

Title: 13/14 PSI - Condensed Matter Review - Lecture 3

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Abstract:

$$H = -h \sum_i^N s_i - J \sum_i^N s_i s_{i+1}$$

$$\boxed{\begin{matrix} \lambda = \beta h \\ K = \beta J \end{matrix}}$$

$$\langle \hat{\sigma}_i^z \hat{\sigma}_j^z \rangle = \frac{1}{Z} \text{Tr} [e^{-H} \hat{\sigma}_i^z \hat{\sigma}_j^z] =$$

$$\langle_i | \hat{\sigma}_i^z | i \rangle = s_i$$

$$Z = \text{Tr} e^{-\beta H} = \text{Tr} (T^N)$$

$$T = \begin{pmatrix} e^{\lambda+K} & e^{-K} \\ e^{-K} & e^{-\lambda+K} \end{pmatrix} = \begin{pmatrix} e^{\lambda} & \\ & e^{-\lambda} \end{pmatrix} \begin{pmatrix} e^K & e^{-K} \\ e^{-K} & e^K \end{pmatrix}$$

$$\begin{cases} T_1 = e^{\lambda \hat{\sigma}_x} \\ T_2 = e^K [1 + e^{-2K} \frac{\hat{\sigma}_x}{\sigma}] \end{cases}$$

$$\begin{matrix} ||| \\ \hline T_1 \\ ||| \end{matrix} \quad \begin{matrix} ||| \\ \hline T_2 \\ ||| \end{matrix}$$

$$H = -h \sum_i^N s_i - J \sum_i^N s_i s_{i+1}$$

$$\boxed{\begin{array}{l} \lambda = \beta h \\ K = \beta J \end{array}}$$

$$\langle \hat{\sigma}_i^z \hat{\sigma}_i^z \rangle = \frac{1}{Z} \text{Tr} [e^{-H} \hat{\sigma}_i^z \hat{\sigma}_i^z]$$

$$\langle \cdot | \hat{\sigma}_i^z | \cdot \rangle = s_i$$

$$Z = \text{Tr} e^{-\beta H} = \text{Tr} (T^N) = \lambda_1^N + \lambda_2^N$$

$$T = \begin{pmatrix} e^{\lambda+K} & e^{-K} \\ e^{-K} & e^{-\lambda+K} \end{pmatrix} = \begin{pmatrix} e^{\lambda} & \\ & e^{-\lambda} \end{pmatrix} \begin{pmatrix} e^K & e^{-K} \\ e^{-K} & e^K \end{pmatrix}$$

$$\begin{cases} T_1 = e^{\lambda+K} \\ T_2 = e^{-K} [1 + e^{-2K} \frac{\lambda}{\sigma}] \end{cases}$$

$$\lambda_{1,2} = e^K \left[\cosh \lambda \pm \sqrt{2 \sinh^2 \lambda + e^{-4K}} \right] \quad \begin{array}{l} ||| \\ T_1 \\ ||| \\ T_2 \end{array}$$

$$\langle \sigma_i^z \sigma_j^z \rangle = \frac{1}{\lambda_1^N + \lambda_2^N} \text{Tr} \left[\sigma_j^z T_2^N \sigma_i^z \right] = \frac{1}{\lambda_1^N + \lambda_2^N} \text{Tr} \left[T_2^{N-i} \sigma_j^z T_2^{i-j} \sigma_i^z T_2^j \right]$$

$$h=0 \Rightarrow \lambda=0$$

$$T_2^N = T_2^{N-i} T_2^{i-j} T_2^j$$

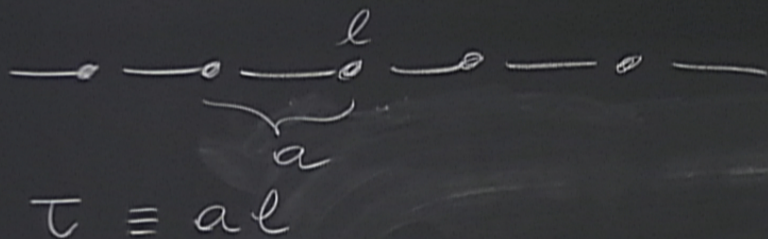
$$T_2 |\pm\rangle = \lambda_{1,2} |\pm\rangle$$

$$\sigma^z = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$= \frac{\lambda_1^{N-i} \lambda_2^{N-i+j} + \lambda_2^{N-i+j} \lambda_1^{i-j}}{\lambda_1^N + \lambda_2^N} \xrightarrow[N \rightarrow \infty]{TDL}$$

$$C_{ij} = \langle \sigma_i^z \sigma_j^z \rangle = \left(\tanh K \right)^{|i-j|}$$

σ^z



$$\tau \equiv al$$

$$\xi^{-1} = \frac{1}{a} \log \coth hK$$

$$C(\tau) = e^{-|\tau|/\xi}$$

$$\frac{a}{\xi} \approx 2e^{-2K} ; L = Na ; \tilde{h} = \frac{h}{a}$$

for large K (large β)
 ξ does diverge

$$L = Na; \quad \tilde{h} = \frac{h}{a}$$

scaling limit

→
 a small
 ξ large

$$\lambda_{1,2} \approx \sqrt{\frac{2\xi}{a}} \left[1 \pm \frac{a}{2\xi} \sqrt{1 + 4\tilde{h}^2 \xi} \right]$$

$$\rightarrow f = \frac{E_0}{a|k|} - \frac{1}{L} \log \left[2 \cosh \left(L \sqrt{\frac{1}{4\xi^2} + \tilde{h}^2} \right) \right]$$

$$L = Na; \quad \tilde{h} = \frac{h}{a}$$

scaling limit

→
 a small
 ξ large

$$\lambda_{1,2} \approx$$

$$\left[\frac{2\xi}{a} \left[1 \pm \frac{a}{2\xi} \sqrt{1 + 4\tilde{h}^2 \xi^2} \right] \right]$$

$$\rightarrow f = E_0 - \frac{1}{L} \log \left[2 \cosh \left(L \sqrt{\frac{1}{4\xi^2} + \tilde{h}^2} \right) \right]$$

$$= \frac{1}{2} \frac{1}{K}$$

$$\rightarrow C(T) =$$

$$\frac{e^{-L/2\xi} + e^{-L/2\xi}}{1 + e^{-L/\xi}}$$

$$f = \overline{\Phi}_f \left(\frac{L}{\sqrt{3}}, L\tilde{h} \right)$$

$$C(\tau) = \overline{\Phi}_\sigma \left(\frac{L}{\sqrt{3}}, \tau \frac{L}{\sqrt{3}} \right)$$

Universality: $\overline{\Phi}_f, \overline{\Phi}_\sigma$ are universal functions

$$e^{-2K} \approx \frac{a}{2\zeta}$$

$$E_0 \approx -\frac{\hbar^2 K^2}{2m}$$

$$\lambda = \frac{h}{\hbar K}$$

Mapping Classical-Quantum

$$T \approx e^{\frac{1}{\hbar} \int_{x_1}^{x_2} a \sqrt{-E_0 + \frac{1}{2m} \hbar^2 k^2} dx}$$

$$e^{a\theta_1} e^{a\theta_2} = e^{a(\theta_1 + \theta_2)} + \mathcal{O}(a^2)$$

$$e^{-2K} \approx \frac{a}{2\zeta}$$

$$E_0 \approx -\frac{\hbar^2 K^2}{2m}$$

$$\lambda = \frac{h}{\hbar K}$$

Mapping Classical-Quantum

$$T \approx e^{\frac{1}{\hbar} \int_{T_1}^{T_2} a(-E_0 + \frac{1}{2m} \hbar^2 k^2) dx} \approx e^{a \left[\frac{\hbar^2 k^2}{2m} - E_0 + \frac{1}{2m} \hbar^2 k^2 \right]}$$

$$e^{a\theta_1} e^{a\theta_2} = e^{a(\theta_1 + \theta_2)} + \mathcal{O}(a^2)$$

$$e^{-2K} \approx \frac{a}{2\xi}$$

$$E_0 \approx -\frac{\hbar^2 K^2}{2m}$$

$$\lambda = \frac{a}{\hbar K}$$

Mapping Classical-Quantum

$$T \approx e^{-\frac{1}{\hbar} \int_{T_1}^{T_2} \sqrt{2m(V(x) - E_0)} dx} \approx e^{-\frac{1}{\hbar} \int_{T_1}^{T_2} \sqrt{2m(V(x) - E_0 + \frac{\hbar^2}{2m\xi^2})} dx}$$

$$\xi^{-1} \equiv \Delta$$

$$H_Q = \frac{\hbar^2}{2m} \frac{d^2}{dx^2} - E_0 + \frac{\Delta}{2} \frac{d}{dx}$$

$$e^{-2K} \approx \frac{a}{2\zeta}$$

$$E_0 \approx -\frac{\hbar^2 k^2}{2m}$$

$$\lambda = a\hbar$$

$$T_{S_1 S_2}$$



Mapping Classical-Quantum

$$T \approx e^{\frac{1}{T_1} \hbar \sigma^z} e^{\frac{1}{T_2} a(-E_0 + \frac{1}{2\zeta} \sigma^x)}$$

$$\approx e^{-a \left[\hbar \sigma^z + E_0 + \frac{1}{2\zeta} \sigma^x \right]}$$

$$\frac{1}{\zeta} \equiv \Delta$$

$$a \sim \tau$$

$$T = e^{-\tau H_Q}$$

$$H_Q = -\hbar \sigma^z + E_0 - \frac{\Delta}{2} \sigma^x$$

$$e^{-2K} \approx \frac{a}{2\xi}$$

$$E_0 \approx -\frac{\hbar^2 k^2}{2m}$$

$$\lambda = a\hbar$$

Mapping Classical-Quantum

$$T \approx e^{-\int_{T_1}^{T_2} \sqrt{2m(E_0 - V(x))} dx} \approx e^{-a \left[\hbar \frac{1}{2} + E_0 - \frac{1}{2} \frac{\hbar^2 x^2}{2m} \right]}$$

$$\xi^{-1} \equiv \Delta$$

$T_S: S_{cl}$

$$a \sim \tau$$

$$T = e^{-\tau H_Q}$$

$$H_Q = -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} + E_0 - \frac{\Delta}{2} \frac{\partial^2}{\partial x^2}$$



$$Z = \text{Tr} T^N \approx \text{Tr} e^{-N H_Q} = \text{Tr} e^{-\frac{H_Q}{T}}$$

$$\frac{1}{T} = L$$

partitioning

$$f = -T \log Z = E_0 - T \log \left[2 \cosh \sqrt{\left(\frac{\Delta}{2}\right)^2 + \tilde{h}^2} / T \right]$$
$$E_{1,2} = E_0 \pm \sqrt{\tilde{h}^2 + \left(\frac{\Delta}{2}\right)^2}$$

parabolas

$$f = -T \log Z = E_0 - T \log \left[2 \cosh \left(\frac{\Delta}{2} \sqrt{\frac{m}{\hbar}} \right) / T \right]$$

$$E_{1,2} = E_0 \pm \sqrt{\hbar^2 + \left(\frac{\Delta}{2}\right)^2}$$

1D classical
Stat-mech



0-D Quantum
Im. time evl.

$\frac{1}{T}$

$$H = -Jg \sum_i^M \sigma_i^x - J \sum_i^M \sigma_i^z \sigma_{i+1}^z$$

$$L_{\text{space}} = Ma$$

$$a \quad U_a = e^{-aH} = T_1 T_2 + \mathcal{O}(a^2)$$

$$\left\{ \begin{aligned} T_1 &= \exp [Jga \sum_i \sigma_i^x] \\ T_2 &= \exp [Ja \sum_i \sigma_i^z \sigma_{i+1}^z] \end{aligned} \right.$$

$$\frac{1}{T} = e^{-\frac{H}{T}}$$

$$\frac{1}{T} \equiv L_{\tau} = Na$$

$\frac{1}{T}$

$$H = -Jg \sum_i^M \sigma_i^x - J \sum_i^M \sigma_i^z \sigma_{i+1}^z$$

$$L_{\text{space}} = Ma$$

$$a \quad \mathcal{U}_a = e^{-aH} = T_1 T_2 + \mathcal{O}(a^2)$$

$$T_1 = \exp [Jga \sum_i \sigma_i^x]$$

$$T_2 = \exp [Ja \sum_i \sigma_i^z \sigma_{i+1}^z]$$

$$\frac{1}{T} = e^{-\frac{H}{T}}$$

$$= \text{Tr} (e^{-aH})^N \quad \frac{1}{T} \equiv L_{\tau} = Na$$

$$Z = \text{Tr} (T_1 T_2)^N$$

$$T_2 | \{ \sigma_i^z \} \rangle = e^{-J a \sum_i \sigma_i^z \sigma_{i+1}^z} | \{ \sigma_i^z \} \rangle$$

$$Z = \sum_{\{ \sigma_i^z(l) \}} \prod_{l=1}^N \langle \{ \sigma_i^z(l) \} | T_1 T_2 | \{ \sigma_i^z(l+1) \} \rangle = \sum_{\{ \sigma_i^z(l) \}} \exp \left[\sum_{i,l} J a \sigma_i^z \sigma_{i+1}^z \right]$$

$$\langle \{ \sigma_i^z(l) \} | T_1 | \{ \sigma_i^z(l+1) \} \rangle = A e^{B \sum_i \sigma_i^z(l) \sigma_i^z(l+1)}$$

$$A = \frac{1}{2} \cosh(2Jga)$$

$$e^{-2B} = \tanh(Jga)$$

$$\exp \left[\sum_{iq} \left(J_a \sigma_i^z(l) \sigma_{i+1}^z(l) + B \sigma_i^z(l) \sigma_i^z(l+1) \right) \right]$$

$$\exp \left[\sum_{iq} \left(J_a \sigma_i^z(l) \sigma_{i+1}^z(l) + B \sigma_i^z(l) \sigma_i^z(l+1) \right) \right]$$

$- H_{2D}$