

Title: On background-independent renormalization in state-sum model

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Abstract: In this talk we discuss some notion of coarse graining in state-sums, most notably a class of spin foam models in their holonomy representations. We discuss the notion of scale in this context, and how diffeomorphism-invariance ties into the existence of a continuum limit. We close with an example and muse about the interplay between diffeomorphism-invariance and non-renormalizability.

On background-independent renormalization in state-sum models

Benjamin Bahr
University of Hamburg

Workshop
„Renormalization Group Approaches to Quantum Gravity“
Perimeter Institute, 24th April 2014



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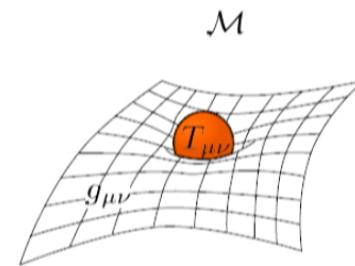


- Introduction
- (Holonomy) Spin Foam Models
- Coarse Graining and continuum limit: Measure theory perspective
- Scale and RG flow in background-independent context
- Easy examples
- Summary

Introduction

Gravity \iff Curvature of space-time metric $g_{\mu\nu}$

$$S_{\text{EH}}[g_{\mu\nu}] = \frac{1}{16\pi G_N} \int d^4x \sqrt{-g} R$$



Quantum gravity: make sense of

$$Z = \int \mathcal{D}g_{\mu\nu} e^{iS_{\text{EH}}[g_{\mu\nu}]}$$

or rather:

$$\langle \mathcal{O} \rangle = \frac{1}{Z} \int \mathcal{D}g_{\mu\nu} \mathcal{O}[g_{\mu\nu}] e^{iS_{\text{EH}}[g_{\mu\nu}]}$$

Introduction

Perturbative ansatz: $g_{\mu\nu} = \eta_{\mu\nu} + \underline{h}_{\mu\nu} + \dots$

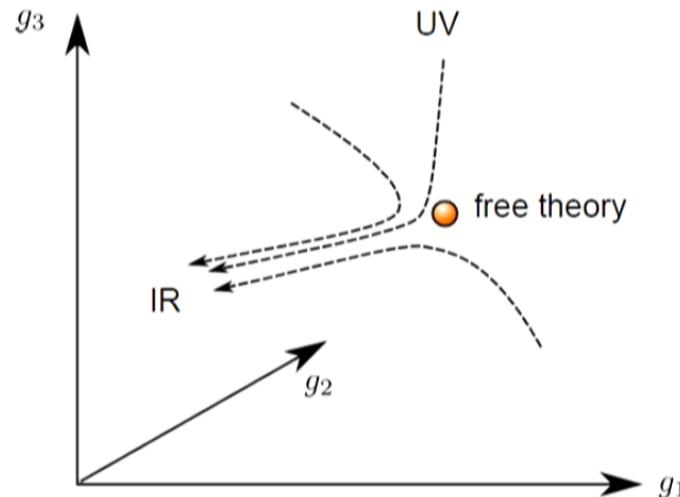
renormalization: scale-dependent couplings $g_i(k)$

scale-dependent effective action

$$S_{\text{EH}} \rightarrow S_{\text{eff}}^{(k)}[h_{\mu\nu}] = S_{\text{EH}}[\eta_{\mu\nu} + h_{\mu\nu}] + \sum_i g_i(k) \mathcal{O}_i[h_{\mu\nu}]$$

non-renormalizable: need infinitely many coupling constants to UV complete theory

[Goroff, Sagnotti '85]



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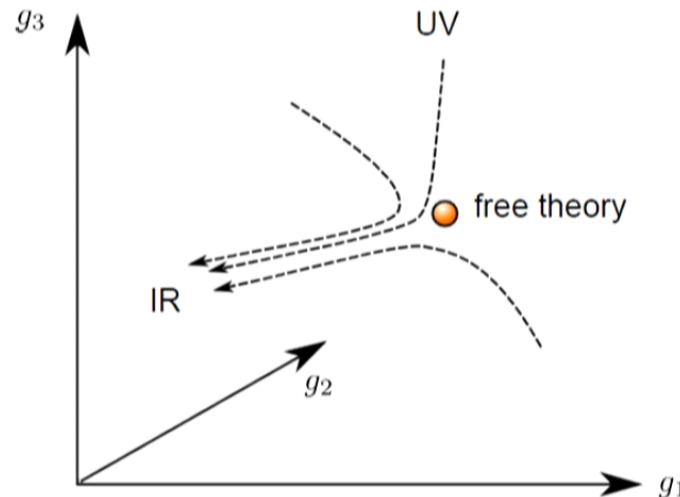
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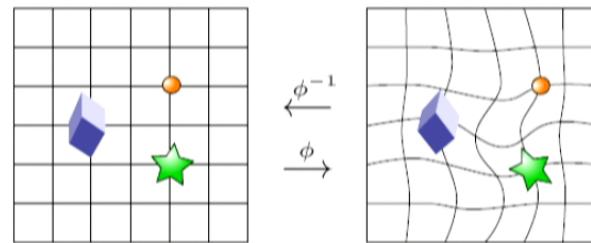
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Diffeomorphism invariance

Diffeomorphism: $\phi: \mathcal{M} \rightarrow \mathcal{M}$

$$g_{\mu\nu}(x) \rightarrow \frac{\partial\phi^{\mu'}}{\partial x^\mu} \frac{\partial\phi^{\nu'}}{\partial x^\nu} g_{\mu'\nu'}(\phi(x))$$



Einstiens „hole argument“:

points in space-time have no meaning, just relations between them

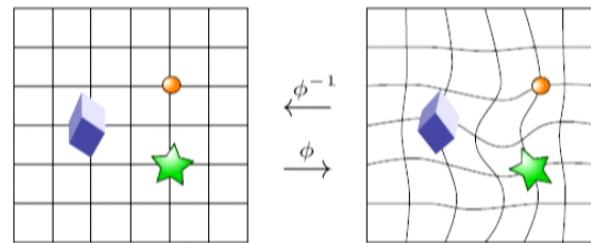
diffeomorphisms = gauge symmetry of GR

\Rightarrow background-independence \Rightarrow no a priori scale available

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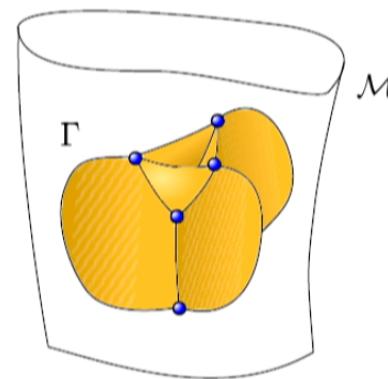
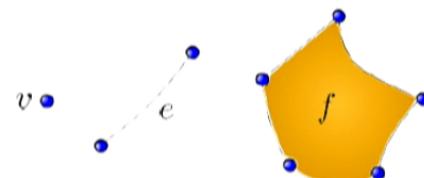
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(Holonomy) spin foam models

[Pfeiffer '01, Pfeiffer, Oeckl '02,
Magliaro, Perini, '10, BB, Dittrich,
Hellmann, Kamiński '12]

Large class of models defined on an oriented, embedded 2-complex $\Gamma \subset \mathcal{M}$



For each Γ a configuration space $\mathcal{A}_\Gamma \simeq G^n$

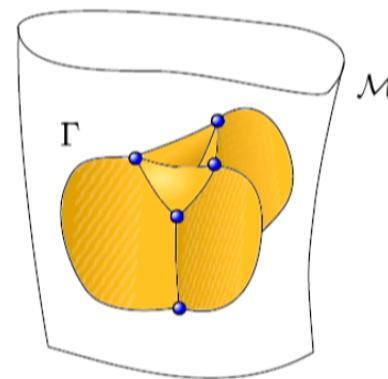
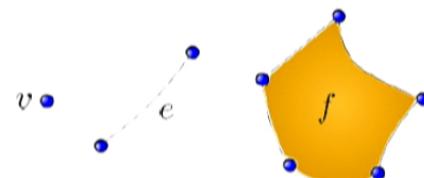
and a measure („amplitude function“) on it $\langle \cdot \rangle_\Gamma : C^0(\mathcal{A}_\Gamma) \rightarrow \mathbb{C}$

$$\langle \mathcal{O} \rangle_\Gamma = \int_{G^n} d\mu_\Gamma(h) \mathcal{O}(h)$$

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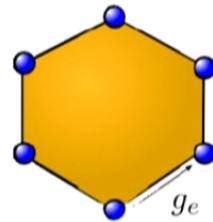
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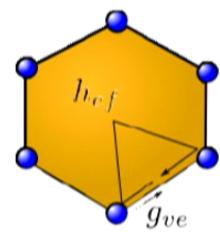
\mathcal{A}_Γ and μ_Γ depend on your model of choice:



lattice gauge theory

[Wilson '74, ...]

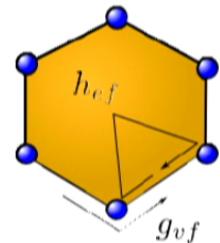
$$d\mu_\Gamma = \prod_f e^{-S[H_f]} dg_e$$



spin foam models

[Reisenberger '94, Barrett, Crane '99, Livine, Speziale '07, Engle, Pereira, Rovelli, Livine '07, Freidel, Krasnov '07, Oriti Baratin '11, ...]

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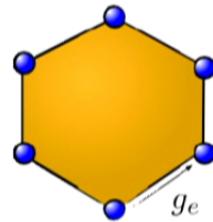
tensor field theories

[DePietri, Freidel, Kransov, Rovelli '00, Oriti '06, Gurau '09, Gurau, Rivasseau '11, ...]

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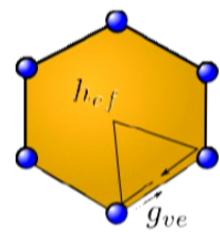
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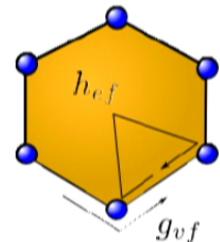
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Conceptual limitations of this formulation:

- a priori no Lorentzian signature: G compact!
- no sum over topologies (\mathcal{M} is fixed)
- embedding excess baggage?

However:

- natural notion of continuum limit
- gives rise to natural notion of scale and RG flow
- allows discussion of diffeo-invariance

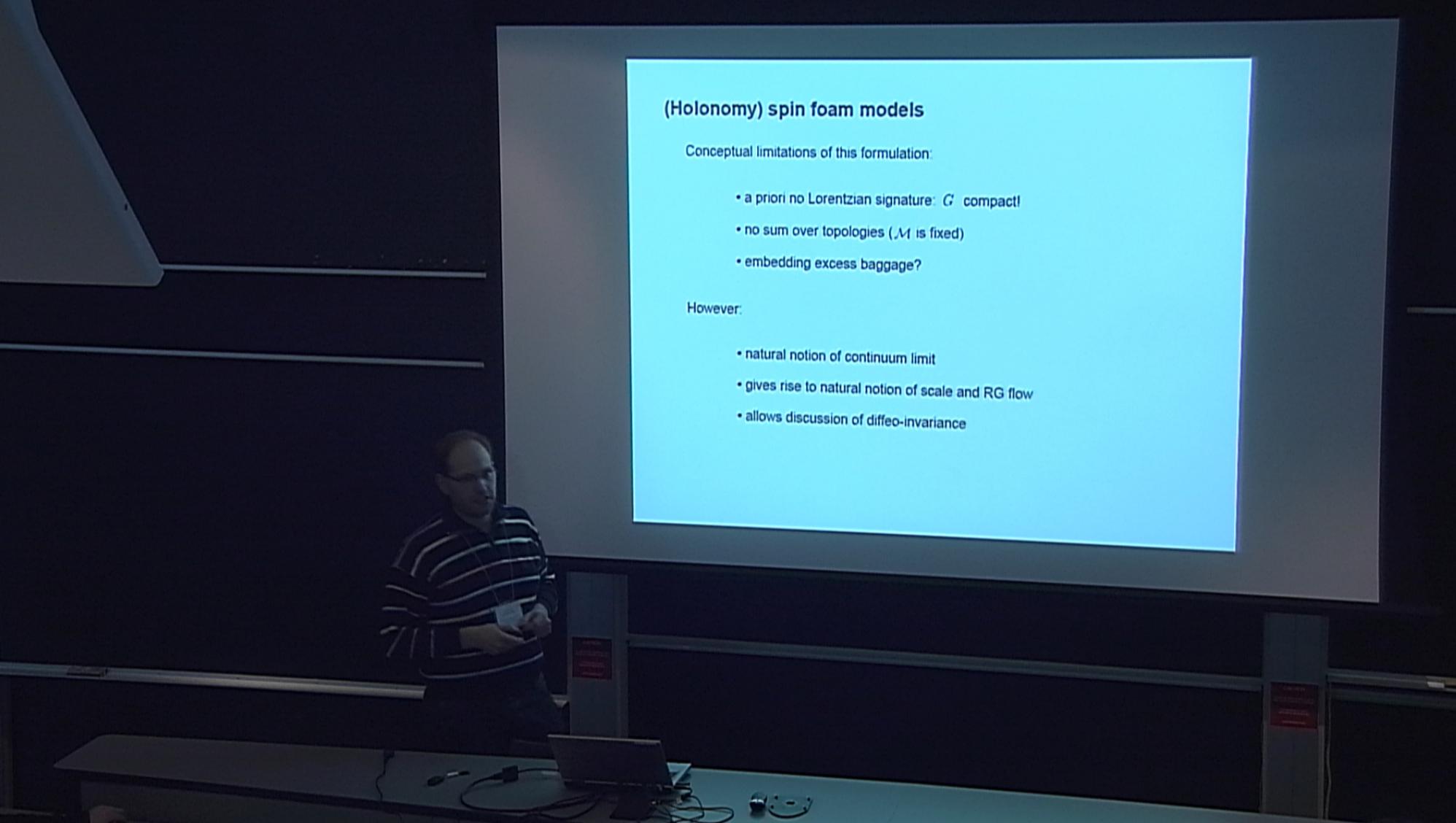
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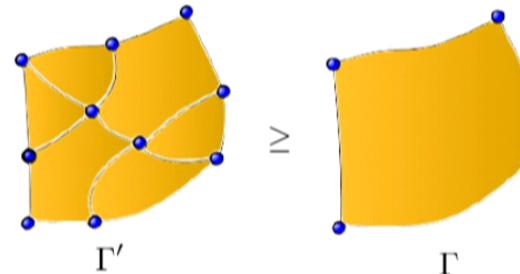


Continuum limit

Natural partial ordering of Γ 's (in semi-analytic category)

[Łojasiewicz, '64, Fleischhack '08]

$$\Gamma \leq \Gamma'$$



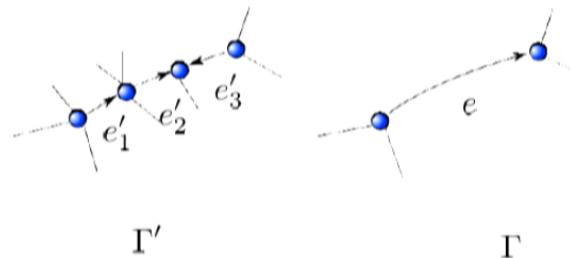
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Not every two Γ 's can be compared, but: for each two Γ , Γ' there is a finer one:

input: projection („coarse graining map“) $\pi_{\Gamma'\Gamma} : \mathcal{A}_{\Gamma'} \rightarrow \mathcal{A}_{\Gamma}$

example from LGT:

$$\pi_{\Gamma'\Gamma}(h_{e'}) = (\cdots, \underbrace{h_{e'_3}^{-1} h_{e'_2} h_{e'_1}}_{=h_e}, \cdots)$$



$$\Gamma'$$

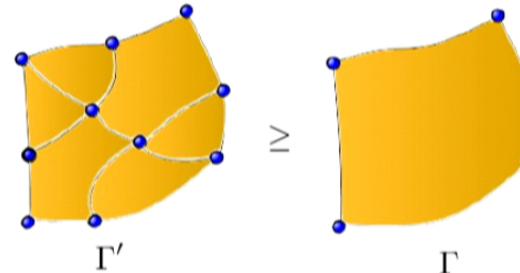
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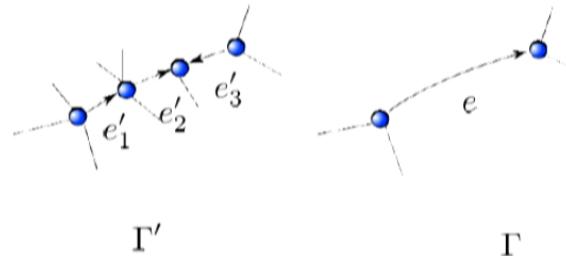


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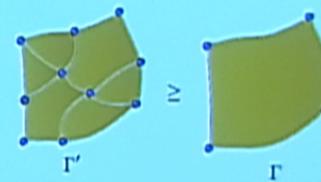


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[Ashtekar, Isham '92,
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Consider all (or sufficiently many) Γ at the same time:

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space of (generalized) continuum connections: compact Hausdorff space

$$\mathcal{A}_{\mathcal{M}} \subset \overline{\mathcal{A}}$$

[cfr. Wiener measure]

condition for continuum measure μ on $\overline{\mathcal{A}}$:

$$\langle \mathcal{O} \rangle_\Gamma = \langle \mathcal{O} \circ \pi_{\Gamma' \Gamma} \rangle_{\Gamma'}$$

„cylindrical consistency“

$$\{\mu_\Gamma\}_\Gamma \rightarrow \mu \quad \text{Radon measure on } \overline{\mathcal{A}}$$

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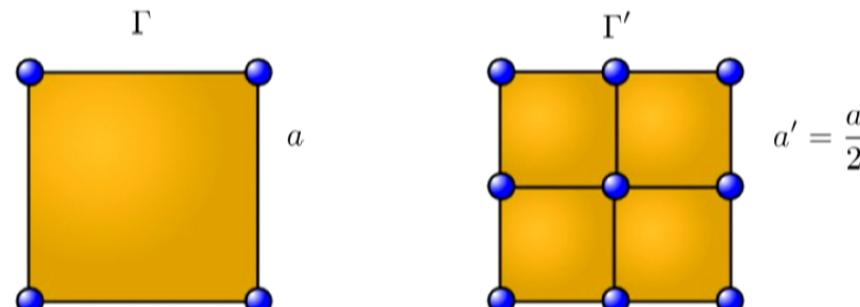
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Where is the physical intuition?

cylindrical consistency: $\langle \mathcal{O} \rangle_{\Gamma} = \langle \mathcal{O} \circ \pi_{\Gamma' \Gamma} \rangle_{\Gamma'}$

is precisely the idea of Wilsonian RG flow!

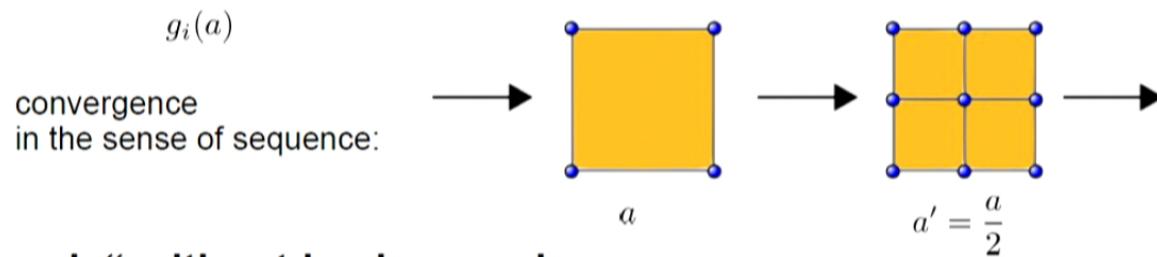


$$d\mu = e^{-S[h_e, g_i(a)]} dh_e$$

measure („action“) parametrised by parameters $g_i(a)$ which depend on a

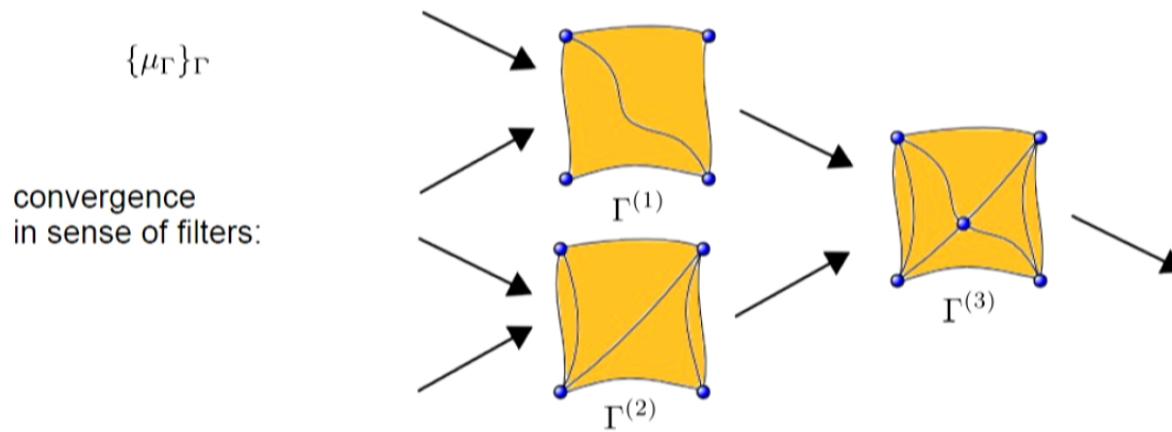
$$e^{-S[h_e, g_i(a)]} = \int dh_{e'} e^{-S[h_{e'}, g_i(a')]} \delta \left(h_e = \overrightarrow{\prod}_{e' \subset e} h_{e'} \right)$$

„scale“ with background:



„scale“ without background:

[Manrique, Oeckl, Weber, Zapata '95
Smerlak, Rovelli '11]



Hamiltonian formulation

Manifold with boundary, e.g.

$$\partial\mathcal{M} = \overline{\Sigma_1} \sqcup \Sigma_2$$

$$\text{boundary graph } \partial\Gamma = \overline{\gamma_1} \sqcup \gamma_2$$

$$\text{boundary holonomies } \mathcal{B}_\gamma = G^{m_\gamma}$$

kinematical boundary Hilbert space:

$$\mathcal{H}_\gamma = L^2(\mathcal{B}_\gamma)$$

boundary observables: $\psi \in C^0(\mathcal{B}_\gamma)$

$$\langle \psi_1 | \psi_2 \rangle_{\Gamma, \text{phys}} := \langle \overline{\psi_1} \otimes \psi_2 \rangle_\Gamma$$

[Ashtekar et al '95,
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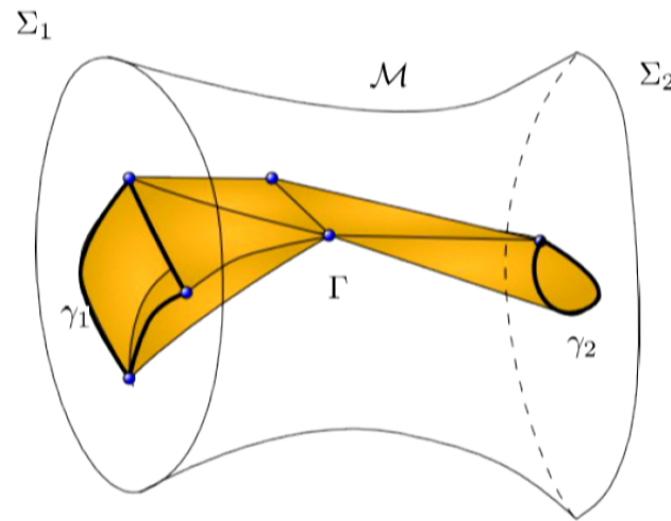
provides physical inner product in sense of rigging map in RAQ

cylindrical consistency of $\langle \cdot \rangle_\Gamma$ guarantees extension to continuum boundary HS

$$\gamma \leq \gamma' \quad \iota_{\gamma\gamma'} : \mathcal{H}_\gamma \hookrightarrow \mathcal{H}_{\gamma'} \quad \langle \psi_1 | \psi_2 \rangle_{\Gamma, \text{phys}} = \langle \iota_{\gamma\gamma'} \psi_1 | \iota_{\gamma\gamma'} \psi_2 \rangle_{\Gamma', \text{phys}}$$

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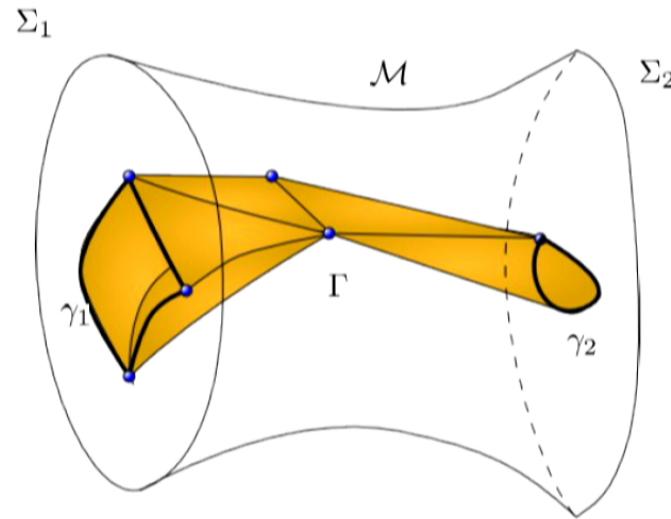
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[Dittrich, Steinhaus '12]



Diffeomorphism group action

Group of (semi-analytic) diffeomorphisms act on classical (continuum) connections $\mathcal{A} \subset \overline{\mathcal{A}}$

$$A \longmapsto (\phi^{-1})^* A$$

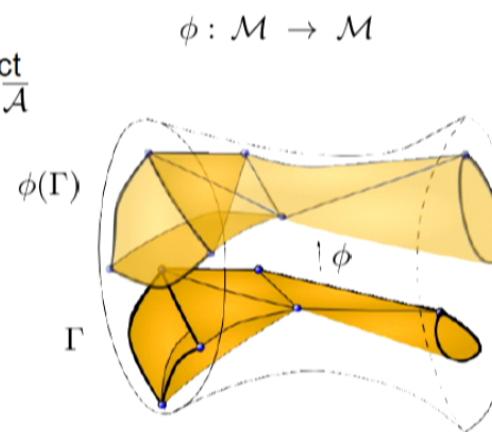
induced action on generalized connections via action on 2-complexes.

$$\phi : \overline{\mathcal{A}} \longrightarrow \overline{\mathcal{A}}$$

invariance of continuum measure under ϕ equivalent to

$$\langle \mathcal{O} \rangle_{\Gamma} = \langle \phi^* \mathcal{O} \rangle_{\phi(\Gamma)} \quad \text{for all } \Gamma \text{ supporting } \mathcal{O}$$

$$\Leftrightarrow \phi_* \mu = \mu$$



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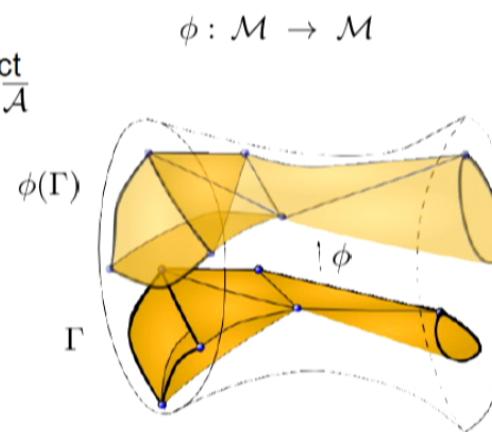
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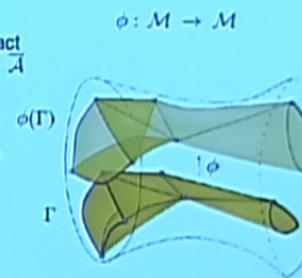
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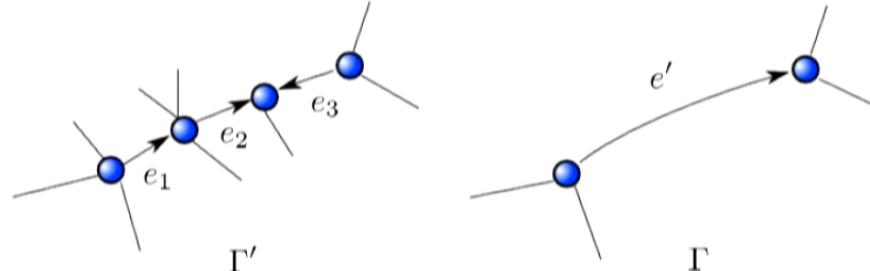
Example:

Near trivial example: $\dim \mathcal{M} = 2$, $G = U(1)$, $\mathcal{A}_\Gamma = U(1)^E$

$$\text{"charge-network functions"} \quad \mathcal{O}(h_e) = \prod_e h_e^{m_e}$$

$$\begin{aligned} \langle \mathcal{O} \rangle_\Gamma &:= \frac{1}{Z} \int_{U(1)^E} dh_e \mathcal{O}(h_e) \prod_f \sum_{n_f \in \mathbb{Z}} \exp(-n_f^2/2a_f + i\theta_f n_f) \left(\prod_{e \subset f} h_e^{[e,f]} \right)^{n_f} \\ &= \int_{U(1)^E} d\mu_{\Gamma}^{\vec{a}, \vec{\theta}}(h_e) \mathcal{O}(h_e) \quad a_f > 0, \theta_f \in \mathbb{R} \end{aligned}$$

coarse graining map:

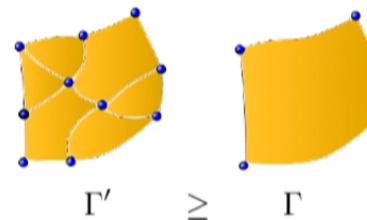


$$\pi_{\Gamma' \Gamma}(h_e) = (\dots, \underbrace{h_{e1} h_{e2} h_{e3}^{-1}}_{=h_{e'}}, \dots)$$

Example:

RG equations:

$$a_f^\Gamma = \sum_{f' \subset f} a_{f'}^{\Gamma'}$$
$$\theta_f^\Gamma = \sum_{f' \subset f} [f', f] \theta_{f'}^{\Gamma'}$$



Obvious solution:

$g \in \text{Sym}^2 T^* M$ (area) metric

$\theta \in \Omega^2(M)$ 2-form

$$a_f = \int_f d\text{vol}_g \quad \theta_f = \int_f \theta$$

Limit solutions:

$\theta = 0$ 2d Yang-Mills

$$g \rightarrow 0 \quad \int_{\mathcal{A}} \mathcal{D}A \delta(F[A] - \theta)$$

Example 2:

RG flow equations:

$$\omega_{(n_1, m_1, p_1), \dots, (n_E, m_E, p_E)}^{(\Gamma)} := \Omega_\Gamma(E^{n_1} F^{m_1} K^{p_1} \otimes \dots E^{n_E} F^{m_E} K^{p_E})$$

e.g. edge subdivision:

$$\omega_{\dots, (0, m, p), \dots}^{(\Gamma)} = \sum_{k=0}^m \left[\binom{m}{k} \right]_q \omega_{\dots, (0, k, p-m+k), (0, m-k, p+k) \dots}^{(\Gamma')}$$

Most obvious diffeo-invariant solution: $\Omega_\Gamma = \epsilon \otimes \dots \otimes \epsilon$

counit: $\epsilon : U_q \mathfrak{sl}_2 \rightarrow \mathbb{C}$

$$\epsilon(E) = \epsilon(F) = 0, \quad \epsilon(K^{\pm 1}) = 1, \quad \epsilon(ab) = \epsilon(a)\epsilon(b)$$

$$\Rightarrow \phi^* \Omega = \Omega$$

Summary:

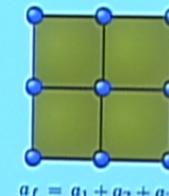
- There is a clear notion of background-independent Wilsonian RG flow
- Scale = discretization of manifold \mathcal{M} given by 2-complex Γ
- RG flow equation result from coarse graining scales: cylindrical consistency flow along partially ordered set!
- Solution to RG flow equations \Leftrightarrow continuum limit
- Clear notion of diff-invariance: strong extra conditions on solutions!

Example:

∞ many coupling constants for continuum measures $\mu_{g,\theta}$



a_f



$a_f = a_1 + a_2 + a_3 + a_4$

≡ information about background structure

fixing coupling constants up to certain scale Γ^e does not fix them for scales $\Gamma^f \geq \Gamma^e$
demanding in addition $\phi_*\mu = \mu$ fixes all of them

