

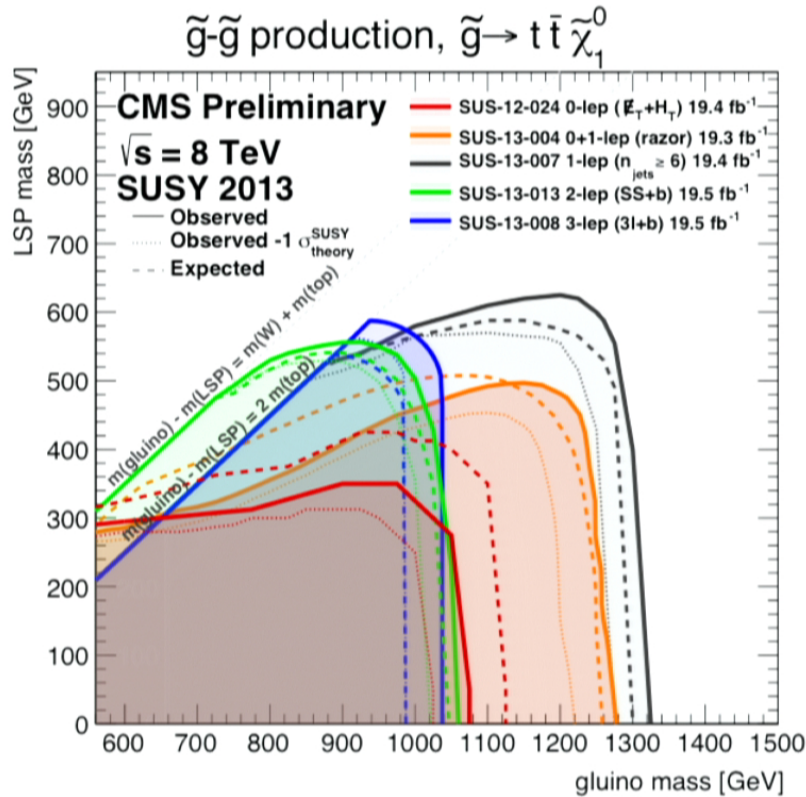
Title: BBN and the Unnatural Weak Scale

Date: Dec 05, 2013 03:00 PM

URL: <http://pirsa.org/13120058>

Abstract: The discovery of a perturbatively-coupled, 125 GeV Higgs, together with the absence of LHC signals for supersymmetry, places the principle of naturalness under tension. In this talk I will discuss the possibility that the weak scale is unnatural, with its value determined environmentally in the landscape. In particular, this environmental selection may be driven by BBN: as the weak scale is increased, the abundance of Hydrogen in the early universe is rapidly depleted. That our own universe contains an  $O(1)$  fraction of primordial Helium arises as a conspiracy among the weak scale, the neutron-proton mass difference, and the planck scale.

# naturalness under siege



$$m_{\tilde{g}} \gtrsim 1 \text{ TeV}$$

(unless no tops, no MET, few high- $p_T$  jets)

Evans, et al, 1310.5758

$$\delta m_{H_u}^2 \sim \frac{2}{\pi^2} y_t^2 \left( \frac{\alpha_s}{\pi} \right) |M_3|^2 \log^2 \left( \frac{\Lambda}{\text{TeV}} \right)$$

$$\implies \Delta \gtrsim 10 - 100$$

# environmental selection

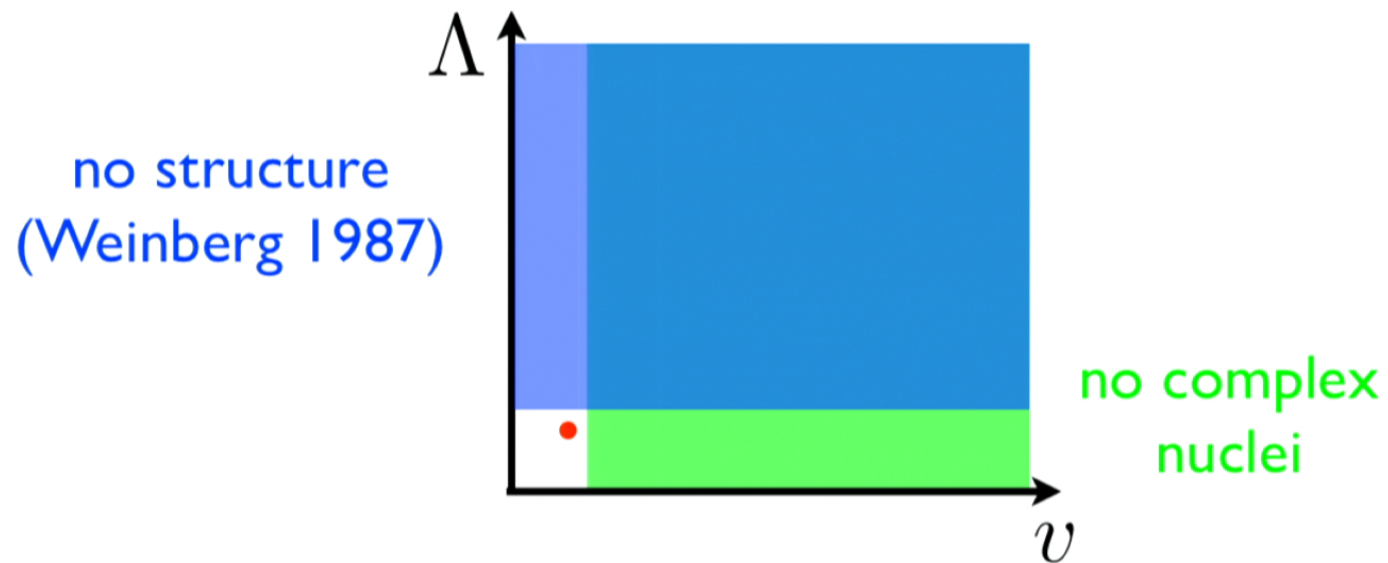
An anthropic “model” consists of:

- scanning parameters
- catastrophic boundaries
- prior distributions

Fixing these, we can ask how likely our universe is in the landscape

# minimal landscape

assume only dimensionful parameters scan:



Arkani-Hamed, Dimopoulos, Kachru 0501082



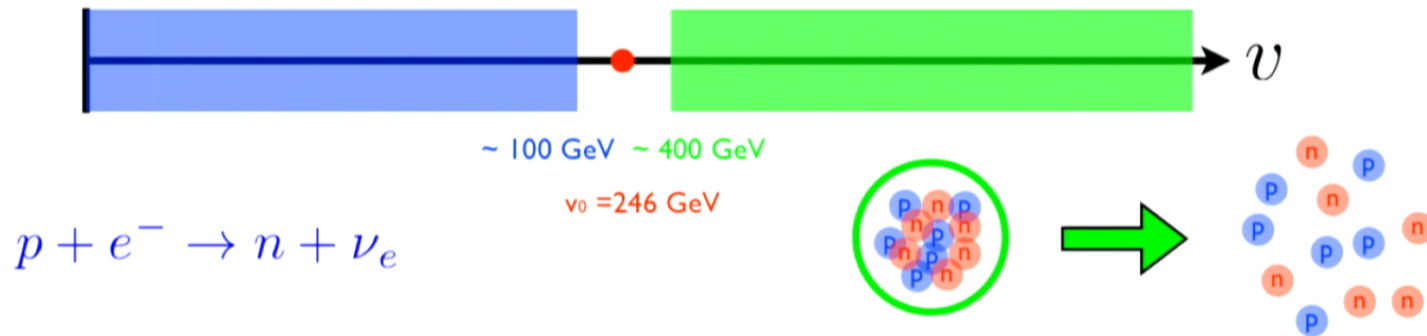
# dangers of a variable weak scale

$$m_u = y_u v$$

$$m_d = y_d v$$

hydrogen unstable

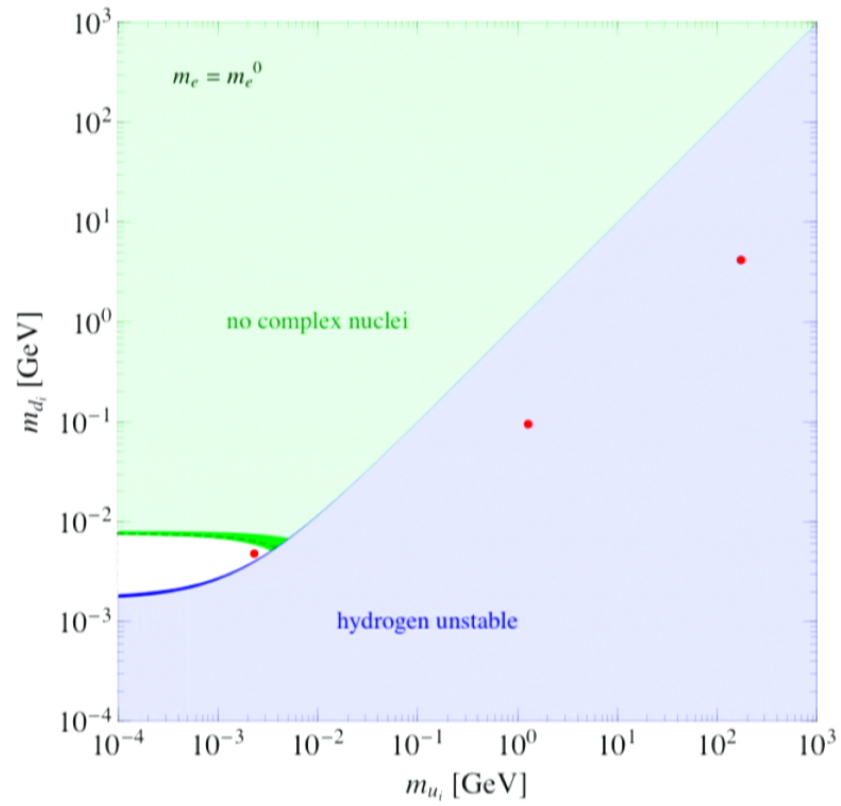
complex nuclei  
unbound



- Agrawal, Barr, Donoghue, Seckel 9707380
- Damour, Donoghue 0712.2968

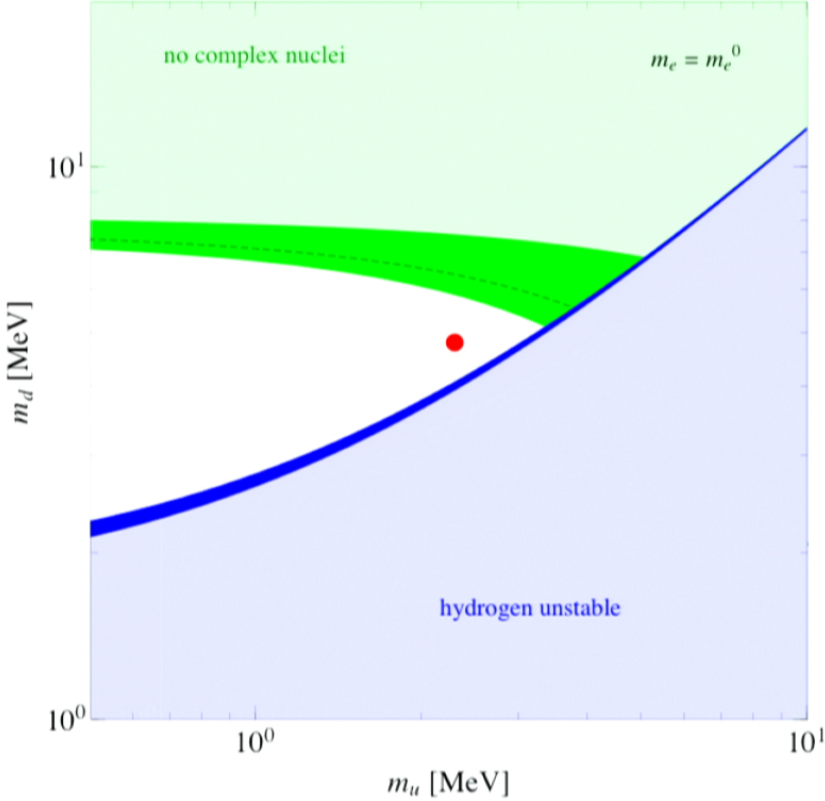
# general landscape

scan:  
 $(y_u, y_d, v)$



# general landscape

scan:  
 $(y_u, y_d, v)$



# runaway to large $v$ ?

scan:

$(y_u, y_d, v)$

- nuclear physics depends on the quark masses
- runaway: increase  $v$ , fixing quark masses,

$$y_{u,d} \rightarrow \frac{m_{u,d}}{v}$$

“Weakless Universe,” Harnik, Kribs, Perez 0604027

# weak-scale physics in our Universe

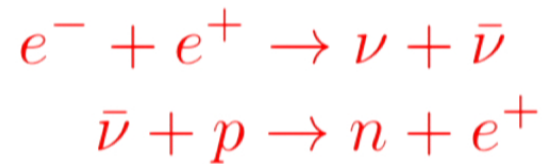
## 1. BBN



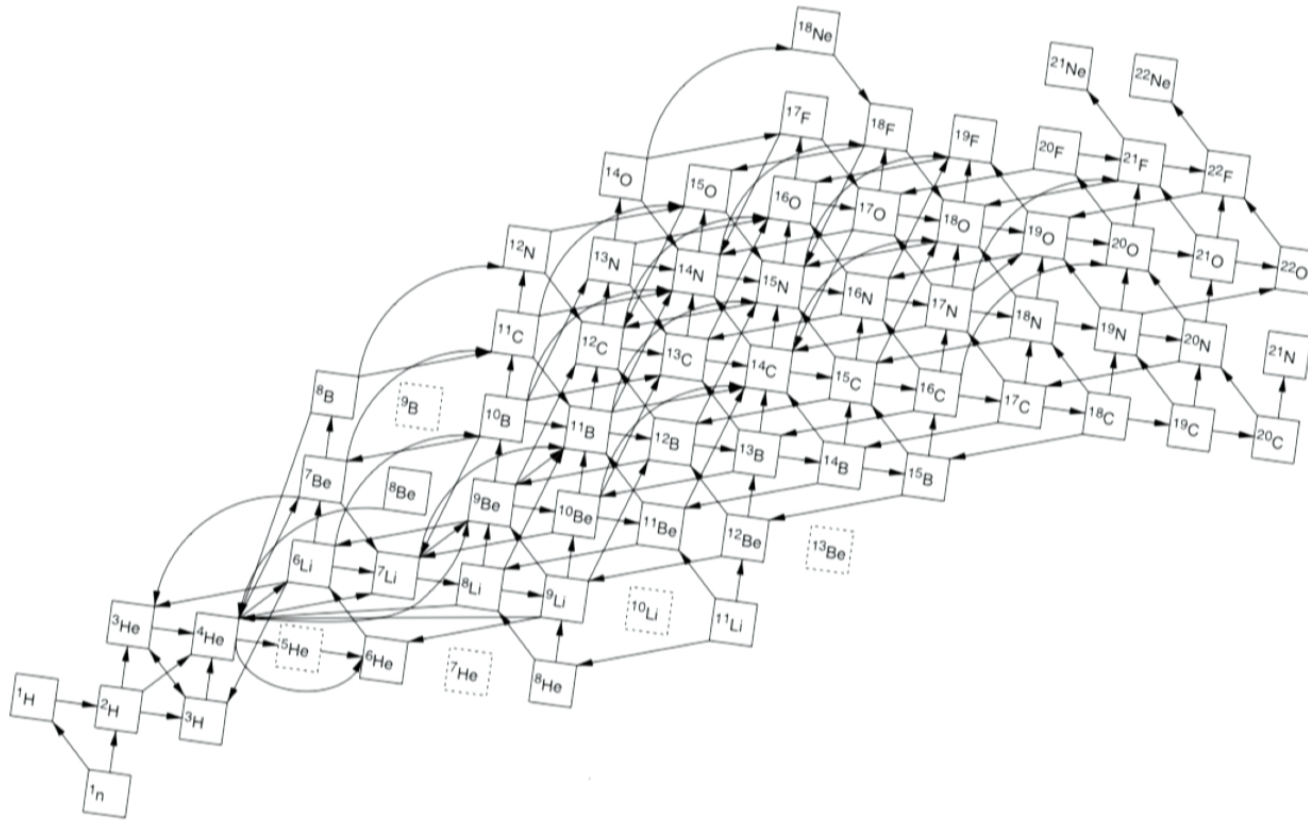
## 2. pp chain in stars



## 3. supernovae



# BBN

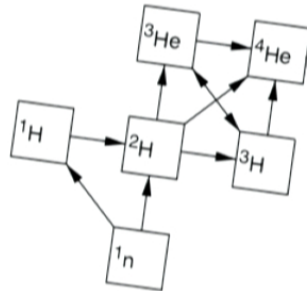


# BBN and He4



decouples:  $T_{\text{fr}} \sim \frac{v^{4/3}}{M_p^{1/3}} \approx 1 \text{ MeV}$

$$\frac{n}{p} = e^{-(m_N - m_P)/T_{\text{fr}}} \quad m_N - m_P \approx 1.3 \text{ MeV}$$



$$Y_4 \approx \frac{2(n/p)}{1 + n/p} \approx 0.25$$

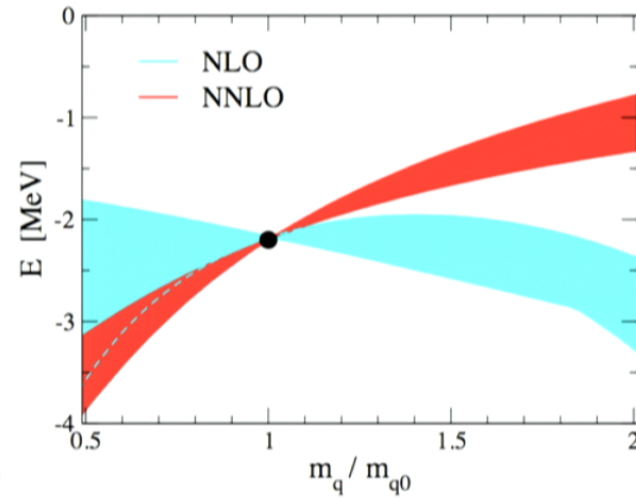
# BBN and He4

- also  $m_n - m_p = \delta_{\text{EM}} + \delta_q(m_d - m_u)$

- $\tau_n \sim \frac{v^4}{\delta m^5}$

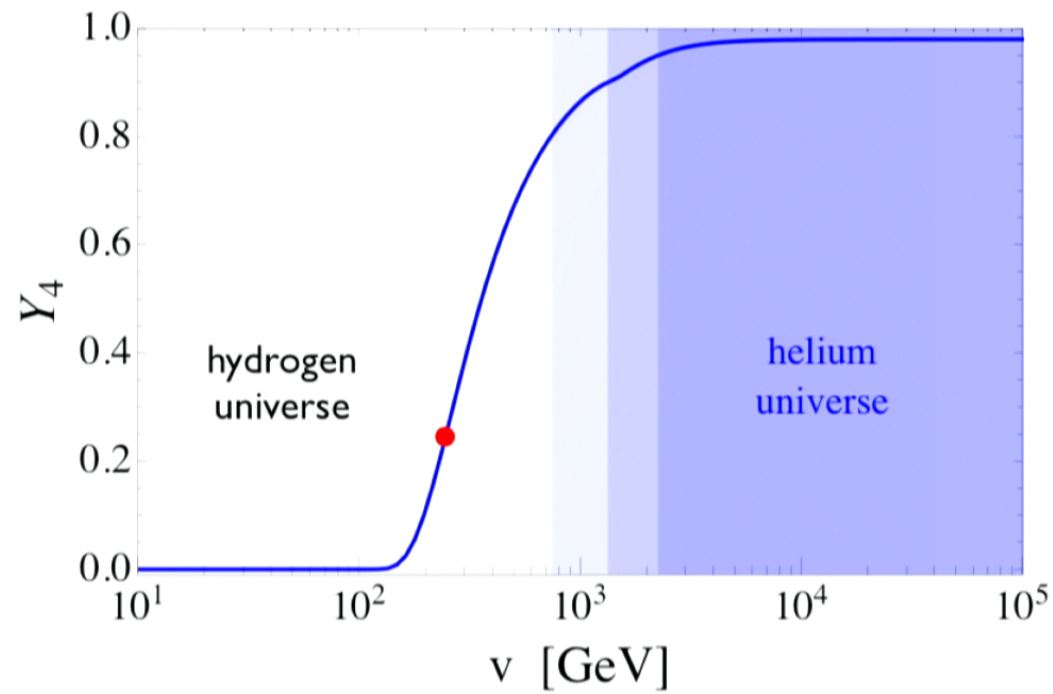
- $B_d, \sigma(p + n \leftrightarrow d + \gamma)$

Berengut, et al, 1301.1738





# BBN and He4



$$(m_N - m_P)^3 M_p \sim v^4$$

# dangers of a helium universe

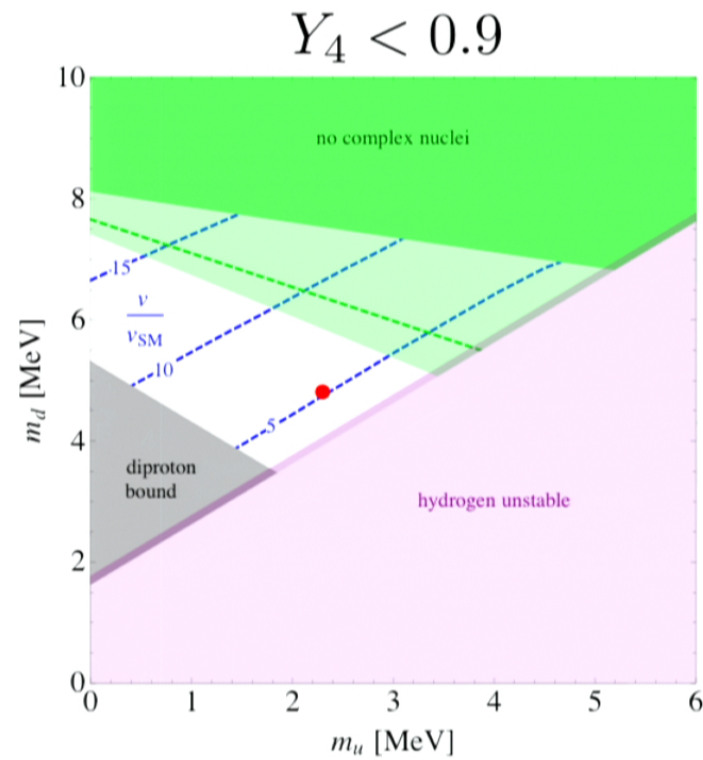
primordial hydrogen is important for:

- galactic halo cooling
- stars powered by pp chain
- water

quantifying how much hydrogen is needed for  
observers we leave for future work...

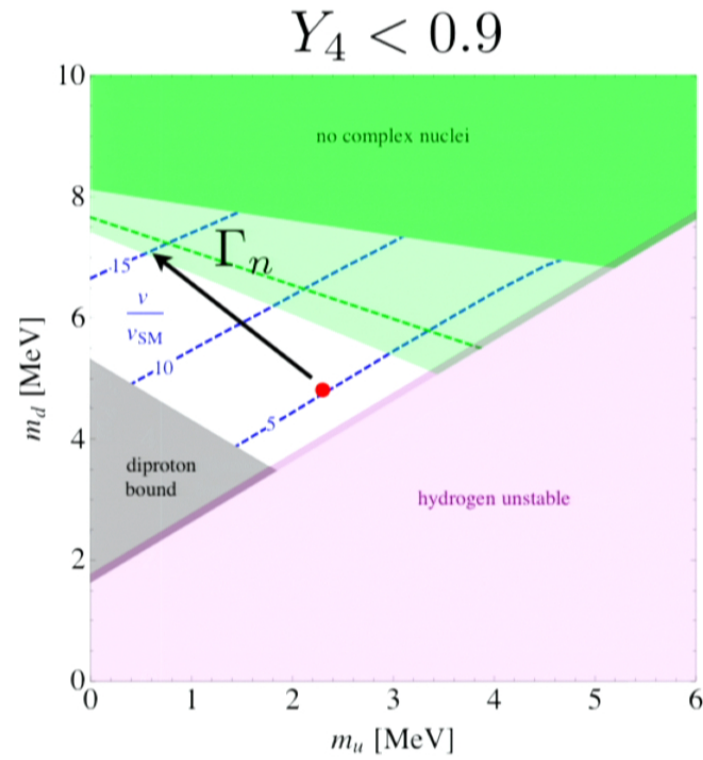
# BBN and the weak scale

scan:  
 $(y_u, y_d, v)$



# BBN and the weak scale

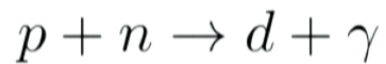
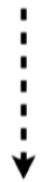
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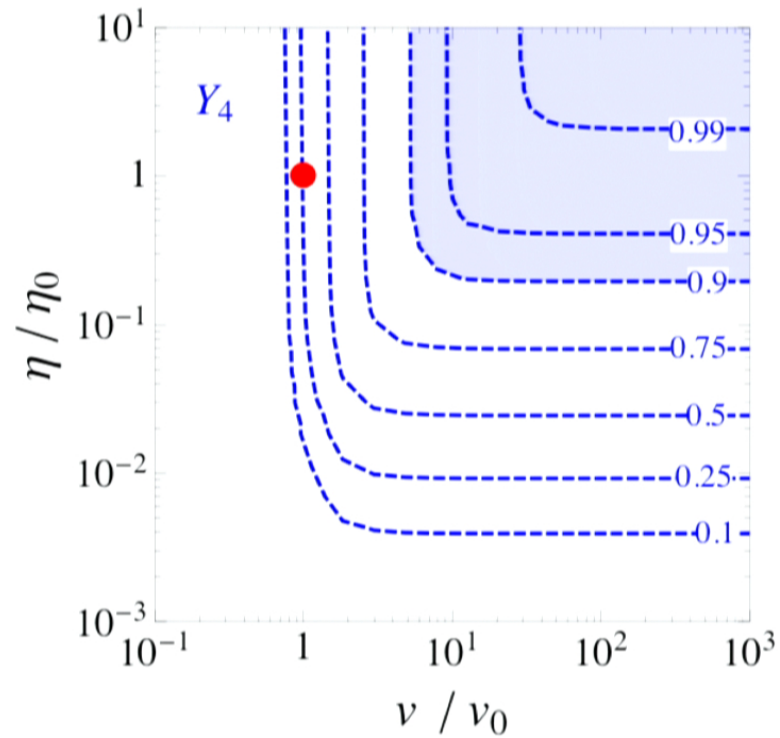
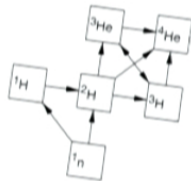
- in a multiverse where  $(y_u, y_d, v)$  scan, all three parameters are bounded by requiring stable Hydrogen, complex nuclei, and not too much Helium from BBN
- but what if other parameters scan too?

# varying the baryon density

scan:  
( $v, \eta$ )



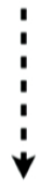
freezes out



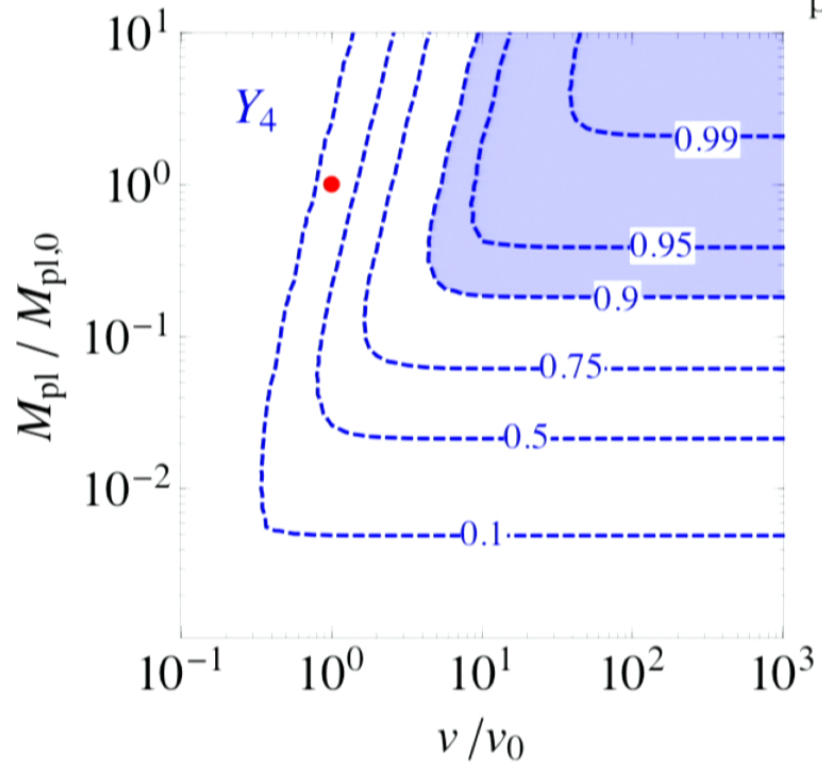
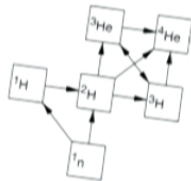
# varying $M_{\text{pl}}$

$$T_{\text{fr}} \sim \frac{v^{4/3}}{M_{\text{pl}}^{1/3}}$$

scan:  
( $v, M_{\text{pl}}$ )



$p + n \rightarrow d + \gamma$   
freezes out



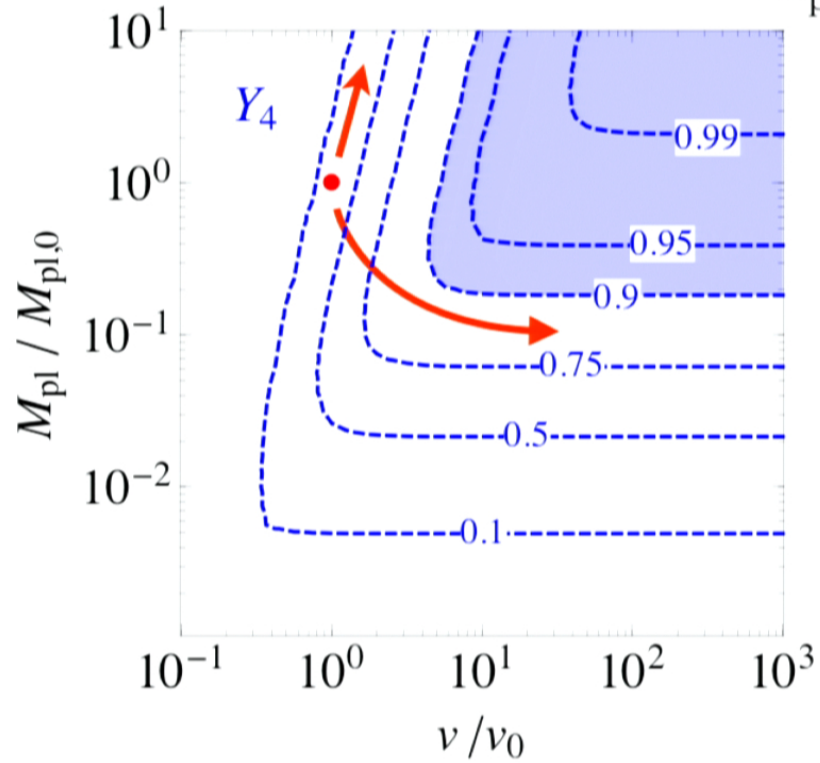
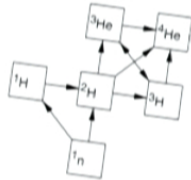
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# Scanning cosmology

- $M_{\text{pl}}, \rho_b, \rho_c$
- structure formation boundaries
- WIMP DM  $\implies \rho_c \propto v^2$

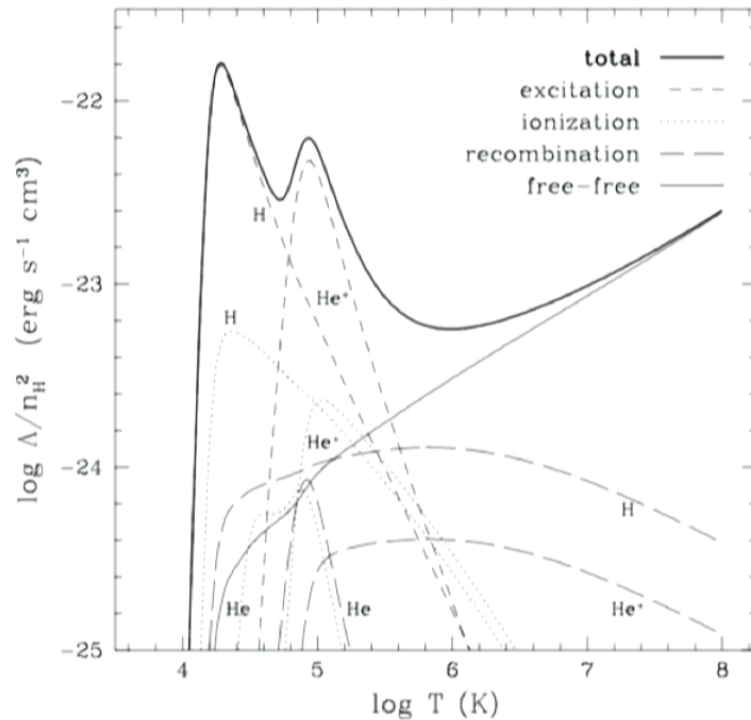
Tegmark, Aguirre, Rees, Wilczek 0511774

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Tegmark, Aguirre, Rees, Wilczek 0511774

# cooling



Katz, Weinberg, Hernquist, 9509107

$$T_{\text{vir}} \sim \frac{M m_p}{M_{\text{pl}}^2 R_{\text{vir}}}$$

$$\sim \frac{m_p}{M_{\text{pl}}^2} (\rho_{\text{vir}} M^2)^{1/3}$$

$$t_{\text{brem}} \sim \frac{nT}{\Lambda} \propto \frac{\sqrt{T_{\text{vir}}}}{f_b \rho_{\text{vir}}}$$

Inefficient cooling unless

$$10^6 M_{\star} \lesssim M \lesssim 10^{13} M_{\star}$$

$$(t_{\text{cool}} > H^{-1})$$

# disk fragmentation

- necessary for star formation
- Jeans instability in the z-direction
- $M_{\text{disk}} \gtrsim 120 M_{\text{pl}}^2 v_p v_c R_{\text{disk}}$

Fall & Efstathiou (1980)

$$\Rightarrow f_b \gtrsim 0.1 \sqrt{\frac{T_{\text{min}}}{T_{\text{vir}}}}$$

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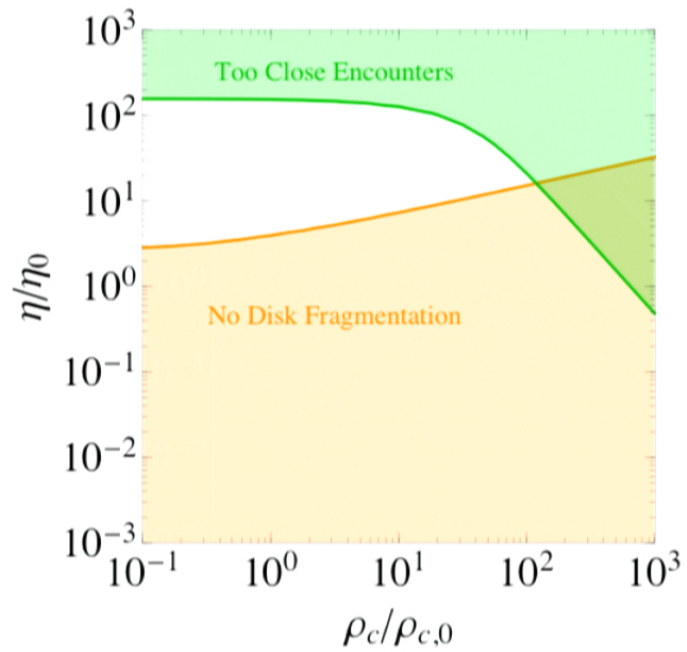
# close encounters

- stellar collision rate
- $\gamma \equiv n_{\star} \sigma v_p \lesssim (10^9 \text{ years})^{-1}$
- $n_{\star} \sim \rho_{\text{disk}}/M_{\star}, \sigma \sim (1 \text{ AU})^2$
- “Goldilocks” relates cross-section to fundamental parameters

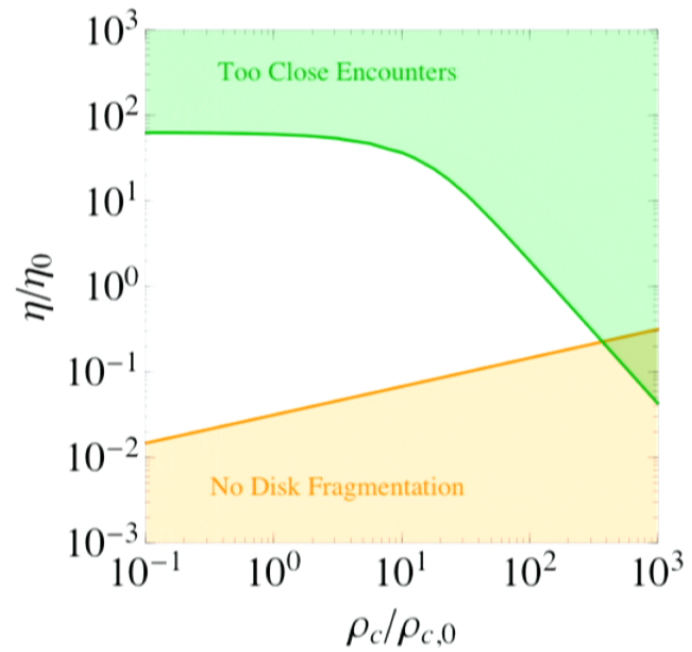
$$\gamma \propto f_b^2 \frac{\rho_{\text{vir}}^{4/3} M^{2/3}}{M_{\text{pl}}^4 m_e}$$

# varying baryon asymmetry

$M = 10^6 M_{\text{sun}}$

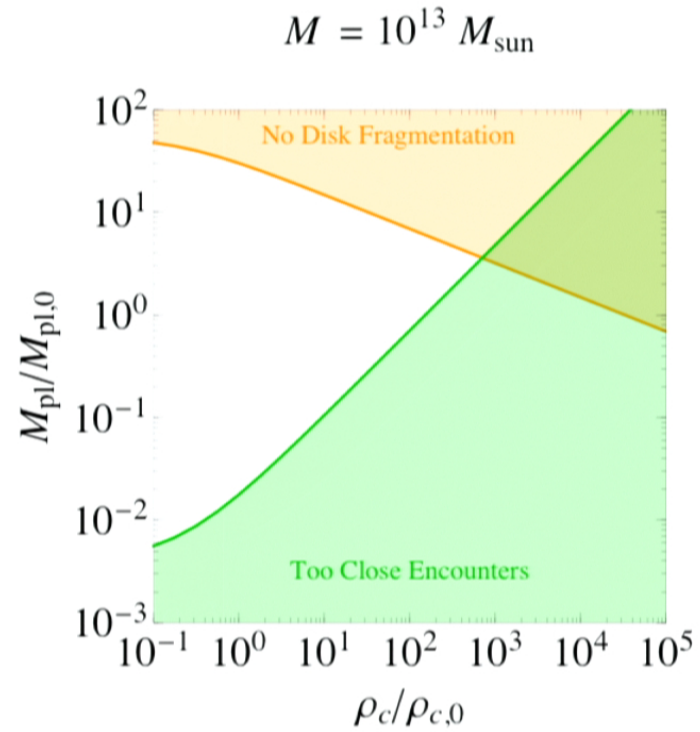
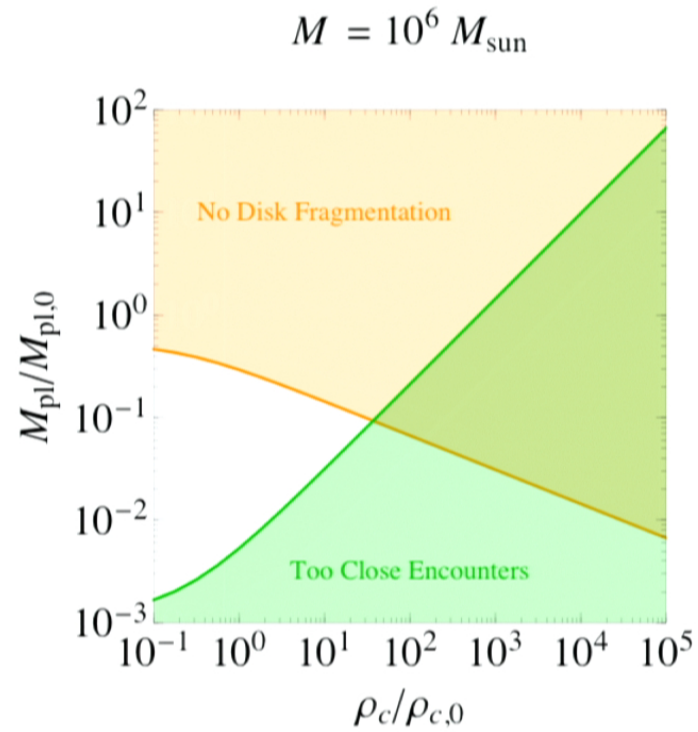


$M = 10^{13} M_{\text{sun}}$

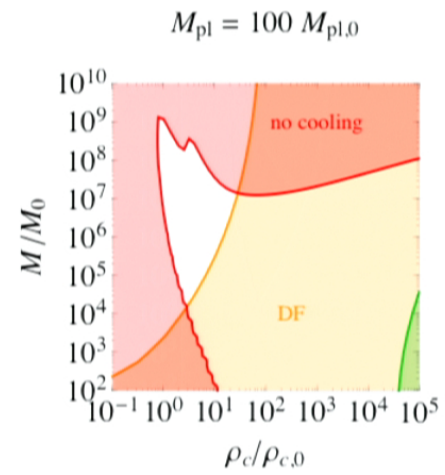
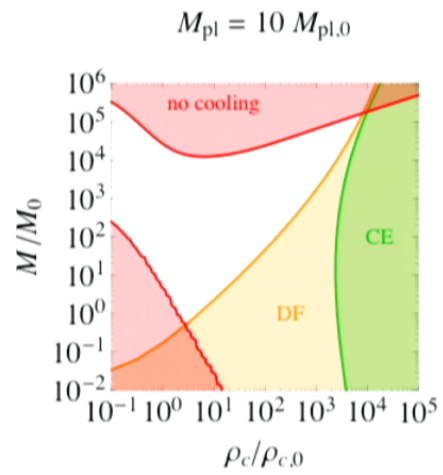
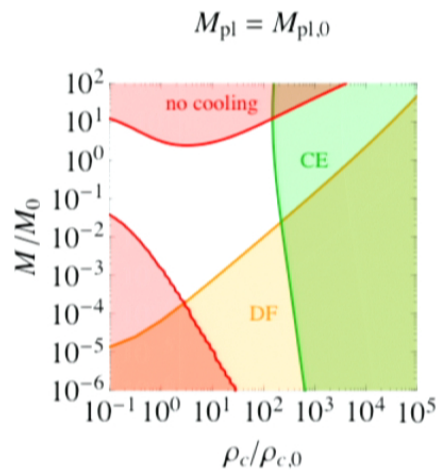




# varying planck mass



# add in cooling



# conclusions

- atomic principle bounds quark masses, not  $v$
- EW fine tuning may be explained by He, unless baryon asymmetry, Mpl scan
- close stellar encounters, disk fragmentation, cooling bound virial density
  - restores Weinberg's solution to CC problem
  - explains EW fine tuning for WIMP DM
- more to do:  $m_e$ ,  $\delta\rho/\rho$ , Silk damping, etc.