

Title: Renormalization Group Analysis of a Non-Fermi Liquid System

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Abstract: We devise a renormalization group analysis for quantum field theories with Fermi surface to study scaling behaviour of non- Fermi liquid states in a controlled approximation. The non-Fermi liquid fixed points are identified from a Fermi surface in $(m+1)$ spatial dimensions, while the co-dimension of Fermi surface is also extended to a generic value. We also study superconducting instability in such systems as a function of dimension and co-dimension of the Fermi surface. The key point in this whole analysis is that unlike in relativistic QFT, the Fermi momentum k_F enters as a dimensionful parameter, thus modifying the naive scaling arguments. The effective coupling constant is found to be a combination of the original coupling constant and k_F .

RG Analysis of a Non-Fermi Liquid System

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Plan of the talk

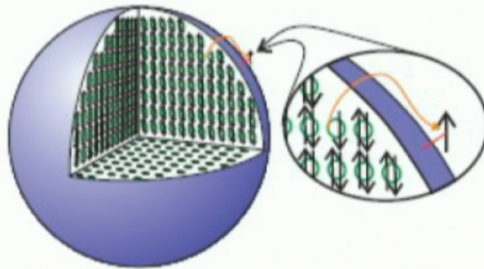
- Prologue → Landau Fermi-Liquid Theory
- Breakdown of FLT → Non-Fermi Liquids
- Ising-Nematic QCP
- Dimensional & Co-dimensional RG
- Critical Dimension
- Beta-Functions & Critical Exponents
- Stable NFL Fixed Point
- Superconducting Instability
- Epilogue

Landau Fermi-Liquid Theory



[*Landau, 50s*]: A finite density of interacting fermions doesn't depend on specific microscopic dynamics of individual systems :-

- **Ground state**: characterized by a sharp Fermi surface (FS) in momentum space
- **Low energy excitations**: weakly interacting quasi-particles around FS



$$(\omega = 0, \quad k_{\perp} \equiv k - k_F = 0)$$



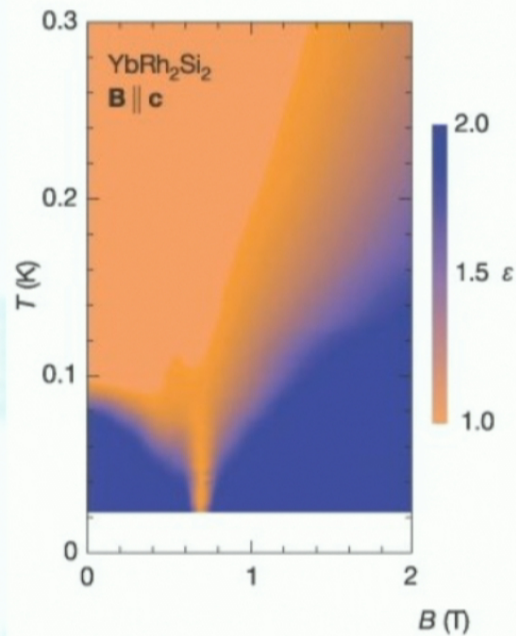
$$G_R(\omega, \vec{k}) = \frac{Z}{\omega - v_F k_{\perp} + i\Gamma}$$

- 1 Quasi-particles have long life time \rightarrow Decay rate $\Gamma \sim \omega^2$
- 2 Electron has a finite over-lap with quasi-particle adiabatically connected to non-interacting Fermi gas \rightarrow quasi-particle wt $Z > 0$

Breakdown of FL Theory



Unusual Scaling Phenomenology



[Custers et al, Nature (2003)]

$$[\rho(T) - \rho_0] \propto T^\epsilon$$

ε = 1 for NFL (yellow)
ε = 2 for FL (blue)

- 1 Calculational framework that replaces FL theory needed.
- 2 QFT of metals • low symmetry + extensive gapless modes need to be kept in low energy theories • less well understood compared to relativistic QFTs.

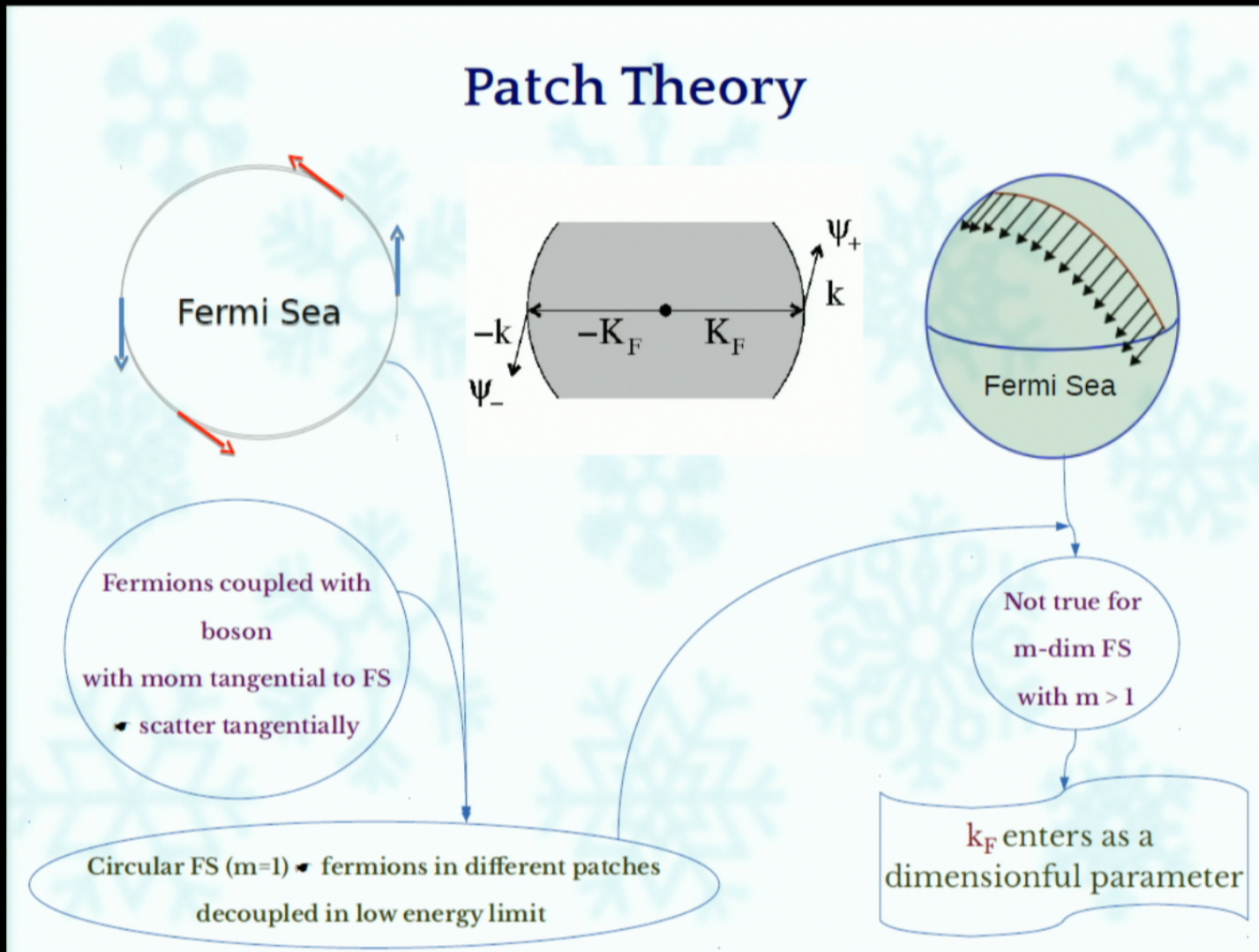
Dimension as A Tuning Parameter

- Couplings irrelevant for $d > \text{upper critical dim } d_c$.
 - For $d < d_c$ couplings become slightly relevant & theories flow into a perturbative fixed point.
 - Modify spacetime dim continuously to gain a controlled access to NFL states
 - Extend
 - ① FS dim [Chakravarty, Norton and Syljuasen (1995) ← issues with k_F]
- or
- ② FS co-dim [Senthil and Shankar (2009); Dalidovich and Lee (2013)]

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Patch Theory

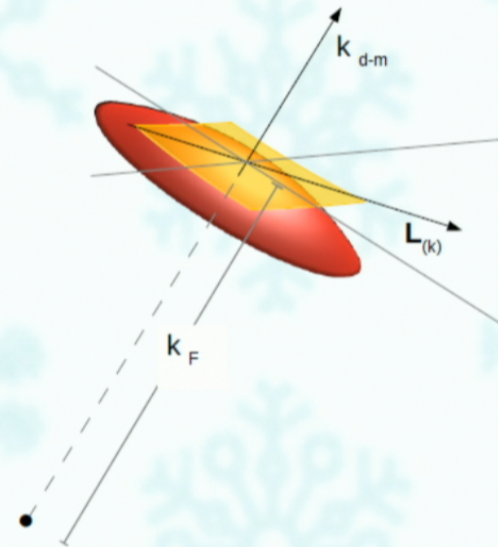


Role of “ k_F ”

- Dalidovich & Lee [Phys. Rev. B 88, 245106 (2013)]
 - DR for circular FS extending **co-dim** to a generic value with $m=1$ fixed
 - $(2+1)$ -d ISN QCP described by a stable NFL state slightly below $d_c = 5/2$.
- We devise DR extending both **dim** & **co-dim** • FS with $m>1$ included naturally.
- k_F enters as a dimensionful parameter modifying naïve scaling arguments for $m>1$.
- The global FS should be considered. But patch coords still allow to extract leading order effect of the global couplings of different patches through k_F .

Generic Fermi Surface

Patch of m -dim FS
of arbitrary shape



- At a chosen point on FS : $k_{d-m} \perp$ local S^m \leftarrow its magnitude measures deviation from k_F .
- $L_{(k)} = (k_{d-m+1}, k_{d-m+2}, \dots, k_d)$ \leftarrow tangential along the local S^m .

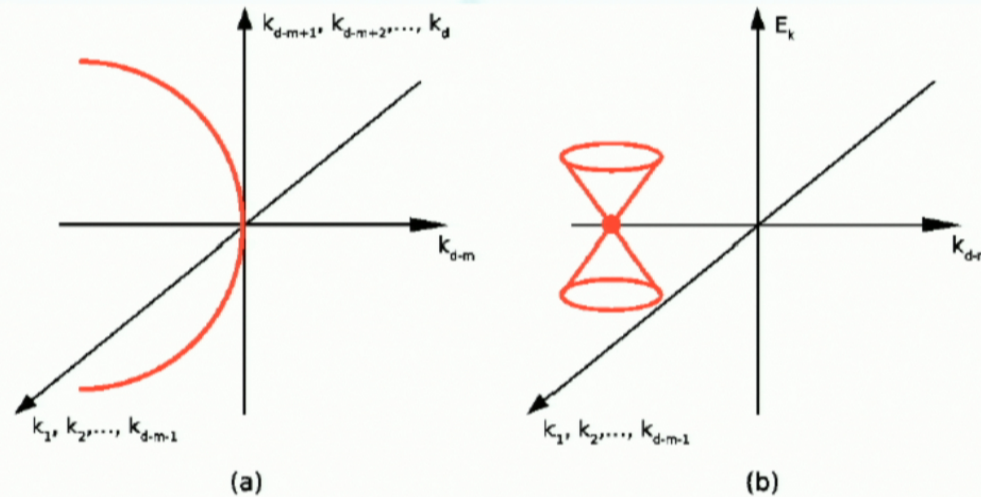
Action

2 patches of m -dim FS
coupled with one critical boson
in $(m+1)$ -space & one time dim:



$$\begin{aligned} S &= \sum_{s=\pm} \sum_{j=1}^N \int \frac{d^{m+2}k}{(2\pi)^{m+2}} \psi_{s,j}^\dagger(k) \left[ik_0 + sk_{d-m} + \vec{L}_{(k)}^2 \right] \psi_{s,j}(k) \\ &+ \frac{1}{2} \int \frac{d^{m+2}k}{(2\pi)^{m+2}} \left[k_0^2 + k_{d-m}^2 + \vec{L}_{(k)}^2 \right] \phi(-k) \phi(k) \\ &+ \frac{e}{\sqrt{N}} \sum_{s=\pm} \sum_{j=1}^N \int \frac{d^{m+2}k d^{m+2}q}{(2\pi)^{2m+4}} \phi(q) \psi_{s,j}^\dagger(k+q) \psi_{s,j}(k) \end{aligned}$$

Line of Dirac Points

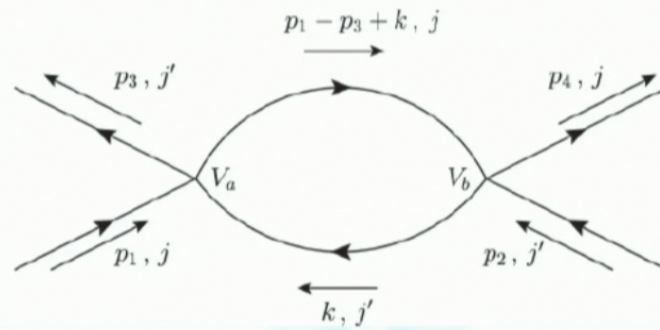
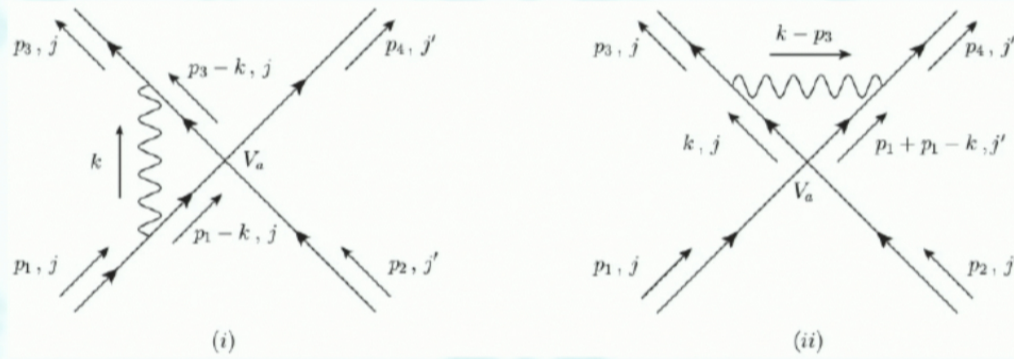


(a) m -dim FS embedded in d -dim mom space.

(b) Spinor has 2 bands:
$$E_k = E_F \pm \sqrt{\sum_{i=1}^{(d-m-1)} k_i^2 + \delta_k^2}$$

For each $\mathbf{L}_{(k)}$ Dirac point $\equiv (k_1=0, k_2=0, \dots, k_{d-m} = -(\mathbf{L}_{(k)})^2)$ around which energy disperses linearly like a Dirac fermion in the $(d-m)$ -dim subspace.

One-Loop Diagrams



Beta-Fns for V_a 's

- One-loop divergent terms proportional to

$$\frac{e_{eff} \mu^{d_v}}{\epsilon} V_b \quad \text{for} \quad d = d_c - \epsilon;$$

$$\frac{k_F^{\frac{m}{2}} \mu^{2d_v}}{\epsilon} V_b V_c \quad \text{for} \quad d - m = 1 - \epsilon.$$

- Both valid only for $d=3 - \epsilon$, $m=2$. Since $d=3$, $m=2$ correspond to real systems, we assume both conditions hold.
- Effective coupling constants:

$$\tilde{V}_a = \tilde{k}_F V_a$$

Epilogue

- RG analysis for QFTs with FS → scaling behaviour of NFL states in a controlled approx.
- m -dim FS with its co-dim extended to a generic value → stable NFL fixed points identified using $\epsilon = d_c - d$ as perturbative parameter.
- SC instability in such systems as a fn of dim & co-dim of FS.
- Key point → k_F enters as a dimensionful parameter unlike in relativistic QFT → modify naive scaling arguments.
- Effective coupling constants → combinations of original coupling constants & k_F .