

Title: Crab Nanoshots, Fast Radio Bursts and Cosmology

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Abstract: <p>We explore the brightness frontier in time domain radio astronomy and its possible usefulness for cosmology. It is argued that the brightest known source of emission, Crab nanoshots, are caused by Schwinger pair production. The same mechanism may be the source of Fast Radio Bursts (FRBs) if this emission is from coalescing neutron stars. It is then shown how using FRBs as triggers can extend the reach of gravitational radiation and neutrino telescopes. Finally we discuss how combining FRB monitoring, large neutrino telescopes, combined with preexisting galaxy catalogs could provide an accurate cosmological distance estimator.</p>

Crab Nanoshots, Fast Radio Bursts and Cosmology

Albert Stebbins
Fermilab
w/ Hojin Yoo (LBNL)

Perimeter Institute
13 January 2015
Waterloo, Canada

The Brightness Frontier

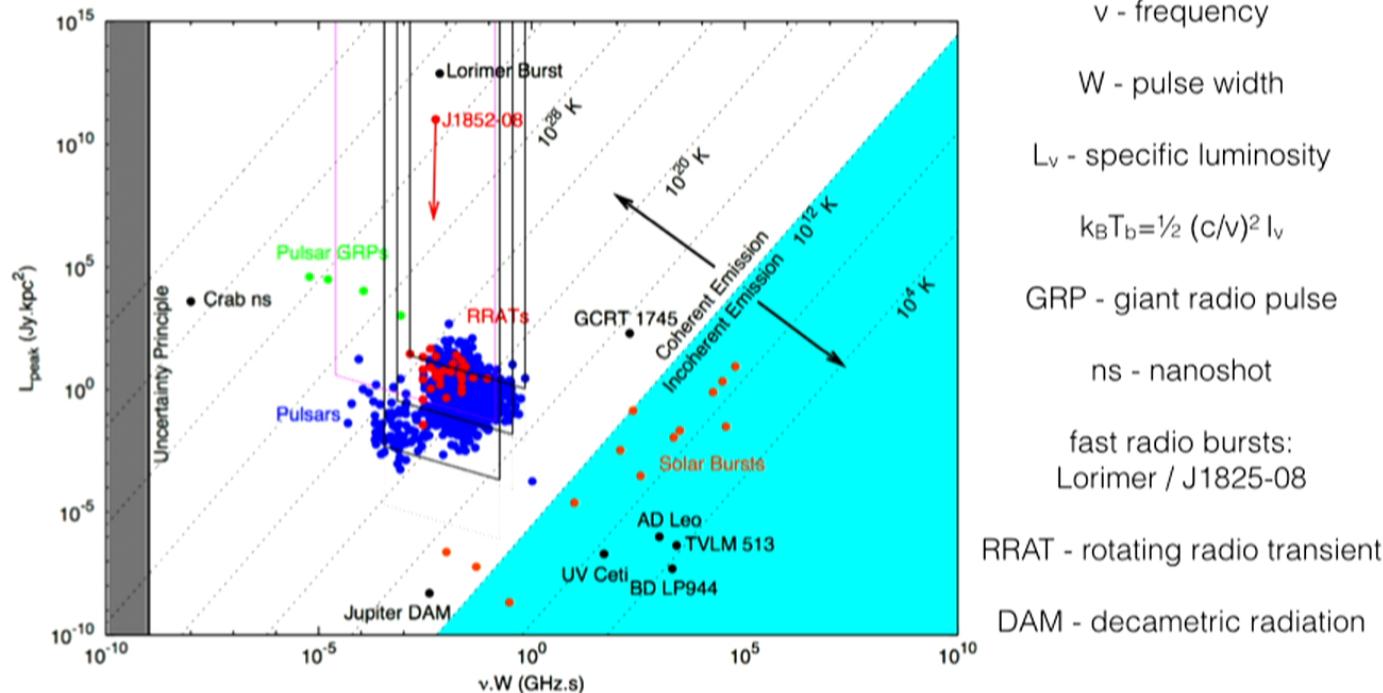
- In HEP-speak we have Energy, Intensity, Cosmic Frontiers
 - Radio astronomy has another interests frontier: Brightness
 - Short radio transients can be extremely bright!

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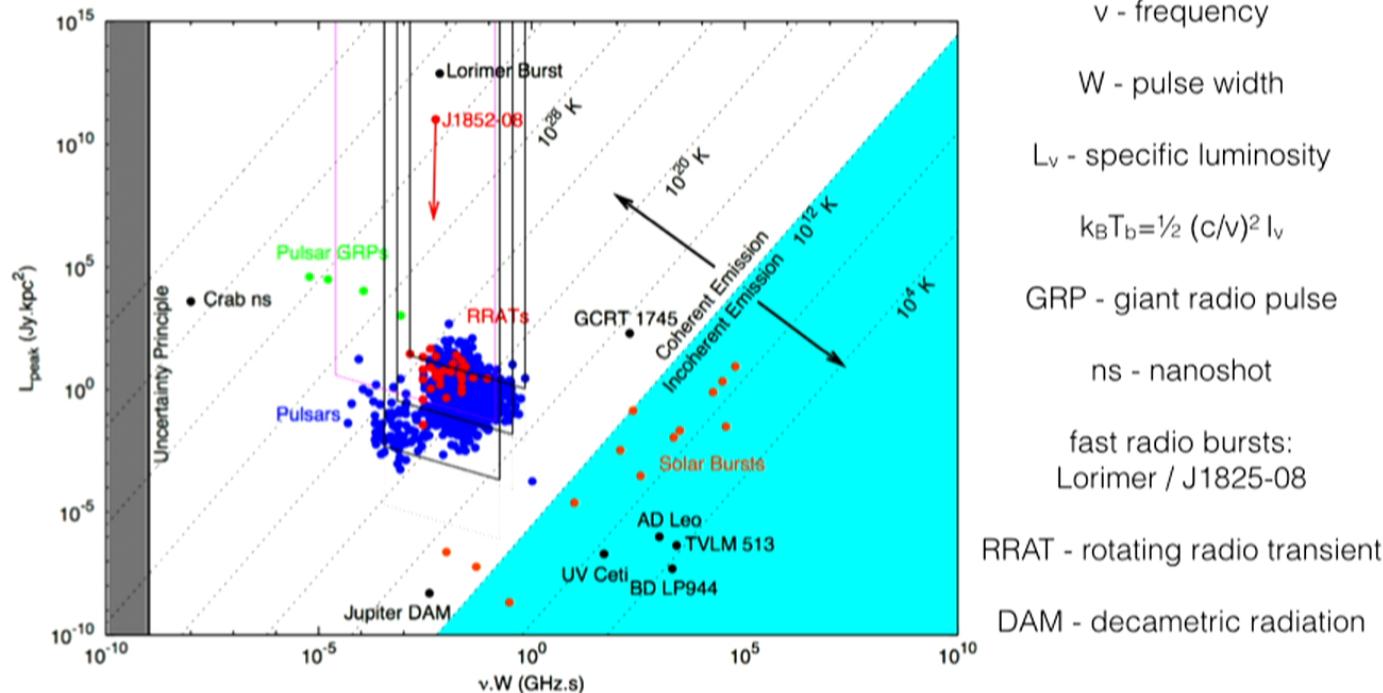
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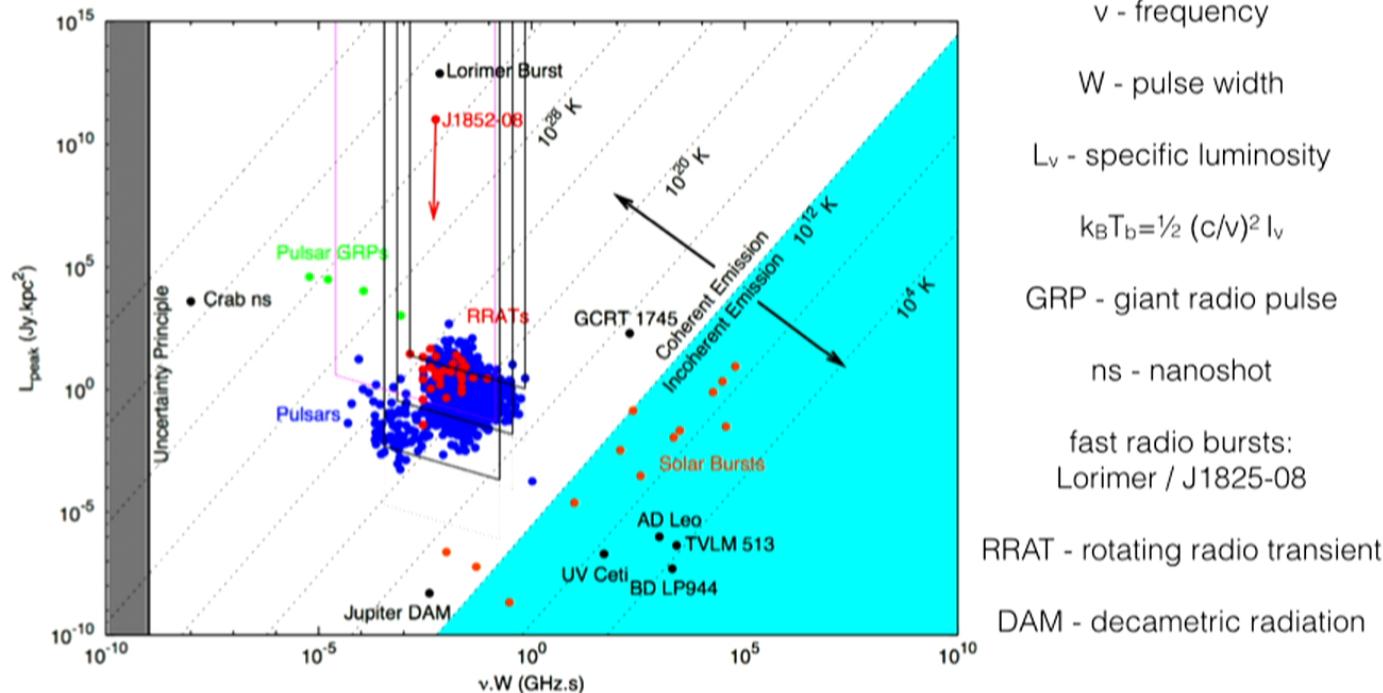
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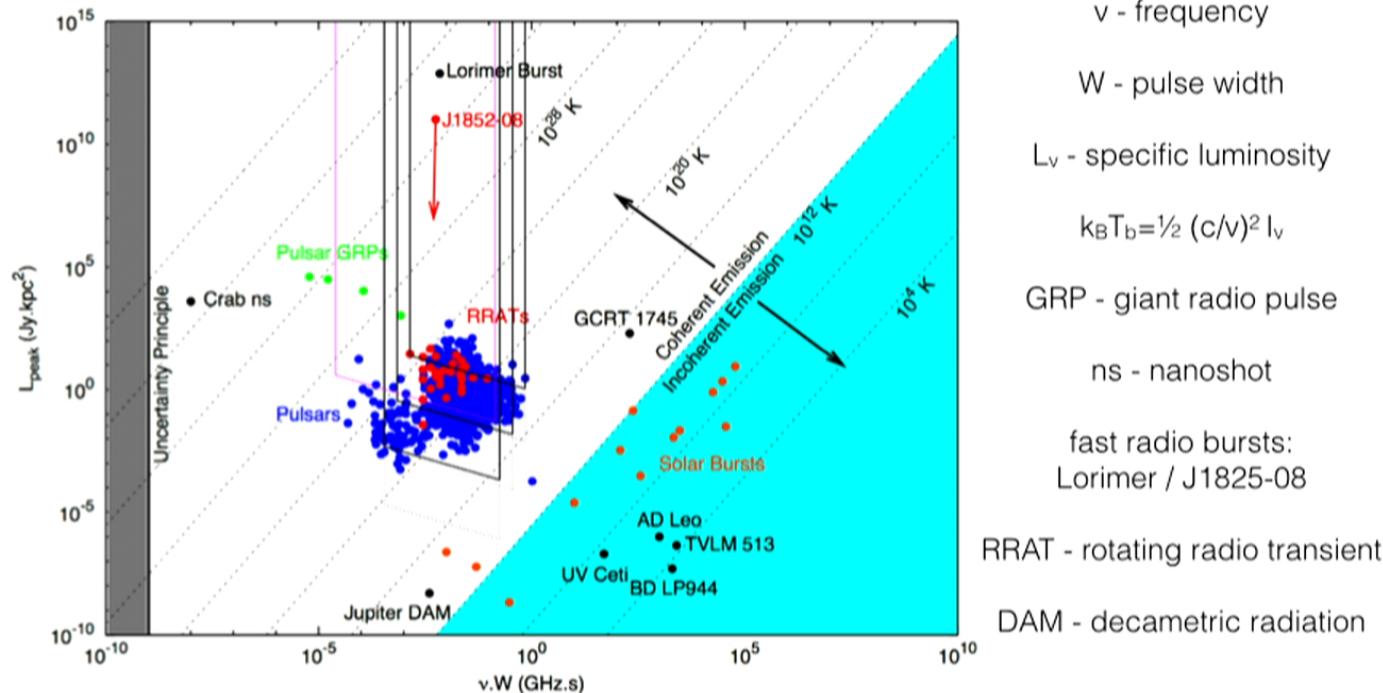
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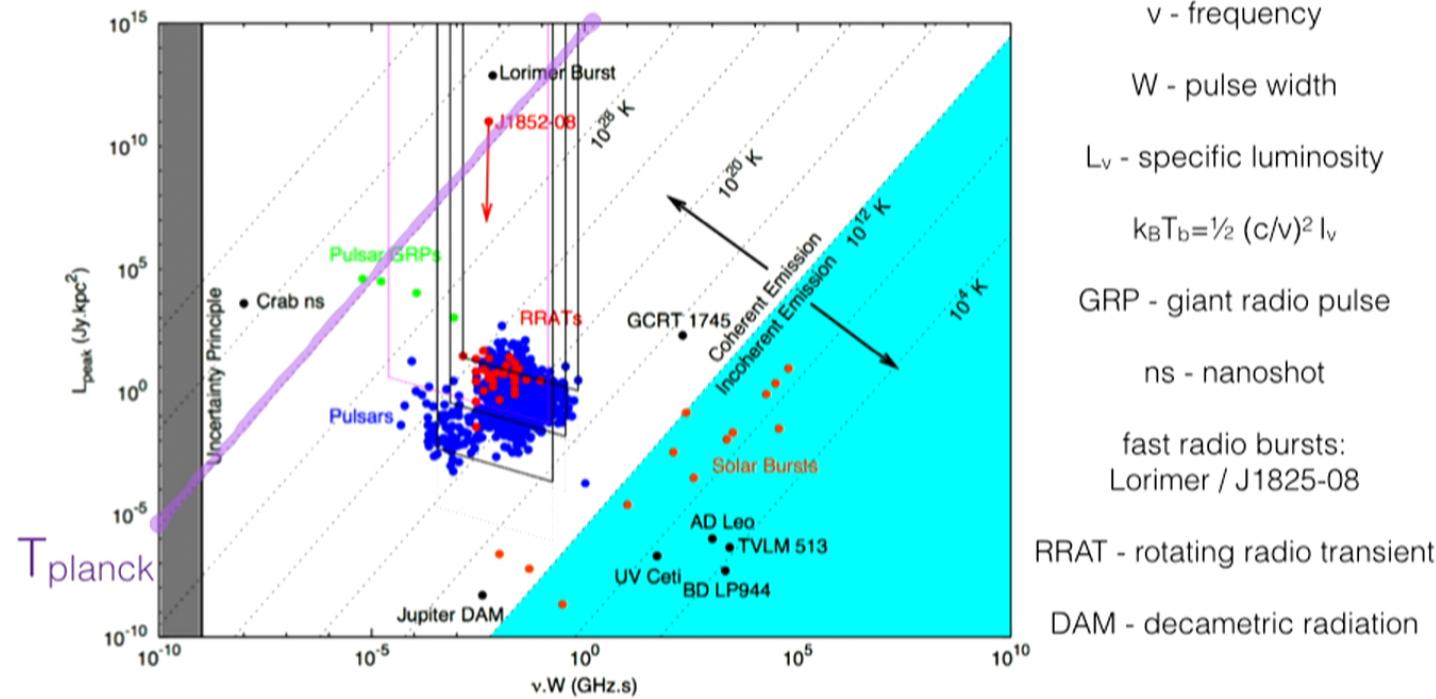
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Coherent vs Incoherent Radiation



Incoherent radiation
charged particles moving independently

$$I_v \propto N$$

synchrotron, free-free, ...

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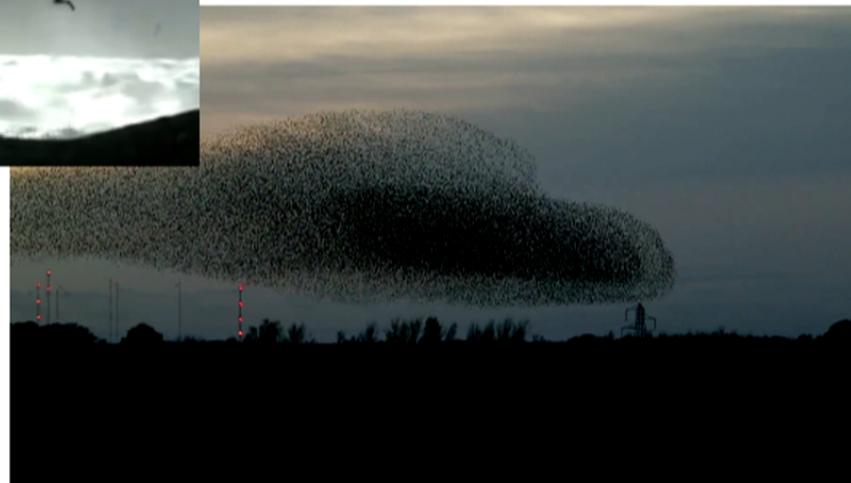


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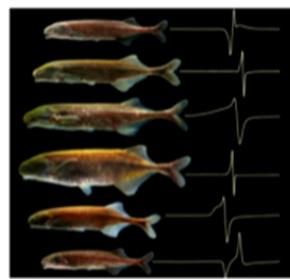
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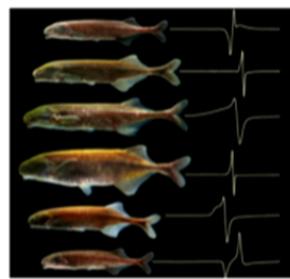
Macroscopic E&M and Thermodynamics

- the large T_b indicates a large amount of energy crammed into a small amount of phase space
- super-Planckian “T” is common for macroscopic motions
- super-Planckian T_b is a manifestation of **macroscopic** electro-magnetic phenomena
- macroscopic modes are usually very out of thermodynamic equilibrium and find various pathways to thermalize this energy.



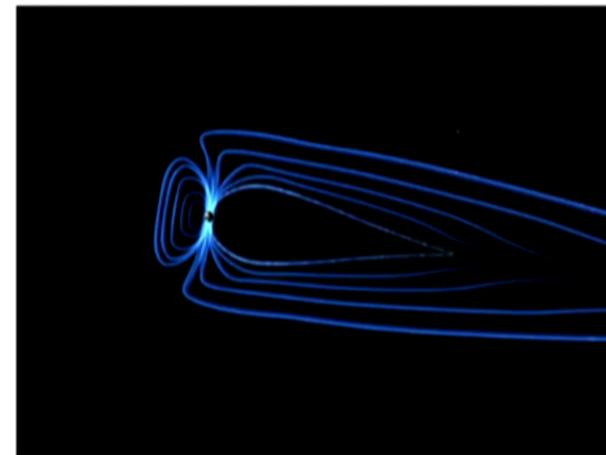
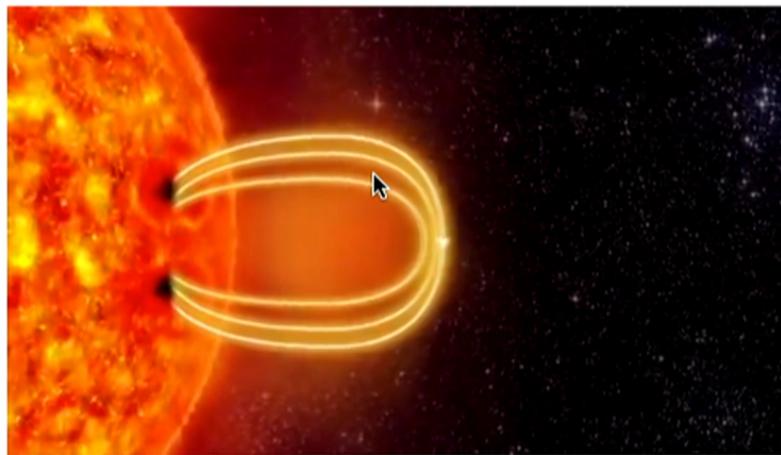
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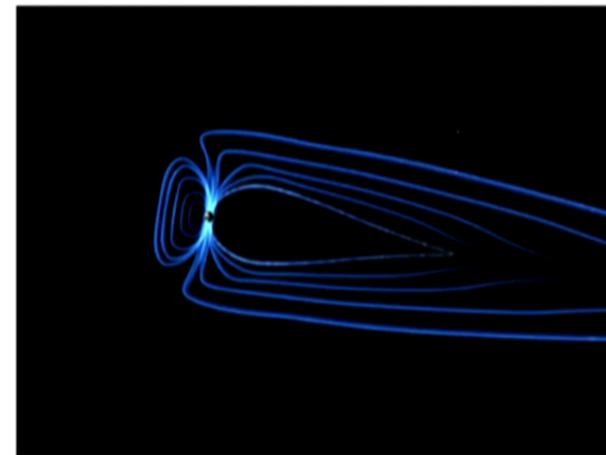
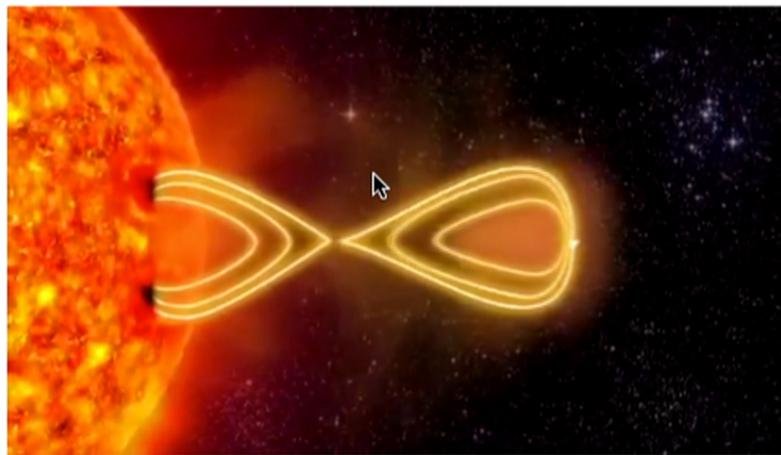
Magnetic Field Reconnection and Astrophysical EMP

- Very large magnetic fields (up to 10^{15} G) are often frozen into astrophysical plasmas but this energy is occasionally violently released in *magnetic field reconnection* events. On short timescales a the B field can oscillate into E field which accelerates and separates charges which also produces large EMP.



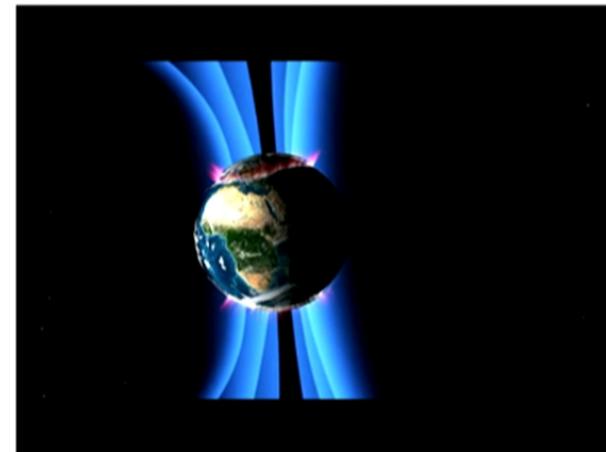
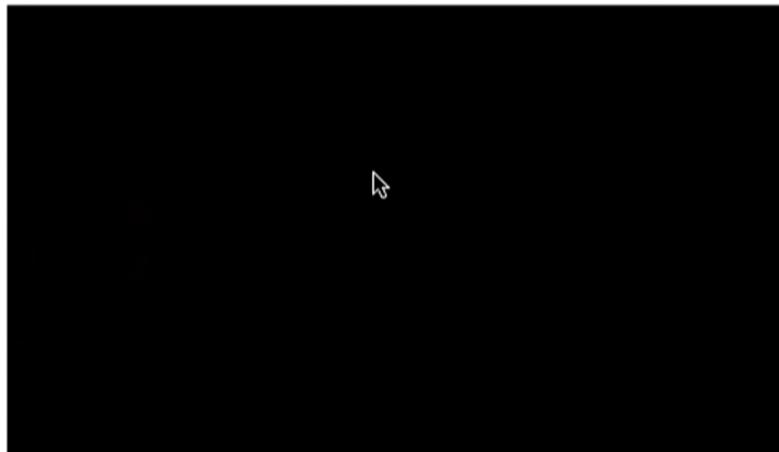
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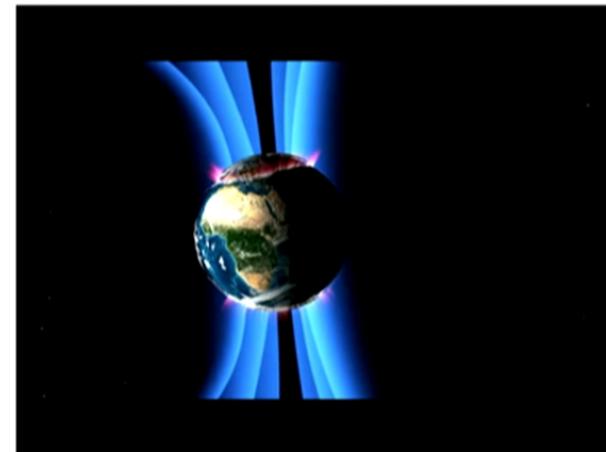
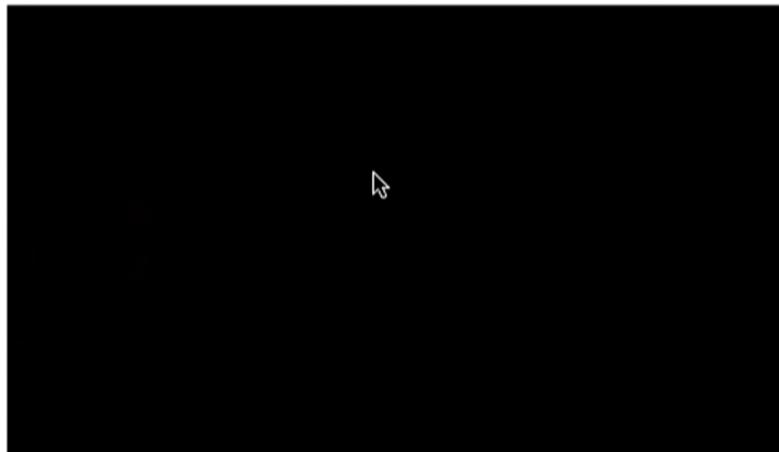
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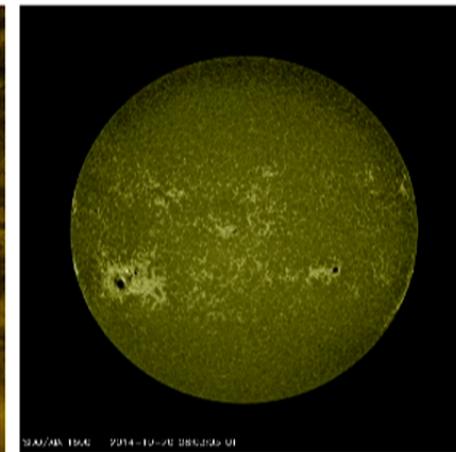
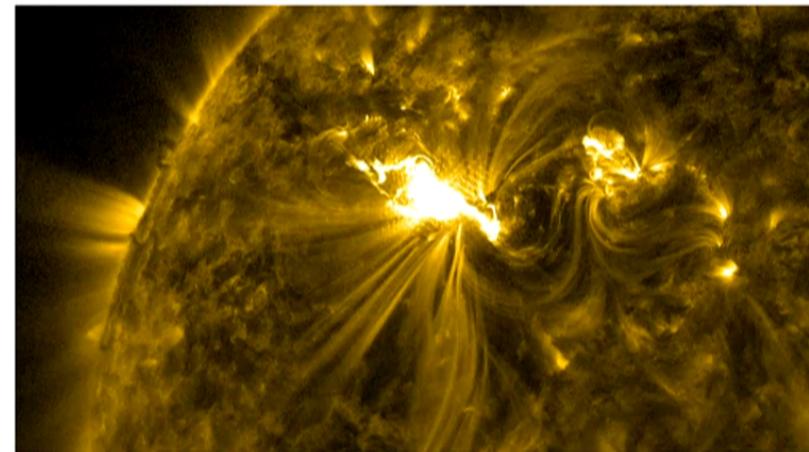


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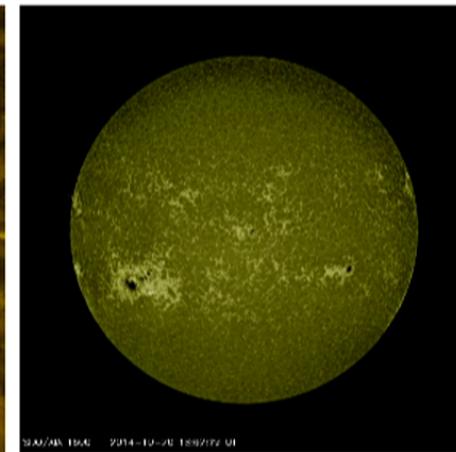
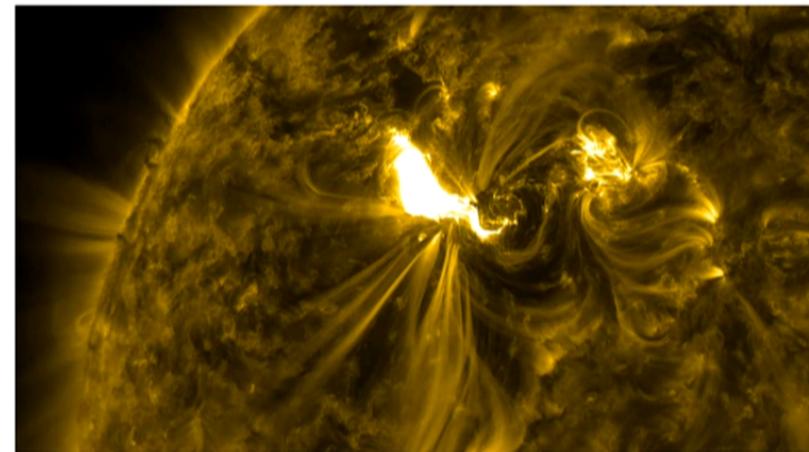
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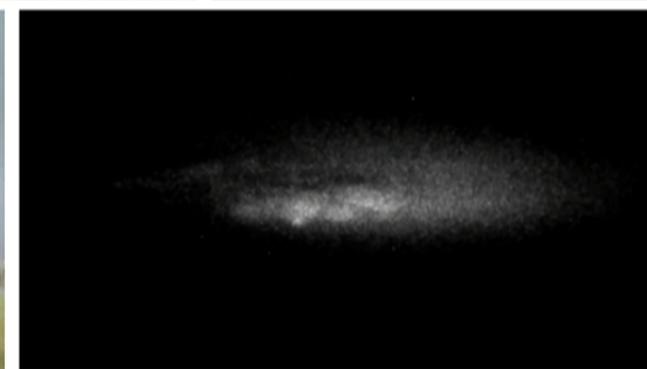
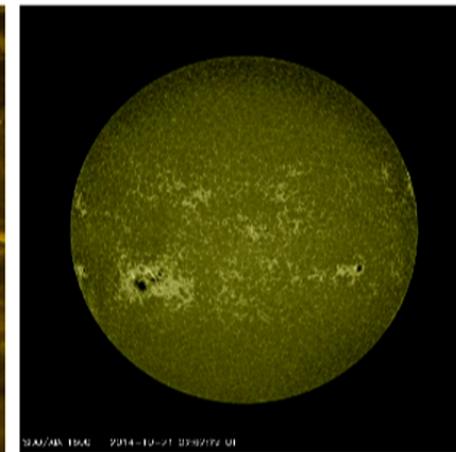
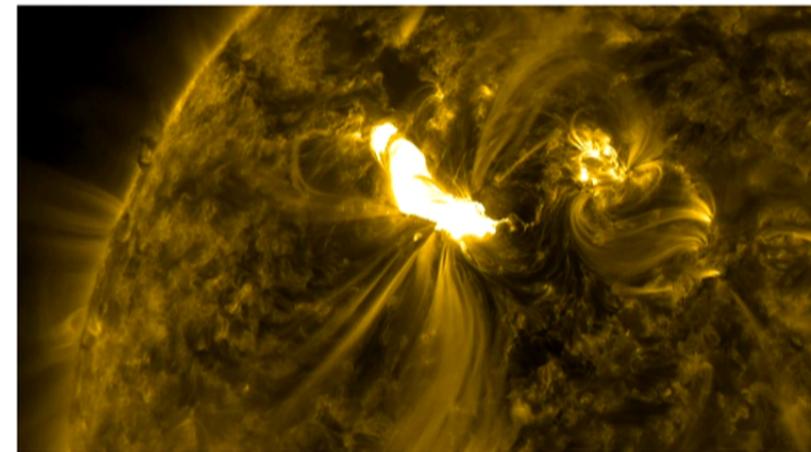
Magnetic Field Reconnections and other types of Sparks are Bright !



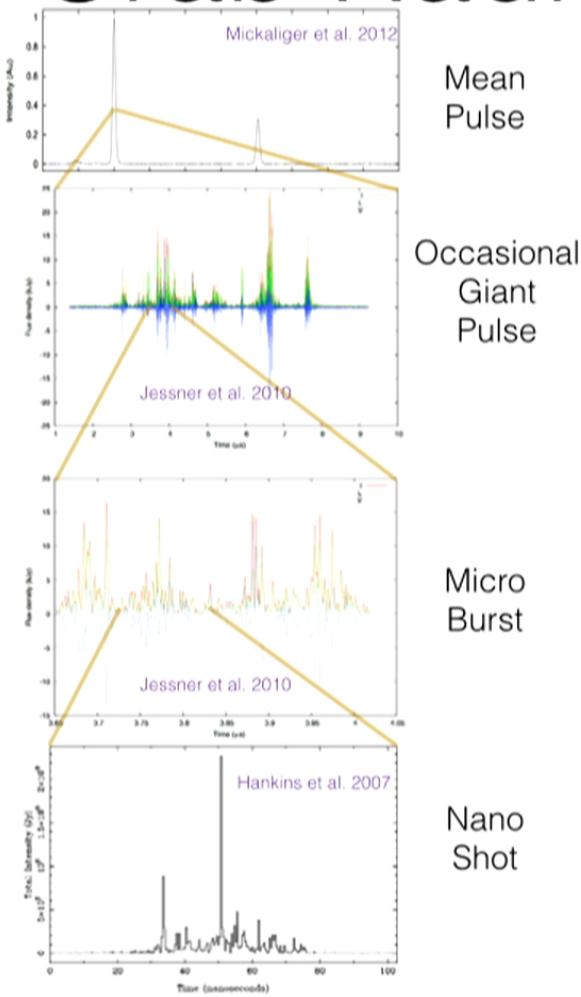
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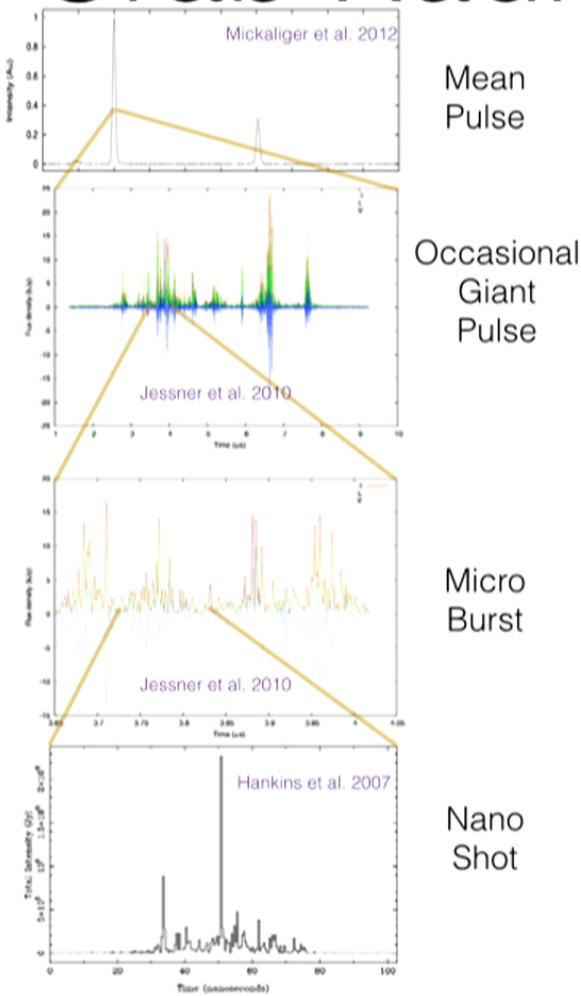
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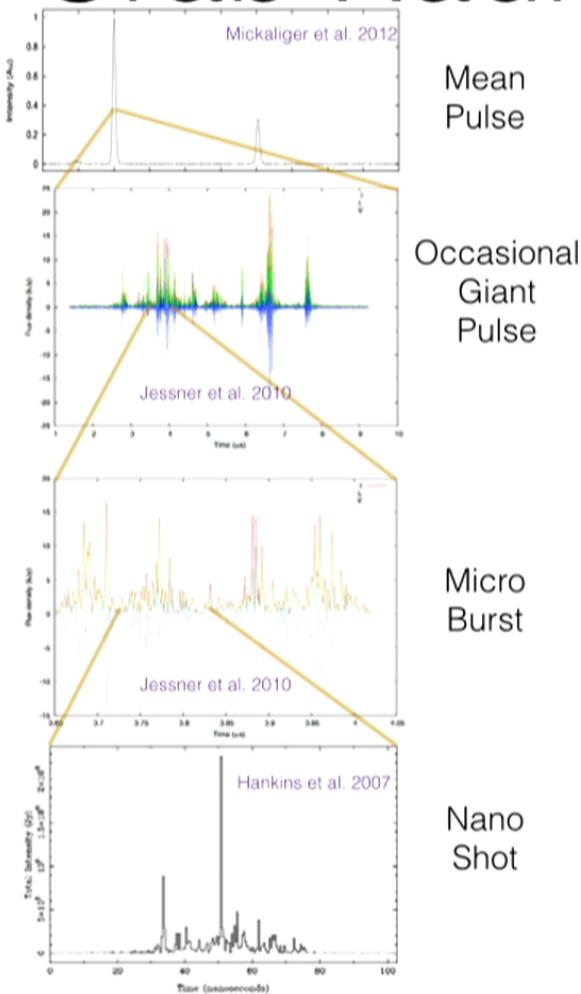
Crab Radio Pulse Structure



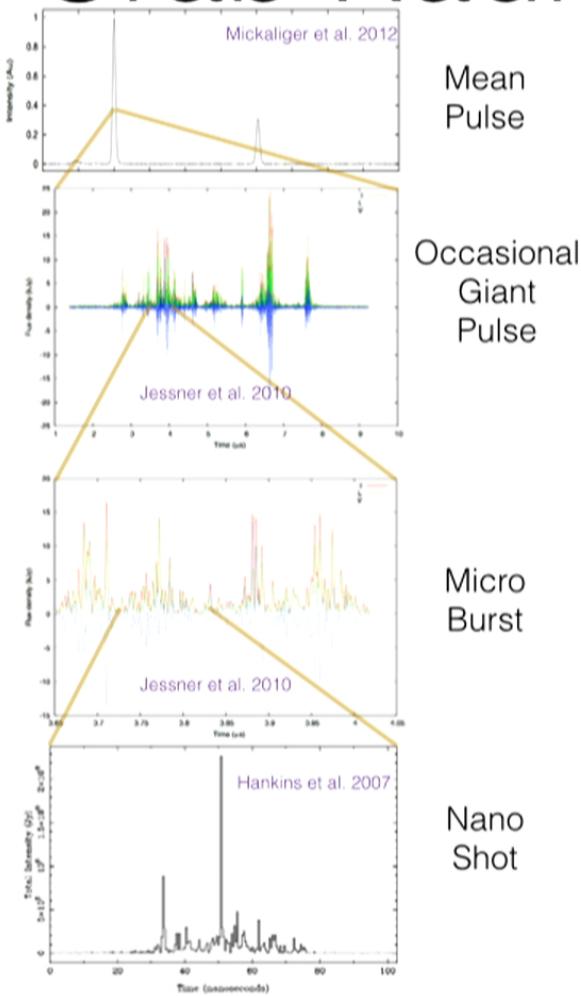
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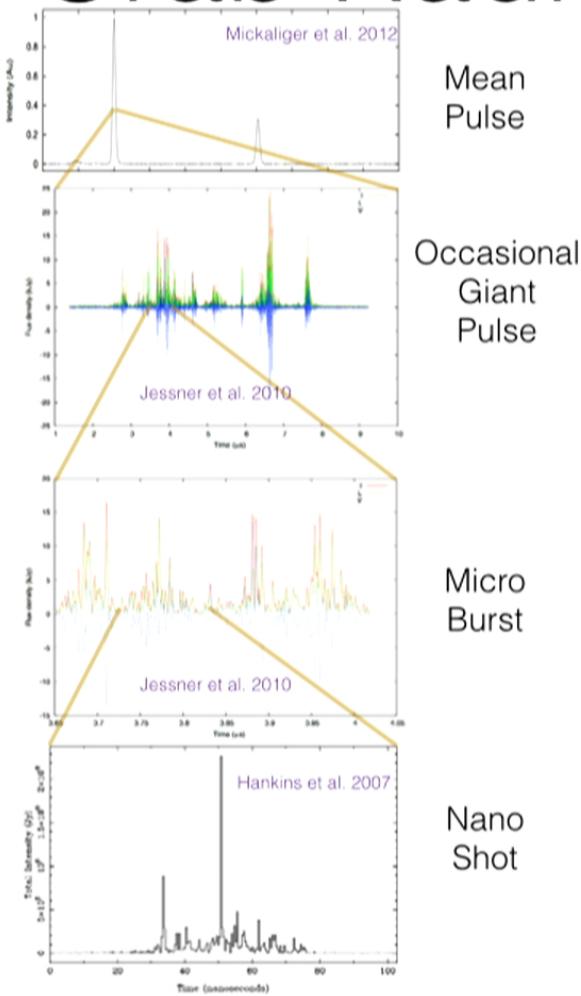
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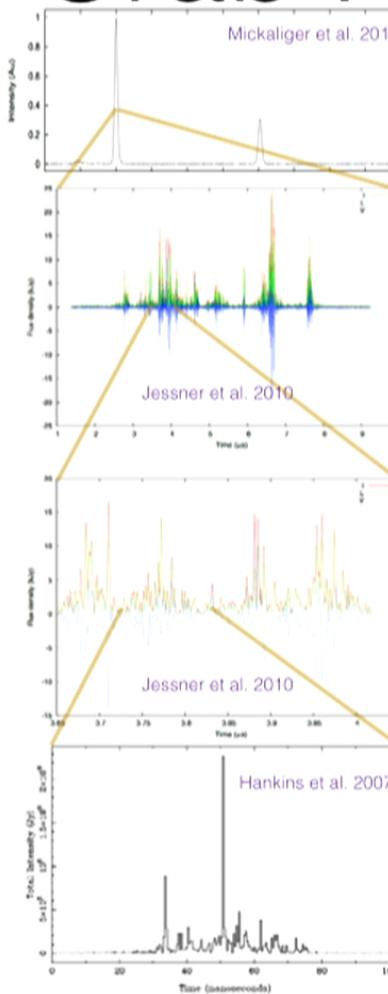
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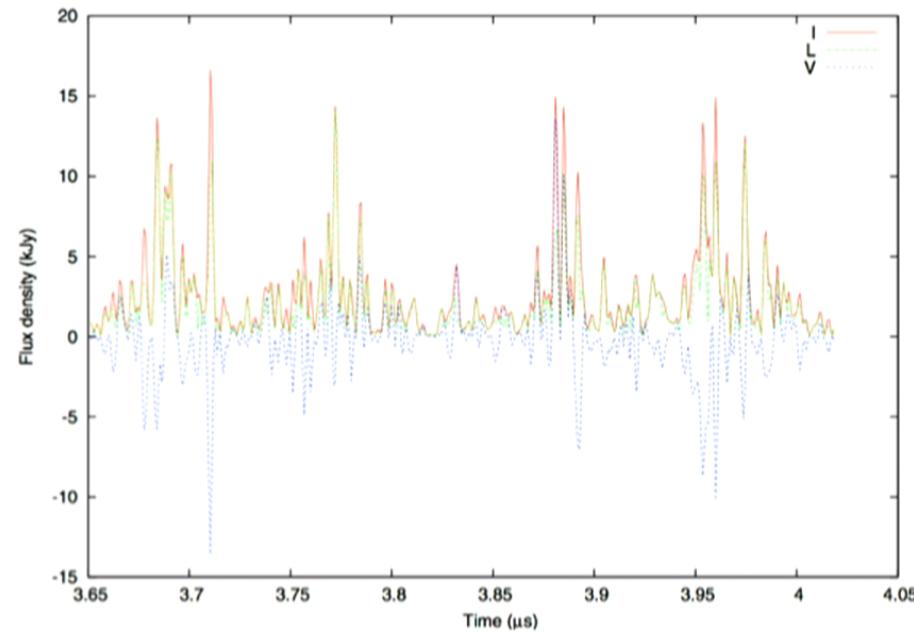


Mean
Pulse

Occasional
Giant
Pulse

Micro
Burst

Nano
Shot



This record nanoshot had peak observed flux of 2MJy, astounding for an object 2kpc away!

Crab Nanoshot Numerology

- (Assuming no significant relativistic beaming)
- size: $\leq 30\text{cm}$ [c ns]
- energy: $\leq 10^{28}\text{erg}$ [4 π D² MJy ns GHz]



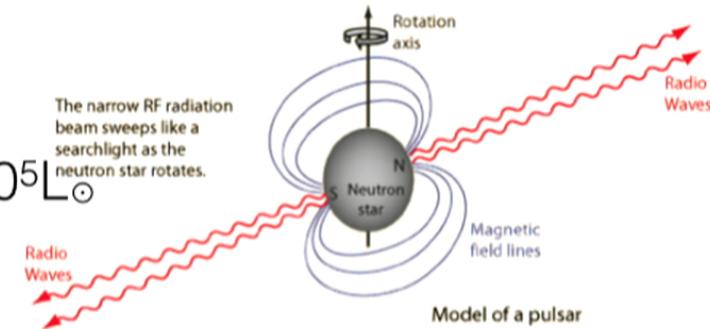
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- peak luminosity: $\leq 1.5 \times 10^3 L_\odot$ [$4\pi D^2 \text{ MJy nsec GHz}$]
- electric field: $\leq 2 \times 10^{12}\text{G}$ [$4\pi D / \text{ns} \sqrt{(\text{MJy GHz}) / c^3}$]
- charge: $\leq 5 \text{ Mole e}$ [$4\pi D \text{ ns} \sqrt{(\text{MJy GHz c})}$]
- e^\pm energy: $\leq 15 \text{ PeV}$ [$4\pi D e \sqrt{(\text{MJy GHz}) / c}$]



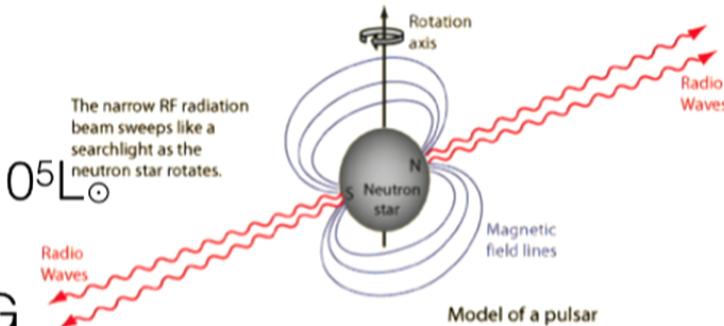
Crab Pulsar Magnetosphere Numerology

- inner radius: 10km
- spin down luminosity: $\sim 10^5 L_\odot$



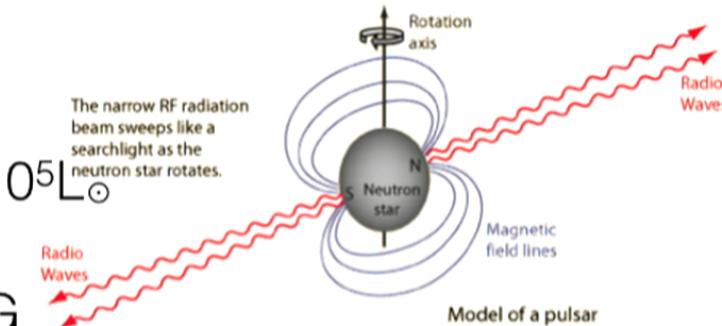
Crab Pulsar Magnetosphere Numerology

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- particle density: $\sim 0.004 \text{ mole/m}^3$



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- plasma frequency: 250 GHz at surface



Radio Transients & QuantumElectroDynamics

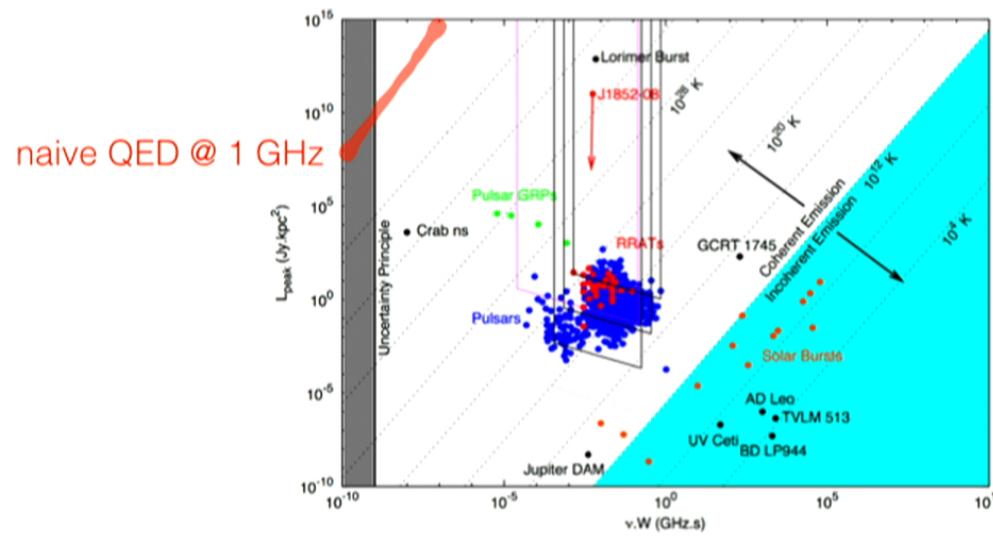
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Radio Transients & Quantum ElectroDynamics

- In QED electric or magnetic fields are said to be *strong* if $E, B \gtrsim m_e c^3 / (e\hbar)$. (in CGS units this is the *Schwinger critical field*)
- Neutron stars with magnetic fields are called *magnetars*. These exhibit interesting QED phenomena such as photon splitting.
- If we know the distance to the source one can estimate the electromagnetic field strength at emission (neglecting relativistic beaming): $E \approx \sqrt{4\pi v L_v \Delta t / (c \Delta t)}$

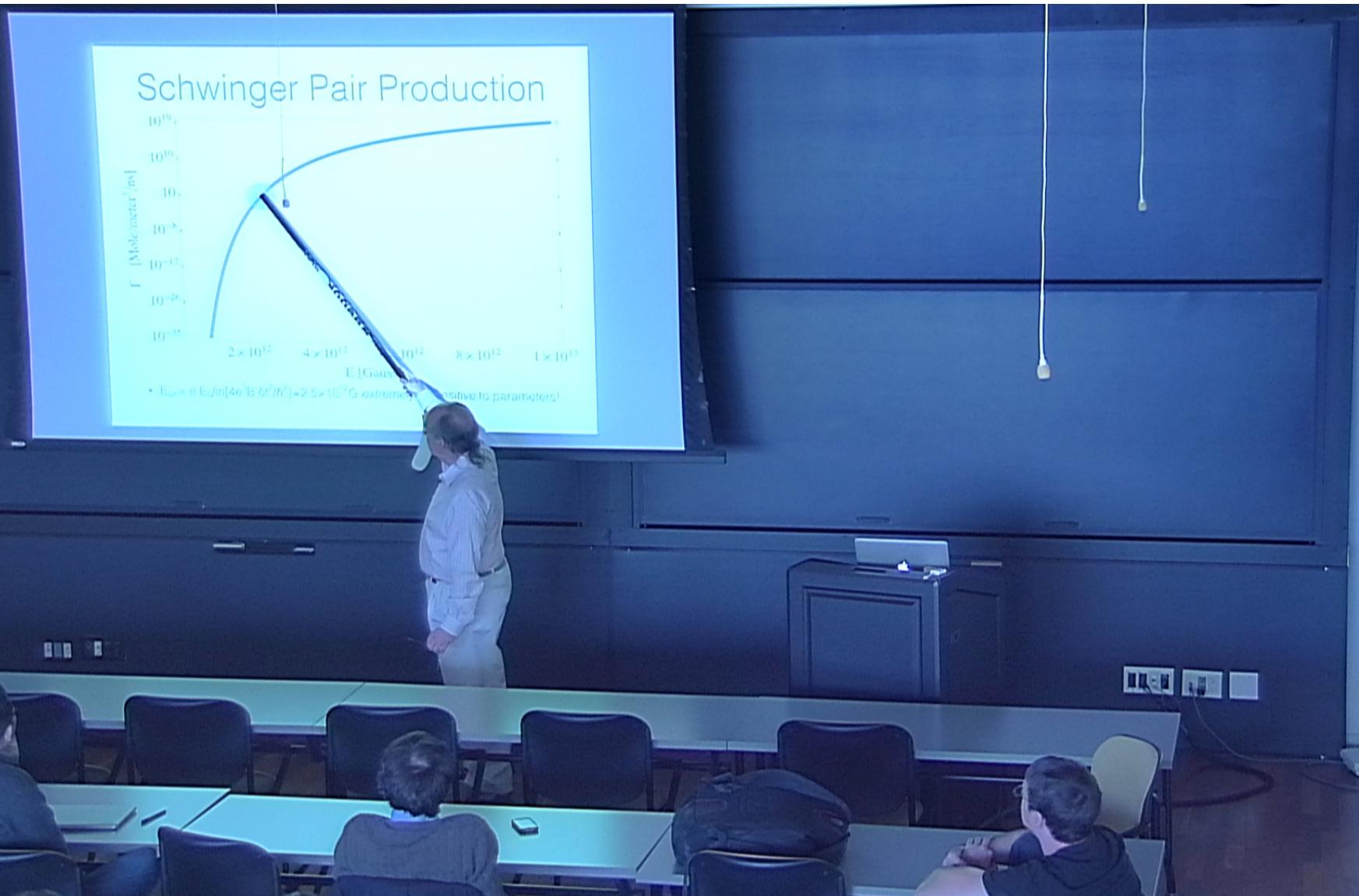


Schwinger Pair Production

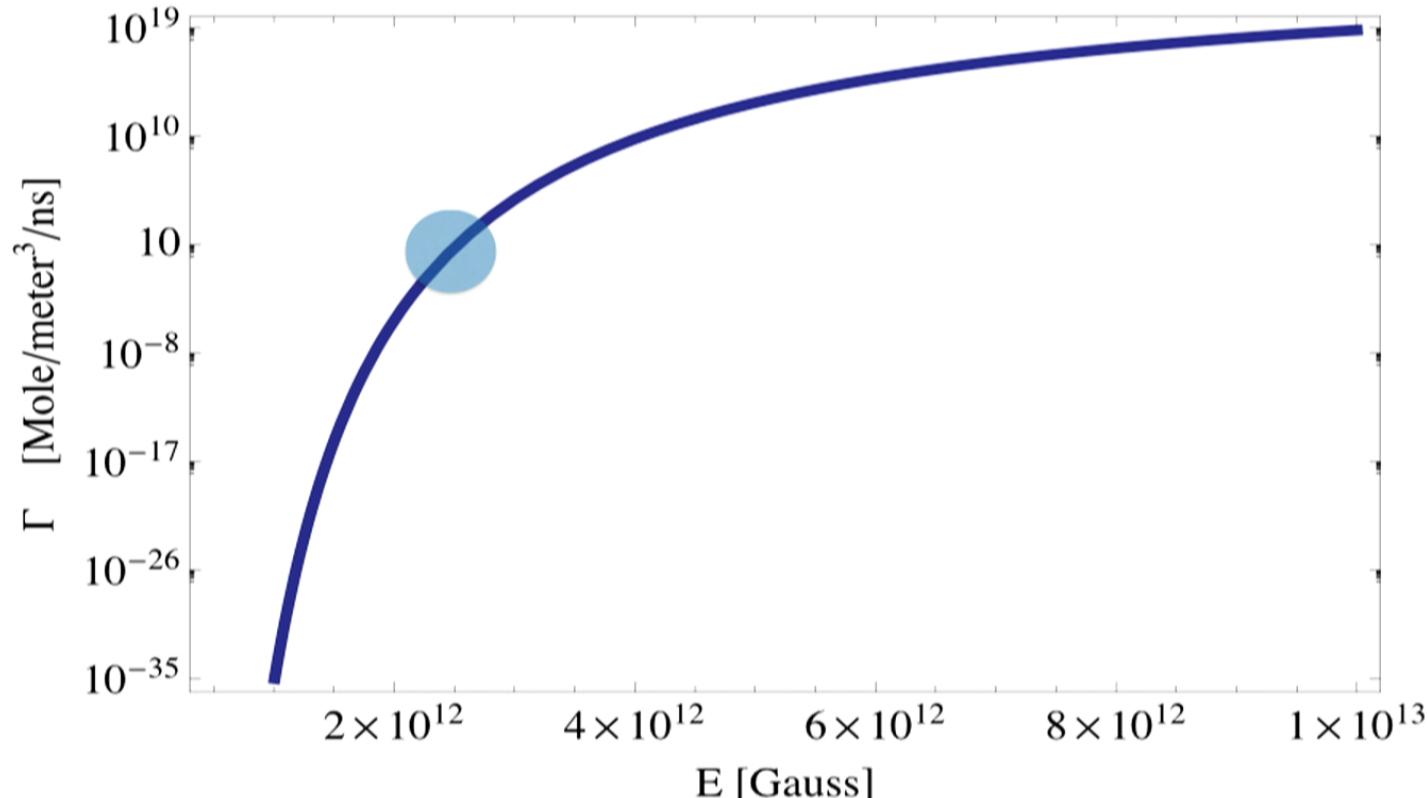
- required pair production rate is $\sim 0.2 \text{ Mole/m}^3/\text{ns}$
- In the presence of a strong electric field pairs are spontaneously produced out of the vacuum.
- in co-linear frame ($\mathbf{E} \parallel \mathbf{B}$) in terms of $E_S \equiv m_e^2 c^3 / (e\hbar) = 4 \times 10^{14} \text{ G}$

$$\Gamma = \frac{e^2}{\pi c \hbar^2} E B \sum_{n=0}^{\infty} \sum_{s=\pm 1} e^{-\pi \frac{e E_S + (2n-s+1)eB}{eE}} \cong \frac{e^2}{\pi c \hbar^2} E B e^{-\pi \frac{E_S}{E}}$$

- Note $(e E_S)^2 / (\pi c \hbar^2) = 7 \times 10^{24} \text{ Mole/m}^3/\text{ns}$ is too large! E never approaches E_S
- electric field will saturate when **shorted**
 - pair production rate exceeds rate of field increase
 - all speeds ultra-relativistic
- $E_{\lim} = \pi E_S / \ln[4e^3 B \delta t^2 / \hbar^2] \cong 2.5 \times 10^{12} \text{ G}$ extremely insensitive to parameters!



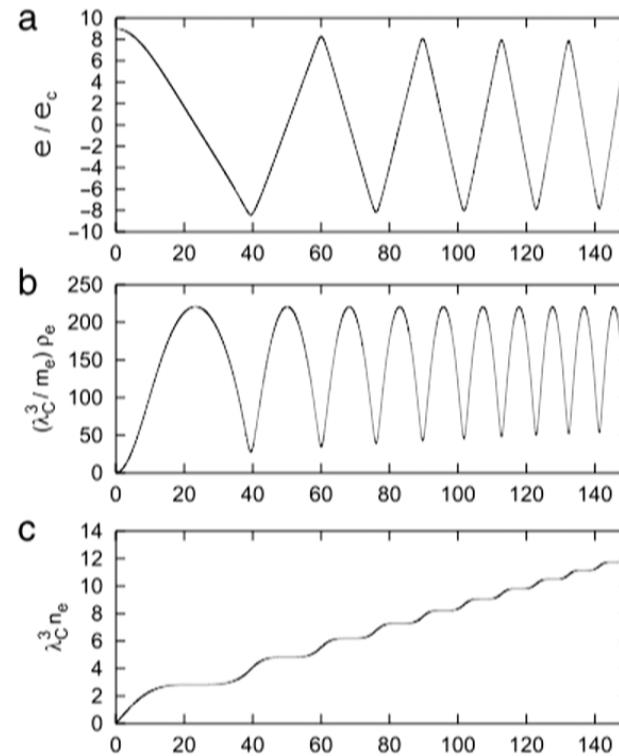
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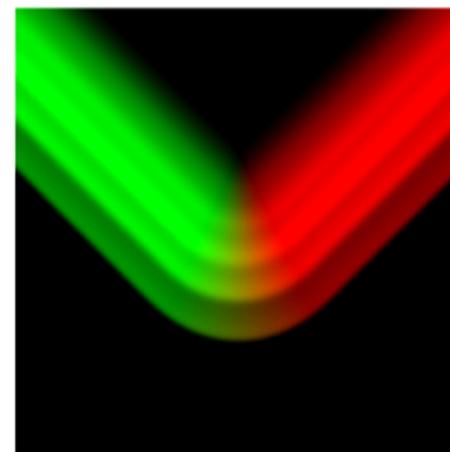
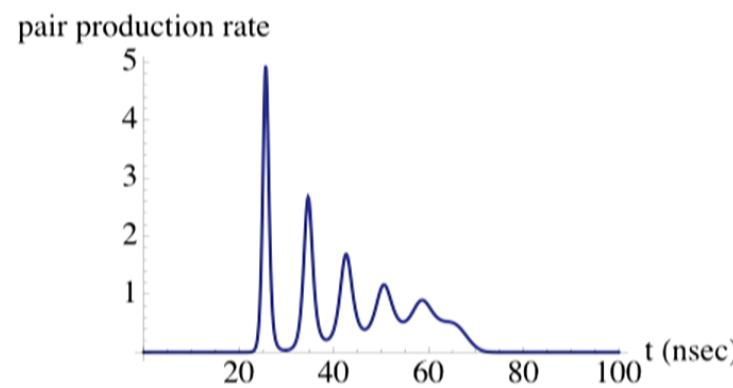
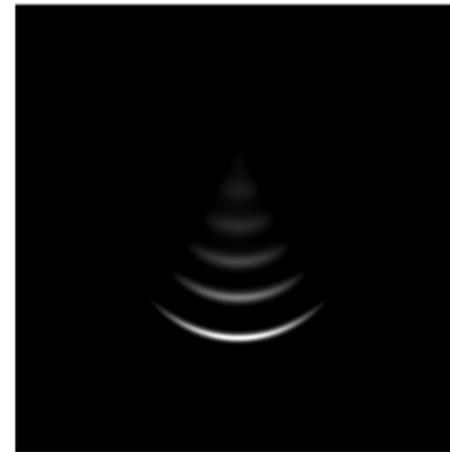
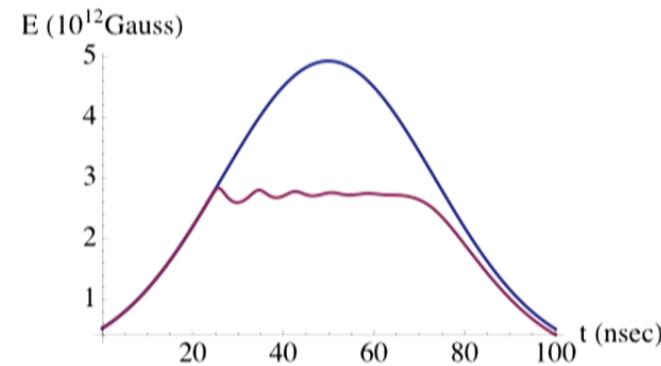
For Large Uniform Fields

- **Plasma Oscillations** (see Ruffini, Vereshchagin, Xue 2010)



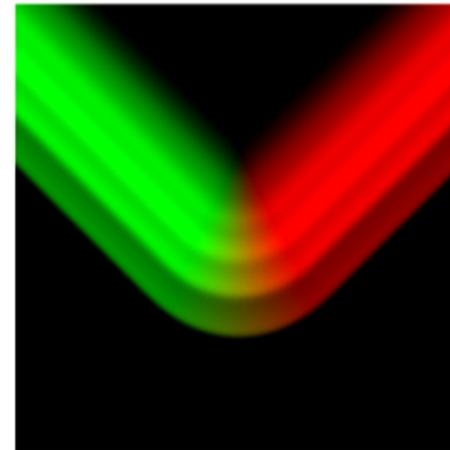
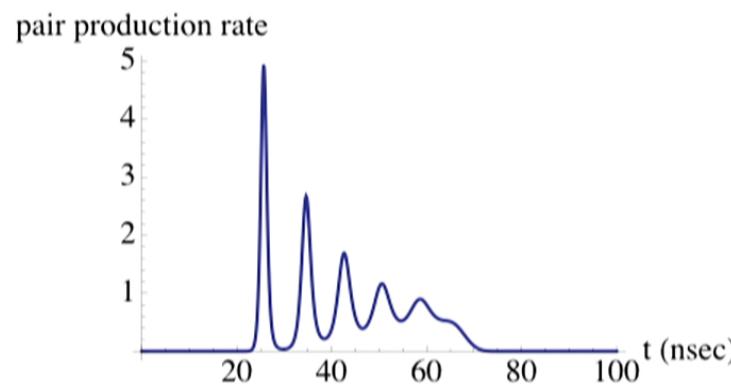
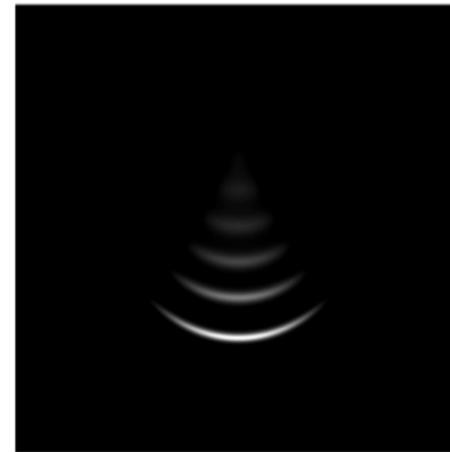
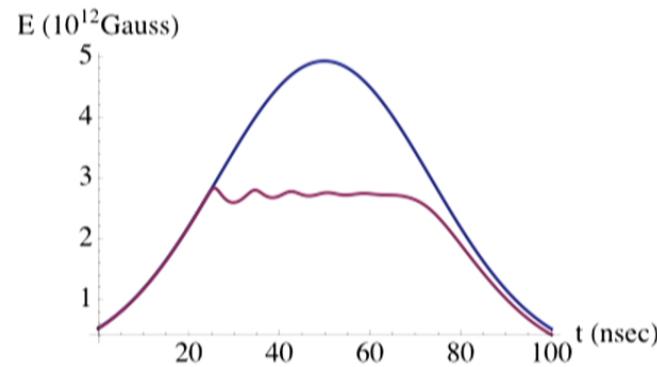
For Smaller Schwinger Sparks

1+D simulation (Stebbins, Yoo 2015)



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Schwinger Spark ElectroMagnetic Pulse

$$\dot{Q} = 2e \dot{N} \text{Sec}[\theta]^2 \hat{z} \quad \dot{N}[\mathbf{x}, t] \equiv \int d^3 \mathbf{x}' \Gamma[\mathbf{x}', t - \frac{|\mathbf{x} - \mathbf{x}'|}{c}]$$

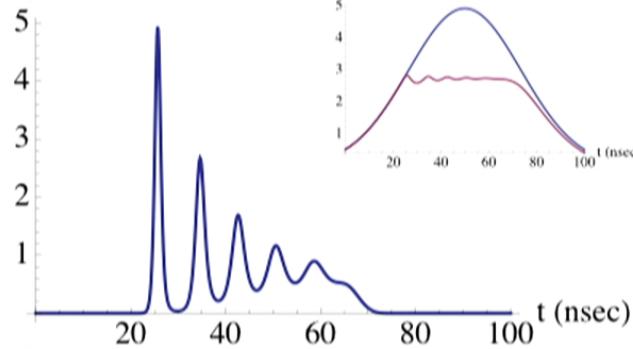
$$S[\mathbf{x}, t] \equiv \frac{c}{4\pi} (\mathbf{E} \times \mathbf{B}) = \frac{\dot{Q} \cdot \dot{Q} - (\hat{r} \cdot \dot{Q})^2}{4\pi D^2 c} \hat{r} = \frac{(2e \dot{N})^2}{4\pi D^2 c} \text{Sec}[\theta]^2 \hat{r}$$

$$F_\nu[\mathbf{x}] \equiv \int_{-\infty}^{\infty} dt S_\nu[\mathbf{x}, t] = \frac{(2e)^2}{4\pi D^2 c} \left| \int_{-\infty}^{\infty} dt e^{-i2\pi\nu t} \dot{N}[\mathbf{x}, t] \right|^2 \text{Sec}[\theta]^2$$

$$F_\nu \rightarrow F_0 = \frac{(2e N \text{Sec}[\theta])^2}{4\pi c D^2} \quad N \equiv \int_{-\infty}^{\infty} dt \int d^3 \mathbf{x} \Gamma[\mathbf{x}, t]$$

$$\mathcal{E}_P = \left(\frac{g_P}{2\pi}\right)^2 (c \delta t_P)^3 E_P^{\text{lim}}{}^2 = 2 \times 10^3 L_\odot \text{nsec} \left(\frac{\delta t_P}{\text{nsec}}\right)^3 \left(\frac{E_P^{\text{lim}}}{2.8 \times 10^{12} \text{Gauss}}\right)^2$$

pair production rate



Stebbins, Yoo 2015

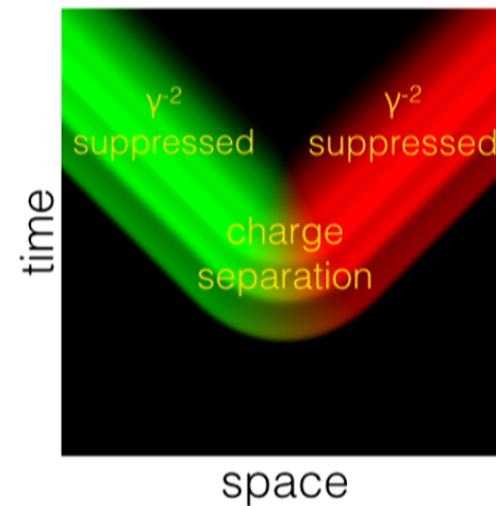
If created pairs are rapidly accelerated to speed of light the EMP just traces the pair production rate.

EMP is "charge separation" radiation

EMP spectrum depends on shape of electric field.

Schwinger sparks up-converts electro-magnetic energy.

Total energy scales with cube of pulse length: v_{max}^{-3} .



Schwinger Spark ElectroMagnetic Pulse

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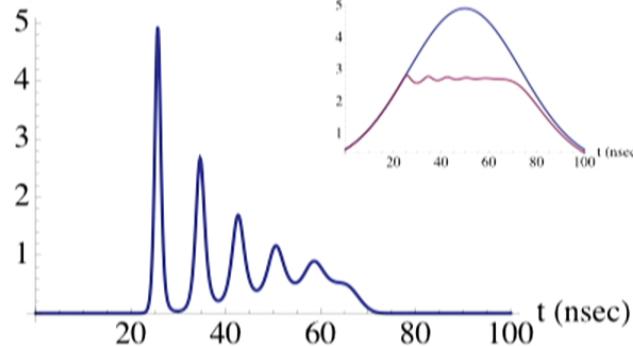
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$$F_\nu \rightarrow F_0 = \frac{(2e N \text{Sec}[\theta])^2}{4\pi c D^2} \quad N \equiv \int_{-\infty}^{\infty} dt \int d^3 \mathbf{x} \Gamma[\mathbf{x}, t]$$

$$\mathcal{E}_P = \left(\frac{g_P}{2\pi}\right)^2 (c \delta t_P)^3 E_P^{\text{lim}}{}^2 = 2 \times 10^3 L_\odot \text{nsec} \left(\frac{\delta t_P}{\text{nsec}}\right)^3 \left(\frac{E_P^{\text{lim}}}{2.8 \times 10^{12} \text{Gauss}}\right)^2$$

pair production rate



Stebbins, Yoo 2015

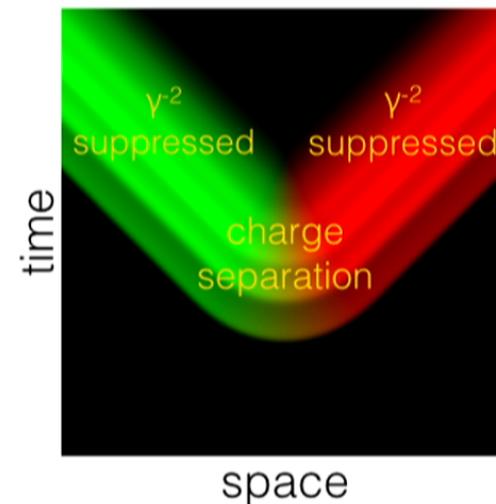
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EMP is "charge separation" radiation

EMP spectrum depends on shape of electric field.

Schwinger sparks up-converts electro-magnetic energy.

Total energy scales with cube of pulse length: v_{max}^{-3} .



Schwinger Spark ElectroMagnetic Pulse

$$\dot{Q} = 2e \dot{N} \text{Sec}[\theta]^2 \hat{z} \quad \dot{N}[\mathbf{x}, t] \equiv \int d^3 \mathbf{x}' \Gamma[\mathbf{x}', t - \frac{|\mathbf{x} - \mathbf{x}'|}{c}]$$

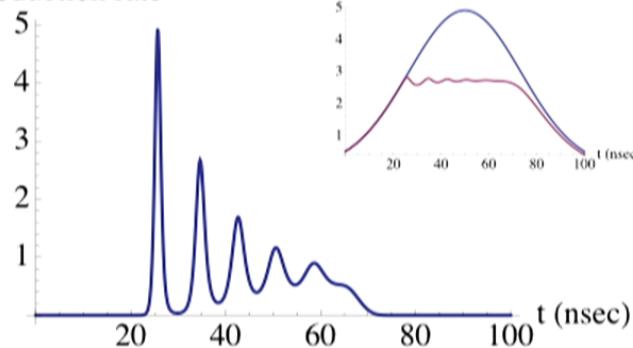
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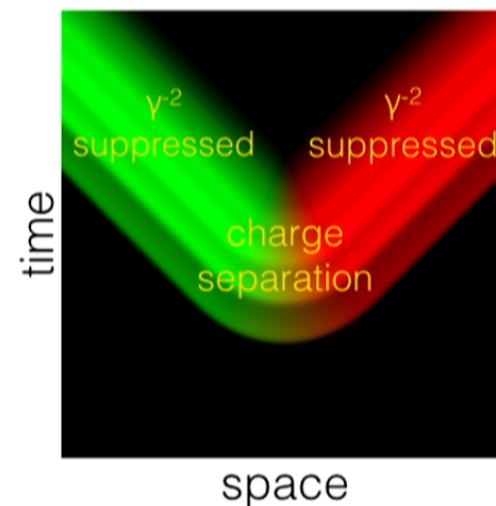
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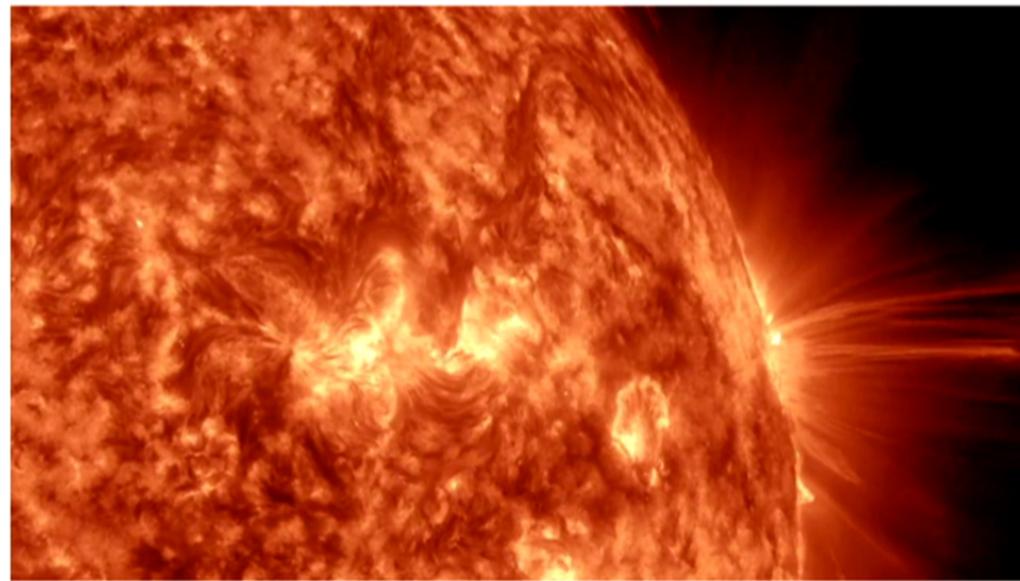
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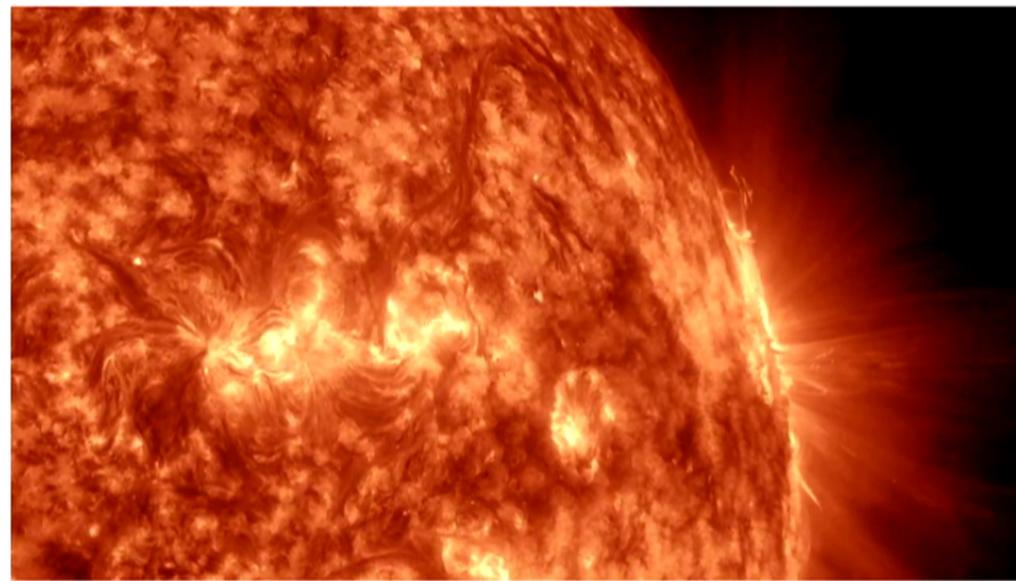
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- EMP energy density higher than that of magnetosphere!
- will accelerate any e^\pm in path to ultra-relativistic energies

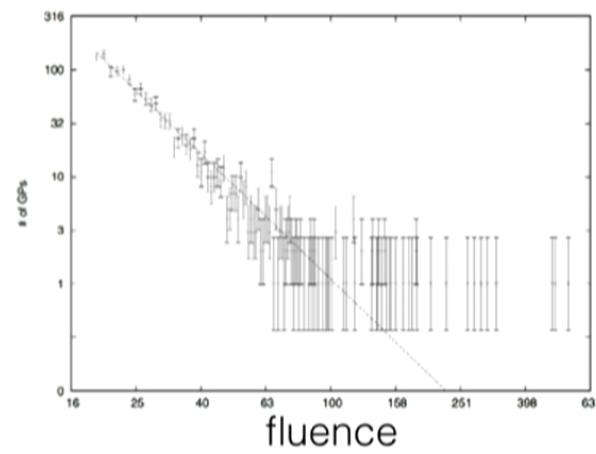
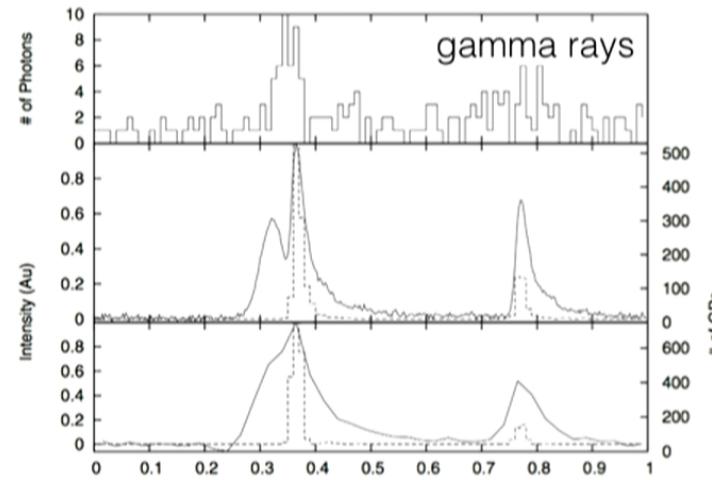
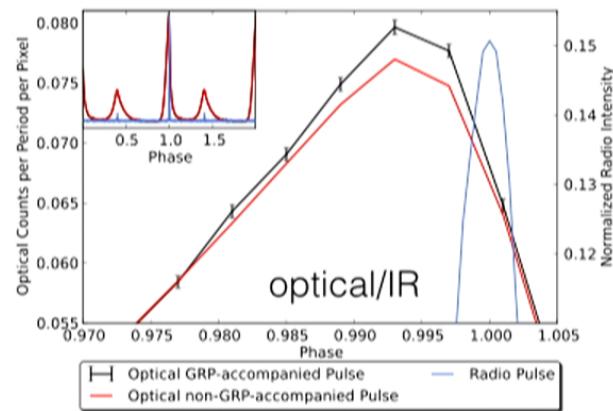
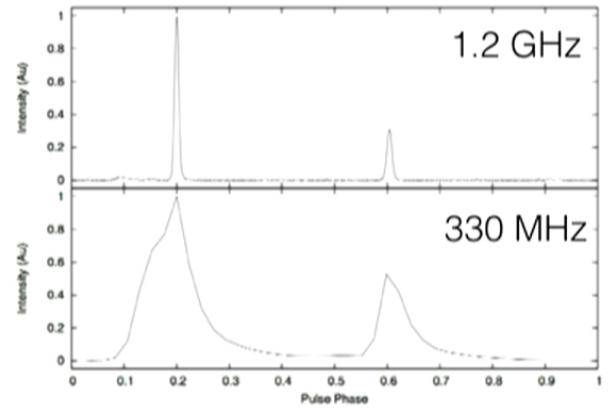


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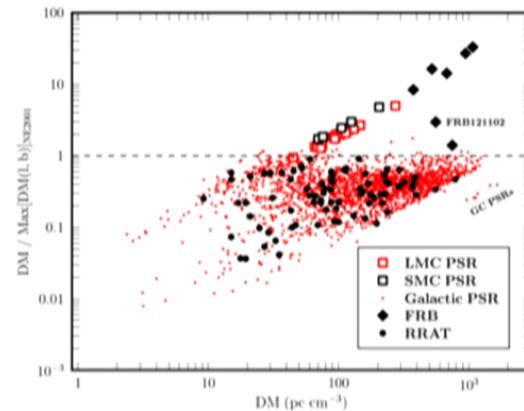
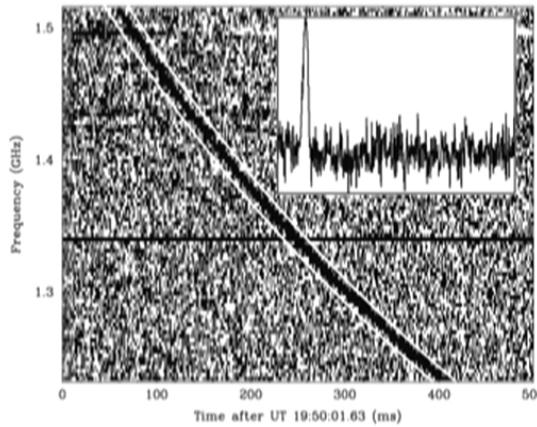
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Crab Giant Pulse Phenomenology



Fast Radio Bursts: Cosmological Dispersion Measure?



$$\Delta t_{\text{delay}} = \frac{e^2}{2\pi m_e c^3} (\lambda^2 \text{DM})$$

$$= 1.5 \times 10^{-24} \text{ s} (\lambda^2 \text{DM})$$

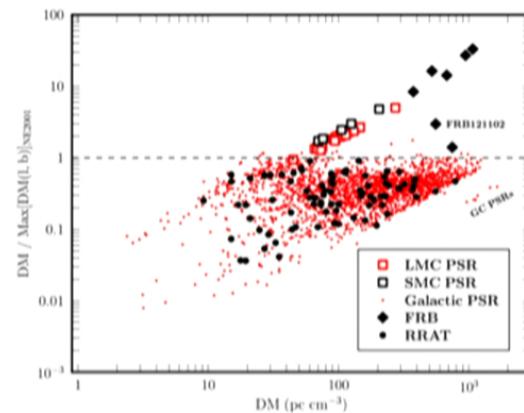
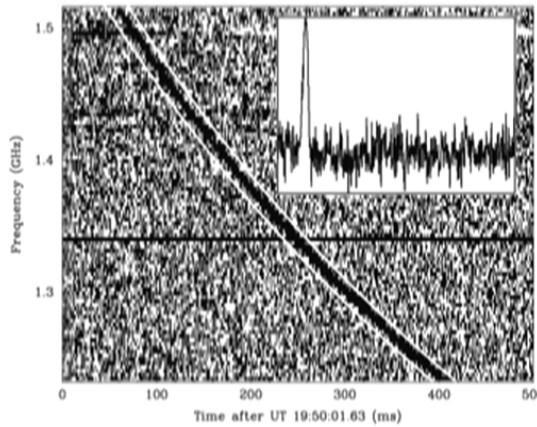
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$$\delta t_{\text{width}} = 4.6 \text{ ms} \left(\frac{\nu}{1.4 \text{ GHz}} \right)^{-4.8 \pm 0.4}$$

$$\int dt I_\nu \simeq 150 \pm 50 \text{ Jy ms} @ 1.4 \text{ GHz}$$

	DM_{IGM}	redshift	$\text{energy}_{\text{GHz}}$	Δt
FRB0102	375- 25-100	~ 0.3	$\sim 10^{40}$ erg	5ms
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FRB1102	944- 34-100	~ 0.81	$\sim 10^{40}$ erg	5.6ms
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FRBs: what are they? reasoning by rates

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$\Rightarrow 10^4/\text{sky/day} \quad (3 \text{ Jy ms}/F)^{(3/2)}$ (assume Euclidean)

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MWEG = Milky Way Equivalent Galaxy

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extragalactic: subclass of CCs (SURON/Blitzar), nearly all CBCs, magnetar masers

galactic: flaring stars, erratic pulsars, white dwarf mergers

Loeb et al. 2013

Kashiyama et al. 2013

Lyubarsky 2014

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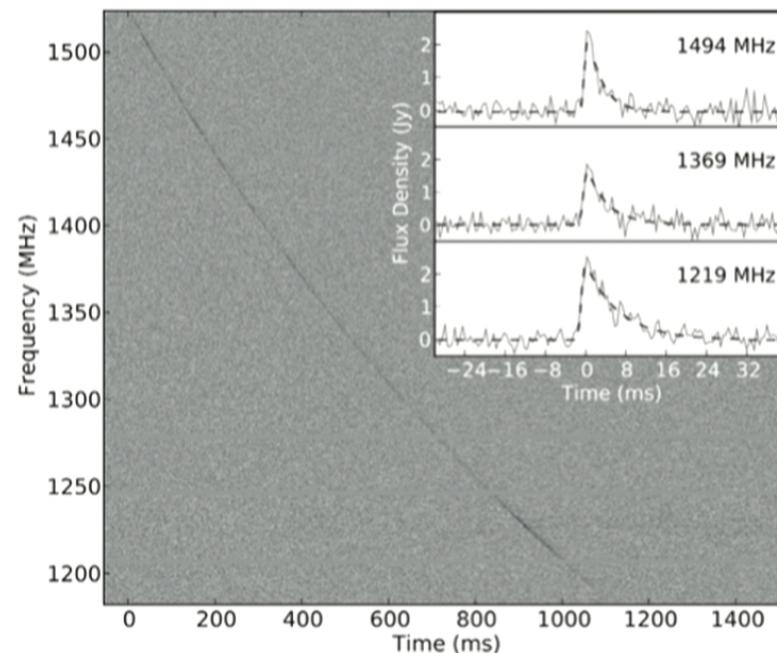
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solar system: colliding mini-asteroids

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Fast Radio Bursts (FRBs)



Lorimer et al. 2007 (1)
Keane et al. 2011 (1)
Thornton et al. 2013 (4)
Spitler et al. 2014 (1)
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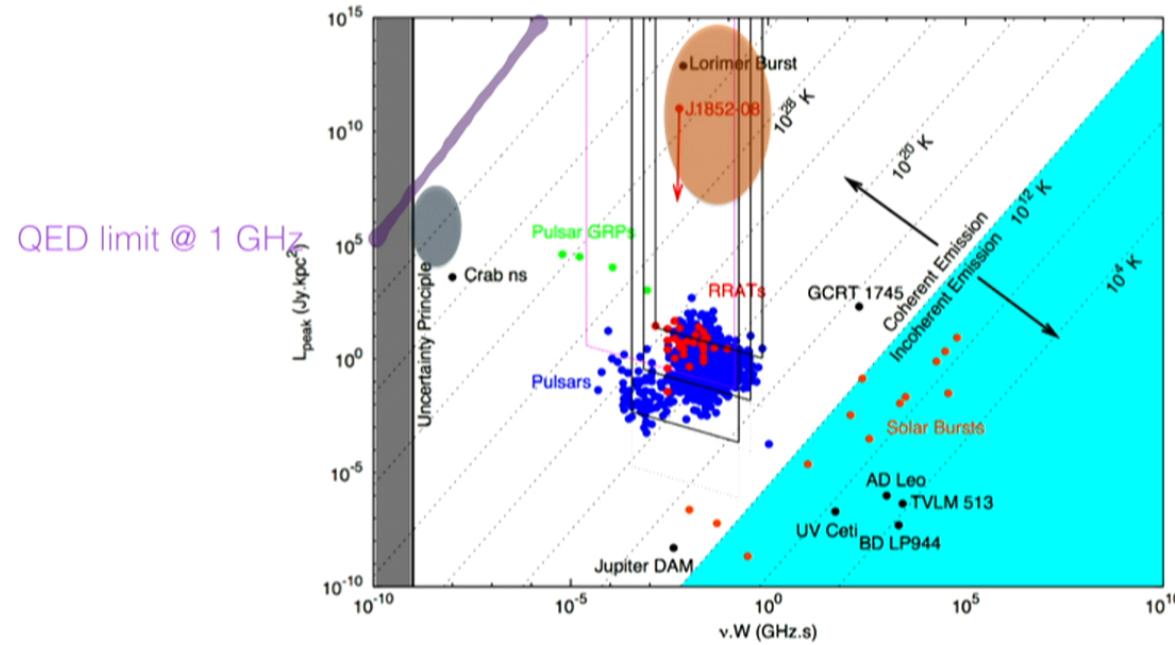
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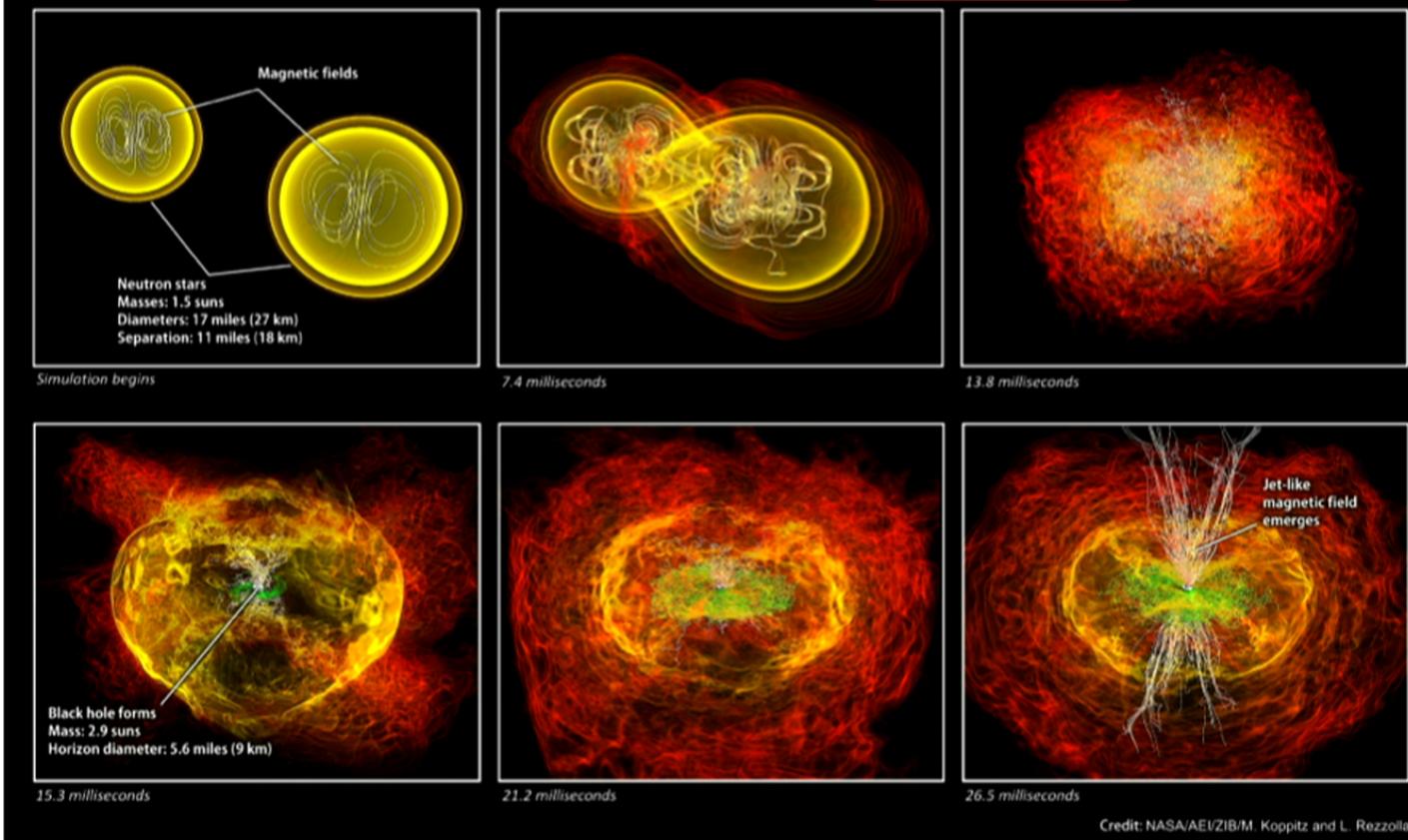
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Schwinger Model of FRBs

- **Blitzar:** ~billion Swinger Sparks within a few msec.

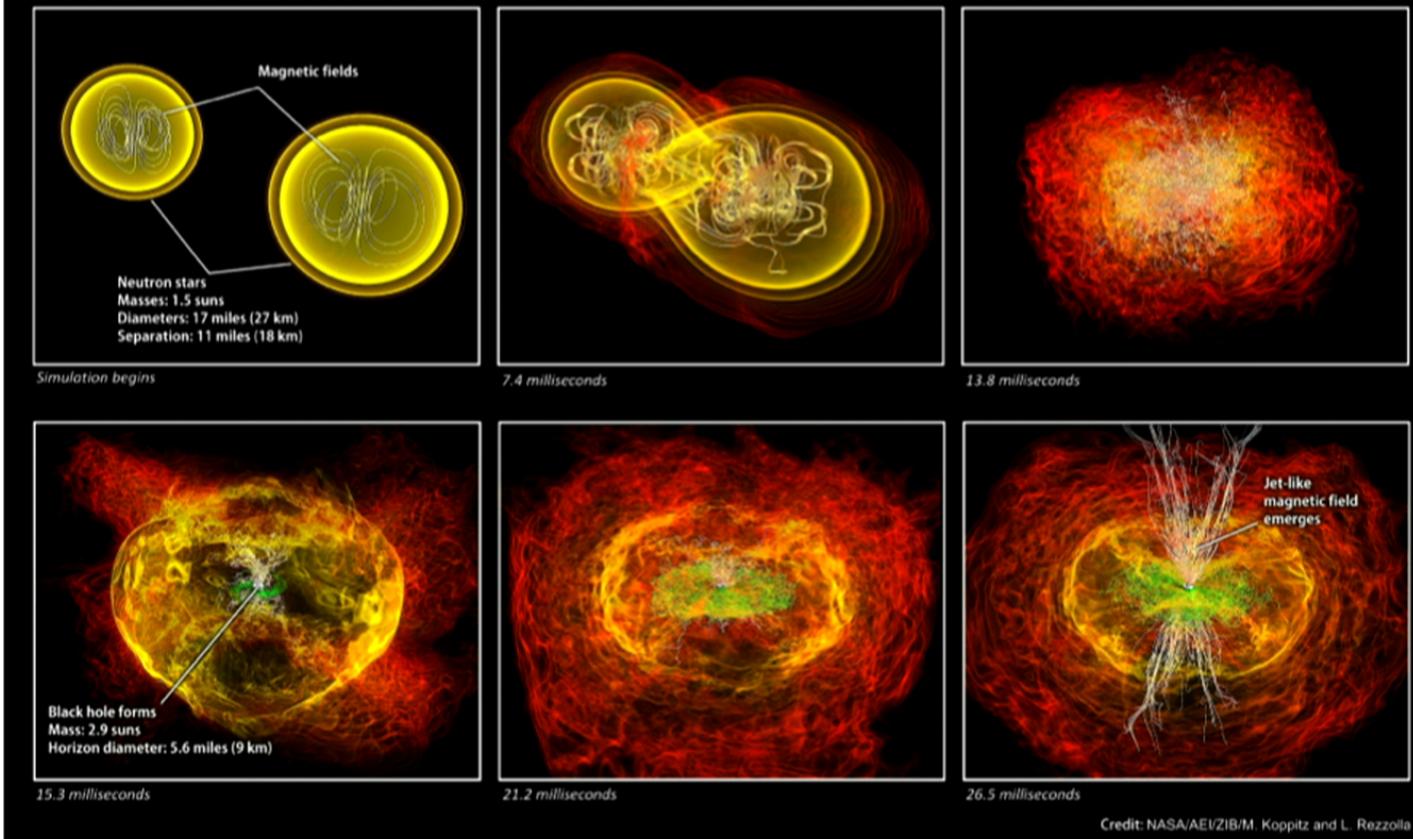


Crashing neutron stars can make **radio** burst jets?



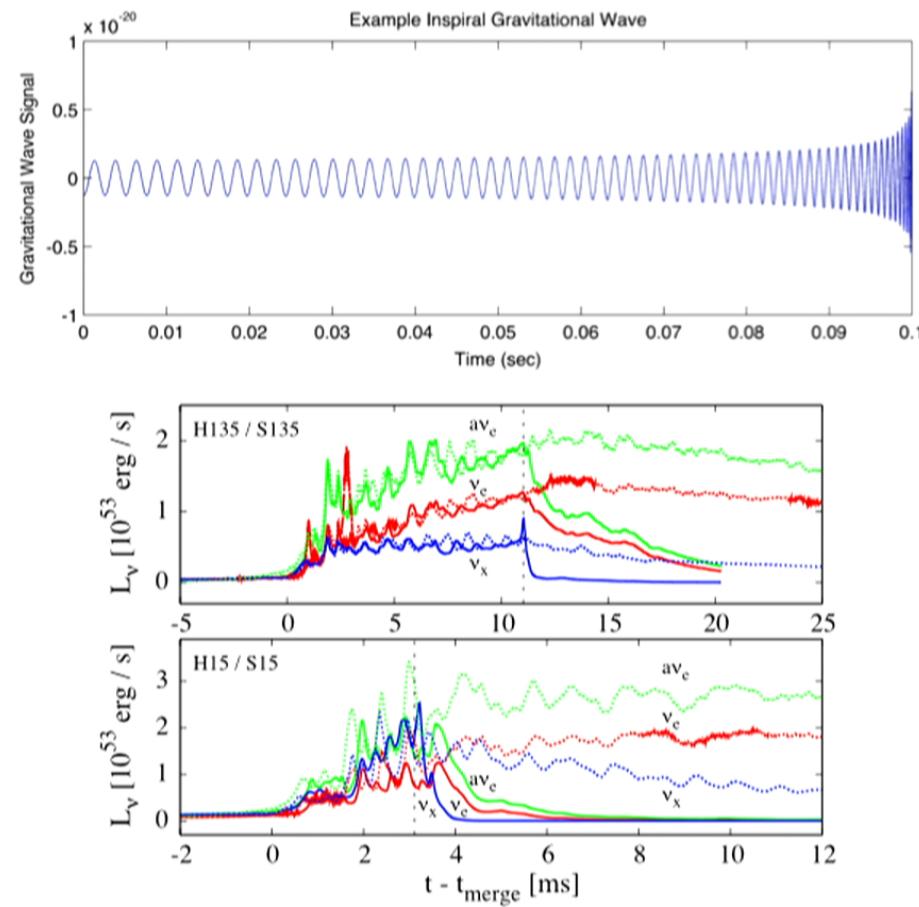
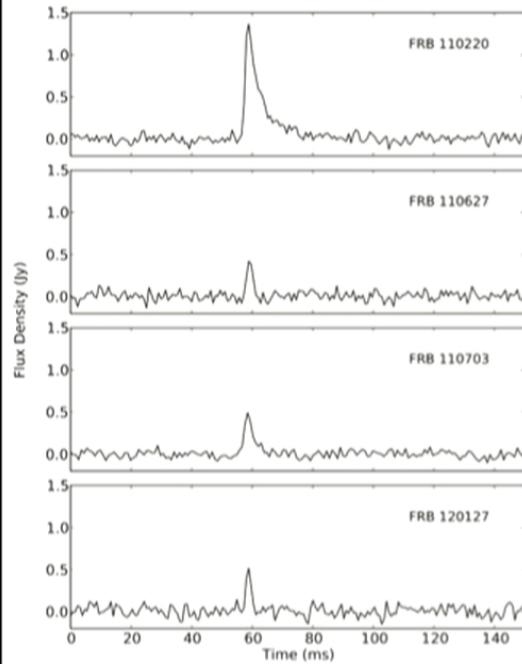
Rapid turn off due to increasing opacity (plasma oscillations)?

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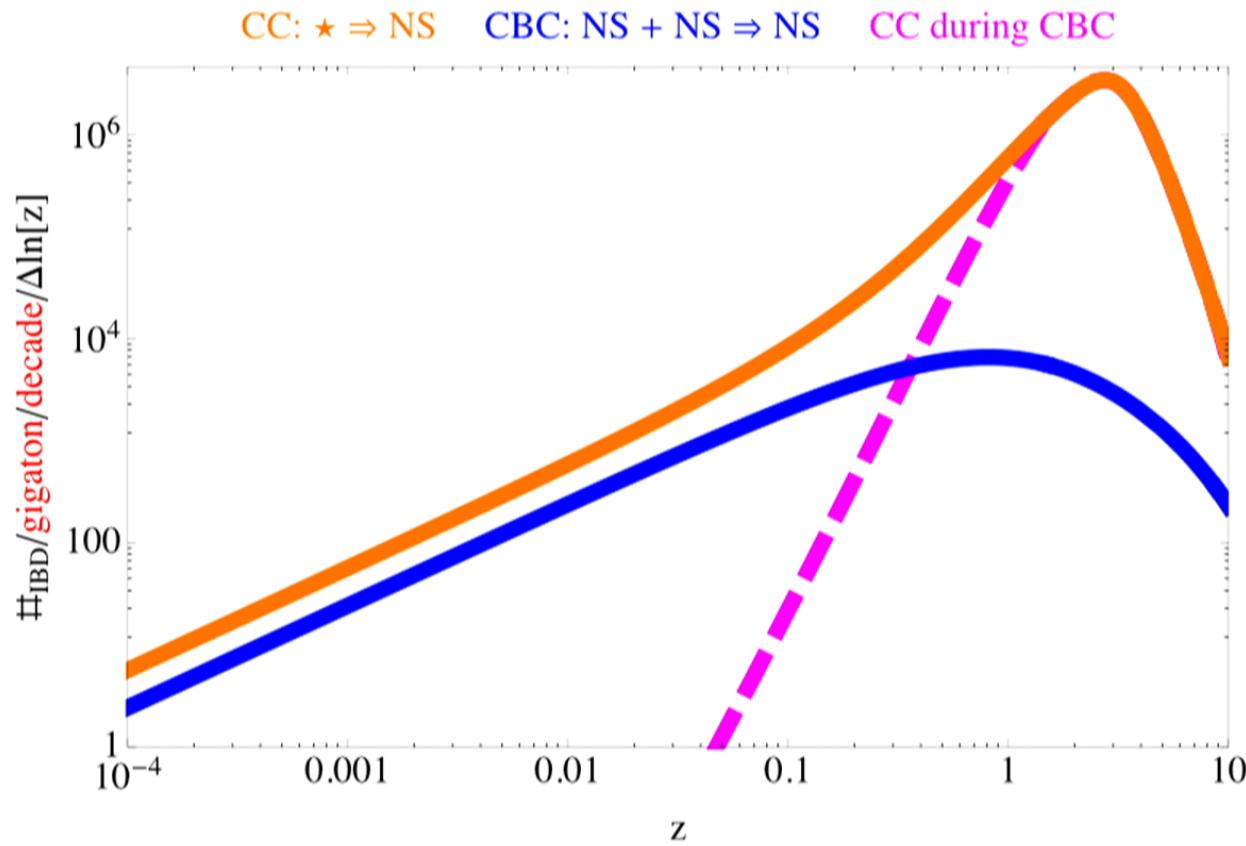
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Coincidence of Timescales for NS+NS \Rightarrow NS CBCs



Kiuchi et al. 2012

FRB Associated Neutrinos



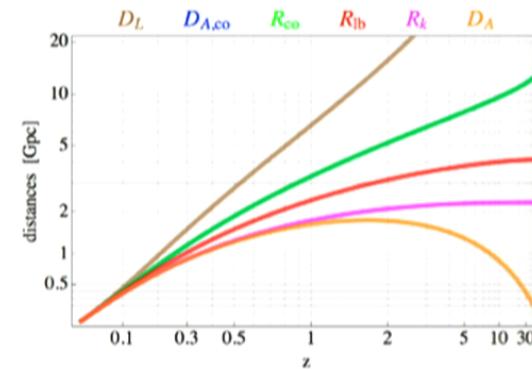
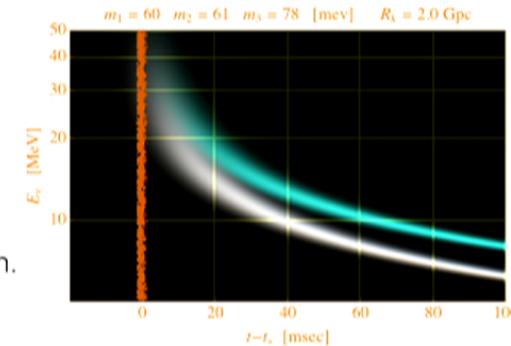
With FRB triggers “negligible” backgrounds for $z \lesssim 0.2$

What to do with 1000s of FRB v's?

- Learn about FRB progenitors.
- Measure the Hubble Constant
 - massive v time of arrival energy dependent
 - we don't know absolute mass scale of v's but we might soon.

$$\Delta t = \frac{1}{2} \left(\frac{m}{E} \right)^2 \frac{R_k}{c} = 10 \text{ sec} \left(\frac{m}{0.1 \text{ eV}} \frac{10 \text{ MeV}}{E} \right)^2 \frac{R_k}{2 \text{ Gpc}}$$

- N.B. For small z since $\Delta t \propto R$ while $\#_v \propto R^{-2}$ the Fisher information for H_0 is contributed equally for each FRB.
- A gigaton detector operating for a decade could in principle yield % level determination of H_0 .



$$R_{lb}[z] = c \int_0^z \frac{dz'}{(1+z')H[z']}$$

$$D_{A,co}[z] = \begin{cases} \frac{\sin[\sqrt{-K} R_{co}[z]]}{\sqrt{-K}} & K > 0 \\ R_{co}[z] & K = 0 \\ \frac{\sinh[\sqrt{-K} R_{co}[z]]}{\sqrt{-K}} & K < 0 \end{cases}$$

$$D_A[z] = \frac{1}{1+z} D_{A,co}[z]$$

$$R_{co}[z] = c \int_0^z \frac{dz'}{H[z']}$$

$$D_L[z] = (1+z) D_{A,co}[z]$$

$$R_k[z] = c \int_0^z \frac{dz'}{(1+z')^2 H[z']}$$

FRB Triggers Increase Reach of GW Telescopes

- It is normally thought that gravitational radiation detectors may provide triggers for other observations but the reverse may be true in the case of FRBs.
- For triggers it is quantity that matters.
 - FRBs occur 10 to 100 times more frequently than GRBs and are shorter.
 - While core collapses SNe occur more frequently than FRBs the time of core collapse can at best be determined to ~ 1 day accuracy and one is in the confusion limit.
- With 10^6 triggers one increases $\langle h^2 \rangle / \langle n^2 \rangle$ of incoherent signal by 10^3 , bringing effective S/N distance of co-added $z \sim 0.7$ events down to $z \sim 0.02$ or 100Mpc.
 - if there were a component to the gravitational radiation waveform correlated with the radio burst trigger one could decrease this much more: down to $z \sim 0.0007$ or 3 Mpc.

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- A heretofore unobserved QED phenomena, the Schwinger mechanism, may play important role in producing some of these phenomena.

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- We propose Schwinger sparks as the emission mechanism of both Crab nanoshots and Fast Radio Bursts.
 - neutron star senility: young neutron star behavior re-exhibited as an adult