

Title: The Asymptotic Safety Program: New results and an inconvenient truth

Date: Apr 25, 2014 09:00 AM

URL: <http://pirsa.org/14040108>

Abstract: We briefly review the various components and their conceptual status of the full Asymptotic Safety Program which aims at finding a nonperturbative infinite-cutoff limit of a regularized functional integral for a quantum field theory of gravity. It is explained why in the continuum formulation based on the Effective Average Action the key requirement of background independence unavoidably results in a "bi-metric" framework, and recent results on truncated RG flows of bi-metric actions are presented. They suggest that the next generation of truncations that must be explored should be of bi-metric type. As an application, a method of characterizing and counting physical states is shown to arise.



- 1. The Asymptotic Safety program**
- 2. Can Asymptotic Safety coexist with Background Independence?**
 - (a) Background Independence, Split Symmetry, bi-metric character of the Effective Average Action
 - (b) New results from a bi-metric truncation

- 1. The Asymptotic Safety program**
- 2. Can Asymptotic Safety coexist with Background Independence?**
 - (a) Background Independence, Split Symmetry, bi-metric character of the Effective Average Action
 - (b) New results from a bi-metric truncation

The fundamental problem:

Give a meaning to ("define", "renormalize",
"take the continuum limit of", "...") a functional
integral over all metrics on a space time $\mathcal{M}$:

$$\int \mathcal{D}\hat{g}_{\mu\nu} e^{-S[\hat{g}_{\mu\nu}]}$$

S : diff $(\mathcal{M})$ -invariant
bare action,

e.g. S_{EH} + counter terms

$$\mathcal{D}\hat{g}_{\mu\nu} \equiv \prod_{x \in \mathcal{M}} \prod_{\mu, \nu} dg_{\mu\nu}(x)$$

↑ requires regularisation (UV cutoff)

The fundamental problem:

Give a meaning to ("define", "renormalize",
"take the continuum limit of", "...") a functional
integral over all metrics on a space time $\mathcal{M}$:

$$\int \mathcal{D}\hat{g}_{\mu\nu} e^{-S[\hat{g}_{\mu\nu}]}$$

S : diff($\mathcal{M}$)-invariant
bare action,

e.g. S_{EH} + counter terms

$$\mathcal{D}\hat{g}_{\mu\nu} \equiv \prod_{x \in \mathcal{M}} \prod_{\mu, \nu} dg_{\mu\nu}(x)$$

↑ requires regularisation (UV cutoff)

The Asymptotic Safety program

- Pick (topological) manifold $\mathcal{M}$
- Pick a set of "carrier fields" Φ :
 $g_{\mu\nu}, \Theta_\mu^a, (e_\mu^a, \omega_\mu^{ab}), \dots$
- Pick symmetry/gauge group G :
 $\text{Diff}(\mathcal{M}), \text{Diff}(\mathcal{M}) \ltimes O(4)_{\text{loc}},$
foliation preserving diffeos, ...

$\Rightarrow$ total set of fields $\Psi \equiv (\Phi, \bar{\Phi}, \text{FP ghosts}),$
 $\bar{\Phi}$ = backgrd. field for Φ

- Pick space of fields:

$$\mathcal{F} = \left\{ \text{fields } \Psi \text{ on } \mathcal{M} \mid (\dots) \right\}$$

boundary conditions, degree of regularity, etc.

(N.B.: b.c. $\longleftrightarrow$ states !)

The Asymptotic Safety program

- Pick (topological) manifold $\mathcal{M}$
 - Pick a set of "carrier fields" Φ :
 $g_{\mu\nu}, \Theta_\mu^a, (e_\mu^a, \omega_\mu^{ab}), \dots$
 - Pick symmetry/gauge group G :
 $\text{Diff}(\mathcal{M}), \text{Diff}(\mathcal{M}) \ltimes O(4)_{\text{loc}},$
 foliation preserving diffeos, ...
- $\Rightarrow$ total set of fields $\Psi \equiv (\Phi, \bar{\Phi}, \text{FP ghosts}),$
 $\bar{\Phi}$ = backgrd. field for Φ

- Pick space of fields:

$$\mathcal{F} = \left\{ \text{fields } \Psi \text{ on } \mathcal{M} \mid (\dots) \right\}$$

boundary conditions, degree of regularity, etc.

(N.B.: b.c. $\longleftrightarrow$ states !)

The Asymptotic Safety program

- Pick (topological) manifold $\mathcal{M}$
 - Pick a set of "carrier fields" Φ :
 $g_{\mu\nu}, \Theta_\mu^a, (e_\mu^a, \omega_\mu^{ab}), \dots$
 - Pick symmetry/gauge group G :
 $\text{Diff}(\mathcal{M}), \text{Diff}(\mathcal{M}) \otimes O(4)_{\text{loc}},$
 foliation preserving diffeos, ...
- $\Rightarrow$ total set of fields $\Psi \equiv (\Phi, \bar{\Phi}, \text{FP ghosts}),$
 $\bar{\Phi}$ = backgrd. field for Φ

- Pick space of fields:

$$\mathcal{F} = \left\{ \text{fields } \Psi \text{ on } \mathcal{M} \mid \text{(\dots)} \right\}$$

boundary conditions, degree of regularity, etc.

(N.B.: b.c. $\longleftrightarrow$ states !)

- Pick space of action functionals $\equiv$ theory space :

$$\mathcal{T} = \left\{ A: \mathcal{F} \rightarrow \mathcal{R}, \psi \mapsto A[\psi] \right\}$$

$$\left\{ \delta_G^B A = 0, (\dots) \right\}$$

degree of (non-) locality, regularity, etc.

- Pick coarse graining (averaging) scheme on $\mathcal{T}$:
vector field β

$$A \xrightarrow[\text{RG step}]{\text{infinite.}} A + \underbrace{\beta(A)}_{\in T_A \mathcal{T}}$$

$\Rightarrow$ RG flow $(\mathcal{T}, \beta)$

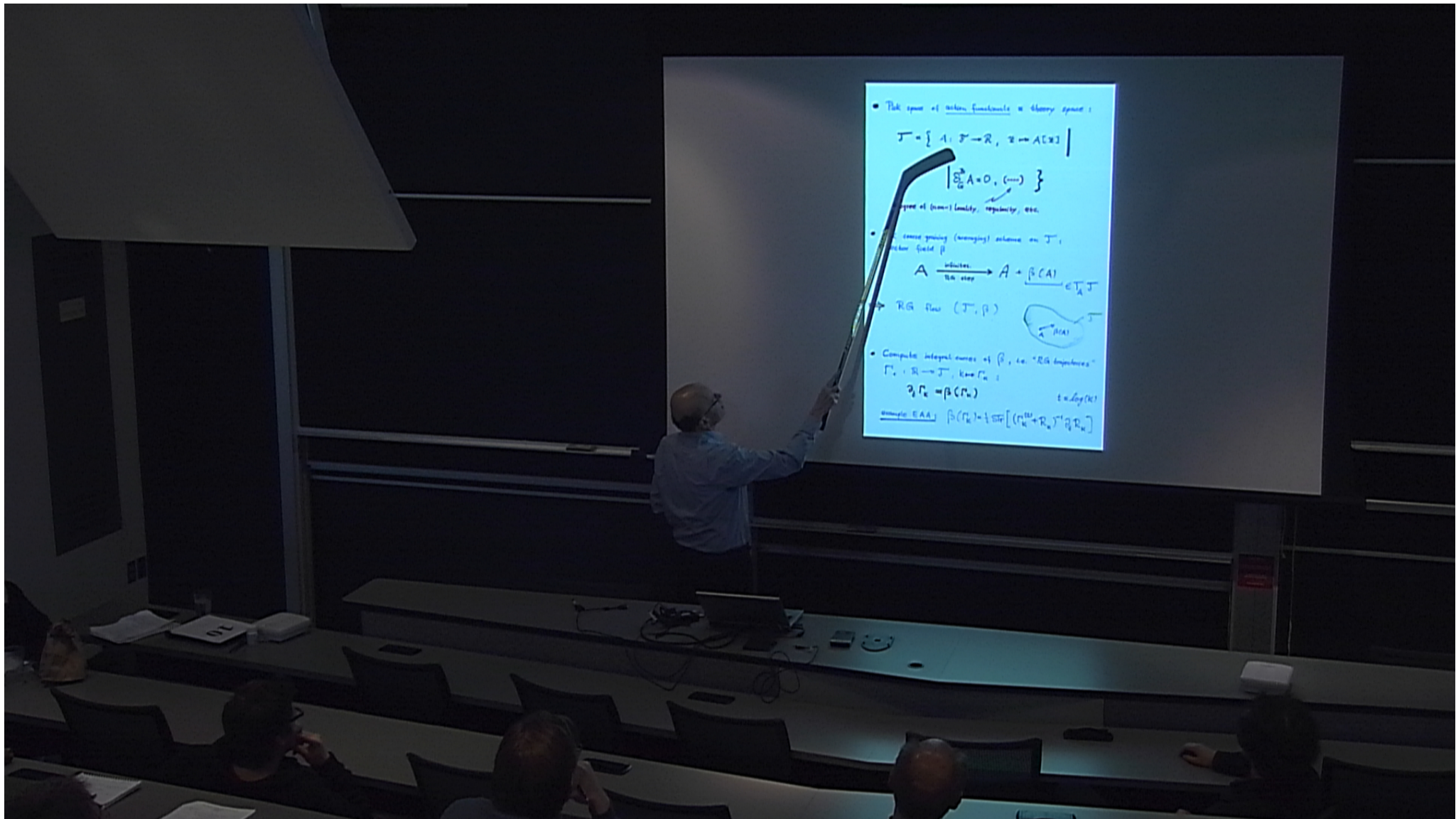


- Compute integral curves of β , i.e. "RG trajectories"

$$\Gamma_\bullet: \mathbb{R} \rightarrow \mathcal{T}, k \mapsto \Gamma_k:$$

$$\partial_t \Gamma_k = \beta(\Gamma_k) \quad t \equiv \log(k)$$

example EAA: $\beta(\Gamma_k) = \frac{1}{2} \text{Str}[(\Gamma_k^{(2)} + R_k)^{-1} \partial_t R_k]$



- Search for fixed points : $\beta(A_*) = 0$
- Determine linearised flow near A_* :
 eigenvectors $\hat{=}$ scaling fields
 eigenvalues $\hat{=}$ critical exponents
- Determine UV critical hypersurface $\mathcal{L}_{UV}$
- Try to find complete RG trajectories,
 having well-behaved limits $k \rightarrow \infty$ and $k \rightarrow 0$.

Example: employ Asymptotic Safety construction w.r.t. A_*

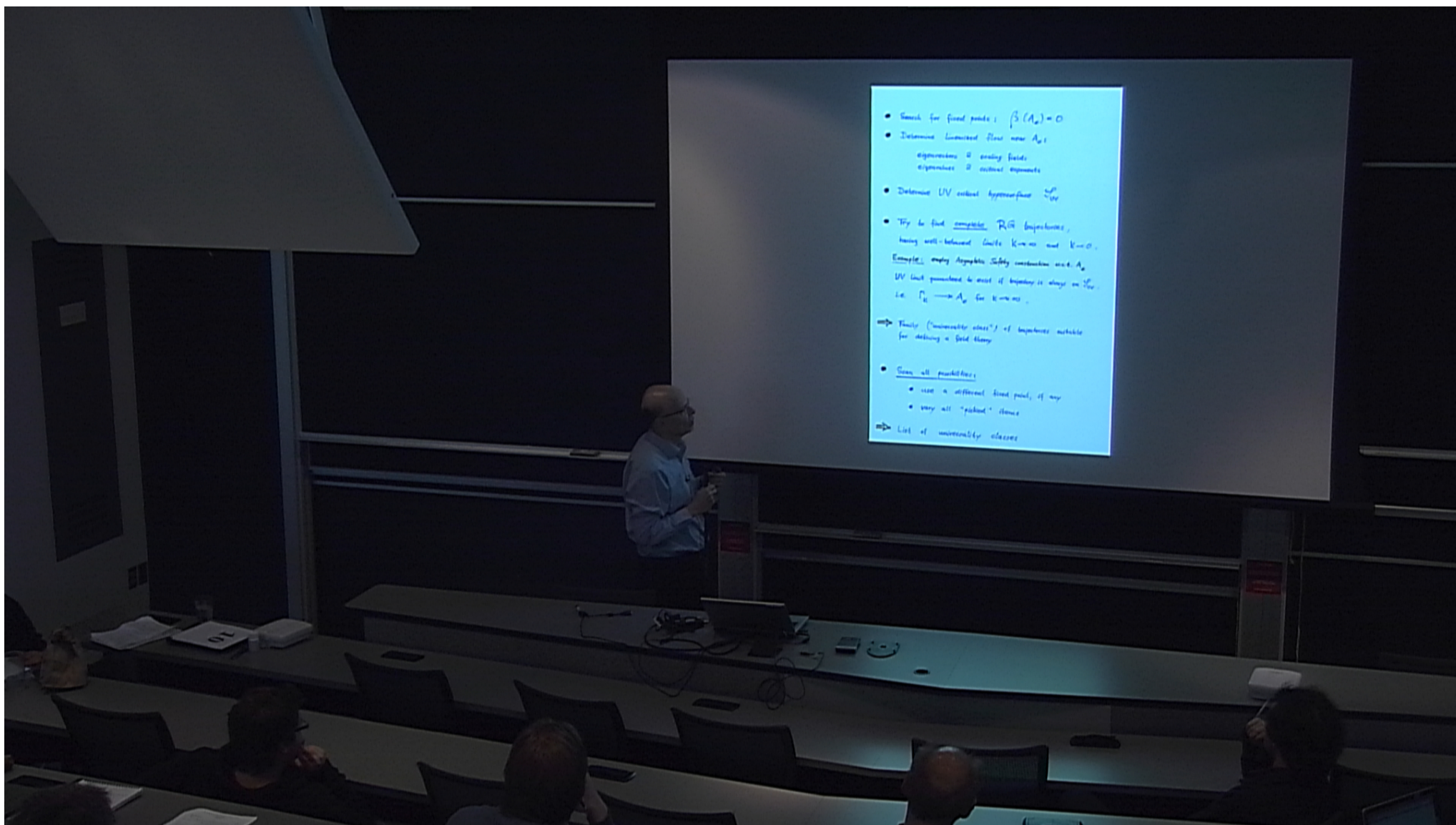
UV limit guaranteed to exist if trajectory is always on $\mathcal{L}_{UV}$,

i.e. $\Gamma_k \rightarrow A_*$ for $k \rightarrow \infty$.

$\Rightarrow$ Family ("universality class") of trajectories suitable
 for defining a field theory

- Scan all possibilities:
 - use a different fixed point, if any
 - vary all "picked" items

$\Rightarrow$ List of universality classes



- Search for fixed points: $\beta(A_\mu) = 0$
- Determine Low-energy flow near A_μ :
 - asymptotically $\Rightarrow$ scaling fields
 - asymptotically $\Rightarrow$ critical exponents
- Determine UV critical hypersurface $\mathcal{H}_{UV}^c$
- Try to find complete RG trajectories,
 - having well-behaved limits $K \rightarrow \infty$ and $K \rightarrow 0$.

Example: empty AdS/CFT. Safety conclusion: A_μ
 UV limit guaranteed to exist if trajectory is always on $\mathcal{H}_{UV}^c$,
 i.e. $\Gamma_K \rightarrow A_\mu$ for $K \rightarrow \infty$.

$\Rightarrow$ Family ("universality class") of hypersurfaces suitable
 for defining a field theory.
- Scan all possibilities:
 - use a different fixed point, if any
 - vary all "padding" flows

$\Rightarrow$ List of universality classes

- Search for fixed points : $\beta(A_*) = 0$
- Determine linearised flow near A_* :
 - eigenvectors $\hat{=}$ scaling fields
 - eigenvalues $\hat{=}$ critical exponents
- Determine UV critical hypersurface $\mathcal{J}_{UV}$
- Try to find complete RG trajectories, having well-behaved limits $k \rightarrow \infty$ and $k \rightarrow 0$.

Example: employ Asymptotic Safety construction w.r.t. A_*

UV limit guaranteed to exist if trajectory is always on $\mathcal{J}_{UV}$,

i.e. $\Gamma_k \rightarrow A_*$ for $k \rightarrow \infty$.

$\Rightarrow$ Family ("universality class") of trajectories suitable for defining a field theory

- Scan all possibilities:
 - use a different fixed point, if any
 - vary all "picked" items

$\Rightarrow$ List of universality classes

- Search for fixed points : $\beta(A_*) = 0$
 - Determine linearised flow near A_* :
 - eigenvectors $\hat{=}$ scaling fields
 - eigenvalues $\hat{=}$ critical exponents
 - Determine UV critical hypersurface $\mathcal{J}_{UV}$
 - Try to find complete RG trajectories, having well-behaved limits $k \rightarrow \infty$ and $k \rightarrow 0$.
- Example: employ Asymptotic Safety construction w.r.t. A_*
- UV limit guaranteed to exist if trajectory is always on $\mathcal{J}_{UV}$,
i.e. $\Gamma_k \rightarrow A_*$ for $k \rightarrow \infty$.

$\Rightarrow$ Family ("universality class") of trajectories suitable for defining a field theory

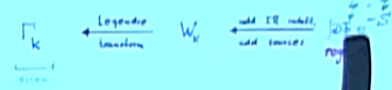
- Scan all possibilities:
 - use a different fixed point, if any
 - vary all "picked" items

$\Rightarrow$ List of universality classes

• Solve the "Reconstruction Problem"

(give a precise definition to the functional integral)

- Pick UV (1) regularization scheme for $\int \mathcal{D}\phi e^{-S}$ $\text{reg } \Lambda$
- Fix a "theory" a complete RG trajectory $\Gamma_K, K \in [0, \infty)$ and construct a functional integral representation of it, for $K < \Lambda$:



(trajectory, UV reg. scheme) with proper choice for taking the "continuum limit", i.e. sending the two parameters in the Λ, reg to ∞ .

$$S_\Lambda = \Gamma_{K=\Lambda} + \text{counterterm}(\Lambda, \text{reg})$$

E. Mouhssine, MR, 2008

- Find the "minimal" interpretation of S_Λ , read off phase-space description of the "elemental dof's"

- Solve the "Reconstruction Problem"

(give a precise definition to the functional integral)

- Pick UV (1) regularization scheme for $\int \mathcal{D}\Phi e^{-S}$
reg, Λ
- Fix a "theory" a complete RG trajectory $\Gamma_K \in \mathcal{K}(\mathcal{D}, \infty)$
and construct a functional integral representation
of it, for $K < \Lambda$:

$$\Gamma_K \xleftarrow[\text{trajectory}]{\text{Legendre}} W_K \xleftarrow[\text{add sources}]{\text{add IR infl.}} \int \mathcal{D}\Phi e^{-S} \xrightarrow{\text{reg, } \Lambda} \int \mathcal{D}\Phi e^{-S}$$

(trajectory, UV reg. scheme) with prescription for
taking the "continuum limit", i.e. for fixing the
two parameters in the limit $\Lambda \rightarrow \infty$:

$$S_\Lambda = \Gamma_{K, \Lambda} + \text{correction term}(\Lambda, \text{reg})$$

E. Mouhssine, MR, 2008

- Find Hamiltonian interpretation of S_Λ ,
read off phase-space description of the
"fundamental dof's"

- Solve the "Reconstruction Problem"

(give a precise definition to the functional integral)

- Pick UV (1) regularization scheme for $\int \mathcal{D}\Phi e^{-S}$ reg, Λ
- Fix a "theory" a complete RG trajectory $\Gamma_K, K_0(\infty)$ and construct a functional integral representation of it, for $K < \Lambda$:

$$\Gamma_K \xleftarrow[\text{bosons}]{\text{fermions}} W_K \xleftarrow[\text{add sources}]{\text{add IR infl.}} \int \mathcal{D}\Phi e^{-S} \xrightarrow[\text{reg, } \Lambda]{\substack{\text{add } \frac{1}{2} \log \Lambda \\ \text{add } \frac{1}{2} \log \Lambda}}$$

(trajectory, UV reg. scheme) with prescription for taking the "continuum limit", i.e. for fixing the bare parameters in the limit $\Lambda \rightarrow \infty$:

$$S_\Lambda = \Gamma_{K=\Lambda} + \text{correction term}(\Lambda, \text{reg})$$

E. Mouhssine, MR, 2008

- Find Hamiltonian interpretation of S_Λ , read off phase-space description of the "fundamental dof's"

Changing the input data (examples)

* Change topology

- S^4 vs. $\mathbb{R}^4$ in conformally reduced gravity,
 $SO(4)$ -dimensional bransons MR, H. Langer, 2009
- DLL at $D=4$: running boundary terms
 J. Becker, MR, 2010

* Change definition of field space $\tilde{\mathcal{F}}$

different h.c. for conf. factors: MR, H. Langer, 2009
 $\dim \mathcal{U}_{UV} = \text{finite} \longleftrightarrow \text{infinite}$

* Change set of carrier fields

$\mathcal{C}_n^A, (\mathcal{C}_n^B, \omega_n^{A+B}), (g_{\mu\nu}, T_{\mu\nu})$ J. Harst, MR, 2010
 MR, G. Schottmeyer

* Change gauge group

$\rightarrow$, foliation preserving diffeomorphism, T. Sauerbrunn et al.
 volume preserving diffeomorphism A. Kalloni

* Change "degree of integrability"

$SO(4)$ -dim. non-local $(\mathcal{C}_n^B)$ -bransons MR, T. Sauerbrunn,
 2002
 $\rightarrow$ infrared fixed point fields $\mathcal{F}_0(V)$

Background Independence

The hallmark of classical GR
and Quantum Gravity !

- Treat all possible spacetimes on an equal footing
- no "vacuum" singled out a priori
- derive rather than put in "by hand"
the arena of all non-gravitational physics
(Minkowski space, ...)

Even more fundamental and momentous
than the renormalizability issue !

Truncating theory space (metric gravity)

$$\overline{\mathcal{T}}_{\text{QEG}} : \quad \Gamma_k[h, c, \bar{c}; \bar{g}] \equiv \Gamma_k[g, \bar{g}, c, \bar{c}] \quad \Big|_{g_{\mu\nu} = \bar{g}_{\mu\nu} + h_{\mu\nu}}$$

"extra $\bar{g}$ -dependence" 

"single metric truncations":

MR, 1996

$$\Gamma_k = \underbrace{\Gamma_k[g]}_{1996 \dots 2009} + \int (\mathcal{F}(\bar{g})h)^2 + \int \bar{c} M(\bar{g}) c$$

Manrique, MR, 2009

Manrique, MR, Smorczys, 2010

"bi-metric truncations":

D. Becker, MR, 2014

$$\Gamma_k = \Gamma_k[g, \bar{g}] + \int (\mathcal{F}(\bar{g})h)^2 + \int \bar{c} M(\bar{g}) c$$

Taylor expansion:

$$\Gamma_k[g, \bar{g}] \equiv \Gamma_k[h; \bar{g}] = \sum_{p=0}^{\infty} \frac{1}{p!} \int \chi_k^{(p)}[\bar{g}] \underbrace{h \dots h}_{p \text{ factors}}$$

$p = \text{"level"}$

Single metric approx.: retain only level $p=0$,
 $g_{\mu\nu}$ and $\bar{g}_{\mu\nu}$ not disentangled

- 6 fixed points on the total theory space:

$$G^{\mathcal{B}} \oplus G^{\text{Dyn}}\text{-FP}, \quad G^{\mathcal{B}} \oplus NG_{-}^{\text{Dyn}}\text{-FP}, \quad G^{\mathcal{B}} \oplus NG_{+}^{\text{Dyn}}\text{-FP}$$

$$NG^{\mathcal{B}} \oplus G^{\text{Dyn}}\text{-FP}, \quad NG^{\mathcal{B}} \oplus NG_{-}^{\text{Dyn}}\text{-FP}, \quad \underline{NG^{\mathcal{B}} \oplus NG_{+}^{\text{Dyn}}\text{-FP}}$$

candidate for the
A.S. construction:

$$(g_{*}^{\text{Dyn}} = 0.703, \lambda_{*}^{\text{Dyn}} = 0.207; \quad g_{*}^{\mathcal{B}} = 8.2, \lambda_{*}^{\mathcal{B}} = -0.01)$$

$$\Theta = (3.6 \pm 3.0 i, \quad 2, \quad 4) \rightsquigarrow \dim \mathcal{F}_{UV} = 4$$

- Background Independence and Asymptotic Safety
can be achieved simultaneously:

There exists a 2-parameter subset of
RG trajectories on the 4-dimensional theory
space which are asymptotically safe w.r.t. the
 $NG^{\mathcal{B}} \oplus NG_{+}^{\text{Dyn}}\text{-FP}$, and which restore
split-symmetry in the physical limit ($K \rightarrow 0$).

Predictivity higher than expected:

$$\dim \mathcal{F}_{UV} - 2 = 4 - 2 = \underline{2 \text{ free param.s}}$$

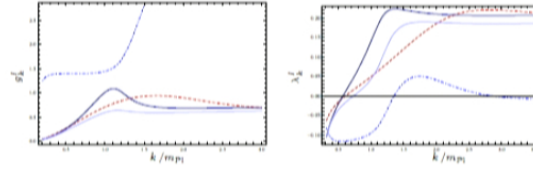


Figure 6. Type (Ia)^{Dyn}-(Attr)^Btrajectory: dimensionless couplings.

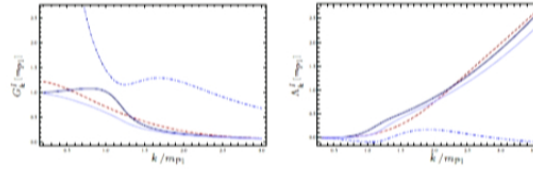


Figure 7. Type (Ia)^{Dyn}-(Attr)^Btrajectory: dimensionful couplings.

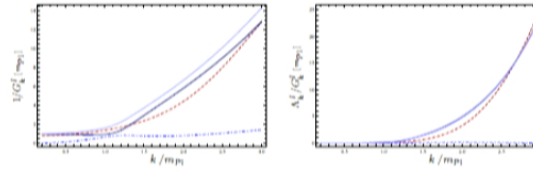


Figure 8. Type (Ia)^{Dyn}-(Attr)^Btrajectory: the coefficients as they appear in the EAA. Note the perfect split-symmetry restoration in the IR: $1/G_k^{\text{B}}$ and $\lambda_k^{\text{B}}/G_k^{\text{B}}$ vanish for $k \rightarrow 0$, implying that $\Gamma_k^{\text{B}aw}$ loses its extra $\hat{g}_{aw}$ dependence.

dashed (red):	---	$I = \text{sm}$ (single-metric)
solid (dark-blue):	—	$I = \text{Dyn} \equiv (p)$ for $p \geq 1$
solid (light-blue):	—	$I = (0)$
dot-dashed (blue):	-.-	$I = \text{B}$

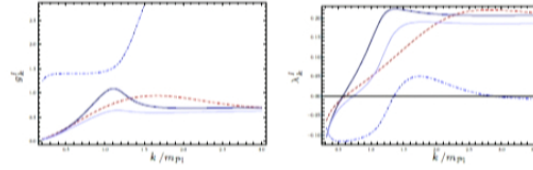


Figure 6. Type (Ia)^{Dyn}-(Attr)^Btrajectory: dimensionless couplings.

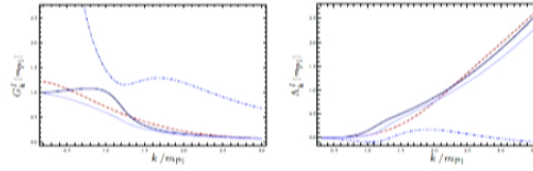


Figure 7. Type (Ia)^{Dyn}-(Attr)^Btrajectory: dimensionful couplings.

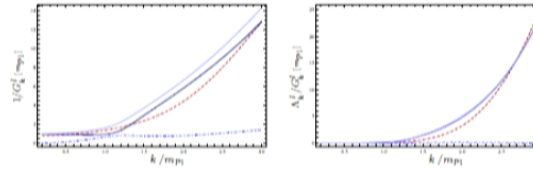


Figure 8. Type (Ia)^{Dyn}-(Attr)^Btrajectory: the coefficients as they appear in the EAA. Note the perfect split-symmetry restoration in the IR: $1/G_k^{\text{B}}$ and $\lambda_k^{\text{B}}/G_k^{\text{B}}$ vanish for $k \rightarrow 0$, implying that $\Gamma_k^{\text{B}av}$ loses its extra $\hat{g}_{\mu\nu}$ dependence.

dashed (red):	---	$I = \text{sm}$ (single-metric)
solid (dark-blue):	—	$I = \text{Dyn} \equiv (p)$ for $p \geq 1$
solid (light-blue):	—	$I = (0)$
dot-dashed (blue):	-.-	$I = \text{B}$

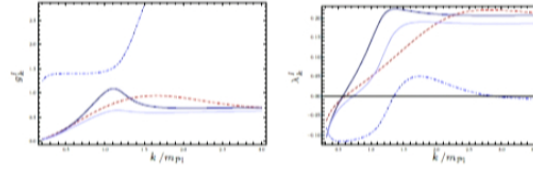


Figure 6. Type (Ia)^{Dyn}-(Attr)^Btrajectory: dimensionless couplings.

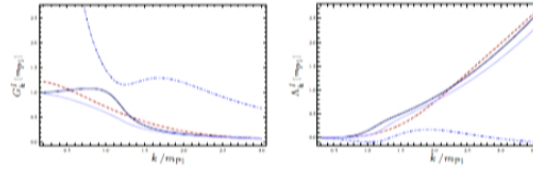


Figure 7. Type (Ia)^{Dyn}-(Attr)^Btrajectory: dimensionful couplings.

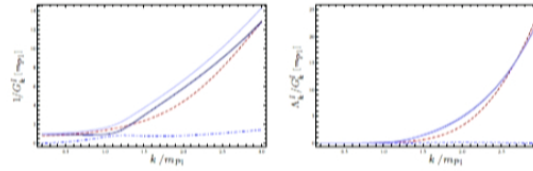


Figure 8. Type (Ia)^{Dyn}-(Attr)^Btrajectory: the coefficients as they appear in the EAA. Note the perfect split-symmetry restoration in the IR: $1/G_k^{\text{B}}$ and $\lambda_k^{\text{B}}/G_k^{\text{B}}$ vanish for $k \rightarrow 0$, implying that $\Gamma_k^{\text{B}aw}$ loses its extra $\hat{g}_{aw}$ dependence.

dashed (red):	---	$I = \text{sm}$ (single-metric)
solid (dark-blue):	—	$I = \text{Dyn} \equiv (p)$ for $p \geq 1$
solid (light-blue):	—	$I = (0)$
dot-dashed (blue):	-.-	$I = \text{B}$

The Main Result:

The 2-parameter family of RG trajectories restoring split symmetry in the IR consists of precisely those which approach $\text{Attr}^{\mathcal{B}}(K \rightarrow 0)$ in the IR.

Summary:

(Potentially) acceptable QFTs can be based on RG trajectories with

initial point $(K \rightarrow \infty)$: $NG^{\mathcal{B}} \oplus NG_+^{\text{Dyn}}$ -FP

final point $(K \rightarrow 0)$: $\text{Attr}^{\mathcal{B}}(K \rightarrow 0)$

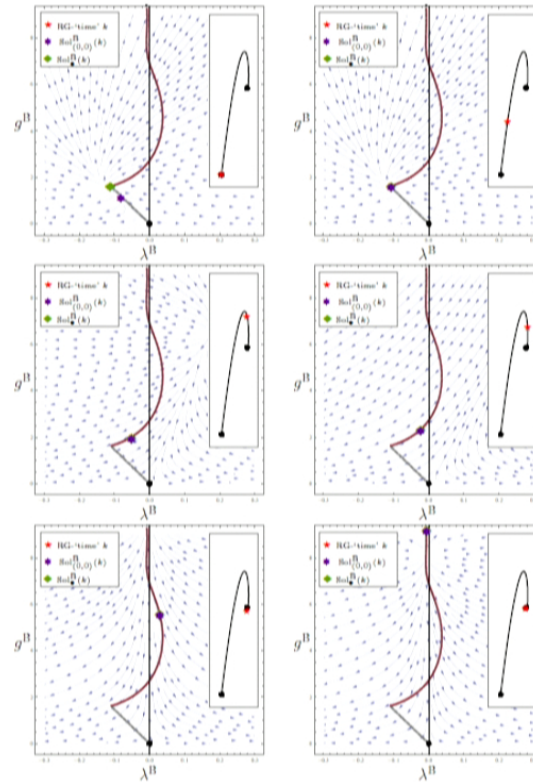


Figure 3. The B-phase portraits at increasing scales k . The underlying separatrix in the Dyn-sector is shown in the inset on the right, and the current RG time is marked with a star therein. The arrows point towards the IR and picture the instantaneous vector field in the B-sector. The (red) solid and the (gray) dashed curve highlight two important solutions in the B-sector, namely $\text{Sol}_+^B(k)$ and $\text{Sol}_{0,0}^B(k)$, respectively. Their current position is indicated by the (green) diamond and the (violet) six-pointed star, respectively.

Towards an EAA-based "C-Theorem"

D. Becker, MR, 2012 ; 2014

$$e^{-\Gamma_k[\langle\phi\rangle_k]} = \int \mathcal{D}\phi e^{-S[\phi]} e^{-\Delta S_k[\phi]} \sim \int_{p>k} \mathcal{D}\phi e^{-S}$$

Γ_k evaluated "on-shell" counts field modes integrated out above the scale k .

gravity: $\Gamma_k[h=0; \bar{g}_k^{\text{self consistent}}] = C_k$

$$\left(\frac{\delta}{\delta h_{\mu\nu}} \Gamma_k \right) \Big|_{h=0, \bar{g} = \bar{g}_k^{\text{self cons.}}} = 0 \quad (\text{tadpole equation})$$

bi-metric EH truncation:

$$\Gamma_k[h=0; \bar{g}] = -\frac{1}{16\pi G_k^{(0)}} \int dx \sqrt{\bar{g}} \{ R(\bar{g}) - 2\Lambda_k^{(0)} \}$$

$$G_{\mu\nu}(\bar{g}_k^{\text{self con.}}) = -\Lambda_k^{(1)} [\bar{g}_k^{\text{self con.}}]_{\mu\nu}$$

$$\leadsto S^4 \text{ with radius } \sim [\Lambda_k^{(1)}]^{-1/2}$$

Conclusion

Split-symmetry $\equiv$ Background Independence

- is crucial to select the correct universality class
- is restrictive, but not too restrictive
- can indeed coexist with Asymptotic Safety, at least in the truncation considered

Truncating theory space (metric gravity)

$$\overline{\mathcal{T}}_{\text{QEG}} : \quad \Gamma_k[h, C, \bar{C}; \bar{g}] \equiv \Gamma_k[g, \bar{g}, C, \bar{C}] \quad \Big|_{g_{\mu\nu} = \bar{g}_{\mu\nu} + h_{\mu\nu}}$$

"extra $\bar{g}$ -dependence" 

"single metric truncations":

MR, 1996

$$\Gamma_k = \underbrace{\Gamma_k[g]}_{1996 \dots 2009} + \int (\mathcal{F}(\bar{g})h)^2 + \int \bar{C} M(\bar{g}) C$$

Manrique, MR, 2009

Manrique, MR, Smorczys, 2010

"bi-metric truncations":

D. Becker, MR, 2014

$$\Gamma_k = \Gamma_k[g, \bar{g}] + \int (\mathcal{F}(\bar{g})h)^2 + \int \bar{C} M(\bar{g}) C$$

Taylor expansion:

$$\Gamma_k[g, \bar{g}] \equiv \Gamma_k[h; \bar{g}] = \sum_{p=0}^{\infty} \frac{1}{p!} \int \chi_k^{(p)}[\bar{g}] \underbrace{h \dots h}_{p \text{ factors}}$$

$p = \text{"level"}$

Single metric approx.: retain only level $p=0$,
 $g_{\mu\nu}$ and $\bar{g}_{\mu\nu}$ not disentangled

Truncating theory space (metric gravity)

$$\overline{\Gamma}_{\text{QEG}} : \quad \Gamma_k[h, c, \bar{c}; \bar{g}] \equiv \Gamma_k[g, \bar{g}, c, \bar{c}] \quad \Big|_{g_{\mu\nu} = \bar{g}_{\mu\nu} + h_{\mu\nu}}$$

"extra $\bar{g}$ -dependence" 

"single metric truncations":

MR, 1996

$$\Gamma_k = \underbrace{\Gamma_k[g]}_{1996 \dots 2009} + \int (\mathcal{F}(\bar{g})h)^2 + \int \bar{c} M(\bar{g}) c$$

Manrique, MR, 2009

Manrique, MR, Smorczys, 2010

"bi-metric truncations":

D. Becker, MR, 2014

$$\Gamma_k = \Gamma_k[g, \bar{g}] + \int (\mathcal{F}(\bar{g})h)^2 + \int \bar{c} M(\bar{g}) c$$

Taylor expansion:

$$\Gamma_k[g, \bar{g}] \equiv \Gamma_k[h; \bar{g}] = \sum_{p=0}^{\infty} \frac{1}{p!} \int \chi_k^{(p)}[\bar{g}] \underbrace{h \dots h}_{p \text{ factors}}$$

$p = \text{"level"}$

Single metric approx.: retain only level $p=0$,
 $g_{\mu\nu}$ and $\bar{g}_{\mu\nu}$ not disentangled

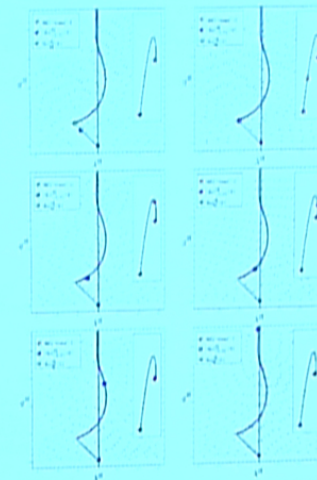


Figure 8. The H-plane profiles at increasing values of k . The underlying equations in the Discretizer is shown in the inset on the right, and the current (BC) time is marked with a star therein. The arrows point towards the BE and picture the instantaneous vector field in the H-vector. The (red), (blue) and the (green) dashed curve highlight two important solutions in the H-vector, namely $\mathbf{H}_0^{\text{red}}(t)$ and $\mathbf{H}_0^{\text{blue}}(t)$, respectively. Their current position is indicated by the (green) diamond and the (red) six-pointed star, respectively.