

Title: Foundations of Quantum Mechanics-13

Date: Jan 20, 2015 11:30 AM

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Abstract:

State Space

fixed states $K-1 = N^2 - 1$

any n normalized states $K-1 = N^2 - 2$

In convex sets

$$P_A, P_B$$

or also ham

$$P = \lambda P_A + (1-\lambda) P_B$$

$$\text{for } 0 \leq \lambda \leq 1$$

State Space

states $K-1 = N^2 - 1$

nonnormalized states $K-1 = N^2 - 2$

In convex sets

$$P_A, P_B$$

then also have

$$P = \lambda P_A + (1-\lambda) P_B$$

$$\text{for } 0 \leq \lambda \leq 1$$

for extremal (pure) states P
cannot write

$$P = \lambda P_A + (1-\lambda) P_B \quad \text{for distinct } P_A, P_B$$

and $0 < \lambda < 1$



Shape of Quantum State Space

$$K = N^2$$

Normalized states

$$K-1 = N^2 - 1$$

Boundary of normalized states

$$K-1 = N^2 - 2$$

In QT we
dimension

that the pure states have

$$|\psi\rangle = \sum_{n=1}^N a_n |a_n\rangle$$

In convex sets

p_A, p_B

then also have

$$p = \lambda p_A + (1-\lambda)p_B$$

for $0 \leq \lambda \leq 1$

for external (pure) states

cannot write

$$p = \lambda p_A + (1-\lambda)p_B$$

✓
In QT we also know that the pure states have
dimension $L = 2N - 2$ $|\psi\rangle = \sum_{n=1}^N a_n |a_n\rangle$

If $N=2$ $K-2 = 2 = L$

Bloch
ball



In QT we also know that the pure states have dimension $L = 2N - 2$

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If $N=2$ $K-2 = 2 = L$

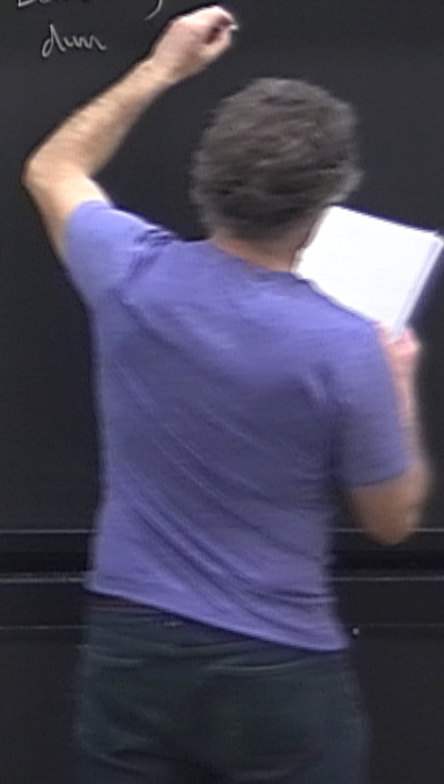
Bloch
ball



all states on boundary
are pure.

$N > 2$

Boundary
dim



Boundary dimension $k-2 = N^2 - 2$
states

In QT we also know that the pure states have
 dimension $L = 2N - 2$ $|\psi\rangle = \sum_{n=1}^N a_n |a_n\rangle$

for $0 \leq \lambda \leq 1$
 for extremal (pure) states ρ
 cannot write $\rho = \lambda \rho_A + (1-\lambda) \rho_B$ for distinct ρ_A, ρ_B

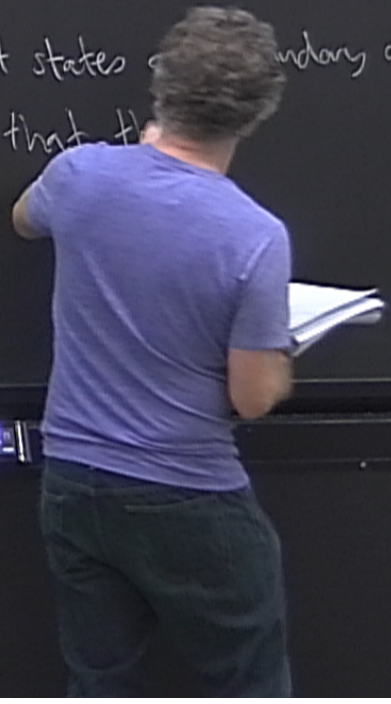
$N=2$ $k-2 = 2 = L$



all states on boundary are pure.

$N > 2$ Boundary dim $k-2 > L$ pure state dim.

Hence most states on boundary are mixed.
 Also know that the



Boundary dimension $k-2 = N^2-2$
states

also know that the pure states have

$$L = 2N-2 \quad |\psi\rangle = \sum_{n=1}^N a_n |a_n\rangle$$

for $0 \leq \eta \leq 1$

for external (pure) states \neq
cannot write

$$\neq = \eta |p_A\rangle + (1-\eta) |p_B\rangle \quad \text{for distinct } |p_A\rangle, |p_B\rangle$$

and $0 < \eta < 1$

$$k-2 = 2 = L$$

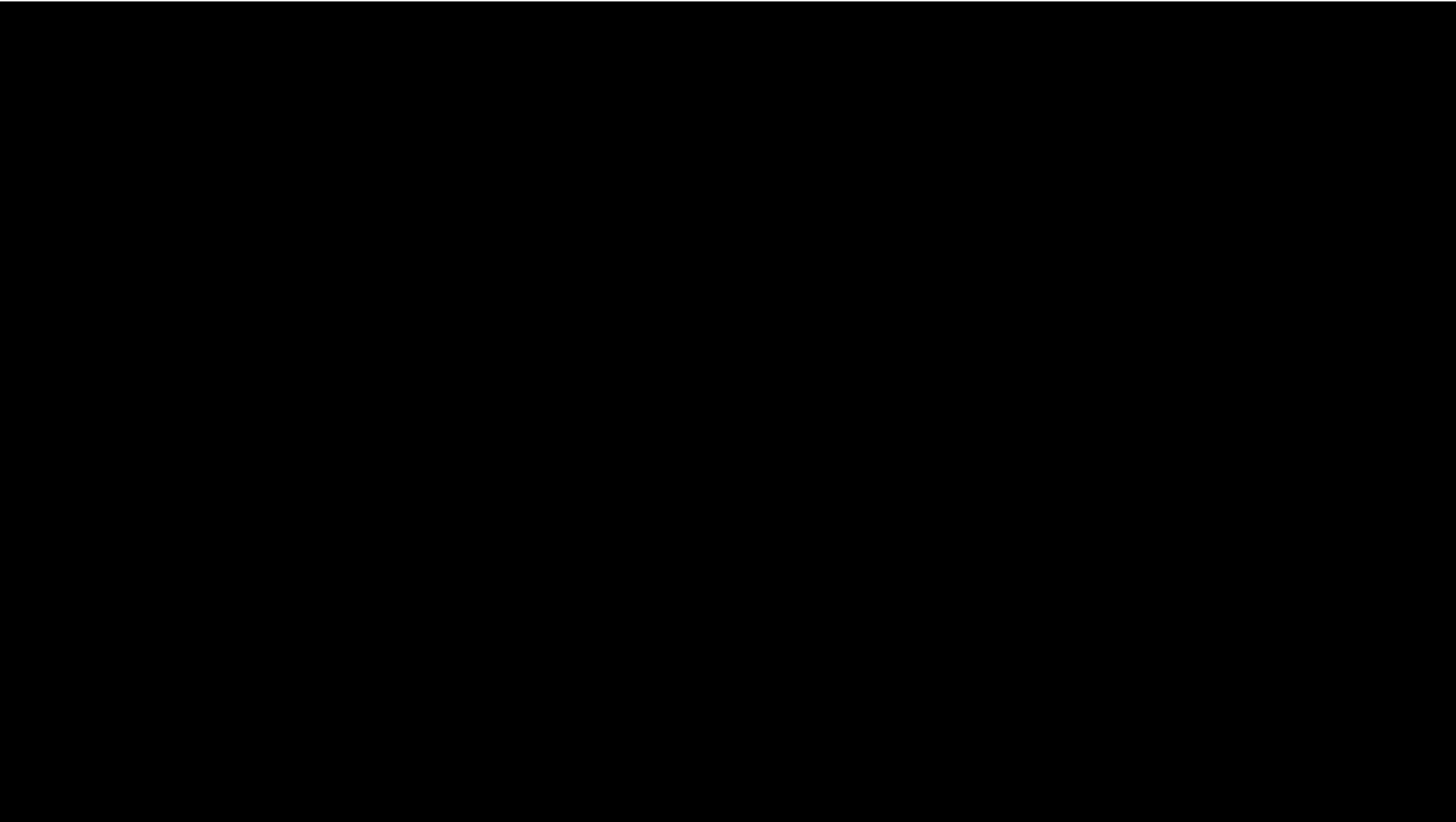
all states on boundary
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Boundary
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dim.

Hence most states on boundary are mixed.

Also know that the pure state space is closed (no edges)



Boundary of normalized states $k-2 = N^2-2$

that the pure states have

$$|\psi\rangle = \sum_{n=1}^N a_n |a_n\rangle$$

for $0 \leq \lambda \leq 1$

for extremal (pure) states ρ
cannot write

$$\rho = \lambda \rho_A + (1-\lambda) \rho_B$$

for distinct ρ_A, ρ_B
and $0 < \lambda < 1$

$$2 = L$$

on boundary
etc.

$N > 2$ Boundary dim $k-2 > L$ pure state dim.

Hence most states on boundary are mixed.

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Pringle (Ben)



Boundary of normalized states $k-2 = N^2-2$

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Boundary dimension $k-2 = N^2-2$
states

In QT we also know that the pure states have
dimension $L = 2N-2$ $|\psi\rangle = \sum_{n=1}^N a_n |a_n\rangle$

for $0 \leq \lambda \leq 1$
for extremal (pure) states ρ
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 $\rho = \lambda \rho_A + (1-\lambda) \rho_B$ for distinct
and

$N=2$ $k-2 = 2 = L$



all states on boundary
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Boundary of normalized states $k-2 = N^2 - 2$

In QT we also know that the pure states have dimension $L = 2N - 2$

$$|\psi\rangle = \sum_{n=1}^N a_n |a_n\rangle$$

for $0 \leq \dots$
for external
cannot write
 \neq

If $N=2$ $k-2 = 2 = L$

Block



all states on boundary are pure.

$$\rho = \begin{pmatrix} P_1 & a \\ a^* & P_2 \end{pmatrix}$$

$$\text{tr } \rho = P_1 + P_2$$



$N > 2$ Boundary dim $k-2 > L$ pure state dim.

Hence most states on boundary are mixed. Also know that the pure state space is closed. Pringle (Ben)





Boundary of normalized states $k-2 = N^2-2$

In QT we also know that the pure states have dimension $L = 2N-2$

$$|\psi\rangle = \sum_{n=1}^N a_n |a_n\rangle$$

for $0 \leq \alpha$
for external
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 \neq

If $N=2$ $k-2 = 2 = L$

Bloch ball



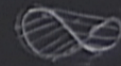
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$N > 2$ Boundary dim $k-2 > L$ pure state dim.

Hence most states on boundary are mixed. Also know that the pure state space is closed. Pringle (Ben)



Boundary of normalized states $k-2 = N^2-2$

but the pure states have

$$|\psi\rangle = \sum_{n=1}^N a_n |a_n\rangle$$

for $0 \leq \lambda \leq 1$

for external (pure) states ρ
cannot write

$$\rho = \lambda \rho_A + (1-\lambda) \rho_B \quad \text{for distinct } \rho_A, \rho_B \text{ and } 0 < \lambda < 1$$

$$2 = L$$

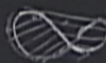
$$N > 2$$

Boundary dim $k-2 > L$ pure state dim.

Hence most states on boundary are mixed.

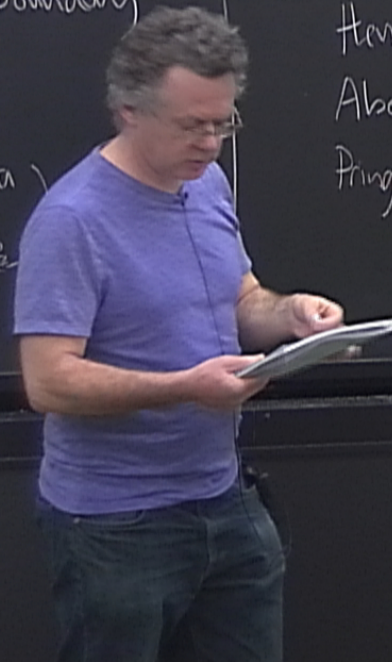
Also know that the pure state space is closed (no edges)

Pringle (Ben)



on boundary

$$\begin{pmatrix} \rho_A & a \\ a^\dagger & \rho_B \end{pmatrix}$$



$$\rho = (1-\lambda) \sum_{n=2}^N a_n |n\rangle\langle n| + \lambda |1\rangle\langle 1|$$

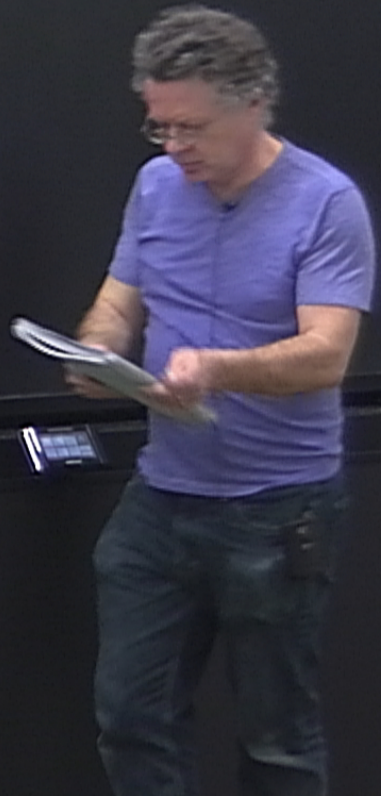
Can't make $\lambda < 0$ as then ρ is not +ve.

fuchs 0906.2187
0910.2750

State is pure $\rho = \rho^2$



$$\Rightarrow \rho = \rho^2$$
$$p_k = \text{quadratic fn of } p_{k+1}$$
$$= q_k(\{p_{k+1}\})$$

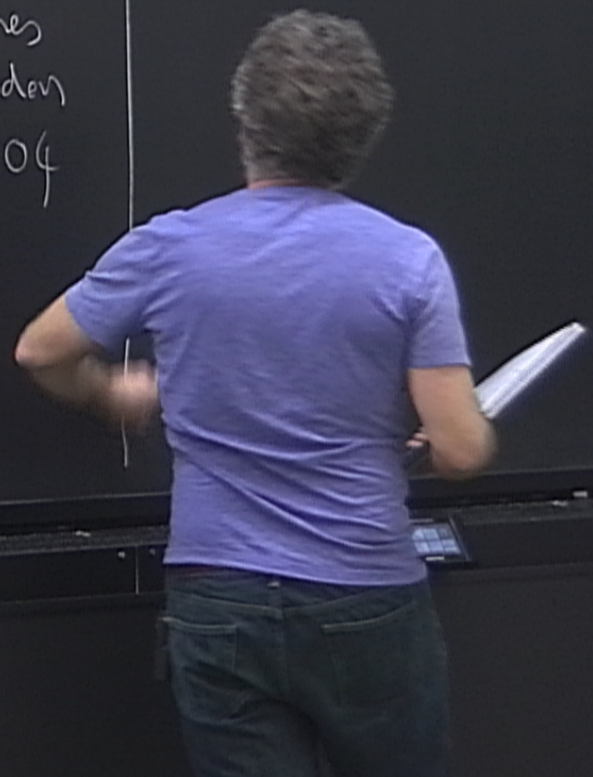


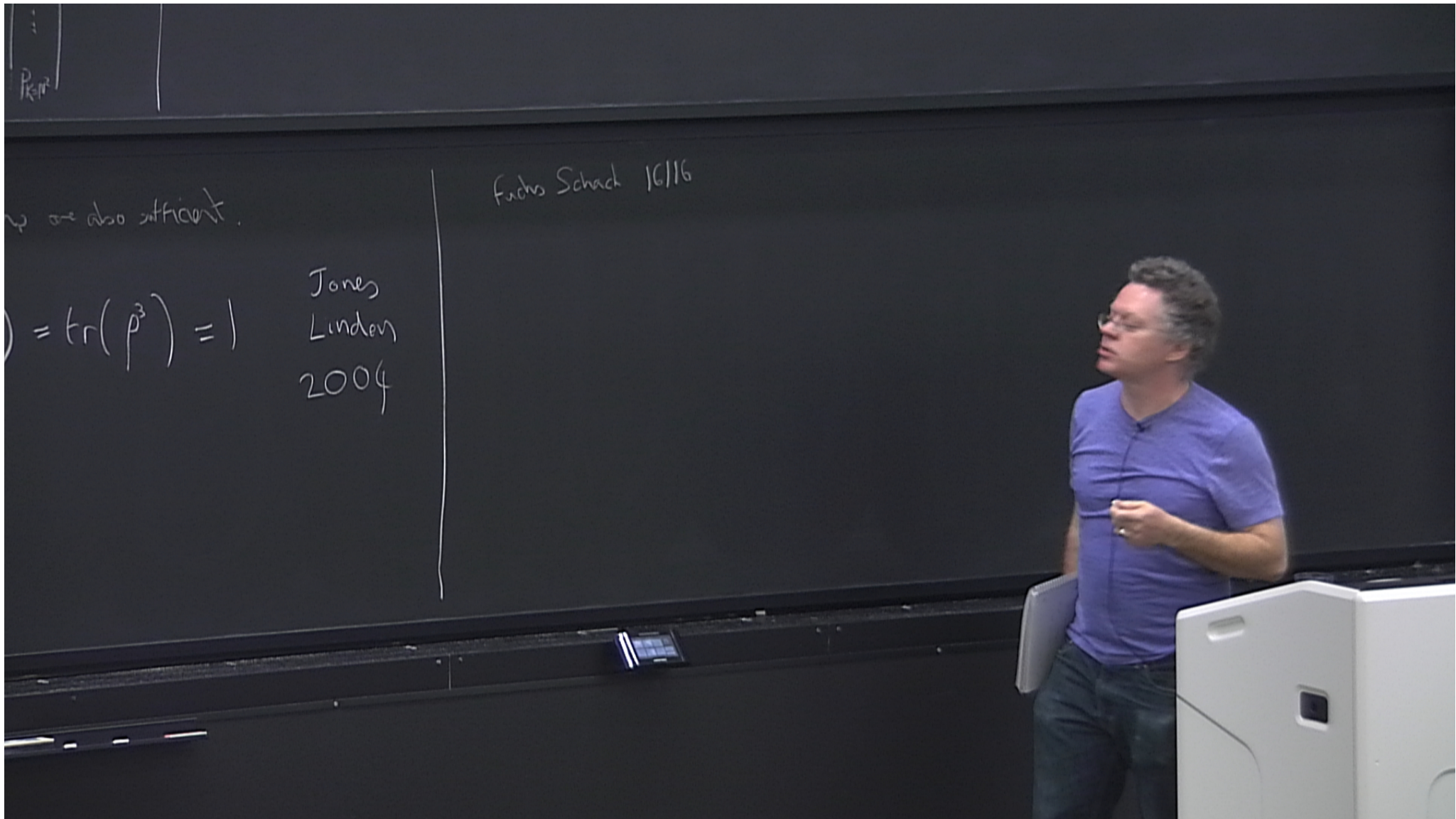
$\rho_{K|K^2}$

Turns out these cond-nums are also sufficient.
State v pnc it's-

$$\text{tr}(\rho^2) = \text{tr}(\rho^3) = 1$$

Jones
Linden
2004





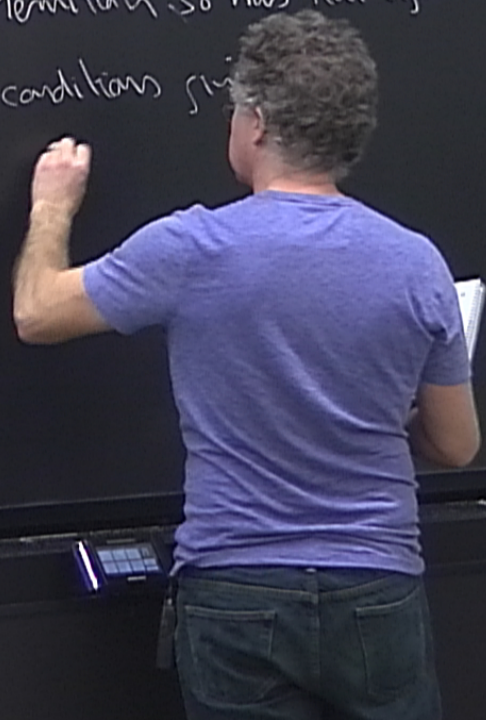
also sufficient.

$$\text{tr}(\rho^3) = 1$$

Jones
Linden
2004

Fuchs Schack 1976 proof

ρ is Hermitian, so has real eigenvalues λ_i
The two conditions give



λ_1
 λ_2
 \vdots
 λ_n

$$\textcircled{1} \quad |\lambda_i| \leq 1 \quad \forall i$$

$$\Rightarrow \lambda_i \geq -1 \quad \forall i$$

$\textcircled{1} - \textcircled{2}$

$$\sum_i \underbrace{\lambda_i^2}_{\geq 0} \underbrace{(1 - \lambda_i)}_{\geq 0} = 0$$

$$\Rightarrow \lambda_i = 1 \text{ or } 0.$$

p_1
 p_2
 \vdots
 p_{k+1}

$$\textcircled{1} \quad |\lambda_i| \leq 1 \quad \forall i$$

$$\Leftrightarrow 1 - \lambda_i \geq 0 \quad \forall i$$

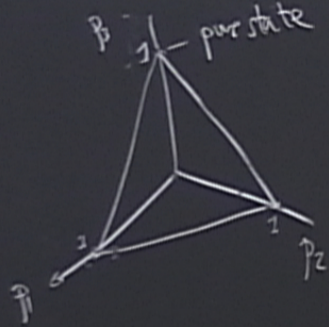
$\textcircled{1} - \textcircled{2}$

$$\sum_i \underbrace{\lambda_i^2}_{\geq 0} \underbrace{(1 - \lambda_i)}_{\geq 0} = 0$$

$$\Leftrightarrow \lambda_i = 1 \text{ or } 0.$$

$\textcircled{1}$ only one λ_i can equal 1

$P =$



$$0 \leq p_i \leq 1$$

$$\sum p_i = 1$$

Pure states have

$$\sum_i p_k = 1$$

$$\sum_i \hat{p}_k = 1$$

$$k=1 \text{ to } N^2 \quad \sum_k \hat{p}_k = 1$$

$$f = \begin{pmatrix} p_{2+} \\ p_{2-} \\ p_{2+} \\ p_{2+} \end{pmatrix}$$

$$\underline{1} = \begin{pmatrix} 1 \\ 1 \\ 0 \\ 0 \end{pmatrix}$$

$\underline{1} \cdot f = 1$
for normalized
stats.

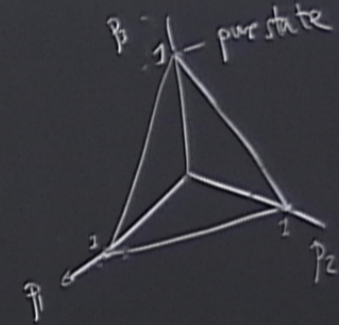
This gives us

$$\text{quadratic fn } (\{p_k\}) = 1$$

$$\text{cubic fn } (\{p_k\}) = 1$$

Two equations.

$$\begin{array}{c} \boxed{p} \\ \downarrow \\ \boxed{A^x} \end{array}$$



$$0 \leq p_k \leq 1$$

$$\sum p_k = 1$$

$$k=1 \text{ to } N^2$$