

Title: Standard Model-2

Date: Jan 08, 2015 09:00 AM

URL: <http://pirsa.org/15010003>

Abstract:

Today:
 Finish: Cast of Characters
 SM: QCD ($SU(3)_c$) $SU(2)_L \times U(1)_Y$
 8 Gluons, 6 quarks x 3 colors Electroweak Sector
 $\begin{pmatrix} u \\ c \\ t \end{pmatrix}$ $\begin{pmatrix} d \\ s \\ b \end{pmatrix}$ - No
 - Ch
 - P
 $\begin{matrix} G \\ G \end{matrix}$ $\begin{matrix} \psi \\ \psi \end{matrix}$

Today:
 Finish: Cast of Characters:

SM: QCD ($SU(3)_c$)
 8 Gluons, 6 quarks x (3 colors)
 $\begin{pmatrix} u \\ c \\ t \end{pmatrix}$ $\begin{pmatrix} d \\ s \\ b \end{pmatrix}$
 - No flavor changing
 - prescribes CP?

$\begin{pmatrix} X \\ W \\ Y \end{pmatrix}$ B_μ
 $SU(2)_L \times U(1)_Y \rightarrow U(1)_{EM}$
 Electroweak Sector
 W^\pm, Z^0, γ (after EWSB)

Today:
 Finish: List of Characters:

SM: QCD ($SU(3)_c$)
 8 Gluons, 6 quarks x (3 colors)
 - No flavor changing
 - preserves CP?

$SU(2)_L \times U(1)_Y \rightarrow U(1)_{EM}$
 Electroweak Sect
 W^\pm, Z^0, γ (all)
 u_i, d_j
 flavor change

leptons: Spin 1/2
 e, μ, τ
 ν_e, ν_μ, ν_τ
 $m=0$ in SM

Today:
 Finish: Cast of Characters:

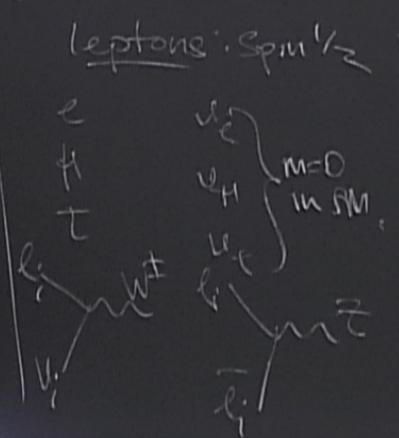
SM: QCD ($SU(3)_C$)
 8 Gluons.
 3 colors
 - No flavor changing processes
 CP?

Electroweak Sector
 $SU(2)_L \times U(1)_Y \rightarrow U(1)_{EM}$
 W^\pm, Z^0, γ (after EWSB)
 Flavor changing: $u_i \rightarrow d_j$ via W^\pm
 No flavor change (tree level): $q_i \rightarrow q_i$ via Z^0, γ

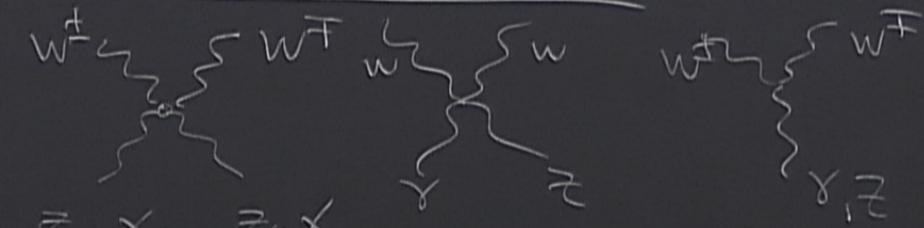
leptons: Spin 1/2
 e, μ, τ
 ν_e, ν_μ, ν_τ
 $m=0$ in SM

$U(1)_Y \rightarrow U(1)_{EM}$
 Sector after (WSB)

No flavor change (tree level)

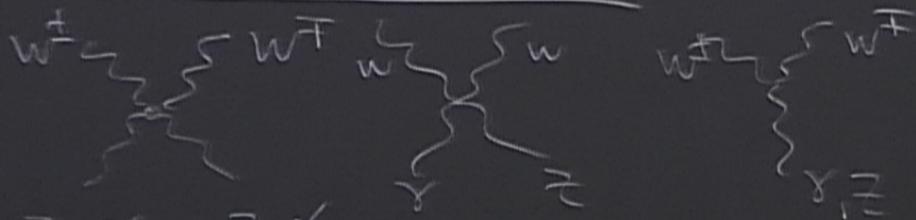


Gauge Boson Self Interactions:



Electroweak Int.
 Break parity

Gauge Boson Self Interactions:

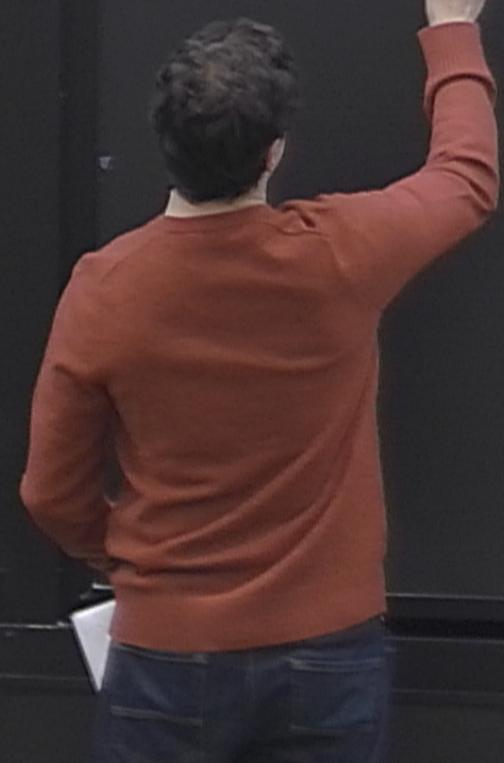


$$= \frac{g^2}{4} \delta$$

Electroweak Int.
 Break parity (LH \neq RH)
 Can also Violate CP

$QED = \gamma$ $LH = RH$

Higgs = (20)

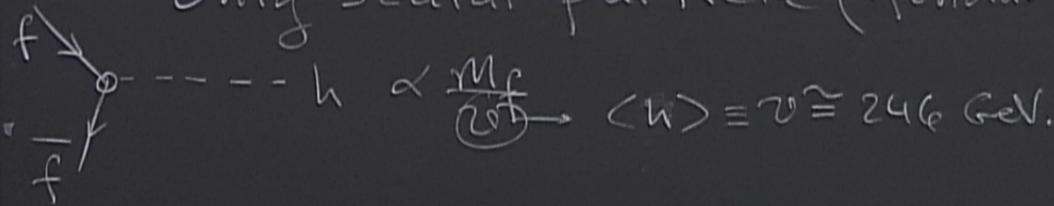


Higgs: (2012) LHC $m_H = 125$ GeV.

Only scalar particle (fundamental) in SM!

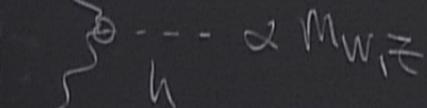
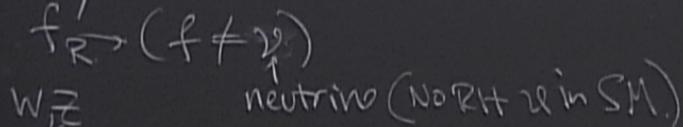
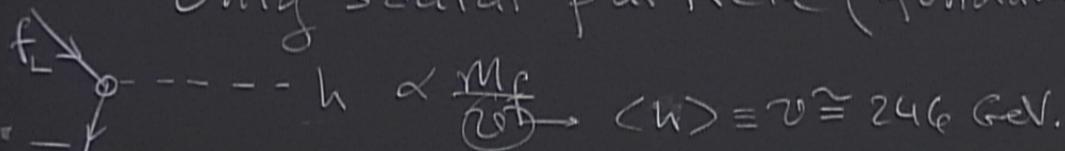
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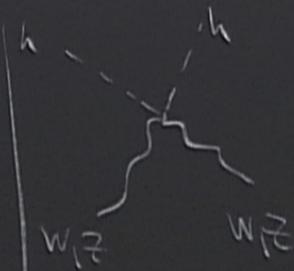


Higgs: (2012) LHC $m_h = 125$ GeV.

Only scalar particle (fundamental) in SM!



W, Z



h is Neutral Under QED

enta) in SM!

is Neutral
nder QED

Machinery:

where we're going:

- 2 component Spinors
(Lorentz group)
- Majorana / Dirac Masses
- 2 vs. 4 Component Spin
- Goldstone theorem



enta) in SM!

is Neutral
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Machinery:

where we're going:

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(Lorentz group)
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Lorentz Group

x^μ (four vectors)

$$x^\mu \rightarrow x'^\mu \equiv \Lambda^\mu{}_\nu x^\nu$$

Lorentz Group:

x^μ (four vectors)

$$x^\mu \rightarrow x'^\mu = \Lambda^\mu{}_\nu x^\nu$$

demand: $x^\mu x_\mu = x'^\mu x'_\mu$

Metric tensor as:

$g_{\mu\nu}$

give

Symbol

contracts Vectors \Rightarrow Scalars.

contracts Vectors \Rightarrow Scalars.

Transformations: $\vec{\Theta} \equiv \Theta \hat{n}$, $\vec{\xi} \equiv \hat{v} \tanh^{-1}(|\vec{v}|)$ boost param.
axis of Rot.

Gen

Contracts Vectors \Rightarrow Scalars.

Transformations: $\vec{\Theta} \equiv \Theta \hat{n}$, $\vec{\Sigma} = \hat{v} \tanh^{-1}(|\vec{v}|)$ boost param.
axis of Rot.

Generators Λ

$$\exp\left(-\frac{i}{2} \omega_{\mu\nu} J_{\mu\nu}\right)$$

Choice of

Contracts Vectors \Rightarrow Scalars.

Transformations: $\vec{\theta} \equiv \theta \hat{n}$ axis Rot, $\vec{\xi} \equiv \hat{v} \tanh^{-1}(\vec{v})$ boost param.

Generic Trans. Λ

$$\Lambda_{\mu\nu}^H \equiv \left[\exp\left(-\frac{i}{2} \omega^{\alpha\beta} J_{\alpha\beta} + \xi^i K_i\right) \right]^H$$

Choice of trans. param.

define $\omega^{\alpha\beta}$ as:

$$\theta^i = \frac{1}{2} \epsilon^{ijk} \omega_{jk}, \quad \xi^i = \omega^{0i} = -\omega^{i0}$$

$$\epsilon^{123} = +1$$

Things we choose

Vectors \Rightarrow Scalars.

relations: $\vec{\theta} \equiv \theta \hat{n}$, $\vec{\xi} \equiv \hat{v} \tanh^{-1}(\vec{v})$ boost param.
axis of Rot.

us. Λ

$$\exp\left(-\frac{i}{2} \omega_{ij} J_{ij} + \frac{1}{2} \omega_{0i} J_{0i}\right)$$

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define $\omega^{\alpha\beta}$ as:

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Things we choose.

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Vectors \Rightarrow Scalars.

Rotations: $\vec{\theta} \equiv \theta \hat{n}$, $\vec{\xi} \equiv \hat{v} \tanh^{-1}(|\vec{v}|)$ boost param.
axis of Rot.

us. Λ

$$\exp\left(-\frac{i}{2} \omega^{jk} J_{jk}\right)$$
Choice of trans. param.
Generators (axis-label Matrix)

define ω^{jk} as:

$$\theta^i \equiv \frac{1}{2} \epsilon^{ijk} \omega_{jk}, \quad \xi^i = \omega^{0i} = -\omega^{i0}$$

$\epsilon^{123} = +1$

Things we choose.

Generators:

$$J^i \equiv \frac{1}{2} \epsilon^{ijk} J_{jk}, \quad K^i \equiv J^{0i} = -J^{i0}$$

vectors \Rightarrow Scalars.

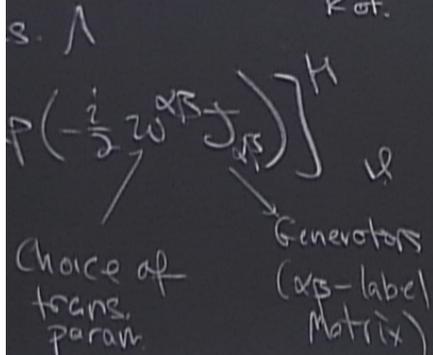
Rotations: $\vec{\Theta} \equiv \Theta \hat{n}$
axis of Rot.

$\vec{\omega} = \hat{v} \tanh^{-1}(|\vec{v}|)$ boost param.

define $w^{\alpha\beta}$ as:

$\Theta^i = \frac{1}{2} \epsilon^{ijk} w_{jk}$, $\vec{\omega} = w^{\alpha\beta} \hat{e}_{\alpha\beta}$
 $\epsilon^{123} = +1$

Things we choose.



Generators:

$J^i = \frac{1}{2} \epsilon^{ijk} J_{jk}$

vectors \Rightarrow Scalars.

$\vec{\theta} \equiv \theta \hat{n}$
 axis of Rot.

$\vec{w} = \hat{v} \tanh^{-1}(\vec{v})$ boost param.

define $w^{\alpha\beta}$ as:

$\theta^i = \frac{1}{2} \epsilon^{ijk} w_{jk}$
 $\epsilon^{123} = +1$

$w^{0i} = -w^{i0}$

Things we choose.

$\Lambda = \exp\left(-\frac{i}{2} w_{\alpha\beta} J^{\alpha\beta}\right)$
 Generators (axis-label Matrix)
 Choice of trans. param.

Generators:

$J^i = \frac{1}{2} \epsilon^{ijk} J^{jk}$

$J^{0i} = -J^{i0}$

$\Lambda_u = \left[\exp(-i \vec{\theta} \cdot \vec{J}) \right]$

vectors \Rightarrow Scalars.

Rotations: $\vec{\theta} \equiv \theta \hat{n}$ axis of Rot., $\vec{\omega} \equiv \hat{v} \tanh^{-1}(\vec{v})$ boost param.

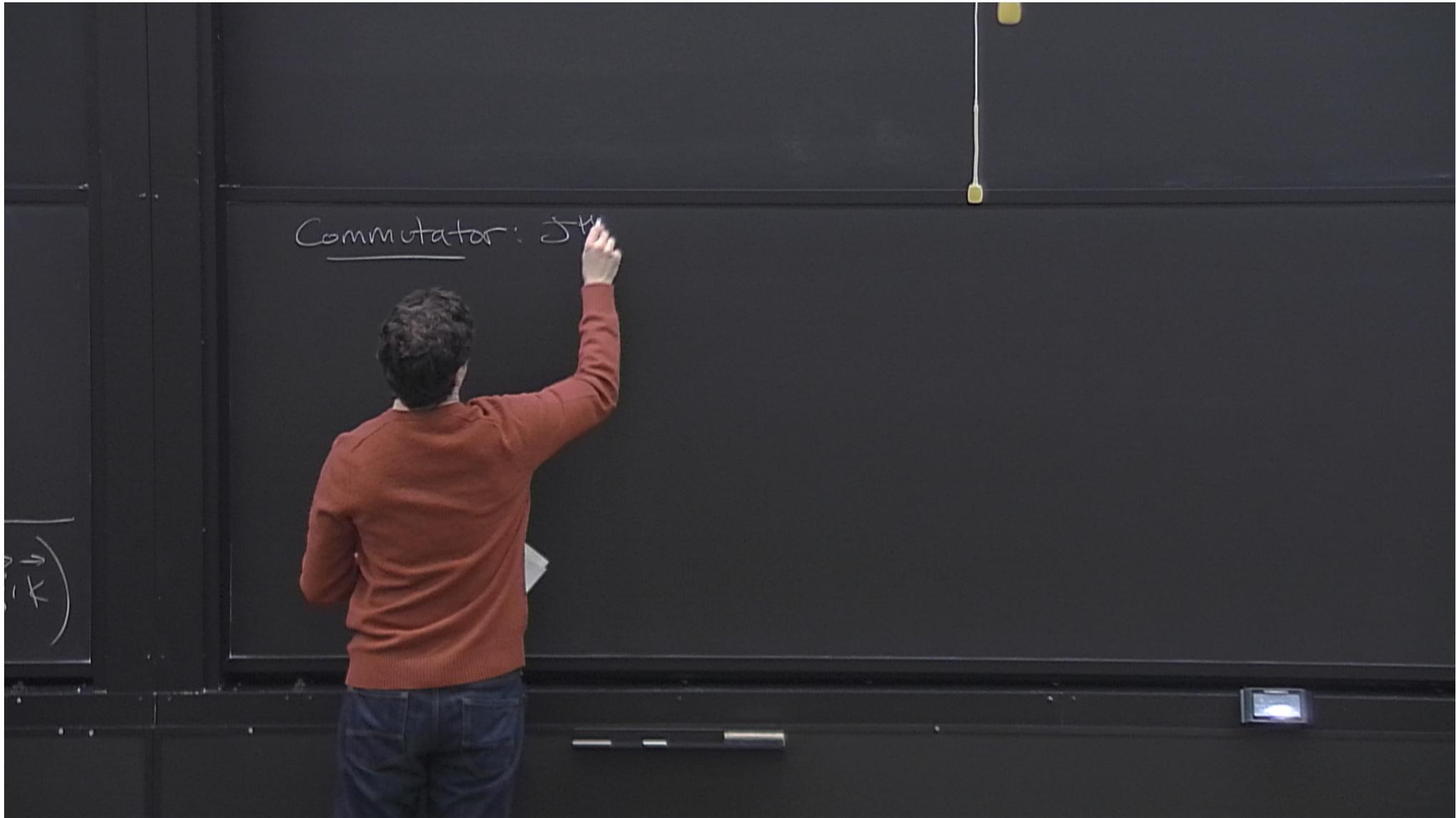
as: w^{AB} as:

$\sum_{i,j,k} \epsilon^{ijk} w_{jk}$, $\vec{\omega} = w^{0i} = -w^{i0}$

Things we choose.

Choice of trans. param. Generators (x^B-lab. Matrix)

$$K^i \equiv J^{0i} = -J^{i0} \quad \Lambda_{\mu}^{\nu} = \left[\exp(-i\vec{\theta} \cdot \vec{J} - i\vec{\xi} \cdot \vec{K}) \right]$$



Commutator: $J_{\mu\nu}$

$$[J_{\mu\nu}, J_{\rho\sigma}] = i \left(g_{\mu\rho} J_{\nu\sigma} - (\mu \leftrightarrow \nu) \right) - (\rho \leftrightarrow \sigma)$$

J 's will differ
depending on what we're
acting on!

Commutator: $J^{\mu\nu}$

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for four vectors:

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J 's will differ depending on what we're acting on!

for four vectors:

$$\left(\begin{array}{c} J \\ \rho\sigma \end{array} \right)_{\mu}^{\nu} = i \left(g_{\rho}^{\mu} g_{\sigma\nu} - g_{\sigma}^{\mu} g_{\rho\nu} \right)$$

labels

Commutator: $J^{\mu\nu}$

$$[J^{\mu\nu}, J^{\rho\sigma}] = i \left(g^{\mu\rho} J^{\nu\sigma} - (\mu \leftrightarrow \nu) \right) - (\rho \leftrightarrow \sigma)$$

J 's will differ depending on what we're acting on!

for four vectors:

$$\left(J_{\rho\sigma}^{\mu} \right)_{\nu} = i \left(g_{\rho}^{\mu} g_{\sigma\nu} - g_{\sigma}^{\mu} g_{\rho\nu} \right)$$

labels

Check Me on sign & indices

operator: $J^{\mu\nu}$

$$J^{\mu\nu} = i \left(g^{\mu\alpha} J^{\nu\alpha} - (u \leftrightarrow v) \right)$$

J 's will differ depending on what is acting on!

vectors:

$$J^{\mu\nu} = i \left(g^{\mu\alpha} g^{\nu\beta} - g^{\nu\alpha} g^{\mu\beta} \right)$$

* Check

indices

operator: $J^{\mu\nu}$

* Check.

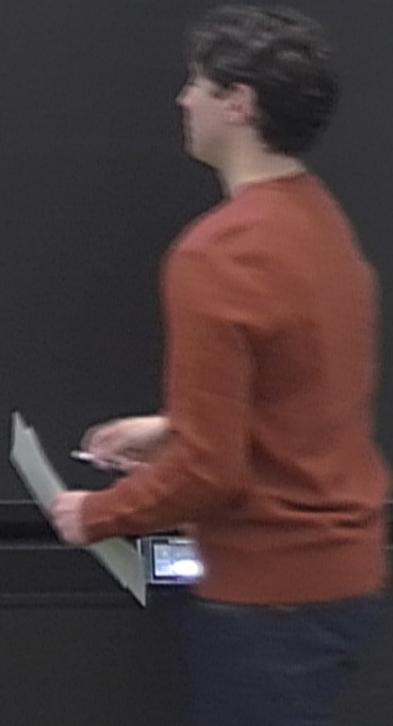
$$J^{\mu\nu} = i \left(g^{\mu\alpha} J^{\nu\alpha} - (\mu \leftrightarrow \nu) \right) - (p \leftrightarrow \sigma)$$

J 's will differ depending on what we're acting on!

Vectors:

Check Me on sign & indices

$$J^{\mu\nu} = i \left(g^{\mu\alpha} g^{\nu\beta} - g^{\nu\alpha} g^{\mu\beta} \right)$$



contracts Vectors \Rightarrow Scalars.

transformations: $\vec{\Theta} \equiv \Theta \hat{n}$,
axis of Rot.

$\vec{\Sigma} \equiv \hat{v} \tanh^{-1}(|\vec{v}|)$ boost param.

generic Trans. Λ

$$\Lambda_{\mu\nu}^H = \left[\exp\left(-\frac{i}{2} \omega_{\alpha\beta} J_{\alpha\beta}^H\right) \right]_{\mu\nu}$$

Choice of trans. param \rightarrow Generators (xp-label matrix)

define $\omega^{\alpha\beta}$ as:

$$\Theta^i \equiv \frac{1}{2} \epsilon^{ijk} \omega_j$$

$$\omega^{\alpha\beta} = -\omega^{\beta\alpha}$$

$$\epsilon^{123} = +1$$

Things we choose.

Generators:

$$J^i \equiv \frac{1}{2} \epsilon^{ijk} J_{jk}$$

$$\Lambda_{\mu\nu}^H = \left[\exp\left(-i \vec{\Theta} \cdot \vec{J} - i \vec{\Sigma} \cdot \vec{K}\right) \right]_{\mu\nu}$$

* Check. (Srednicki Text QFT)

what about spinors?

$$\text{define } \psi_\alpha(x) \xrightarrow{\Lambda} U(\Lambda)^{-1} \psi_\alpha(x) U(\Lambda) \equiv M_\alpha^\beta(\Lambda) \psi_\beta(\Lambda^{-1}x)$$

what is $M_\alpha^\beta(\Lambda)$?

different.

$$M(\Lambda) = \exp\left(-\frac{i}{2} \omega^{\alpha\beta} J_{\alpha\beta}\right)$$

$$\vec{J} \equiv \vec{S}/2$$

$$\vec{K} \equiv -\vec{S}/2$$

$\left. \right) - (p \leftrightarrow \sigma)$

at we're

check me
on sign & indices

$g_{\mu\nu} = \Lambda^\alpha_\mu \Lambda^\beta_\nu g_{\alpha\beta}$
 $g_{\mu\nu}$ is an Invariant Symbol.

Choice of
trans.
param.

Generators
(axis-label
Matrix)

Generators:

$$J^i \equiv \frac{1}{2} \epsilon^{ijk} J_{jk}, \quad K^i \equiv J^0{}^i$$

$\vec{\sigma} \equiv (\sigma^1, \sigma^2, \sigma^3)$ Pauli Matrices: $\sigma^1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$, $\sigma^2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$, $\sigma^3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

LH: $(\frac{1}{2}, 0)$

$g_{\mu\nu} = \Lambda^\alpha_\mu \Lambda^\beta_\nu g_{\alpha\beta}$
 is an invariant symbol.

Choice of trans. param.

Generators (vs-label Matrix)

Generators:

$$J^i \equiv \frac{1}{2} \epsilon^{ijk} J_{jk}, \quad K^i \equiv J^{0i} = -J^{i0}$$

$\vec{\sigma} \equiv (\sigma^1, \sigma^2, \sigma^3)$ Pauli Matrices: $\sigma^1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$, $\sigma^2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$, $\sigma^3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

$\underline{L}_H = (\frac{1}{2}, 0)$, Build Conj Rep. $\vec{J}_{RH} \equiv \frac{\vec{\sigma}^x}{2}$, $\vec{K}_{RH} = -\frac{\vec{\sigma}^x}{2}$ RH Rep acts on χ^a - dotted l



$g_{\mu\nu} = \Lambda^\alpha_\mu \Lambda^\beta_\nu g_{\alpha\beta}$
 $g_{\mu\nu}$ is an Invariant Symbol.

Choice of trans. param

Generators (as-label Matrix)

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LH: $(\frac{1}{2}, 0)$, Build Conj Rep. $\vec{J}_{RH} \equiv \frac{\vec{\sigma}^x}{2}$, $\vec{K}_{RH} = -i\vec{\sigma}^x/2$ RH Rep acts on χ^a

Conjugate On LH b/c $\psi_2 \rightarrow (\psi_2)^\dagger \equiv \psi_2^\dagger$

How to find Invariants

$g_{\mu\nu} = \Lambda^\mu_\alpha \Lambda^\nu_\beta g_{\alpha\beta}$
 $g_{\mu\nu}$ is an Invariant Symbol.

Choice of trans. param

Generators (x₃-label Matrix)

Generators:

$$J^i \equiv \frac{1}{2} \epsilon^{ijk} J_{jk}, \quad K^i \equiv J^{0i}$$

$\vec{\sigma} \equiv (\sigma^1, \sigma^2, \sigma^3)$ Pauli Matrices: $\sigma^1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$, $\sigma^2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$, $\sigma^3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

LH: $(\frac{1}{2}, 0)$, Build Conj Rep. $\vec{J}_{RH} \equiv \frac{\vec{\sigma}^*}{2}$, $\vec{K}_{RH} = -i\vec{\sigma}^*/2$ RH Rep acts on $\chi^{\alpha} - d$

Concentrate on LH b/c $\psi_{\alpha} \rightarrow (\psi_{\alpha})^{\dagger} \equiv \psi_{\dot{\alpha}}$

How to build Invariants?

$g_{\mu\nu} = \Lambda^\alpha_\mu \Lambda^\beta_\nu g_{\alpha\beta}$
 $g_{\mu\nu}$ is an Invariant Symbol.

Choice of trans. param

Generators (as-label Matrix)

Generators:

$$J^i \equiv \frac{1}{2} \epsilon^{ijk} J_{jk}, \quad K^i \equiv J^{0i}$$

$\vec{\sigma} \equiv (\sigma^1, \sigma^2, \sigma^3)$ Pauli Matrices: $\sigma^1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$, $\sigma^2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}$, $\sigma^3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

LH: $(\frac{1}{2}, 0)$, Build Conj Rep. $\vec{J}_{RH} \equiv \frac{\vec{\sigma}^*}{2}$, $\vec{K}_{RH} = -i\vec{\sigma}^*/2$ RH Rep acts on χ^{α}

Concentrate on LH b/c $\psi_\alpha \rightarrow (\psi_\alpha)^\dagger \equiv \psi_\alpha^\dagger$

How to build Invariants?

Note: $\chi_\mu \phi_\nu \rightarrow g^{\mu\nu} \chi_\mu \phi_\nu = \text{Inv.}$

$$g_{\mu\nu} \rightarrow \Lambda^\alpha_\mu \Lambda^\beta_\nu g_{\alpha\beta}$$

$g_{\mu\nu} = \Lambda^\alpha_\mu \Lambda^\beta_\nu g_{\alpha\beta}$
 $g_{\mu\nu}$ is an invariant symbol.

Choice of trans. param

Generators (xp-label Matrix)

Generators:

$J^i \equiv \frac{1}{2} \epsilon^{ijk} J_{jk}, K^i \equiv J^{0i}$

$\vec{\sigma} \equiv (\sigma^1, \sigma^2, \sigma^3)$ Pauli Matrices: $\sigma^1 = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \sigma^2 = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \sigma^3 = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$

LH: $(\frac{1}{2}, 0)$, Build Conj Rep. $\vec{J}_{RH} \equiv \frac{\vec{\sigma}^*}{2}, \vec{K}_{RH} = -i\vec{\sigma}^*/2$ RH Rep acts on χ^α

concentrate on LH b/c $\psi_\alpha \rightarrow (\psi_\alpha)^\dagger \equiv \psi_\alpha^\dagger$

How build invariants?

Note: $\chi \rightarrow g^{\mu\nu} \chi_\mu \chi_\nu = \text{Inv.}$

$g_{\mu\nu} \rightarrow \Lambda^\alpha_\mu \Lambda^\beta_\nu g_{\alpha\beta}$

Recall: Combine $Sp=1/2$ objects \rightarrow Singlet & triplet

$\psi_\alpha \psi_\beta \equiv C_{\alpha\beta} = \epsilon_{\alpha\beta} \underbrace{D(x)}_{\text{Scalar fn.}} + \underbrace{G_{\alpha\beta}}_{\text{Symm.}}$

$g_{\mu\nu} = \Lambda^\alpha_\mu \Lambda^\beta_\nu g_{\alpha\beta}$
 $g_{\mu\nu}$ is an invariant symbol.

Choice of trans. param

Generators (as-label Matrix)

Generators.

$$J^i \equiv \frac{1}{2} \epsilon^{ijk} J_{jk}, \quad K^i \equiv J^{0i}$$

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LH: $(\frac{1}{2}, 0)$, Build Conj Rep: $\vec{J}_{RH} \equiv \frac{\vec{\sigma}^*}{2}$, $\vec{K}_{RH} = -i\vec{\sigma}^*/2$ RH Rep acts on χ^{α}

Concentrate on LH b/c $\psi_\alpha \rightarrow (\psi_\alpha)^\dagger \equiv \psi_\alpha^\dagger$

How to build invariants?

Note: $X_\mu \phi_\nu \rightarrow g^{\mu\nu} X_\mu \phi_\nu = \text{Inv.}$

$$g_{\mu\nu} \rightarrow \Lambda^\alpha_\mu \Lambda^\beta_\nu g_{\alpha\beta}$$

Recall: Combine $Sp = 1/2$ objects \rightarrow Singlet & triplet

$$\psi_\alpha \psi_\beta \equiv C_{\alpha\beta} = \epsilon_{\alpha\beta} \underbrace{D(x)}_{\text{Scalar fn.}} + \underbrace{G_{\alpha\beta}}_{\text{Symm.}}$$

Operators.

$$\begin{pmatrix} \vec{J} \\ \vec{K} \end{pmatrix}$$

for four vectors:

$$\begin{pmatrix} \vec{J} \\ \vec{K} \end{pmatrix}^{\mu} = i \left(g_{\mu\nu}^{\mu} g_{\nu\rho} - g_{\mu\rho}^{\mu} g_{\nu\nu} \right)$$

labels

Check on sign & indices

$$\begin{aligned} \vec{J} &\equiv \vec{\sigma}/2 \\ \vec{K} &\equiv -i\vec{\sigma}/2 \end{aligned}$$

Act on a/s part:

$$U(\Lambda) \in_{\alpha\beta} D(x) U(\Lambda) \equiv M_{\alpha\beta}^K$$

Spinor indices
 $\alpha, \beta = 1 \text{ or } 2$

$$\begin{pmatrix} \vec{J} \\ \vec{K} \end{pmatrix}$$

for four vectors:

$$\left(\begin{matrix} \vec{J} \\ \vec{K} \end{matrix} \right)^\mu = i \left(\begin{matrix} g_{\mu\nu}^+ g_{\rho\sigma} - g_{\mu\rho}^+ g_{\sigma\nu} \end{matrix} \right)$$

labels

Check on sign & indices

$$\begin{aligned} \vec{J} &\equiv \vec{S}/2 \\ \vec{K} &\equiv -i\vec{S}/2 \end{aligned}$$

Act on a/s part.

$$\bar{U}^{-1}(\Lambda) \epsilon_{\alpha\beta} D(x) U(\Lambda) \equiv M_{\alpha}^{\lambda}(\Lambda) M_{\beta}^{\kappa}(\Lambda) \epsilon_{\lambda\kappa} D(\Lambda^{-1}x)$$

Spinor indices
 $\alpha, \beta = 1 \text{ or } 2$

$$= \epsilon_{\alpha\beta} \underbrace{\bar{U}^{-1}(\Lambda) D(x) U(\Lambda)}_{\equiv D(\Lambda^{-1}x)} = M_{\alpha}^{\lambda}(\Lambda) M_{\beta}^{\kappa}(\Lambda) \epsilon_{\lambda\kappa}$$

check \vec{k}
on G on δ indices

$$\vec{J} = \vec{\sigma}/2 \quad \vec{\sigma} = (\sigma^1, \sigma^2, \sigma^3)$$
$$\vec{K} = -i\vec{\sigma}/2$$

$D(\Lambda^{-1}x)$

$$\epsilon_{\alpha\beta} = M_{\alpha}^{\kappa} M_{\beta}^{\lambda} \epsilon_{\kappa\lambda} \quad \text{just like } g_{\mu\nu}!$$

- Raises / Lowers Spinor Indices
- Contracts Spinors \rightarrow Scalars!

$$\psi_{\alpha} \psi_{\beta} \epsilon^{\alpha\beta} = \psi^{\alpha} \psi_{\alpha}$$

$$\epsilon^{12} = +1$$

Choice of trans. param

Generators (x₃-label Matrix)

Generators:

$$J^i \equiv \frac{1}{2} \epsilon^{ijk} J_{jk}, \quad K^i \equiv J^{0i} = -J^{i0}$$

$$\Lambda_{\mu\nu}^H = \left[\exp(-i \vec{\theta} \cdot \vec{J} - i \vec{\xi} \cdot \vec{K}) \right]$$

What's the simplest 2-comp theory?

$$\mathcal{L} = i \psi^\dagger \left(\sigma^H \right) \partial_t \psi$$

Recall: $\gamma^t = \begin{pmatrix} 0 & \sigma^H \\ \sigma^H & 0 \end{pmatrix}$

$$\sigma^H \equiv \left(\mathbb{1}_{2 \times 2}, \vec{\sigma} \right)$$

$$\overline{\sigma^H} = \left(\mathbb{1}_{2 \times 2}, -\vec{\sigma} \right)$$

(for σ^H)
 α_2

Choice of trans. param

Generators (x₃-label Matrix)

Generators:

$$J^i \equiv \frac{1}{2} \epsilon^{ijk} J_{jk}, \quad K^i \equiv J^{0i} = -J^{i0}$$

$$\Lambda_{\vec{v}}^{\mu} = \left[\exp(-i \vec{\theta} \cdot \vec{J} - i \vec{\beta} \cdot \vec{K}) \right]$$

What's the simplest 2-comp theory?

$$\mathcal{L} = i \psi_{\dot{\alpha}}^{\dagger} (\sigma^{\mu})^{\dot{\alpha}\alpha} \partial_{\mu} \psi_{\alpha} + \frac{m}{2} (\psi_{\dot{\alpha}} \psi_{\alpha} + \psi_{\alpha}^{\dagger} \psi_{\dot{\alpha}}^{\dagger})$$

Same 2-comp Spinors Majorana

Recall: $\gamma^{\mu} = \begin{pmatrix} 0 & \sigma^{\mu} \\ \sigma^{\mu} & 0 \end{pmatrix}$

$$\sigma^{\mu} \equiv (\mathbb{1}_{2 \times 2}, \vec{\sigma})$$

$$\bar{\sigma}^{\mu} = (\mathbb{1}_{2 \times 2}, -\vec{\sigma})$$

(for σ^{μ})
 $\alpha_{\dot{\alpha}}$

$$\psi_{\dot{\alpha}} \psi_{\alpha} = \chi^{\alpha} \psi_{\dot{\alpha}}$$

$$H: \psi_{\dot{\alpha}}^{\dagger} \chi^{\alpha} \rightarrow$$

$$g_{\mu\nu} = \Lambda^\alpha_\mu \Lambda^\beta_\nu g_{\alpha\beta}$$

$$g_{\mu\nu} \text{ is an Invariant Symbol.}$$

Choice of trans. param

Generators (x³-label Matrix)

Generator $J^i \equiv \frac{1}{2} \epsilon^{ijk} J^j J^k$

So what's wrong? Cant couple to spin 1 particle!
 U(1) local gauge Symm: $\psi_\alpha \rightarrow e^{iQ_\alpha f(x)} \psi_\alpha(x)$

Promote kinetic term $\rightarrow D_\mu \equiv \partial_\mu - ig_\alpha Q_\alpha A_\mu$, $A_\mu \rightarrow A_\mu - \frac{1}{g} \partial_\mu f(x)$

$$i \psi_\alpha^\dagger (\overleftrightarrow{\sigma}^{\mu\nu})_{\alpha\beta} D_\mu \psi_\beta + \frac{m}{2} (\psi_\alpha^\dagger \psi_\alpha + \text{h.c.})$$

$g_{\mu\nu} = \Lambda^\mu_\alpha \Lambda^\nu_\beta g^{\alpha\beta}$
 $g_{\mu\nu}$ is an invariant symbol.

Choice of trans. param.

Generators (sp-label Matrix)

Generators:
 $J^i \equiv \frac{1}{2} \epsilon^{ijk} J_{jk}$, $K^i \equiv J^{0i} = -J^{i0}$

So what's wrong? Can't couple to spin-1 particle!

$U(1)$ local gauge symm: $\psi_\alpha \rightarrow e^{iQ_\alpha f(x)} \psi_\alpha(x)$

Promote kinetic term

$\partial_\mu \rightarrow D_\mu \equiv \partial_\mu - ig_\psi Q_\psi A_\mu$, $A_\mu \rightarrow A_\mu - \frac{1}{g} \partial_\mu f(x)$

$\mathcal{L}_0 = i \underbrace{\psi_\alpha^\dagger (\overleftrightarrow{\partial}_\mu)^{\alpha\beta}}_{\text{ok!}} D_\mu \psi_\beta + \frac{m}{2} (\bar{\psi} \psi + \text{h.c.})$, Mass term: $\bar{\psi}^\alpha \psi_\alpha \rightarrow e^{i2Q_\psi f(x)} \bar{\psi}^\alpha \psi_\alpha$!

for four vectors:

$$\left(\frac{\vec{J}}{\vec{K}}\right)_{\mu}^{\nu} = i \left(g_{\mu}^{\alpha} g_{\alpha\nu} - g_{\mu\nu} \right)$$

labels

Check on sign & indices

$$\vec{J} = \vec{\sigma}/2 \quad \vec{K} = -i\vec{\sigma}/2$$

fix problem: Add second field! χ w/ $Q_{\chi} = -Q_{\psi}$

$$\mathcal{L}_{\chi} = i \chi_{\alpha}^{\dagger} (\bar{\sigma}^{\mu})^{\dot{\alpha}\alpha} \mathcal{D}_{\mu} \chi_{\alpha} + i \psi_{\dot{\alpha}}^{\dagger} (\bar{\sigma}^{\mu})^{\dot{\alpha}\alpha} \mathcal{D}_{\mu} \psi_{\alpha} + m \left(\psi_{\alpha} \chi^{\alpha} + \psi_{\dot{\alpha}}^{\dagger} \chi^{\dagger\dot{\alpha}} \right)$$

Chose χ, ψ to both be LH,

Dirac Mass term

for four vectors:

$$\left(\frac{\vec{J}}{\vec{K}}\right)_{\mu\nu} = i(g_{\mu}^{\alpha} g_{\alpha\nu} - g_{\mu\nu}^{\alpha} g_{\alpha}^{\nu})$$

labels

Check on sign & indices

$$\vec{J} = \vec{\sigma}/2 \quad \vec{K} = -i\vec{\sigma}/2$$

fix problem: Add second field! χ w/ $Q\chi = -Q\psi$

$$= i\chi_{\alpha}^{\dagger} (\bar{\sigma}^{\mu})^{\dot{\alpha}\alpha} \partial_{\mu} \chi_{\dot{\alpha}} + i\psi_{\dot{\alpha}}^{\dagger} (\bar{\sigma}^{\mu})^{\dot{\alpha}\alpha} \partial_{\mu} \psi_{\alpha} + m \underbrace{(\psi_{\alpha} \chi^{\dagger\dot{\alpha}} + \psi_{\dot{\alpha}}^{\dagger} \chi^{\alpha})}_{\text{Dirac Mass term}}$$

chose χ, ψ to both be LH,

$$\bar{\Psi}_D \equiv \begin{pmatrix} \psi_{\alpha} & \chi^{\dagger\dot{\alpha}} \end{pmatrix} \begin{matrix} Q\psi \\ Q\chi = -Q\psi \end{matrix}$$

for four vectors:

$$\left(\frac{\vec{J}}{\vec{K}}\right)_{\mu\nu} = i \left(g_{\mu}^{\alpha} g_{\alpha\nu} - g_{\mu\nu} \right)$$

labels

check on sign & indices

$$\vec{J} = \vec{\sigma}/2 \quad \vec{K} = -i\vec{\sigma}/2$$

for problem: Add second field! χ w/ $Q\chi = -Q\psi$

$$L_{\chi} = i\chi_{\alpha}^{\dagger} (\sigma^{\mu})^{\alpha\beta} \partial_{\mu} \chi_{\beta} + i\psi_{\dot{\alpha}}^{\dagger} (\bar{\sigma}^{\mu})^{\dot{\alpha}\beta} \partial_{\mu} \psi_{\beta} + m \left(\psi^{\dagger} \chi + \chi^{\dagger} \psi \right)$$

Dirac Mass term

chose χ, ψ to both be LH,

define $\bar{\Psi}_D \equiv \begin{pmatrix} \psi_{\dot{\alpha}} \\ \chi^{\dagger\beta} \end{pmatrix} \begin{matrix} Q\psi \\ Q\chi = -Q\psi \end{matrix}$

define $\bar{\Psi}_D \equiv \bar{\Psi}_D^{\dagger} \gamma^0 \rightarrow i\bar{\Psi}_D \gamma^{\mu} \partial_{\mu} \Psi_D + m\bar{\Psi}_D \Psi_D$

$$\begin{pmatrix} 0 & 1 & 2 & 2 \\ 1 & 0 & 2 & 2 \\ 1 & 2 & 0 & 2 \\ 1 & 2 & 2 & 0 \end{pmatrix} \quad \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

for four vectors:

$$\left(\frac{\sigma}{\vec{\sigma}}\right)^\mu_{\nu} = i(g^\mu_{\nu} g_{\rho\sigma} - g^\mu_{\rho} g_{\sigma\nu})$$

labels

check on sign & indices

$$\vec{J} = \vec{\sigma}/2 \quad \vec{S} = (\sigma^1, \sigma^2, \sigma^3)$$

$$\vec{K} = -i\vec{\sigma}/2$$

for problem: Add second field! χ w/ $Q_\chi = -Q_\psi$

$$L_e = i\chi^\dagger_\alpha (\sigma^\mu)^\alpha_\beta \not{D}_\mu \chi_\beta + i\psi^\dagger_{\dot{\alpha}} (\bar{\sigma}^\mu)^{\dot{\alpha}\beta} \not{D}_\mu \psi_\beta + m(\psi^\dagger \chi + \chi^\dagger \psi)$$

Dirac Mass term

chose χ, ψ to both be LH,

define $\bar{\Psi}_D \equiv \begin{pmatrix} \psi_\alpha \\ \chi^\dagger_{\dot{\alpha}} \end{pmatrix}^{Q_\psi}$

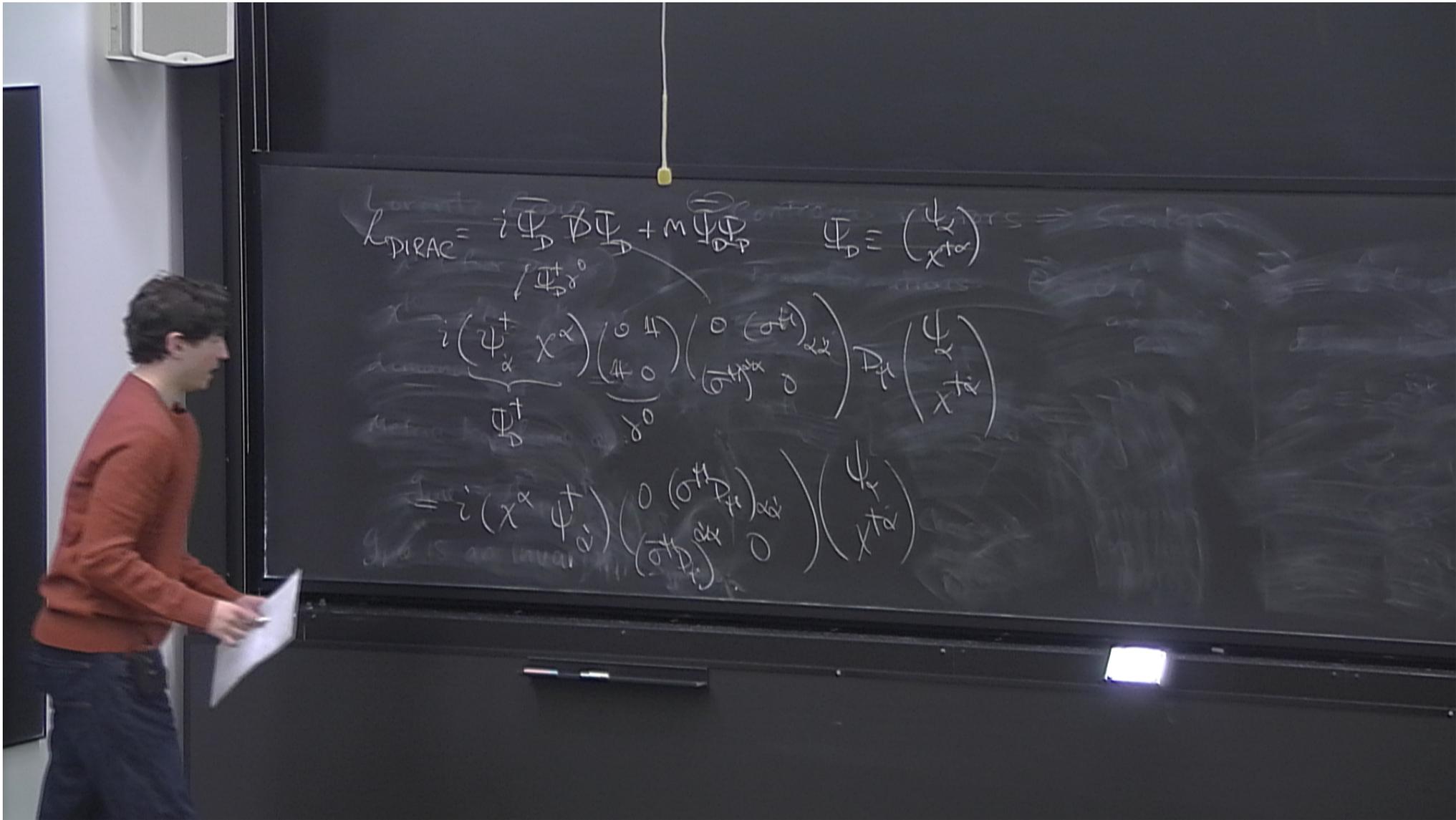
$Q_\psi = -Q_\chi$

define $\bar{\Psi}_D \equiv \begin{pmatrix} \psi^\dagger_{\dot{\alpha}} \\ \chi_\alpha \end{pmatrix}^{Q_\psi}$

$\rightarrow i\bar{\Psi}_D \not{D} \Psi_D + m\bar{\Psi}_D \Psi_D$

$\begin{pmatrix} 0 & 1 & 2 & 2 \\ 1 & 0 & 2 & 2 \\ 1 & 2 & 0 & 2 \end{pmatrix}$

$\begin{pmatrix} 0 & 1 & 1 \\ 1 & 0 & 1 \\ 1 & 1 & 0 \end{pmatrix}$



$$\begin{aligned}
 \mathcal{L}_{\text{DIRAC}} &= i \bar{\Psi}_D \not{\partial} \Psi_D + m \bar{\Psi}_D \Psi_D & \Psi_D &= \begin{pmatrix} \psi_\alpha \\ \chi^{\dot{\alpha}} \end{pmatrix} \\
 &= i \begin{pmatrix} \psi_\alpha^\dagger & \chi^{\dot{\alpha}} \end{pmatrix} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \begin{pmatrix} 0 & (\sigma^\mu)_{\alpha\dot{\beta}} \\ (\sigma^\mu)^{\dot{\alpha}\beta} & 0 \end{pmatrix} \begin{pmatrix} \psi_\alpha \\ \chi^{\dot{\alpha}} \end{pmatrix} \\
 &= i \begin{pmatrix} \chi^{\dot{\alpha}} & \psi_\alpha^\dagger \end{pmatrix} \begin{pmatrix} 0 & (\sigma^\mu)_{\alpha\dot{\beta}} \\ (\sigma^\mu)^{\dot{\alpha}\beta} & 0 \end{pmatrix} \begin{pmatrix} \psi_\alpha \\ \chi^{\dot{\alpha}} \end{pmatrix}
 \end{aligned}$$