

Title: Nuclear spin precession of noble gases in ultra low magnetic fields

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Abstract: In the low energy regime, precision measurements of spin precession have gained increased attention as an alternative pathway to physics beyond the standard model. These measurements aim at the detection of minute frequency changes superimposed on low Larmor frequencies at extremely weak magnetic fields. Such measurements require an effective shielding against the magnetic field of the Earth and other perturbations. For measuring the precession frequency with high precision, a long lifetime of the precessing nuclear magnetization is required, thus the homogeneity of the applied field is a crucial parameter. In addition, criteria are needed that unambiguously distinguish magnetic artifacts from the non-magnetic exotic interactions that we search for. This can be accomplished by the concept of co-magnetometry, i.e., by simultaneous measuring the precession of two nuclear species such as ^3He and ^{129}Xe . Yet another kind of co-magnetometry is the use of SQUIDs for monitoring the spin precession. SQUIDs are magnetic field detectors of their own kind, which can measure the oscillating magnetic field generated by the precessing nuclear magnetic moment as well as the magnetic dc background field. In this presentation, I will report on the current state of the art in our lab in measurements of nuclear spin precession of noble gases.

Outline



Experimental Set-up

- Magnetic shielding
- SQUIDs
- Field homogeneity
- ^3He - ^{129}Xe -Comagnetometry

Search for non-magnetic interactions

- XeEDM
- Short range interaction
- Lorentz-Invariance Violation

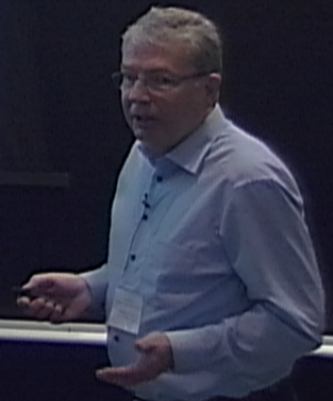
Axion wind

- SQUID based Comagnetometry

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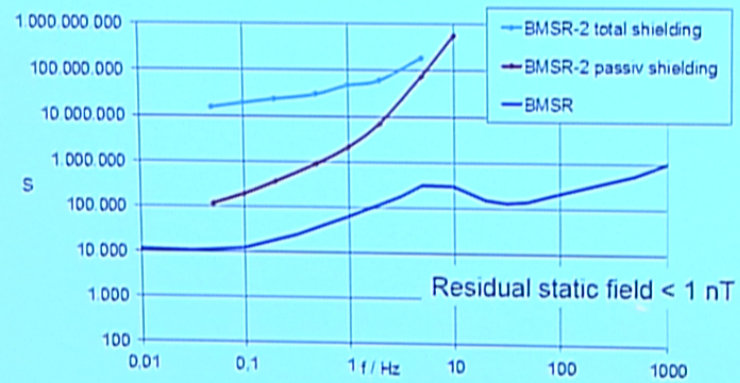
page 2



Experimental: Magnetic shielding



Berlin Magnetically Shielded Room - 2



Shielding factor vs. frequency

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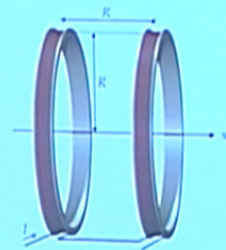
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page 6

Experimental: Magnetic shielding



Helmholtz-Coils inside BMSR-2



+



$\sim 1\mu\text{T}$

$< 1\text{nT}$

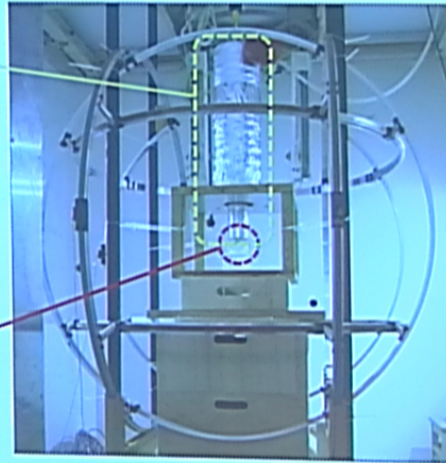
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page 7

Helmholtz-Coils inside BMSR-2

Dewar
filled with liquid He

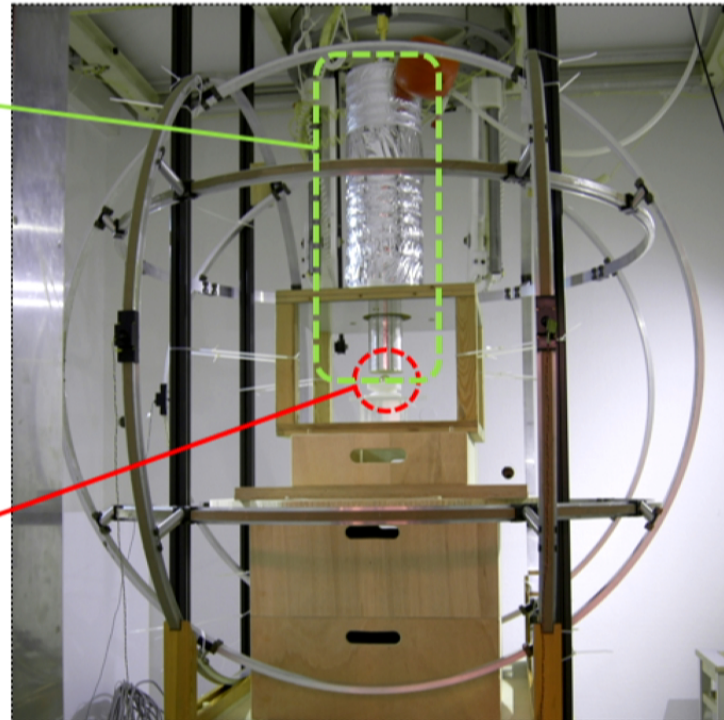


Glass bulb
filled with
hyperpolarized ^3He
and/or ^{129}Xe

Helmholtz-Coils inside BMSR-2

Dewar
filled with liquid He

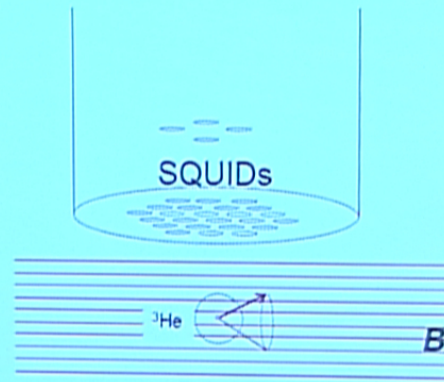
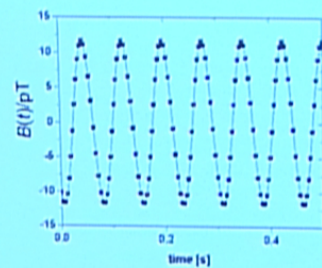
Glass bulb
filled with
hyperpolarized ^3He
and/or ^{129}Xe



Experimental: Field homogeneity



^3He precession as measured by SQUID

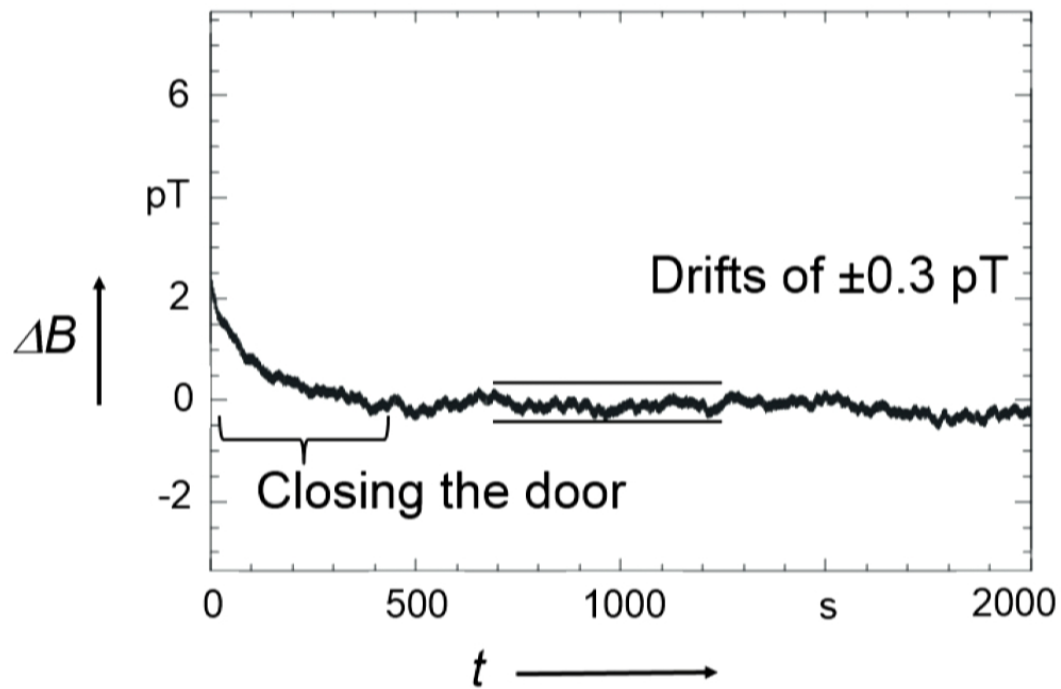


June 19th, 2014

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page 13

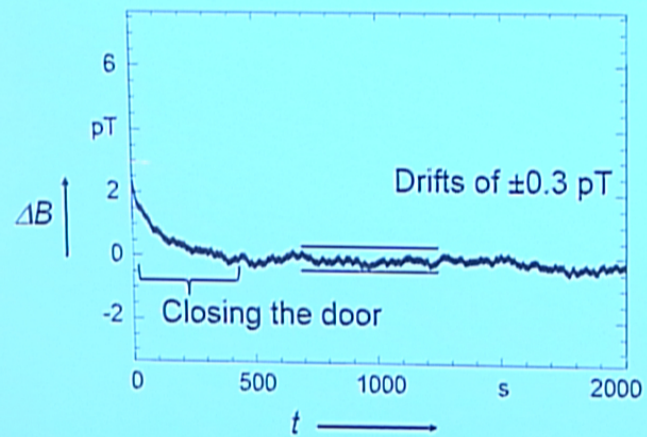
SQUID signal in a field of $\sim 2\mu\text{T}$
after closing the door of the BMSR-2



Experimental: Co-magnetometry



SQUID signal in a field of $\sim 2\mu\text{T}$
after closing the door of the BMSR-2



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page 19

Co-magnetometer:

$$\Delta\omega = \omega_{He} - \frac{\gamma_{He}}{\gamma_{Xe}} \omega_{Xe} = 0 \quad ?$$

In real experiments $\Delta\omega$ is not zero, because:

- uncertainties in γ_{He}/γ_{Xe} :

$$\gamma_{He}/\gamma_{Xe} = 2.7540816 \pm 0.0000002$$

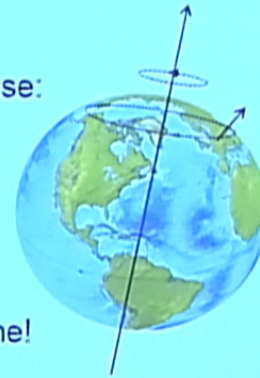
For $\nu_{Xe} = 4$ Hz this may cause a constant drift of up to $\Delta\omega_{\gamma} = 3$ μ rad/s

Experimental: Co-magnetometry



In real experiments $\Delta\omega$ is not zero, because:

- uncertainties in $\gamma_{\text{He}}/\gamma_{\text{Xe}}$
- Earth rotation:



The lab frame is a rotating frame!

$$\Delta\omega_{\text{rot}} = \Omega_{\text{Earth}} \cdot \left(1 - \frac{\gamma_{\text{He}}}{\gamma_{\text{Xe}}}\right) \cos \rho \cos \theta$$

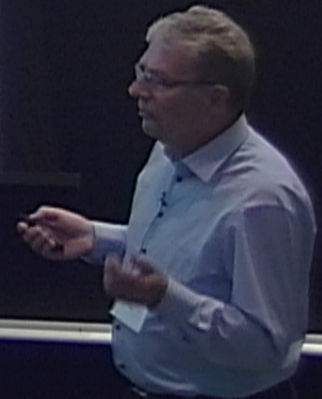
$= 68.7263 \mu\text{rad/s}$

ρ : geograph. latitude
 θ : bearing of B

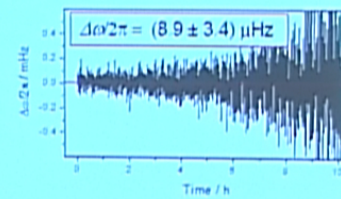
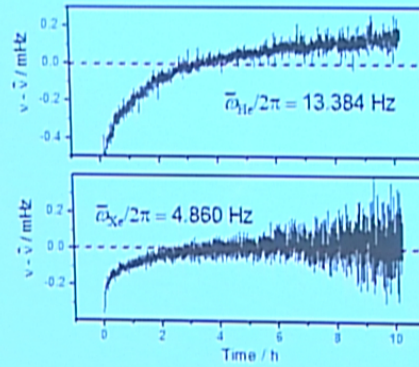
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page 22



Experimental: Co-magnetometry



$$\Rightarrow \Delta\omega = \omega_{He} - \frac{\gamma_{He}}{\gamma_{Xe}} \omega_{Xe} \neq 0$$

\Rightarrow Search for a variation in $\Delta\omega$

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page 23

Search for exotic interactions



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Axion wind

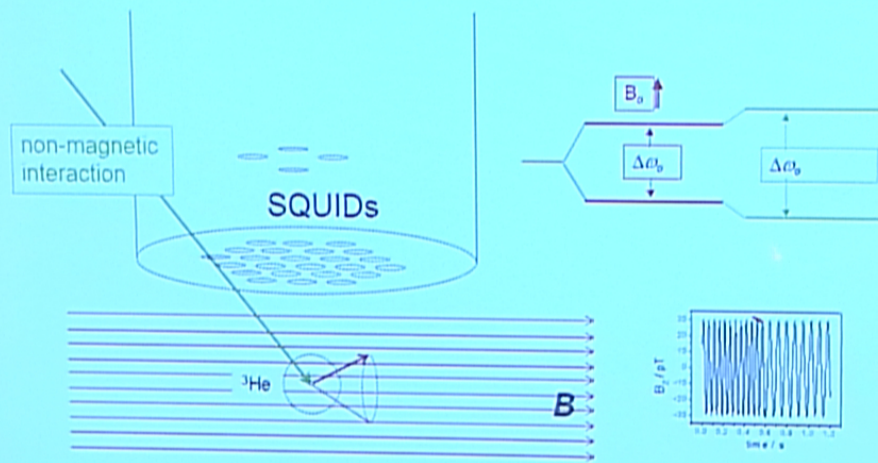
- SQUID based Comagnetometry

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page 24

Search for exotic interactions

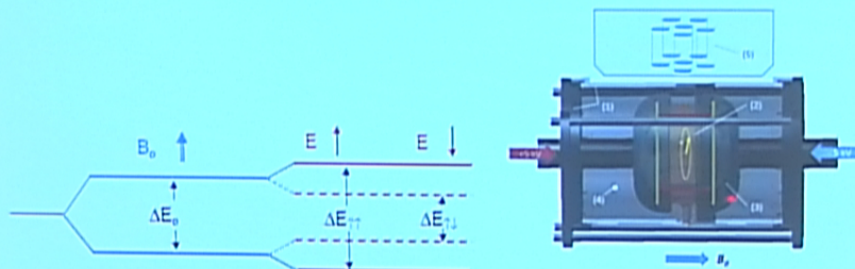


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page 25

Search for ^{129}Xe EDM



$$h \cdot \Delta\nu = h \cdot (\Delta\nu_{\uparrow\uparrow} + \Delta\nu_{\downarrow\downarrow}) = 4 \cdot E \cdot |d_{Xe}|$$

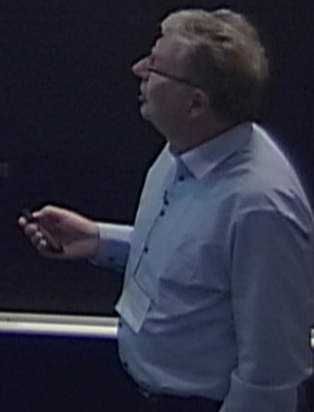
$$\Delta\nu_{\uparrow\uparrow} = \Delta\nu_{\uparrow\uparrow}^{He, EDM} - (\gamma_{He} / \gamma_{Xe}) \cdot \Delta\nu_{\uparrow\uparrow}^{Xe, EDM}$$

$$\Delta\nu_{\downarrow\downarrow} = \Delta\nu_{\downarrow\downarrow}^{He, EDM} - (\gamma_{He} / \gamma_{Xe}) \cdot \Delta\nu_{\downarrow\downarrow}^{Xe, EDM}$$

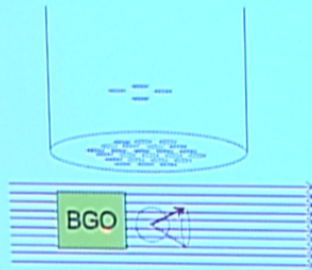
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page 26



Search for axions



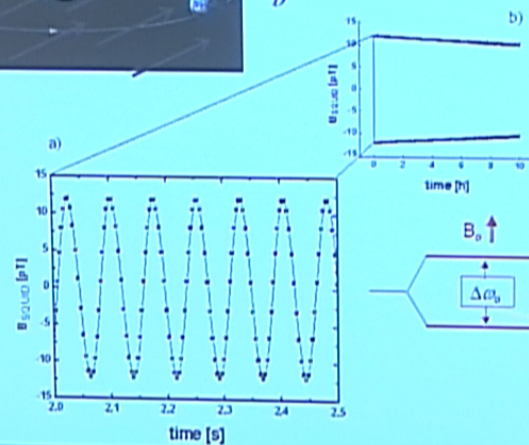
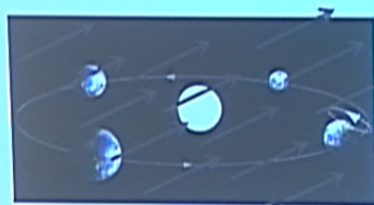
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Tullney et al. (2013) PRL

page 27

Search for Lorentz Invariance Violation



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page 28

Axion wind



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page 29

Experimental Set-up

Magnetic shielding

SQUIDs

Field homogeneity

^3He - ^{129}Xe -Comagnetometry

Search for non-magnetic interactions

XeEDM

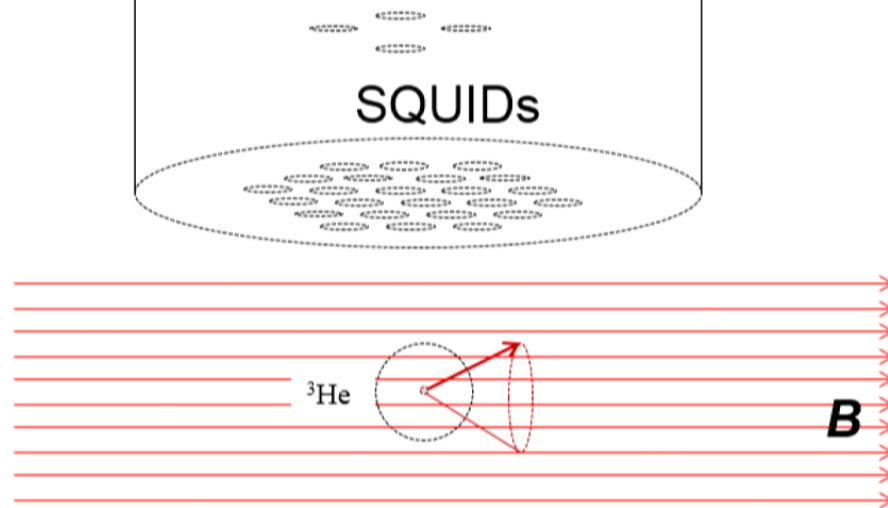
Short range interaction

Lorentz-Invariance Violation

Axion wind

SQUID based Comagnetometry

Axion wind



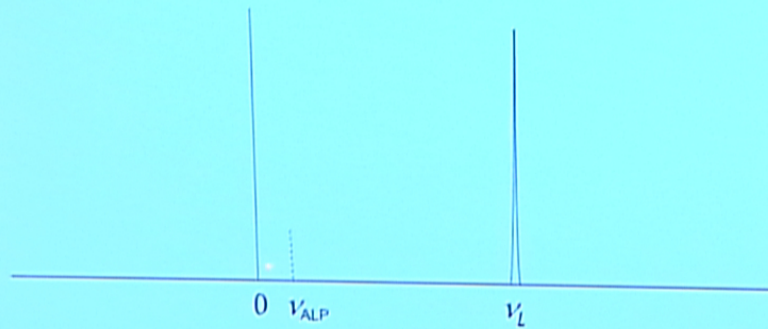
$$v = v_{\text{Larmor}}$$

P. Graham and S. Rajendran, PRD 2011, 2013

Axion wind



Larmor frequency of the noble gas nucleus (^3He or ^{129}Xe): ν_L
Frequency of the axion like particle: ν_{ALP}



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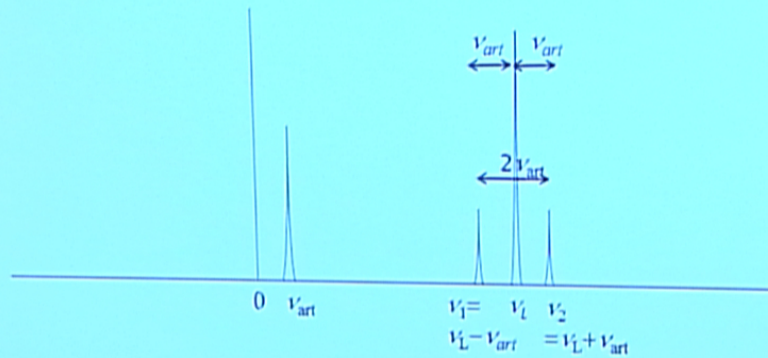
page 33

Axion wind



Signatur of a magnetic artifact (if $v_{art} < v_L$):

→ fourth peak at v_{ext}



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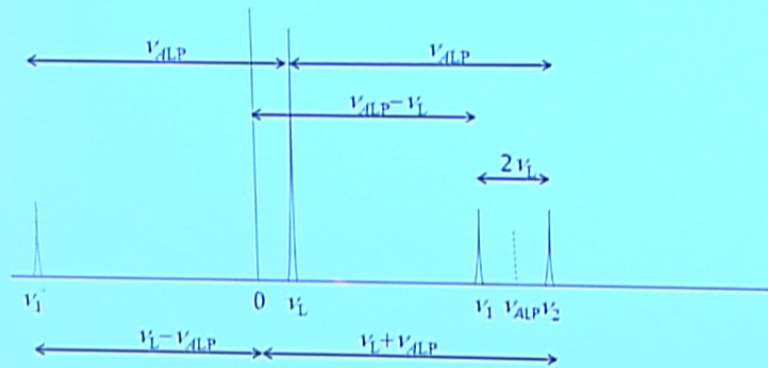
page 35

Axion wind



Alternative Situation: $v_L < v_{ALP}$

Reflection at zero results in doublet $\pm v_L$ around v_{ALP}



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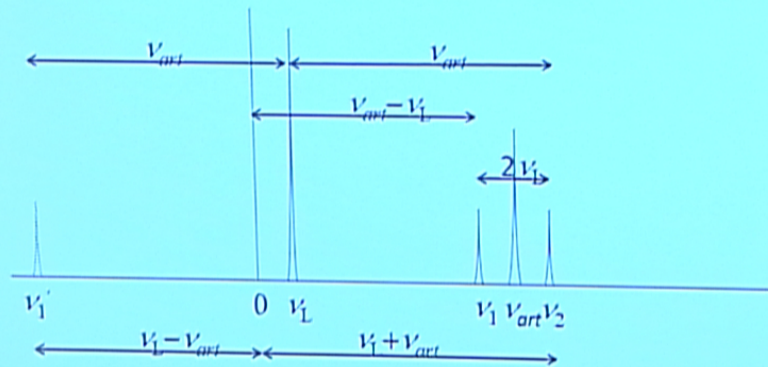
page 38

Axion wind



Signatur of a magnetic artifact

→ fourth peak at v_{art}

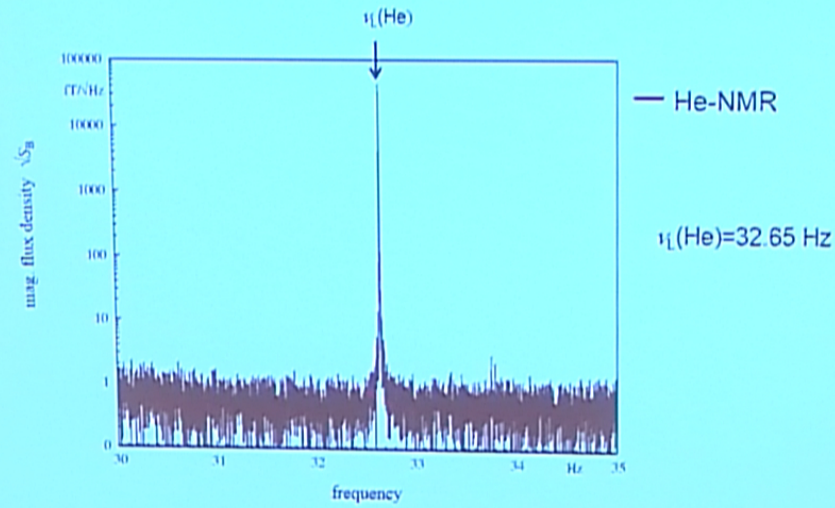


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page 39

Axion wind



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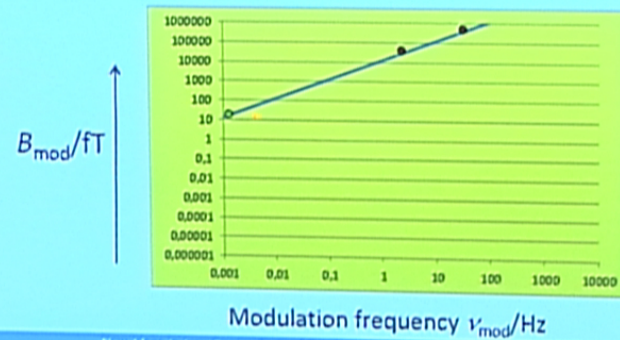
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page 40

Axion wind



From our experimental data we can estimate the range of amplitudes and frequency where interactions can be detected by our current set-up:



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page 47

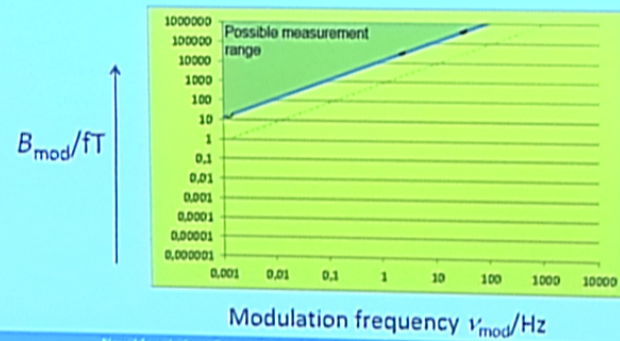
Axion wind



Optimistic outlook:

Increase He-Amplitude
Reduce noise
Decrease distance

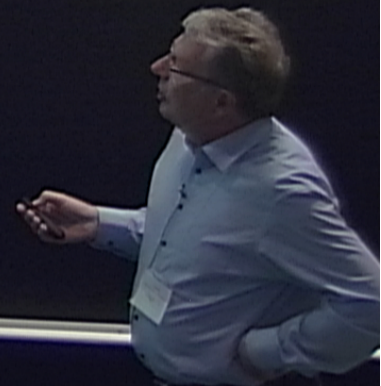
7 pT → 100 pT
2 fT/√Hz → 0.5 fT/√Hz
58 mm → 28 mm
→ ~ x 200



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page 48



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June 19th, 2014

New Ideas in Low-Energy Tests of Fundamental Physics

page 49