Title: Divergences in Spinfoam Quantum Gravity

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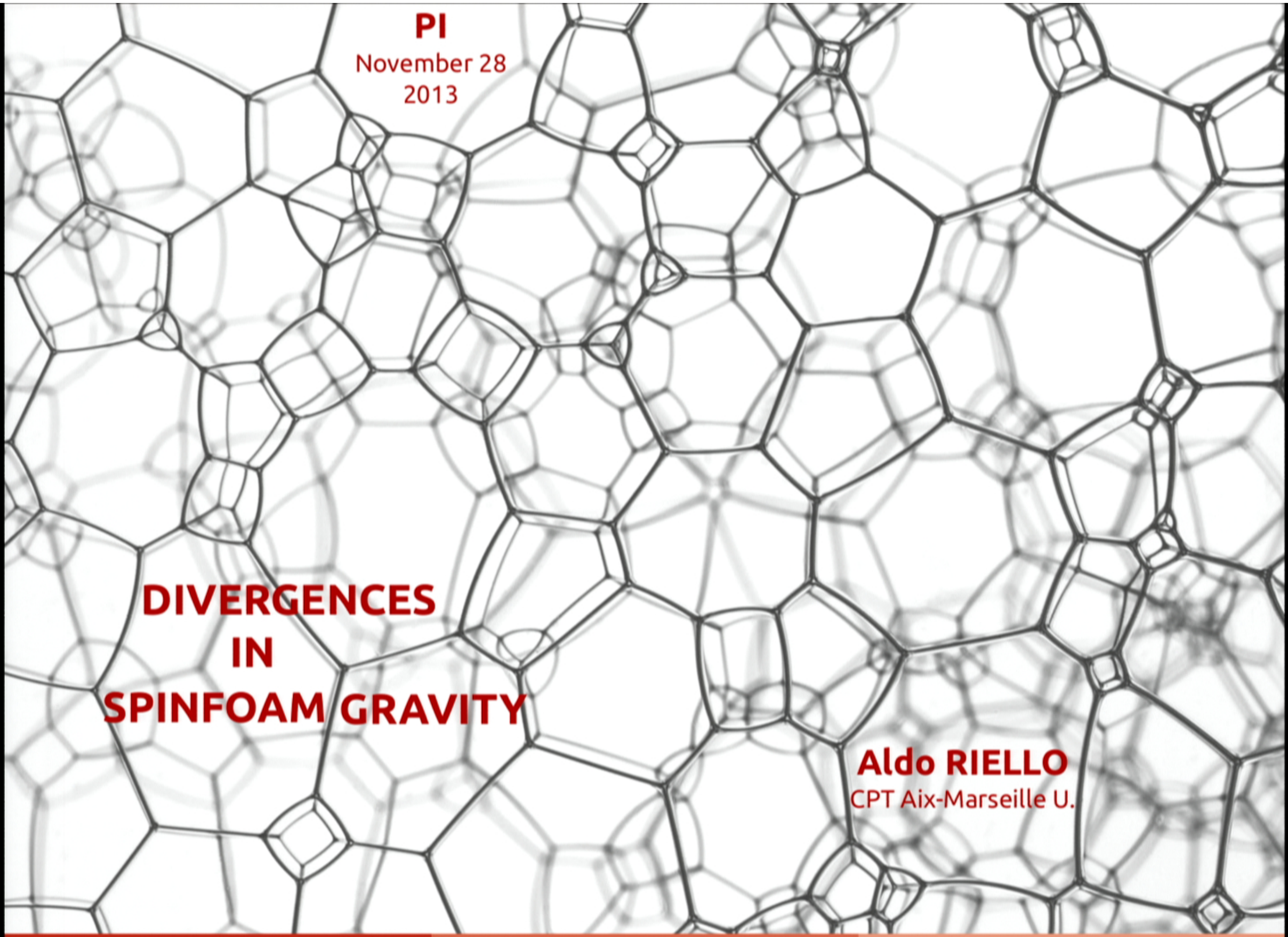
Abstract: The most relevant evidences in favour of the Lorentzian EPRL-FK spinfoam model come from its capibility of reproducing the expected semiclassical limit in the large spin regime. The main examples of this are the large spin limit of the vertex amplitude, later extended to arbitrary triangulations, and that of the spinfoam graviton propagator, which was calculated on the simplest possible two complex. These results are very promising. Nonetheless, their relevance may be endangered by the effects associated to radiative corrections. In this seminar, I will focus on the role played by the simplest diverging graph, the so called 'melon graph', which is known to play a fundamental role in tensorial group field theories. In particular, I will discuss its most divergent part and its geometrical interpretation. I will finally comment on the result, with particular attention to its physical consequences, especially in relation with the semiclassical limit of the spinfoam graviton propagator.



PI
November 28
2013

DIVERGENCES IN SPINFOAM GRAVITY

Aldo RIELLO
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Plan

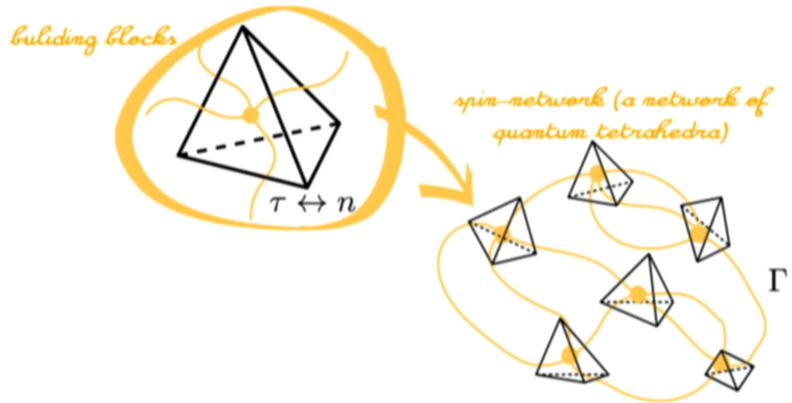
- *Spinfoam gravity overview* [Part 0]
- *Why studying spinfoam divergences?* [Part I]
- *The melon graph: why, what, and therefore?* [Part II]
  [AR, PRD88 (2013), arXiv:1302.1781]
- *Not just the amplitude: correlations* [Part III]
  [AR, arXiv:1310.2174]

Quantum geometry

Building blocks: quantum states of geometry named **spin networks** $\Psi_{\Gamma j_l \iota_n}$

Labelled by:

- ▷ an abstract graph Γ (say 4-valent)
- ▷ an $SU(2)$ spin j_l at every link
- ▷ an $SU(2)$ (4-valent) invariant tensor ι_n at every node (intertwiner)



Nodes are dual to quantum tetrahedra (of a twisted geometry), whose volume is fixed by the intertwiner ι_n, and whose face areas are

$$A_l = 8\pi G\hbar\gamma\sqrt{j_l(j_l+1)}$$

[The group $SU(2)$ encodes the symmetries of the quantized space]

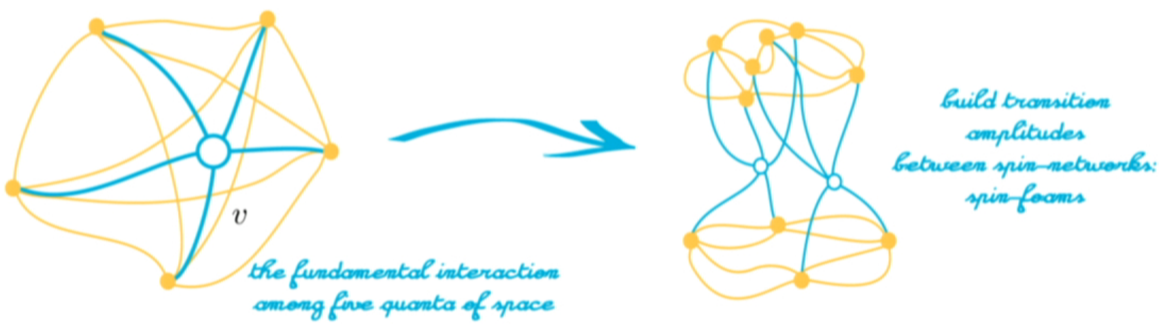
Quantum processes

Spinfoams

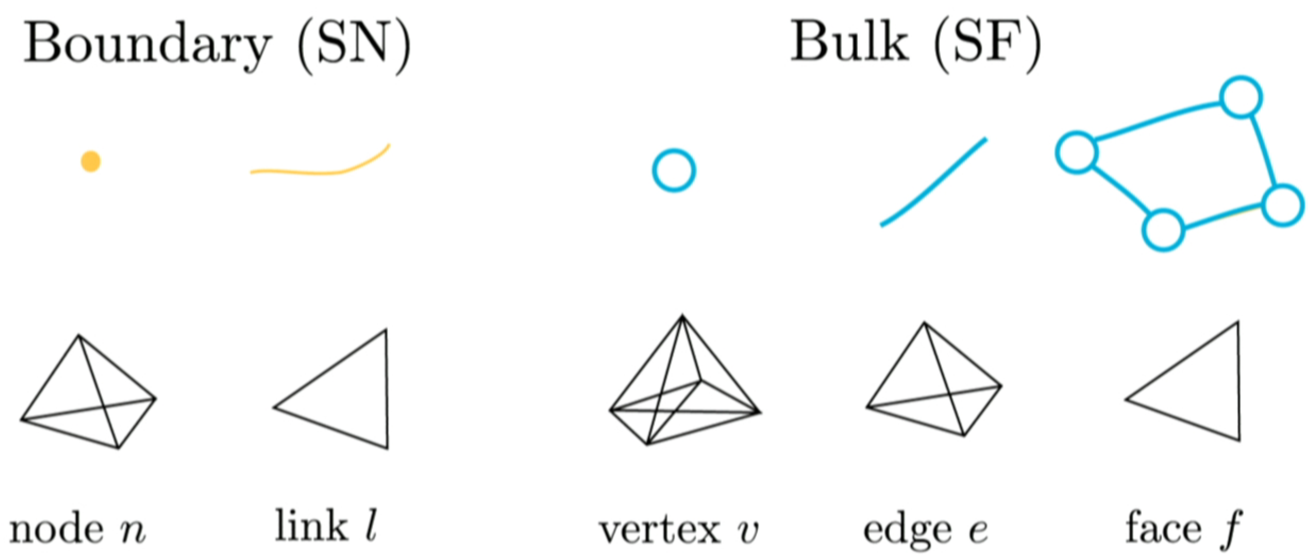
Quantum gravitational processes between spin network states

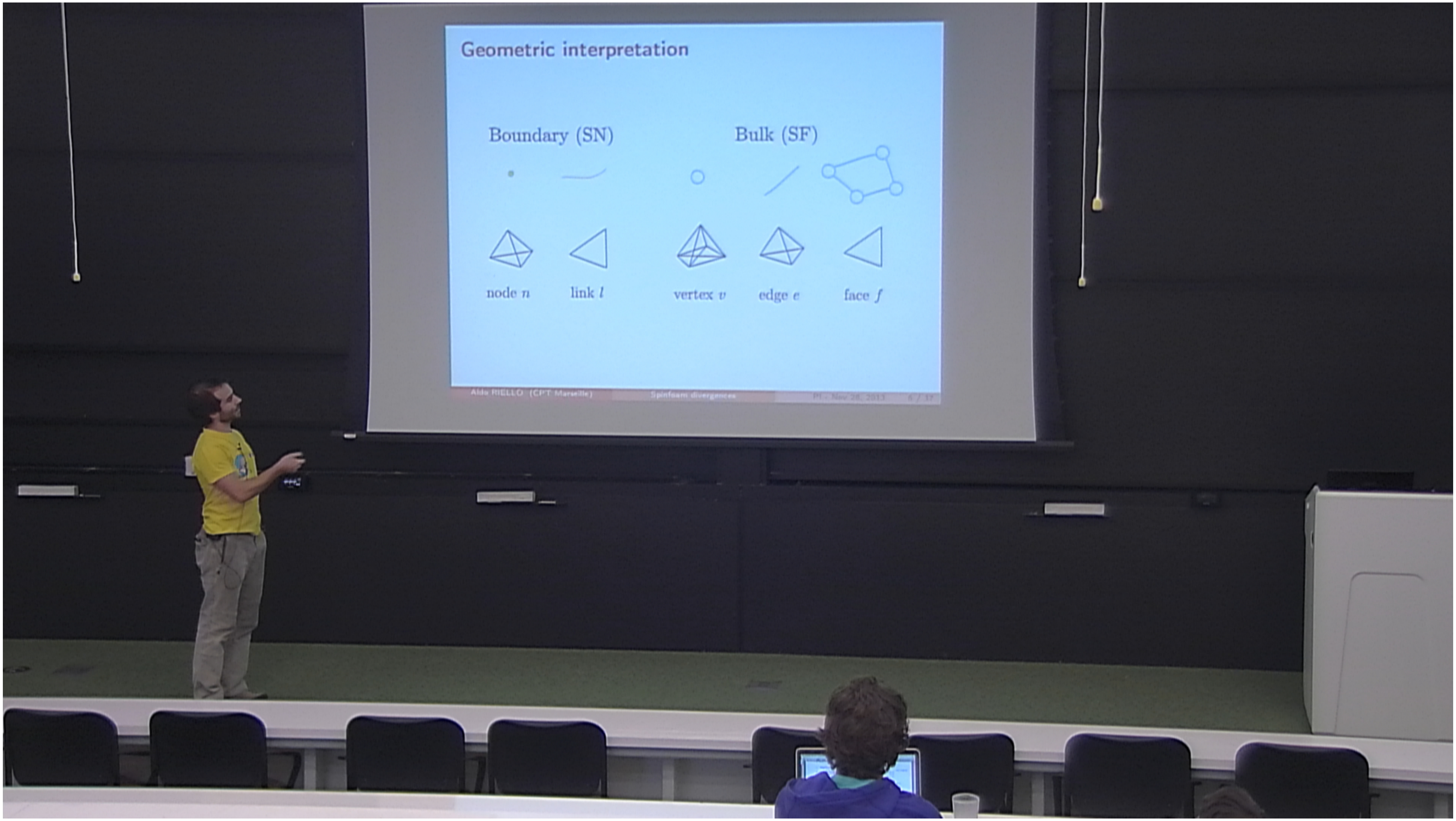
As Feynman graphs, they are built out of fundamental units:

- ▷ interaction vertices among quanta of space
- ▷ gluing of interaction vertices via their boundaries



Geometric interpretation




Geometric interpretation
Boundary (SN)
Bulk (SF)
node n
link l
vertex v
edge e
face f
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Spinfoam divergences
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Spinfoam amplitude

Spinfoam model

A prescription to assign an amplitude ($\in \mathbb{C}$) to every such process.

Local spinfoam Ansatz:
the amplitude is built out of local amplitudes for faces f, edges e and vertices v depending only on local representations and intertwiners:

$$\mathscr{A}_{SF}[\Psi] = \sum_{\text{colourings}} \prod_f A_f[j_f] \prod_e A_e[\iota_e, j_{f\ni e}] \prod_v A_v[\iota_{e\ni v}, j_{f\ni v}]$$

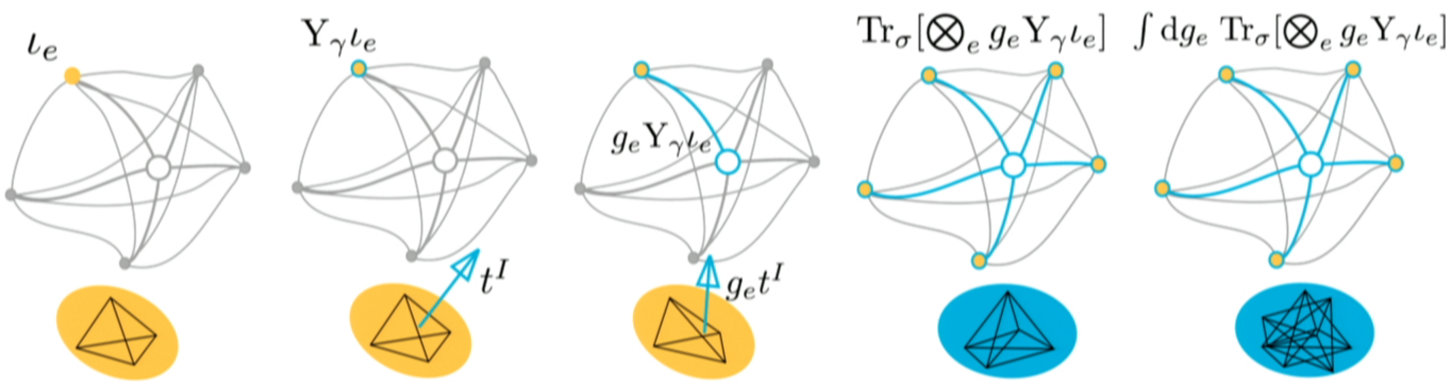
It is expected to be a regularization of the integral over histories:

$$\mathscr{A}_{SF} \sim \int \mathscr{D}g_{\mu\nu}\, e^{iS[g_{\mu\nu}]}$$

EPRL-FK

Four-dimensional Lorentzian spinfoam gravity: the **EPRL-FK model**

Vertex amplitude:



Mathematically:

$$\mathrm{Y}_\gamma : \mathscr{V}^{j}_{SU(2)} \to \mathscr{V}^{(\rho=\gamma j, k=j)}_{SL(2,\mathbb{C})}, \quad |j;m\rangle \mapsto |\gamma j, j; j, m\rangle$$

[Engle, Pereira, Rovelli, Livine, Freidel, Krasnov, Speziale]

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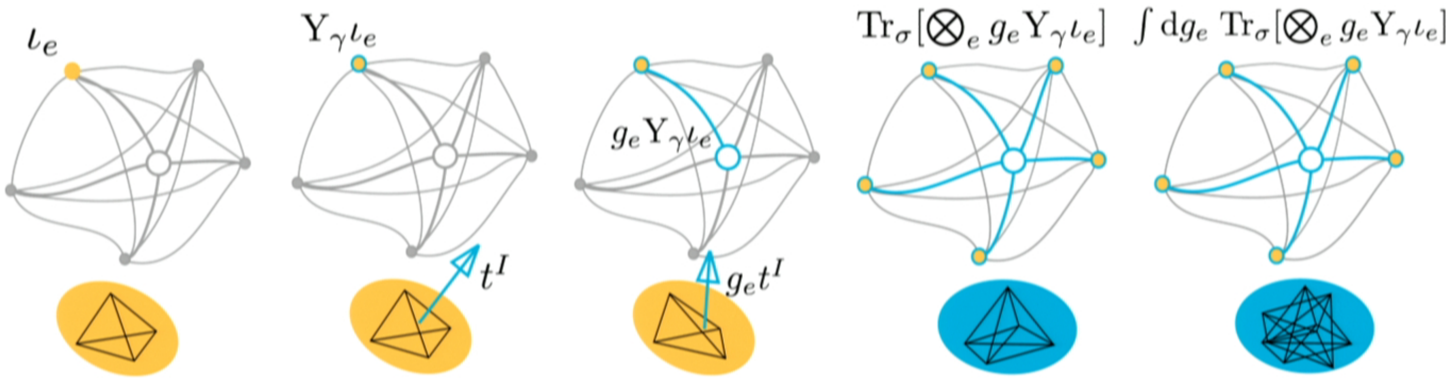

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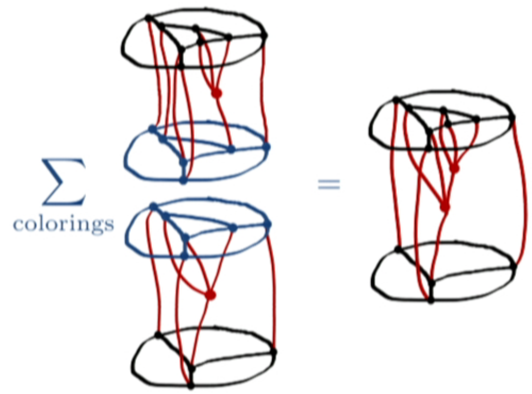
Choice of face and edge amplitudes

In previous definition of EPRL-FK:

$$A_f = (2j+1) \quad \text{and} \quad A_e = 1$$

▷ this face amplitude is selected by a criterion of composition of spinfoams amplitudes. [Bianchi, Regoli, Rovelli] However, the following analysis can be straightforwardly applied to any face weight



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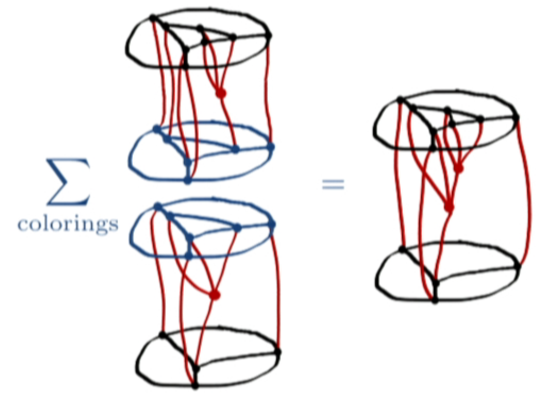

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Indications of viability

Most convincing indications of viability come from **semiclassics**:

$$A_v[\text{coherent state}] \quad \xrightarrow{j\gg 1} \quad N \times \left(e^{iS_{\text{Regge}}} + e^{-iS_{\text{Regge}}} \right),$$

where

$$S_{\text{Regge}} = \sum_{\text{triangles}} A_t \Theta_t$$

[Freidel,Conrady;Barrett,Dowdall,Gomes,Hellmann,Fairbairn,Pereira;Han,Zhang,...]

is a well-defined discretization of Einsten-Hilbert action

$\rightsquigarrow$ This leads (in some proper limit) to the expected graviton propagator [Bianchi,Magliaro,Modesto,Perini,Rovelli,Speziale...]

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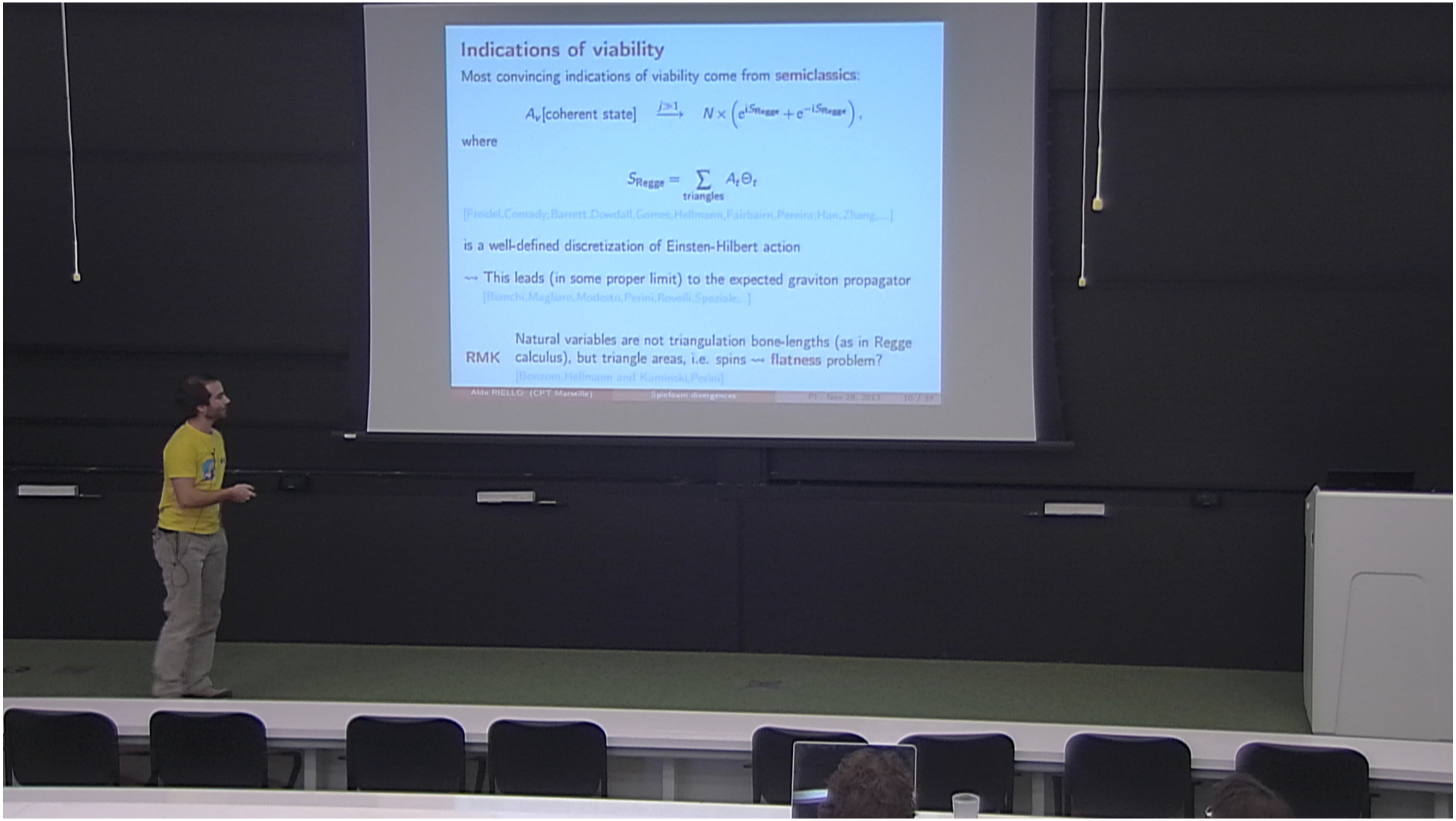

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Spinfoam divergences

Where the wild things are

In Feynman diagrams, divergences are associated to loops. This is because, loops involve integrals on unconstrained variables (momenta)

* * *

Every spinfoam local amplitude A_f, A_e, A_v is finite [Engle and Pereira]

Divergences arise from summations over unconstrained internal colourings

$$\mathscr{A}_{SF}[\Psi] = \sum_{\text{colourings}} \prod_e A_e \prod_f A_f \prod_v A_v$$

Unconstrained colourings are associated to **bubbles**, i.e. to topological structures dual to lower dimensional submanifolds in the triangulation (sides and points) [Perez and Rovelli, Freidel and Louapre, Bonzom and Smerlak, ...]

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Why studying divergences?

(I) Model consistency

Spinfoam continuous limit $\rightsquigarrow$ triangulation refinement $\rightsquigarrow$ bubbles

Therefore, bubbles are a natural piece of the theory. What is their rôle?

- In a renormalization/corarse graining procedure, divergences drive the flow of the theory
  $\rightsquigarrow$ renormalized amplitudes may not have the desired semiclassics
- Divergences may indicate the presence of residual diffeo symmetry (cf. topological 3d gravity) [Freidel,Louapre,...]

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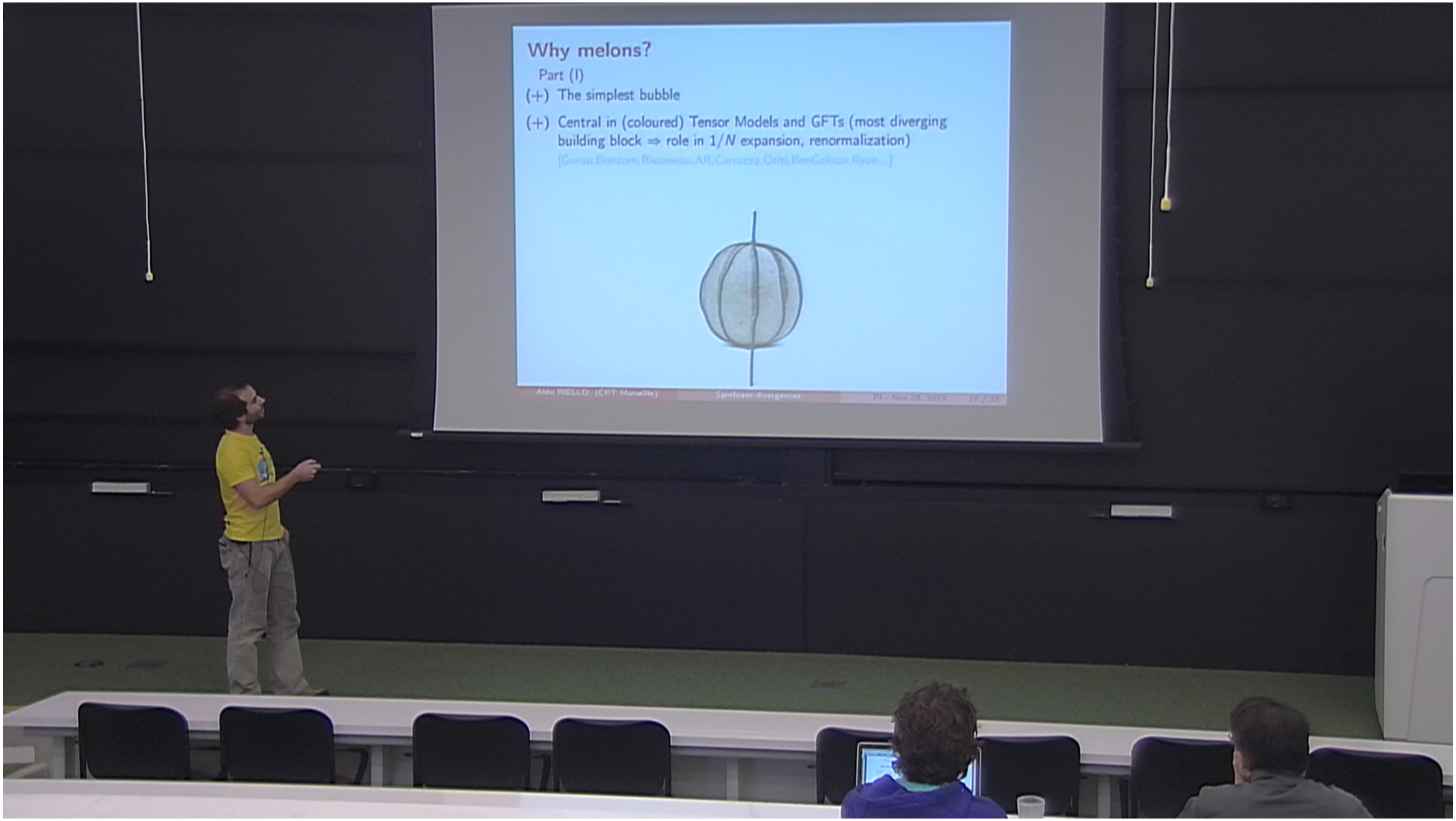
Plan

- *Spinfoam gravity overview* [Part 0]
- *Why studying spinfoam divergences?* [Part I]
  Naturally present. Related to the study of the continuous limit. Even finite models contain a very large number $\sim \lambda^{-1}$. Crucial for the consistency of the model under renormalization/coarse graining. Possibly related with unfixed non-compact gauge symmetries. Possibly related to new physics.
- *The melon graph: why, what, and therefore?* [Part II]
- *Not just the amplitude: correlations* [Part III]


Why melons?
Part (I)
(+) The simplest bubble
(+) Central in (coloured) Tensor Models and GFTs (most diverging building block ⇒ role in $1/N$ expansion, renormalization)
[Gurau,Bonzom,Rivasseau,AR,Carrozza,Oriti,BenGeloun,Ryan...]
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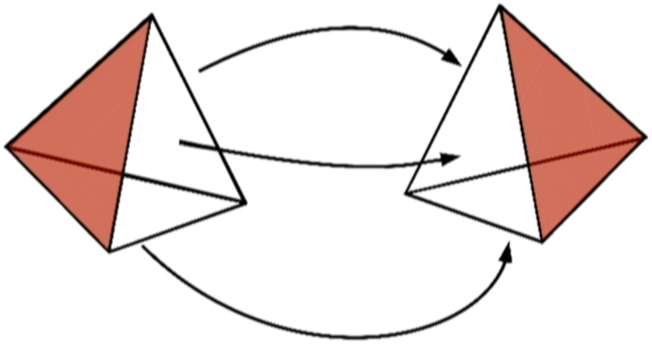
(−) Topological sphere, but dual to a degenerate triangulation (less natural from a geometric perspective)



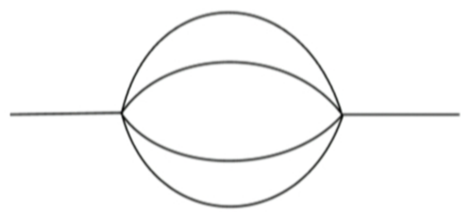
Why melons?

Part (II)

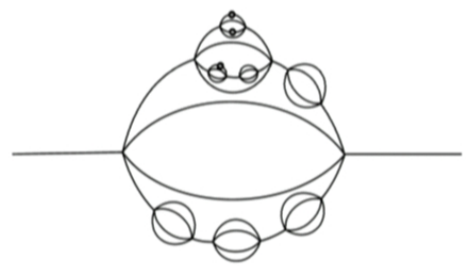
Relevant summations are associated to spinfoam **faces** (spins)
Therefore, at a given order in vertex expansion:

maximize divergence degree $\rightsquigarrow$ maximize # [unconstrained] faces

The melonic family of graphs is the one satisfying this condition
[Actually, additional hypothesis are needed. These are verified in "coloured" models]

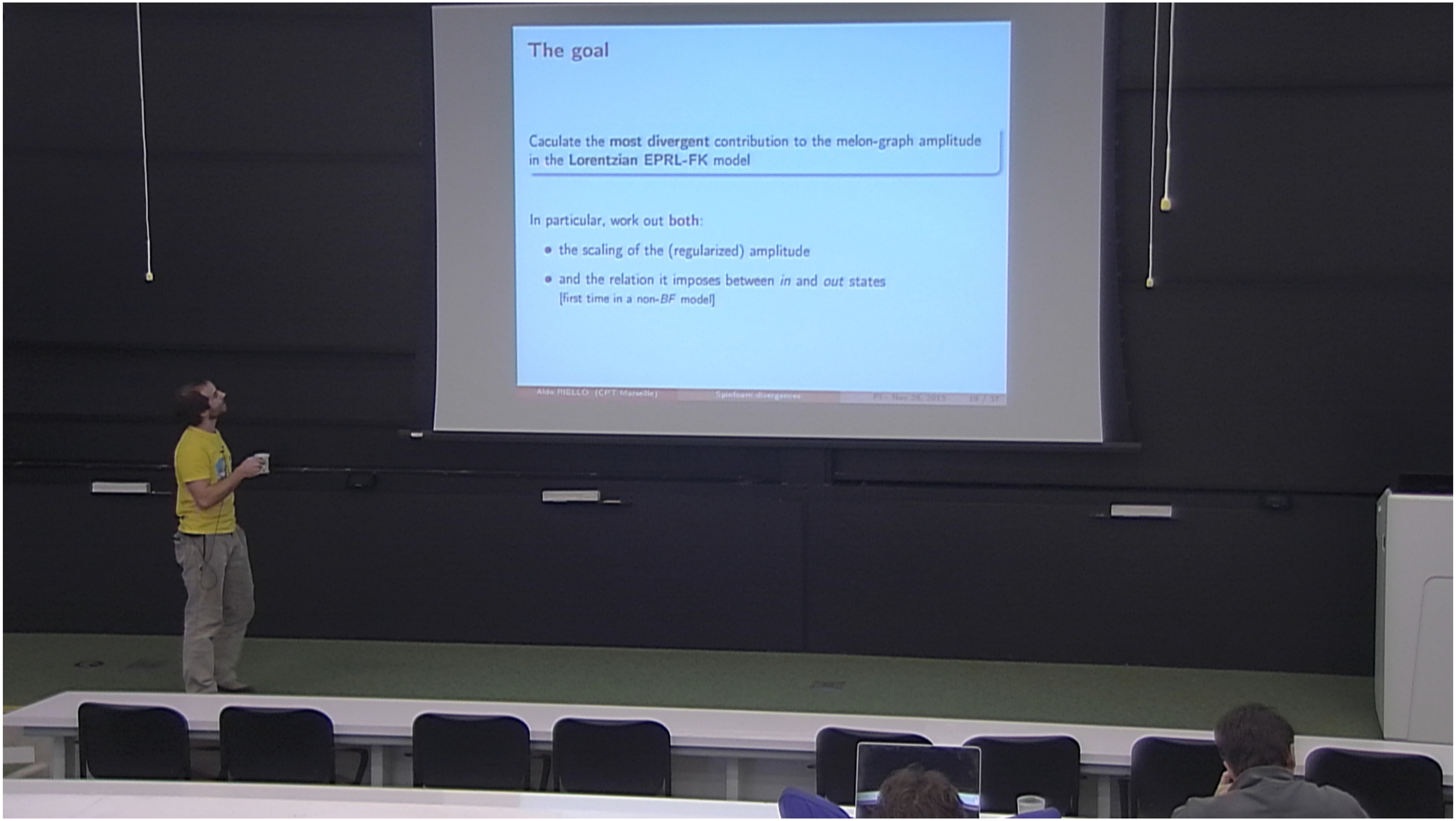


the melon graph



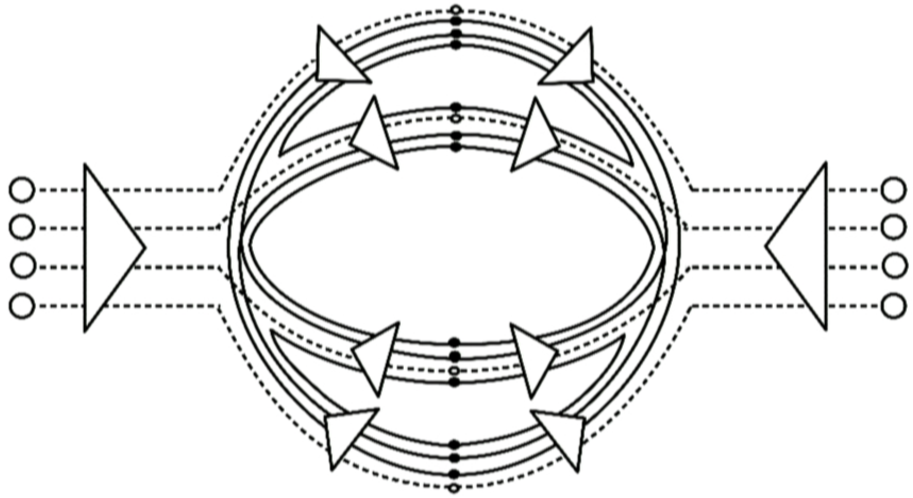
a melonic graph

[first proof in coloured tensor models: Bonzom, Gurau, AR, Rivasseau]



The goal
Caculate the **most divergent** contribution to the melon-graph amplitude in the **Lorentzian EPRL-FK** model
In particular, work out **both**:
• the scaling of the (regularized) amplitude
• and the relation it imposes between *in* and *out* states
[first time in a non-*BF* model]

The structure of the melon



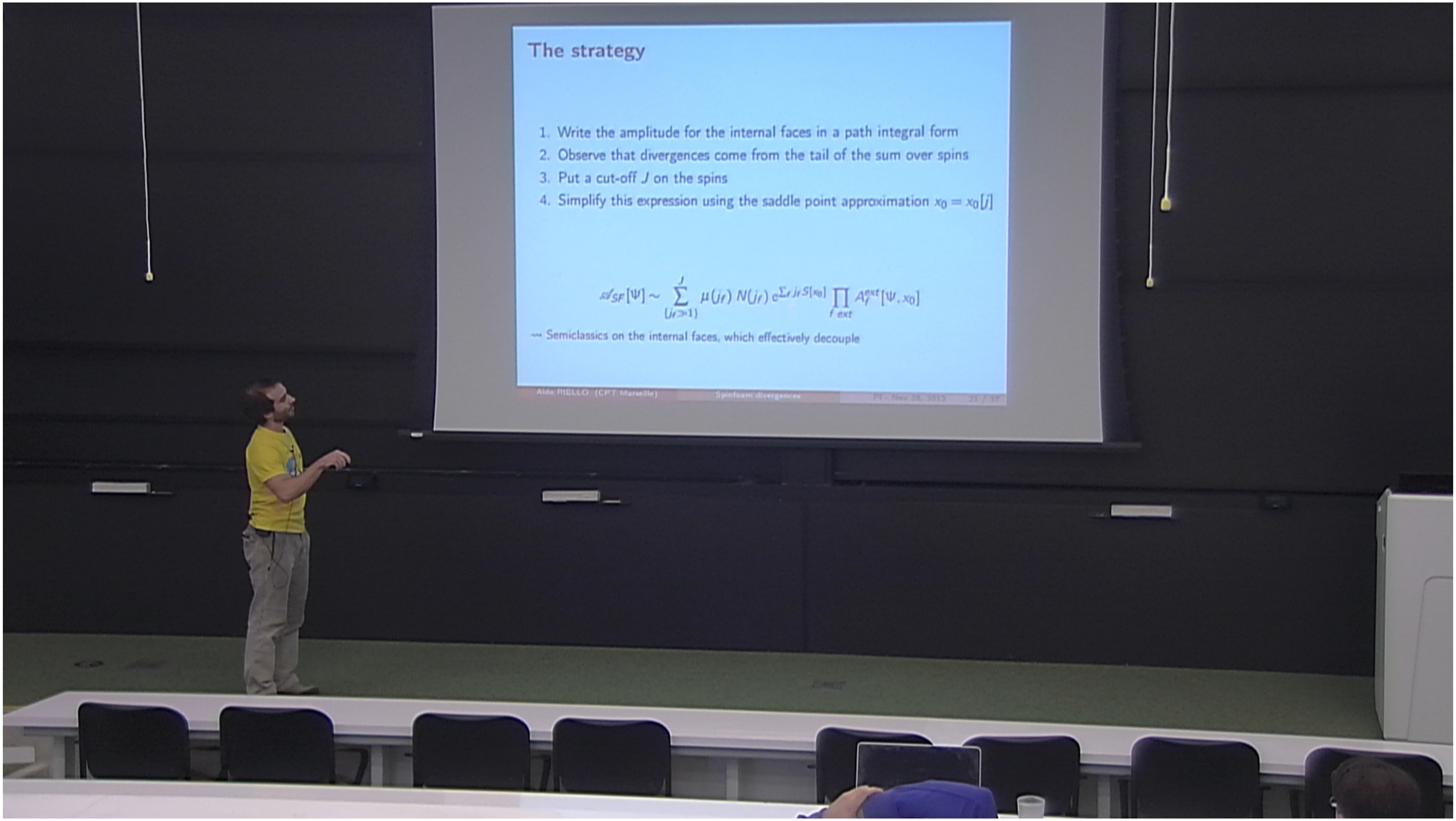
- ▷ dotted lines represent external faces (spin fixed by boundary data)
- ▷ solid lines represent internal faces (spins summed over)
- ▷ triangles represent integrations over $SL(2,\mathbb{C})$ elements

The strategy

1. Write the amplitude for the internal faces in a path integral form
2. Observe that divergences come from the tail of the sum over spins
3. Put a cut-off J on the spins
4. Simplify this expression using the saddle point approximation $x_0 = x_0[j]$

$$\mathscr{A}_{SF}[\Psi] \sim \sum_{\{j_f \gg 1\}}^{J} \mu(j_f)\, N(j_f)\, \mathrm{e}^{\sum_f j_f S[x_0]} \prod_{f\ ext} A_f^{ext}[\Psi, x_0]$$

⇝ Semiclassics on the internal faces, which effectively decouple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4. Simplify this expression using the saddle point approximation $x_0 = x_0[j]$
5. Perform an order of magnitude evaluation of the divergence degree, by taking into account all the symmetries of the action

$$\mathscr{A}_{SF}[\Psi] \sim J^k \prod_{f\ ext} A_f^{ext}[\Psi]$$

Keep the sector dominating the sum over spins: $\sum_f j_f S_0 \equiv 0$ and $A_f^{ext}[\Psi, x_0] \equiv A_f^{ext}[\Psi]$

Approximations

Downsides

- ▷ All spin are let scale together, a priori to maximize divergence, but maybe other effects enter the game
- ▷ No rigorous proof to pick the dominating sector
- ▷ Neglected **degenerate sector** of the sum over spins

  [There are indications that this sector could endanger the result. However, there are difficulties in doing a complete analysis]

Upsides

- ▶ EPRL-FK with cut-off $\rightsquigarrow$ toy model for cosmological constant

$$J \sim \lambda^{-1} \gg 1$$

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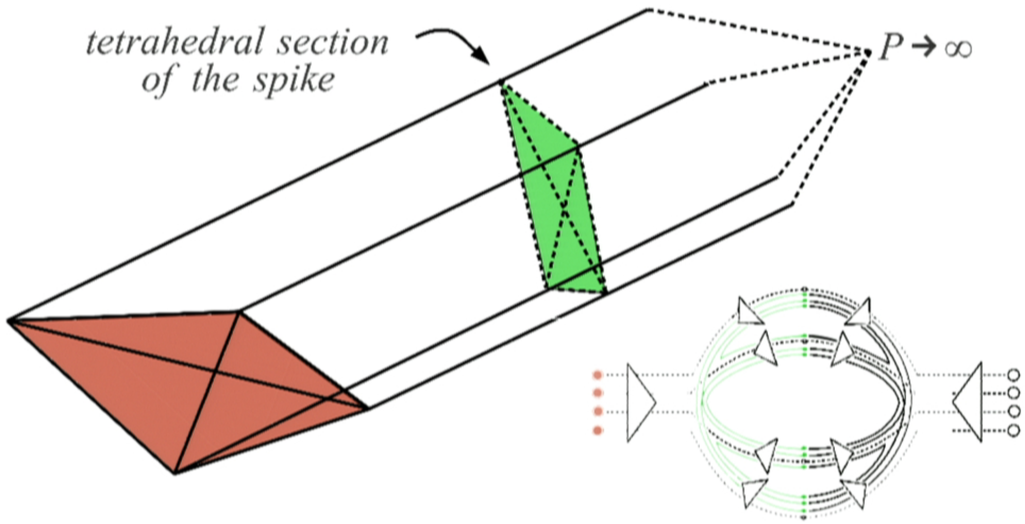
Geometrical interpretation of the calculation



The leading order term comes from a contribution where one spacetime and one “anti”- spacetime interfere, in such a way that their phases (e^{iS_0} and e^{-iS_0}, resp.) cancel. The total action is zero, without flatness

The result

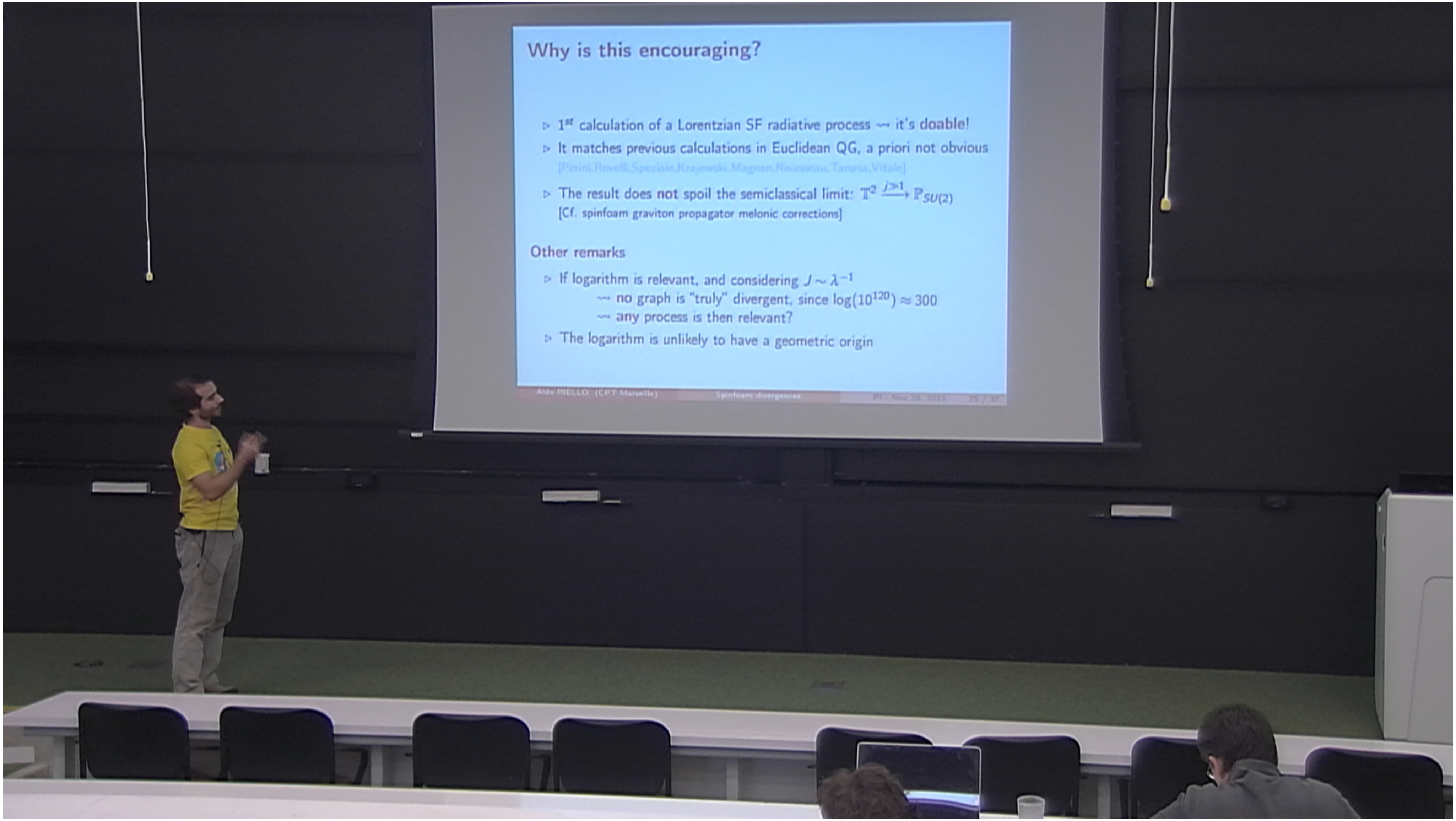
$$\mathscr{A}_{SF}[\Psi_{out}, \Psi_{in}] \sim \log\left(\frac{J}{j_{ext}}\right) \langle \Psi_{out} | \mathbb{T}^2 | \Psi_{in} \rangle$$

where the 1-bubble renormalized gluing operator $\mathbb{T}^2$ is the square of

$$\mathbb{T} \triangleright (\bullet) := \int_{SL(2,\mathbb{C})} \mathrm{d}g \; \left[\mathrm{Y}_\gamma{}^\dagger g \mathrm{Y}_\gamma \triangleright (\bullet) \right]$$

RMK $\mathbb{T}^2 \neq \mathbb{T}$, but $\mathbb{T} \xrightarrow{j_{ext} \gg 1} \langle j_{ext} \rangle^{-3/2} \mathbb{P}_{SU(2)}$, i.e. the "bare" gluing

RMK The scaling depends on the choice of face and edge amplitude


Why is this encouraging?
▷ 1^{st} calculation of a Lorentzian SF radiative process ⇝ it's doable!
▷ It matches previous calculations in Euclidean QG, a priori not obvious
[Perini,Rovelli,Speziale,Krajewski,Magnen,Rivasseau,Tanasa,Vitale]
▷ The result does not spoil the semiclassical limit: $\mathbb{T}^2 \xrightarrow{j \gg 1} \mathbb{P}_{SU(2)}$
[Cf. spinfoam graviton propagator melonic corrections]
Other remarks
▷ If logarithm is relevant, and considering $J \sim \lambda^{-1}$
⇝ no graph is "truly" divergent, since $\log(10^{120}) \approx 300$
⇝ any process is then relevant?
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  - $\rightsquigarrow$ **any** process is then relevant?
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Lessons

- "Anti"-spacetimes play a prominent role
- IR divergences mean large virtual geometries
  [though not necessarily 4d]
- If bubble divergences are not due just to symmetries, which physical meaning to IR divergences? Twist in renormalization intuition at the QG scale?
  [Is this compatible with usual QFT renormalization?]

Question

Since $\mathbb{T}^2 \xrightarrow{j \gg 1} \#\mathbb{P}_{SU(2)}$, then a melonic insertion has *no consequences* in this regime? [Apart from that of modifying the edge weight]

Indeed, it seems that amplitudes (semiclassically) should not be influenced by melonic insertions.

However, is this the end of the story?

The result

$$\mathscr{A}_{SF}[\Psi_{out}, \Psi_{in}] \sim \log\left(\frac{J}{j_{ext}}\right) \langle \Psi_{out} | \mathbb{T}^2 | \Psi_{in} \rangle$$

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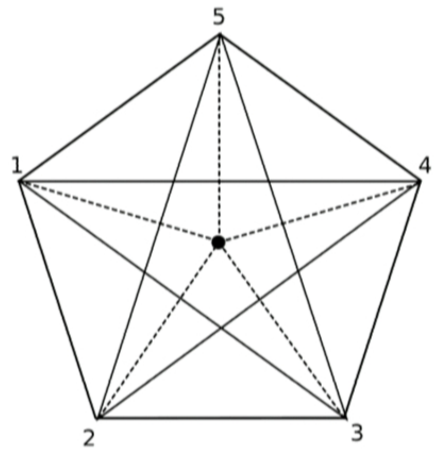
Graviton propagator

At lowest order, analyse the graviton propagator on a single four simplex
[All faces are external, there is no sum over spins associated to the dynamics]



$$\langle h^{ab}_{\mathbf{n}} h^{cd}_{\mathbf{m}} \rangle_{conn} \overset{j \gg 1}{\approx} \left[\left(S'' \right)^{-1} \right]^{ij} (h^{ab}_{\mathbf{n}})'_i (h^{cd}_{\mathbf{m}})'_j \xrightarrow{\gamma \to 0} G^{abcd}_{\mathbf{nm}}$$

RMK $\gamma \to 0$ is needed to kill correlations carried by variables different from spins, not present in Regge calculus [There are indications that this could be needed to select the right semiclassical regime in general]

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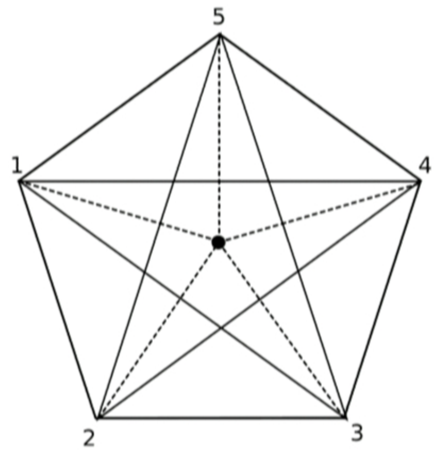
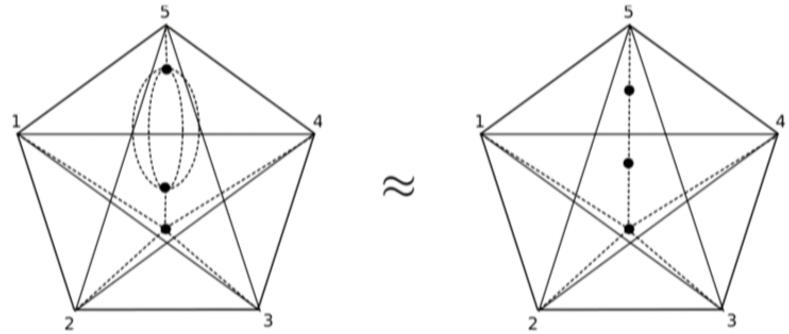


$$\langle h^{ab}_{\mathbf{n}} h^{cd}_{\mathbf{m}} \rangle_{conn} \overset{j \gg 1}{\approx} \left[\left(S'' \right)^{-1} \right]^{ij} (h^{ab}_{\mathbf{n}})'_i (h^{cd}_{\mathbf{m}})'_j \xrightarrow{\gamma \to 0} G^{abcd}_{\mathbf{nm}}$$

RMK $\gamma \to 0$ is needed to kill correlations carried by variables different from spins, not present in Regge calculus [There are indications that this could be needed to select the right semiclassical regime in general]

Melonic graviton

Now, let's put a melonic insertion in the previous calculation



[Relation valid at leading order in cut-off, thanks to previous study of the melon graph]

$$\langle h^{ab}_{\mathbf{n}} h^{cd}_{\mathbf{m}} \rangle^{melon}_{conn} \overset{j \gg 1}{\approx} \left[\left(S''_T \right)^{-1} \right]^{ij} (h^{ab}_{\mathbf{n}})'_i (h^{cd}_{\mathbf{m}})'_j$$

This expression is **different** from the previous one:

- ▷ the matrix S''_T is larger ⇝ there are more d.o.f. one can excite
- ▷ taking the inverse mixes the entries of S''_T
  - ⇝ hard to isolate new contributions
  - ⇝ *a priori* all correlations are modified

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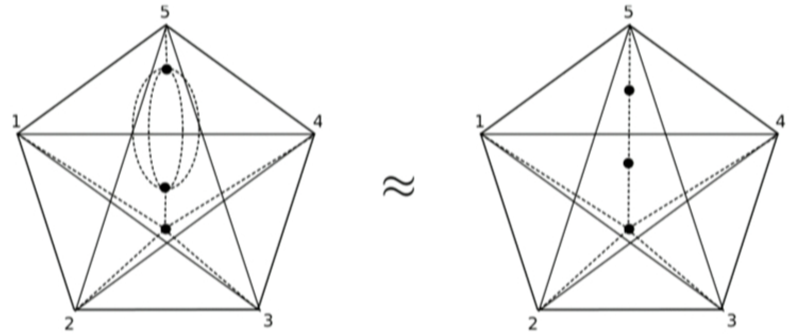

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Comment

Considering the previous graph, even if the presence of the bubble is essentially irrelevant for its amplitude (in the large spin limit), one cannot say the same for the correlations on the that graph

This is a feature of EPRL-FK, where the dominating contribution to the melon graph is *not* proportional to the bare gluing projector

Avoiding this feature would be mostly desirable, if one wanted to identify the divergences with purely topological effects related to residual symmetries

Summary

- ▷ Spin networks, spinfoams, and EPRL-FK model
- ▷ Indications of viability of EPRL-FK: semiclassics
- ▷ Origin of divergences: bubbles
- ▷ Implications divergences may have: invalidate indications of viability, source of new physics
- ▷ The melon graph: simple, crucial in cTGFT, but degenerate

Summary

- Spin networks, spinfoams, and EPRL-FK model
- Indications of viability of EPRL-FK: semiclassics
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- Implications divergences may have: invalidate indications of viability, source of new physics
- The melon graph: simple, crucial in cTGFT, but degenerate
- The EPRL-FK melon: with $SU(2)$ face weights it diverges as $\log(J)$ [vs. J^9 for $SU(2)$ BF on the same graph], it effectively modifies the gluing ($\mathbb{T}^2$) and the edge weights
- Geometrically it can be interpreted via spikes and "anti"-spacetimes

Correlations

Many interesting physical questions concern **correlations** between fluctuations on the top of some (classical) solution

In this sense, they do not directly concern the amplitude itself

In pure gravity, the best example are metric correlations on a background
$\rightsquigarrow$ **"graviton propagator"**