

Title: Fractional statistics in two-dimensions: Anyon there?

Date: Nov 06, 2013 02:00 PM

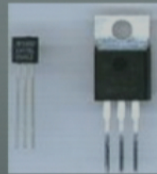
URL: <http://pirsa.org/13110062>

Abstract: A fascinating aspect of the two dimensional world is the possible existence of anyons, particles which obey 'fractional' statistics different from fermionic and bosonic statistics. In this colloquium, following an introduction to fractional particles in the context of quantum Hall systems, some of the tantalizing experiments for detecting the fractional charge of these particles will be described. Probes of fractional statistics in these systems will be discussed, drawing from analogies with the bosonic behavior of light studied by Hanbury Brown and Twiss in the 1950's as well as in the more recent beam-splitter settings in quantum optics. Finally, the exciting prospect of detecting fractional statistics in the context of superconductors will be explored.

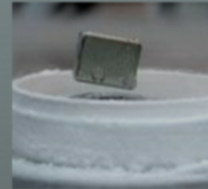
Quantum Statistics



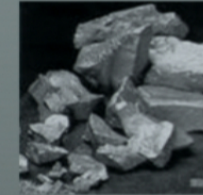
Subatomic particles



Transistors



Superconductors,
condensates



Metals



Terrestrial



Stellar and
galactic objects

Periodic Table of Elements

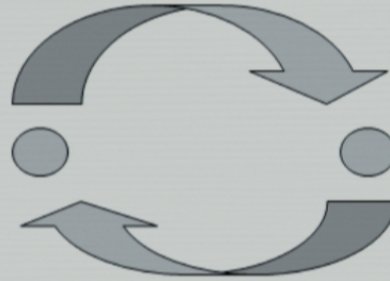
For elements with no stable isotopes, the mass number of the isotope with the longest half-life is in parentheses.

1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18
1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18
19	20	21	22	23	24	25	26	27	28	29	30	31	32	33	34	35	36
37	38	39	40	41	42	43	44	45	46	47	48	49	50	51	52	53	54
55	56	57	58	59	60	61	62	63	64	65	66	67	68	69	70	71	72
73	74	75	76	77	78	79	80	81	82	83	84	85	86	87	88	89	90
91	92	93	94	95	96	97	98	99	100	101	102	103	104	105	106	107	108
109	110	111	112	113	114	115	116	117	118	119	120	121	122	123	124	125	126

Atoms, Elements

Quantum Statistics - Bosons and Fermions

3 spatial dimensions: Only identical particles that are symmetric (bosons) or anti-symmetric (fermions) under exchange allowed.



$$\psi(\vec{r}_1, \vec{r}_2) = \pm \psi(\vec{r}_2, \vec{r}_1)$$

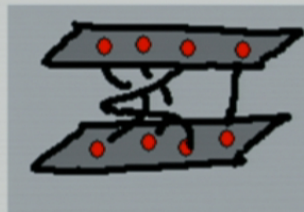
(+ bosons ; - fermions)

Two-dimensions: Anyons



Abelian

$$\psi(\vec{r}_1, \vec{r}_2) = e^{\pm i\pi\alpha} \psi(\vec{r}_2, \vec{r}_1)$$



Non-Abelian

Focus Physics Focus
[From Fermions](#) [To Bose-Einstein](#) [Superfluids](#) [To Bose-Einstein](#)
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[Phys. Rev. B 73, 070402](#)
 (issue of August 2006)
[Phys. Rev. Lett. 95, 170402](#)
 (issue of 21 October 2005)
[Titles and Authors](#)

2 November 2005

Anyon There?

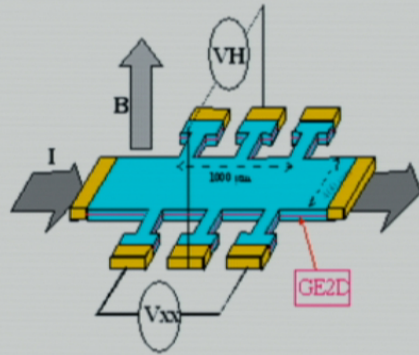
Ordinarily, every particle in quantum theory is neatly classified as either a boson—a particle happy to fraternize with any number of identical particles in a single quantum state—or a fermion, which insists on sole occupancy of its state. Now, in the August *Physical Review B*, a team shows that certain “quasiparticles” associated with electrons are somewhere in between. In the 21 October *PRL*, another team proposes a different experimental technique that should reveal these properties in an independent way. The new data confirms previous suspicions that strange spin excit in the quantum Hall system exist, and the proposed experiments could make the case even stronger.

Fermions must be alone in their quantum states, while bosons can be among an unlimited number of peers. Almost 30 years ago researchers proposed a third category, “anyons,” where a limited number of particles could inhabit a single state [1]. No one has observed this property directly, but many researchers believed that the strange state of electrons observed in the 1980s in the so-called fractional quantum Hall (FQH) effect qualifies as an anyon.

Hall of Quantum Effects Four voltage “gates” in the laboratory setup (inset) enclose a central disk with “quasiparticles” having one-fifth of an electron’s charge (red) surrounded by a ring of one-third charge quasiparticles (blue). Measurements revealed that the quasiparticles are neither fermions nor bosons.

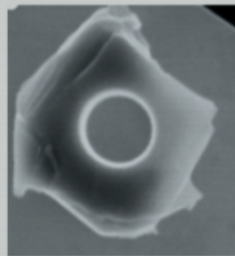
Quantum Hall System: Playground for anyons?

In this talk....



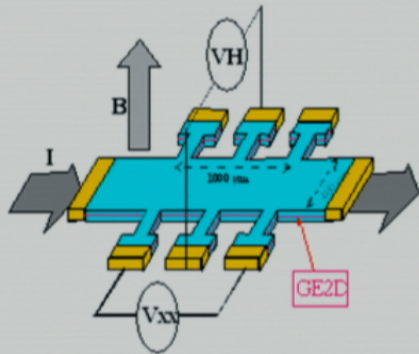
*Fractional particles in
Quantum Hall systems*

*Signatures of quantum statistics
HBT correlators; Beam splitter*

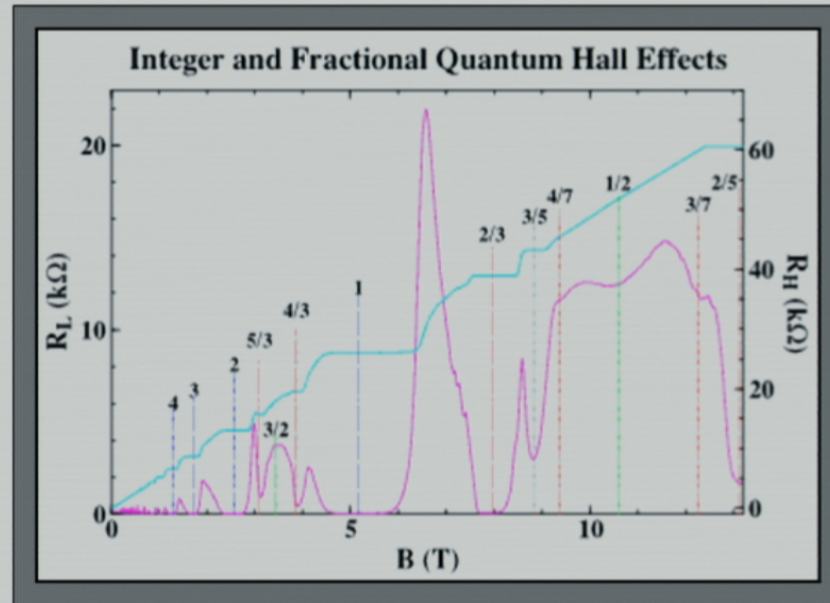


Anyons in superconducting rings

Fractional quantum Hall system



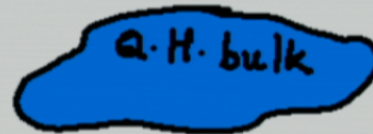
For e.g.,
D. C. Tsui et al, 1982



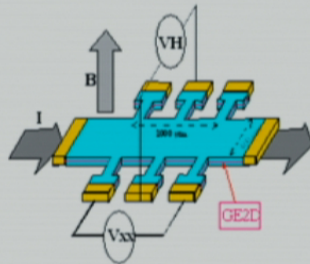
Fractional quantum Hall system

2-dimensional electron gas in
strong magnetic field;
Interactions + disorder

Incompressible
dissipationless
bulk fluid



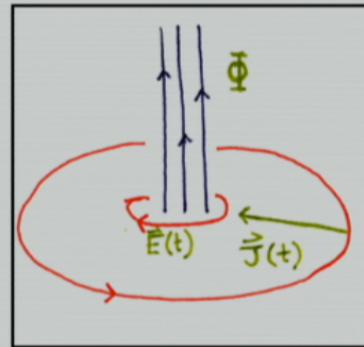
Laughlin state: Filling $\nu = 1/(2n+1)$
Quantized Conductance – $\nu e^2 / h$



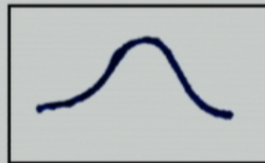
Breakdown of Fermi-liquid theory: Absence of Landau quasiparticles

Fractional quasiholes/particles

Threading extra flux quantum h/e



Fractional Charge

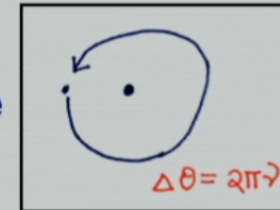


$$e^* = \nu e$$

Fractional Statistics

$$e^{\pm i\pi\nu}$$

under exchange

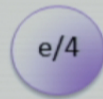


*B. I. Halperin, 1984; D. Arovas et al, 1984
E. D. M. Haldane, 1991; G. S. Jeon et al, 2003*

Non-Abelian quasiparticles

Quantum Hall ($\nu=5/2$)

$e/4$ fractional
quasiparticles



Majorana mode

Upon
exchange

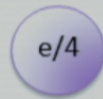


$$\begin{aligned}\nu_1 &\rightarrow \nu_2 \\ \nu_2 &\rightarrow -\nu_1\end{aligned}$$

Non-Abelian quasiparticles

Quantum Hall ($\nu=5/2$)

$e/4$ fractional
quasiparticles



Majorana mode

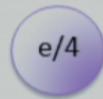
Upon
exchange




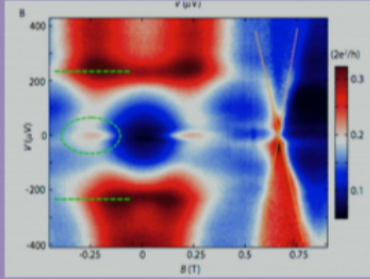
$$\begin{aligned}\nu_1 &\rightarrow \nu_2 \\ \nu_2 &\rightarrow -\nu_1\end{aligned}$$

Quantum Hall ($\nu=5/2$)

$e/4$ fractional
quasiparticles



Majorana particle glimpsed in lab
April 2012



*Mourik et al.,
Science (2010)
and more.....*

**Topological
Superconductor**

Majorana mode

Upon
exchange



$$\nu_1 \rightarrow \nu_2$$

$$\nu_2 \rightarrow -\nu_1$$

Quantum Hall edge states

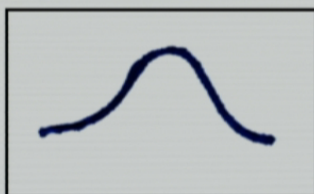
Gapless Edge Excitation

Chiral Luttinger liquid

$$L = -\frac{1}{4\pi v} \partial_x \phi (\partial_t \phi - \partial_x \phi)$$

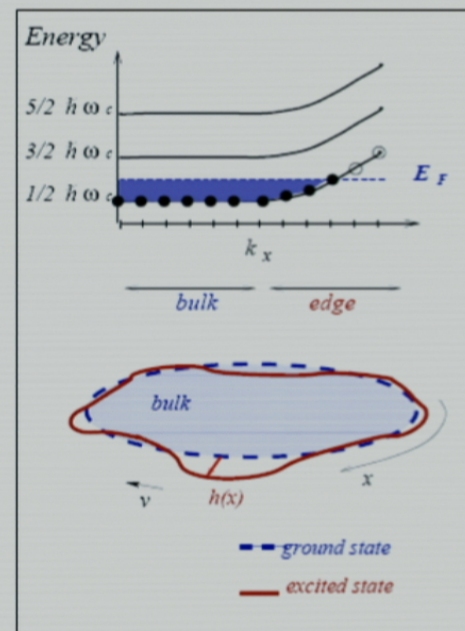
$$\delta\rho \sim \partial_x \phi$$

Edge-state Quasiparticles



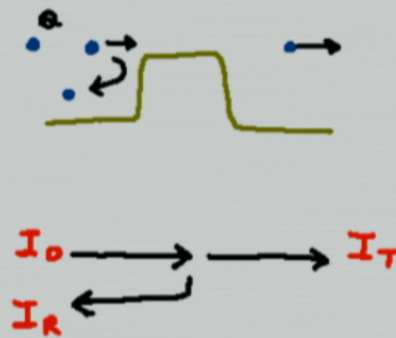
Charge νe , fractional statistics

$$\psi^\dagger(x) = \kappa e^{-i\phi(x)}$$



For e.g. X. G. Wen, PRB 41, 12838 (1990); Adv. Phys. 44, 405 (1995)

Measuring fractional charge: Shot noise

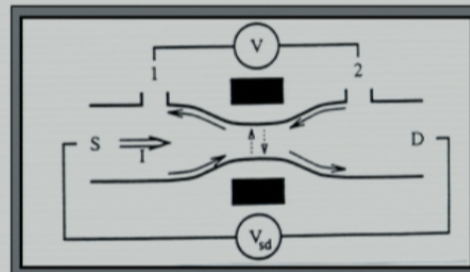


Noise:
Current-current correlation

$$S_I(\omega \rightarrow 0) = Q \langle I_B \rangle$$

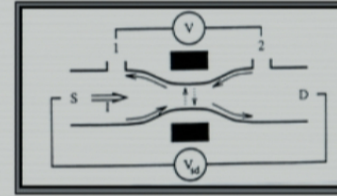
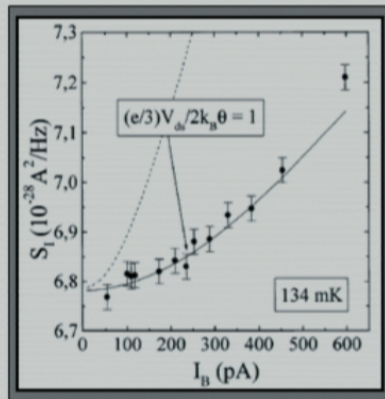
Q – Particle quantum (m,e, etc)
(Weak backscattering, T=0)

In the Hall system:



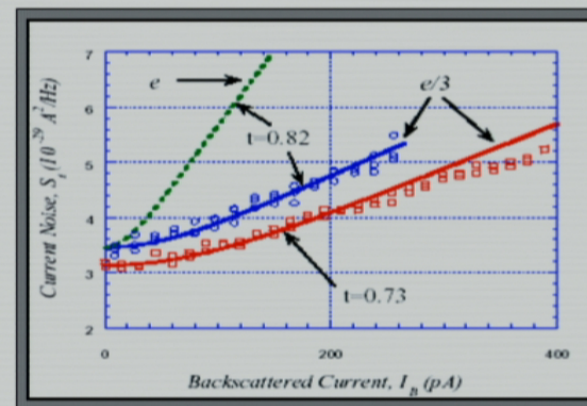
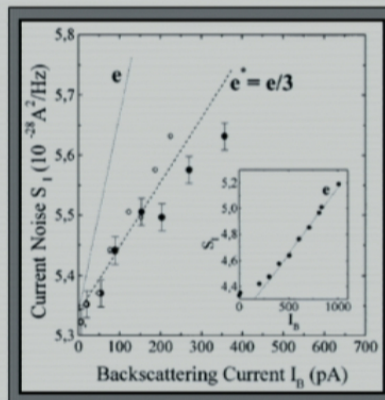
C. Kane and M. P. A. Fisher,
PRL 72, 724, 1994

Shot Noise: Experiment



$$S_I(\omega \rightarrow 0) = e^* \coth(\beta e^* V / 2) \langle I_B \rangle$$

$$e^* = \nu e$$



M. Reznikov et al., Nature 399, 238, 1999

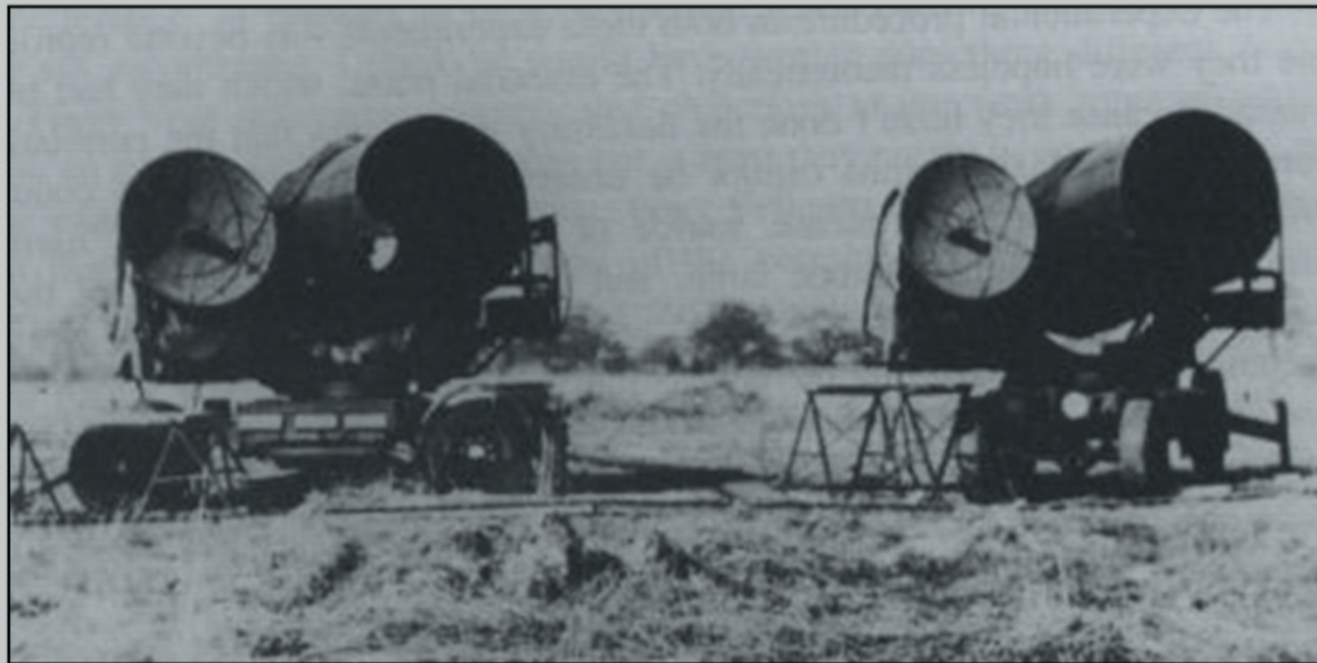
L. Saminadayar et al., PRL 79, 2526, 1997

Signatures of Quantum Statistics



Hanbury-Brown Twiss ideas

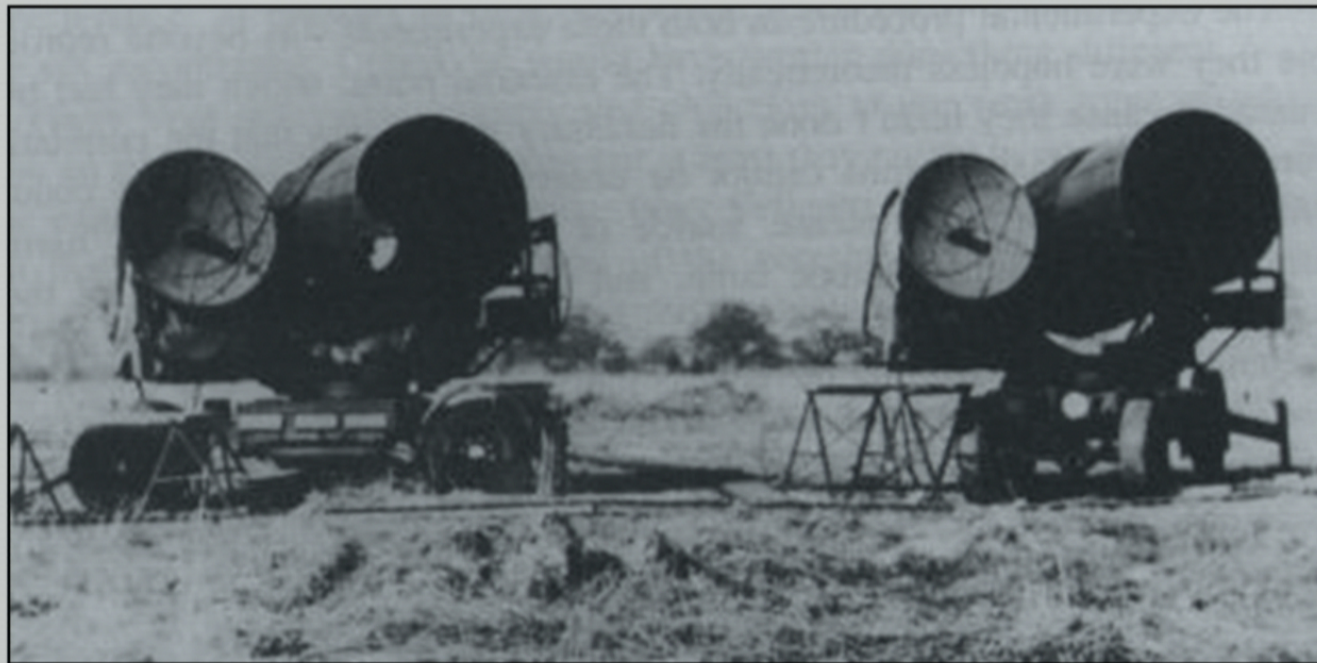
Radio Interferometry : **Intensity- intensity correlations**



Boffin – Hanbury Brown

Hanbury-Brown Twiss ideas

Radio Interferometry : **Intensity- intensity correlations**

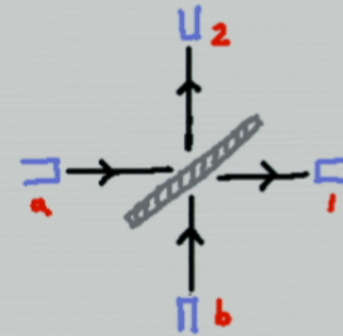


Boffin – Hanbury Brown

The Hanbury Brown-Twiss set-up - Quantum Optics

Prob. of 1 particle at detector 1:

$$P(1, i)$$

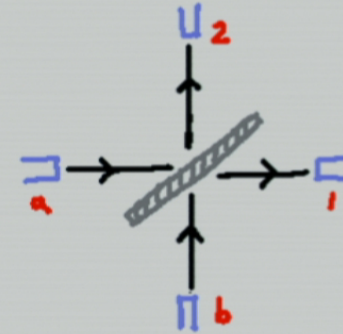


(HBT, *Nature* 177, 27 (1956); E. Purcell, *Nature* 178, 1449; G. Baym, *Act.Phys. Pol B* 29,1, etc)

The Hanbury Brown-Twiss set-up - Quantum Optics

Prob. of 1 particle at detector 1:

$$P(1, i)$$



Prob. of 1 particle in detector 1, and one in detector 2

$$P(1, 2; i, j) = P(1, i)P(2, j) + P(2, i)P(1, j) + \text{exchange}$$

Two-particle correlation function (propagator)

$$\langle \psi_i^\dagger(x) \psi_j^\dagger(x'') \psi_j(x''') \psi_i(x') \rangle$$

(HBT, Nature 177, 27 (1956); E. Purcell, Nature 178, 1449; G. Baym, Act.Phys. Pol B 29,1,etc)

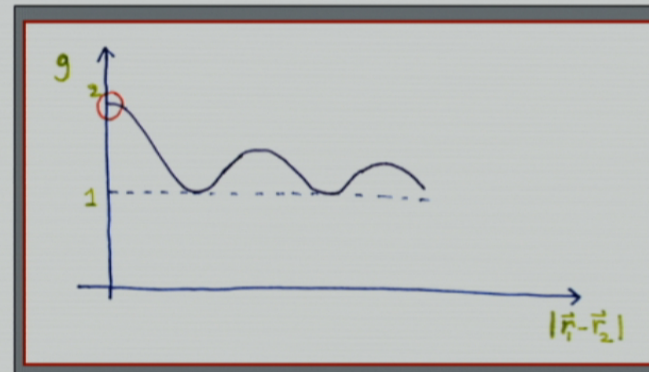
2-particle correlation function

$$g(r_1, r_2) = N(N-1) \int dr_3 \dots dr_N |\Psi(r_1, r_2, \dots, r_N)|^2$$

Statistics

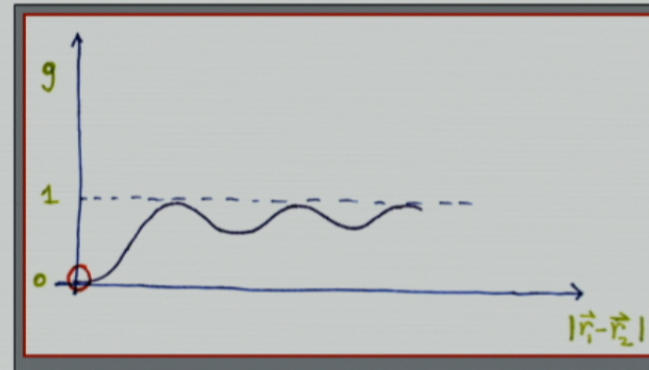
Bosons

E.g. Photons from different sources



Fermions

E.g. Spinless non-interacting electrons



2-particle correlation function

$$g(r_1, r_2) = N(N-1) \int dr_3 \dots dr_N |\Psi(r_1, r_2, \dots, r_N)|^2$$

Statistics

Characteristics of System

Oscillations

Position/Momentum
Energy/Time

Decay

E.g. Spinless non-interacting electrons

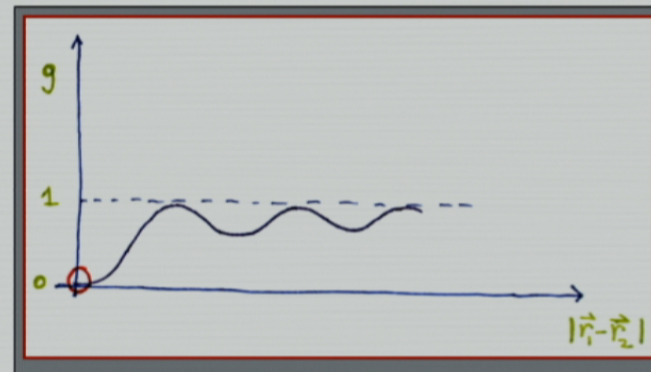
$$g \sim \left[1 - \left(\frac{\sin k_F r}{k_F r} \right)^2 \right] \quad (\text{for one-dimension})$$

HBT Correlation Regime

$$\Delta x \ll (\Delta k)^{-1}$$

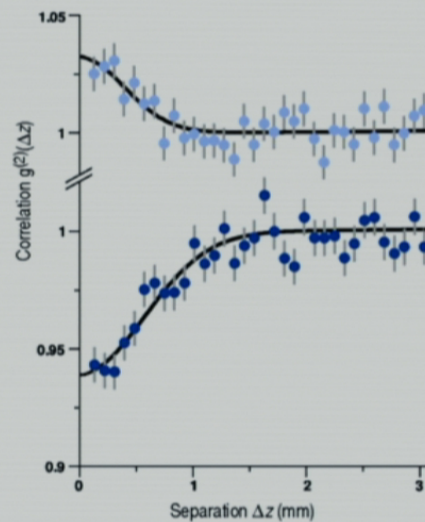
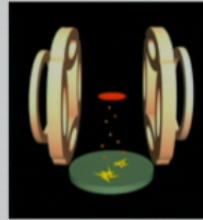
$$\Delta t \ll \omega^{-1}$$

$$kT \ll \Delta E$$



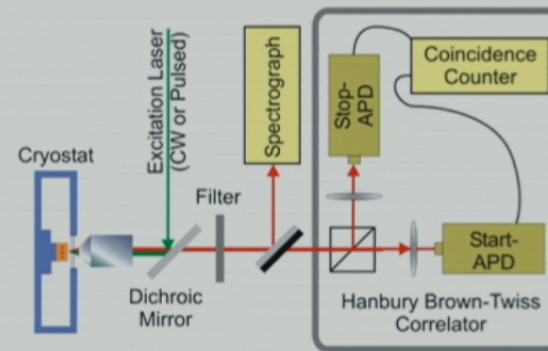
Manifestations

Correlations in trapped helium



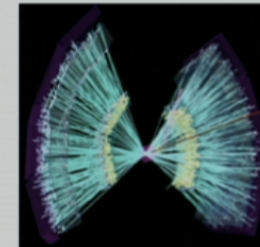
Gruppe d'Optique Atomique, France

Quantum Optics



Nano Optics, Humboldt Univ., Germany

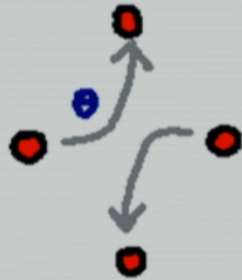
Scattering cross-sections



Nuclear physics group, UIUC

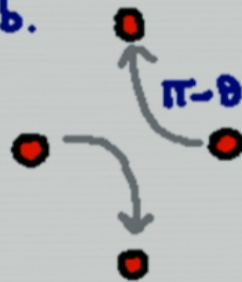
Scattering cross-sections

a.

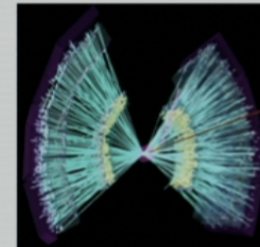


Indistinguishable Particles:
Paths a and b are both present.
Cannot be distinguished

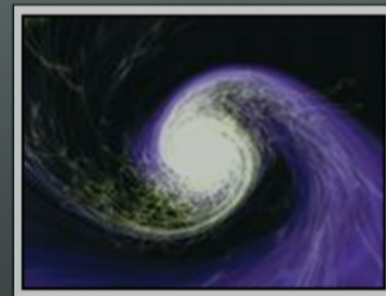
b.



Identical fermions:
Scattering cross-section $|f(\theta) - f(\pi - \theta)|^2$
Vanishes at 90°

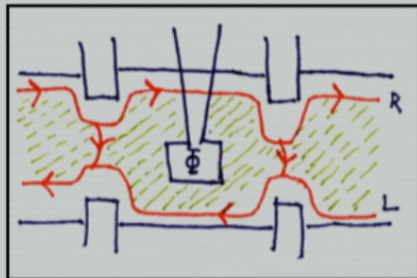


Quantum Statistics - Anyon there?



Proposals for measuring fractional statistics

Anomalous Aharonov-Bohm period

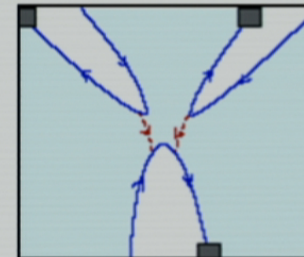


Chamon et al., 1997

Partitioning Noise

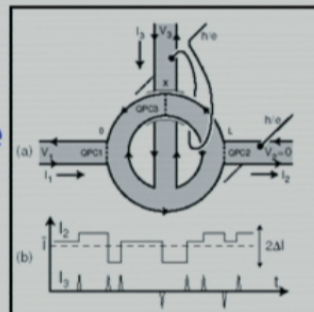
Safi et al., 2001; SV, 2003

E.Kim et al., 2005;

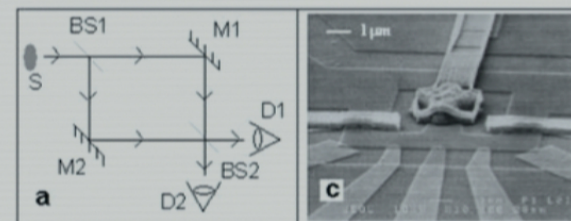


Telegraph Noise

C. L. Kane, 2003



Mach-Zehnder Interferometry



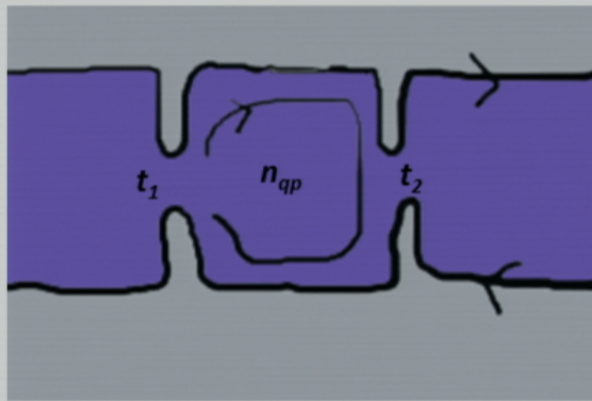
Ji et al., 2003

Some related experiments

Y.C. Chung et al., 2002; F.F. Comino et al., 2005

Proposal for detecting non-Abelian statistics

Quantum Hall $\nu=5/2$



$$\begin{aligned}\nu_1 &\rightarrow \nu_2 \\ \nu_2 &\rightarrow -\nu_1\end{aligned}$$

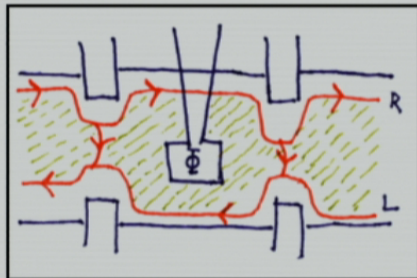
Interferometry

$$\sigma_{xx}^{int} \propto \text{Re}[t_1^* t_2 e^{i\phi} \langle \psi_0 | U_1^{-1} U_2 | \psi_0 \rangle]$$

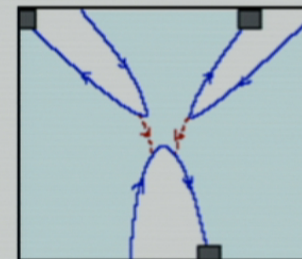
For e.g., E. Fradkin et al., 1998; S. Das Sarma et al., 2005;
Steen & Halperin, 2006; B. Bandyopadhyay et al., 2006

Proposals for measuring fractional statistics

Anomalous Aharonov-Bohm period



Chamon et al., 1997



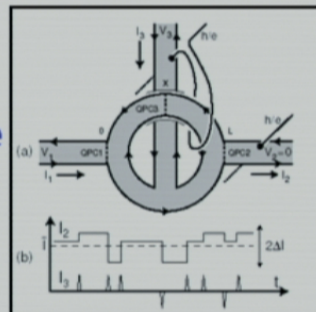
Partitioning Noise

Safi et al., 2001; SV, 2003

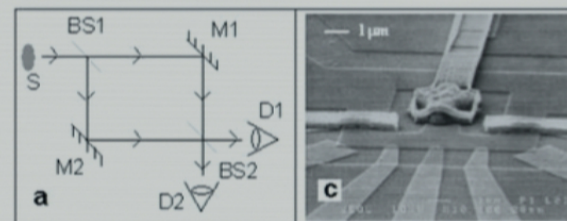
E.Kim et al., 2005;

Telegraph Noise

C. L. Kane, 2003



Mach-Zehnder Interferometry



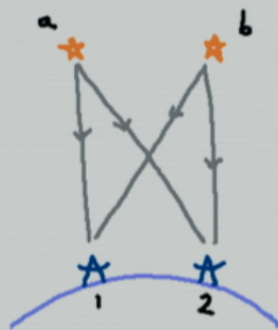
Ji et al., 2003

Some related experiments

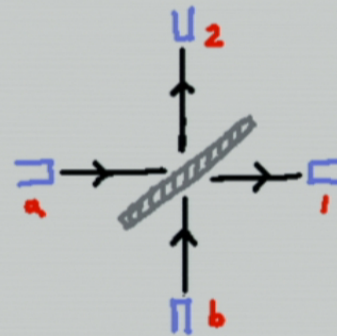
Y.C. Chung et al., 2002; F.F. Comino et al., 2005

Quantum statistics: two-particle properties

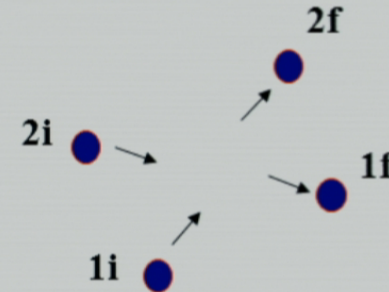
Hanbury-Brown Twiss



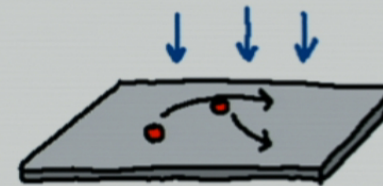
Beam-splitters



Scattering processes

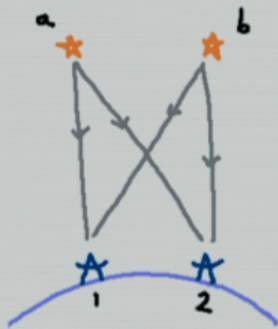


Signatures of statistics in bunching behavior,
interference and angular dependence

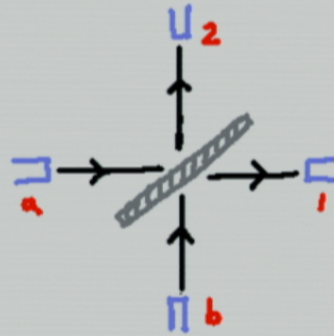


Quantum statistics: two-particle properties

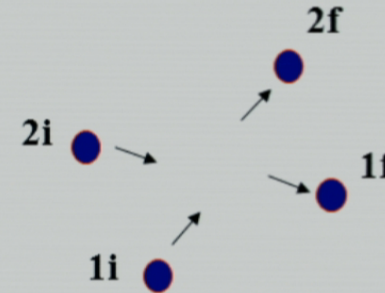
Hanbury-Brown Twiss



Beam-splitters

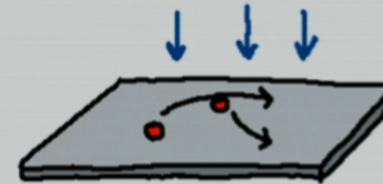


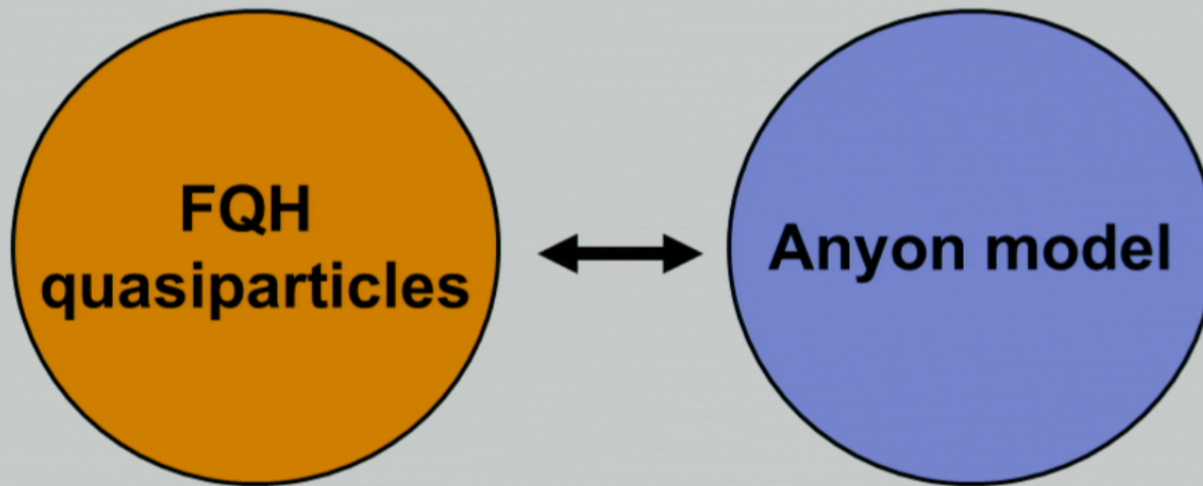
Scattering processes



Signatures of statistics in bunching behavior,
interference and angular dependence

Manifestations in QH bulk?





*B. I. Halperin, 1984; R. B. Laughlin, 1987
J. Myerheim in Les Houches series (1999);
J. Jain, Composite Fermions (2007)*

Two-anyon model

$$H = \frac{1}{4\mu} \left(P_x + \frac{qB}{c} Y \right)^2 + \frac{1}{4\mu} \left(P_y - \frac{qB}{c} X \right)^2 \\ + \frac{1}{\mu} \left(p_x + \frac{qB}{4c} y \right)^2 + \frac{1}{\mu} \left(p_y - \frac{qB}{4c} x \right)^2$$

Center of mass: (\vec{R}, \vec{P})

Magnetic field \mathbf{B} perpendicular to plane

Relative co-ordinates: (\vec{r}, \vec{p})



Anyonic feature

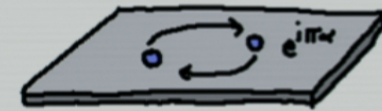
$$\psi(-\vec{r}) = e^{\pm i\pi\alpha} \psi(\vec{r})$$

Two-anyon model

Wavefunctions

$$\psi_p(\vec{r}) = \frac{(4\pi m)^{-1/2}}{\sqrt{\Gamma(2p+\nu+1)l}} \left(\frac{z}{2\sqrt{m}}\right)^{2p+\nu} \exp\left[-\frac{|z|^2}{8ml^2}\right];$$

$$\psi_n(\vec{R}) = \frac{1}{l\sqrt{m\pi n!}} \left(\frac{Z}{\sqrt{m}}\right)^n \exp\left[-\frac{|Z|^2}{2m}\right]$$

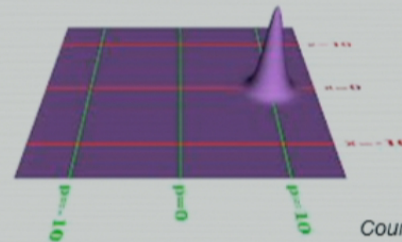


(Symmetric gauge; lowest Landau level; $\alpha = \nu = 1/m$)

Center of mass: (\vec{R}, \vec{P}) Relative co-ordinates: (\vec{r}, \vec{p})

Analogy: Quantum Optics

Coherent States in
Phase Space



Courtesy Denker

Two-anyon LLL Hilbert space

Center of mass:

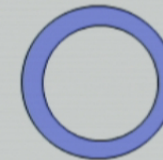
Guiding center coordinates

$$\hat{X} = l(\hat{A} + \hat{A}^\dagger)/2, \quad \hat{Y} = il(\hat{A} - \hat{A}^\dagger)/2$$

Angular momentum eigenstates

$$\hbar A^\dagger A |n\rangle_c = n\hbar |n\rangle_c \quad \psi_n(\vec{R}) \propto \frac{Z^n}{\sqrt{n!}}$$

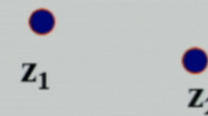
$$[A, A^\dagger] = 1$$



Localized (coherent) states

$$|Z\rangle_c = e^{-|Z|^2/2} \sum_{n=0}^{\infty} \frac{(Z^*)^n}{\sqrt{n!}} |n\rangle_c$$

$$Z = (z_1 + z_2)/2$$



T. H. Hansson, L. M. Leinaas, L. Murubei, 1992 (44M02); H. Kiersberg and L. M. Leinaas, 1997 (46107)

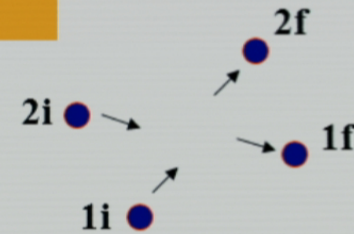
Objects of attention

Single particle

$$K_1(\vec{R}_f; \vec{R}_i) = \sum_n \psi_n(\vec{R}_f) \psi_n^*(\vec{R}_i) e^{-E_n \tau / \hbar}$$

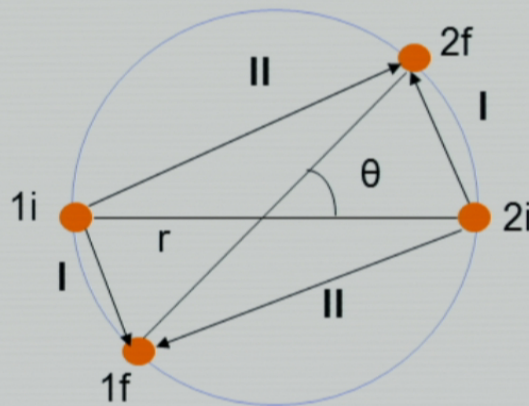
Two particles

$$K_2(\vec{r}_{1f}, \vec{r}_{2f}; \vec{r}_{1i}, \vec{r}_{2i}) = \sum_n \psi_n(\vec{R}_f) \psi_n^*(\vec{R}_i) e^{-E_n \tau / \hbar} \\ \times \sum_p \psi_p(\vec{r}_f) \psi_p^*(\vec{r}_i) e^{-E_p \tau / \hbar}$$



SV et al., 2007

Two-particle propagator in LLL* - fermions and bosons



Destructive Interference

Magnitudes of I and II equal:

$$\theta = \pi/2$$

Relative phase $0/\pi$:

$$2Br^2/\Phi_0 = (2n+0/1)/2$$

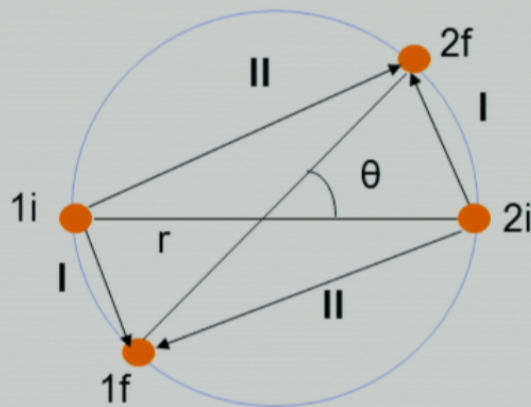
**Two-particle Aharonov-Bohm
+ quantum statistics**

$$K_2(z_{1f}, z_{2f}; z_{1i}, z_{2i}) = K_1(z_{1f}; z_{1i})K_1(z_{2f}, z_{2i}) \mp K_1(z_{2f}; z_{1i})K_1(z_{1f}, z_{2i})$$



* LLL - Lowest Landau level

Two-particle propagator in LLL* - fermions and bosons



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* LLL - Lowest Landau level

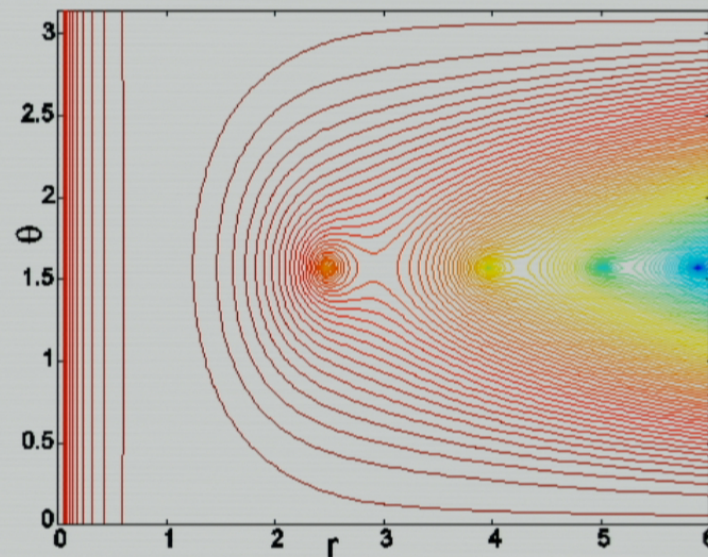
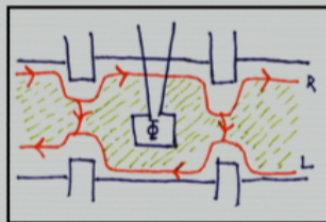
Two-particle propagator - anyons (A ν angle on fractional statistics)

Destructive Interference

Similar geometric picture:

Also at $\theta = \pi/2$

$$2Br^2/\Phi_0 \approx (2n+1+\nu)/2$$



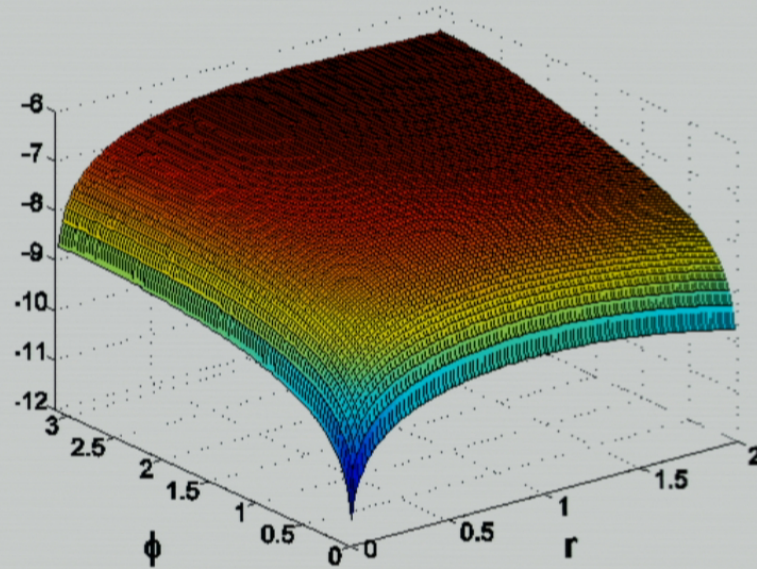
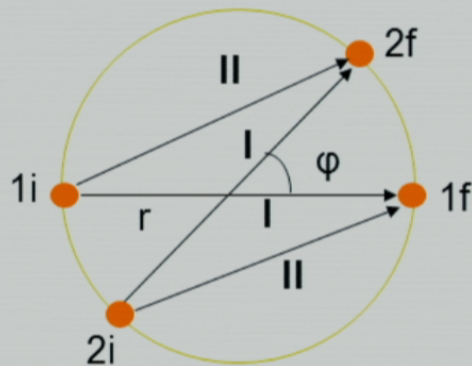
Two-particle Aharonov-Bohm
+ Fractional statistics

Two-particle propagator – anyons (A ν angle on fractional statistics)

Exclusion behavior

Haldane, 1991

$$K_2 \sim |\varphi|^{2\nu}, \varphi \rightarrow 0$$



$$\text{For } 1i = 2i = 0, 1f = 2f = r, K_2 \sim |r|^{2\nu}, r \rightarrow 0$$

Bunching properties

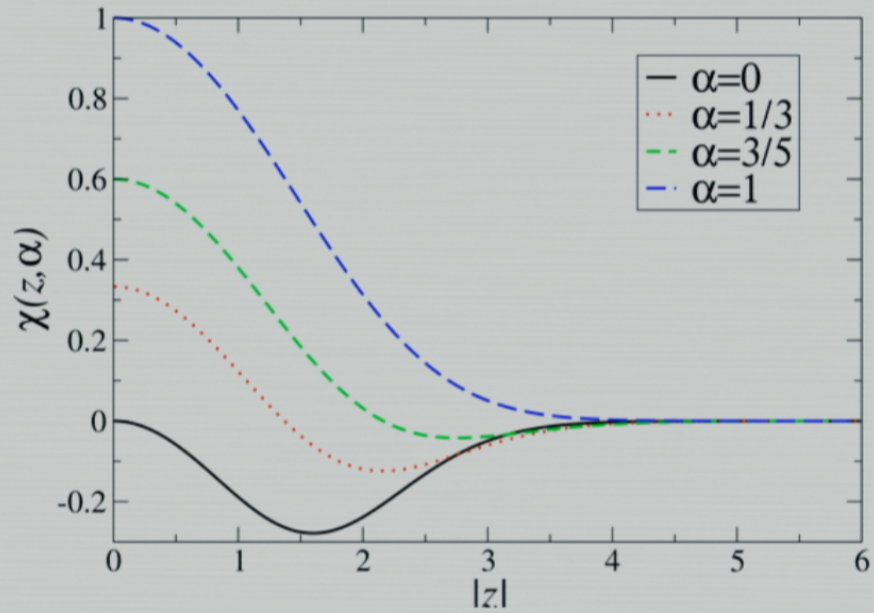


Bunching parameter

$$\chi(|z|, \alpha) \equiv \frac{1}{4l^2} [\langle z | \hat{r}^2 | z \rangle_\alpha - \langle z | \hat{r}^2 | z \rangle_d]$$

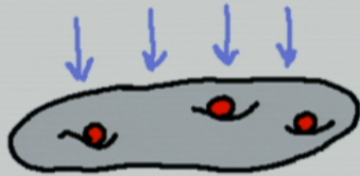
$$\langle \hat{r}^2 \rangle \equiv \langle \hat{x}^2 + \hat{y}^2 \rangle$$

Bunching parameter



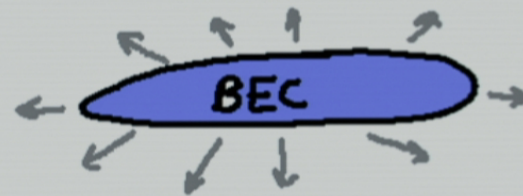
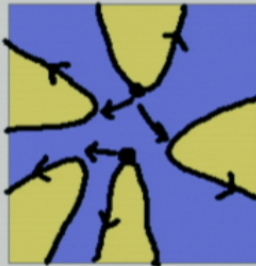
$$\chi = (|z|^2/4)[\coth(|z|^2/4) - 1] \quad \text{Fermions}$$
$$= (|z|^2/4)[\tanh(|z|^2/4) - 1] \quad \text{Bosons}$$

Probing Correlations



Localized quantum Hall bulk quasiparticles

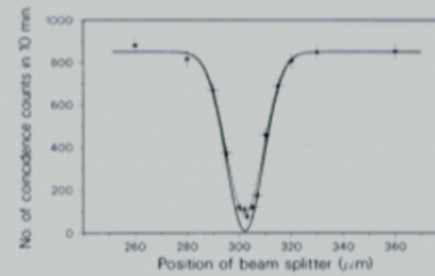
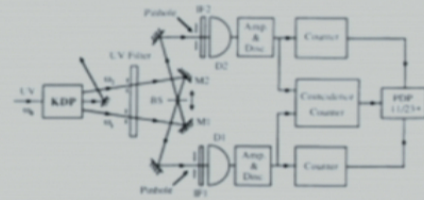
Correlations in ultracold gases



Tunneling into multiple edge states

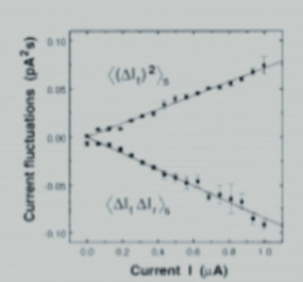
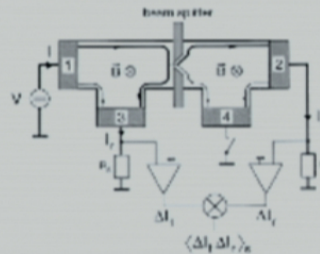
Beam Splitters and Statistics

Coincidence measurements for photons

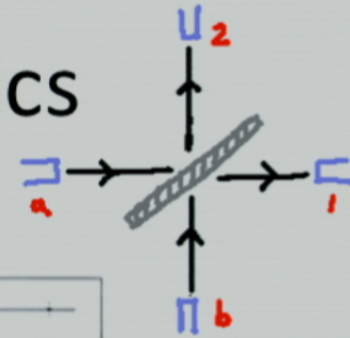


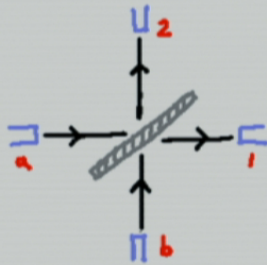
C. K. Hong et al., 1987

Current correlation measurements for electrons

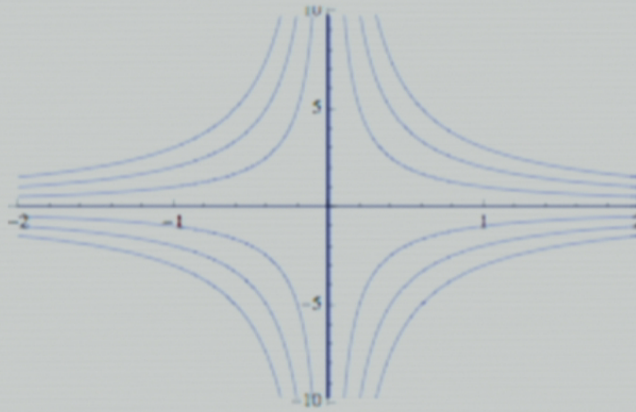


M. Henny et al; W. D. Oliver et al., 1999



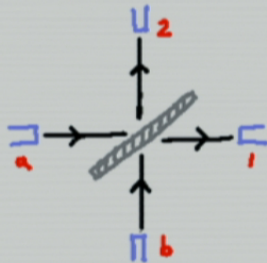


QH beam splitter – Saddle potential

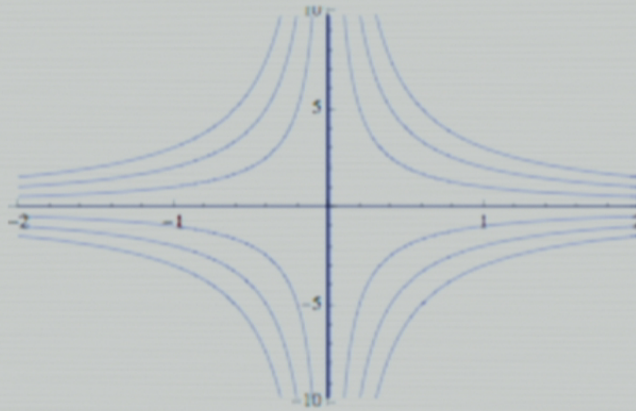


$$\hat{H}_s = \sum_{\gamma=1,2} U \hat{x}_\gamma \hat{y}_\gamma$$

Separable,
Preserves anyon b.c.



QH beam splitter – Saddle potential



$$\hat{H}_s = \sum_{\gamma=1,2} U \hat{x}_\gamma \hat{y}_\gamma$$

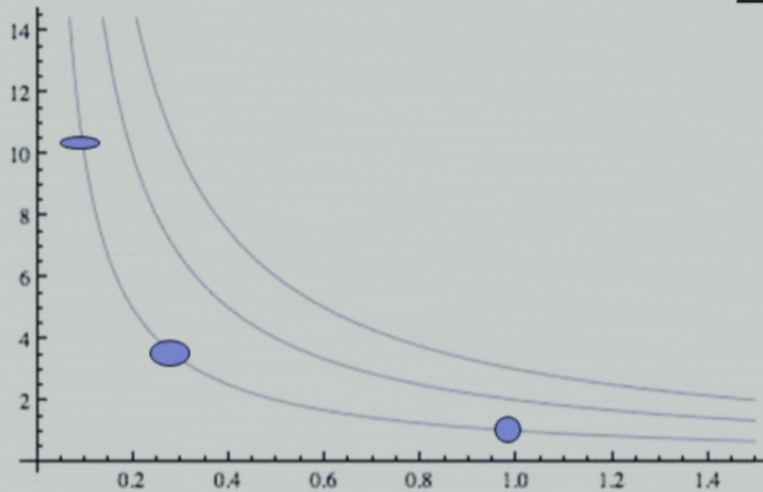
Separable,
Preserves anyon b.c.

LLL projected Hamiltonian

$$\hat{H}_s^P = \frac{1}{2} i U l^2 [\hat{A}^2 - (\hat{A}^\dagger)^2] + 2 U l^2 \hat{c}$$



Saddle dynamics



C.o.m motion

$$|Z\rangle_c = \exp(Z\hat{A}^\dagger - Z^*\hat{A})|0\rangle_c$$
$$|Z(t)\rangle_c = e^{-i\hat{H}_s^P t/\hbar}|Z\rangle_c$$

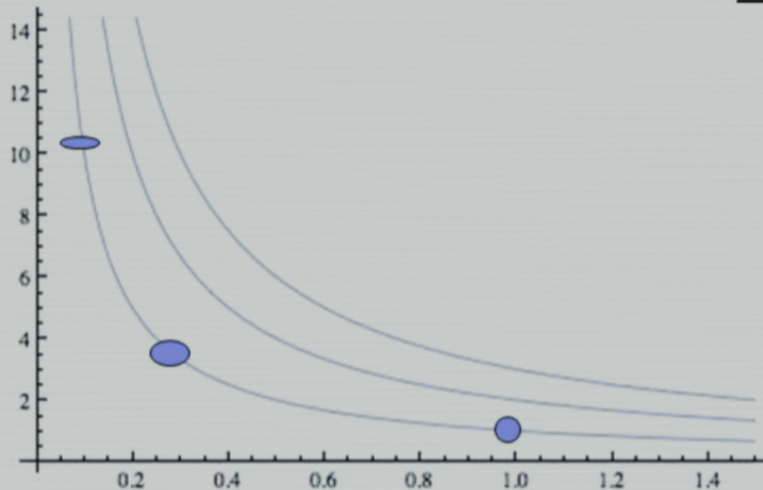
Analogy from qtm optics

Squeeze along $re^{i\phi} \equiv Ut^2/\hbar$

Follow

$$(Xe^{-Ut^2/\hbar}, Ye^{Ut^2/\hbar})$$

Saddle dynamics



C.o.m motion

$$|Z\rangle_c = \exp(Z\hat{A}^\dagger - Z^*\hat{A})|0\rangle_c$$

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Analogy from qtm optics

Squeeze along $re^{i\phi} \equiv Ut^2/\hbar$

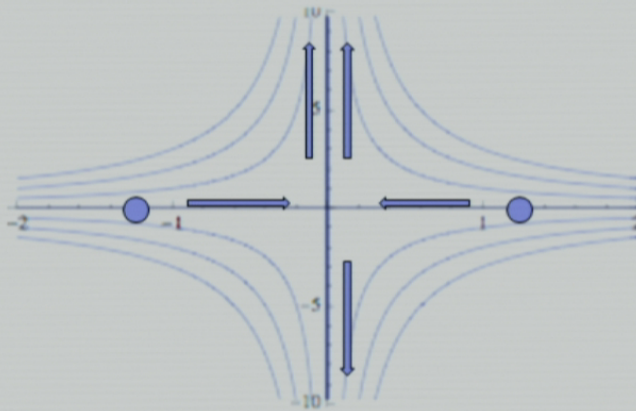
Follow

$$(Xe^{-Ut^2/\hbar}, Ye^{Ut^2/\hbar})$$

Relative motion (Initial amplitude peaked at +z and -z)

$$\hat{x}^2(t) = e^{-2Ut^2/\hbar}\hat{x}^2(0), \hat{y}^2(t) = e^{2Ut^2/\hbar}\hat{y}^2(0)$$

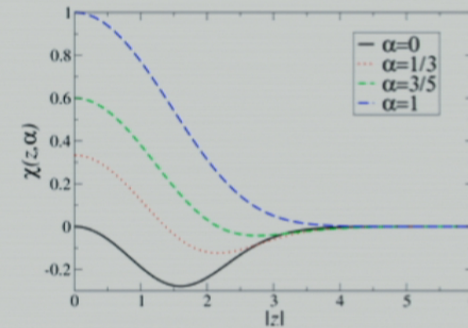
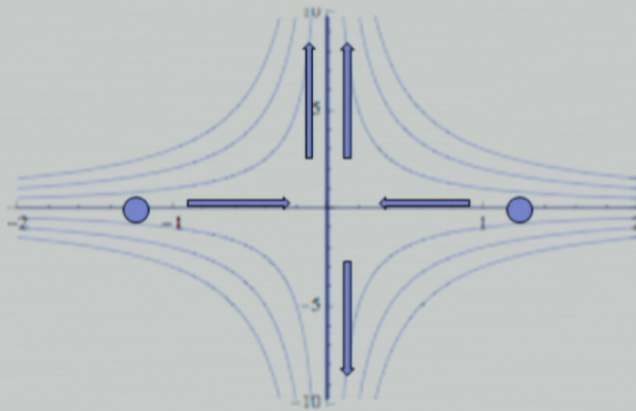
Beam splitter properties



Did the particles go in the same direction or different ones?

Behavior of $\langle y_1 y_2 \rangle$

Beam splitter properties



Did the particles go in the same direction or different ones?

Depends on statistics and bunching parameter

Behavior of $\langle y_1 y_2 \rangle$

$$\langle \hat{y}_1 \hat{y}_2 \rangle = l^2 e^{2Utl^2/\hbar} \left[\text{Im}[Z]^2 - \frac{1}{4} \text{Im}[z]^2 - \frac{1}{2} \chi + \delta \right]$$

QH bulk non-Abelian quasiparticles?

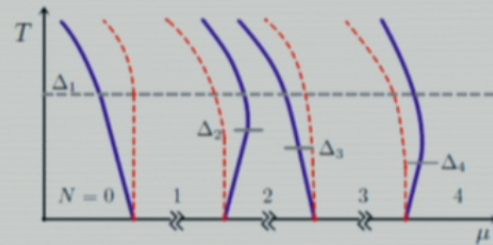
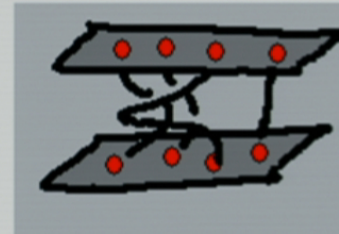
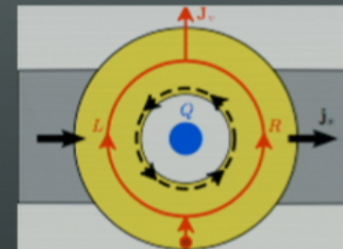


FIG. 1. Cartoon charge stability diagram, showing only peak centres (no broadening). Vertical axis is temperature T ; horizontal axis μ is the chemical potential for charged QPs, controlled in experiment by a gate potential. Red dashed lines are for Abelian particles. Blue solid lines correspond to non-Abelian QPs in a tightly confining well. Δ_N is the gap to excited states for N particles, and sets scale for other entropic effects. Notice even-odd effect for non-Abelian anyons.

Ben-Shach et al., 2013



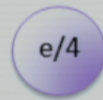
Non-Abelian statistics in superconductor rings



Non-Abelian quasiparticles

Quantum Hall ($\nu=5/2$)

$e/4$ fractional quasiparticles



Chiral p-wave superconductors

Zero energy state in vortex core



$h/(2e)$ or $h/(4e)$ (One or two component sc)

Majorana mode

Upon exchange

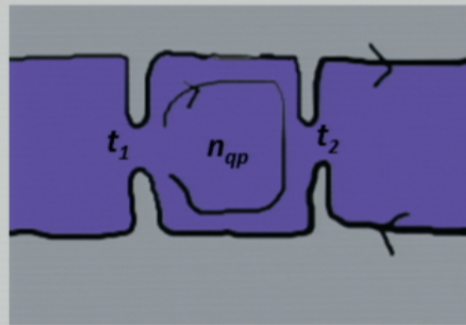


$$\nu_1 \rightarrow \nu_2$$

$$\nu_2 \rightarrow -\nu_1$$

Non-Abelian quasiparticles: detection

Quantum Hall ($\nu=5/2$)



Interferometry

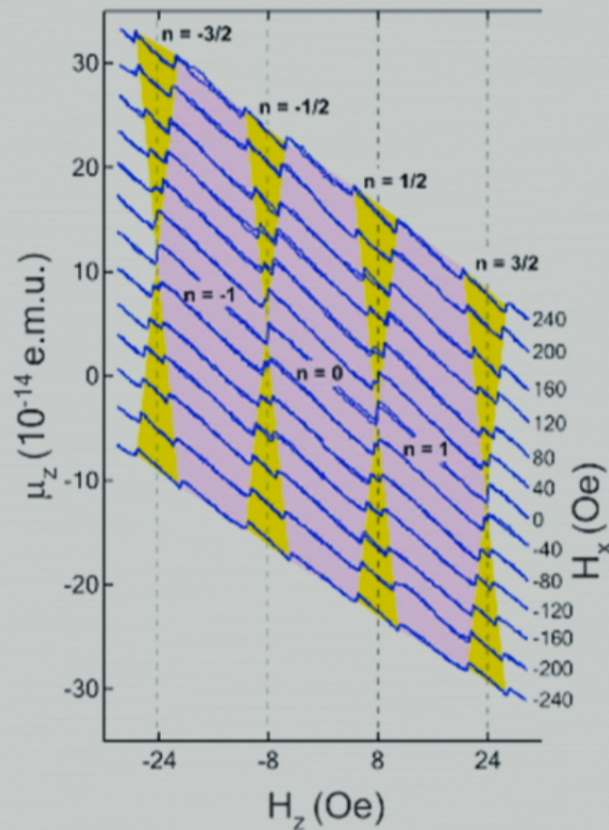
$$\sigma_{xx}^{int} \propto \text{Re}[t_1^* t_2 e^{i\phi} \langle \psi_0 | U_1^{-1} U_2 | \psi_0 \rangle]$$

For e.g., E. Fradkin et al., 1998; S. Das Sarma et al, 2005;
Stern & Halperin, 2006; P. Bonderson et al., 2006

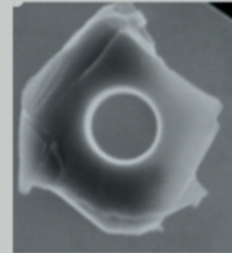
Chiral p-wave superconductors (CpSC)

DUALITY!!?

Measurements from the lab of Raffi Budakian



Strontium Ruthenate (SRO):
Candidate CpSc material



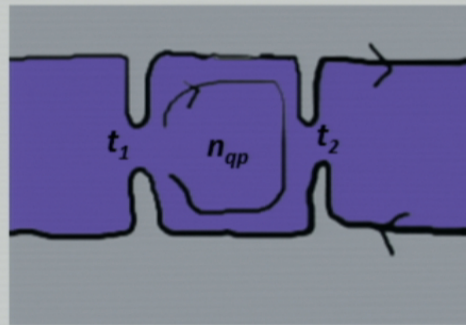
SRO sample –
Mesoscopic ring
Order 1500nm

Magnetic response: Splitting
as a function of H_x indicates
Half Quantum Vortices (HQV)

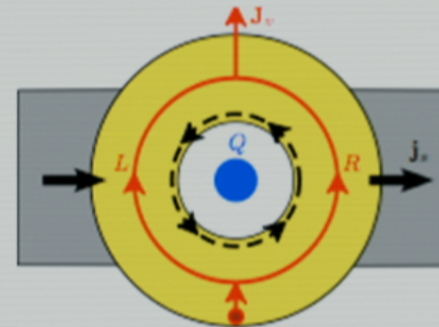
*Courtesy R. Budakian
J. Jang et al, 2011*

Non-Abelian quasiparticles: detection

Quantum Hall ($\nu=5/2$)



Chiral p-wave superconductors



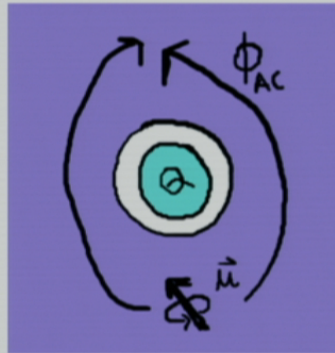
Interferometry

$$\sigma_{xx}^{int} \propto \text{Re}[t_1^* t_2 e^{i\phi} \langle \psi_0 | U_1^{-1} U_2 | \psi_0 \rangle]$$

$$J_v^{int} \propto |t_L| |t_R| [e^{i\phi} \langle \psi_L | \psi_R \rangle]$$

F. Greenfold et al., 2011

Aharonov-Casher effect



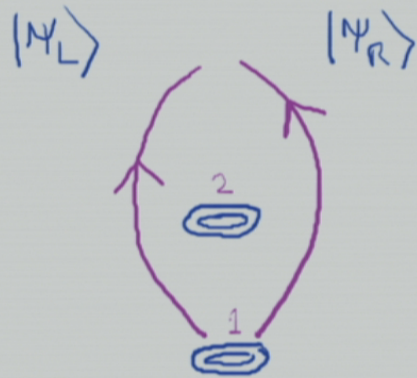
Magnetic moment moving in static electric field – geometric phase:

$$\phi_{AC} = \frac{1}{\hbar c^2} \oint \vec{\mu} \times \vec{E} \cdot d\vec{l}$$

2d vortex encircling charge Q : $\phi_{AC} = \Phi_V Q / \hbar = 2\pi Q / e^*$

Aharonov and Casher, '84; Reznik and Aharonov, '89

Encircling Majorana vortices



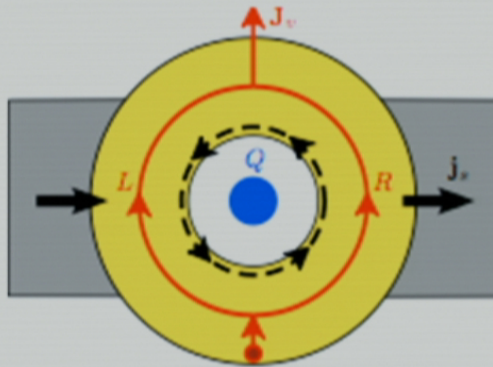
Majorana structure:

$$\begin{aligned} \nu_1 &\rightarrow \nu_2 \rightarrow \langle \psi_L | \psi_R \rangle = 0 \\ \nu_2 &\rightarrow -\nu_1 \end{aligned}$$

Two-component CpSc:

Holds for HQV around HQV as well as around FQV (FQV: can consider one Majorana in each spin component)

Proposed Experiment



$$J_v^{int} \propto |t_L||t_R| [e^{i\varphi} \langle \psi_L | \psi_R \rangle]$$

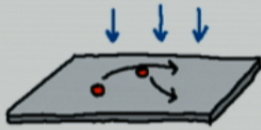
- In zero field, drive j_s
- Vary charge Q , see dI/dV oscillation of period $2e$ or $4e$
- Turn on field to nucleate HQV
- Repeat AC procedure
- Oscillations should disappear

*Many thanks to R. Budakian and D. VanHarlingen
E. Grosfeld et al, 2011*

Alternative: Josephson vortices

*Also: E. Grosfeld and Stern, 2011
Boonakker 2011; Alicia 2012*

Thanks to....



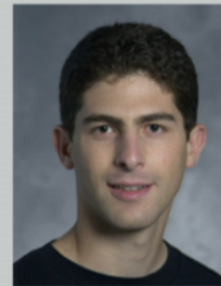
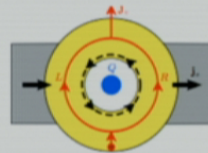
Diptiman Sen
IISc



Michael Stone
UIUC



Nigel Cooper
Univ. of Cambridge



Eytan Grosfeld
Ben Gurion Univ.



Babak Seradjeh
Indiana Univ.



In summary,

- Quantum Hall systems are expected to display the fascinating phenomenon of fractional statistics

- Two-particle correlations exhibit signatures of quantum statistics – interference, bunching, spatial profiles

- Non-Abelian statistics in quantum Hall systems and topological superconductors



Anyon there!!?