

Title: Many-body localization: Local integrals of motion, area-law entanglement, and quantum dynamics

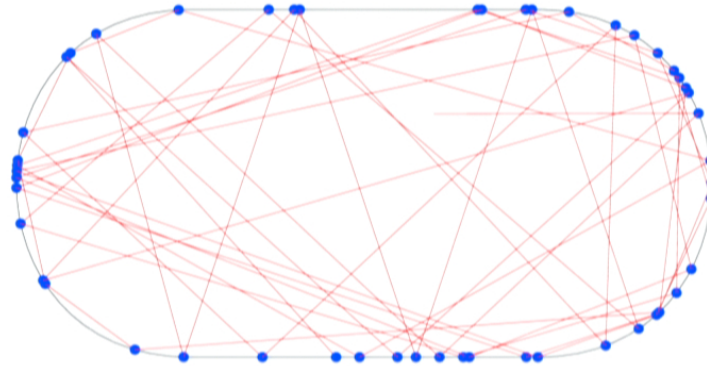
Date: Feb 12, 2014 09:45 AM

URL: <http://pirsa.org/14020127>

Abstract: We demonstrate that the many-body localized phase is characterized by the existence of infinitely many local conservation laws. We argue that many-body eigenstates can be obtained from product states by a sequence of nearly local unitary transformation, and therefore have an area-law entanglement entropy, typical of ground states. Using this property, we construct the local integrals of motion in terms of projectors onto certain linear combinations of eigenstates [1]. The local integrals of motion can be viewed as effective quantum bits which have a conserved z-component that cannot decay. Thus, the dynamics is reduced to slow dephasing between distant effective bits. For initial product states, this leads to a characteristic slow power-law decay of local observables, which is measurable experimentally, as well as to logarithmic in time growth of entanglement entropy [2,3]. We support our findings by numerical simulations of random-field XXZ spin chains. Our work shows that the many-body localized phase is locally integrable, reveals a simple entanglement structure of eigenstates, and establishes the laws of dynamics in this phase.
[1] M. Serbyn, Z. Papic, D. A. Abanin, Phys. Rev. Lett. 111, 127201 (2013).
[2] Jens H. Bardarson, Frank Pollmann, and Joel E. Moore, Phys. Rev. Lett. 109, 017202 (2012).
[3] M. Serbyn, Z. Papic, D. A. Abanin, Phys. Rev. Lett. 110, 260601 (2013)

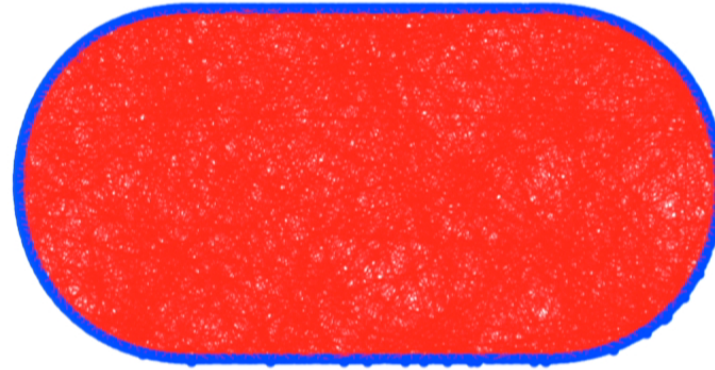
Ergodicity and its breaking

System explores full phase space



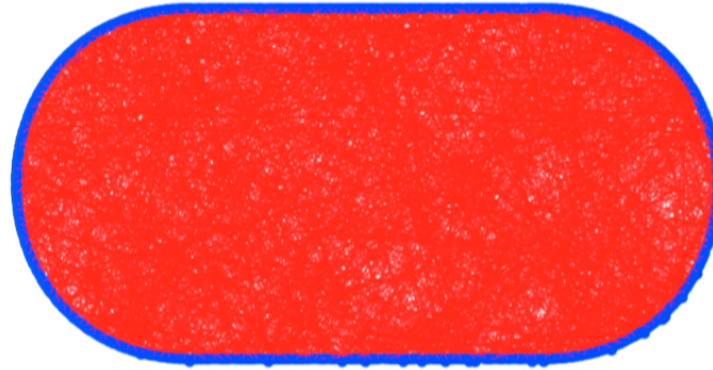
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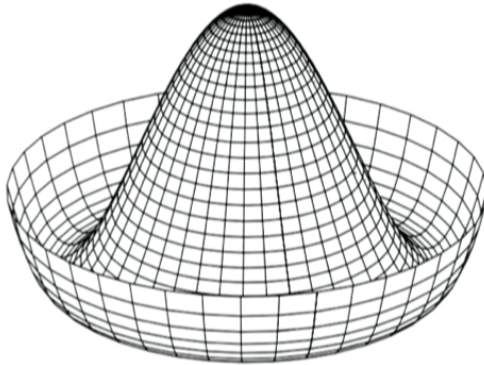
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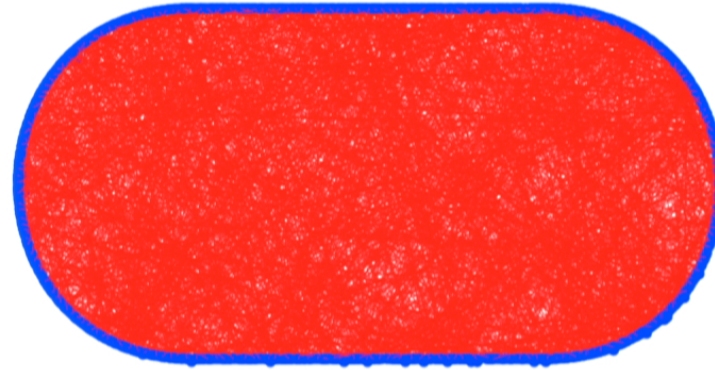
Ergodicity breaking

In phase transitions



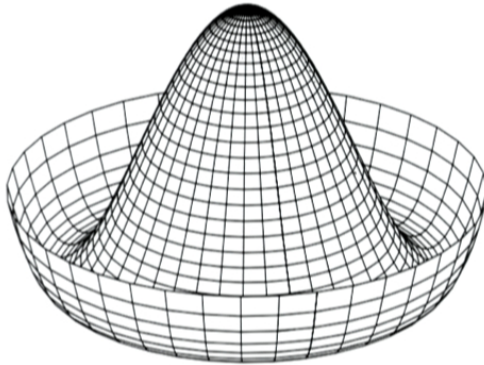
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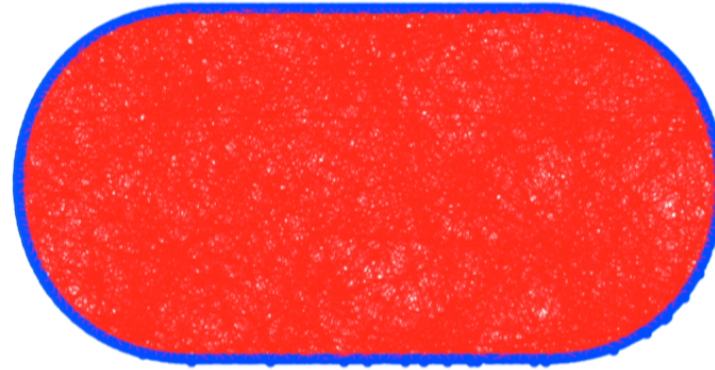
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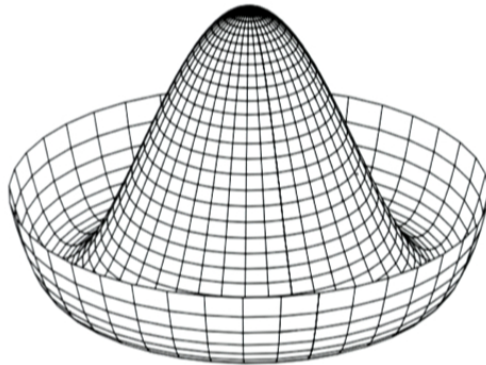
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Ergodicity breaking

In phase transitions



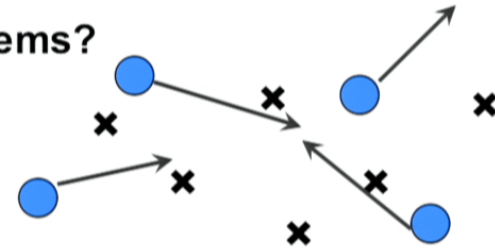
In bear habitats



Many-body localization problem

When/how does ergodicity break in many-body systems?

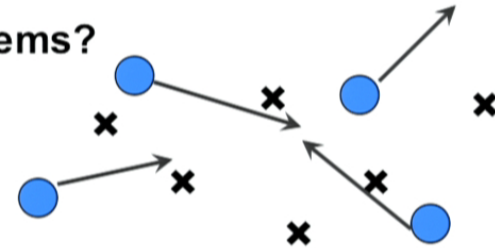
Do interactions destroy localization?



Many-body localization problem

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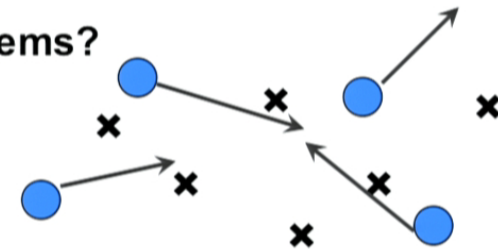
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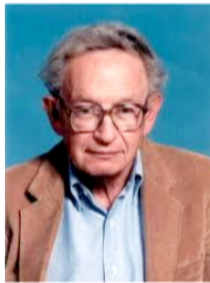
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YES



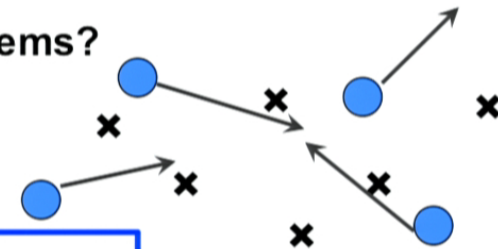
NO



Many-body localization problem

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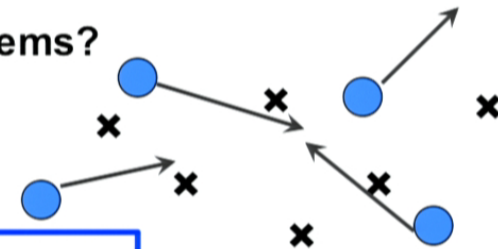
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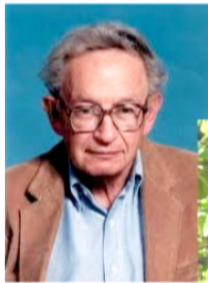
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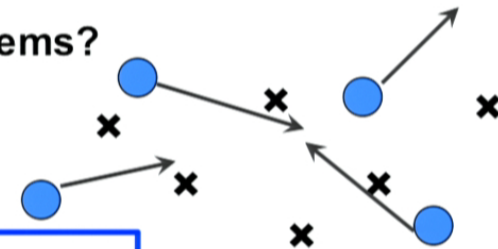
NO



Many-body localization problem

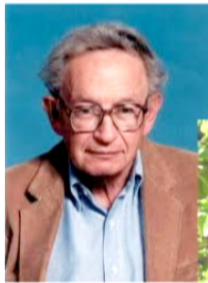
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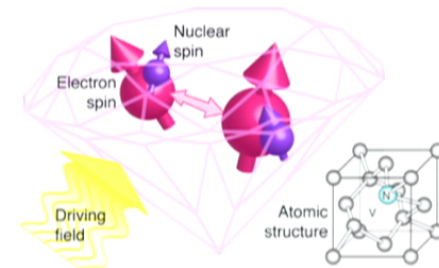
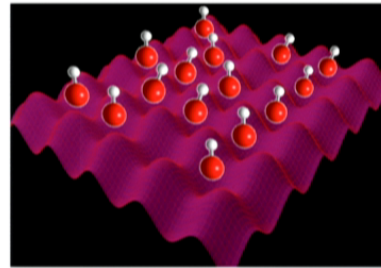
New experimental systems

Isolated & quantum-coherent. Tunable interactions and disorder

-Cold atoms, optical lattices

-Polar molecules

-Spin systems (NV-centers in diamond)



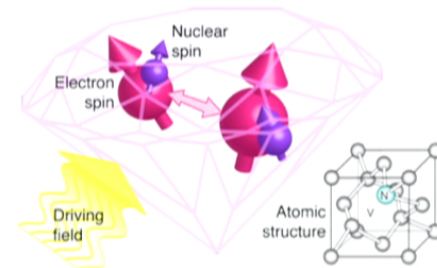
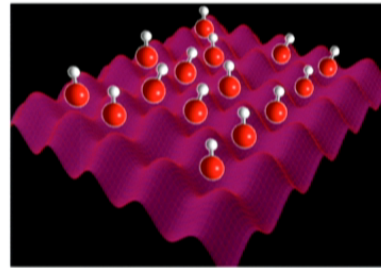
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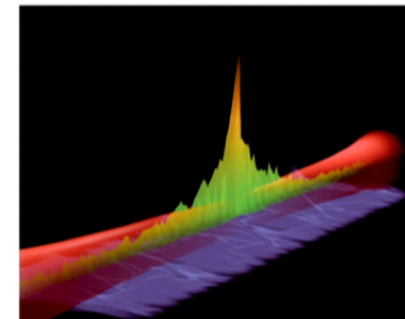
-Spin systems (NV-centers in diamond)



REPORTS

Three-Dimensional Anderson Localization of Ultracold Matter

S. S. Kondov, W. R. McGehee, J. J. Zirbel, B. DeMarco*



EXP: Ecole Normale,
Florence, Urbana

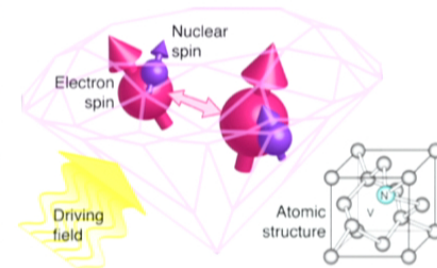
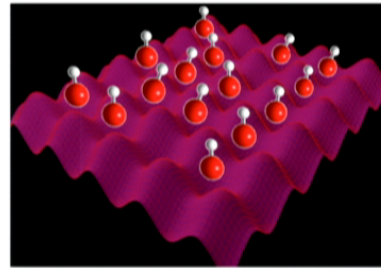
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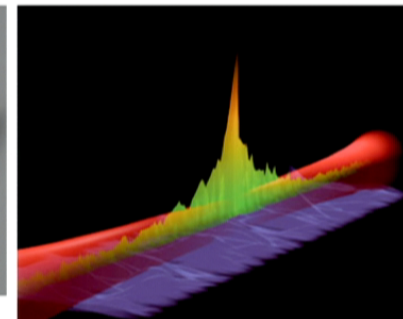
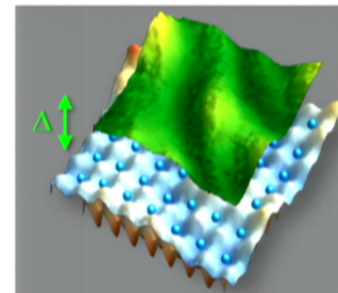
Interplay of disorder and interactions in an optical lattice

Hubbard model

S. S. Kondov, W. R. McGehee, B. DeMarco¹

(Dated: May 28, 2013)

Studying many-body localization experimentally now possible

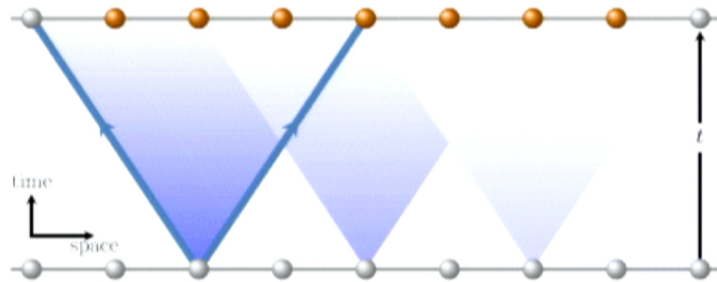


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Entanglement propagation in ergodic systems

Delocalized systems

Light-cone-like spreading of correlations



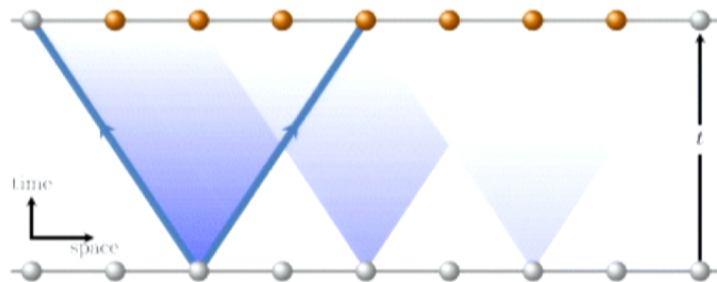
Lieb,Robinson'72, Hastings'04, Calabrese,Cardy'05

Kim, Huse'13

Entanglement propagation in ergodic systems

Delocalized systems

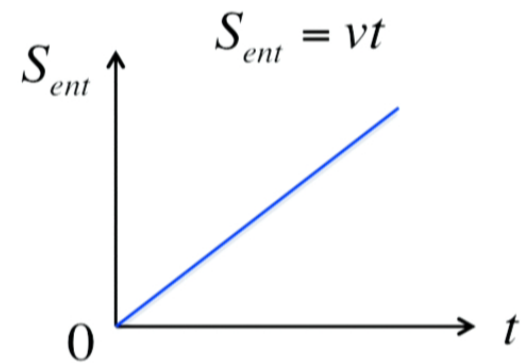
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Initial product state

Linear growth of entanglement entropy

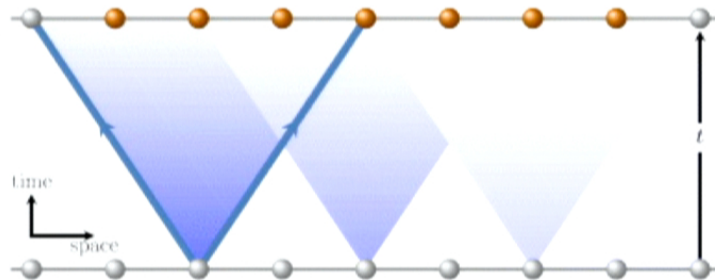


(even when transport is diffusive)

Entanglement propagation in ergodic systems

Delocalized systems

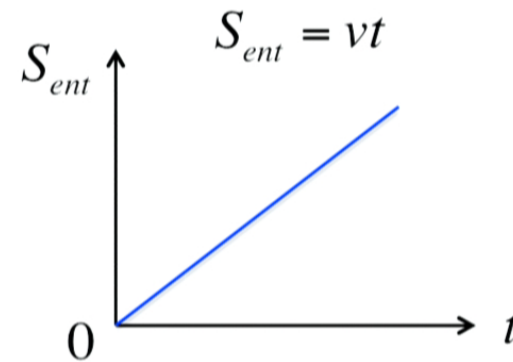
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Initial product state

Linear growth of entanglement entropy



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Experimental observation: Cold atoms [Cheneau et al'12](#) Trapped ions [Jurcevic et al'14](#)



A quantum many-body system

-Max-



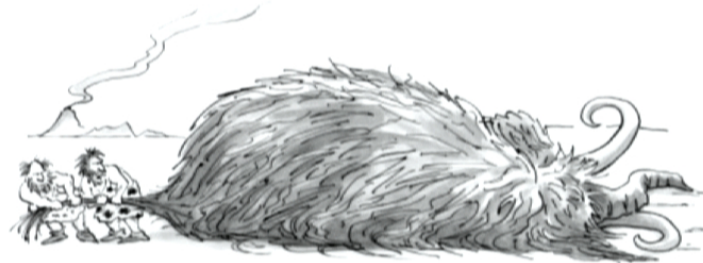
Condensed matter physicists



A quantum many-body system



Condensed matter physicists



System in a **ground state**

But there is much more...



This talk: Highly excited many-body states

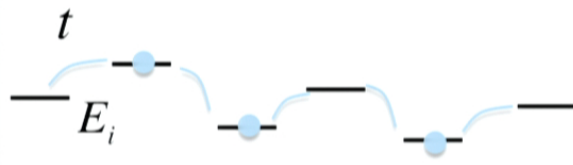
A simple model of many-body localization

Jordan-Wigner

Spinless interacting 1D fermions



Random-field XXZ spin-1/2 chain



$$H = \sum_i E_i c_i^+ c_i + t \sum_i c_i^+ c_{i+1} + h.c. + V \sum_i n_i n_{i+1}$$

$$H = \sum_i h_i S_i^z + J_{\perp} \sum_i (S_i^+ S_{i+1}^- + h.c.) + J_z \sum_i S_i^z S_{i+1}^z$$

Use disorder as tuning parameter $h_i \in [-W; W]$

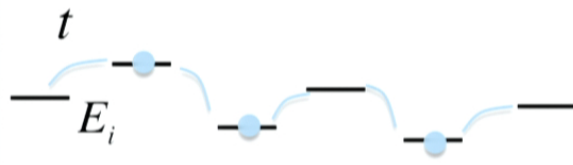
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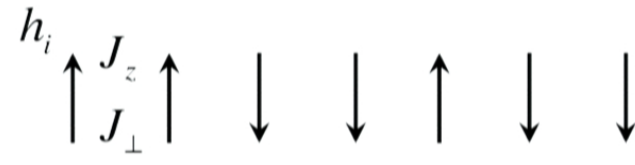
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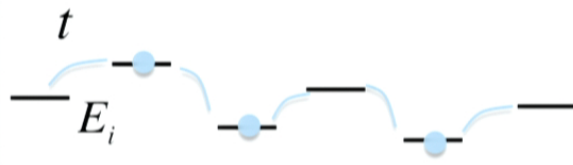
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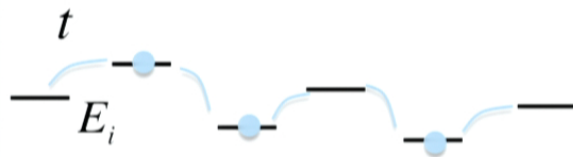
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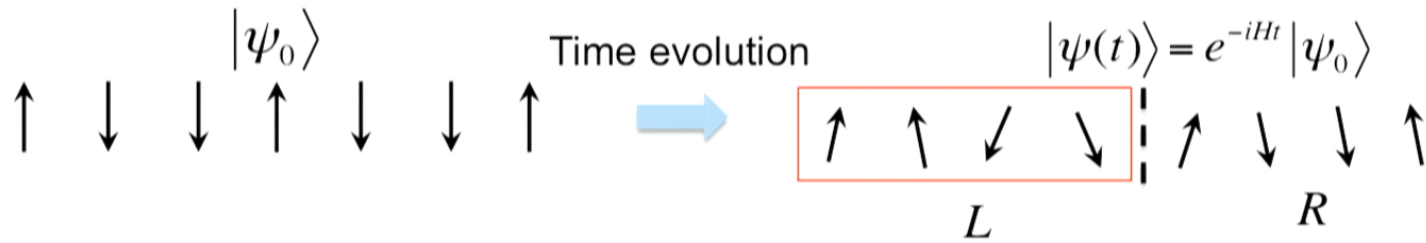
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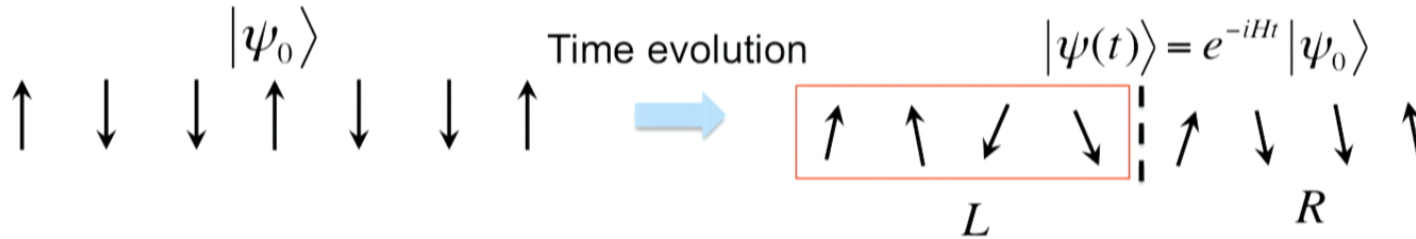
Many-body localization at strong disorder (numerics)

Oganesyan, Huse'07, Prosen et al'08, Pal, Huse'10, Monthus, Garel'10,..

Entanglement propagation in localized systems



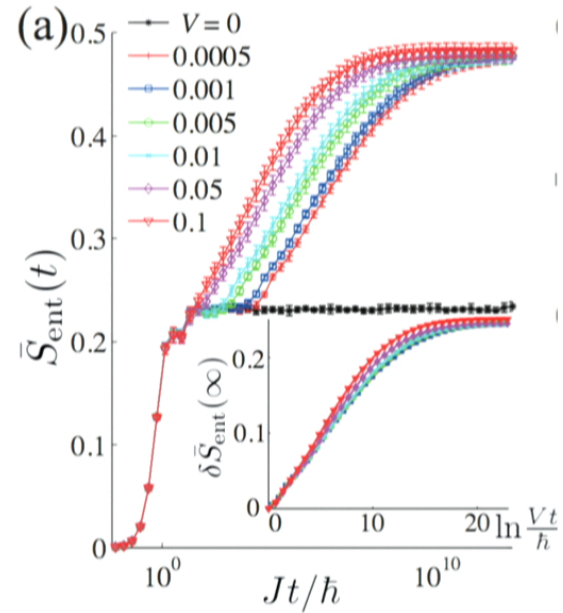
Entanglement propagation in localized systems



-Anderson-localized: $S_{ent}(t) \leq const$

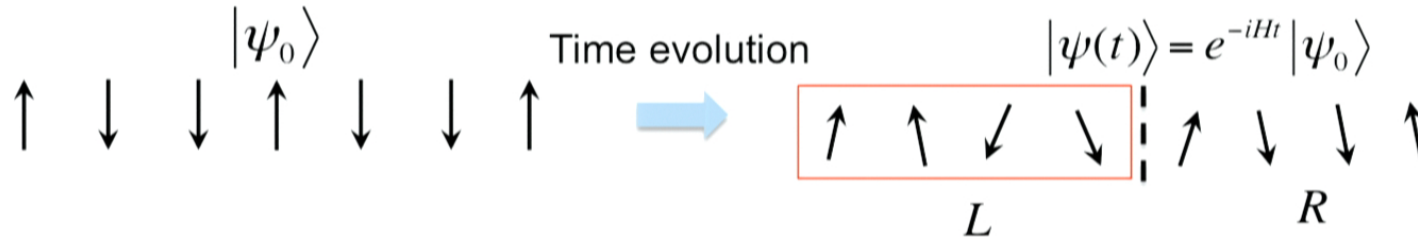
-Many-body localized: slow growth of entanglement

$$S_{ent}(t) \propto \log t$$



Bardarson, Pollmann, Moore'12

Entanglement propagation in localized systems



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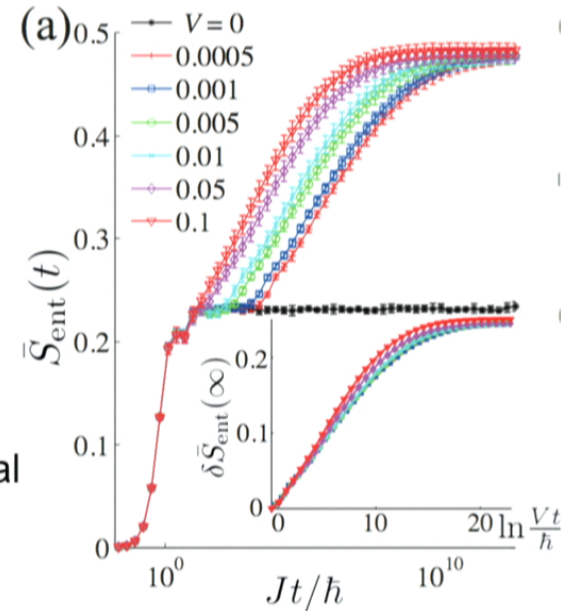
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$$S_{ent}(t) \propto \log t$$

-"Glassy" dynamics, extremely long time scales

-Entanglement extensive in system size, non-thermal

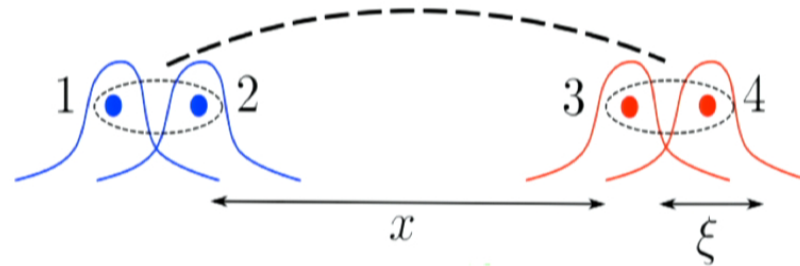
Ergodicity restored? Slow particle transport??



Bardarson, Pollmann, Moore'12

The mechanism of entanglement growth: Toy model

$$|\psi_0\rangle = \frac{1}{2}(c_1^\dagger + c_2^\dagger)(c_3^\dagger + c_4^\dagger)|0\rangle$$



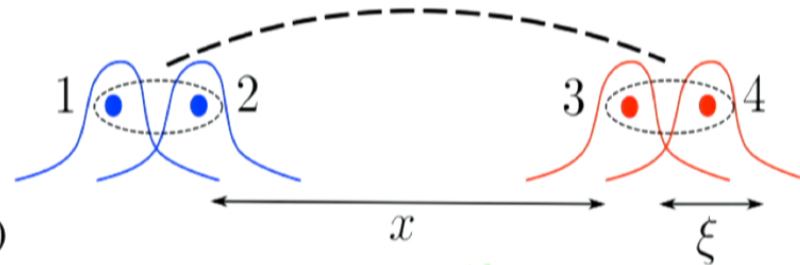
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Assume weak interactions

Eigenstate $|\alpha\beta\rangle = c_\alpha^+ c_\beta^+ |0\rangle + O(e^{-x/\xi})$

Energy: $E_{\alpha\beta} = E_\alpha + E_\beta + \boxed{C_{\alpha\beta} V e^{-x/\xi}}$



Reduced density matrix

$$\rho(t) = \frac{1}{2} \begin{bmatrix} 1 & \cos \omega t \\ \cos \omega t & 1 \end{bmatrix}$$

$$\omega \sim \frac{V}{\hbar} e^{x/\xi}$$

Serbyn, Papic, Abanin PRL '13

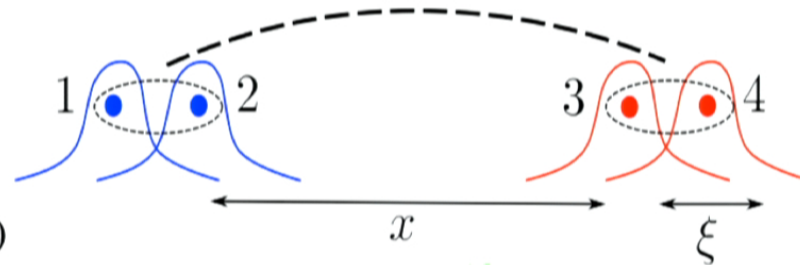
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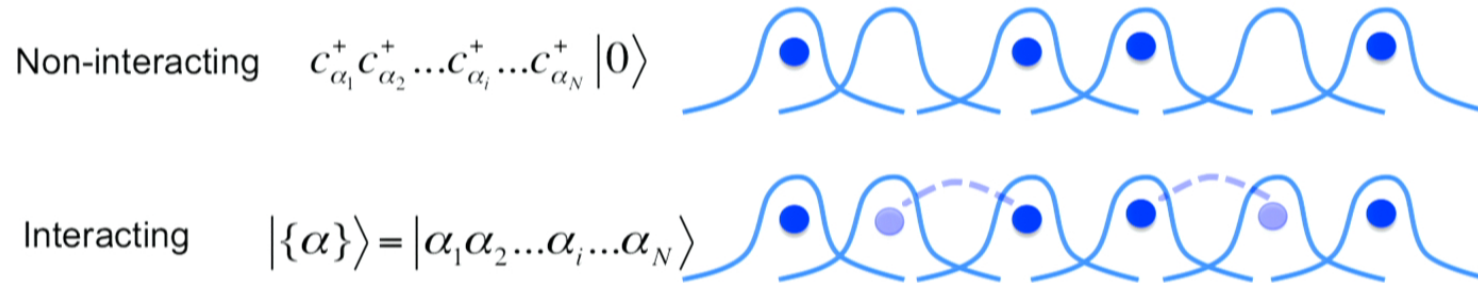
$$\omega \sim \frac{V}{\hbar} e^{x/\xi} \quad t_{deph} \sim \frac{\hbar}{V} e^{x/\xi}$$

Interaction-induced dephasing \rightarrow entanglement generation

Serbyn, Papic, Abanin PRL '13

Case of many particles

Hypothesis: Eigenstates at **small V** are “close”* to non-interacting eigenstates



Energy: perturbation theory in V

$$E_{\{\alpha\}} = \sum E_{\alpha_i} + V \sum C_{\alpha_i \alpha_j} e^{-\frac{|R_i - R_j|}{\xi}} + \dots$$

1-body
energy

2-body
interactions

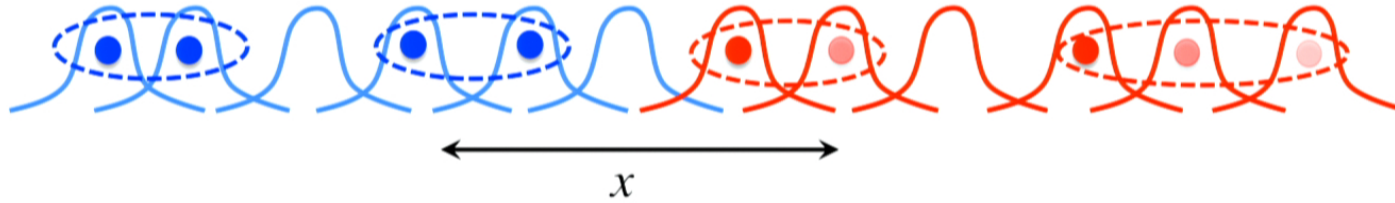
3,4...-body
interactions

Interactions of far-away particles are exponentially small

The laws of entanglement growth

Serbyn, Papic, Abanin PRL'13

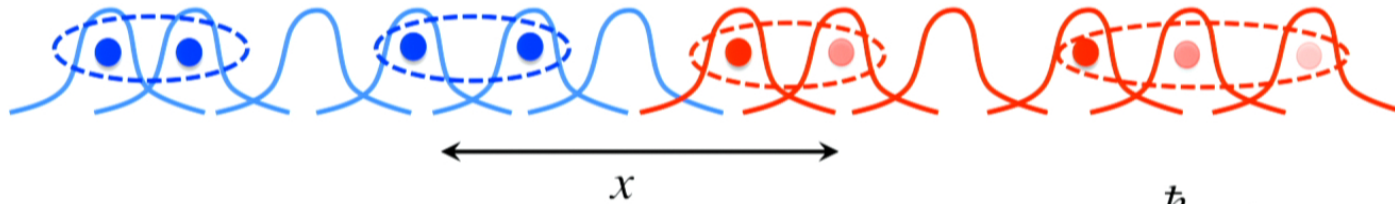
Initial product state is a superposition of many eigenstates



The laws of entanglement growth

Serbyn, Papić, Abanin PRL'13

Initial product state is a superposition of many eigenstates



Degrees of freedom a distance x away get entangled at $t_{deph} \sim \frac{\hbar}{V} e^{x/\xi}$

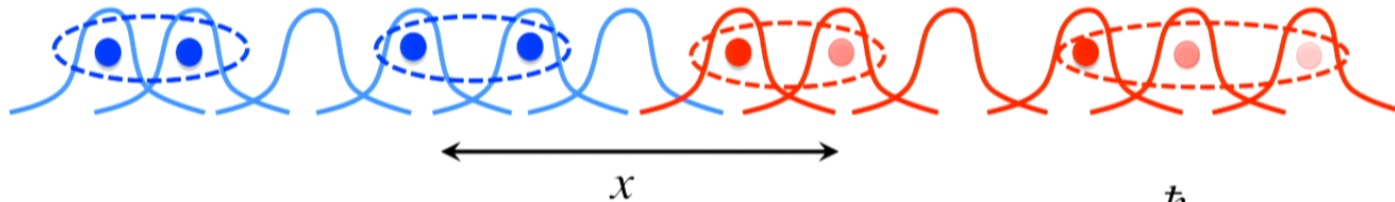
$$S_{ent}(t) = C \xi \log \frac{Vt}{\hbar}$$

C depends on diagonal entropy of initial state

The laws of entanglement growth

Serbyn, Papic, Abanin PRL'13

Initial product state is a superposition of many eigenstates



Degrees of freedom a distance x away get entangled at $t_{deph} \sim \frac{\hbar}{V} e^{x/\xi}$

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C depends on diagonal entropy of initial state

Saturated value: set by diagonal entropy

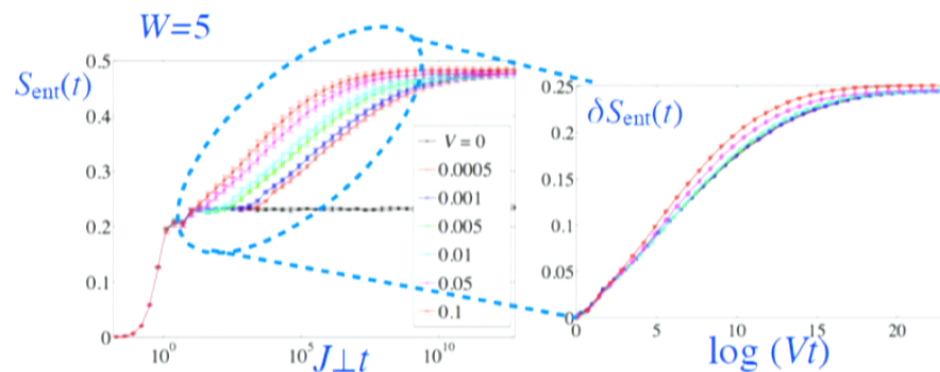
$$S_{ent}(\infty) = S_{diag} \ll S_{Thermal}$$

Predicted: propagation speed, dependence on disorder, interactions, initial state

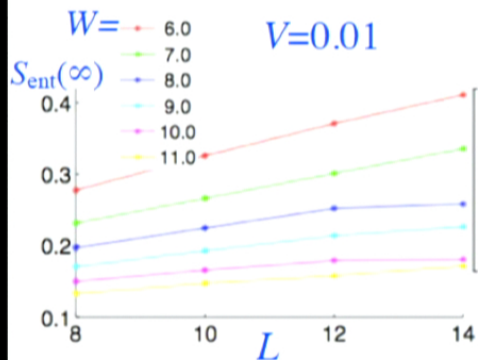
Confirmed by numerics

Numerical analysis

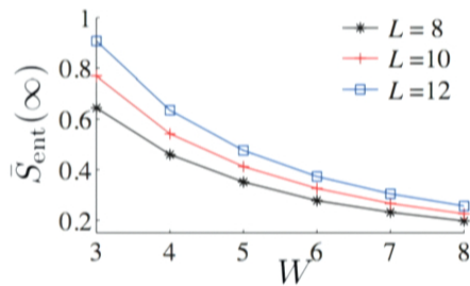
$$S_{ent}(t) \propto \xi \log \frac{Vt}{\hbar}$$



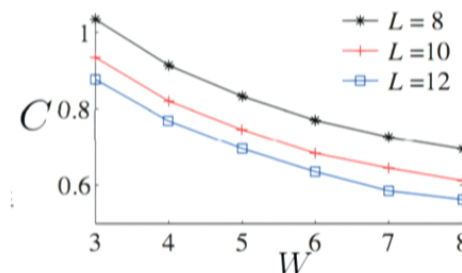
Interaction dependence



Saturated entanglement extensive in system size



Disorder dependence



$$C = \frac{S_{ent}(\infty)}{S_{diag}}$$

Localized phase at **strong interactions** (e.g., spins)?
Is entanglement growth universal?

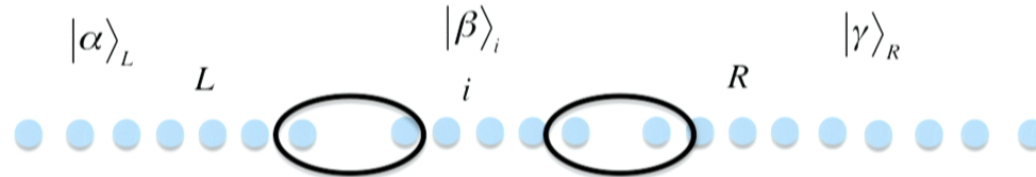
Localized phase at **strong interactions** (e.g., spins)?
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Localized phase at **strong interactions (e.g., spins)**?
Is entanglement growth universal?

YES!

The key: MBL phase is integrable. Integrals of motion are local.

Constructing integrals of motion



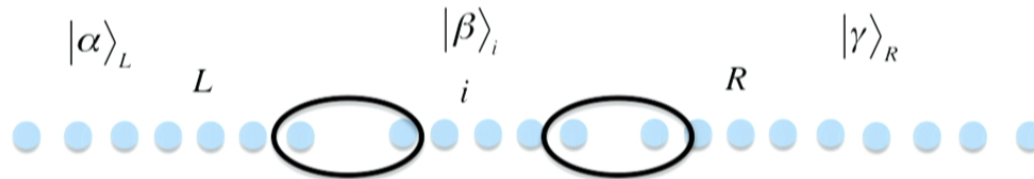
A local integral of motion in terms of projectors onto eigenstates

$$\hat{P}_i(\beta) = \sum_{\alpha, \gamma} |\alpha\beta\gamma\rangle \langle \alpha\beta\gamma|$$

$$|\alpha\beta\gamma\rangle = \hat{O}_{Ri} \hat{O}_{Li} |\alpha\rangle_L \otimes |\beta\rangle_i \otimes |\gamma\rangle_R$$

$$\hat{P}_i(\beta) = \sum_{\alpha, \gamma} |\alpha\beta\gamma\rangle \langle \alpha\beta\gamma| = \hat{O}_{Ri} \hat{O}_{Li} 1_L \otimes |\beta\rangle \langle \beta| \otimes 1_R \hat{O}_{Li}^+ \hat{O}_{Ri}^+$$

Constructing integrals of motion



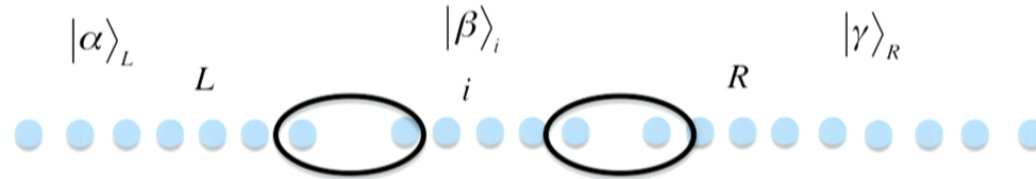
A local integral of motion in terms of projectors onto eigenstates

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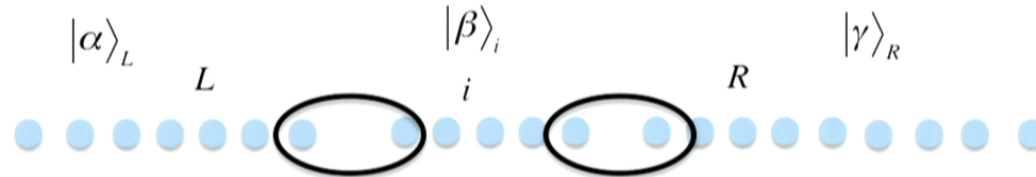
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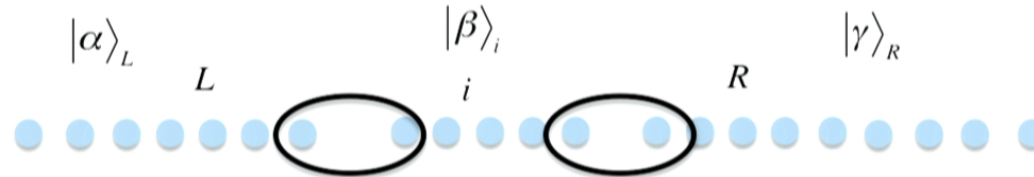
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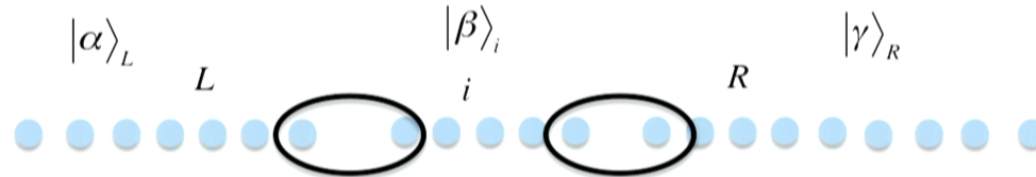
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These operators are local. Integrals of motion by construction.

Constructing integrals of motion



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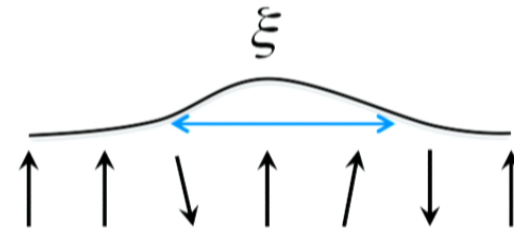
Universal Hamiltonian of MBL phase

"Effective" spins $\frac{1}{2}$ τ_z^i with conserved z-projection (defined via projectors)

$$[\tau_z^i, H] = 0$$

τ_z^i support in a region of size $\sim \xi$

Hamiltonian depends only on τ_z^i



$$H = \sum_i H_i \tau_z^i + \sum_{ij} H_{ij} \tau_z^i \tau_z^j + \sum_{ijk} H_{ijk} \tau_z^i \tau_z^j \tau_z^k + \dots$$

$$H_{ij} \propto \exp(-|i - j|a/\xi)$$

Universal dynamics and entanglement growth

-No relaxation \rightarrow no thermalization

$$\langle \tau_z^i(t) \rangle = \text{Const}$$

-Interaction-induced dephasing

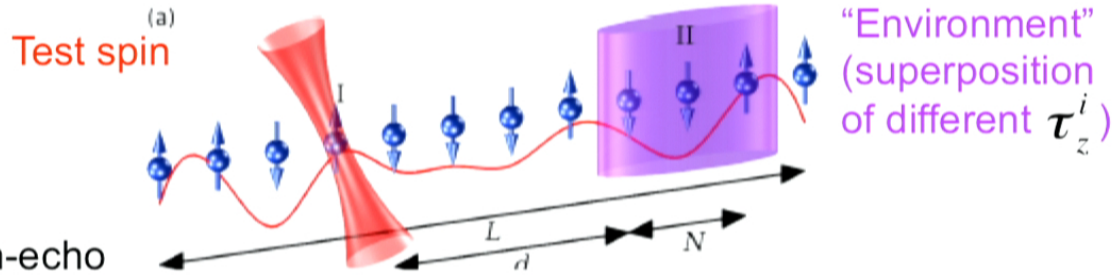
$$\langle \tau_x^i(t) \rangle \rightarrow 0 \quad t \rightarrow \infty$$

-Steady non-thermal state

-Universal logarithmic growth of entanglement entropy (hard to measure)

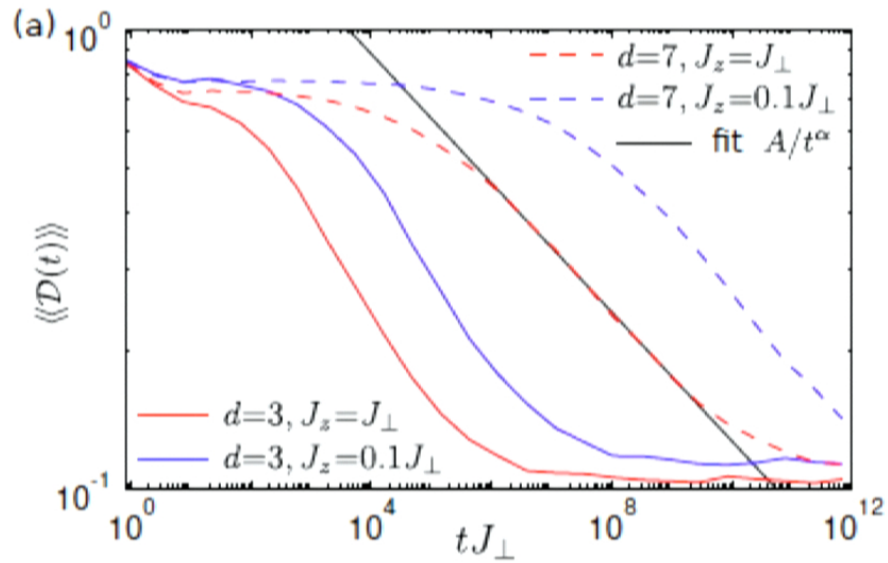
-Any local observable would show oscillations and a power-law decay (easy to measure)

Interferometric signatures of many-body localization



Dephasing \rightarrow power-law spin-echo decay $\bar{D}(t) \sim t^{-\xi \ln 2}$

A way to directly probe MBL
in spin systems, cold atoms



Serbyn, Papić, Abanin+Lukin, Demler'14

The structure of many-body localized states

- Eigenstates obtained from product states by a **sequence of local unitaries** (quantum circuit of finite depth)
- Excited MBL eigenstates obey **area law-entanglement** (similar to ground states of gapped systems!)

Explicitly constructing integrals of motion

Effective spins: good for quantum information processing?

Yes, but how do they relate to physical spins?

Approach 1: Connect eigenstates to product states by local unitary U

$$\tau_z^i = U s_z^i U^\dagger$$

Expand τ_z^i via physical operators

$$\tau_z^i = C_1 s_z^i + C_2 s_z^{i+1} + C_3 s_z^{i-1} + D_1 s_z^i s_x^{i+1} + \dots$$

with A. Chandran, G. Vidal, in progress

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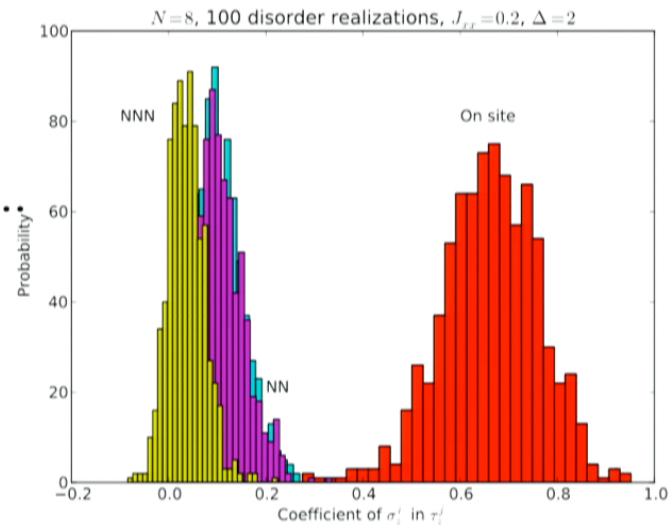
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Approach 2: Extension of a strong disorder renormalization group



Extract statistics of couplings and properties of integrals of motion

with A. Chandran, G. Vidal, in progress

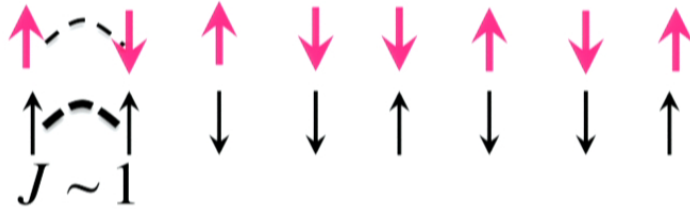
Localization without quenched disorder?

(Muller'13, Huveneers, de Roeck'13, Fisher, Grover'13..)

Two coupled spin chains, "fast" σ_i^z and "slow" S_i^z

Papic, Stoudenmire, DA '14

$$\lambda \ll 1$$



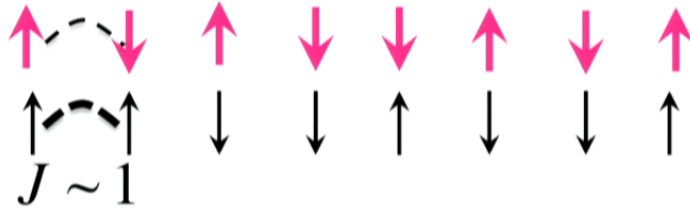
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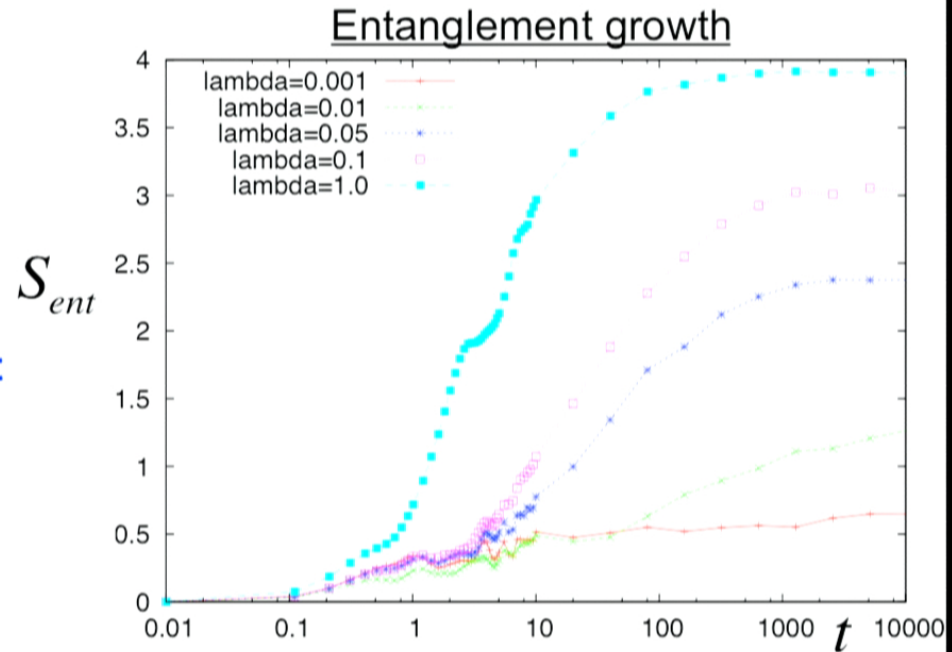
Interaction
$$H_{\text{int}} = W \sum_i \sigma_i^z S_i^z$$

Localization at $\lambda \ll 1$?

Preliminary:

Log-growth of entanglement

Signals localization??



Summary

- Ergodicity breaking in disordered systems: local conservation laws
- Universal dynamics: slow dephasing
- Logarithmic growth of entanglement entropy
- More to come...

DETAILS:

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Phys. Rev. Lett. **111**, 127201 (2013)+To appear

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(Perimeter)



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