

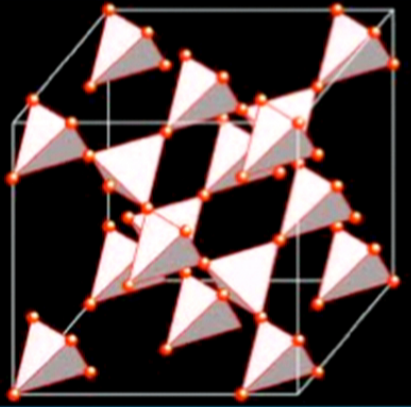
Title: Topological Phases in Transition Metal Oxides

Date: Nov 07, 2013 02:00 PM

URL: <http://pirsa.org/13110072>

Abstract: Certain varieties of transition metal oxides possess both significant interactions and strong spin-orbit coupling. In this talk I will describe materials-motivated models that predict topological phases in heterostructured and bulk transition metal oxides. We find Z_2 topological insulators, Chern insulators, topological crystalline insulators, and interaction-driven topological phases not adiabatically connected to non-interacting topological phases.

Topological Phases in Pyrochlore Oxides: TCI, TCMI, and TI* (and TCI*)



Mehdi Kargarian



Victor Chua



Joseph Maciejko

Gregory A. Fiete

University of Texas at Austin

M. Kargarian and GAF, *Phys. Rev. Lett.* (2013)

J. Maciejko, V. Chua, and GAF arXiv:1307.5566



Outline

- Brief review of predicted topological phases in $A_2M_2O_7$ (mainly M =iridium, A =rare earth)
- Minimal models and band structure considerations
- Topological crystalline insulator (TCI)
 - Interactions can drive a “TCMI”, $U(1)$ gauge theory
- The fractionalized TI^* phase, Z_2 gauge theory
 - Low-energy described by $(3+1)D$ BF theory
- Topological phases in oxide thin films:
 - Pyrochlore Iridates and Perovskite Nickelates

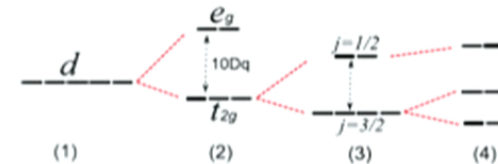
Predictions for topological phases in pyrochlore iridates (partial)

- $A_2Ir_2O_7$ (undistorted): D. Pesin, L. Balents, Nat. Phys. (2010)
“Topological Mott Insulator” (TMI)
- $A_2Ir_2O_7$ (distorted): B.-J. Yang, Y.-B. Kim, PRB (2010)
Single-particle type TI in non-interacting model
- $A_2Ir_2O_7$ (distorted+interaction): M. Kargarian, J. Wen, GAF PRB (2011)
“Weak Topological Mott Insulator” (WTMI)
- $Y_2Ir_2O_7$ (magnetically ordered): Wan *et al.* PRB (2011)
Weyl semi-metal
- $A_2Ir_2O_7$: Go, Witczak-Krempa, Jeon, Park, Y.-B. Kim PRL (2012)
Axion insulator (found via cellular DMFT)
- $A_2Ir_2O_7$: M. Kargarian, GAF PRL (2013)
Topological crystalline insulator, TCMI
- $A_2Ir_2O_7$: J. Maciejko, V. Chua, GAF arXiv:1307.5566
TI*, SM*

Minimal model for $A_2Ir_2O_7$?

- Key energy scales:

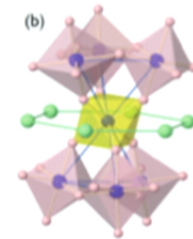
$t \sim 1.5$ eV, $E_{\text{SOC}} \sim 0.5$ eV, $U \sim 2$ eV, $E_{t_{2g}-e_g} \sim 2$ eV,
 $\Delta \sim 0.3-0.5$ eV, $J_H \sim 0.5$ eV



Moreover, extended 5d orbitals may feel significant crystal fields beyond the local oxygen environment, e.g., Hozio *et al.*, arXiv:1212.4009

Elements of minimal not entirely clear at present.

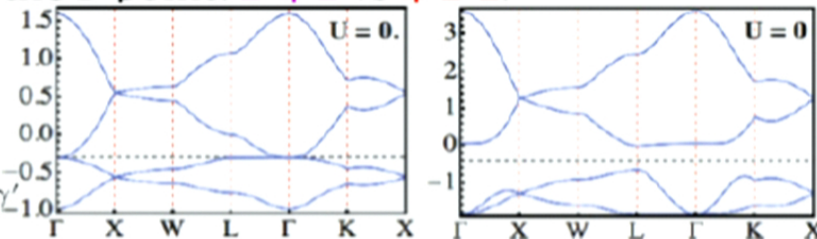
W. Witczak-Krempa, G. Chen, Y.-B. Kim, L. Balents arxiv:1305.2193



- A central issue for topological physics is the $j=1/2$ band structure, namely **order of degeneracies** at the Γ point: **2-4-2** vs **4-2-2**.

$$H_0 = \sum_{\langle ij \rangle} (T_{o,ij}^{\gamma\gamma'} + T_{d,ij}^{\gamma\gamma'}) d_{i\gamma}^\dagger d_{j\gamma'}$$

$$= \sum_{\langle ij \rangle} (t_1 \delta_{\gamma\gamma'} + it_2 \mathbf{d}_{ij} \cdot \vec{\sigma}_{\gamma\gamma'}) d_{i\gamma}^\dagger d_{j\gamma'}$$

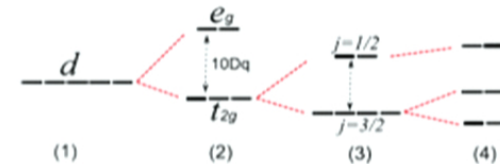


W. Witczak-Krempa and Y.-B. Kim, PRB (2012)

Minimal model for $A_2Ir_2O_7$?

- Key energy scales:

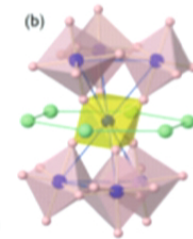
$t \sim 1.5$ eV, $E_{\text{SOC}} \sim 0.5$ eV, $U \sim 2$ eV, $E_{t_{2g}-e_g} \sim 2$ eV,
 $\Delta \sim 0.3-0.5$ eV, $J_H \sim 0.5$ eV



Moreover, extended 5d orbitals may feel significant crystal fields beyond the local oxygen environment, e.g., Hozio *et al.*, arXiv:1212.4009

Elements of minimal not entirely clear at present.

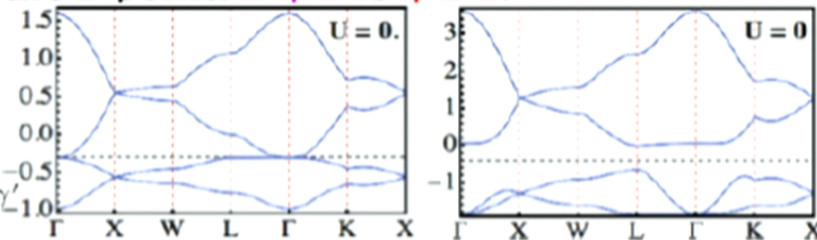
W. Witczak-Krempa, G. Chen, Y.-B. Kim, L. Balents arxiv:1305.2193



- A central issue for topological physics is the $j=1/2$ band structure, namely **order of degeneracies** at the Γ point: **2-4-2** vs **4-2-2**.

$$H_0 = \sum_{\langle ij \rangle} (T_{o,ij}^{\gamma\gamma'} + T_{d,ij}^{\gamma\gamma'}) d_{i\gamma}^\dagger d_{j\gamma'}$$

$$= \sum_{\langle ij \rangle} (t_1 \delta_{\gamma\gamma'} + it_2 \mathbf{d}_{ij} \cdot \vec{\sigma}_{\gamma\gamma'}) d_{i\gamma}^\dagger d_{j\gamma'}$$

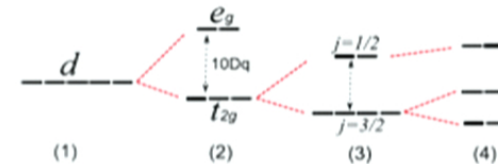


W. Witczak-Krempa and Y.-B. Kim, PRB (2012)

Minimal model for $A_2Ir_2O_7$?

- Key energy scales:

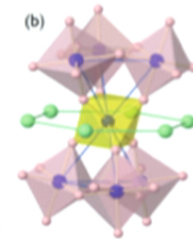
$t \sim 1.5$ eV, $E_{\text{SOC}} \sim 0.5$ eV, $U \sim 2$ eV, $E_{t_{2g}-e_g} \sim 2$ eV,
 $\Delta \sim 0.3-0.5$ eV, $J_H \sim 0.5$ eV



Moreover, extended 5d orbitals may feel significant crystal fields beyond the local oxygen environment, e.g., Hozio *et al.*, arXiv:1212.4009

Elements of minimal not entirely clear at present.

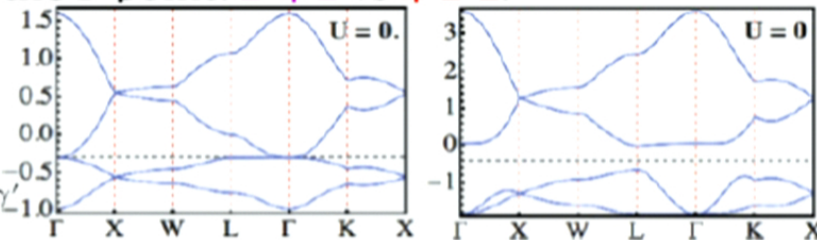
W. Witczak-Krempa, G. Chen, Y.-B. Kim, L. Balents arxiv:1305.2193



- A central issue for topological physics is the $j=1/2$ band structure, namely **order of degeneracies** at the Γ point: **2-4-2** vs **4-2-2**.

$$H_0 = \sum_{\langle ij \rangle} (T_{o,ij}^{\gamma\gamma'} + T_{d,ij}^{\gamma\gamma'}) d_{i\gamma}^\dagger d_{j\gamma'}$$

$$= \sum_{\langle ij \rangle} (t_1 \delta_{\gamma\gamma'} + it_2 \mathbf{d}_{ij} \cdot \vec{\sigma}_{\gamma\gamma'}) d_{i\gamma}^\dagger d_{j\gamma'}$$

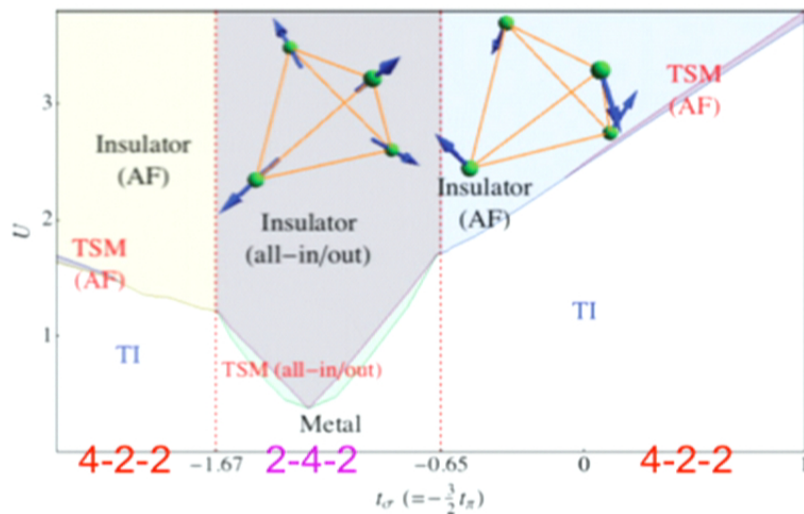


W. Witczak-Krempa and Y.-B. Kim, PRB (2012)

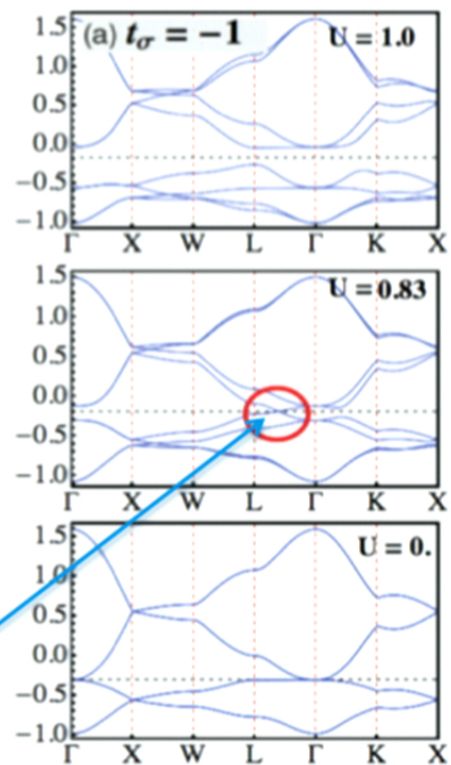
Topological phases in pyrochlore Iridate systems: $A_2Ir_2O_7$

Hartree-Fock Calculations for $j=1/2$

Witczak-Krempa and Kim, PRB (2012)



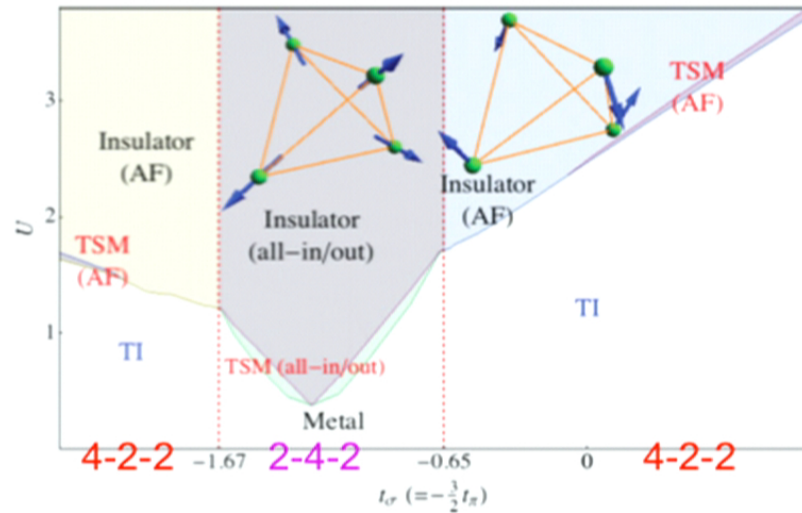
$$H = H_{d-d} + H_{d-o-d} + H_U \quad \text{Weyl point}$$



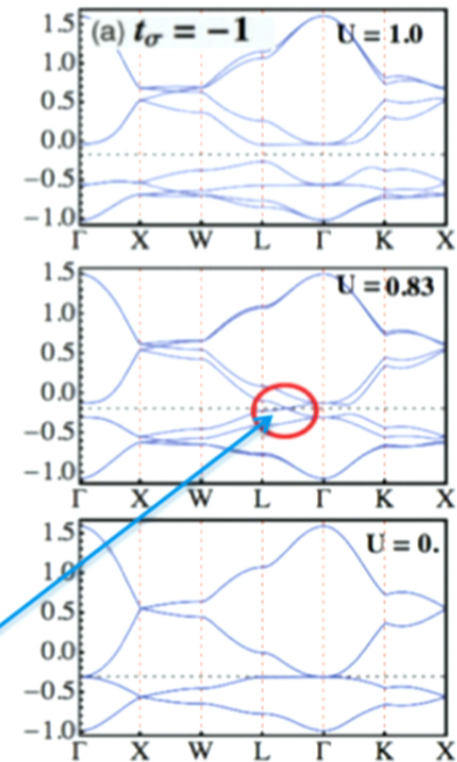
Topological phases in pyrochlore Iridate systems: $A_2Ir_2O_7$

Hartree-Fock Calculations for $j=1/2$

Witczak-Krempa and Kim, PRB (2012)



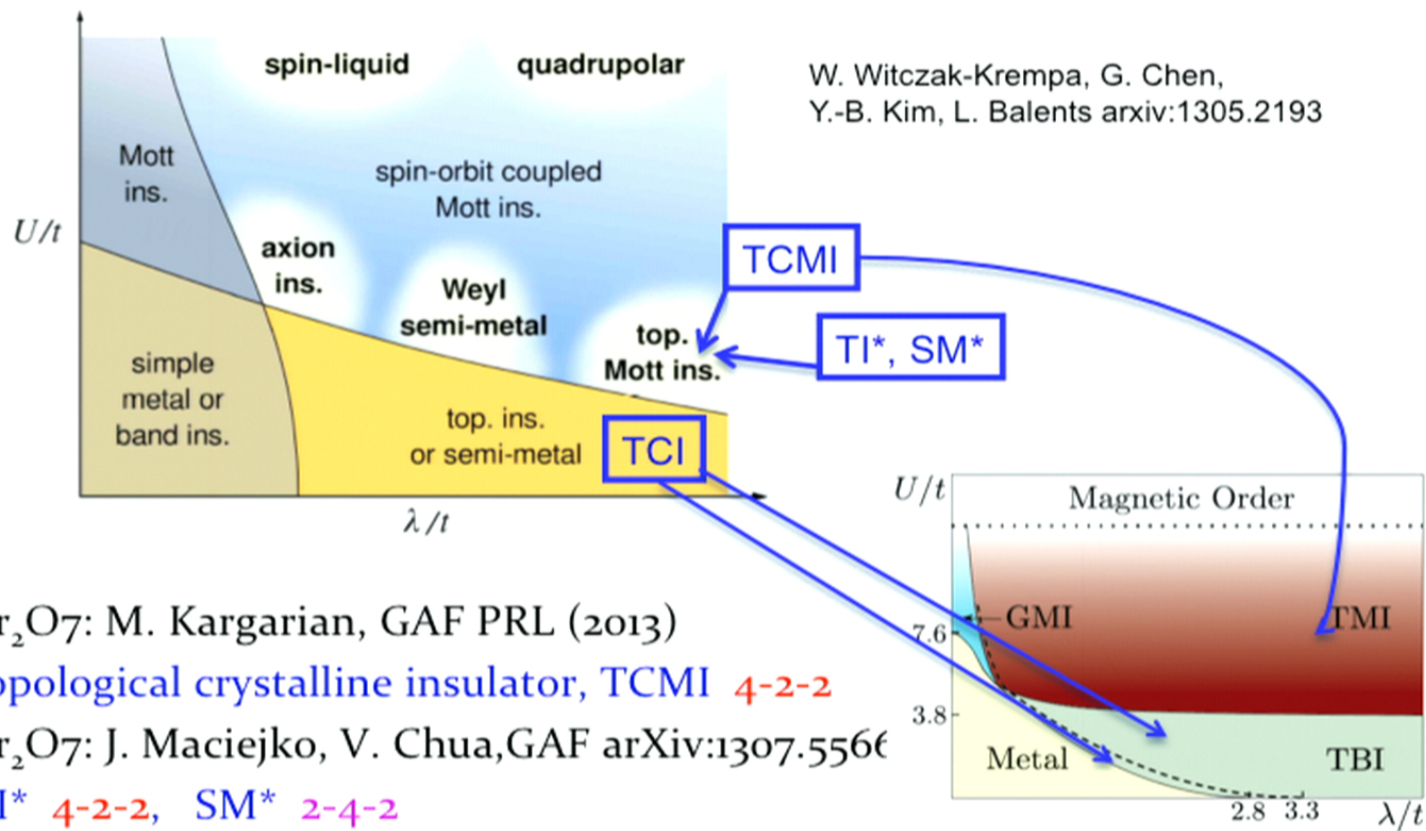
$$H = H_{d-d} + H_{d-o-d} + H_U \quad \text{Weyl point}$$



Predictions for topological phases in pyrochlore iridates with Γ -point degeneracies given

- $A_2Ir_2O_7$ (undistorted): D. Pesin, L. Balents, Nat. Phys. (2010)
“Topological Mott Insulator” (TMI) 4-2-2
- $A_2Ir_2O_7$ (distorted): B.-J. Yang, Y.-B. Kim, PRB (2010)
Single-particle type TI in non-interacting model 4-2-2
- $A_2Ir_2O_7$ (distorted+interaction): M. Kargarian, J. Wen, GAF PRB (2011)
“Weak Topological Mott Insulator” (WTMI) 4-2-2
- $Y_2Ir_2O_7$ (magnetically ordered): Wan *et al.* PRB (2011)
Weyl semi-metal 2-4-2 4-2-2
- $A_2Ir_2O_7$: Go, Witczak-Krempa, Jeon, Park, Y.-B. Kim PRL (2012)
Axion insulator (found via cellular DMFT) 4-2-2
- $A_2Ir_2O_7$: M. Kargarian, GAF PRL (2013)
Topological crystalline insulator, TCMi 4-2-2
- $A_2Ir_2O_7$: J. Maciejko, V. Chua, GAF arXiv:1307.5566
TI* 4-2-2, SM* 2-4-2

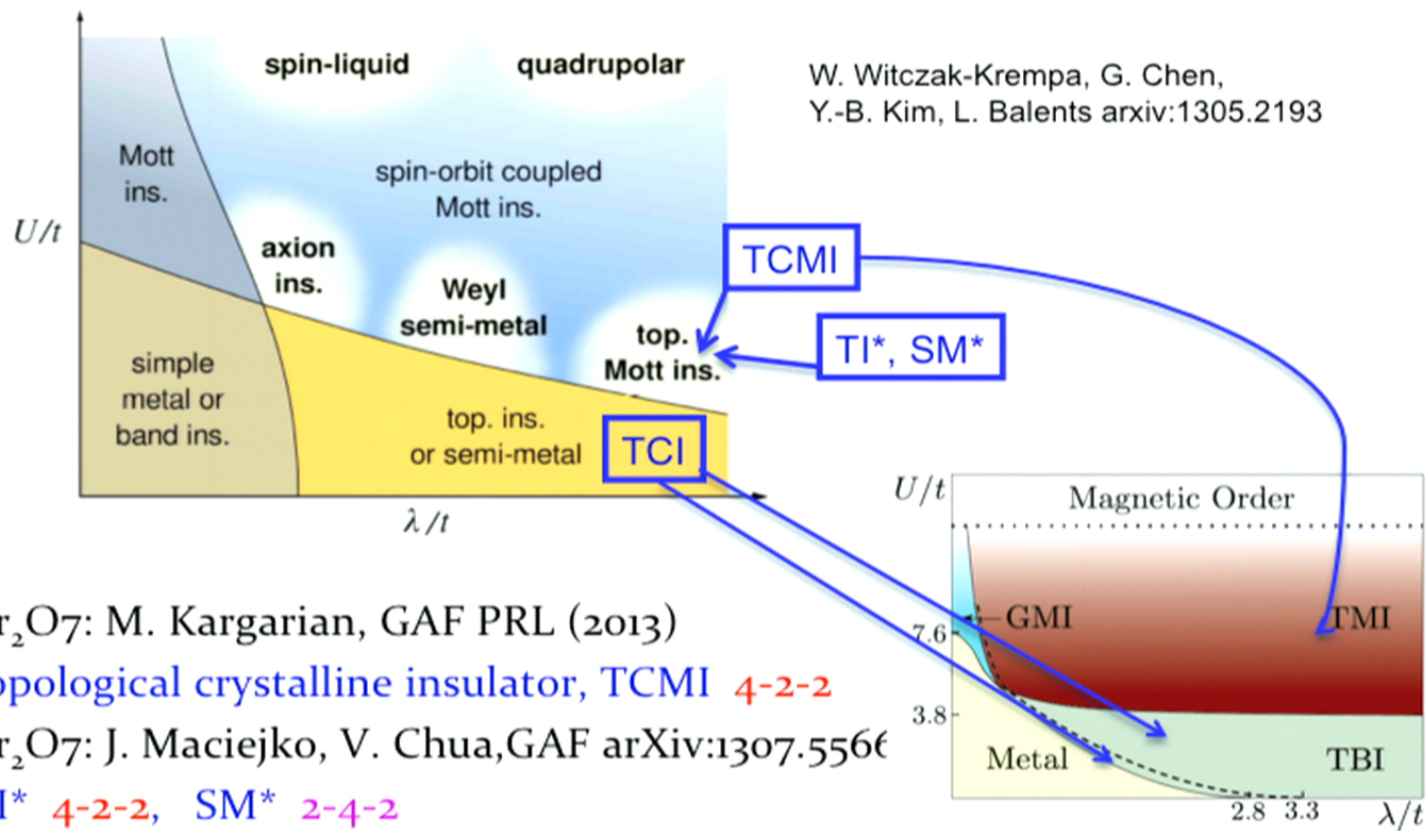
Schematic Phase Diagram



- $A_2Ir_2O_7$: M. Kargarian, GAF PRL (2013)
 Topological crystalline insulator, TCMi 4-2-2
- $A_2Ir_2O_7$: J. Maciejko, V. Chua, GAF arXiv:1307.5566
 TI* 4-2-2, SM* 2-4-2

Pesin, Balents Nat. Phys. (2010)

Schematic Phase Diagram

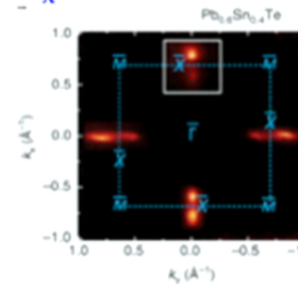
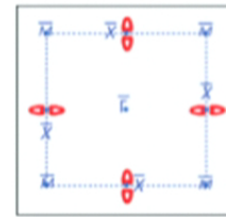
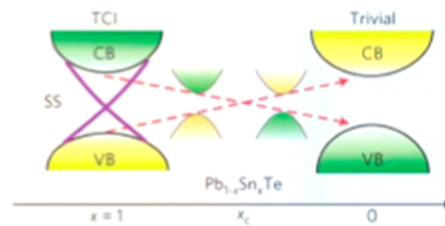
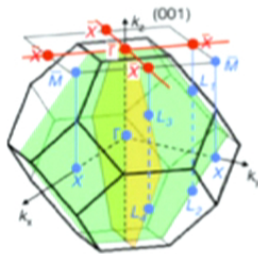


- $A_2Ir_2O_7$: M. Kargarian, GAF PRL (2013)
Topological crystalline insulator, TCI $4-2-2$
- $A_2Ir_2O_7$: J. Maciejko, V. Chua, GAF arXiv:1307.5566
 $TI^* 4-2-2$, $SM^* 2-4-2$

Pesin, Balents Nat. Phys. (2010)

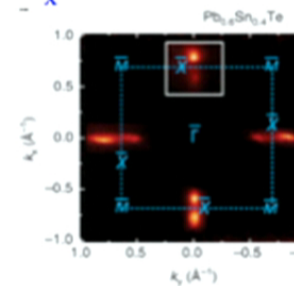
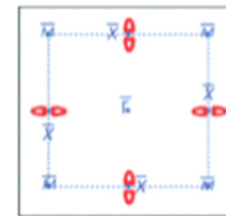
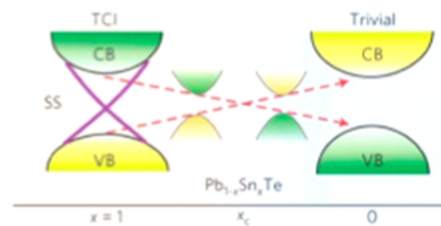
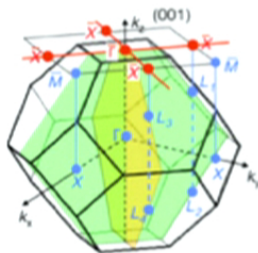
Topological Crystalline Insulators: Mirror Chern Number

- Predicted by Liang Fu and collaborators:
 - T. H. Hsieh, H. Lin, J. Liu, W. Duan, A. Bansil, L. Fu, Nat Commun, 3, 982 (2012). [SnTe Material class](#)
 - Also recent, J. Liu, W. Duan, L. Fu arXiv:1304.0430.
- Realized essentially simultaneously by three groups:
 - S.-Y. Xu, ... M. Z. Hasan, Nat Commun, 3, 1192 (2012) [Pb_{1-x}Sn_xTe](#)
 - Y. Tanaka, ... Y. Ando, Nat Phys, 8, 800 (2012). [SnTe](#)
 - P. Dziawa, ... T. Story, Nat Mater, 11, 1023 (2012). [Pb_{1-x}Sn_xSe](#)



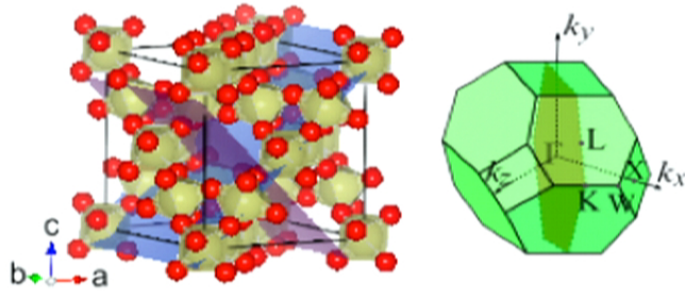
Topological Crystalline Insulators: Mirror Chern Number

- Predicted by Liang Fu and collaborators:
 - T. H. Hsieh, H. Lin, J. Liu, W. Duan, A. Bansil, L. Fu, Nat Commun, 3, 982 (2012). [SnTe Material class](#)
 - Also recent, J. Liu, W. Duan, L. Fu arXiv:1304.0430.
- Realized essentially simultaneously by three groups:
 - S.-Y. Xu, ... M. Z. Hasan, Nat Commun, 3, 1192 (2012) [Pb_{1-x}Sn_xTe](#)
 - Y. Tanaka, ... Y. Ando, Nat Phys, 8, 800 (2012). [SnTe](#)
 - P. Dziawa, ... T. Story, Nat Mater, 11, 1023 (2012). [Pb_{1-x}Sn_xSe](#)



Computing TCI invariant for $A_2Ir_2O_7$

M. Kargarian and GAF PRL (2013)



$$H_0 = \sum_i t_i^{\gamma\gamma'} d_{i\gamma}^\dagger d_{i\gamma'} + \sum_{\langle ij \rangle} (T_{o,ij}^{\gamma\gamma'} + T_{d,ij}^{\gamma\gamma'}) d_{i\gamma}^\dagger d_{j\gamma'}$$

$$t_i = \varepsilon_d - \lambda \mathbf{l} \cdot \mathbf{s} \quad 4-2-2$$

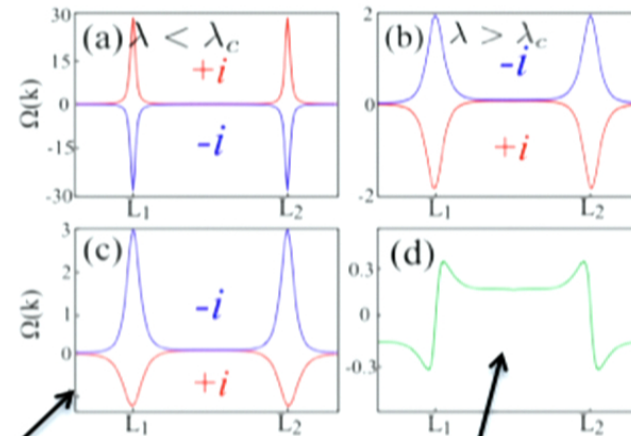
$$n_M = (n_{+i} - n_{-i})/2$$

$$n_M = -1 \text{ for } \lambda > \lambda_c$$

$$n_M = +1 \text{ for } \lambda < \lambda_c$$

$$\Omega(\mathbf{k}) = \nabla \times \mathbf{A}$$

$$\mathbf{A} = i \sum_n \langle u_n(\mathbf{k}) | \nabla | u_n(\mathbf{k}) \rangle$$

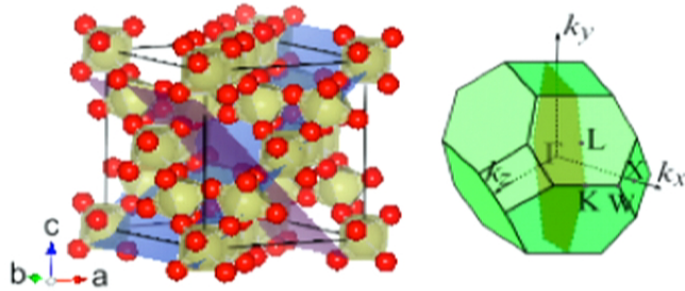


TRS Breaking, Mirror Preserved

TRS Breaking, Mirror Broken

Computing TCI invariant for $A_2Ir_2O_7$

M. Kargarian and GAF PRL (2013)



$$H_0 = \sum_i t_i^{\gamma\gamma'} d_{i\gamma}^\dagger d_{i\gamma'} + \sum_{\langle ij \rangle} (T_{o,ij}^{\gamma\gamma'} + T_{d,ij}^{\gamma\gamma'}) d_{i\gamma}^\dagger d_{j\gamma'}$$

$$t_i = \varepsilon_d - \lambda \mathbf{l} \cdot \mathbf{s} \quad 4-2-2$$

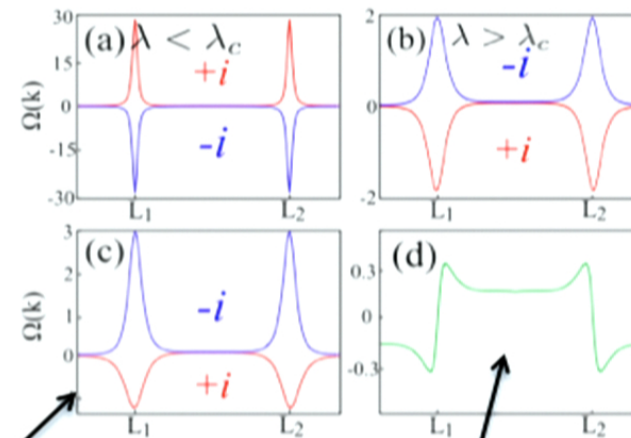
$$n_M = (n_{+i} - n_{-i})/2$$

$$n_M = -1 \text{ for } \lambda > \lambda_c$$

$$n_M = +1 \text{ for } \lambda < \lambda_c$$

$$\Omega(\mathbf{k}) = \nabla \times \mathbf{A}$$

$$\mathbf{A} = i \sum_n \langle u_n(\mathbf{k}) | \nabla | u_n(\mathbf{k}) \rangle$$

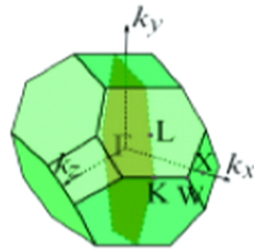
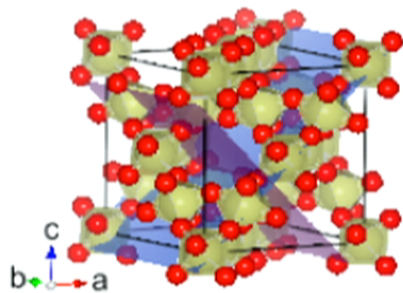


TRS Breaking, Mirror Preserved

TRS Breaking, Mirror Broken

TCI edge states for $A_2Ir_2O_7$: Slab geometry [010] Direction Mirror Planes: (101) and $(10\bar{1})$

M. Kargarian and GAF PRL (2013)

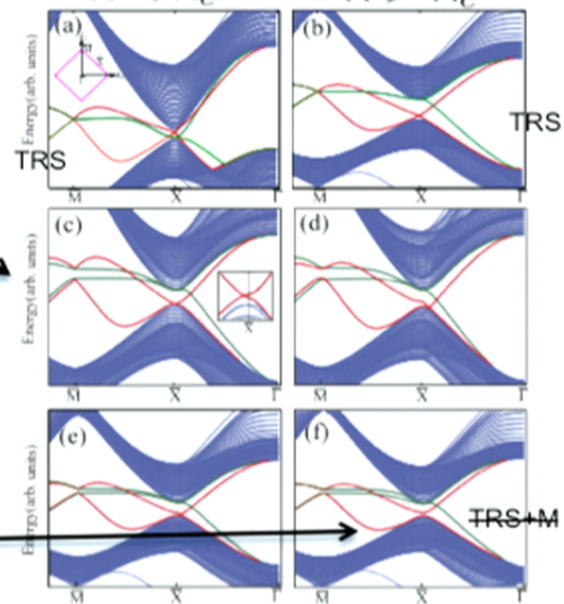


$$H_0 = \sum_i t_i^{\gamma\gamma'} d_{i\gamma}^\dagger d_{i\gamma'} + \sum_{\langle ij \rangle} (T_{o,ij}^{\gamma\gamma'} + T_{d,ij}^{\gamma\gamma'}) d_{i\gamma}^\dagger d_{j\gamma'}$$

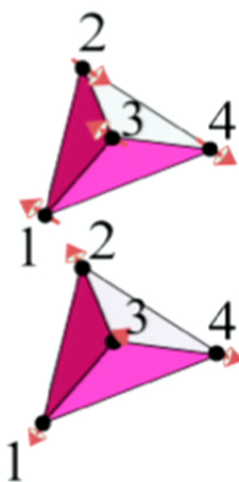
$$t_i = \varepsilon_d - \lambda l \cdot \hat{s} \quad \text{4-2-2}$$

$\lambda < \lambda_c$

$\lambda > \lambda_c$



$$H' = \sum_i \mathbf{B}_i \cdot \mathbf{S}_i$$

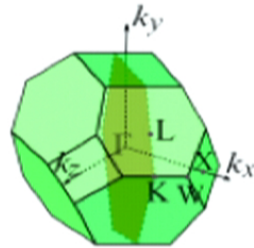
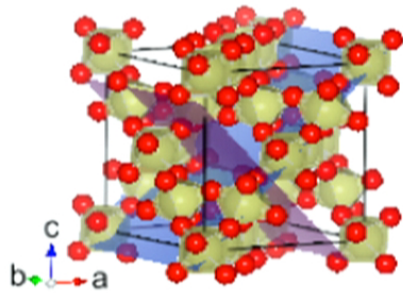


← TRS Breaking, Mirror Preserved

← TRS Breaking, Mirror Broken

TCl edge states for $A_2Ir_2O_7$: Slab geometry [010] Direction Mirror Planes: (101) and $(10\bar{1})$

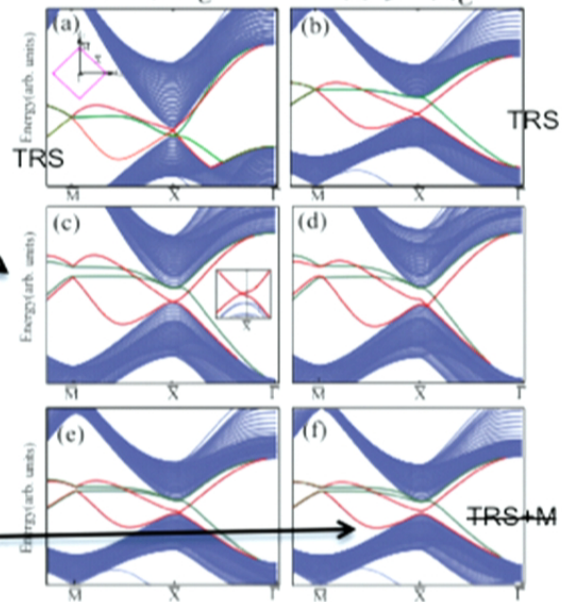
M. Kargarian and GAF PRL (2013)



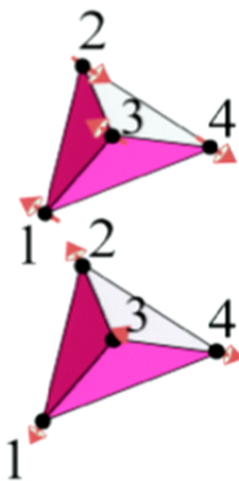
$$H_0 = \sum_i t_i^{\gamma\gamma'} d_{i\gamma}^\dagger d_{i\gamma'} + \sum_{\langle ij \rangle} (T_{o,ij}^{\gamma\gamma'} + T_{d,ij}^{\gamma\gamma'}) d_{i\gamma}^\dagger d_{j\gamma'}$$

$$t_i = \varepsilon_d - \lambda l \cdot \hat{s} \quad \text{4-2-2}$$

$\lambda < \lambda_c$ $\lambda > \lambda_c$



$$H' = \sum_i \mathbf{B}_i \cdot \mathbf{S}_i$$

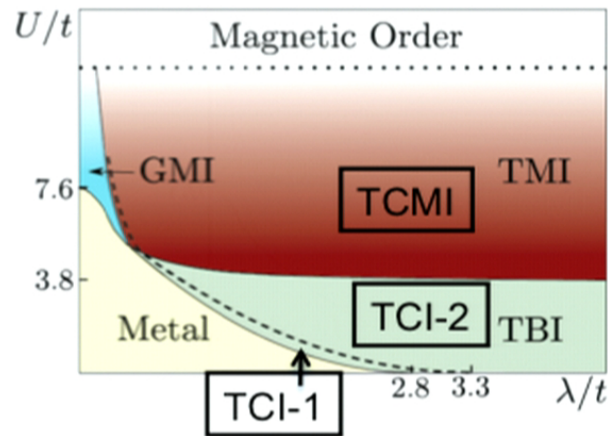


← TRS Breaking, Mirror Preserved

← TRS Breaking, Mirror Broken

“TCMI” from slave rotors: Refined Phase Diagram

M. Kargarian and GAF PRL (2013)



Pesin and Balents Nat. Phys. (2010)

$$H_0 = \sum_i t_i^{\gamma\gamma'} d_{i\gamma}^\dagger d_{i\gamma'} + \sum_{\langle ij \rangle} (T_{o,ij}^{\gamma\gamma'} + T_{d,ij}^{\gamma\gamma'}) d_{i\gamma}^\dagger d_{j\gamma'}$$

$$t_i = \varepsilon_d - \lambda l \cdot \hat{s} \quad 4-2-2$$

$$H_U = U \sum_i \left(\sum_\gamma n_{i\gamma} - n_d \right)^2$$

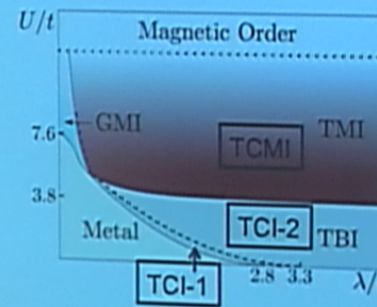
$$d_{j\gamma} = e^{i\theta_j} f_{j\gamma}$$

charge ↗ ↖ spin

Find that TMI is also TCMI—a spin liquid with topological band structure protected by both time-reversal and mirror symmetries. TBI is two “flavors” of TCI.

"TCMI" from slave rotors: Refined Phase Diagram

M. Kargarian and GAF PRL (2013)



$$H_0 = \sum_i t_i^{\gamma\gamma'} d_{i\gamma}^\dagger d_{i\gamma'} + \sum_{\langle ij \rangle} (T_{o,ij}^{\gamma\gamma'} + T_{d,ij}^{\gamma\gamma'}) d_{i\gamma}^\dagger d_{j\gamma'}$$

$$t_i = \varepsilon_d - \lambda \langle \mathbf{s} \cdot \mathbf{s} \rangle \quad 4-2-2$$

$$H_U = U \sum_i (\sum_\gamma n_{i\gamma} - n_d)^2$$

$$d_{j\gamma} = e^{i\theta_j} f_{j\gamma}$$

↑ charge ↑ spin

Pesin and Balents Nat. Phys. (2010)

Find that TMI is also TCM1—a spin liquid with topological band structure protected by both time-reversal and mirror symmetries. TBI is two “flavors” of TCI.

TI* from slave spins: Topological Field Theory & Mean-field Theory

J. Maciejko, V. Chua, GAF arXiv:1307.5566

- Idea is to generalize QSH* phase to three-dimensions

A. Ruegg and GAF PRL (2012)

$$H = \sum_{rr'} \sum_{\alpha\beta} t_{\alpha\beta}^{rr'} c_{r\alpha}^\dagger c_{r'\beta} + \frac{U}{2} \sum_r \left(\sum_{\alpha=\uparrow,\downarrow} n_{r\alpha} - 1 \right)^2$$

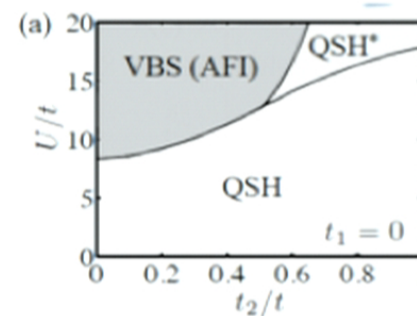
S. Huber and A. Ruegg PRL (2009)

A. Ruegg, S. Huber, M. Sigrist PRB (2010)

R. Nandkishore, M. A. Metliski, T. Senthil PRB (2012)

Zhong, Wang, Luo arXiv:1304.2099

-Exactly solvable models



$$c_{r\alpha} = f_{r\alpha} \tau_r^x$$

**J. Maciejko and A. Ruegg arXiv:1305.1290 CI* Phase

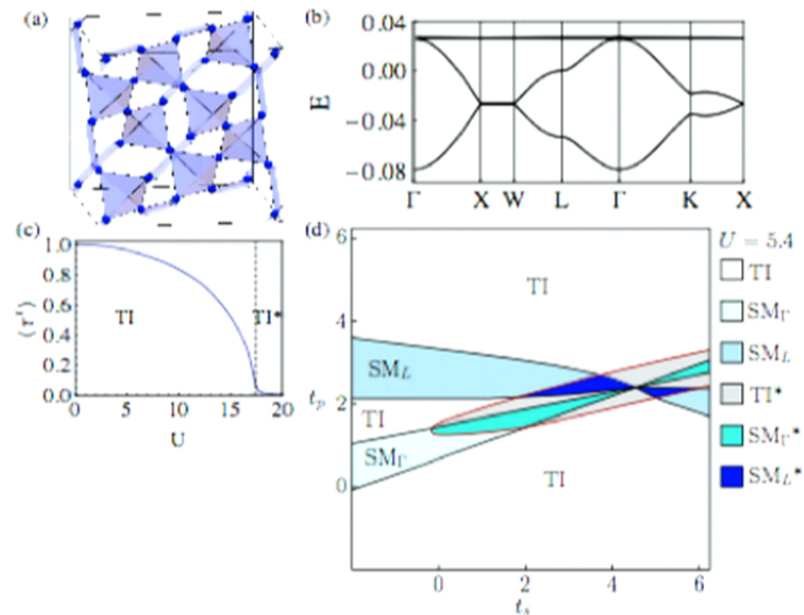
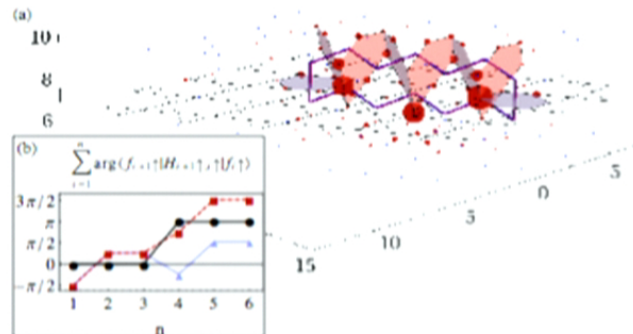
Mean-Field Theory for TI*: Phase Diagram and Braiding

J. Maciejko, V. Chua, GAF arXiv:1307.5566

$$H_{\text{MF}} = \sum_{rr'} \sum_{\alpha\beta} \chi_{\alpha\beta}^{rr'} \sigma_{rr'} f_{r\alpha}^\dagger f_{r'\beta} + \sum_{rr'} J_{rr'} \sigma_{rr'} \tau_r^x \tau_{r'}^x + \frac{U}{4} \sum_r (\tau_r^z + 1)$$

$$\chi_{\alpha\beta}^{rr'} = t_{\alpha\beta}^{rr'} \sigma_{rr'} \langle \tau_r^x \tau_{r'}^x \rangle_{\text{MF}},$$

$$J_{rr'} = \sum_{\alpha\beta} t_{\alpha\beta}^{rr'} \sigma_{rr'} \langle f_{r\alpha}^\dagger f_{r'\beta} \rangle_{\text{MF}}$$



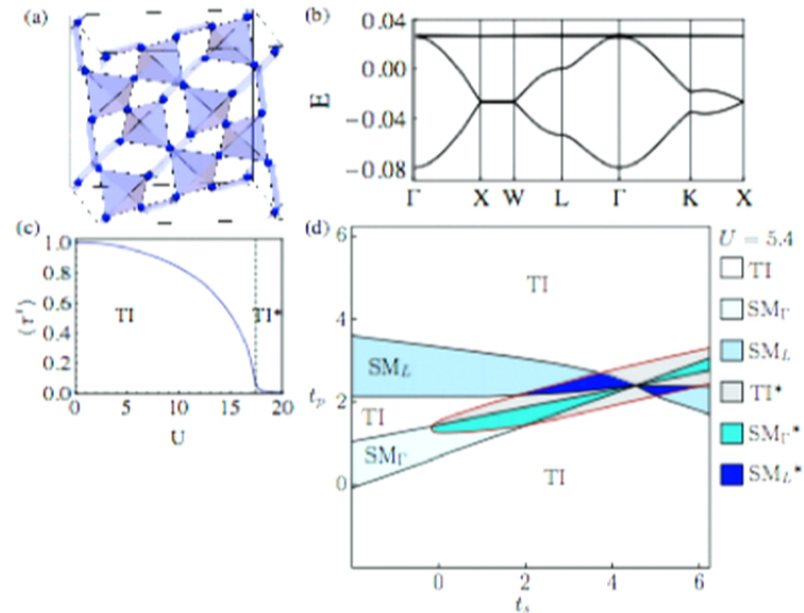
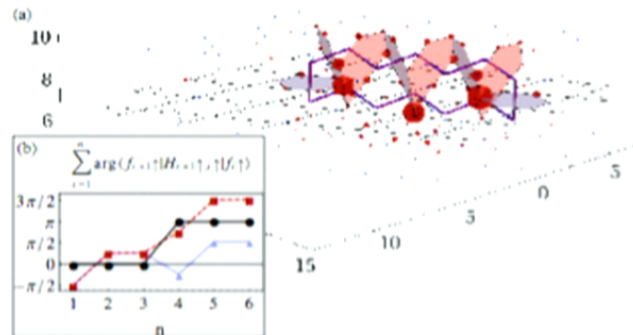
Mean-Field Theory for TI*: Phase Diagram and Braiding

J. Maciejko, V. Chua, GAF arXiv:1307.5566

$$H_{\text{MF}} = \sum_{rr'} \sum_{\alpha\beta} \chi_{\alpha\beta}^{rr'} \sigma_{rr'} f_{r\alpha}^\dagger f_{r'\beta} + \sum_{rr'} J_{rr'} \sigma_{rr'} \tau_r^x \tau_{r'}^x + \frac{U}{4} \sum_r (\tau_r^z + 1)$$

$$\chi_{\alpha\beta}^{rr'} = t_{\alpha\beta}^{rr'} \sigma_{rr'} \langle \tau_r^x \tau_{r'}^x \rangle_{\text{MF}},$$

$$J_{rr'} = \sum_{\alpha\beta} t_{\alpha\beta}^{rr'} \sigma_{rr'} \langle f_{r\alpha}^\dagger f_{r'\beta} \rangle_{\text{MF}}$$



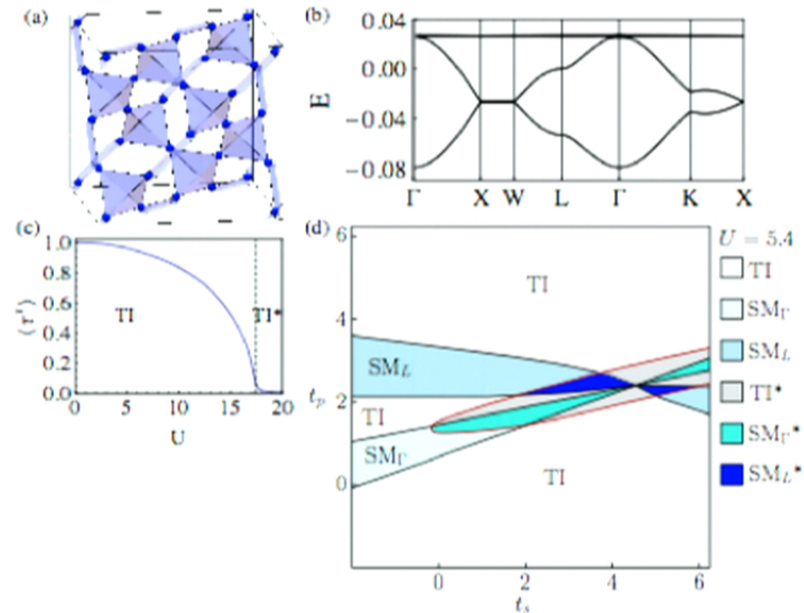
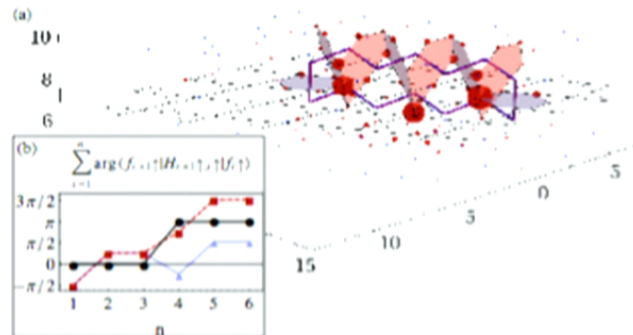
Mean-Field Theory for TI*: Phase Diagram and Braiding

J. Maciejko, V. Chua, GAF arXiv:1307.5566

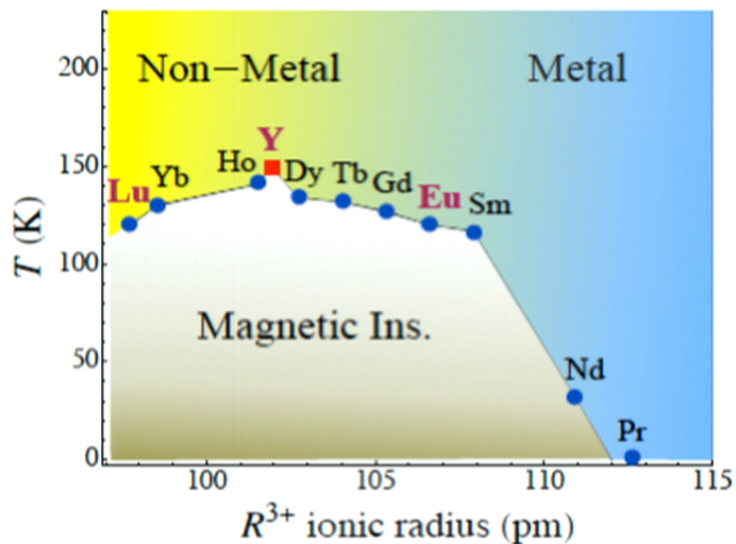
$$H_{\text{MF}} = \sum_{rr'} \sum_{\alpha\beta} \chi_{\alpha\beta}^{rr'} \sigma_{rr'} f_{r\alpha}^\dagger f_{r'\beta} + \sum_{rr'} J_{rr'} \sigma_{rr'} \tau_r^x \tau_{r'}^x + \frac{U}{4} \sum_r (\tau_r^z + 1)$$

$$\chi_{\alpha\beta}^{rr'} = t_{\alpha\beta}^{rr'} \sigma_{rr'} \langle \tau_r^x \tau_{r'}^x \rangle_{\text{MF}},$$

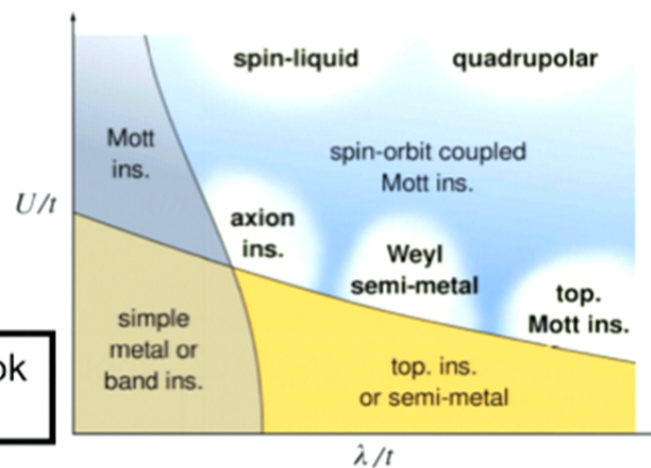
$$J_{rr'} = \sum_{\alpha\beta} t_{\alpha\beta}^{rr'} \sigma_{rr'} \langle f_{r\alpha}^\dagger f_{r'\beta} \rangle_{\text{MF}}$$



Experimental Phase Diagram for Bulk Pyrochlore Iridates

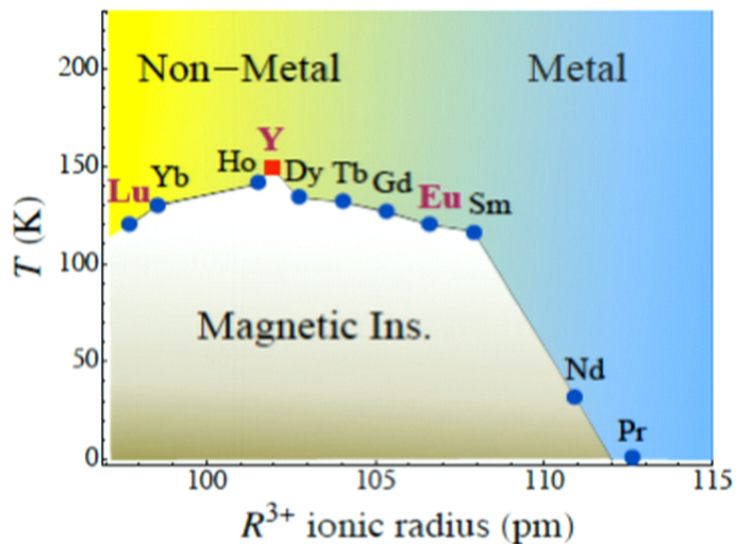


W. Witczak-Krempa, G. Chen,
Y.-B. Kim, L. Balents arxiv:1305.2193

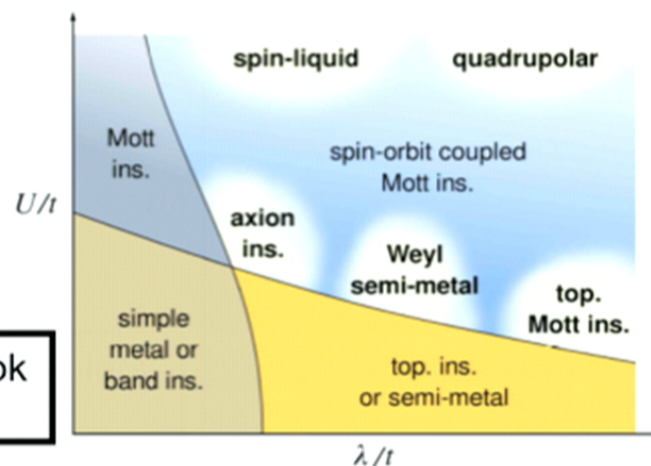


Overall, these particular materials do not look promising for topological phases.

Experimental Phase Diagram for Bulk Pyrochlore Iridates



W. Witczak-Krempa, G. Chen,
Y.-B. Kim, L. Balents arxiv:1305.2193



Overall, these particular materials do not look promising for topological phases.