

Title: Hofstadter's Butterfly and interaction driven quantum Hall ferromagnetism in graphene

Date: Nov 07, 2013 09:05 AM

URL: <http://pirsa.org/13110070>

Abstract: Electrons moving in a periodic electric potential form Bloch energy bands where the mass of electrons are effectively changed. In a strong magnetic field, the cyclotron orbits of free electrons are quantized and Landau levels forms with a massive degeneracy within. In 1976, Hofstadter showed that for 2-dimensional electronic system, the intriguing interplay between these two quantization effects can lead into a self-similar fractal set of energy spectrum known as "Hofstadter's Butterfly". Experimental efforts to demonstrate this fascinating electron energy spectrum have continued ever since. Recent advent of graphene, where its Bloch electrons can be described by Dirac fermions, provides a new opportunity to investigate this half century old problem experimentally. In this presentation, I will discuss the experimental realization Hofstadter's Butterfly via substrate engineered graphene under extremely high magnetic fields controlling two competing length scales governing Dirac-Bloch states and Landau orbits, respectively. In addition, the strong Coulomb interactions and approximate spin-pseudo spin symmetry are predicted to lead to a variety of integer quantum Hall ferromagnetic and fractional quantum Hall states and the quantum phase transition between them in graphene. I will discuss several recent experimental evidences to demonstrate the role of the electron interaction in single and bilayer graphene.

# Hofstadter's Butterfly in Graphene

and electron interaction driven quantum Hall ferromagnetism



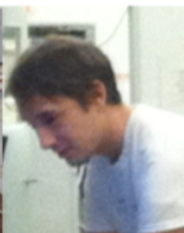
**Philip Kim**

**Physics Department, Columbia University**

# Acknowledgment



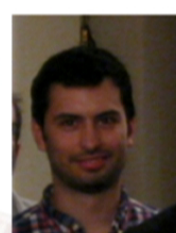
Cory Dean  
Now at CCNY



Patrick Maher



Lei Wang



Carlos Forsythe



Andrea Young  
Now at MIT



Fereshte Ghahari



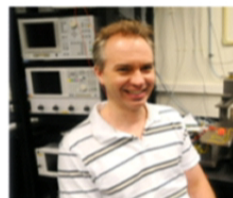
Changyao Chen



Vikram Desphande  
Now at U. Utah



Prof. Jim Hone



Prof. Ken Shepard

Theory: P. Moon & M. Koshino (Tohoku); A. Macdonald (Austin)

hBN samples: T. Taniguchi & K. Watanabe (NIMS); UHV AFM: M. Ishigami (UCF)

**Funding**



# Hofstadter's Butterfly



PHYSICAL REVIEW B

VOLUME 14, NUMBER 6

15 SEPTEMBER 1976

## Energy levels and wave functions of Bloch electrons in rational and irrational magnetic fields\*

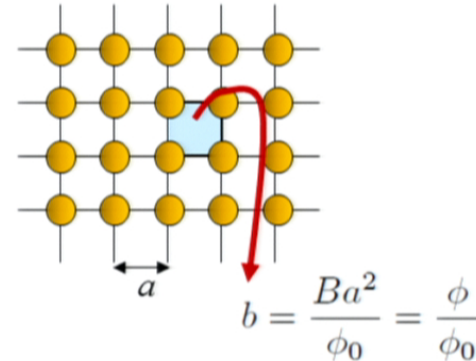
Douglas R. Hofstadter<sup>†</sup>

*Physics Department, University of Oregon, Eugene, Oregon 97403*

(Received 9 February 1976)

Harper's Equation

$$2\psi_l \cos(2\pi lb - \kappa) + \psi_{l+1} + \psi_{l-1} = E\psi_l$$



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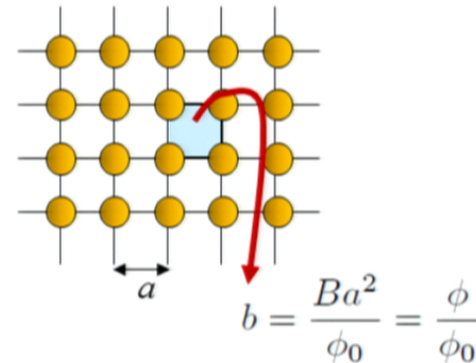
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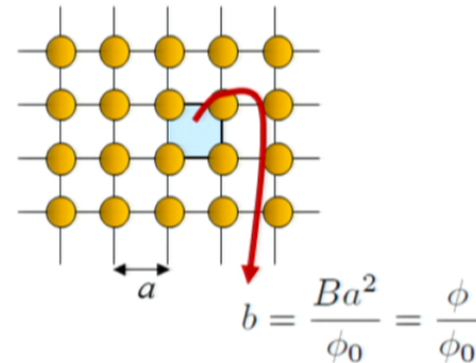
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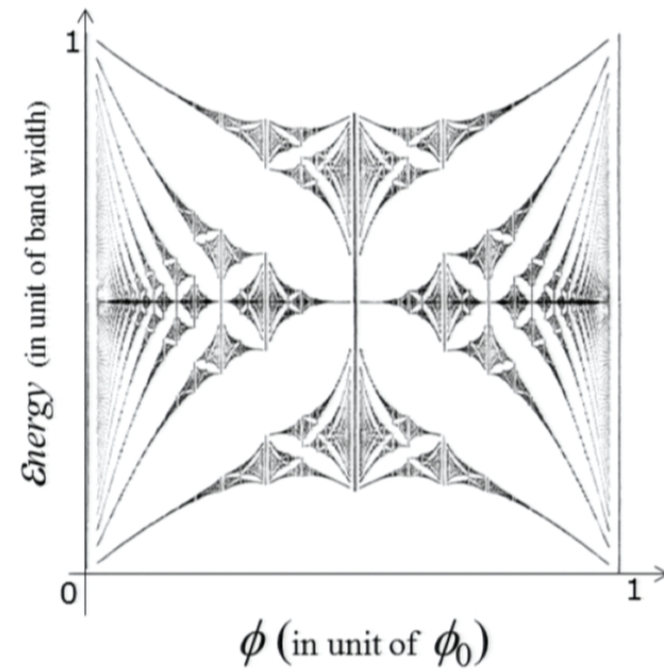
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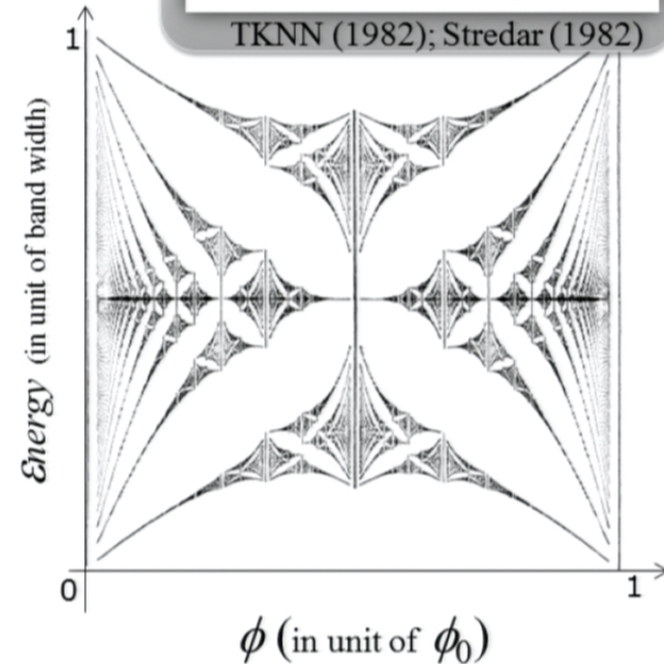
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$$(n/n_0) = t(\phi/\phi_0) + s \quad t, s \in \mathbb{Z}$$

Wannier (1978)





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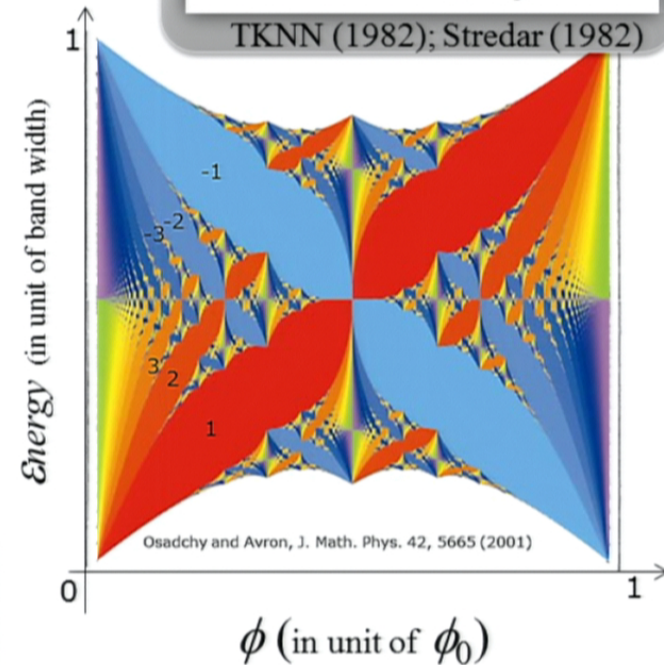
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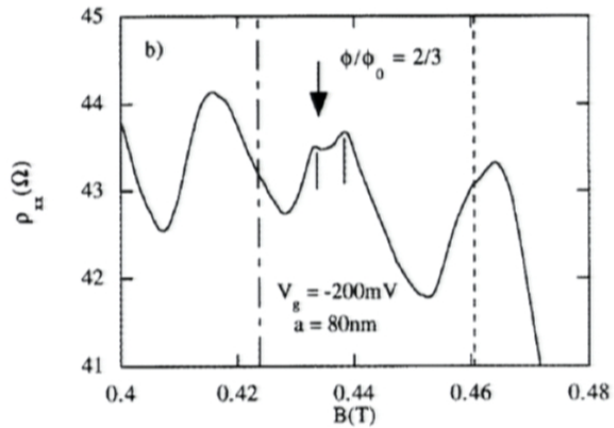
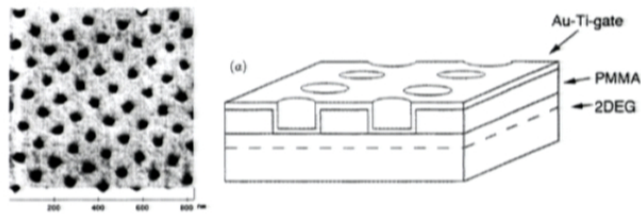
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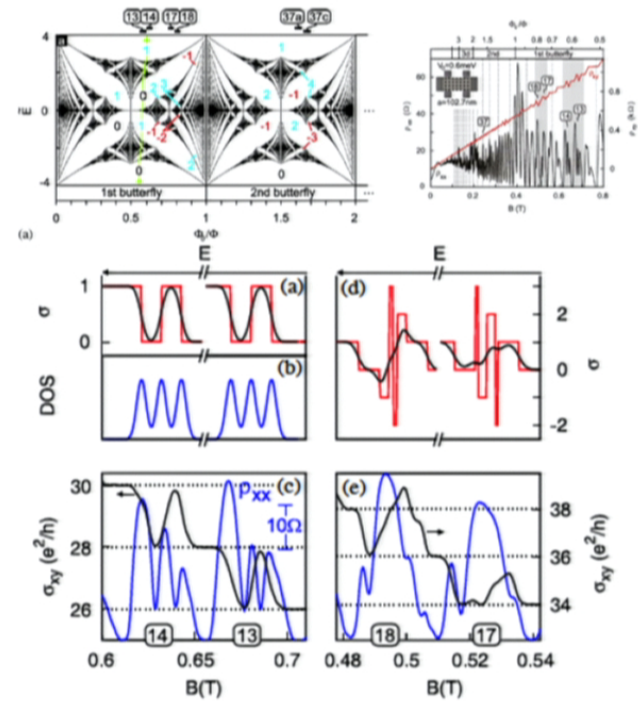
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# Experimental Search For Butterfly



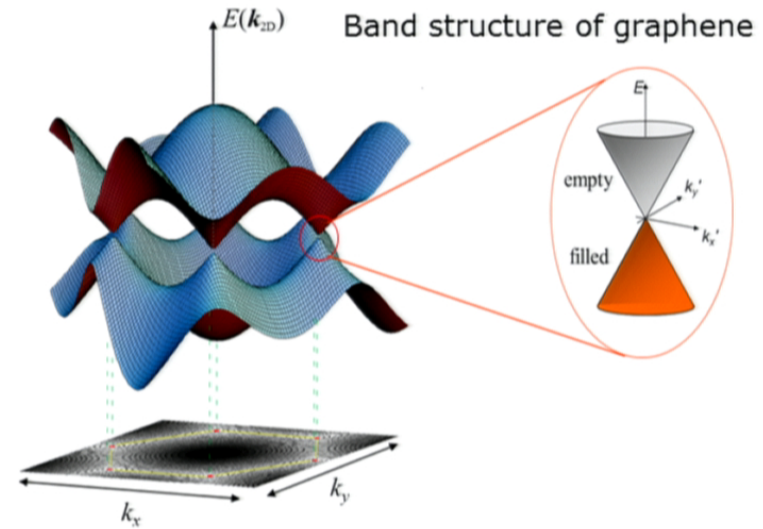
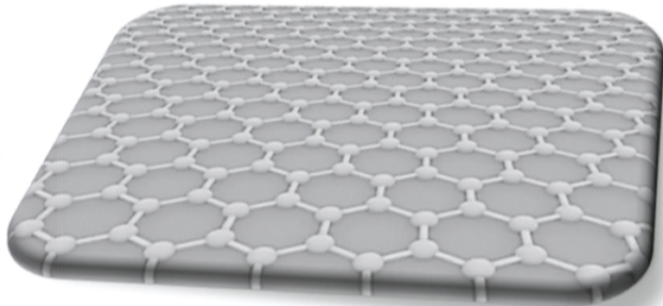
-Schlosser et al, Semicond. Sci. Technol. (1996)



Albrecht et al, PRL. (2001);  
Geisler et al, PRL (2004)

# Electrons in Graphene: Effective Dirac Fermions

Graphene,  
ultimate 2-d conducting system



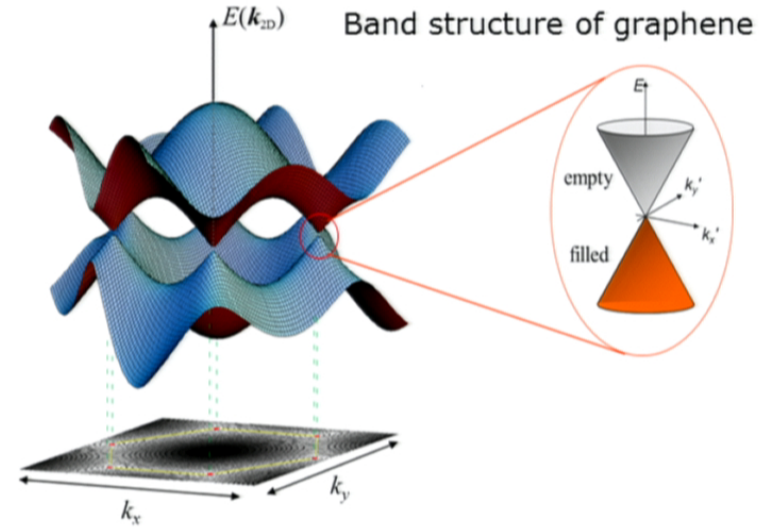
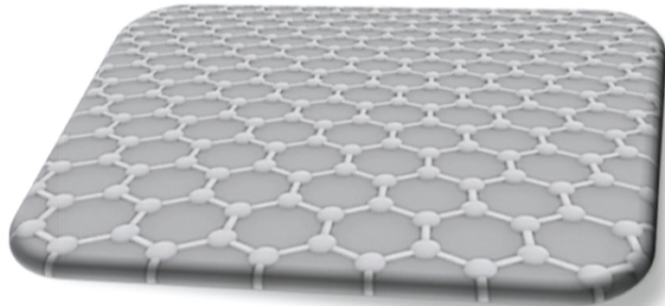
## Effective Dirac Equations

$$H_{eff} = \pm \hbar v_F \begin{pmatrix} 0 & k_x - ik_y \\ k_x + ik_y & 0 \end{pmatrix} = \pm \hbar v_F \vec{\sigma} \cdot \vec{k}_\perp$$

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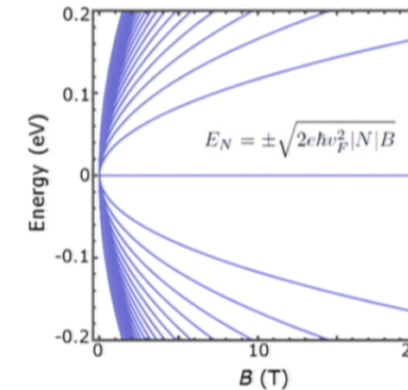


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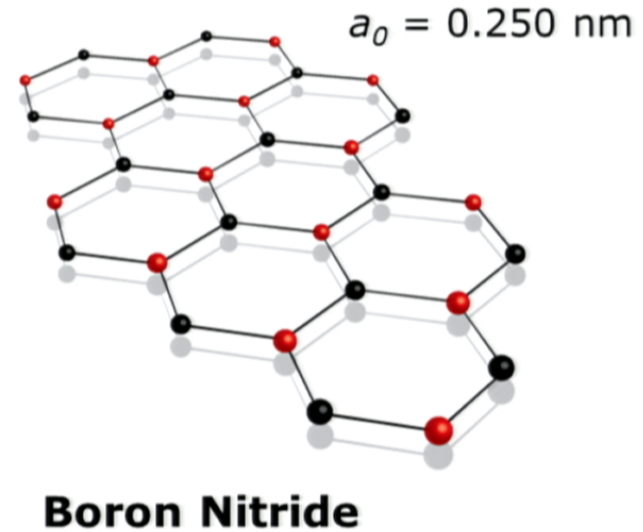
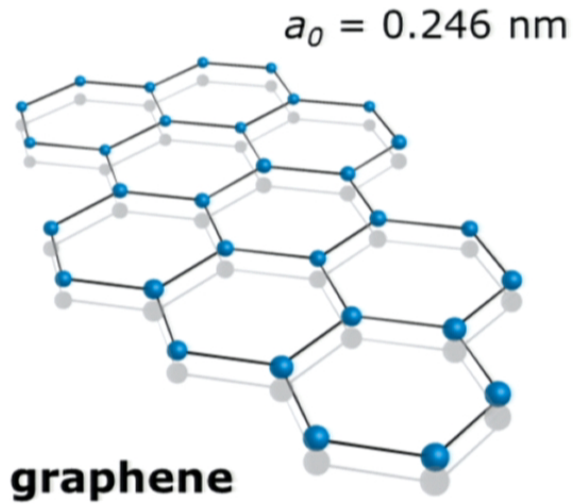
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## Landau Levels in Graphene



# Hexa Boron Nitride: Polymorphic Graphene



## Comparison of h-BN and SiO<sub>2</sub>

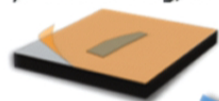
	Band Gap	Dielectric Constant	Optical Phonon Energy	Structure
BN	5.5 eV	~4	>150 meV	Layered crystal
SiO <sub>2</sub>	8.9 eV	3.9	59 meV	Amorphous

# Stacking graphene on hBN

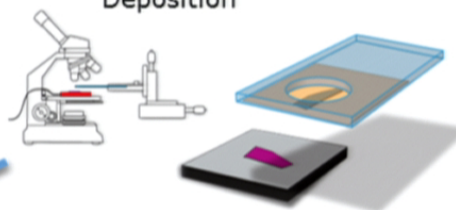
Dean et al. Nature Nano (2009)

(Hone, Kim and Shepard groups @ Columbia)

Polymer coating/cleaving/peeling



Micro-manipulated  
Deposition



Remove polymer  
Annealing



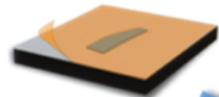
- Co-lamination techniques
- Submicron size precision
- Atomically smooth interface

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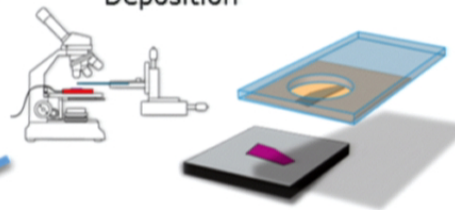
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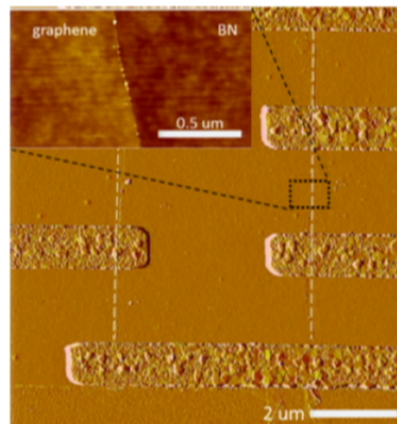
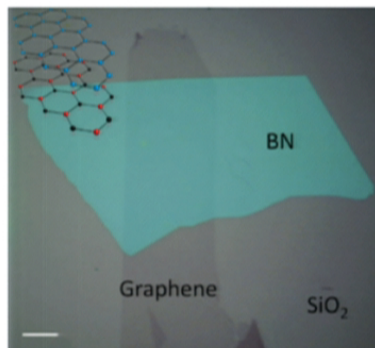
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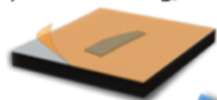
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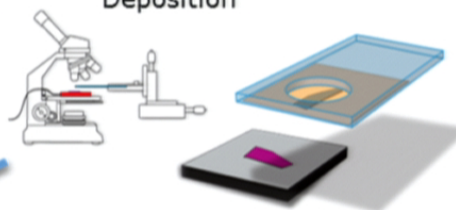
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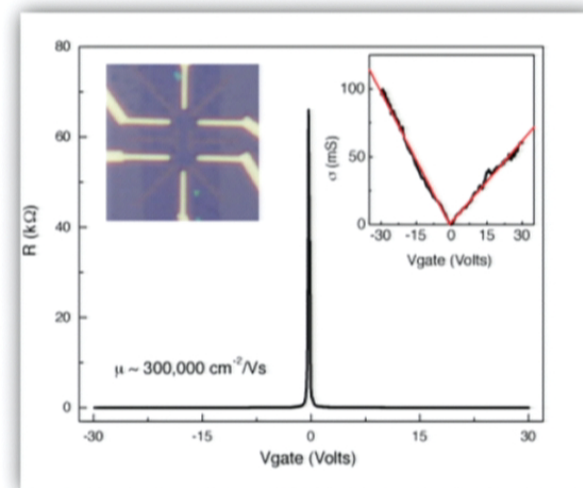
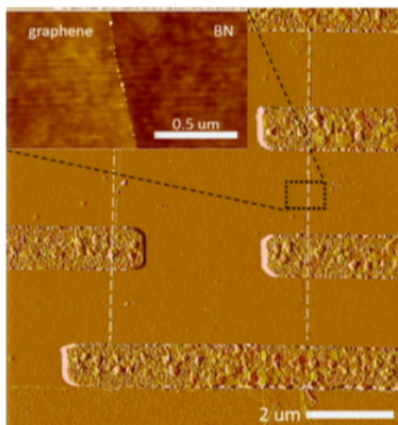
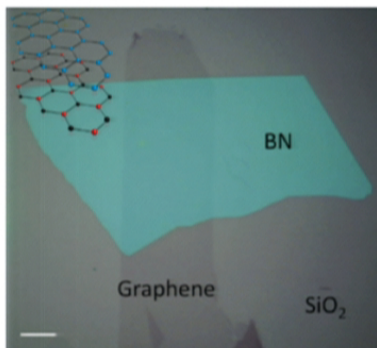
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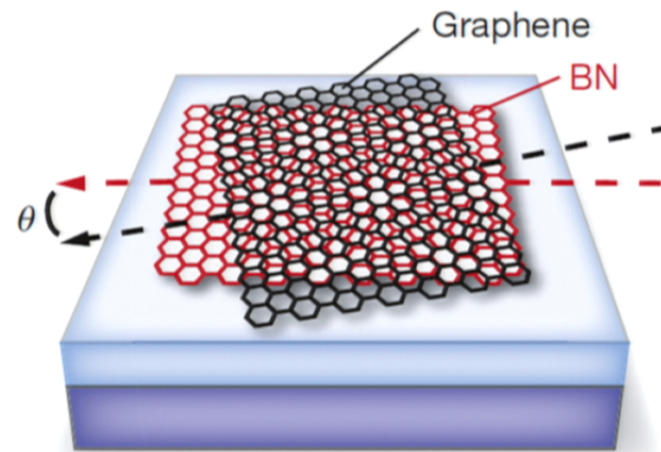


*Mobility > 100,000 cm<sup>2</sup>V<sup>-1</sup>s<sup>-1</sup>*



# Bilayer Graphene on hBN: Moire Pattern

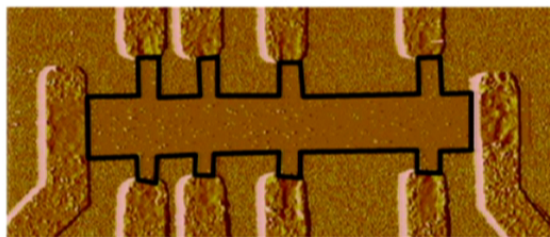
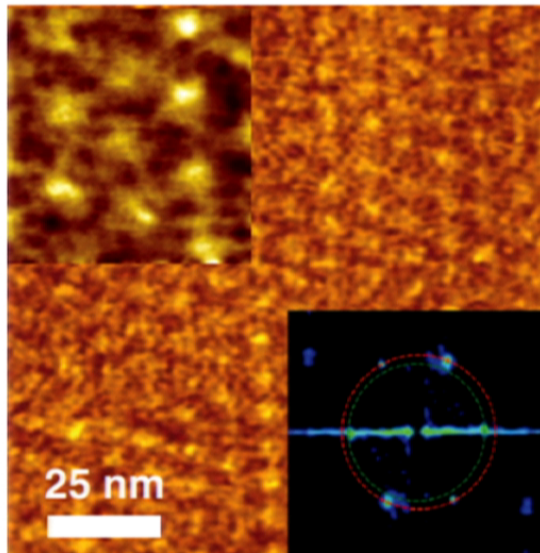
Misaligned graphene on hBN: Moire pattern



Moire wavelength  $\lambda = \frac{(1 + \delta)a}{\sqrt{2(1 + \delta)(1 - \cos(\phi)) + \delta^2}}$

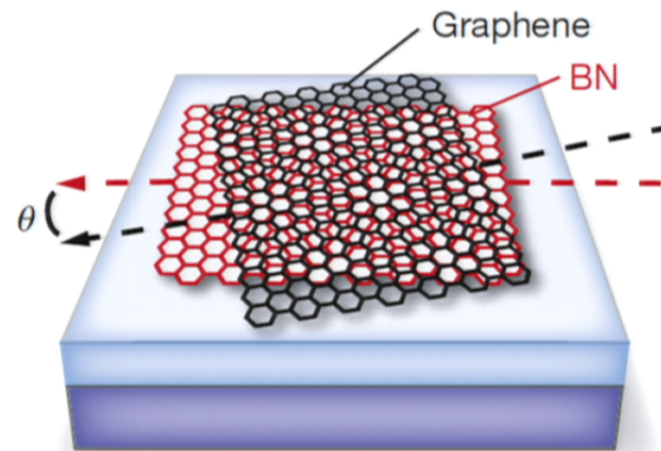
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UHV AFM (Ishigami group)



**Moiré  $\lambda_{AFM} = 15.5 \text{ nm}$**

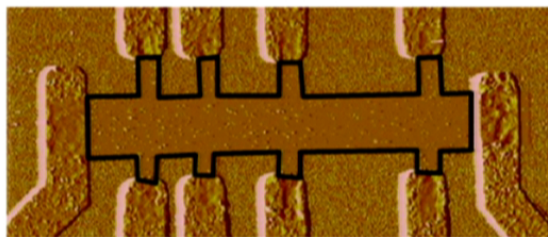
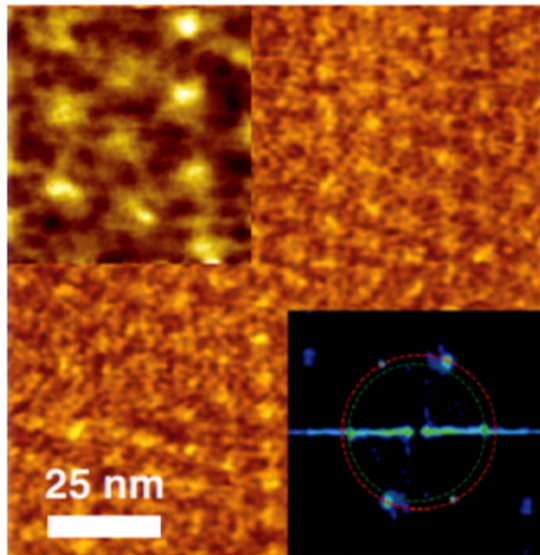
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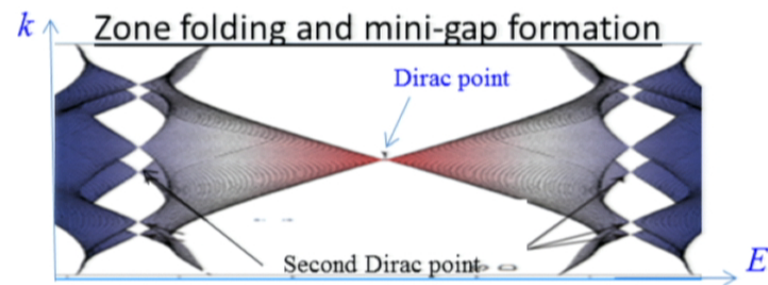
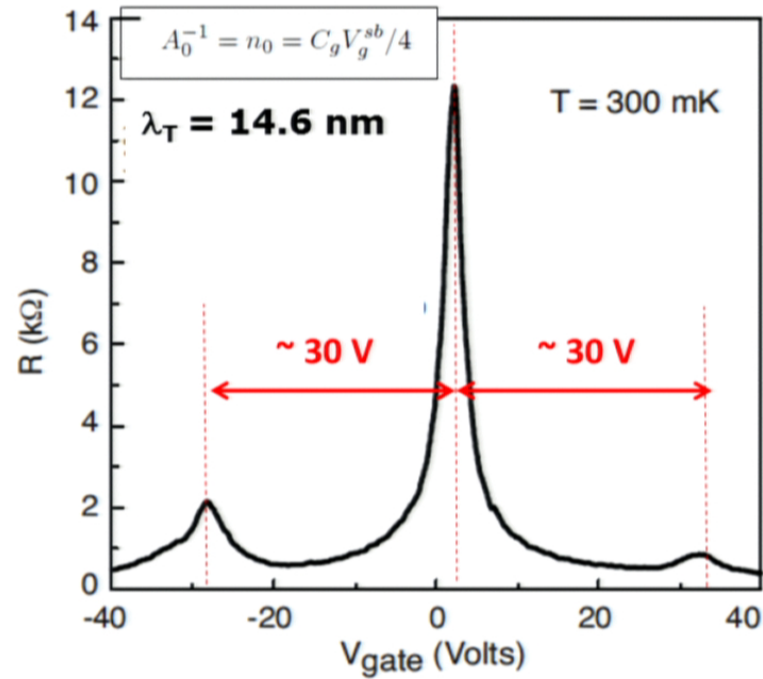
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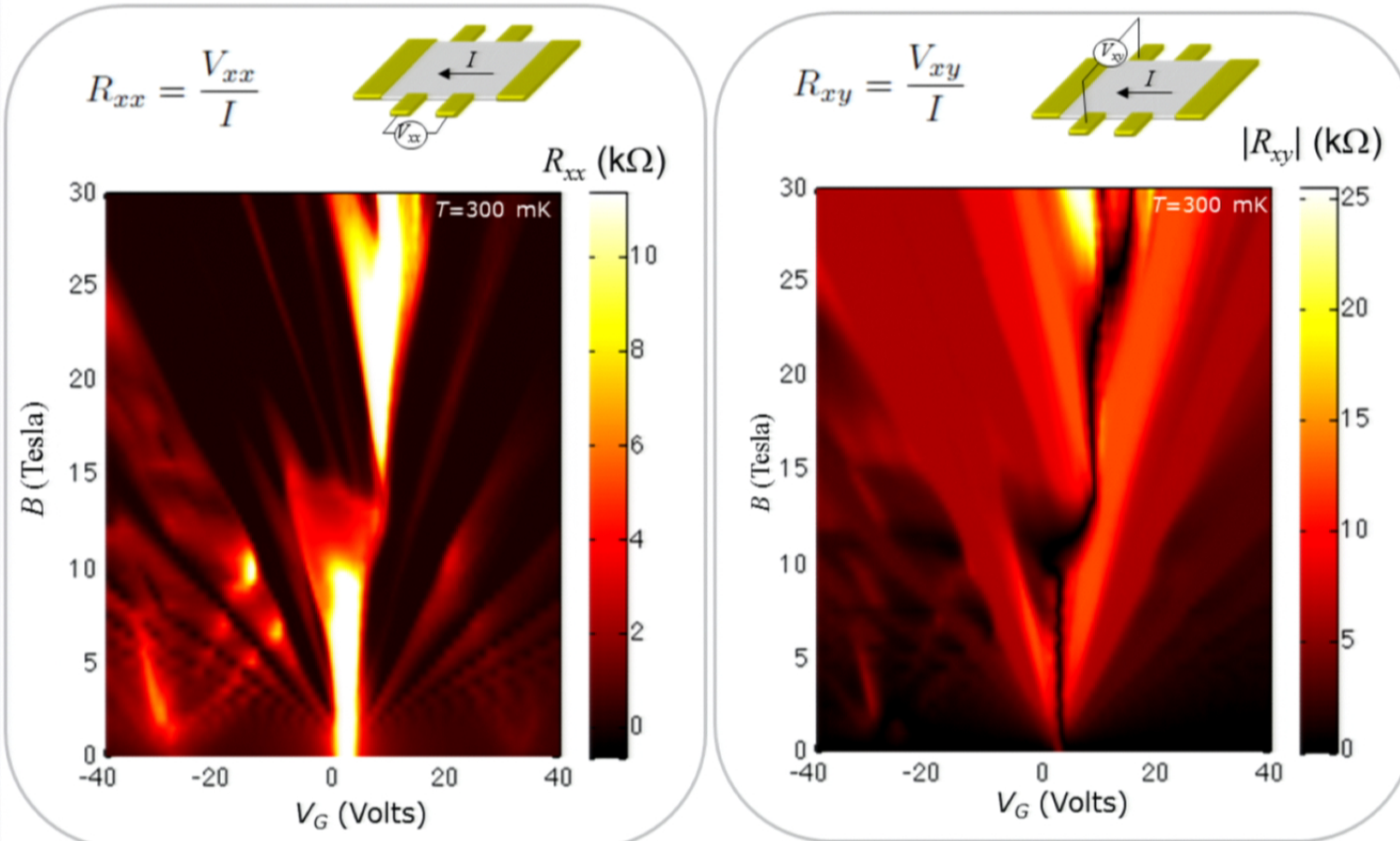


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# Abnormal Landau Fan Diagram in Graphene on hBN

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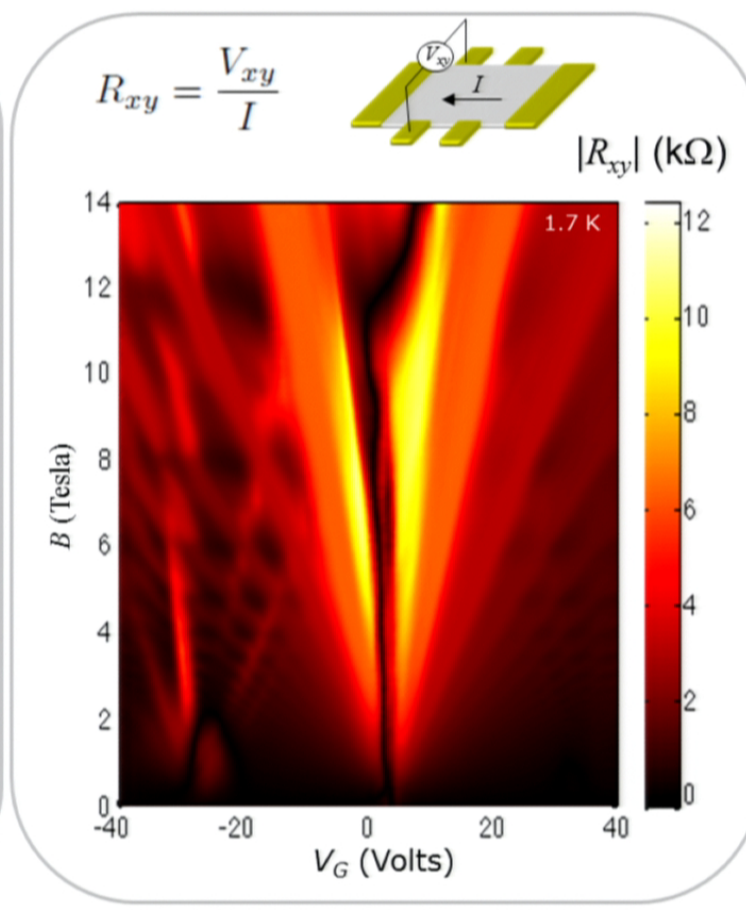
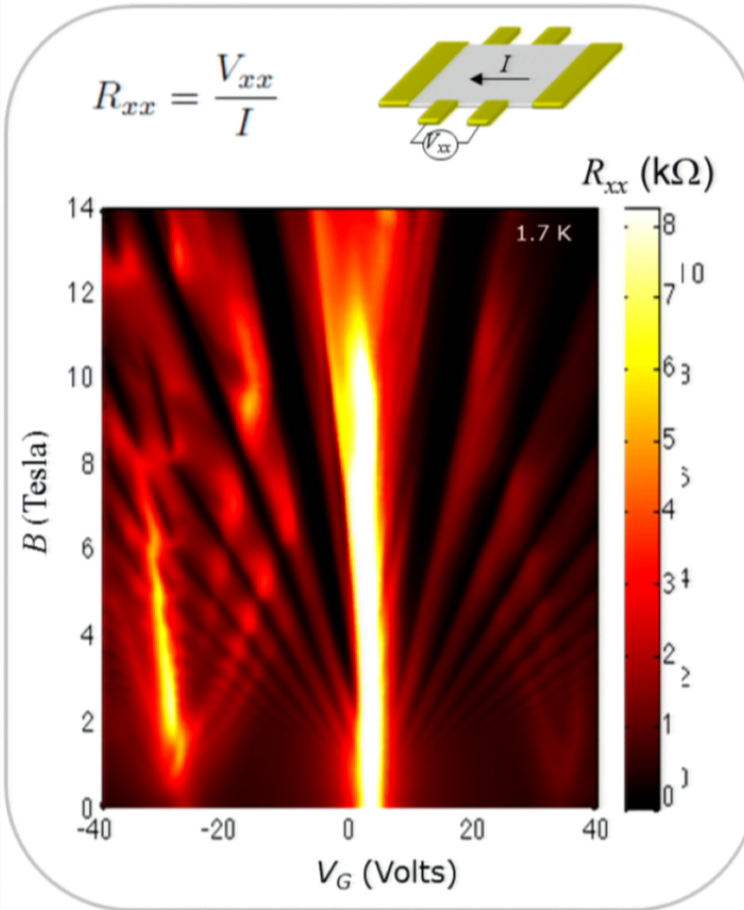


Similar results were reported on single layer on hBN by Gorbachev *et al*. *Nature* (2013) and Hunt *et al*. *Science* (2013)

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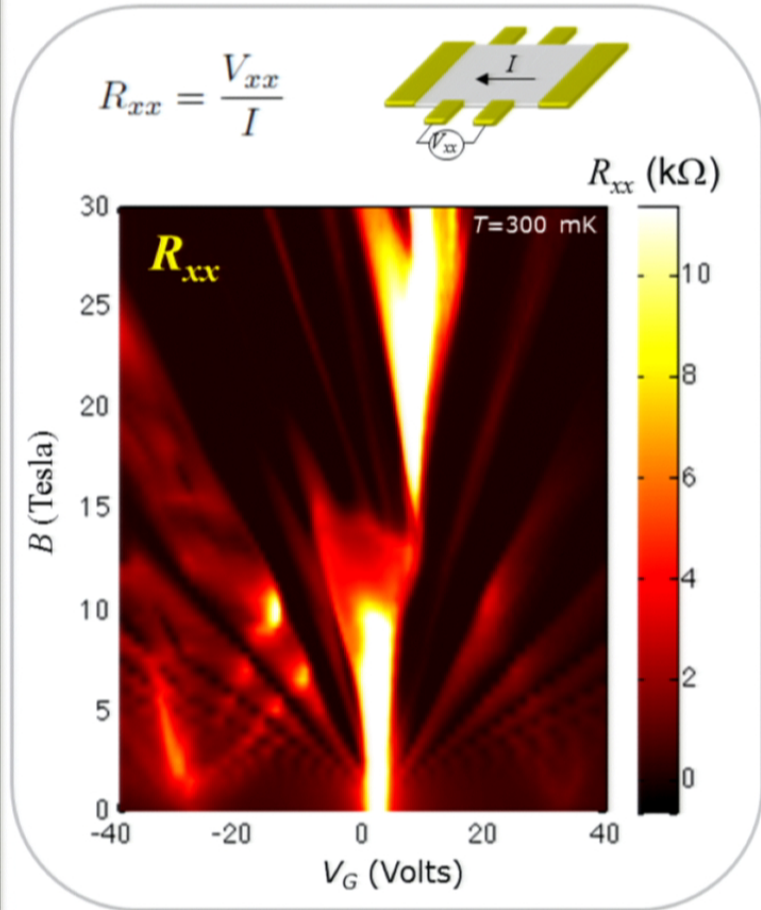
Low Magnetic field regime



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# Quantum Hall Effect in Graphene Moire

Quantum Hall-like Transport



Landau level  
filling factor

$$\nu = \frac{\phi_0}{B} n$$

Quantum Hall  
conductance

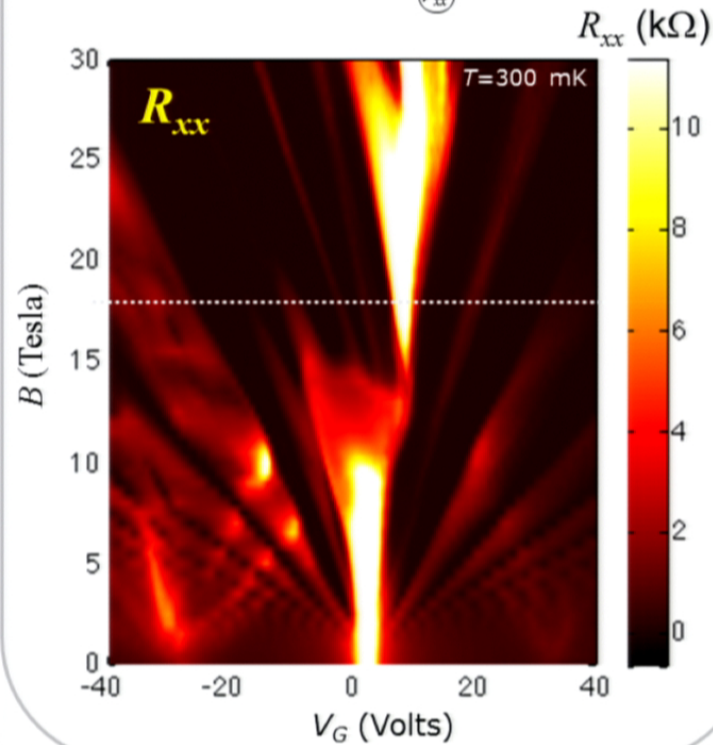
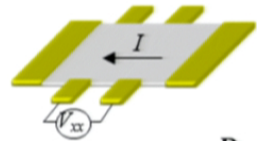
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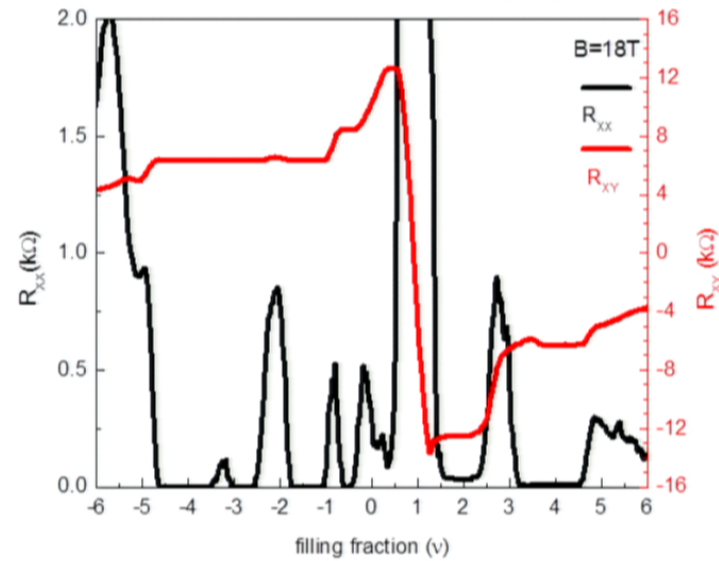
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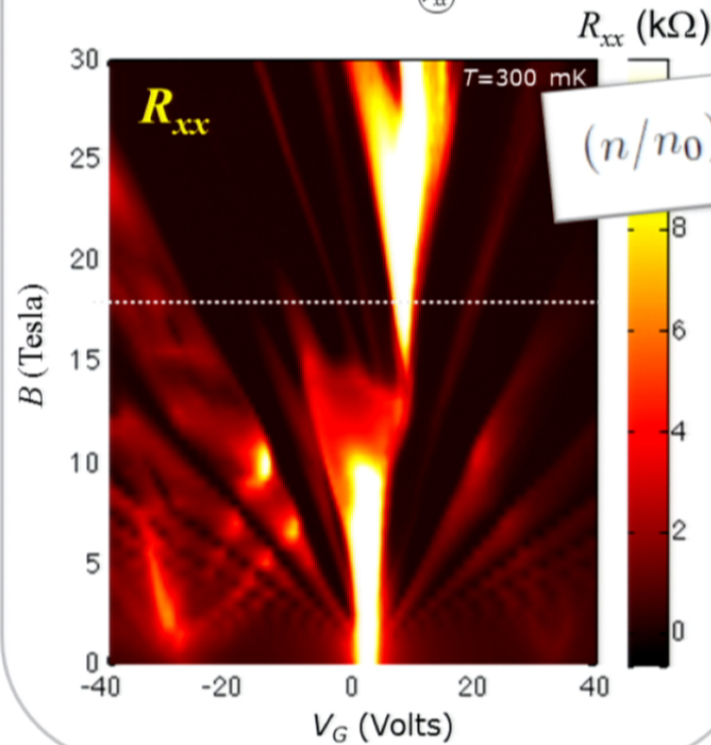
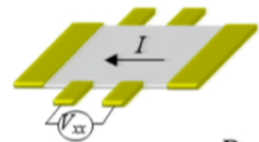
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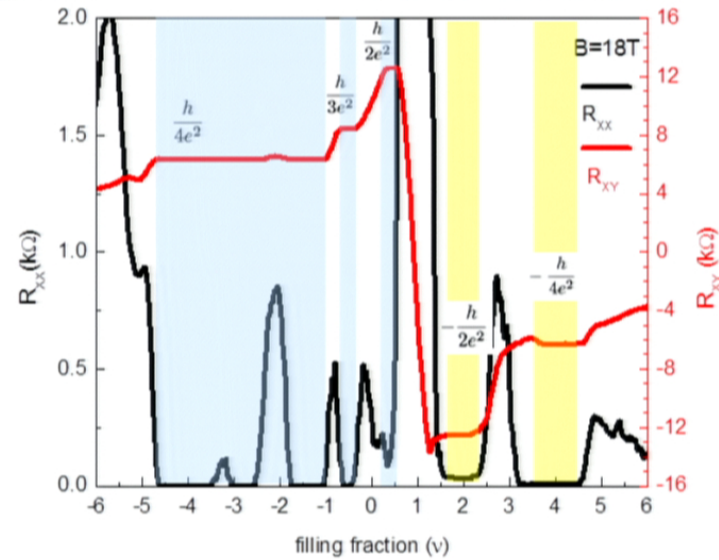
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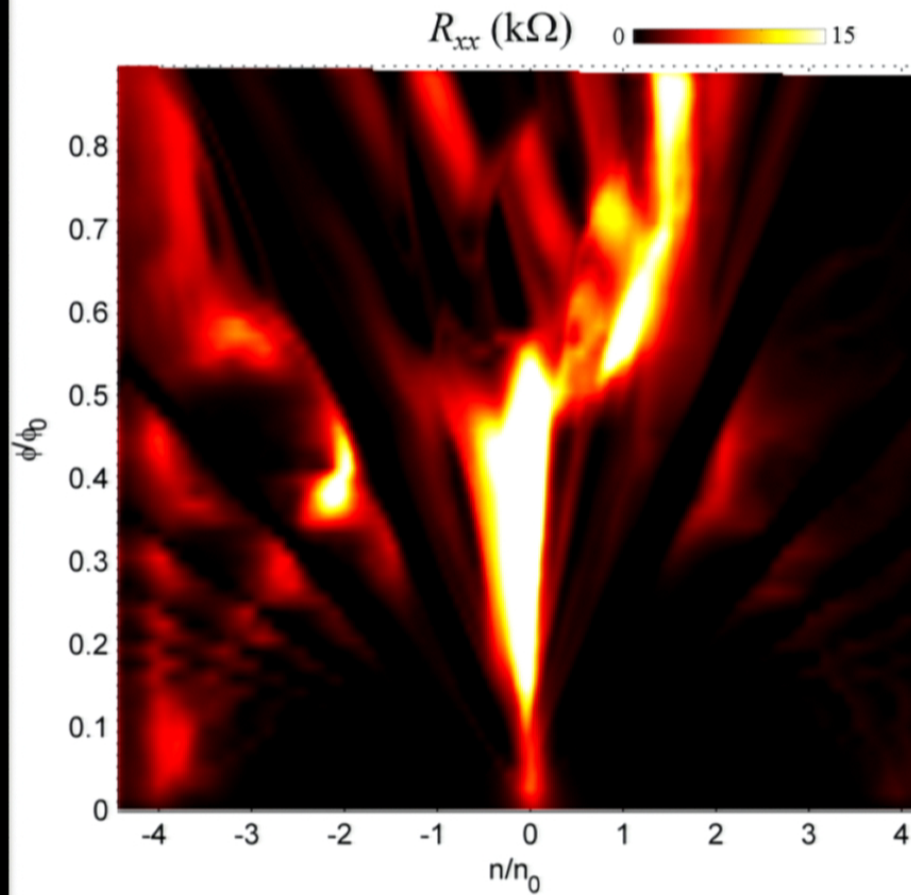
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# Quantum Hall Effect with Two Integer Numbers

$n/n_0$ : density per unit cell;  $\phi$ : flux per unit cell



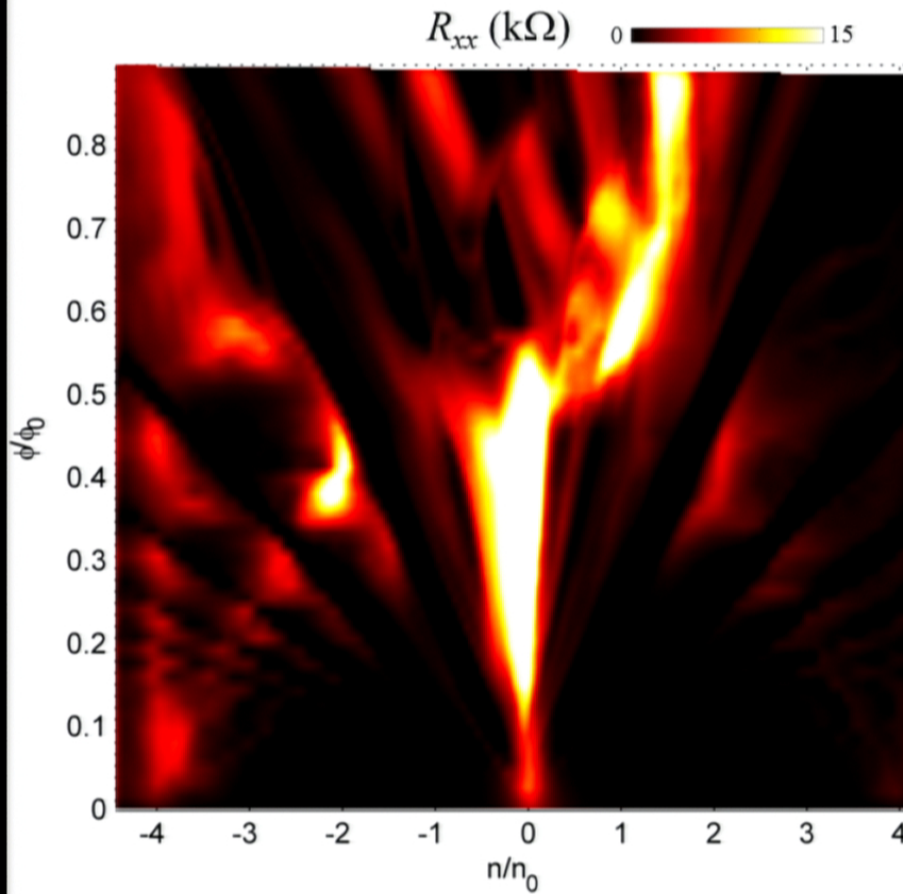
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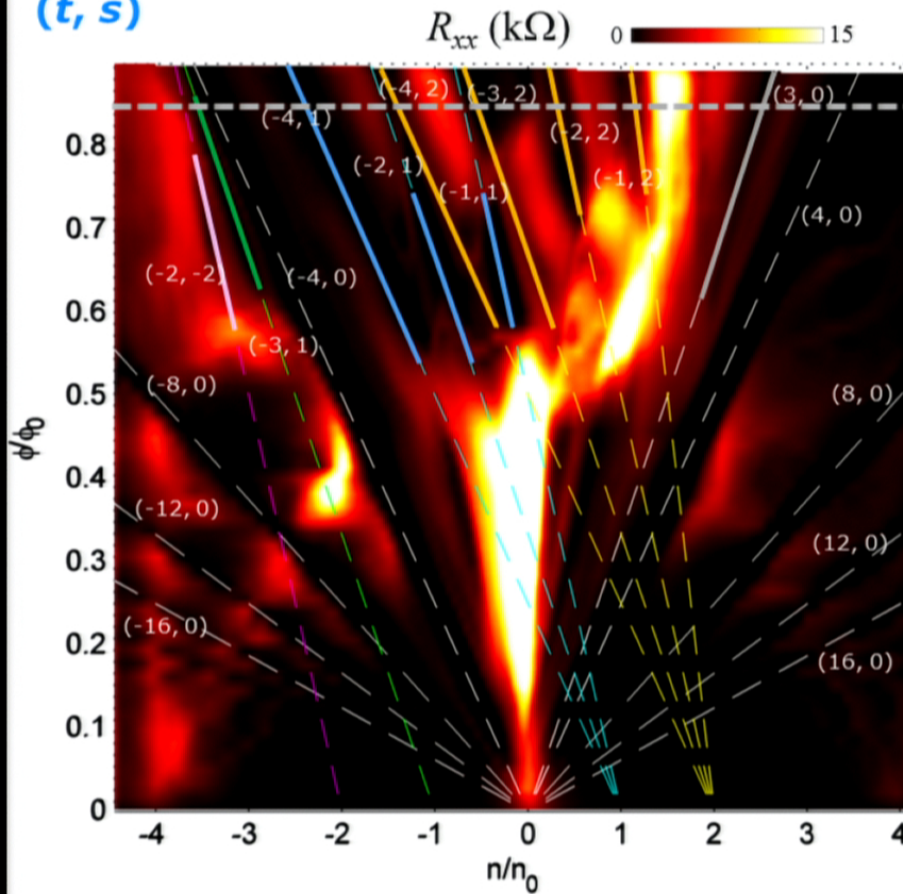
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$(t, s)$

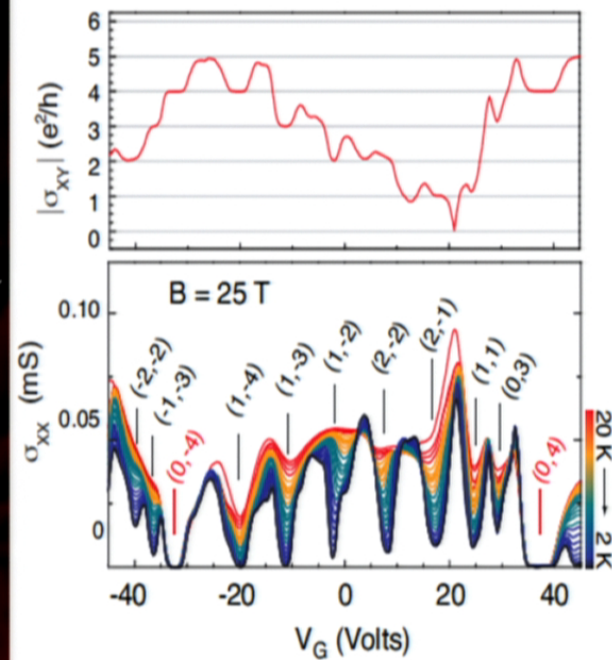
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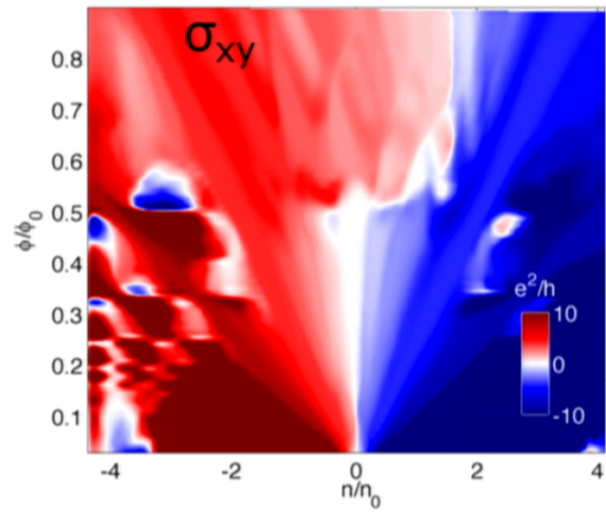


Quantization of  $\sigma_{xx}$  and  $\sigma_{xy}$



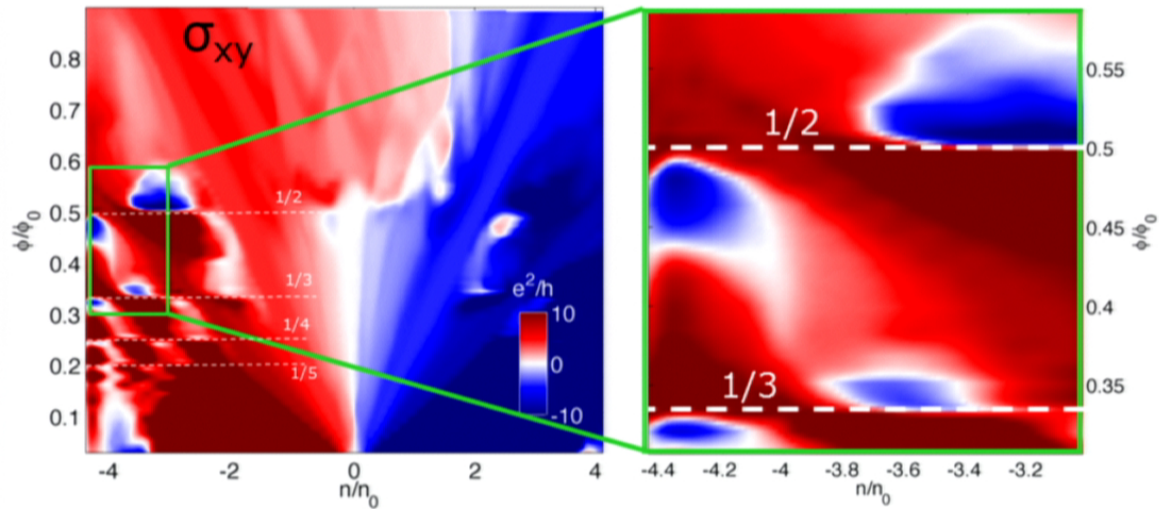
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Higher quality sample with lower disorder



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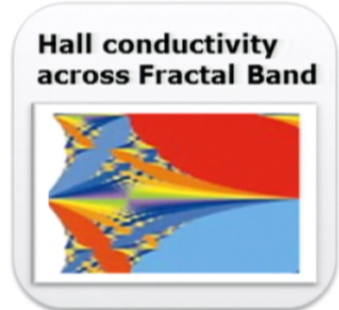
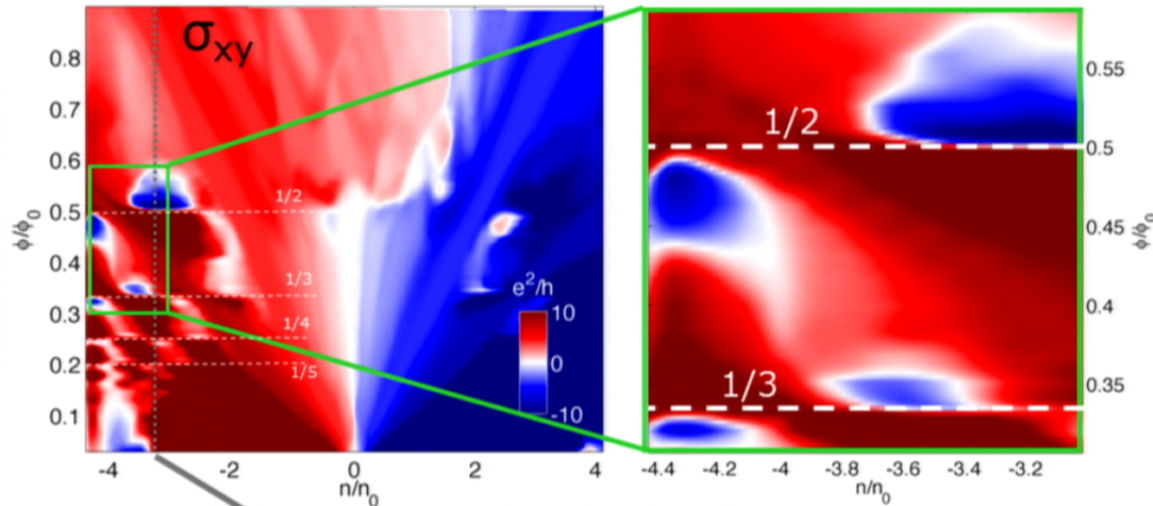


Hall conductivity  
across Fractal Band



# Recursive QHE near the Fractal Bands

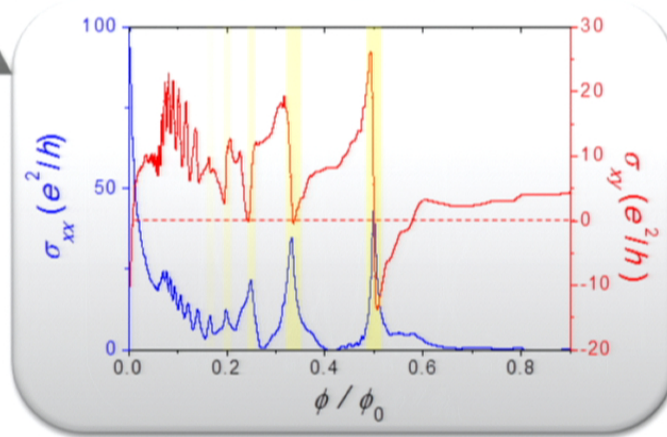
Higher quality sample with lower disorder



At the Fractal Bands

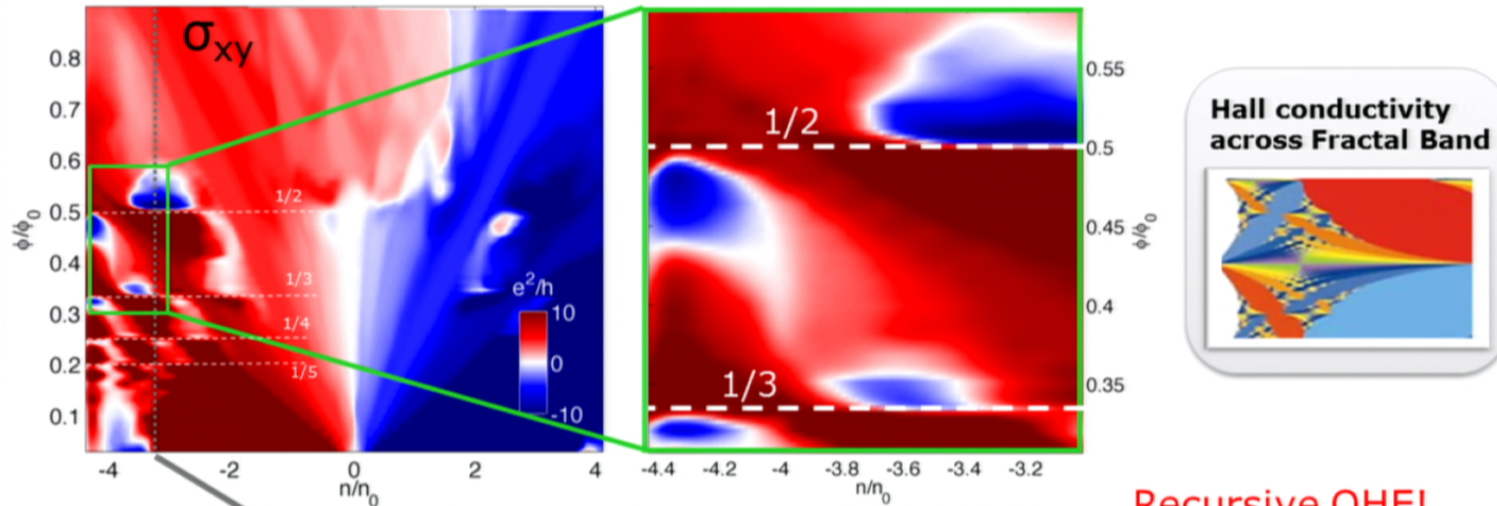
Sign reversal of  $\sigma_{xy}$

Large enhancement of  $\sigma_{xx}$



# Recursive QHE near the Fractal Bands

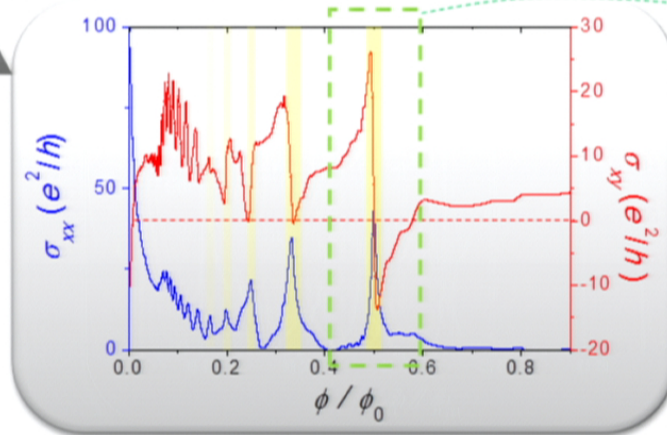
Higher quality sample with lower disorder



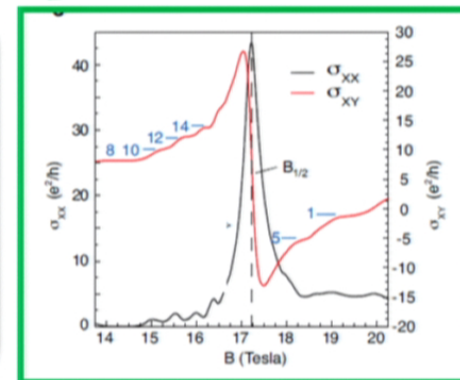
At the Fractal Bands

Sign reversal of  $\sigma_{xy}$

Large enhancement of  $\sigma_{xx}$

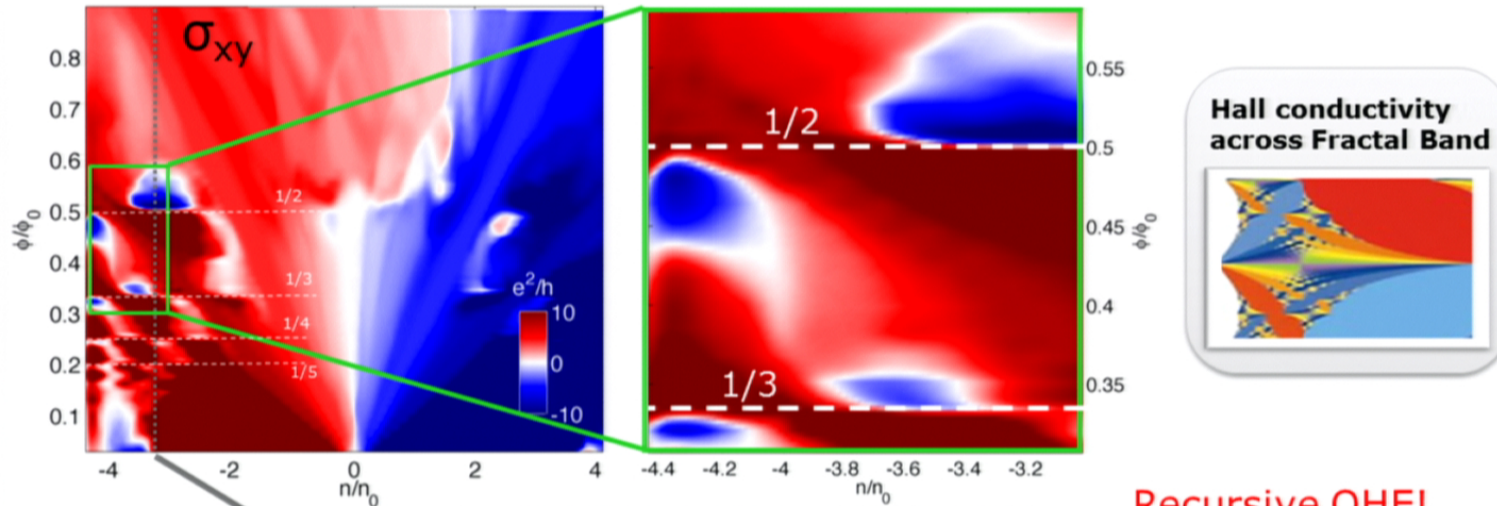


Recursive QHE!



# Recursive QHE near the Fractal Bands

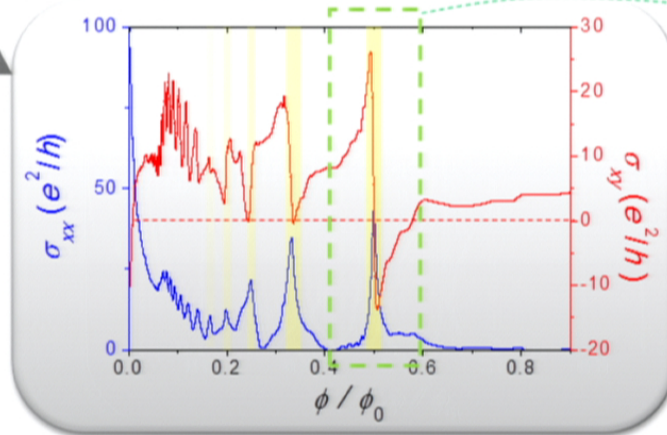
Higher quality sample with lower disorder



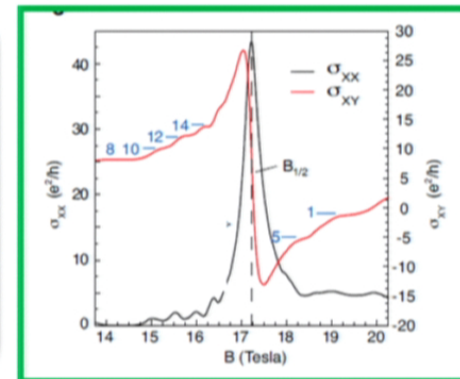
At the Fractal Bands

Sign reversal of  $\sigma_{xy}$

Large enhancement of  $\sigma_{xx}$

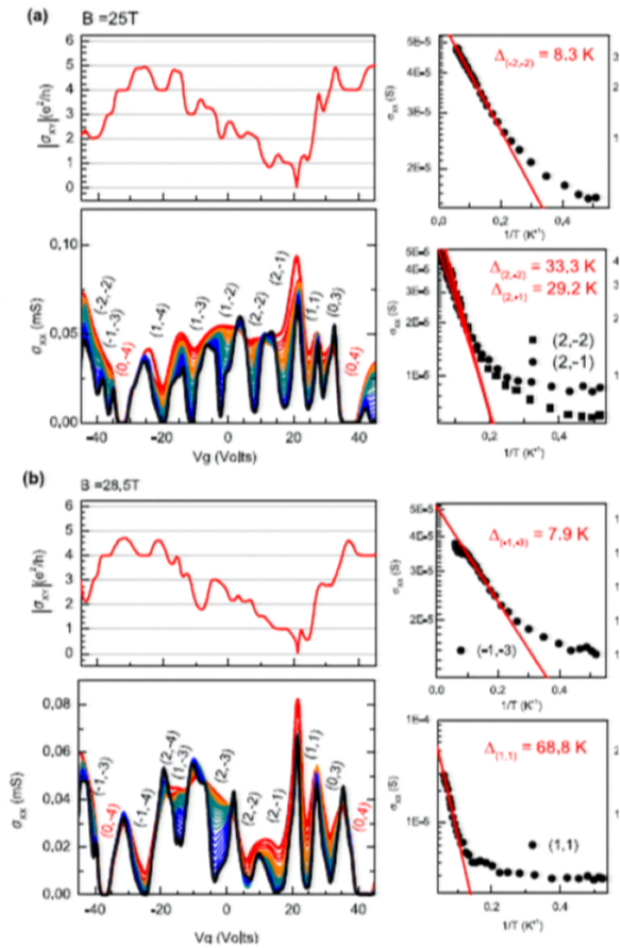


Recursive QHE!

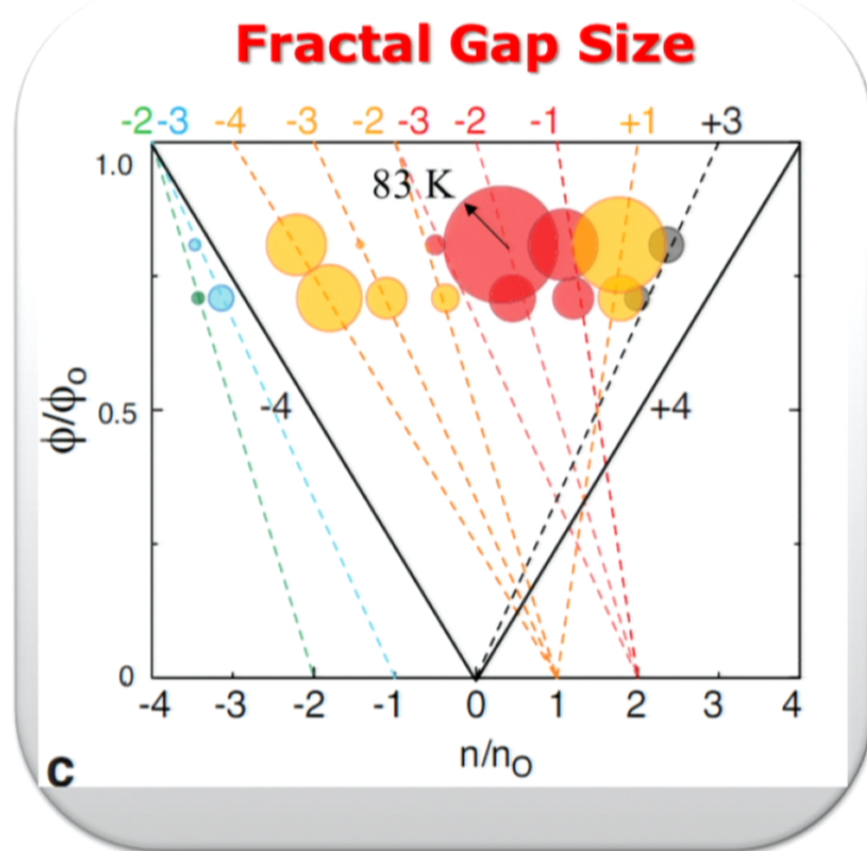
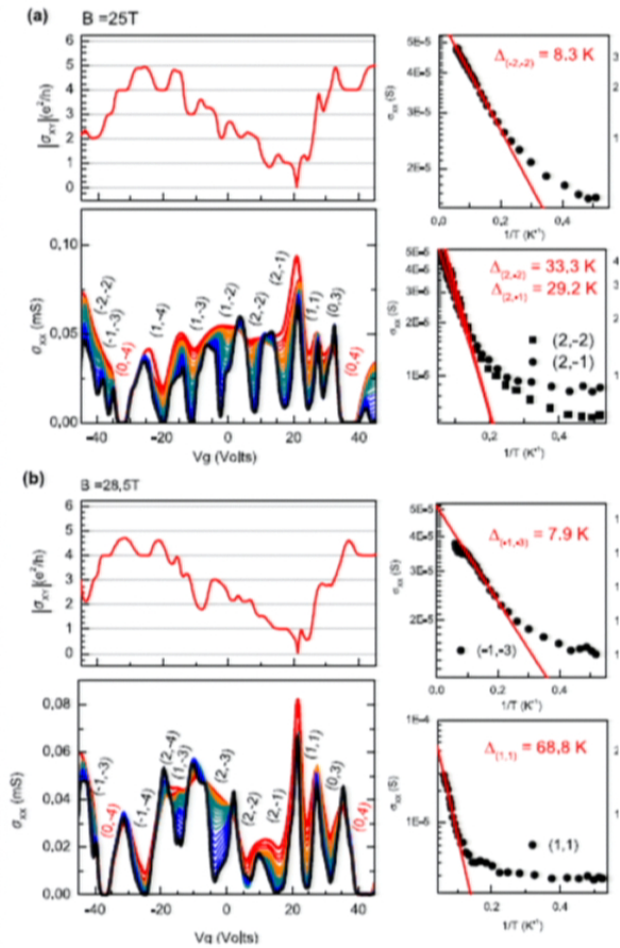




# Fractal Gaps: Energy Scales

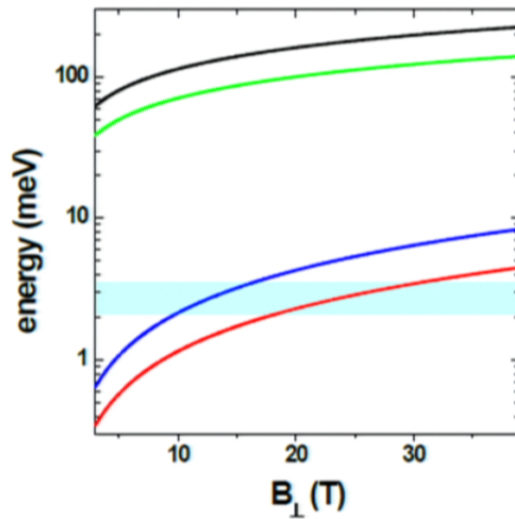


# Fractal Gaps: Energy Scales



# Graphene SU(4): Energy Scales

## Monolayer Graphene



—	$E_0$	$\frac{\hbar v_F \sqrt{2}}{l_B}$
—	$E_C$	$\frac{e^2}{\epsilon l_B}$
—	$E_{SR}$	$\frac{a e^2}{l_B \epsilon l_B}$
—	$E_Z$	$g \mu_B B_T$
—	$\Gamma$	<i>exp.</i>

Cyclotron Energy

Coulomb Energy

Short Range Lattice Interaction  
(Pseudo/valley spin anisotropy)

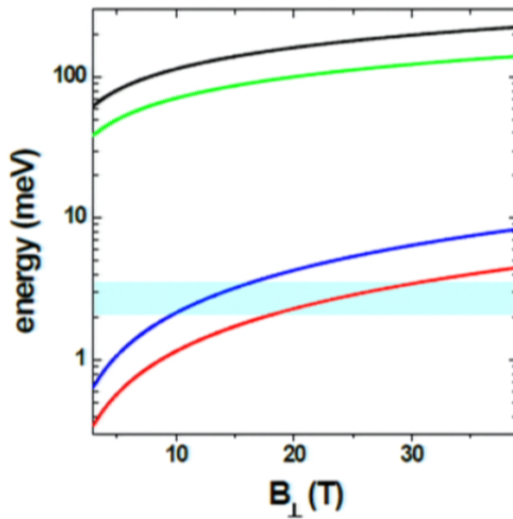
Zeeman Energy  
(Spin anisotropy)

Disorder

Goerbig, RMP (2011)

# Graphene SU(4): Energy Scales

## Monolayer Graphene



— $E_0$	$\frac{\hbar v_F \sqrt{2}}{l_B}$
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Cyclotron Energy

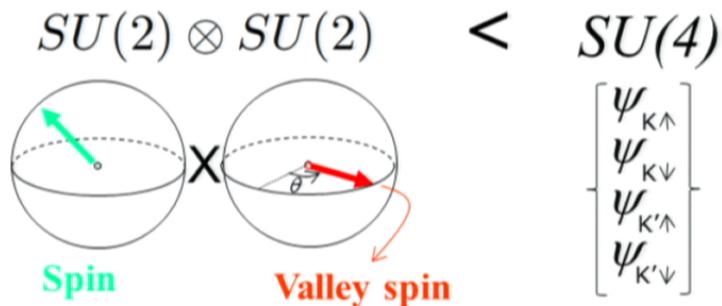
Coulomb Energy

Short Range Lattice Interaction  
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Zeeman Energy  
(Spin anisotropy)

Disorder

Goerbig, RMP (2011)

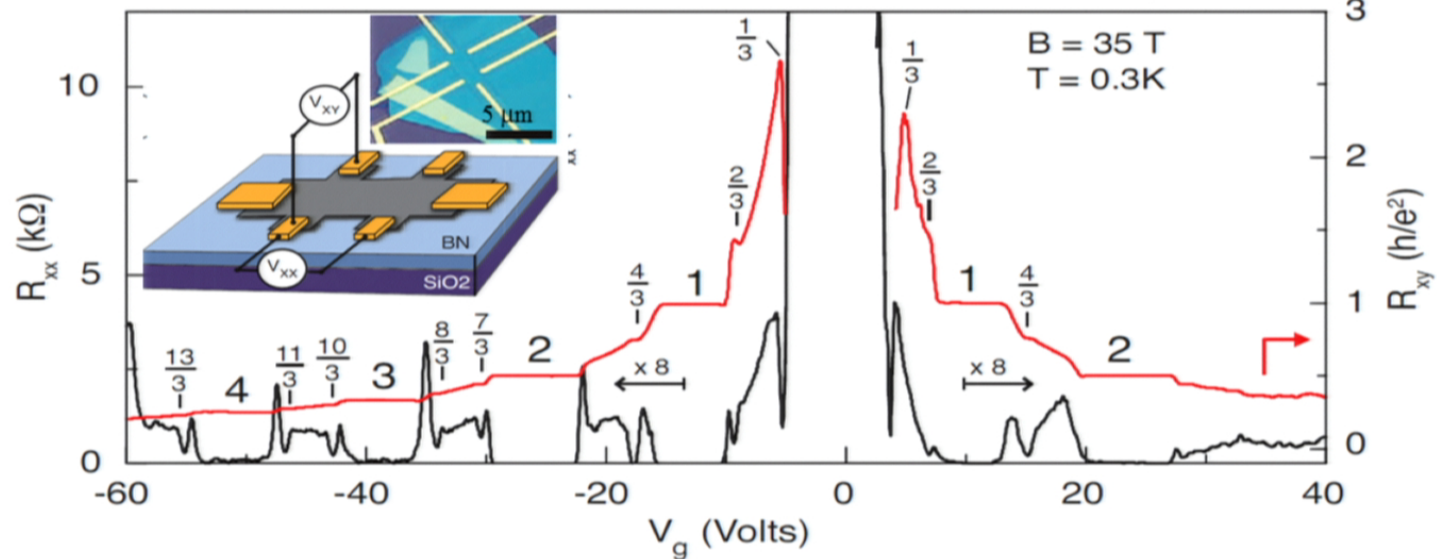


**$SU(4)$  symmetry is conjectured to produce rich new set of ground state wavefunctions in graphene**

Yang, Das Sarma and MacDonald, PRB (2006);

# Spin & pseudo spins: many body physics in graphene

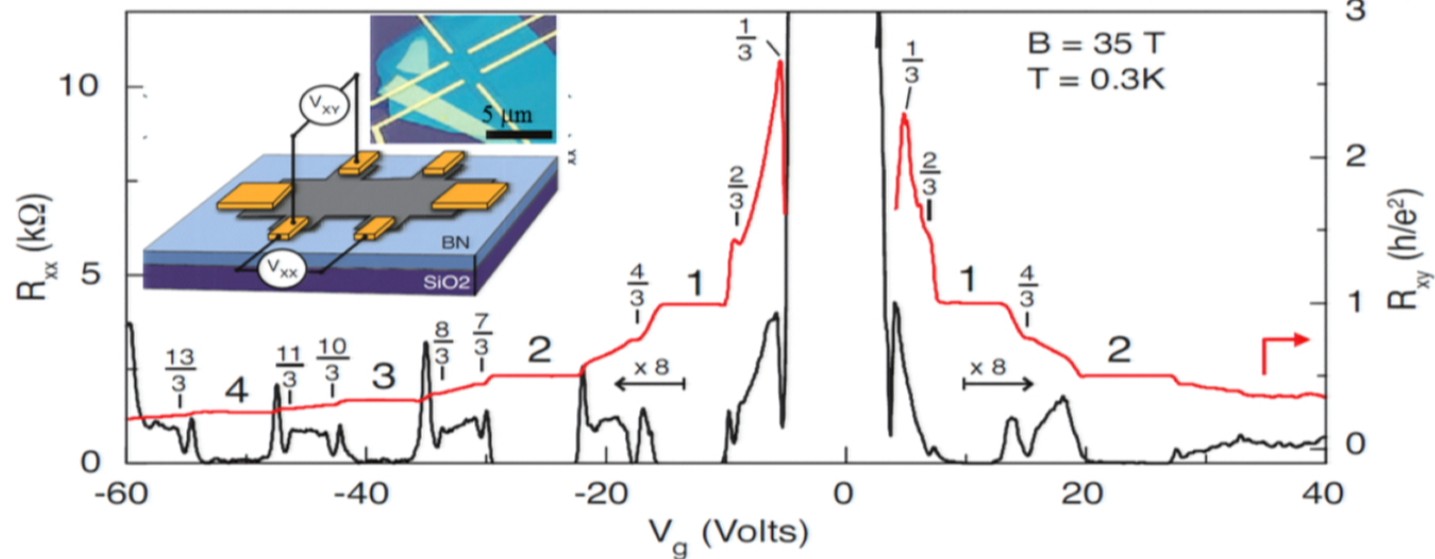
Dean *et al.*, Nature Physics (2011)  
Young *et al.*, Nature Physics (2012)



Integer quantum Hall Ferromagnetism and Fractional Quantum Hall Effect in single layer graphene

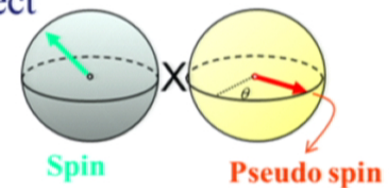
# Spin & pseudo spins: many body physics in graphene

Dean *et al.*, Nature Physics (2011)  
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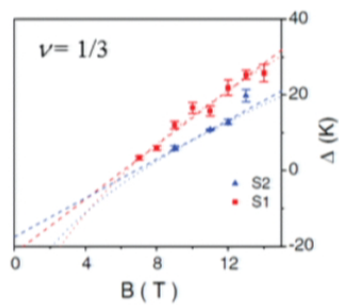


Integer quantum Hall Ferromagnetism and Fractional Quantum Hall Effect in single layer graphene

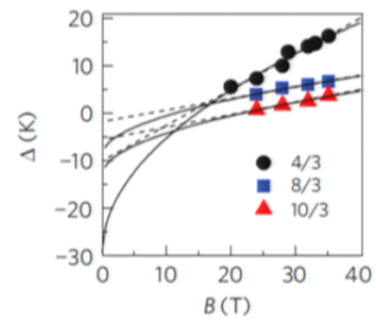
- SU(4) hierarchical Integer and Fractional Quantum Hall Effect
- Spin and Pseudospin Ferromagnetic Quantum Hall Effect
- Spin Skyrmion and Valley Skyrmions



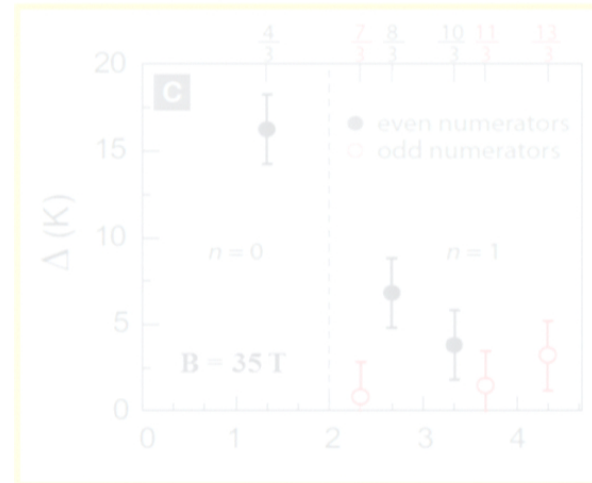
## Hierarchical Structure in Fractional Quantum Hall Gaps



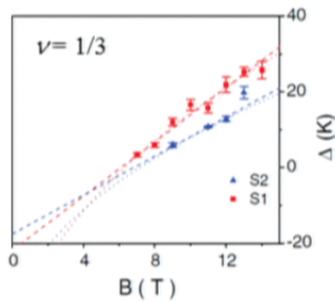
Ghahari et al. PRL (2011)



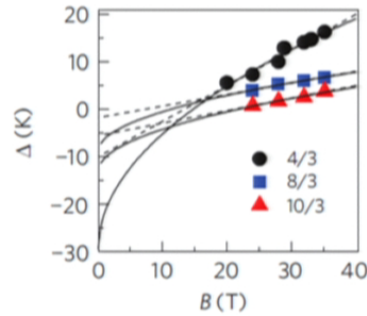
Young et al. Nature Phys. (2012)



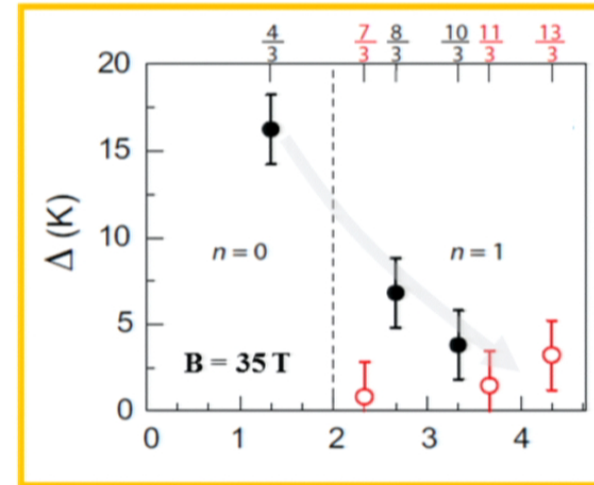
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Ghahari et al. PRL (2011)

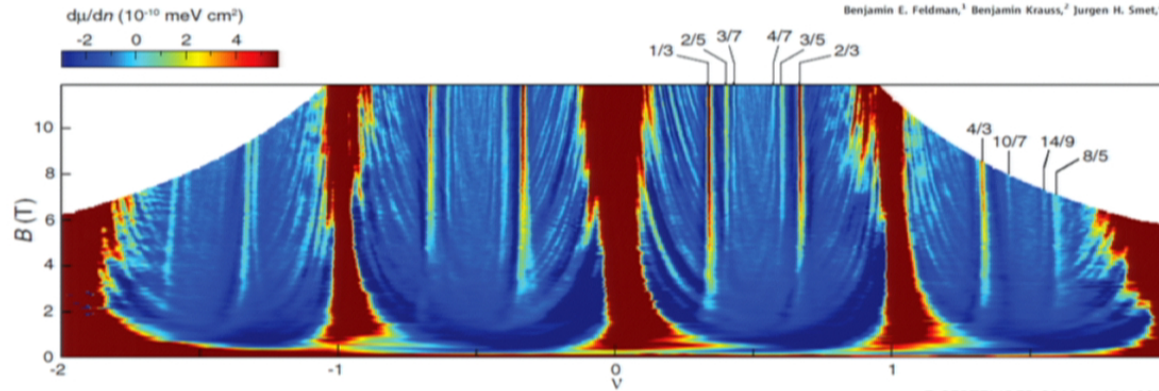


Young et al. Nature Phys. (2012)



## Unconventional Sequence of Fractional Quantum Hall States in Suspended Graphene

Benjamin E. Feldman,<sup>1</sup> Benjamin Krauss,<sup>2</sup> Jurgen H. Smet,<sup>2</sup> Amir Yacoby<sup>1\*</sup>

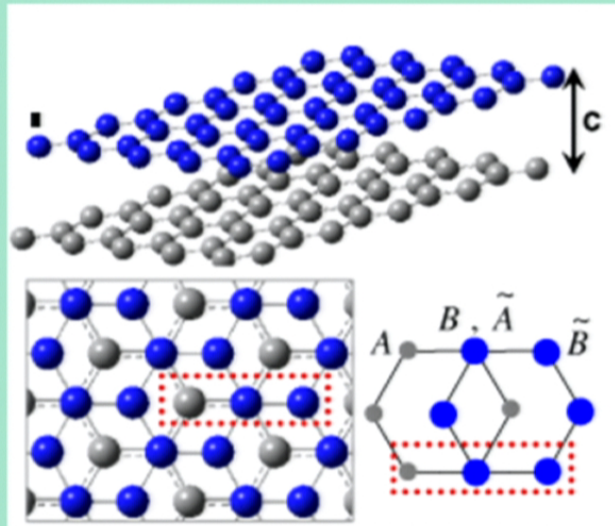


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# Pseudo Spin/ Valley Spin in Bilayer Graphene

Bilayer graphene and pseudo spin

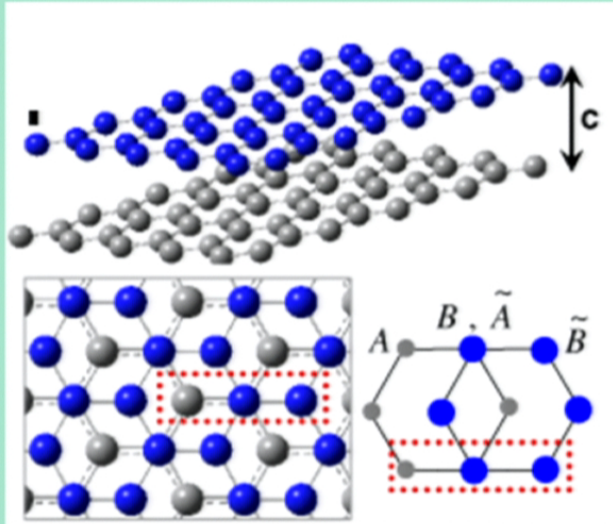


F. Guinea, A. H. C. Neto, & N. M. R. Peres, Phys. Rev. B 73, 245426 (2006)

E. McCann, Phys. Rev. B 74, 161403 (2006)

# Pseudo Spin/ Valley Spin in Bilayer Graphene

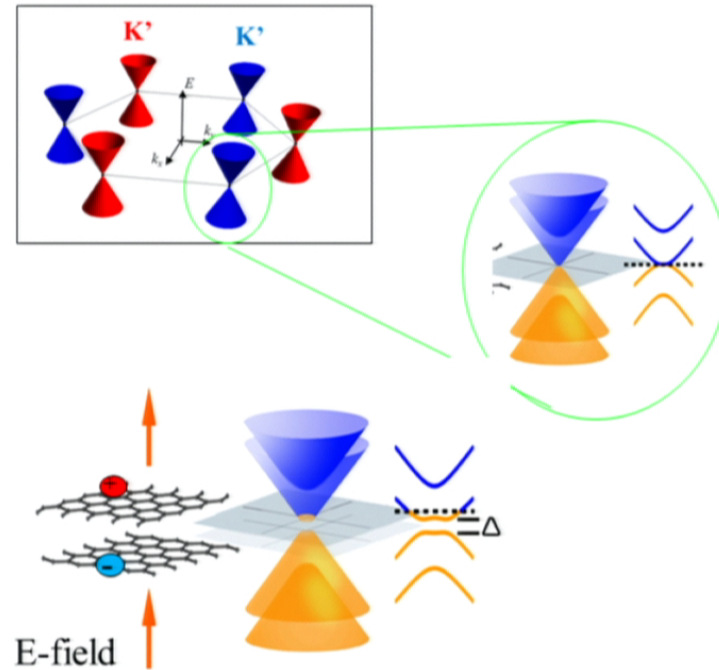
Bilayer graphene and pseudo spin



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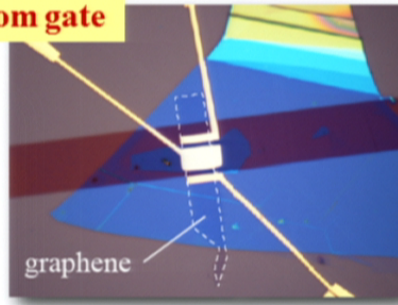
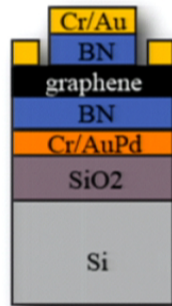
Low energy approximation in 1<sup>st</sup> BZ



Perpendicular electric field breaks layer symmetry, polarize pseudo spin, and open up a gap

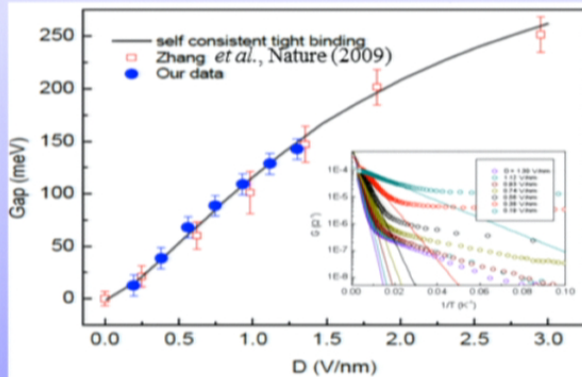
# Gap Opening in Bilayer Graphene

## BN top and bottom gate



Wang et al. (Hone, Kim, Shepard) (2012)

## Electric Field Induced Layer Polarization Induced Gaps

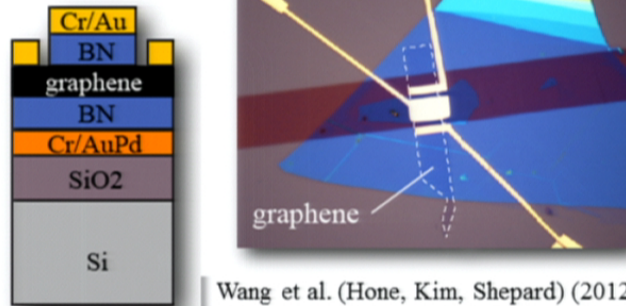


**>200 meV transport gaps may be possible!**

see also Yan et al (2010), Cheng et al (2010), Taychatanapat (2010)

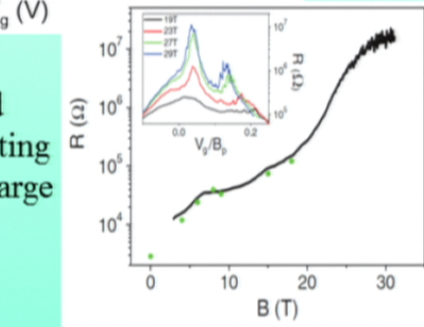
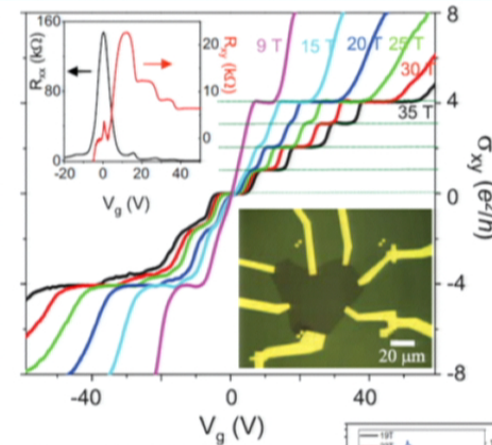
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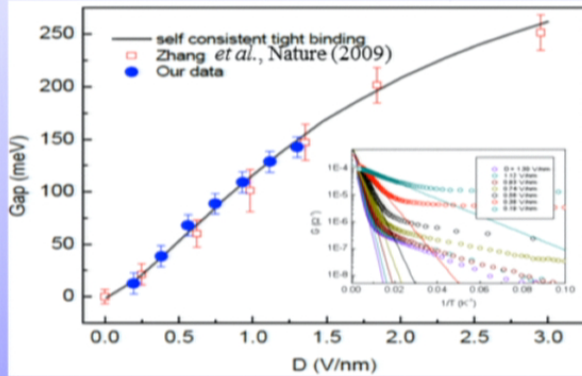
Broken Symmetry QHE in Zero Energy LL in Bilayer Graphene:  
magnetically induced insulating state



Magnetic Field Induced Insulating State at the Charge Neutrality

Zhao et al PRL (2010);  
See also Feldman et al (2009)

Electric Field Induced Layer Polarization Induced Gaps



**>200 meV transport gaps may be possible!**

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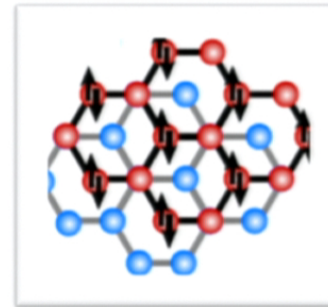
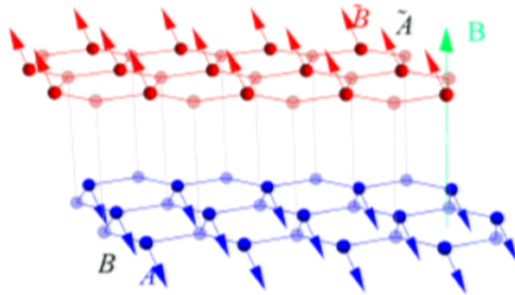
# Nature of Magnetically Induced Insulating States

## Partial list of references

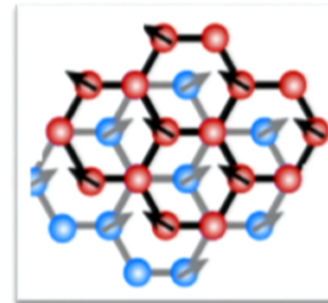
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# Nature of Magnetically Induced Insulating States

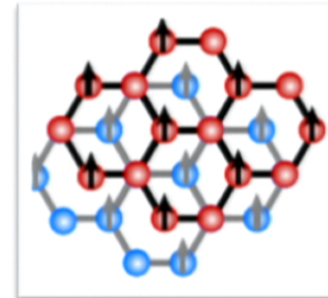
Half filled Bilayer Graphene Degree of freedom:  
**Layer-sublattice Pseudospin / Spin**



**Layer Polarized**



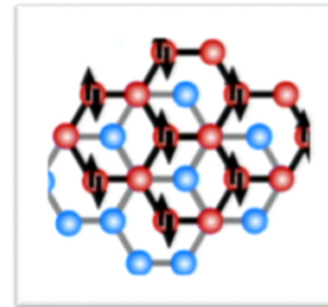
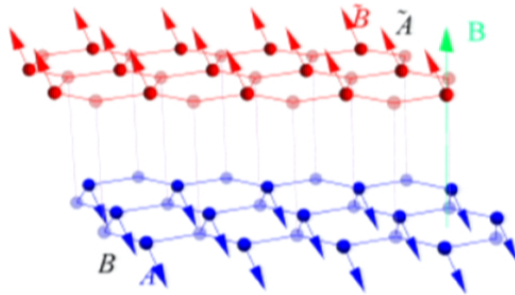
**(canted)  
Anti Ferromagnet**



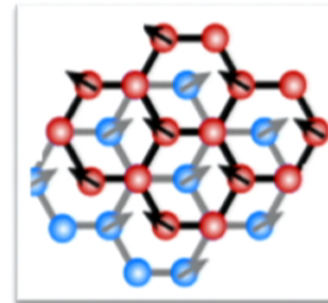
**Ferromagnet**

# Nature of Magnetically Induced Insulating States

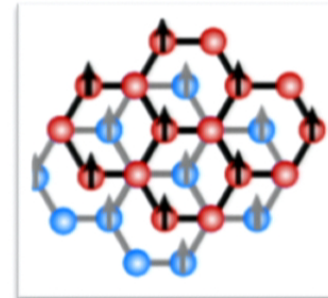
Half filled Bilayer Graphene Degree of freedom:  
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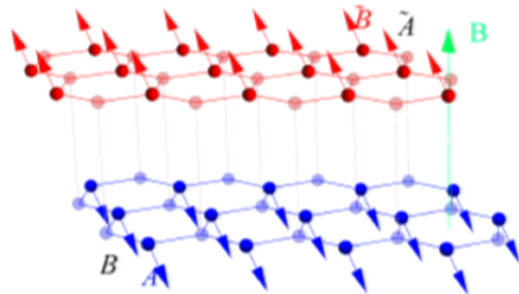
**(canted)  
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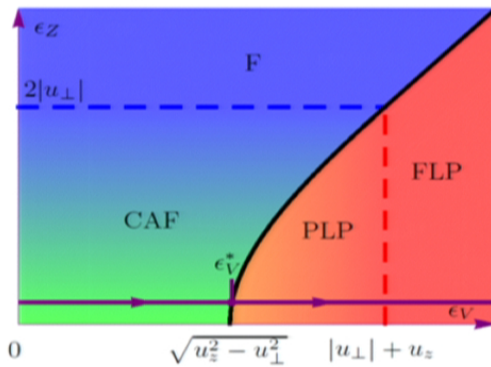
**Ferromagnet**

# Nature of Magnetically Induced Insulating States

Half filled Bilayer Graphene Degree of freedom:  
**Layer-sublattice Pseudospin / Spin**

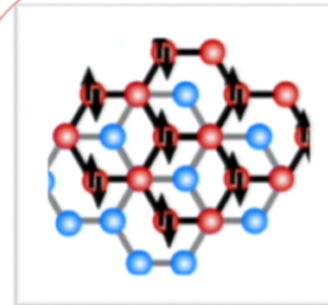


**Phase Diagram**

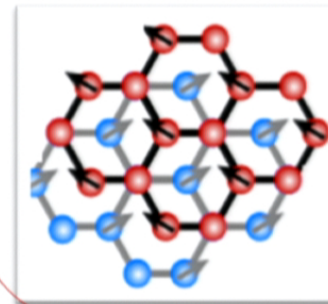


Kharitonov, arXiv:1105.5386 (2011)

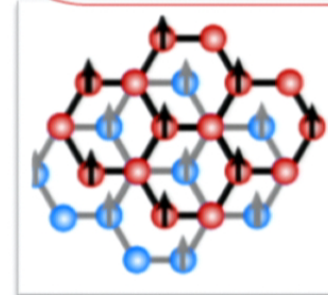
**Spin-upolarized**



**Layer Polarized**



**(canted)  
 Anti Ferromagnet**

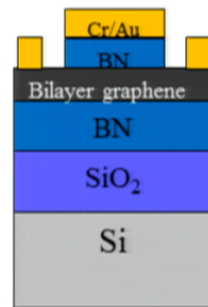
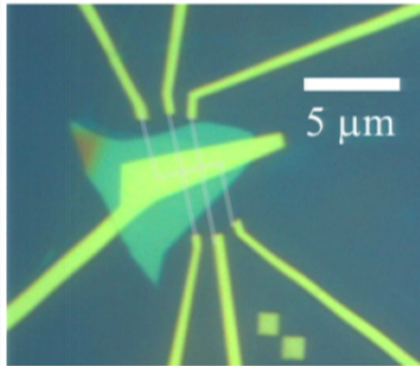


**Ferromagnet**



# Controlling Spin and Pseudo Spin in Bilayer

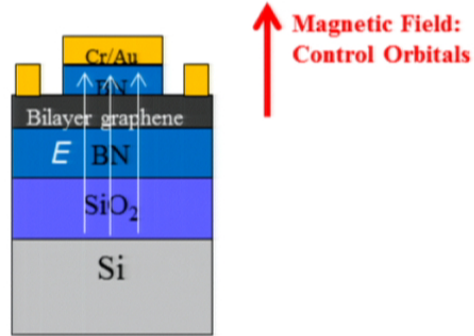
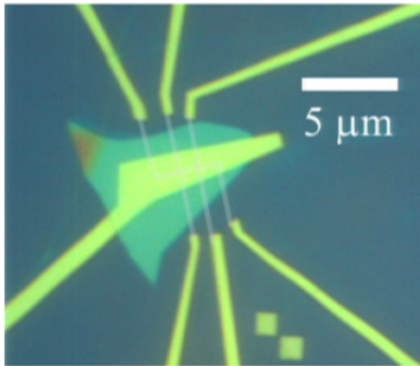
Doubly gated Bilayer graphene with BN dielectric



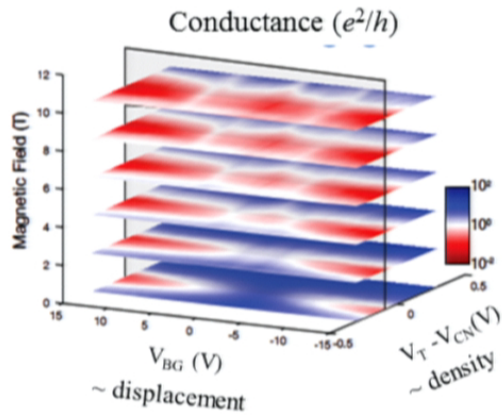
*Maier et al. (2012)*

# Controlling Spin and Pseudo Spin in Bilayer

Doubly gated Bilayer graphene with BN dielectric



Electric field:  
controlling layer pseudo spin

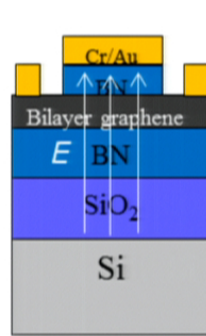
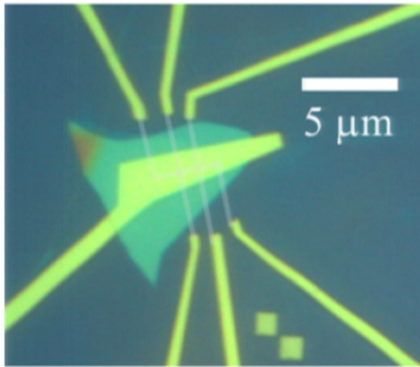


$T = 1.5$  K

Maher *et al.* (2012)

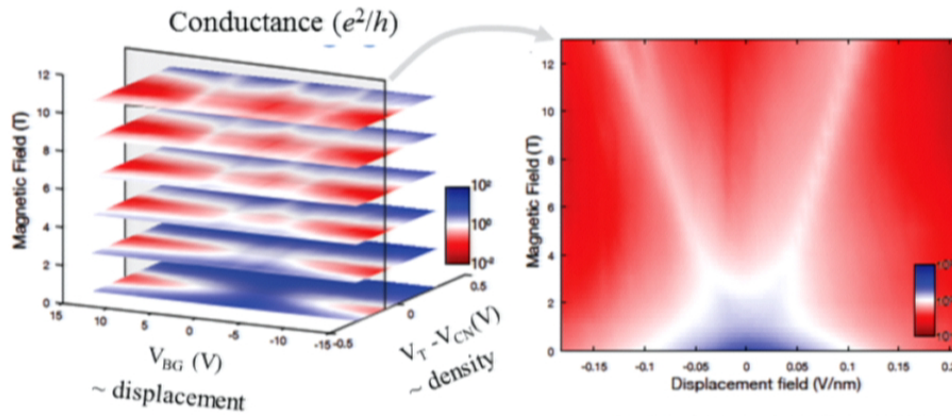
# Controlling Spin and Pseudo Spin in Bilayer

Doubly gated Bilayer graphene with BN dielectric



Magnetic Field:  
Control Orbitals

Electric field:  
controlling layer pseudo spin



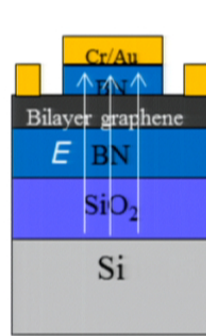
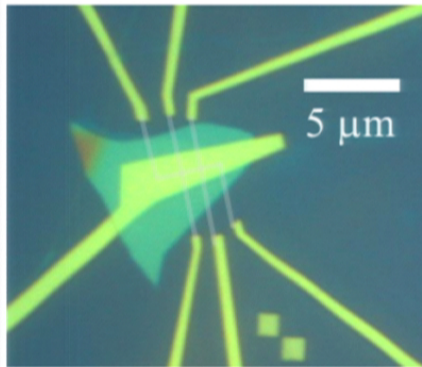
$T = 1.5$  K

See also Previous work: Weitz et al (2010)  
and also work from Manchester & UCR

Maher *et al.* (2012)

# Controlling Spin and Pseudo Spin in Bilayer

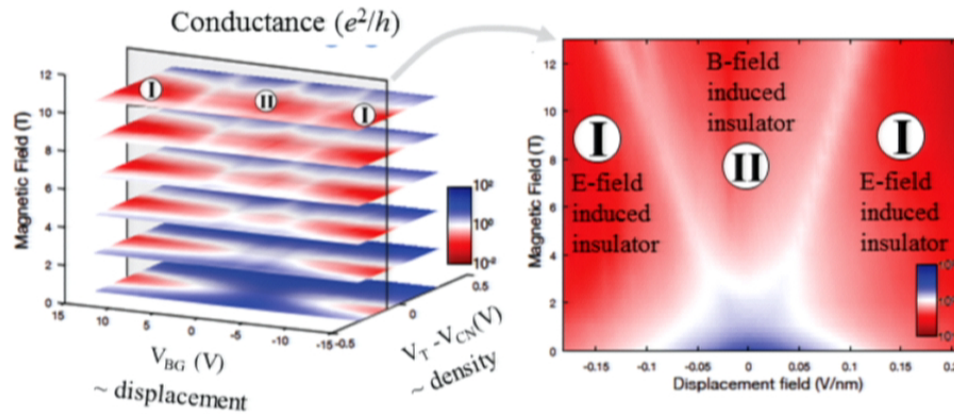
Doubly gated Bilayer graphene with BN dielectric



Magnetic Field:  
Control Orbitals

Parallel magnetic field:  
controlling spin

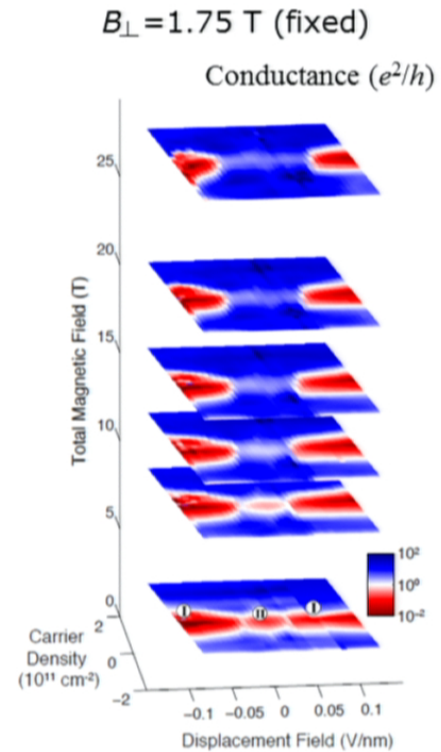
Electric field:  
controlling layer pseudo spin



$T = 1.5 \text{ K}$

See also Previous work: Weitz et al (2010)  
and also work from Manchester & UCR

## Parallel Magnetic Field

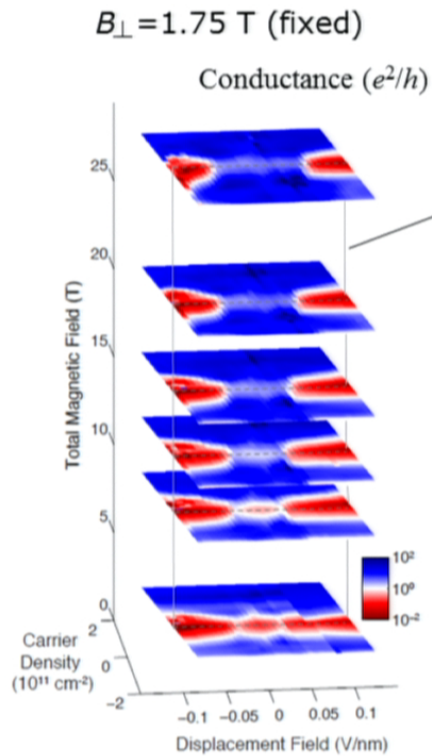


$T = 0.35 \text{ K}$

Maher *et al.* (2012)

# Transition from Insulating State to Metallic State

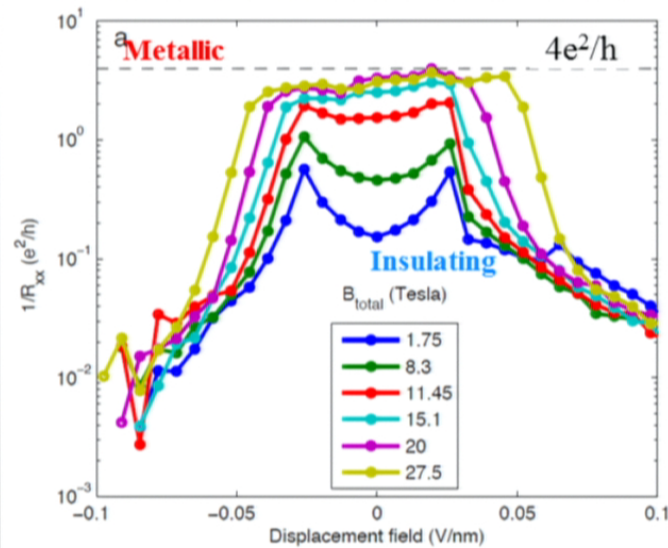
## Metal to Insulator Transition



$T = 0.35$  K

Maher *et al.* (2012)

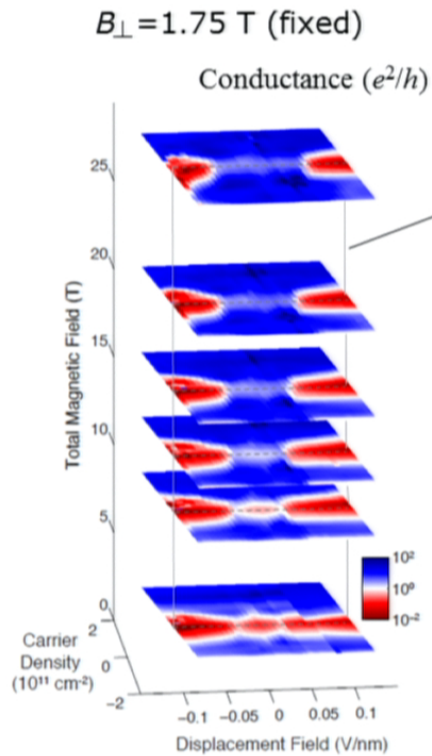
## Conductance versus Displacement



- Low  $B_{\text{tot}}$  : insulating
- High  $B_{\text{tot}}$  : metallic ( $\sim 4e^2/h$ )

# Transition from Insulating State to Metallic State

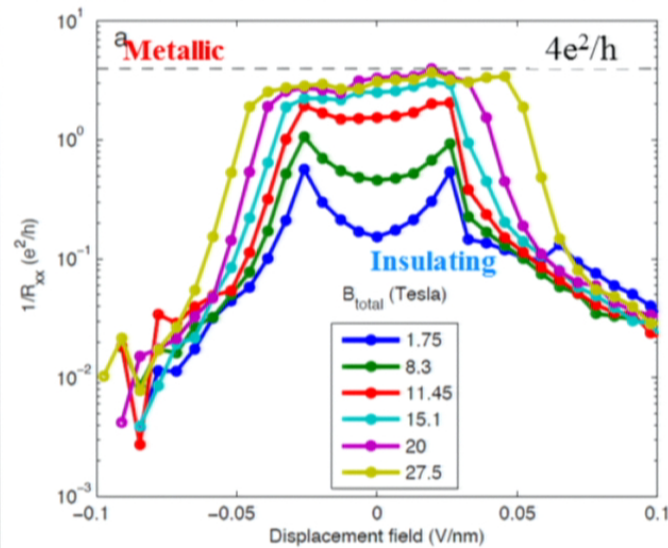
## Metal to Insulator Transition



$T = 0.35$  K

Maher *et al.* (2012)

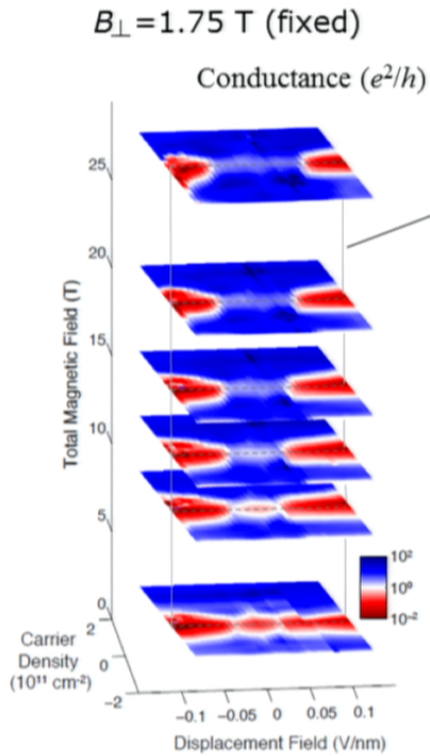
## Conductance versus Displacement



- Low  $B_{\text{tot}}$  : insulating
- High  $B_{\text{tot}}$  : metallic ( $\sim 4e^2/h$ )

# Transition from Insulating State to Metallic State

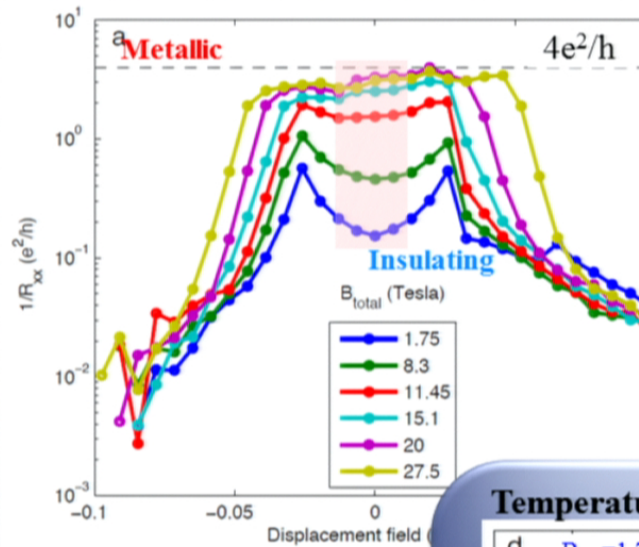
## Metal to Insulator Transition



$T = 0.35$  K

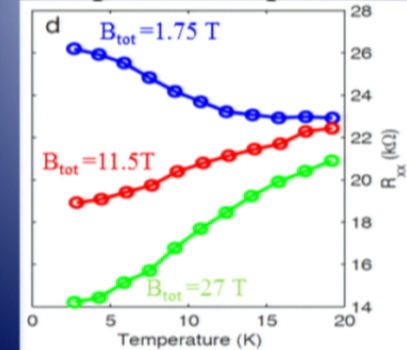
Maher *et al.* (2012)

## Conductance versus Displacement



- Low  $B_{tot}$  : insulating
- High  $B_{tot}$  : metallic ( $\sim 4e^2/h$ )

## Temperature Dependence

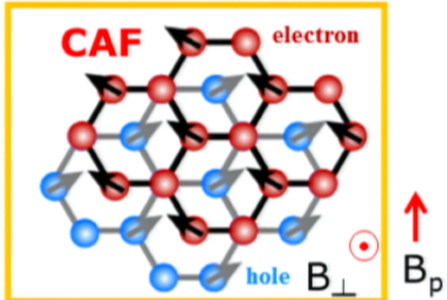


# Insulator to Metallic: Edge State Configuration

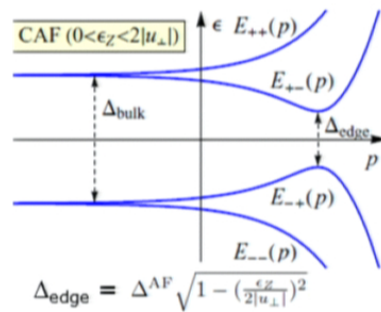
M. Kharitonov (2012) and others

$$\varepsilon_z = -\mu_B B_{tot}$$

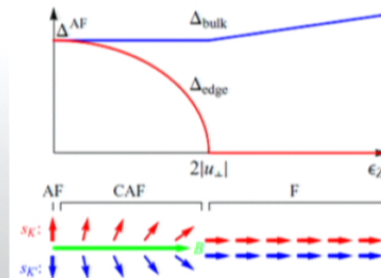
## Canted Anti-Ferromagnetic



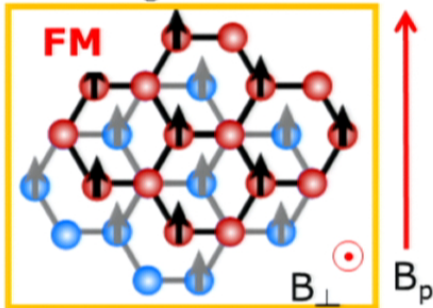
## Edge Configuration



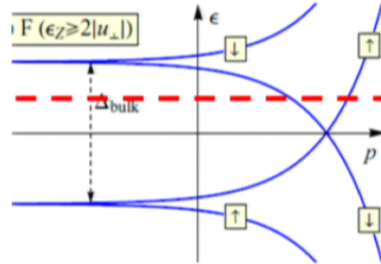
## Energy Gap



## Ferromagnetic

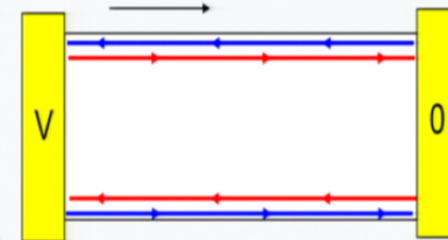


## Edge Configuration



## Counter Propagating Spin polarized edge states

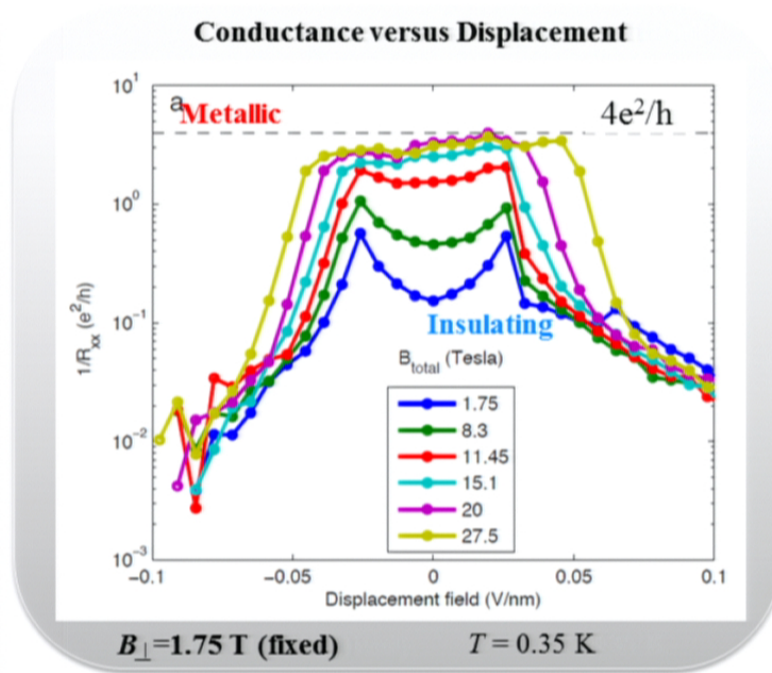
$$I = 4(e^2/h) V$$



Experimental results are reported by Hunt *et al.* *Science* (2013) in SLG

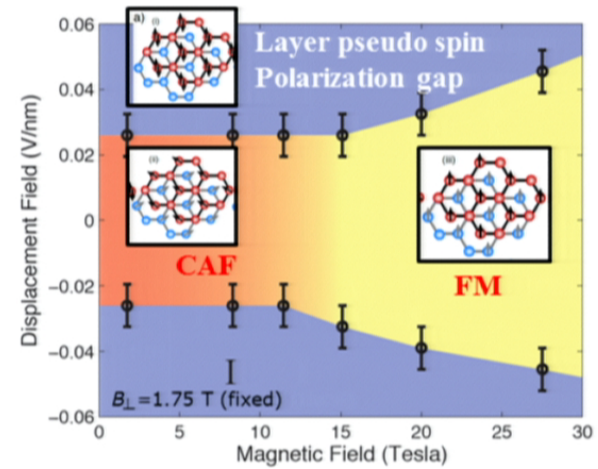
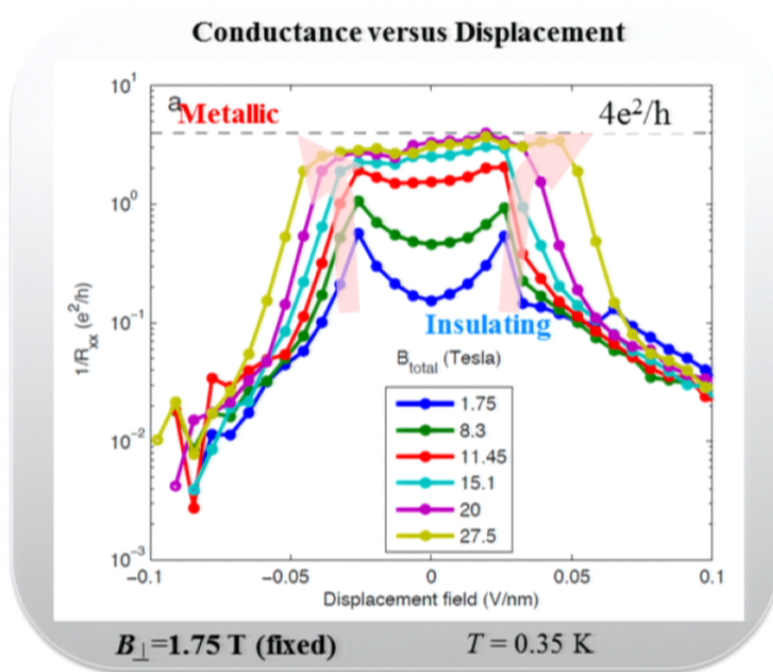


# Energy Scale of Phase Boundaries



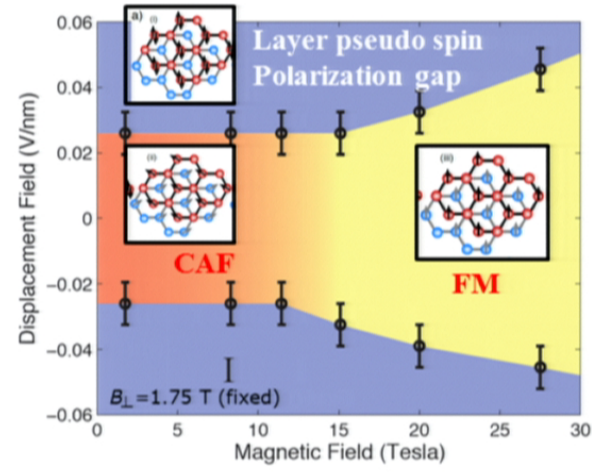
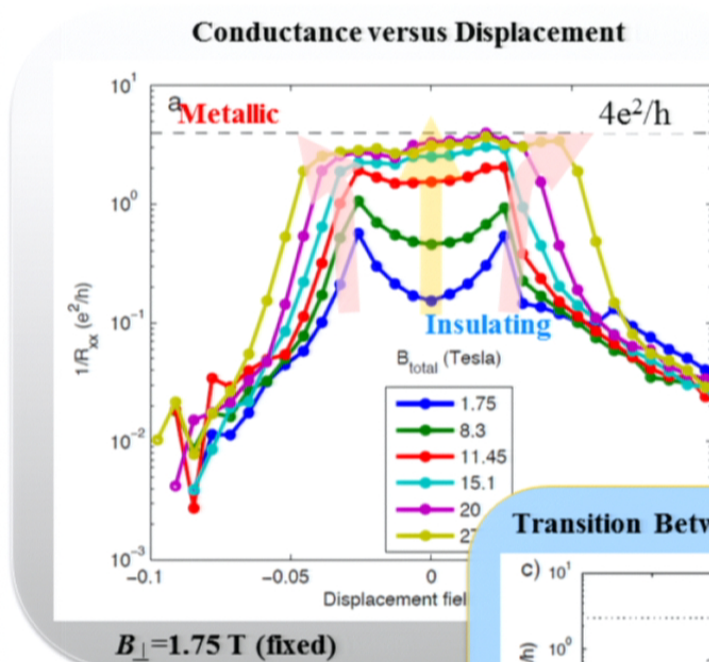
Maher *et al.* (2013)

# Energy Scale of Phase Boundaries

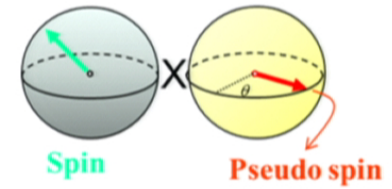
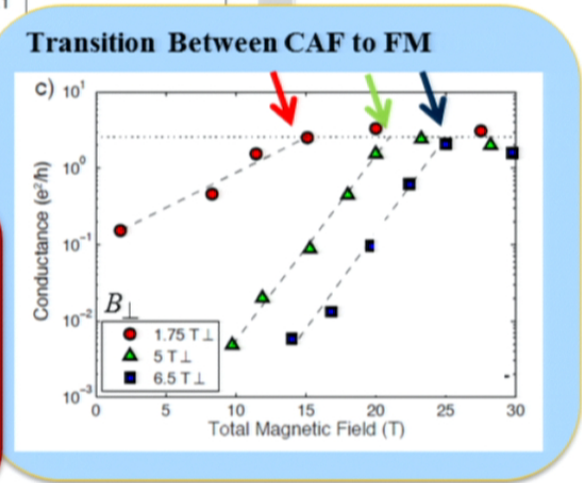
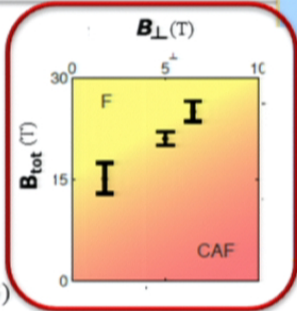


Maher *et al.* (2013)

# Energy Scale of Phase Boundaries



**B<sub>⊥</sub> = 1.75 T (fixed)**



Maher *et al.* (2013)

## Thermodynamic properties of an interacting two-dimensional electron gas in a strong magnetic field

A. H. MacDonald, H. C. A. Oji, and K. L. Liu\*

*National Research Council of Canada, Ottawa, Canada K1A 0R6*

(Received 29 January 1986)

Hartree Fock calculation of

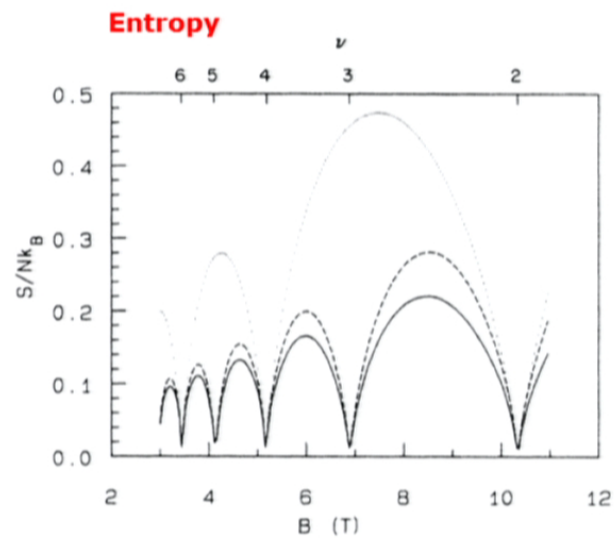


FIG. 2. Entropy versus magnetic field for ideal electron gas (dotted curve) interacting electron gas without disorder (dashed curve), and with disorder ( $\Gamma[\text{meV}] = 0.34\sqrt{B[\text{T}]}$ ).  $n = 5.0 \times 10^{11} \text{ cm}^{-2}$  and  $T = 4.2 \text{ K}$ . The main effect of interaction is to split the peak values through the gigantic exchange enhancement as discussed in the text.

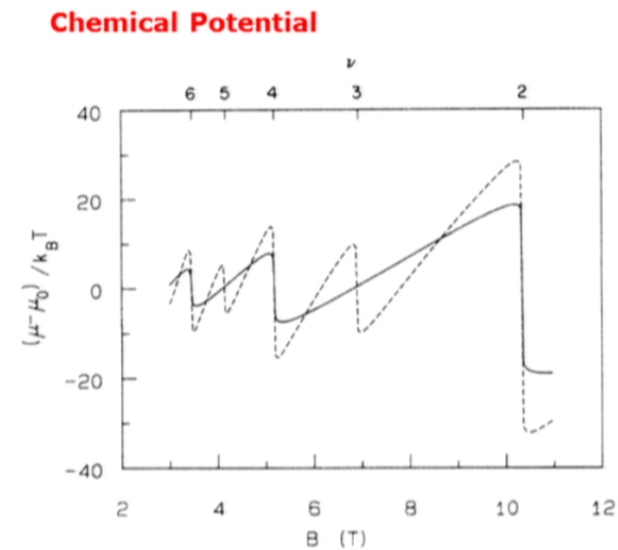
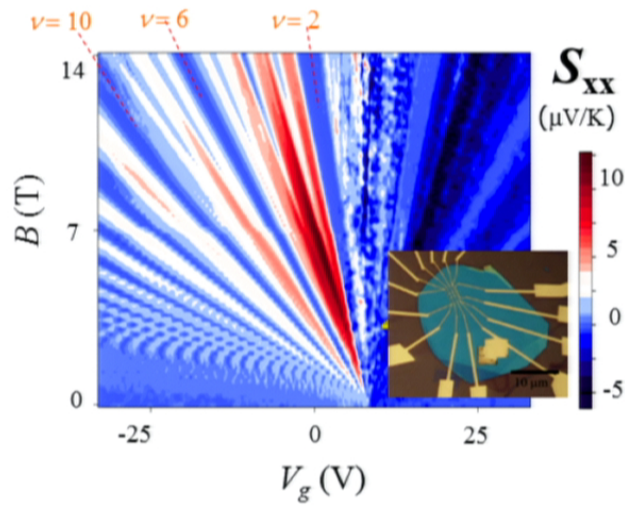
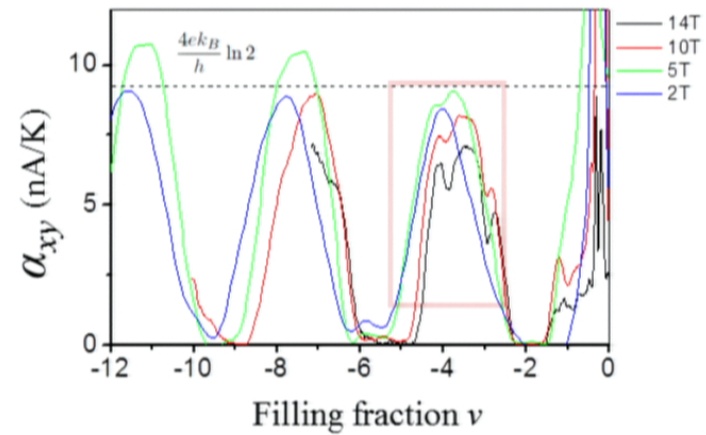


FIG. 3. Shift of the chemical potential from the zero-field value  $\mu_0$ .  $\Gamma[\text{meV}] = 0.34\sqrt{B[\text{T}]}$ ,  $n = 5 \times 10^{11} \text{ cm}^{-2}$ , and  $T = 4.2 \text{ K}$ . The dashed curve shows the effect of interaction. The discontinuities in  $\mu$  which occur at integer filling factors is a measure of the gaps in the quasiparticle energy spectrum (see text).

# Magneto Thermoelectric Power in Broken Symmetry Quantum Hall States



Peltier Conductivity  $\alpha_{ij} = \sum_k S_{ik} \sigma_{kj}$

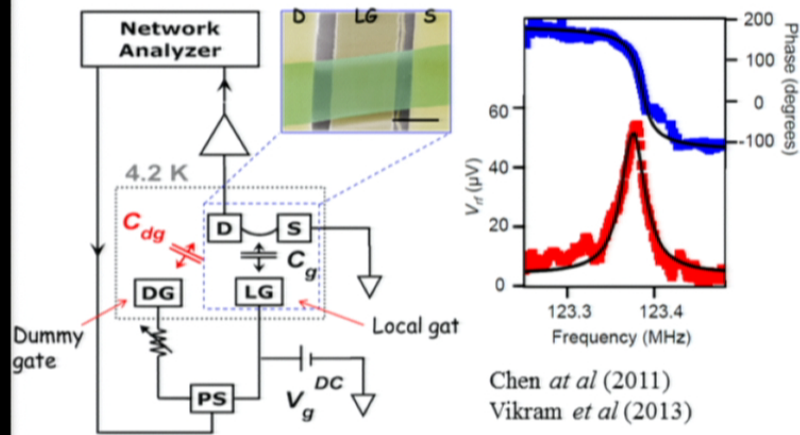


F Ghahari et al (2013)

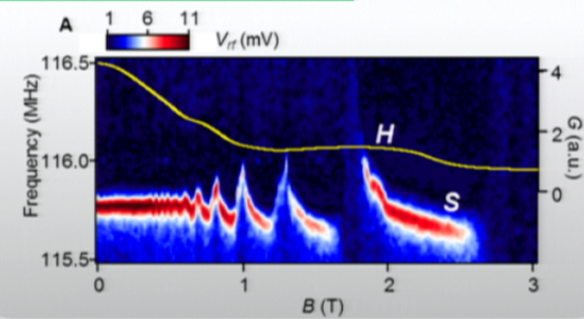
# Mechanical Electromagnetic Measurements

(Deshpande et al, with Hone group collaboration)

## Electromechanical Devices: suspended graphene

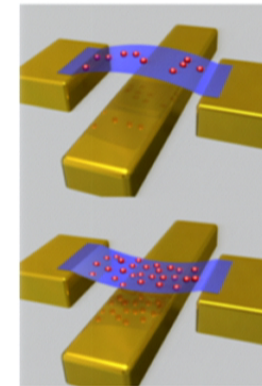


## Magnetization measurement



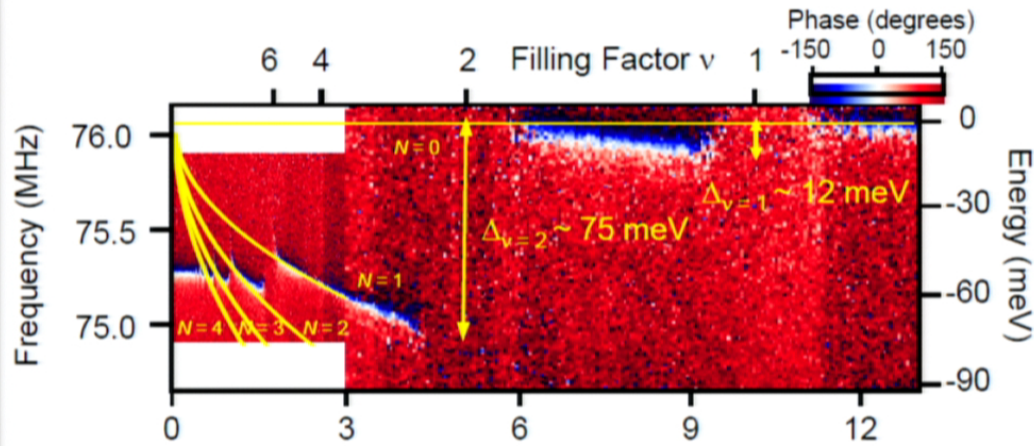
$$F_{mag} = -\frac{\partial U}{\partial z} = -\frac{\partial U(n)}{\partial n} \frac{\partial n}{\partial z} = -\mu V_g \frac{\partial C_g}{\partial z}$$

Direct experimental access of chemical potential!



# Chemical Potential Measurement of Broken Symmetry QH States

(Deshpande *et al* (2013), with Hone group collaboration)

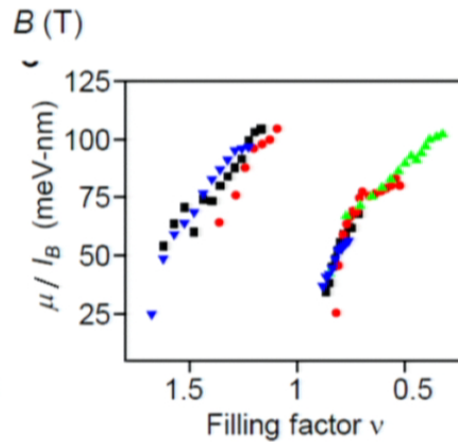
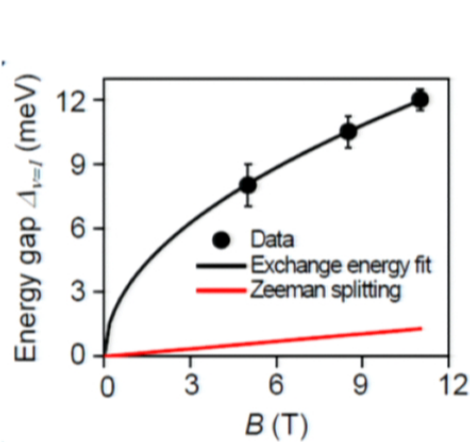


- Broken Symmetry State ( $\nu=1$ ) are observed
- $\Delta_{\nu=1}$  is measured to be scale as  $B^{1/2}$

- Partially filled broken symmetry LL provides chemical potential universal scaling of

$$\mu = \hbar v_F / \ell_B F(\nu)$$

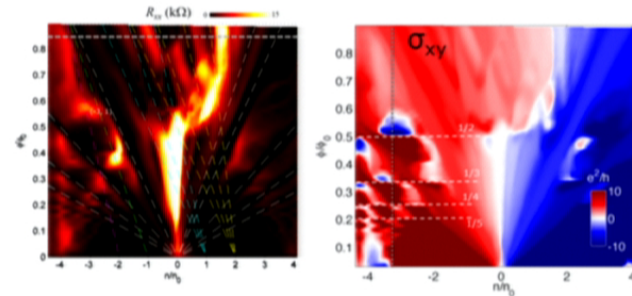
Kravchenko *et al.* PRB (1990)



# Summary

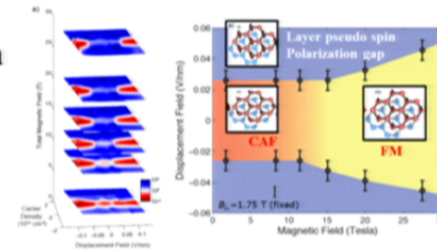
- **Graphene/ hBN heterostructures:**  
**Observation of Hofstadter's Butterfly**

Dean *et al*, *Nature* (2013);  
similar results were reported by Gorbachev *et al*. *Nature* (2013)  
and Hunt *et al*. *Science* (2013)



- **Bilayer graphene:**  
**e-e interaction and spin and pseudo-spin SU(4) phase transition**

Maher *et al*, *Nature Phys.* (2013);  
similar results were reported by Hunt *et al*. *Science* (2013)



- **Magneto-thermoelectric effect and Magneto-electromechanical effect:**

Reduction of transport entropy in broken symmetry integer quantum Hall effect  
Universal scaling of chemical potential in QHF

