

Title: Flux Compactifications Grow Lumps

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Abstract: The simplest flux compactifications are highly symmetricâ€”a q-form flux is wrapped uniformly around an extra-dimensional q-sphere. I will discuss a family of solutions that break the internal $SO(q+1)$ symmetry of these solutions down to $SO(q)\tilde{\times}Z_2$, and show that often at least one of them has lower vacuum energy, larger entropy, and is more stable than the symmetric solution. I will describe the phase diagram of lumpy solutions and provide an interpretation in terms of an effective potential. Finally, I will provide evidence that the perturbatively stable vacua have a non-perturbative instability to spontaneously sprout lumps; generically this new decay is exponentially faster than all other known decays of the model.

Flux Compactifications Grow Lumps

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Outline

- The Toy Model
- Symmetric Solutions
 - Properties
 - Stability
- Warped Solutions
- Effective Potential Description
- Non-Perturbative Decays
- Future Work

The Toy Model Freund, Rubin (1980)

$$S = \int d^D x \sqrt{-g} \left[\frac{1}{2} M_D^{D-2} \mathcal{R} - \frac{1}{2q!} F_q^2 - \Lambda_D \right] \quad (p \geq 3, q \geq 2)$$

- q -dimensional internal manifold wrapped by a q -form flux F_q
- $p = D - q$ extended dimensions

The Toy Model: Motivations

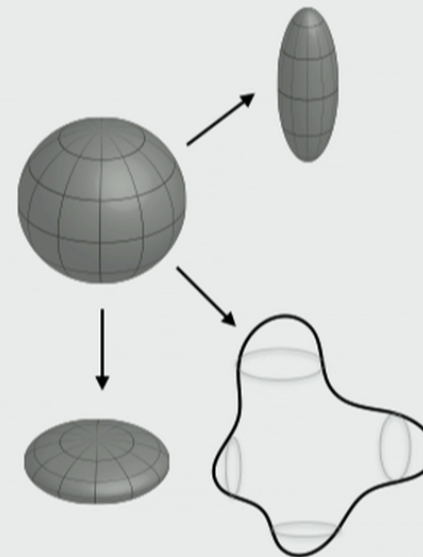
“Simplest model of stabilized extra dimensions”
(Repulsive flux + attractive curvature can form stable minimum)

- One of several toy models often invoked to gain intuition about the much more complicated string landscape
- We should fully study this simple model, then try to understand if its physical implications carry over to the more general case
- An extremely simple physical system (around for > 50 years!), with some rich and tractable corners of phase space that remain to be understood!

The Toy Model

The take-home message:

Even this “simplest” model is actually far from simple. There are symmetry-breaking effects whereby a symmetrically shaped internal manifold generically develops lumps, either perturbatively or non-perturbatively.



Symmetric Solution

$$S = \int d^D x \sqrt{-g} \left[\frac{1}{2} M_D^{D-2} \mathcal{R} - \frac{1}{2q!} F_q^2 - \Lambda_D \right]$$

Direct product total space \Rightarrow both components are Einstein spaces

$$\mathcal{M}_p \times S^q$$

“Freund-Rubin model”

Symmetric Solution

Metric: $ds^2 = L^2 ds_p^2 + R^2 d\Omega_q^2$

Effective c.c.: $\frac{\Lambda_{\text{eff}}}{M_p^{p-2}} = \frac{(p-1)(p-2)}{2L^2}$

Conserved flux: $n \equiv \int_{S^q} F_q = \rho \text{Vol}_{S^q}$

flux number (quantized) flux density

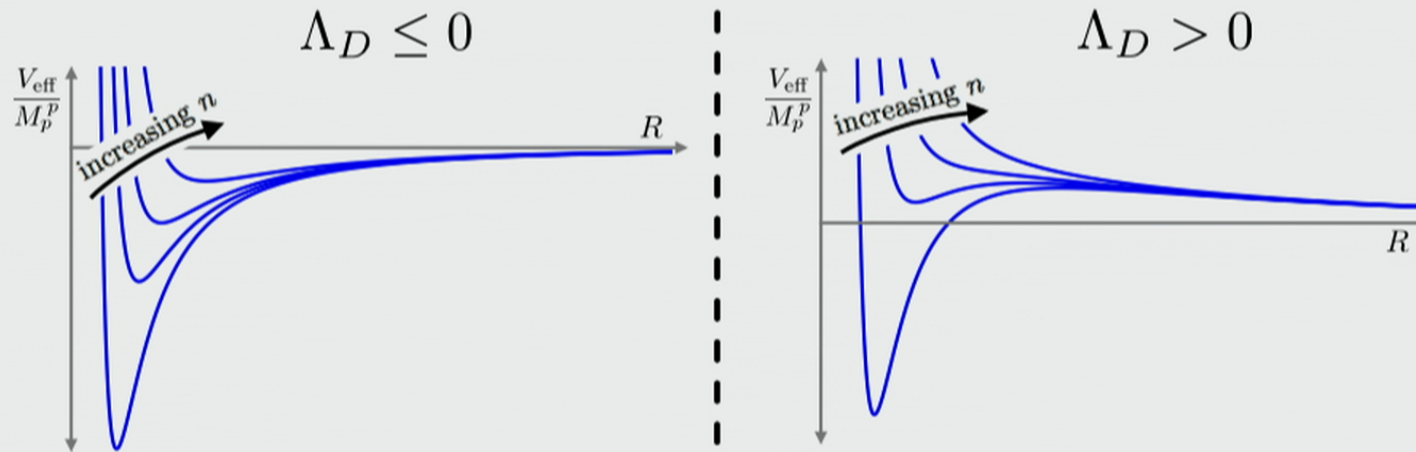
Symmetric Solution: Effective Potential

Treat sphere as fixed shape and vary radius as field R :

$$\frac{V_{\text{eff}}(R)}{M_p^p} \sim \left(\frac{1}{R}\right)^{2q/(p-2)} \left[\frac{n^2}{R^{2q}} - \frac{1}{R^2} + \Lambda_D \right]$$

conversion from D- to p-
dimensional Planck mass

flux (repulsive) **curvature** (attractive) **c.c.** (repulsive)



Symmetric Solution: Stability

Metric: $g_{MN} + h_{MN}$

Flux: $A_{\alpha_1 \dots \alpha_{q-1}} + \delta A_{\alpha_1 \dots \alpha_{q-1}}$

\nearrow \nwarrow
 Background Fluctuation

The procedure:

- Expand fluctuations in complete basis of forms
- Expand action to quadratic order
- Integrate over internal manifold
- Diagonalize, then read off spectrum and stability

Shape mode sector: Expand in spherical harmonics $Y_{lm}(\theta, \phi)$

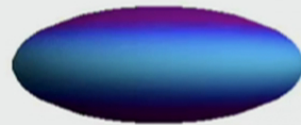
Symmetric Solution: Stability

h_ℓ and a_ℓ mix! \rightarrow physical shape mode is linear combination $\psi_\pm^{(\ell)}$

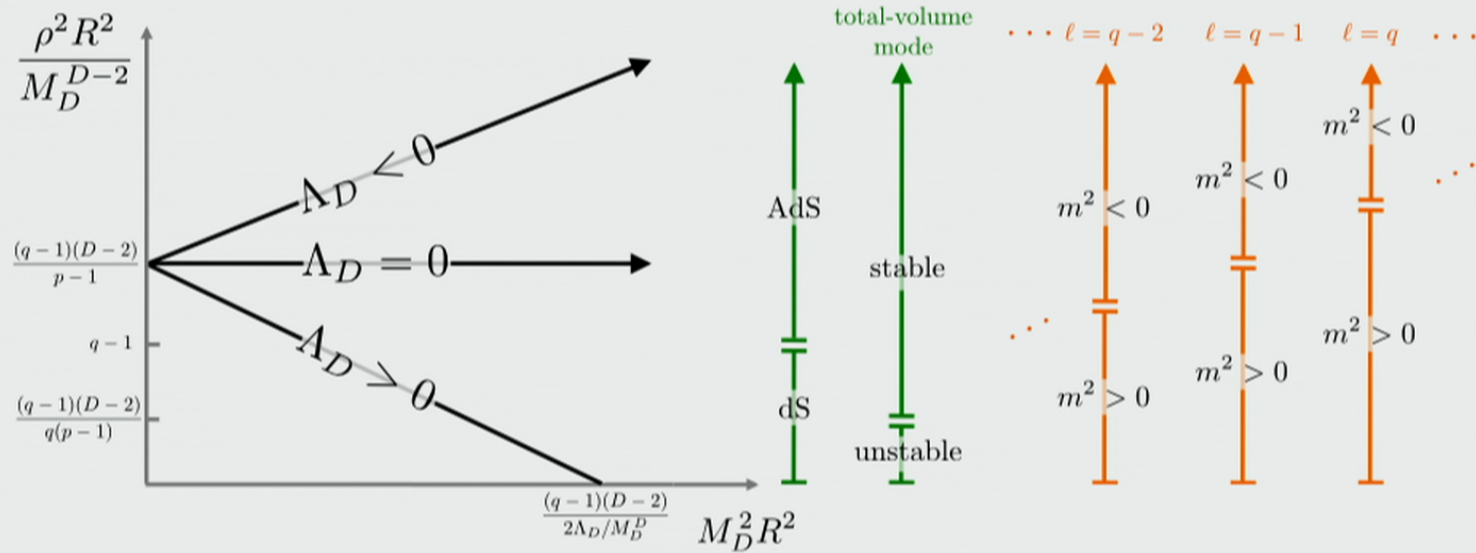
$\psi_-^{(\ell)}$ is the “danger mode” that sometimes has negative mass squared

\Rightarrow **shape mode instability**

For $\psi_-^{(\ell)}$, flux density larger where radius is larger
(opposite for $\psi_+^{(\ell)}$)



Stability



Danger mode has negative mass squared iff:

$$\rho^2 R^2 = \frac{1}{p-1} [(q-1)(D-2) - 2\Lambda_D R^2] > \frac{D-2}{2(p-1)(q-2)} [\ell(\ell+q-1) - 2(q-1)]$$

Warped Solutions

Kinoshita (2007)

Kinoshita, Mukohyama (2012)

$$S = \int d^D x \sqrt{-g} \left[\frac{1}{2} M_D^{D-2} \mathcal{R} - \frac{1}{2q!} F_q^2 - \Lambda_D \right]$$

A new set of lumpy, warped solutions can be found numerically for $q \geq 3$!

$\mathcal{M}_p \times \text{warped}$ 

Let's explore their physics...

Warped Solutions

Metric: $ds^2 = \Phi(\theta)^2 ds_p^2 + R(\theta)^2 d\Omega_q^2$

Flux ansatz: $F_{\alpha_1 \dots \alpha_q} = Q \Phi^{-p}(\theta) \epsilon_{\alpha_1 \dots \alpha_q}$

To find these solutions, shoot numerically from the equator varying Q , Λ_D and imposing regularity at poles, evenness about equator

$$\mathrm{SO}(q+1) \rightarrow \mathrm{SO}(q) \times \mathbb{Z}_2$$

Warped Solutions: Sample Solutions

Symmetric

$R(\theta)$



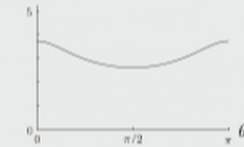
$\Phi(\theta)$



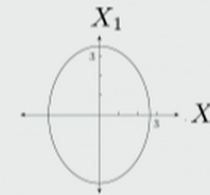
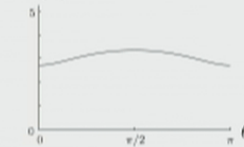
0 extrema

Football

$R(\theta)$



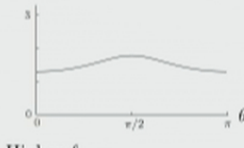
$\Phi(\theta)$



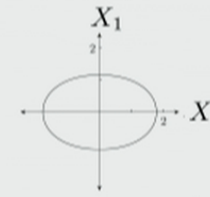
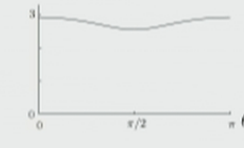
1 extremum

M&M

$R(\theta)$



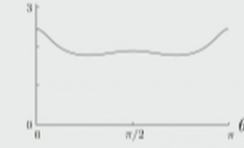
$\Phi(\theta)$



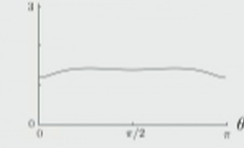
1 extremum

Higher ℓ

$R(\theta)$



$\Phi(\theta)$



3 extrema

Warped Solutions: Game Plan

- Consider the number of extrema as a classifier.
- For simplicity, we restrict to the 1 extremum solutions.
- To study these solutions, we use numerical techniques to map out their phase space. Later, we will return to motivate this picture using an effective potential description.
- First we parametrize the solutions using a convenient order parameter, the ellipticity.

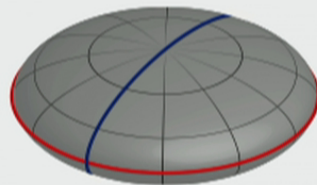
Warped Solutions: Ellipsoidal Case

(1 extremum)

Order parameter: Ellipticity

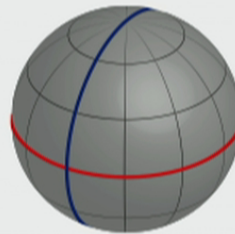
$$\varepsilon \equiv \frac{2 \times (\text{distance from north to south pole})}{\text{distance around equator}} = \frac{2 \int_0^\pi R(\theta) d\theta}{2\pi R(\theta = \pi/2)}$$

M&M



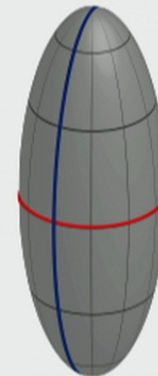
$\varepsilon < 1$

Symmetric



$\varepsilon = 1$

Football

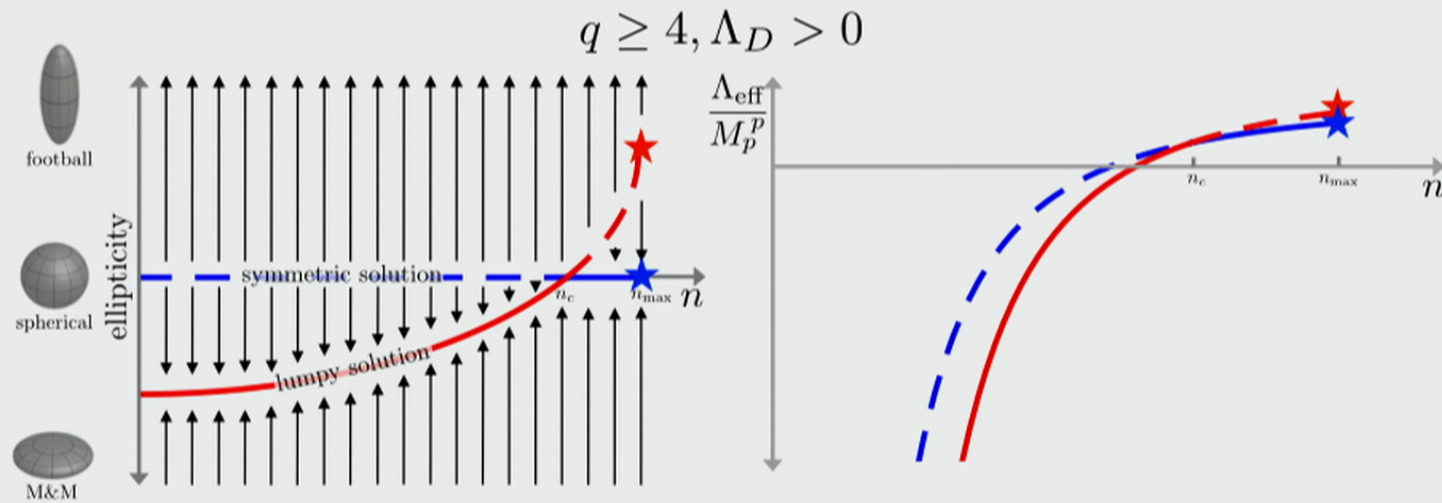


$\varepsilon > 1$

Warped Solutions: Ellipsoidal Case

(1 extremum)

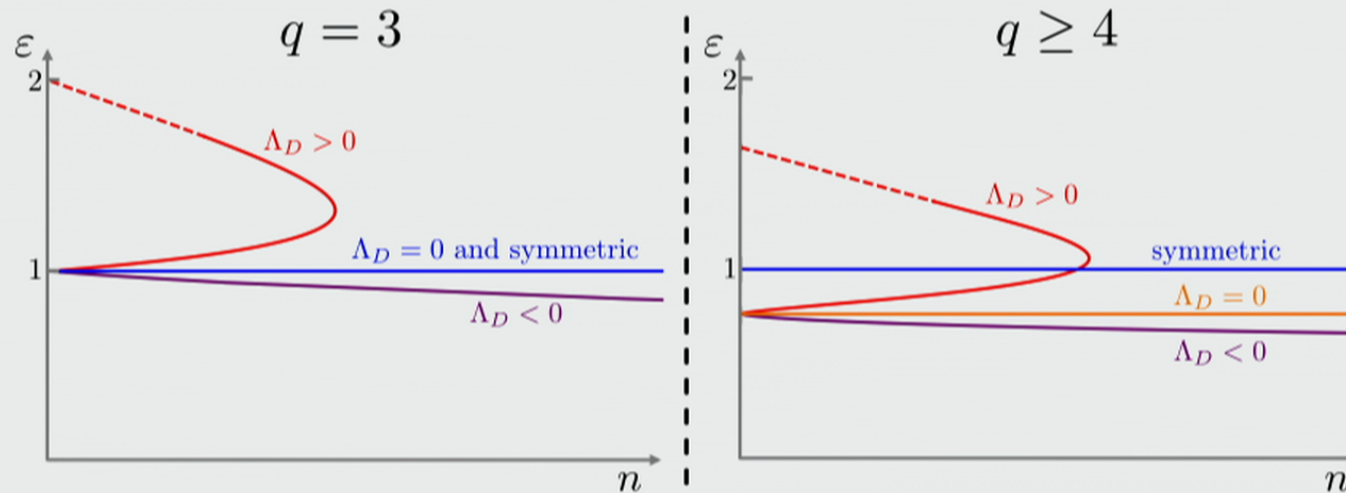
Cartoon Distillation:



Warped Solutions: Ellipsoidal Case

(1 extremum)

Full phase diagram:



Scaling properties of EOM:

$$\Lambda_D \rightarrow \alpha^{-2} \Lambda_D,$$

$$Q \rightarrow \beta^{p-1} Q \quad (\Lambda_D = 0)$$

Warped Solutions: Ellipsoidal Case

(1 extremum)

The underlying principles:

- When the $\ell = 2$ ‘danger mode’ has a negative mass squared, the ellipsoidal solution is M&M shaped and energetically favored. When the mode has a positive mass squared, the ellipsoidal solution is football-shaped and energetically disfavored.
- The $n \rightarrow 0$ behavior is independent of Λ_D
- Each Freund-Rubin solution is accompanied by a single ellipsoidal solution with the same flux n

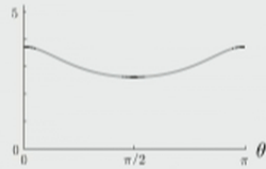
Warped Solutions: Ellipsoidal Case

(1 extremum)

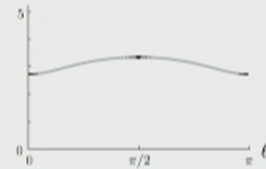
1 extremum solution $\Leftrightarrow \ell = 2$ instability of the symmetric solution

Football

$R(\theta)$

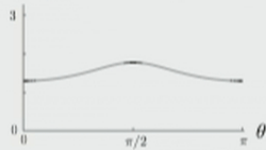


$\Phi(\theta)$

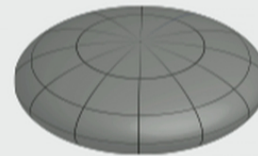
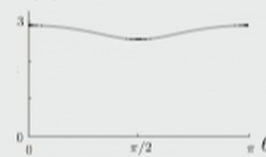


M&M

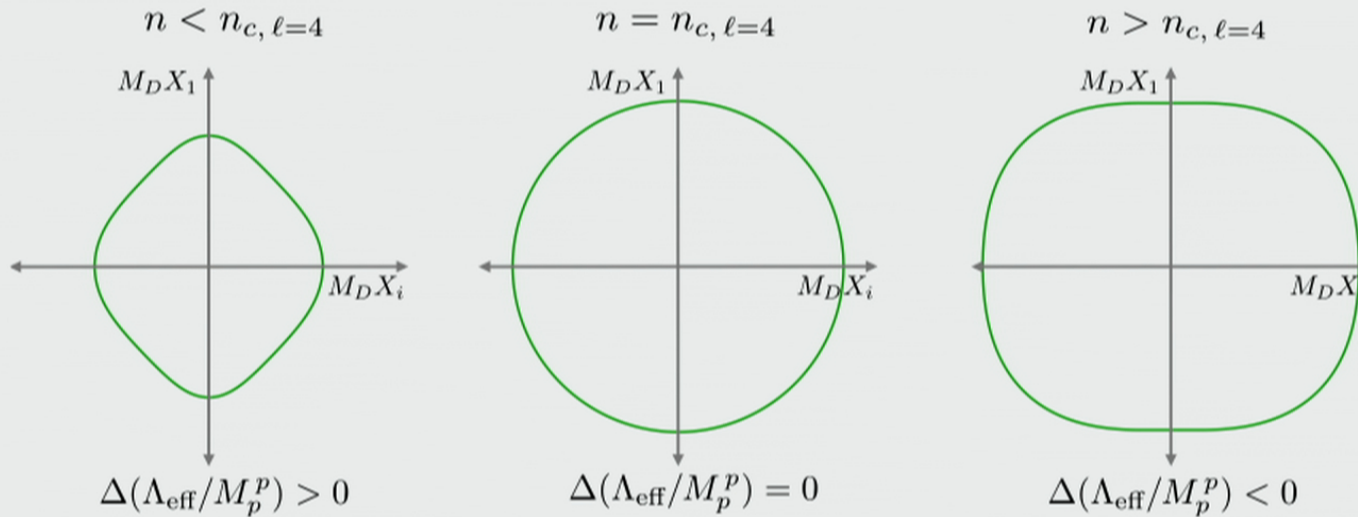
$R(\theta)$



$\Phi(\theta)$



Warped Solutions: Higher- ℓ



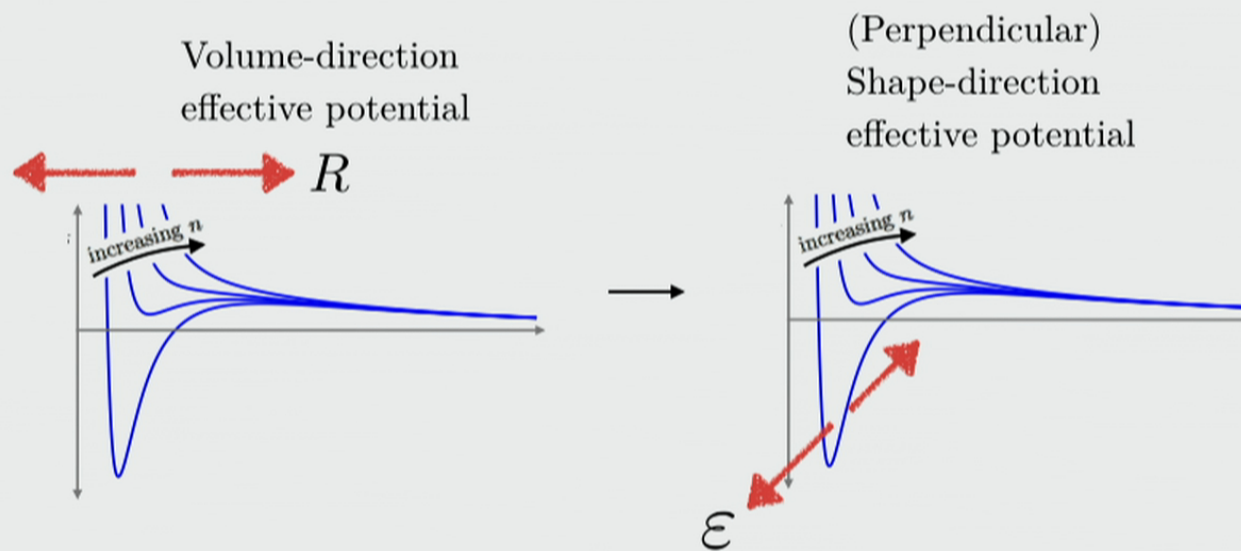
Same principles hold!

Conjecture:

$\ell - 1$ extremum \Leftrightarrow ℓ -mode instability of the solution

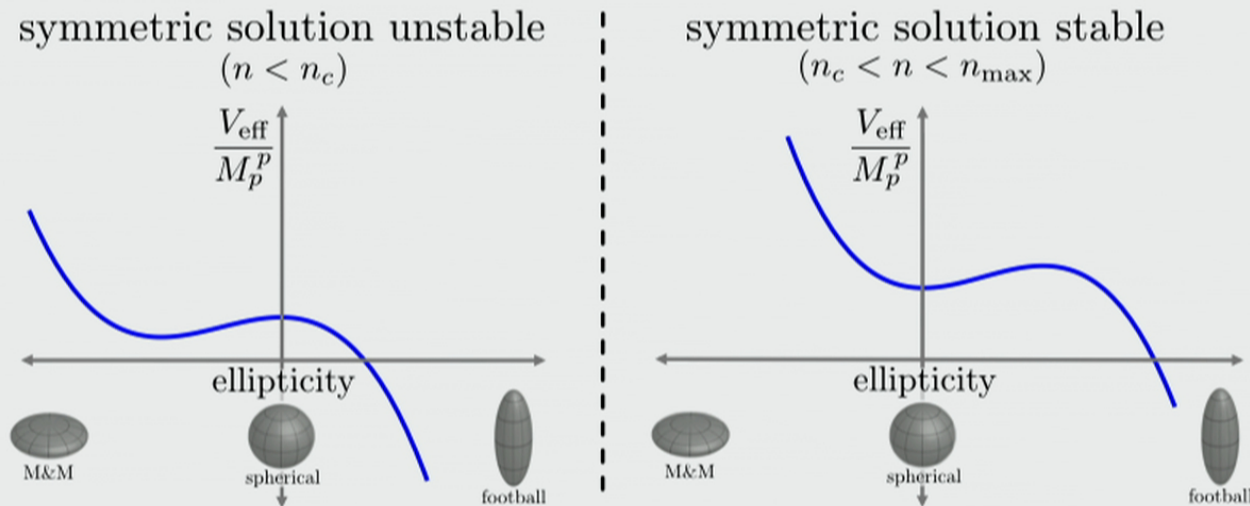
Effective Potential Derivation

Now that we've mapped out the phase space of some warped solutions, we will argue for a neat effective potential picture that encapsulates the physics.



Warped Solutions: Effective Potential

Effective potential in shape direction:



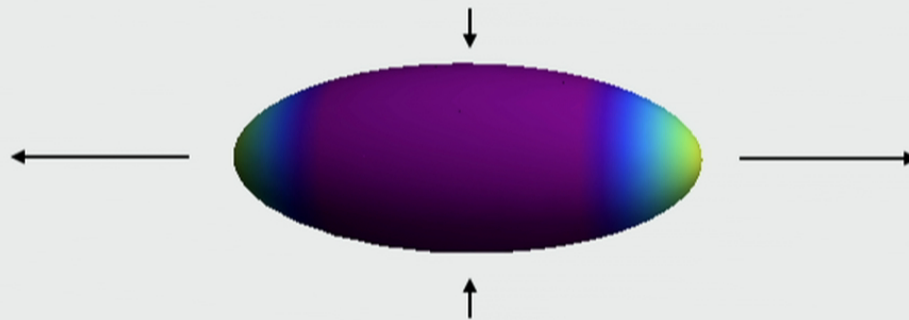
Critical value of n where symmetric flips from positive to negative mass squared + fixed asymptotics

\Rightarrow New (for us, lumpy) solutions are “ejected”

Warped Solutions: Effective Potential

Asymptotics:

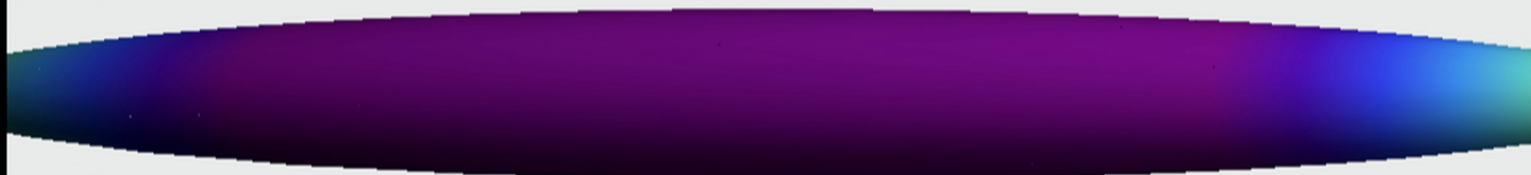
- $V \rightarrow +\infty$ (M&M) and $V \rightarrow -\infty$ (football)
- Turns around in either or both directions
If so, should have found evidence of additional solutions
- Finite asymptote in either or both directions
Long flat section with no symmetry to protect it, we find unlikely



Warped Solutions: Effective Potential

Asymptotics:

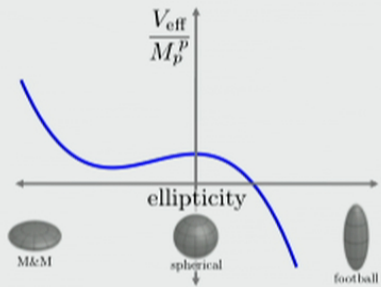
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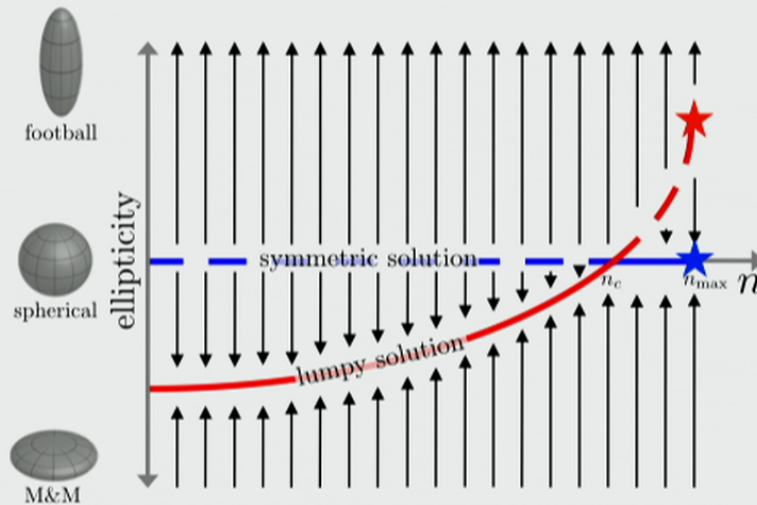
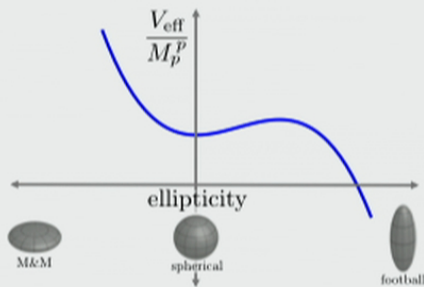
Warped Solutions: Effective Potential

Summary: Consistent effective potential description that captures all the physics of our phase diagram

symmetric solution unstable
($n < n_c$)



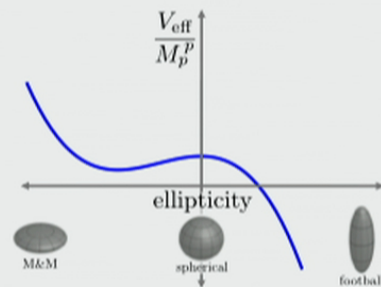
symmetric solution stable
($n_c < n < n_{\text{max}}$)



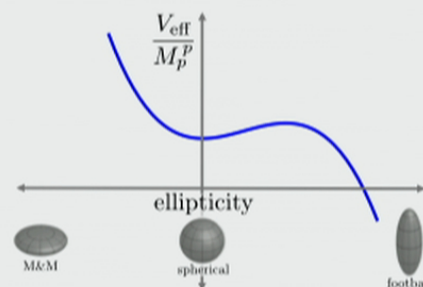
Warped Solutions: Effective Potential

The classical stability of the symmetric solution determines the form of the lumpy solutions: their shape, stability and whether they are dominant or subdominant.

symmetric solution unstable
($n < n_c$)



symmetric solution stable
($n_c < n < n_{\text{max}}$)



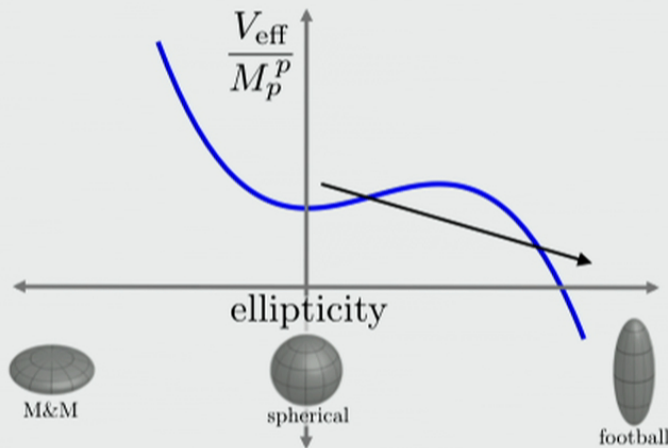
Non-Perturbative Instabilities

New non-perturbative instability: Tunnel to infinite football!

Typical mass squareds in shape direction are a couple orders of magnitude smaller than in total-volume direction.

symmetric solution stable
($n_c < n < n_{\max}$)

\Rightarrow Shape-mode tunneling is fastest decay (vs. e.g. decompactification).



$p = q = 4, \Lambda_D > 0 :$

- 97% of FR perturbatively unstable
- 2.97% decay to footballs

Future Work

- Study spontaneous symmetry breaking along different axes at each point in space-time, and whether this leads to topologically nontrivial configurations
- Study the evolution in this shape direction potential further from sphericity
- Generalize to some more realistic models
- Do the full stability calculation around the warped product solutions.

Conclusions

- There exist additional lumpy solutions to the Freund-Rubin equations that can be more stable than the symmetric branch.
- There is a symmetry-breaking effect where the internal sphere generically rolls to these new solutions.
- Even when the symmetric solution is stable, we found new non-perturbative decays to lumpiness, including to the infinite football-shaped direction. These are the fastest-yet-known decays for this model.
- We derived a phase diagram and argued all physics is neatly encapsulated by a shape-direction effective potential.