15 AM

URL: <http://pirsa.org/15030041>

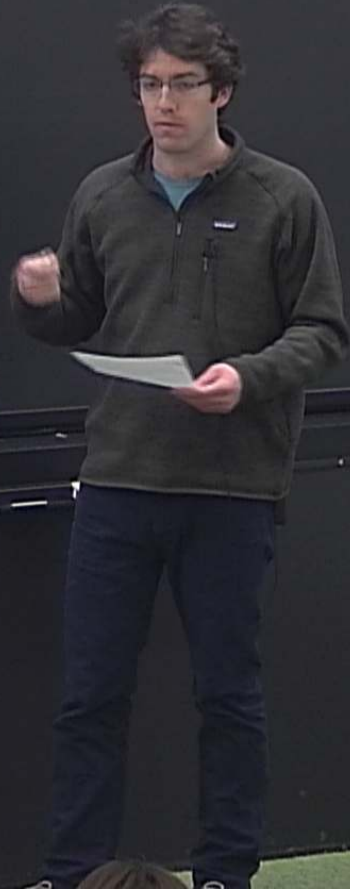
Abstract:



Lattice models



Lattice models

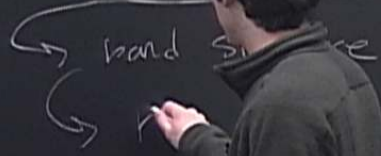


Lattice models

correlation (strong)  $\rightarrow$  interactions dominate kinetic + potential

Lattice models

correlation (strong)  $\rightarrow$  interactions dominate kinetic + potential



Lattice models

correlation (strong)  $\rightarrow$  interactions dominate kinetic

$\rightarrow$  band structure

$\rightarrow$  KS-DFT

build in interactions into bands

Lattice models

correlation (strong)  $\rightarrow$  mt is dominate kinetic + potential

$\hookrightarrow$  band structure

$\hookrightarrow$  KS-DFT

build in

into bands

tight-binding



Lattice models

correlation (strong)  $\rightarrow$  interactions dominate kinetic energy

$\hookrightarrow$  band structure

$\hookrightarrow$  KS-DFT

tight-binding



build in interactions into bands

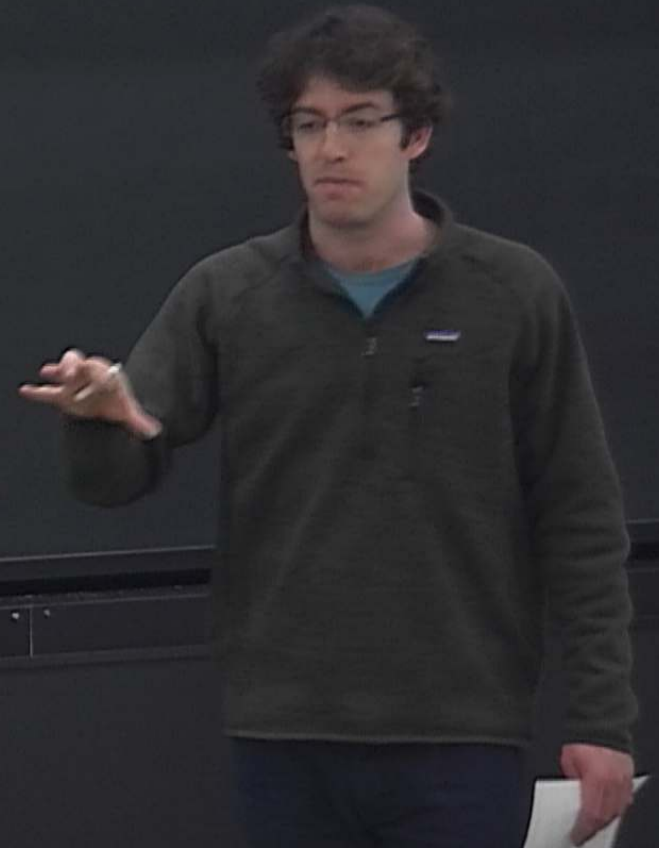
$$t \sim \int \phi_i^*(r) \left[ -\frac{1}{2} \nabla^2 + U(\vec{r}) \right]$$

interactions dominate kinetic + potential

in interactions into bands  
 $T \sim \int \phi_1^*(r) \left[ -\frac{1}{2} \nabla^2 + U(\vec{r}) \right] \phi_2(r) \rightarrow$   
 $U \sim \int V_{\text{ext}}(r-r') |\phi_1(r)|^2 |\phi_1(r')|^2$



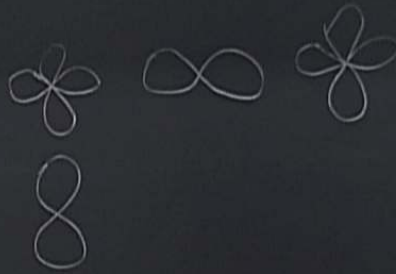
strong correlations  $\rightarrow$  d & f bands

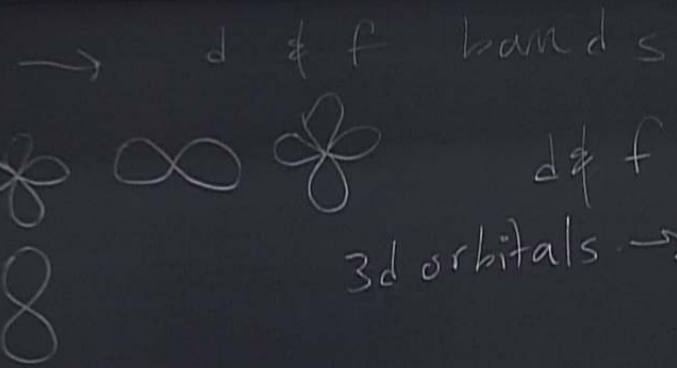
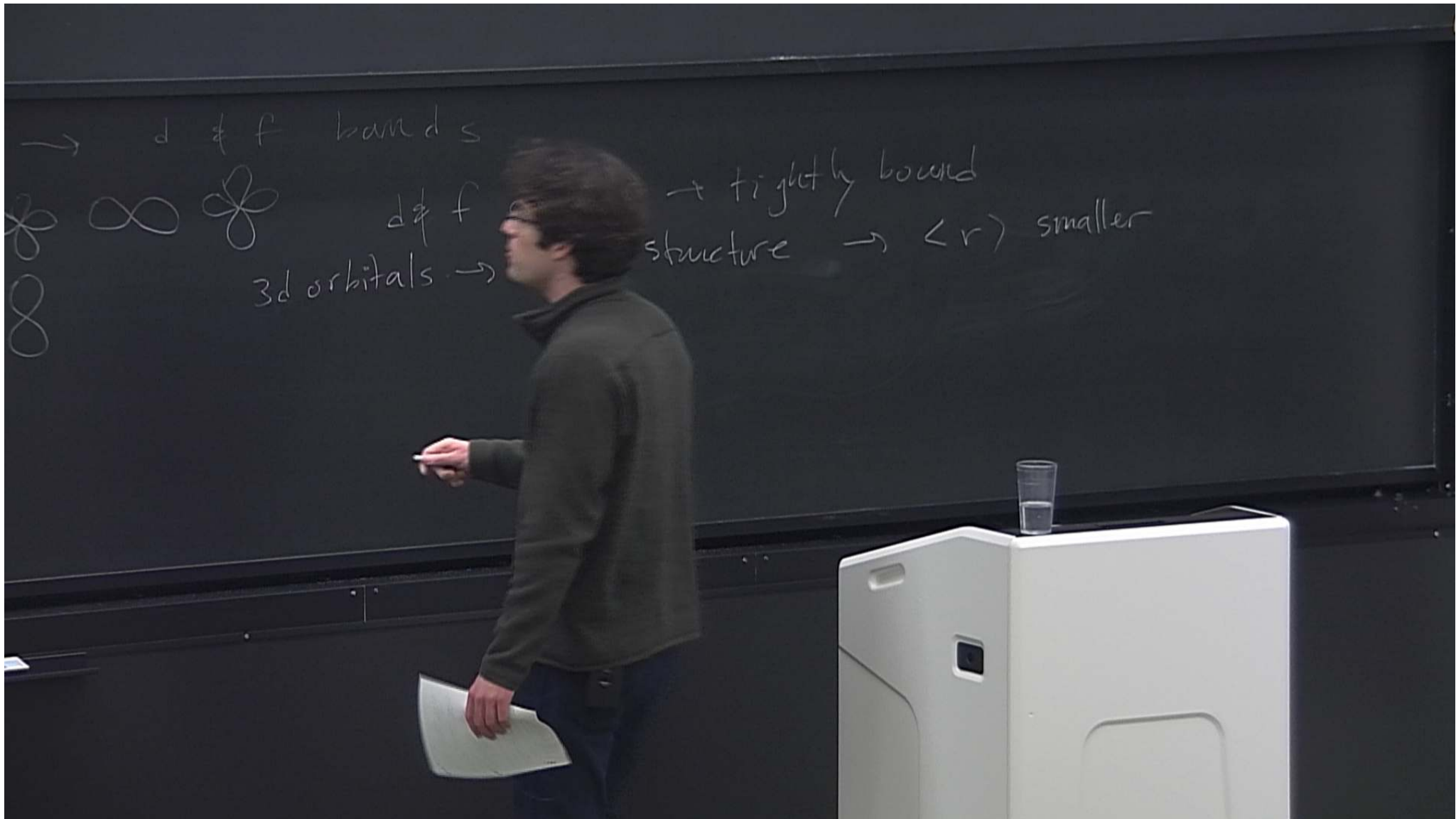


strong correlations  $\rightarrow$  d & f bands

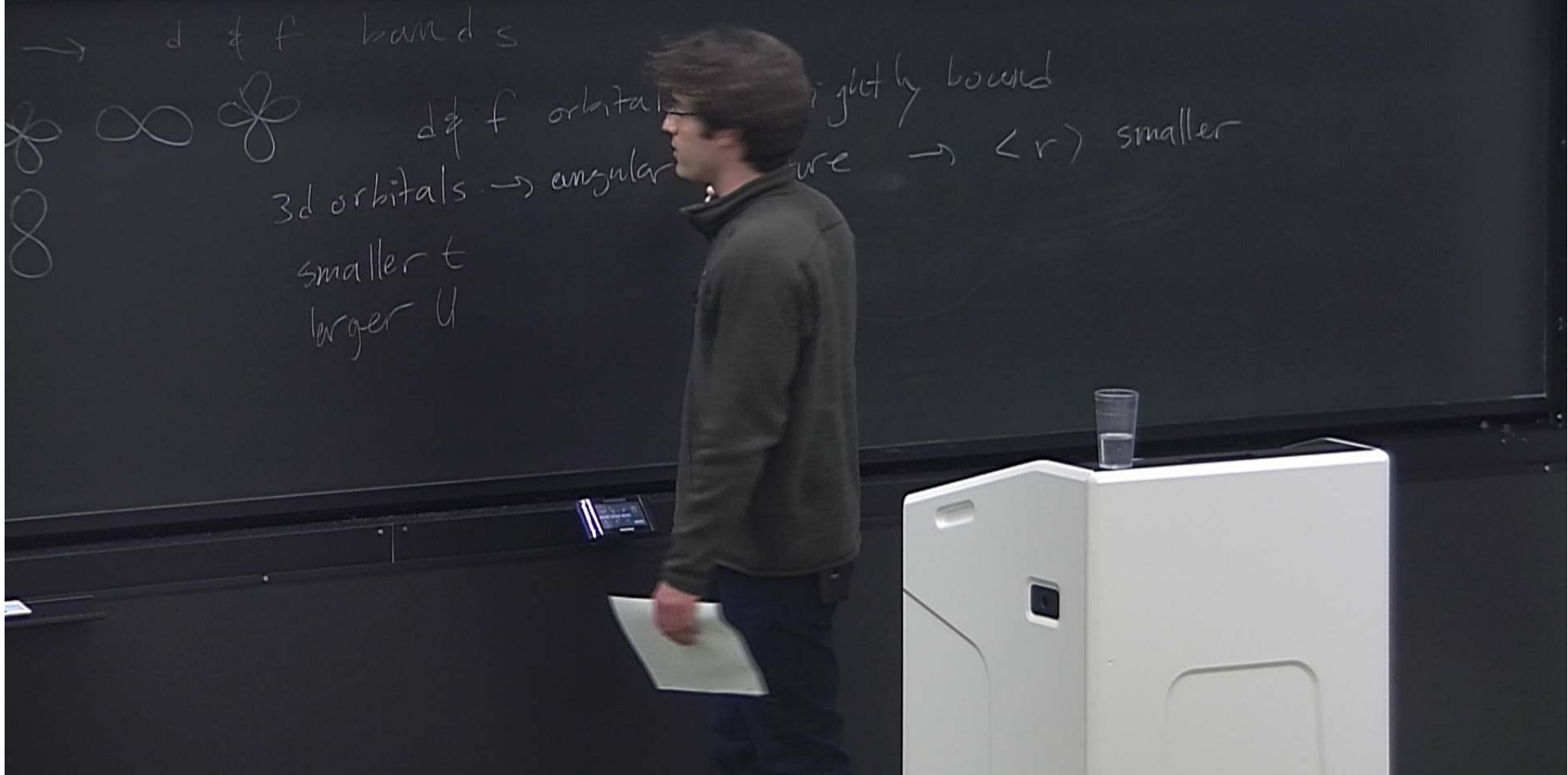


strong correlations  $\rightarrow$  d & f bands





→ tightly bound structure →  $\langle r \rangle$  smaller

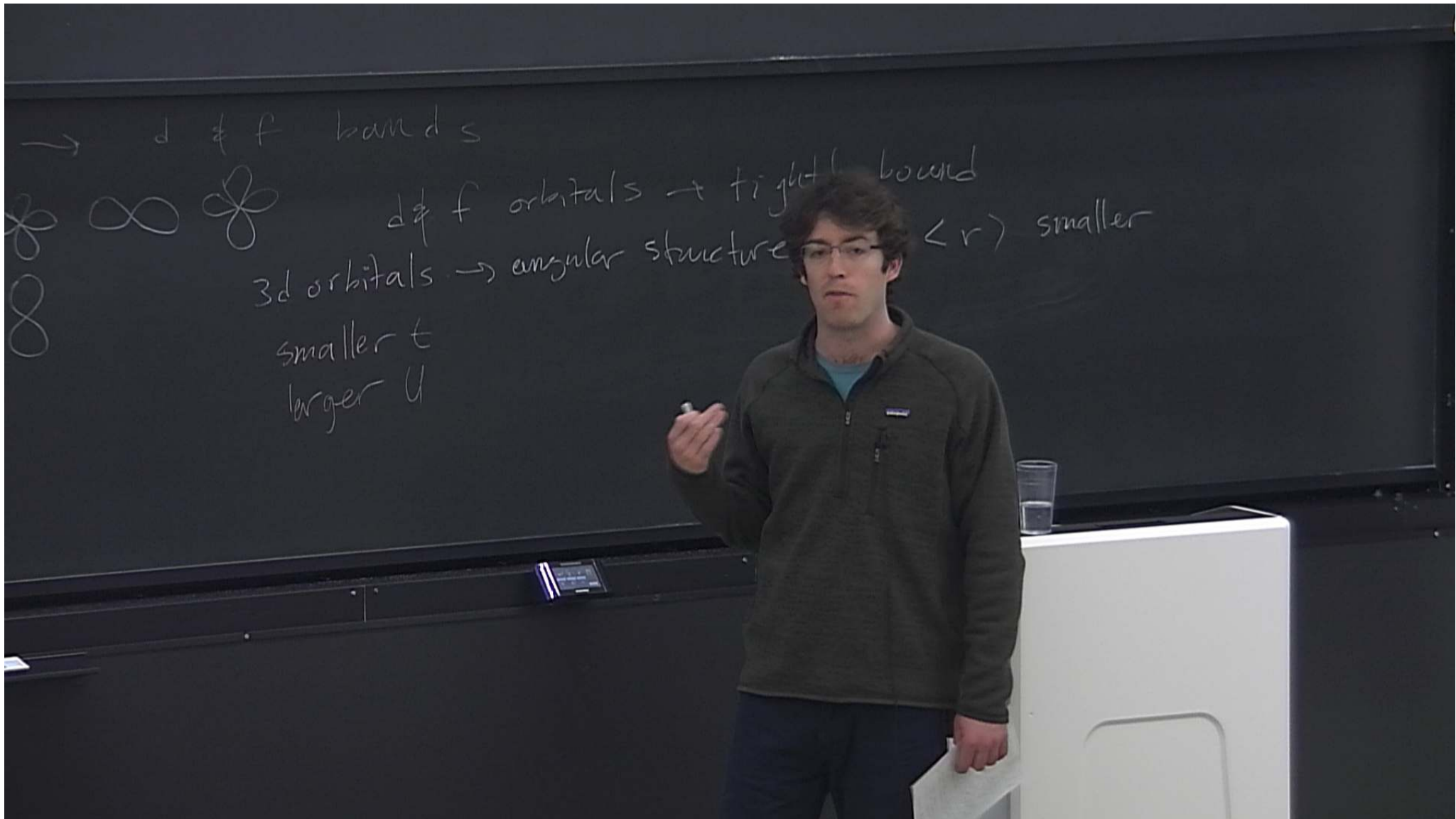


→ d & f bands



d & f orbitals are tightly bound  
3d orbitals → angular momentum →  $\langle r \rangle$  smaller

smaller  $\epsilon$   
larger U



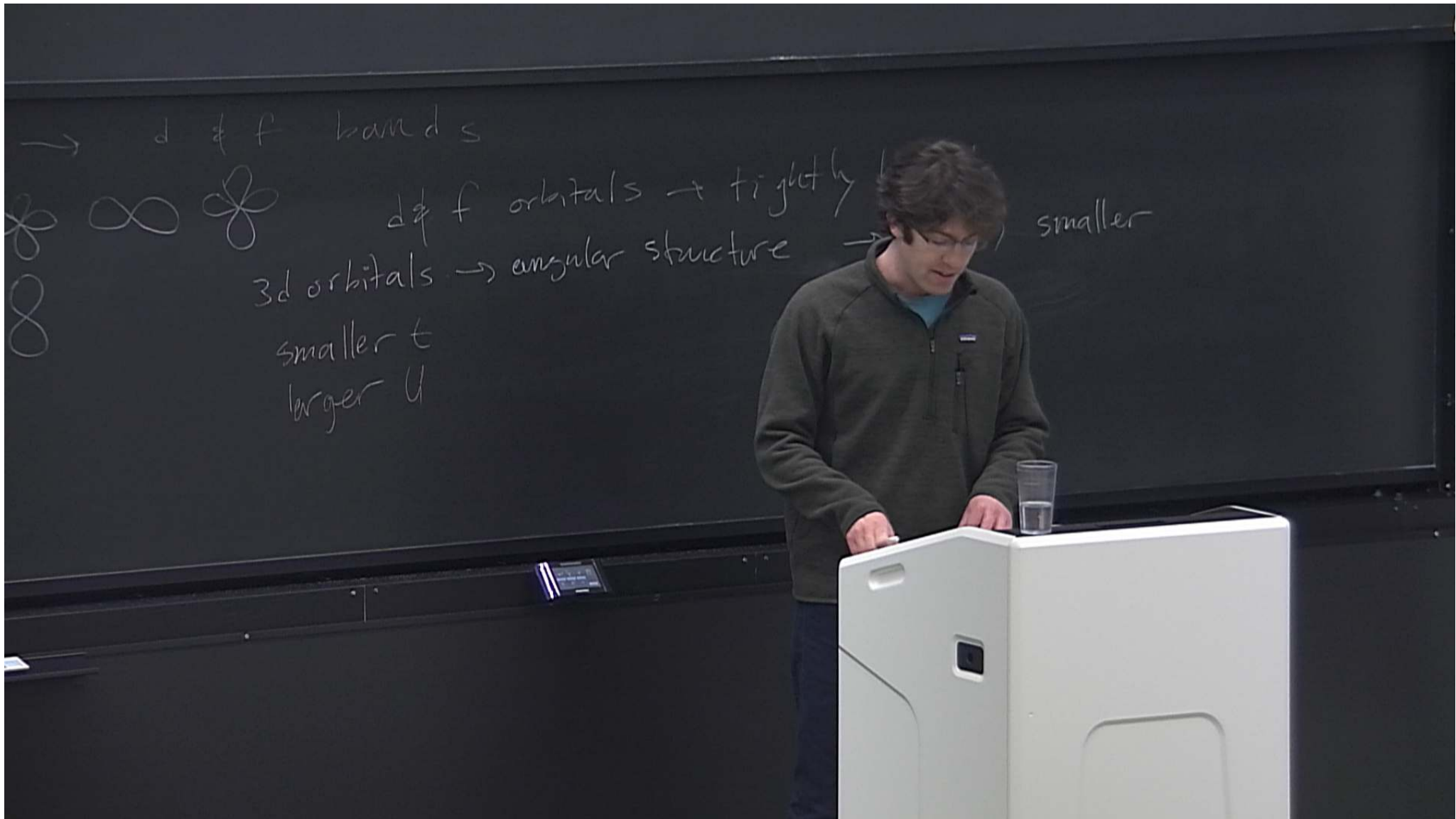
→ d & f bands



d & f orbitals → tightly bound

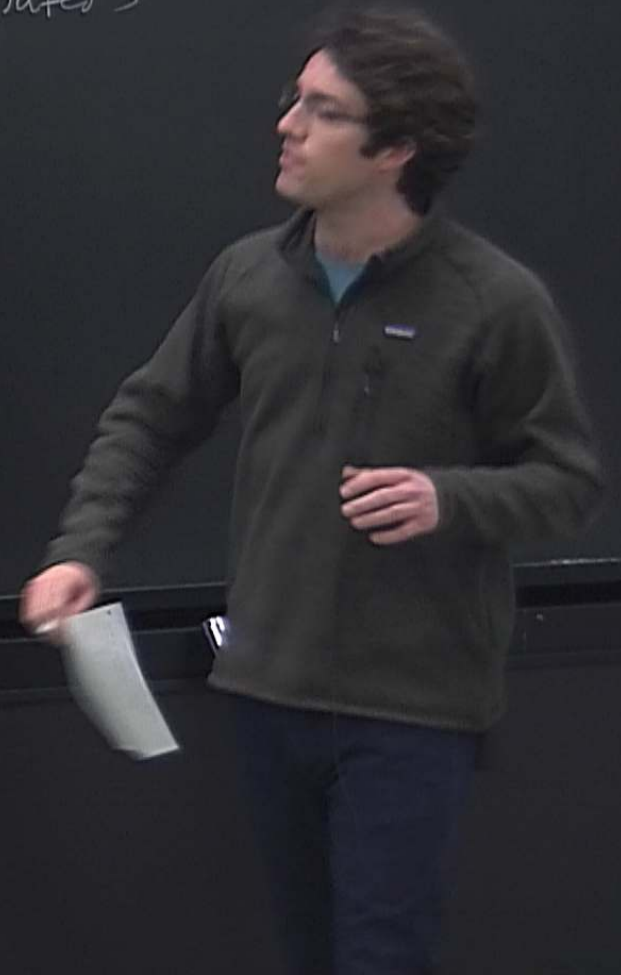
3d orbitals → angular structure < r > smaller

smaller t  
larger U



lattice models  $\begin{cases} \rightarrow \\ \rightarrow \end{cases}$  simple for

lattice models  $\begin{cases} \rightarrow \text{simple for people} \\ \rightarrow \text{simple for computers} \end{cases}$



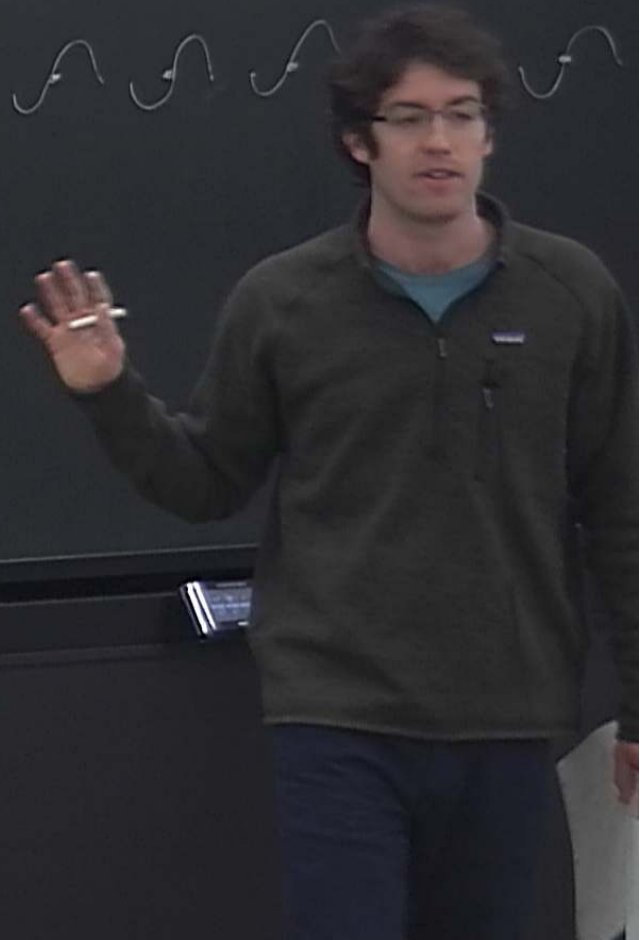
lattice models  $\begin{cases} \rightarrow \text{simple for people } \checkmark \\ \rightarrow \text{simple for computers } \checkmark \end{cases}$

bands Bloch states  $\phi_{\alpha\vec{k}}(\vec{r})$



lattice models  $\begin{cases} \rightarrow \text{simple for people } \checkmark \\ \rightarrow \text{simple for computers } \checkmark \end{cases}$

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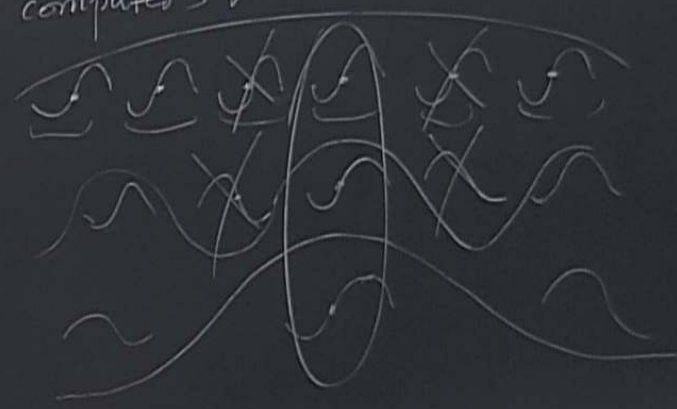
$k \approx 0$

$k \approx \pi$

lattice models  $\begin{cases} \rightarrow \text{simple for people } \checkmark \\ \rightarrow \text{simple for computers } \checkmark \end{cases}$

bands Bloch states  $\psi_{\vec{k}}(\vec{r})$

$$\phi_{\alpha j}(\vec{r}) = \frac{1}{\sqrt{N}}$$



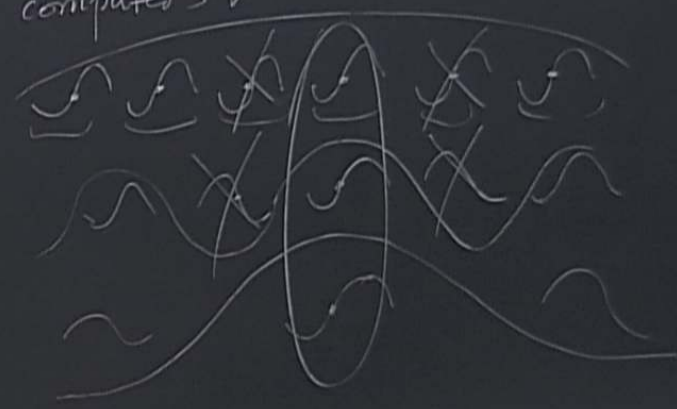
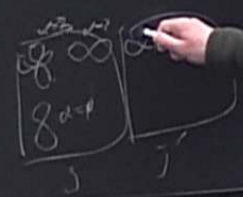
$k=0$

$k=\pi$

lattice models  $\begin{cases} \rightarrow \text{simple for people } \checkmark \\ \rightarrow \text{simple for computers } \checkmark \end{cases}$

bands Bloch  $\phi_{\alpha\vec{k}}(\vec{r})$

$$\phi_{\alpha j}(\vec{r}) = \frac{1}{\sqrt{N}} \sum_{\vec{k} \in \text{BZ}} e^{i\vec{k}\cdot\vec{r}} \phi_{\alpha\vec{k}}(\vec{r})$$



$k \approx 0$

$k \approx \pi$

$$\hat{c}_{i\sigma}^+ = \int_{\vec{r}} \phi_i(\vec{r}) \hat{\psi}_{\vec{r}\sigma}^+$$

$(\alpha i) \rightarrow j$

$$\hat{c}_{i\sigma}^{\dagger} = \int_{\vec{r}} \phi_i(\vec{r}) \hat{\psi}_{\vec{r}\sigma}^{\dagger} \quad (\Rightarrow \hat{\psi}_{\vec{r}\sigma}^{\dagger} = \sum_i \phi_i(\vec{r}) \hat{c}_{i\sigma}^{\dagger})$$

$(\alpha i) \rightarrow j$  (orbitals)

$\hat{H}$

$$\hat{C}_{i\sigma}^\dagger \stackrel{\text{def}}{=} \int_{\vec{r}} \phi_i(\vec{r}) \hat{\psi}_{\vec{r}\sigma}^\dagger \quad (\Rightarrow \hat{\psi}_{\vec{r}\sigma}^\dagger = \sum_i \phi_i(\vec{r}) \hat{C}_{i\sigma}^\dagger = \sum_i \phi_i(\vec{r}) \phi_i^\dagger)$$

$(\propto i) \rightarrow j$  (orbitals)

$$\hat{H} = \sum_{\sigma} \int_{\vec{r}} \hat{\psi}_{\vec{r}\sigma}^\dagger \left[ -\frac{1}{2} \nabla^2 + U(\vec{r}) \right] \hat{\psi}_{\vec{r}\sigma} + \frac{1}{2} \int_{\vec{r}, \vec{r}'} U_E(\vec{r} - \vec{r}')$$

$$\hat{C}_{i\sigma}^\dagger \stackrel{\text{def}}{=} \int_{\vec{r}} \phi_i(\vec{r}) \hat{\psi}_{\vec{r}\sigma}^\dagger \quad (\Rightarrow) \quad \hat{\psi}_{\vec{r}\sigma}^\dagger = \sum_i \phi_i(\vec{r}) \hat{C}_{i\sigma}^\dagger = \sum_i \int_{\vec{r}'} \phi_i(\vec{r}) \phi_i^*(\vec{r}') \hat{\psi}_{\vec{r}'\sigma}^\dagger$$

$(\propto i) \rightarrow j$  (orbitals)

$$\hat{H} = \sum_{\sigma} \int_{\vec{r}} \hat{\psi}_{\vec{r}\sigma}^\dagger \left[ -\frac{1}{2} \nabla^2 + v(\vec{r}) \right] \hat{\psi}_{\vec{r}\sigma} + \frac{1}{2} \int_{\vec{r}, \vec{r}'} U(\vec{r}, \vec{r}') \hat{\psi}_{\vec{r}\sigma}^\dagger \hat{\psi}_{\vec{r}'\sigma} \hat{\psi}_{\vec{r}'\sigma}^\dagger \hat{\psi}_{\vec{r}\sigma}$$

$$\hat{C}_{i\sigma}^+ \stackrel{\text{def}}{=} \int_{\vec{r}} \phi_i(\vec{r}) \hat{\psi}_{\vec{r}\sigma}^+ \quad \left( \Rightarrow \hat{\psi}_{\vec{r}\sigma} = \sum_i \phi_i(\vec{r}) \hat{C}_{i\sigma} = \sum_i \underbrace{\int_{\vec{r}'} \phi_i(\vec{r}) \phi_i^*(\vec{r}')}_{\delta(\vec{r}-\vec{r}')} \hat{\psi}_{\vec{r}'\sigma} = \hat{\psi}_{\vec{r}\sigma} \right)$$

$(\alpha i) \rightarrow j$  (orbitals)

$$\hat{H} = \sum_{\sigma} \int_{\vec{r}} \hat{\psi}_{\vec{r}\sigma}^+ \left[ -\frac{1}{2} \nabla^2 + U(\vec{r}) \right] \hat{\psi}_{\vec{r}\sigma} + \frac{1}{2} \int_{\vec{r}, \vec{r}'} U_E(\vec{r}-\vec{r}')$$

$$\hat{c}_{i\sigma}^\dagger \stackrel{\text{def}}{=} \int_{\vec{r}} \phi_i(\vec{r}) \hat{\psi}_{\vec{r}\sigma}^\dagger \quad \left( \Rightarrow \hat{\psi}_{\vec{r}\sigma} = \sum_i \phi_i(\vec{r}) \hat{c}_{i\sigma} = \sum_i \underbrace{\int_{\vec{r}'} \phi_i(\vec{r}) \phi_i^*(\vec{r}')}_{\delta(\vec{r}-\vec{r}')} \hat{\psi}_{\vec{r}'\sigma} = \hat{\psi}_{\vec{r}\sigma} \right)$$

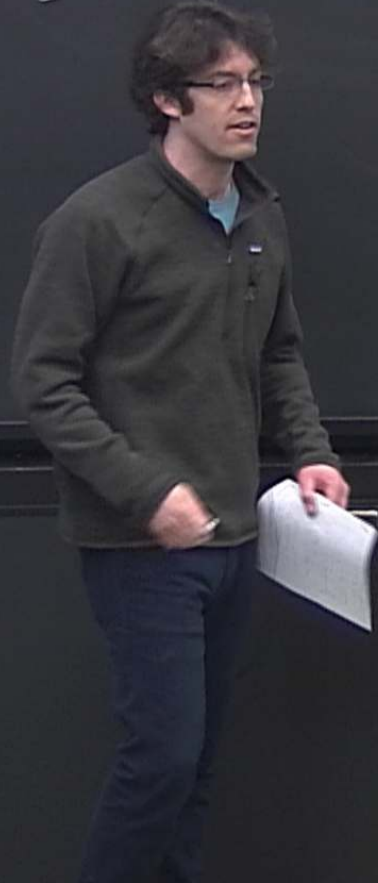
(\(\alpha\)) \(\rightarrow\) orbitals

$$\hat{H} = \sum_{\vec{r}\sigma} \left[ -\frac{1}{2} \nabla^2 + v(\vec{r}) \right] \hat{\psi}_{\vec{r}\sigma} + \frac{1}{2} \int_{\vec{r}, \vec{r}'} U_E(\vec{r}-\vec{r}')$$



(tunneling  
hopping)

$$t_{ij} = -\langle \phi_i | \hat{H}_0 | \phi_j \rangle = -\int_{\mathbb{R}} \phi_i^*(\mathbb{R}) \left[ -\frac{1}{2} \nabla^2 + V(\mathbb{R}) \right] \phi_j(\mathbb{R})$$



(tunneling  
hopping)

$$t_{ij} = -\langle \phi_i | \hat{H}_0 | \phi_j \rangle = -\int \phi_i^*(\mathbf{r}) \left[ -\frac{1}{2} \nabla^2 + V(\mathbf{r}) \right] \phi_j(\mathbf{r})$$



(tunneling  
hopping)

$$t_{ij} = -\langle \phi_i | \hat{H}_0 | \phi_j \rangle = -\int_{\vec{r}} \phi_i^*(\vec{r}) \left[ -\frac{1}{2} \nabla^2 + v(\vec{r}) \right] \phi_j(\vec{r})$$

$$V_{ijkl} = \frac{1}{2} \int v_{ee}(\vec{r}-\vec{r}') \phi_i^*(\vec{r}) \phi_j(\vec{r}) \phi_k^*(\vec{r}') \phi_l(\vec{r}')$$

→ exact transformation / change of basis

→ hope that  $t_{ij}, V_{ijkl}$  short ranged

→ hope that  $t_{ij}, V_{ijke}$  short ranged

What  $t$ 's and  $V$ 's are there?

- crystal f:  $t_{ij}$   $E_{\text{exp}} = t_{(ij)(\beta\gamma)}$

→ hope that  $t_{ij}, V_{ijke}$  short ranged

What  $t$ 's and  $V$ 's are there?

crystal field matrix  $E_{\alpha\beta} = t_{\alpha i}(\beta i)$

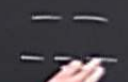
→ symmetries, group theory

→ hope that  $t_{ij}, V_{ij}$  short ranged

What  $t$ 's and  $V$ 's are there?

- crystal field matrix  $E_{\alpha\beta} = t_{\alpha\beta}$   
→ symmetries, group theory

$3d e^-$ , octahedral environment, phenomenological potential (x-tal field)



→ hope that  $t_{ij}, V_{ijke}$  short ranged

What  $t$ 's and  $V$ 's are there?

- crystal field matrix  $E_{\alpha\beta} = t_{\alpha\beta}(\beta\epsilon)$   
→ sym, group theory

$3d^5$ , octahedral environment, phenomenological potential (crystal field)  
→  $3g$   
→  $3t_{2g}$



→ hope that  $t_{ij}, V_{ij}$  short ranged

What  $t$ 's and  $V$ 's are there?

- crystal field matrix  $E_{\alpha\beta} = t_{\alpha\beta}$   
→ symmetries, group theory

$3d e^-$ , octahedral environment, phenomenological potential (crystal field)

→  $-\text{---} 3g$   
 $-\text{---} 3t_{2g}$  → Jahn-Teller distortions

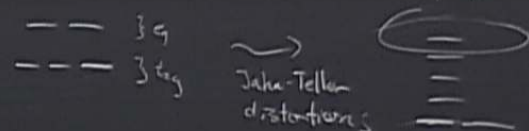
→ hope that  $t_{ij}$ ,  $V_{ij}$  short ranged

What  $t$ 's and  $V$ 's are there?

- crystal field matrix  $E_{sp} = t_{ij}(p_i)$   
 → symmetry theory

- next-nearest neighbor hopping  
 $t \stackrel{\text{def}}{=} t$
- 

$3d^5$ , octahedral environment, phenomenological parameters



→  $3e_g$   
 $3t_{2g}$   
 Jahn-Teller distortions

• Most important interaction

$$\frac{U}{Z} \stackrel{\text{def}}{=} V_{\text{coul}} = \frac{1}{2} \int_{\vec{r}, \vec{r}'} V_{\text{ec}}(\vec{r} - \vec{r}') |\phi_c(\vec{r})|^2 |\phi_c(\vec{r}')|^2 \quad (\sim \sigma^2)$$

• Most important interaction

$$\frac{U}{2} \stackrel{\text{def}}{=} V_{cccc} = \frac{1}{2} \int_{\vec{r}, \vec{r}'} v_{cc}(\vec{r} - \vec{r}') |\phi_c(\vec{r})|^2 |\phi_c(\vec{r}')|^2 \quad \left( \sim \overbrace{c_{cc}^{\uparrow} c_{cc}^{\downarrow}}^{\uparrow} \overbrace{c_{cc}^{\downarrow} c_{cc}^{\uparrow}}^{\downarrow} \right)$$

further-neighbor  $V$ 's

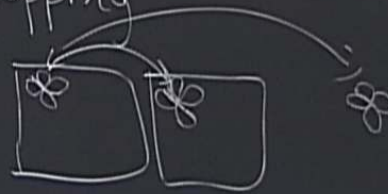
$$V_{cc_{ij}} = \frac{1}{2} \int v_{cc}(\vec{r} - \vec{r}') |\phi_c(\vec{r})|^2 |\phi_c(\vec{r}')|^2$$

→ hope that  $t_{ij}, V_{ij}$  short ranged

What  $t$ 's and  $V$ 's are there?

- crystal field matrix  $E_{\alpha\beta} = t(\alpha i)(\beta i)$   
 → symmetries, group theory

- next-neighbor "hopping"  
 $t = t(\alpha i)(\alpha j)$   
 further-neighbor  
 $t$



$3d^5$ , octahedral environment phenomena  
 ~) ---  $3g$   
 ---  $3t_{2g}$

## "Exchange" mechanisms

- Coulomb exchange → short range, ferromagnetic  $\leftrightarrow$  int.
- Kinetic exchange → neighboring sites, antiferromagnetic spin

## "Exchange" mechanisms

- Coulomb exchange → short range, ferromagnetic spin int.
- Kinetic exchange → neighboring sites, magnet spin

$$\sum_{ij} V_{ij} j_i$$

"Exchange" mechanisms (spin interactions)

◦ Coulomb exchange → short range, ferromagnetic spin int.

◦ Kinetic exchange → neighboring sites, antiferromagnetic spin

$$\sum_{\sigma\sigma'} V_{ijj'j''} c_{i\sigma}^\dagger c_{j\sigma} c_{j'\sigma'}^\dagger c_{j''\sigma'}$$



## "Exchange" mechanisms (spin interactions)

◦ Coulomb exchange → short range, ferromagnetic spin int.

◦ Kinetic exchange → neighboring sites, antiferromagnet spin

$$\sum_{\sigma\sigma'} V_{ijji} \underbrace{c_{i\sigma}^\dagger c_{j\sigma} c_{j\sigma'}^\dagger c_{i\sigma'}}_{\text{Coulomb exchange}} = -V_{ijji} \sum_{\sigma\sigma'} (c_{i\sigma}^\dagger c_{i\sigma'}) (c_{j\sigma'}^\dagger c_{j\sigma}) + (\sim \vec{p}^2)$$

# "Exchange" mechanisms (spin interactions)

o Coulomb exchange → short range, ferromagnetic spin int.

o Kinetic exchange → neighboring sites, ferromagnetic spin

$$\sum_{\sigma\sigma'} V_{ijji} c_{i\sigma}^\dagger c_{j\sigma} c_{j\sigma'}^\dagger c_{i\sigma'} = -V_{ijji} \sum_{\sigma\sigma'} (c_{j\sigma'}^\dagger c_{j\sigma}) + (\sim \hat{n})$$

$$S_i^z \stackrel{\text{def}}{=} \frac{1}{2} \sum_{\alpha\beta} \hat{\sigma}_{\alpha\beta} c_{i\alpha}^\dagger c_{i\beta}$$