

Title: Asymptotic Dynamics: Spin Foam Partition Functions in an Asymptotic Regime

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Abstract: <span>Spin foam models are models for space time built from discrete chunks of quantized geometry. In the asymptotic regime the classical geometry is regained.<br>In the last year we have seen rapid developments in our understanding of this geometry at the level of the entire partition function. In particular it was found that the geometries that contribute to the partition function in the asymptotic regime satisfy accidental curvature constraints.<br>I will discuss the classic results and role of asymptotics, the recent results and their impact on the interpretation of these models.</span>

# LOOPS 13

## Spin foams

Asymptotic geometry of the partition function

July 25<sup>th</sup> 2013  
Frank Hellmann

## Spin foam models

are

- discrete (2-complex: vertices, edges and faces)
- with quantized geometry

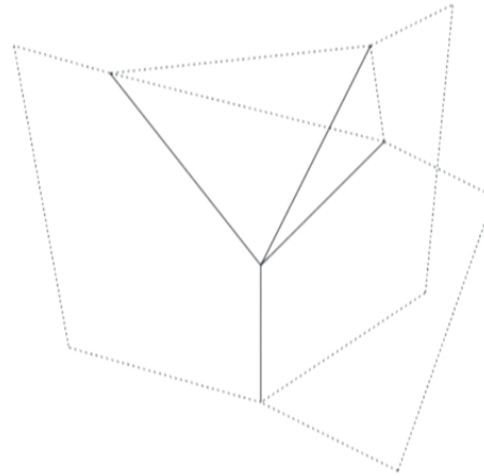
to provide covariant dynamics for network states.

Many contributors over the years:

Alesci, Alexandrov, Baez, Baratin, Barbieri, Barrett, Bahr, Bianchi, Bonzom, Carrozza, Cianfrani, Conrady, Crane, DePietri, Dittrich, Dupuis, Ding, Engle, Fairbairn, Freidel, Gambini, Girelli Gomes, Hal, Han, Hnybida, Immirzi, Kaminski, Kisielowski, Krasnov, Lewandowski, Livine, Louapre Magliaro, Martin-Benito, Meusburger, Mikovic, Modesto, Noui, Oeckl, Oriti, Perez, Pereira, Perini, Pfeiffer, Pullin, Raasakka, Reisenberger, Riello, Roche, Rovelli, Ryan, Smerlak, Speziale, Steinhaus, Thiemann, Vidotto, Vojinovic, Wieland, Wilson-Ewing, Zapata, Zhang, Zipfel ...

## Spin foam models

Space time picture  
for operators on  
graph Hilbert spaces  
(Loop quantum gravity)



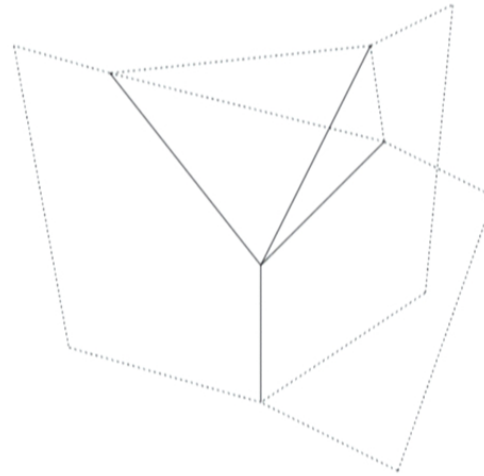
Feynman diagrams for  
a type of non local  
Interaction (Group field theory)

Fiducial structure for  
the construction of  
topological quantum field theories

Lattice gauge gravity

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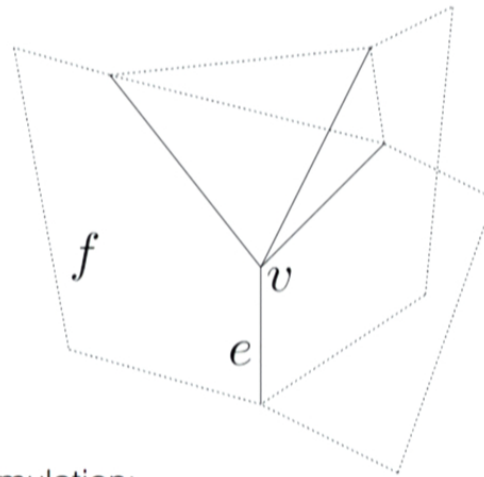


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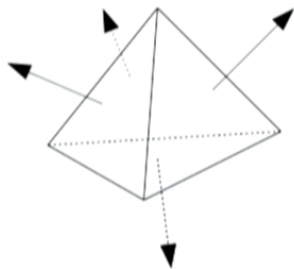


State sum formulation:

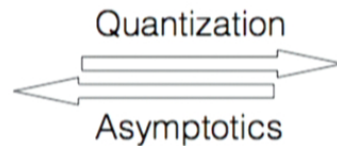
$$\mathcal{Z} = \sum_{j_f, i_e} \prod_f A_f(j_f) \prod_e A_e(i_e) \prod_v A_v(j_f, i_e)$$

Vertex asymptotics:

$$\mathcal{Z} = \sum_{j_f, i_e} \prod_f A_f(j_f) \prod_e A_e(i_e) \prod_v A_v(j_f, i_e)$$



Phase space of 3d geometry



$$i_e \in \text{Inv}(j_f \otimes j_{f'} \otimes j_{f''} \otimes j_{f'''})$$

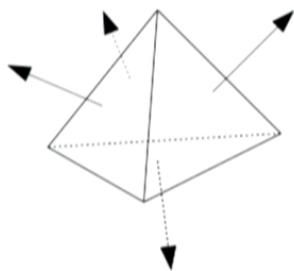
SU(2) representation theoretic data

E.g. coherent tetrahedral states  
ala Livine Speziale '07

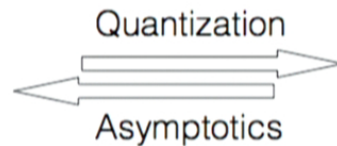
$$\psi_e(j_f, n_f)$$

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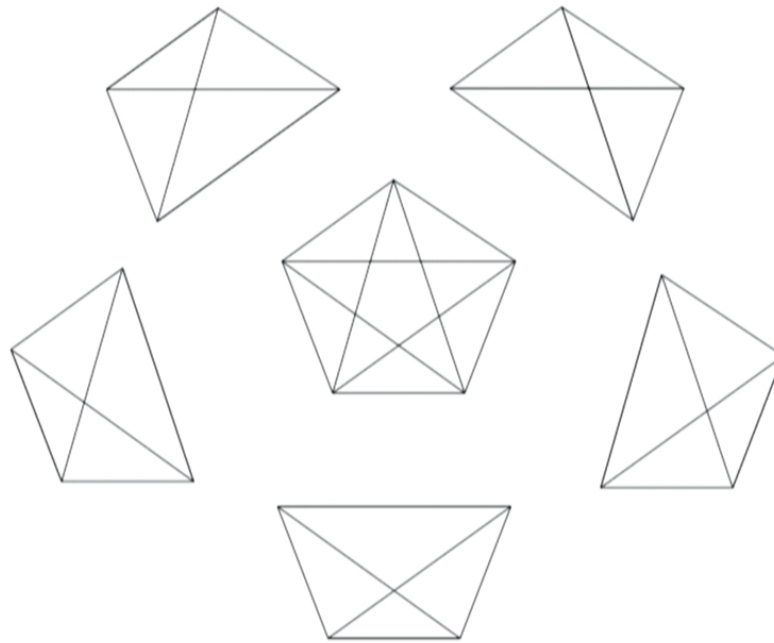
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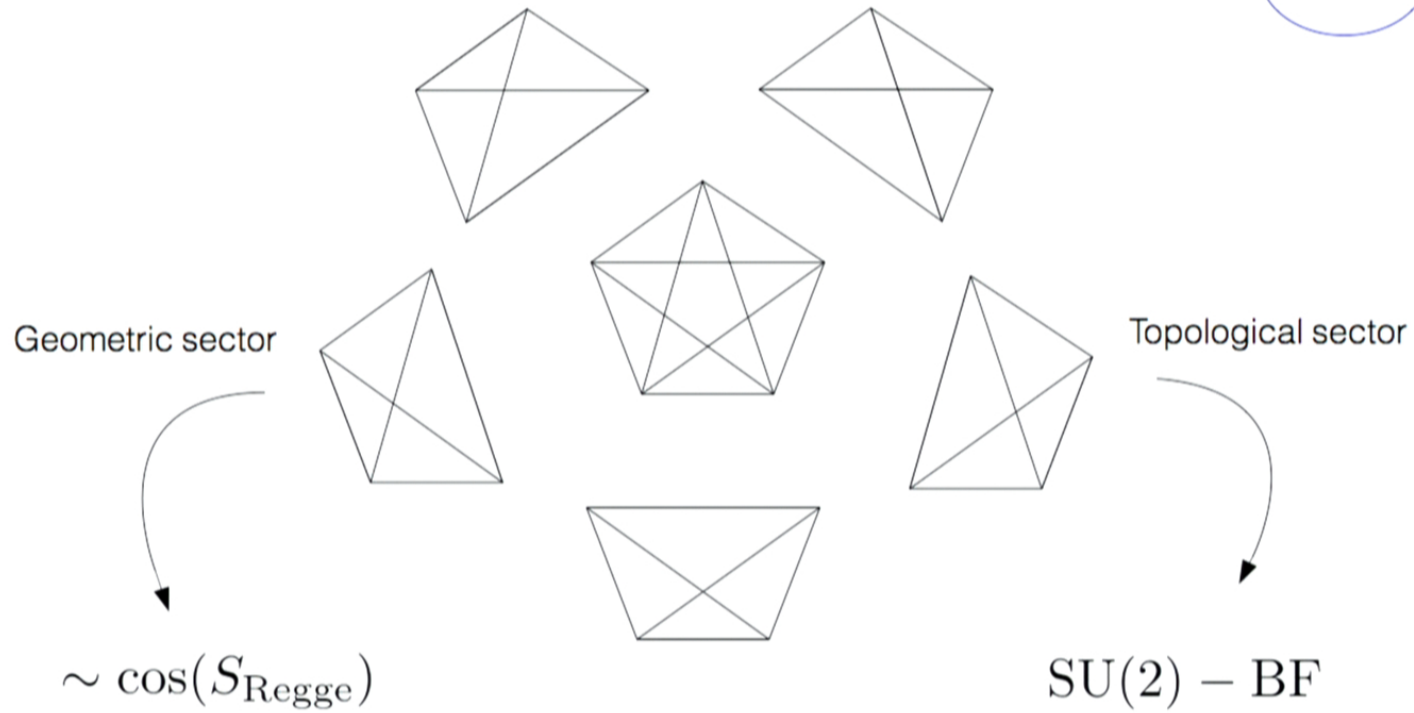
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Based on these were calculated observables on the single simplex (and similar simple triangulations):

Propagator: Alesci, Bianchi, Ding, Magliaro, Modesto, Perini, Rovelli, Speziale

Dipole boundaries: Rovelli, Vidotto, Bianchi, FH, Kisielowski, Lewandowski, Puchta

Effective action: Mikovic, Vojinovic

Radiative corrections: Riello

This tests: The geometricity. The boundary action. The operators.

This does not test: The dynamics. Continuum physics.

It is unclear how to treat the sum over representation.

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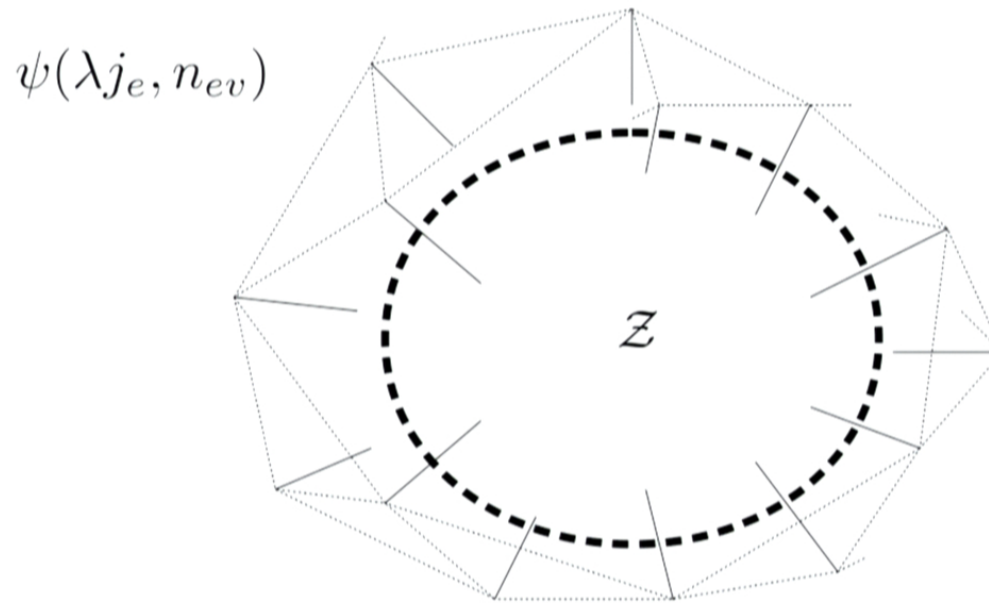
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Asymptotic dynamics:

$$\langle \psi(\lambda j_e, n_{ev}) | \mathcal{Z} \rangle$$

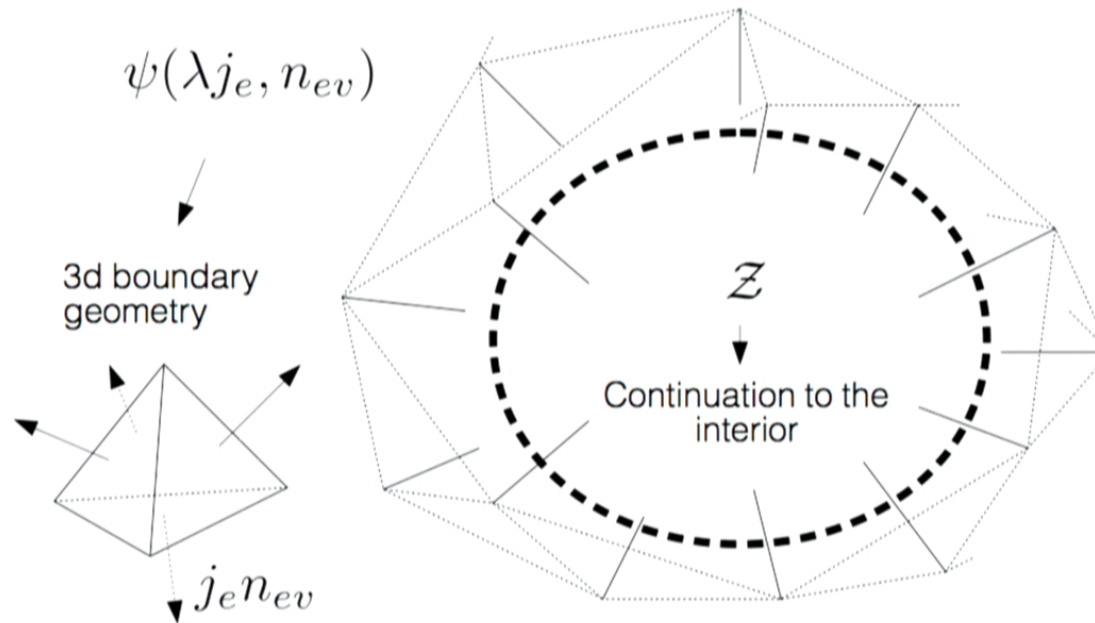


Asymptotic boundary state, no assumption on the interior.



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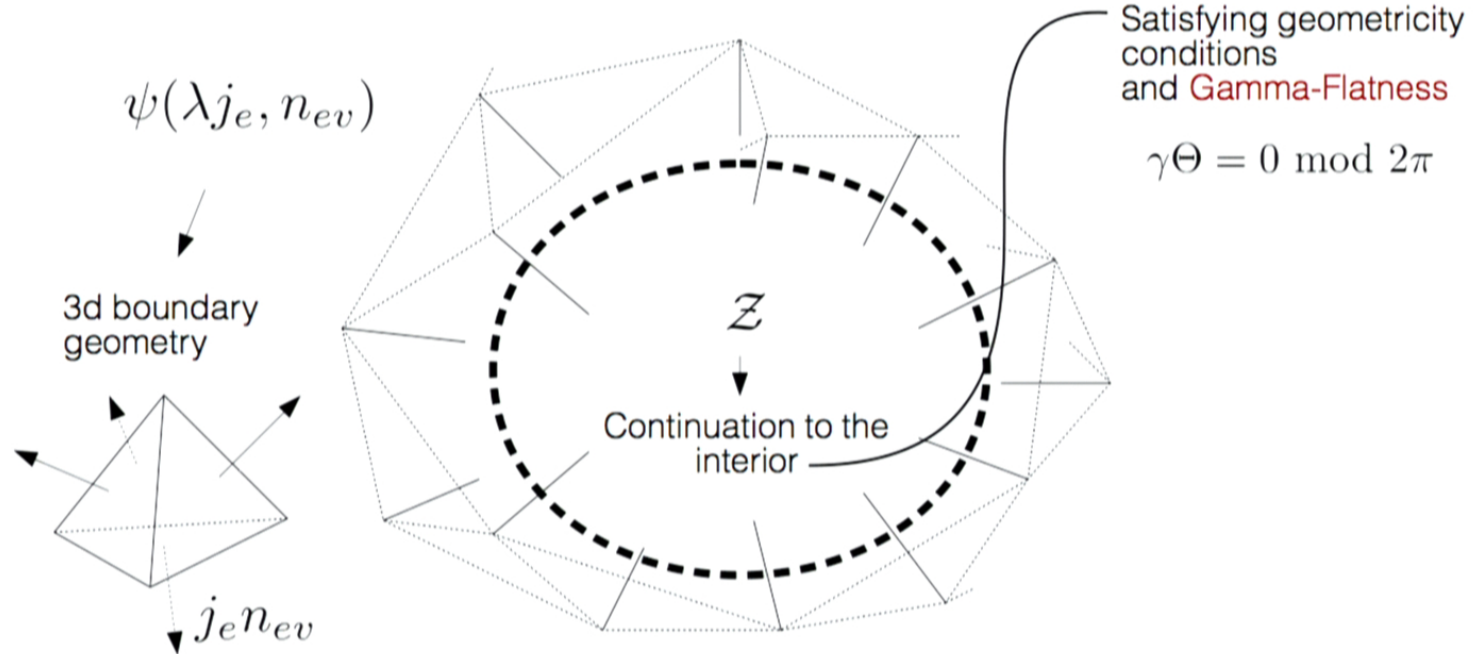
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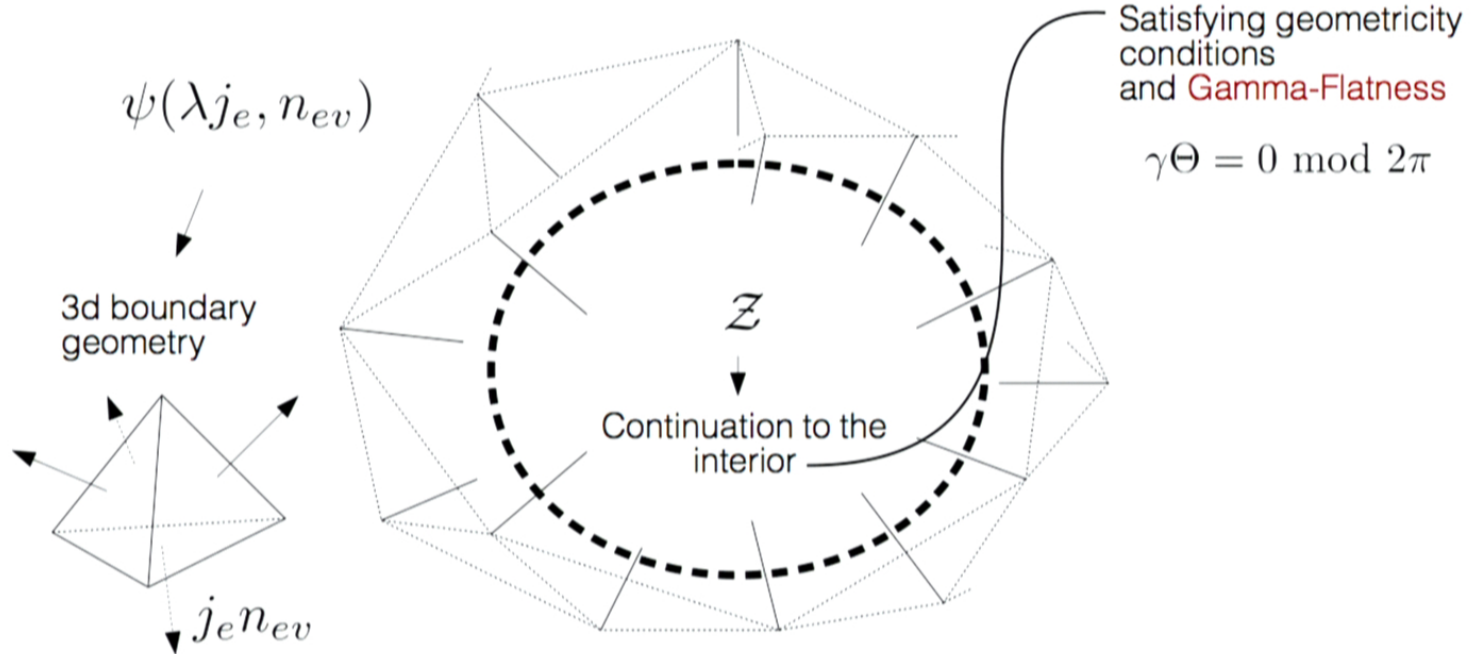
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Theorem (FH, Kaminski):

Partition function exists if the interior continuation is bounded whenever it exists, is asymptotically suppressed if the boundary data allows no continuation.

Interior continuation is given by

$$\{p_{ee'}^v, p_{fv}^e, g_{ev}, g_{ef}\}$$

satisfying geometricity conditions:

$$p_{ee'}^v = -p_{e'e}^v$$

$$p_{fv}^e = -p_{fv}^e$$

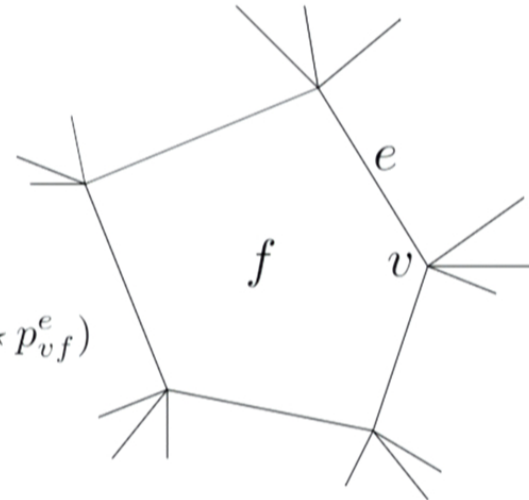
$$p_{fv}^e = g_{ev} p_{ee'}^v$$

$$N^0 \cdot (1 + \gamma \star) p_{fv}^e = 0$$

$$g_{ef} = \exp(\xi_{ef} \star p_{vf}^e)$$

$$\sum_f p_{fv}^e = 0$$

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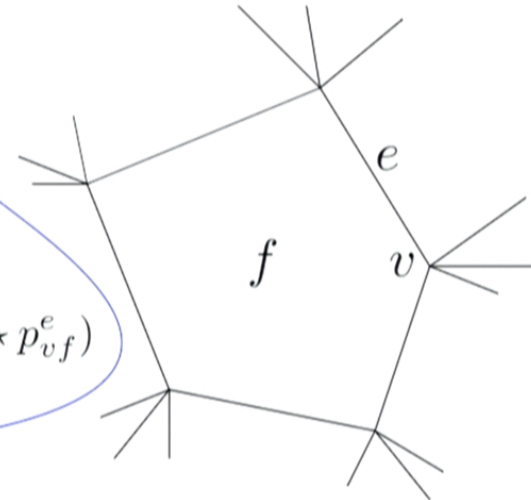
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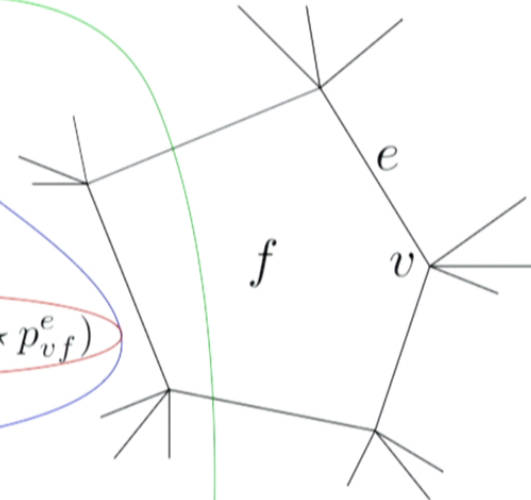
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Iff? Extra equations? (Regge?)

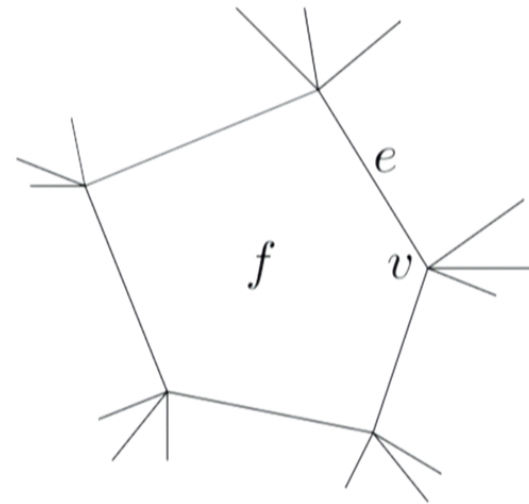


## Holonomy spin foams. The dual formulation. $\langle \psi(\lambda_{j_e}, n_{ev}) | \mathcal{Z} \rangle$

Idea: Perform the spin sum explicitly, obtain formulation in terms of group elements.

$$\mathcal{Z} = \int dg_{ef} dg_{ev} \prod_{ef} E(g_{ef}) \prod_f \omega(\tilde{g}_f)$$

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Bahr, Dittrich, FH, Kaminski '12-13.  
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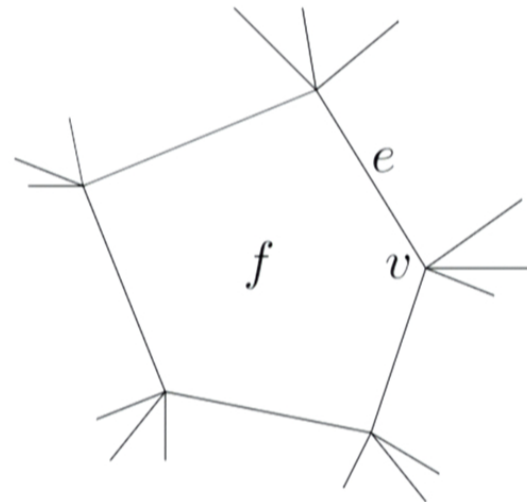
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The important part lives here.

Measure factors live here.



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## Method: Wave front sets

How does having functions on the group help with the large  $j$  limit? Take  $U(1)$ :

$$f(e^{i\phi}) = \sum_n f^{(n)} e^{i\phi/n}$$

large representations  $\sim$  short wavelength on the group  $\sim$  local structure  $\sim$  smoothness

Wave front sets are a refinement of the notion of the singular support of a distribution to the cotangent bundle (phase space) of the space (configuration space, the group manifold) on which the distribution is defined. They probe the same regime as the large representation limit does.

This is why the holonomy formulation allows us to study the large  $j$  limit:

The large  $j$  limit is captured by the failure of distributions to be smooth.  
The failure of smoothness propagates in a natural way under convolution.

(Really: large covector limit for a very general class of coherent states.)

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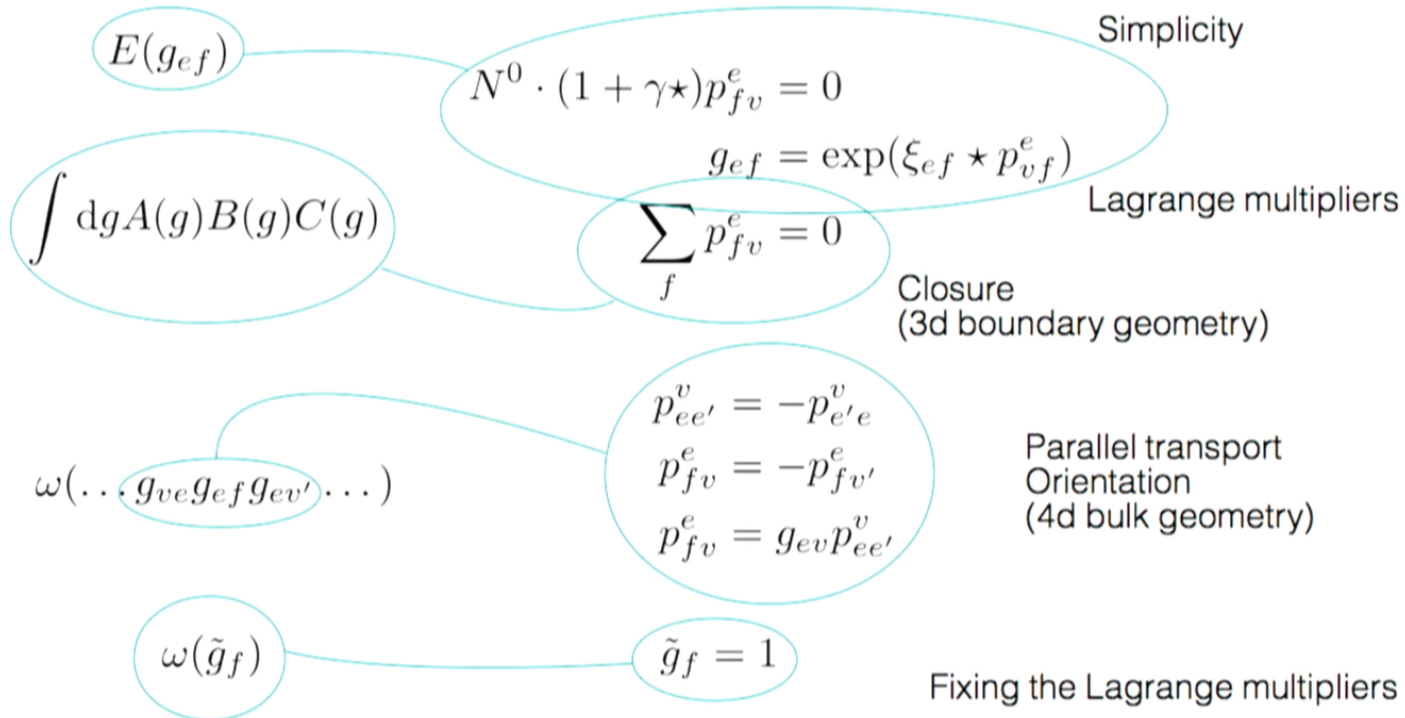
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### Flatness:

First reported by Bonzom '11 (Path integral, no gamma)  
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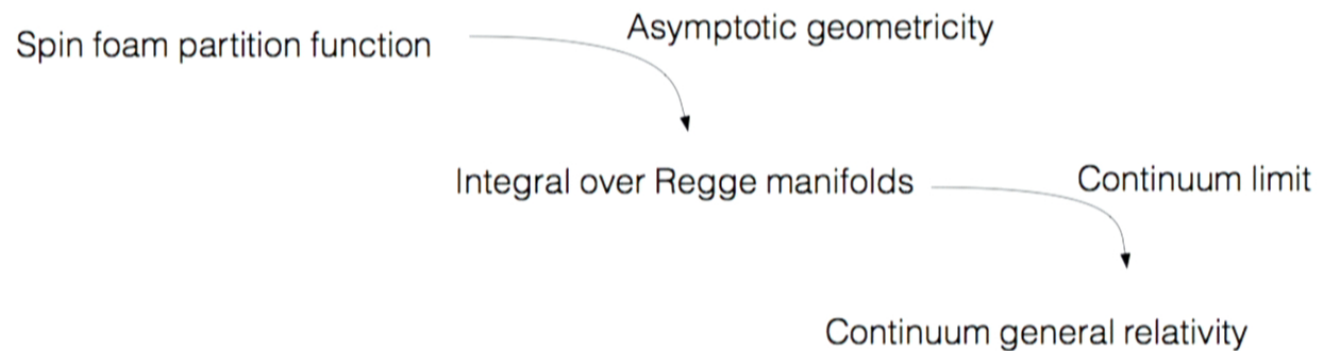
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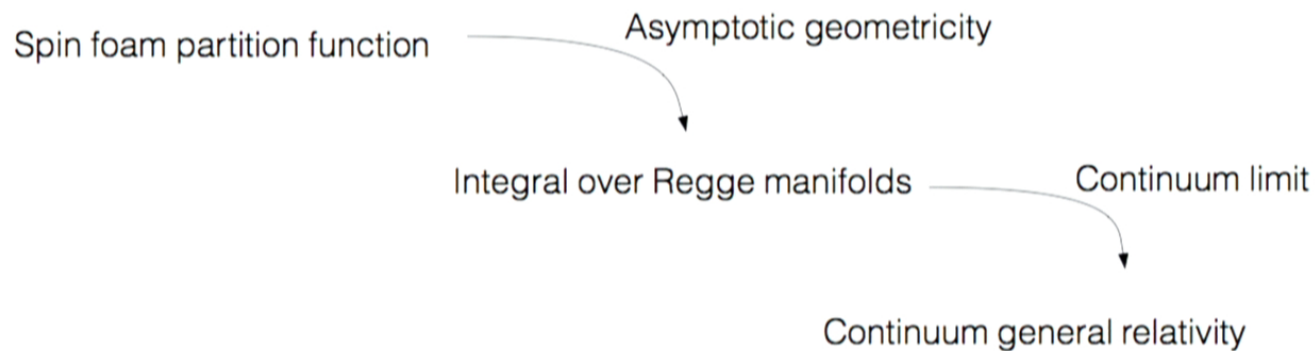


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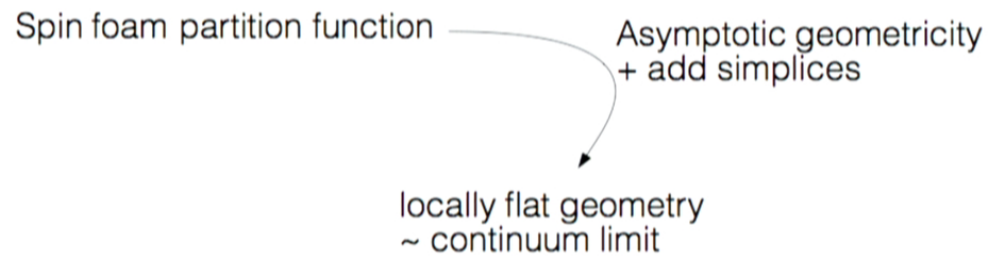
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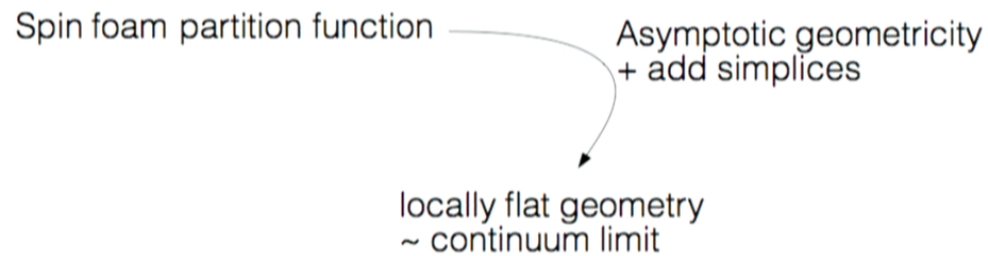
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## What's next:

We have:

Sufficient condition for existence of the partition function.

Necessary condition for non-suppressed-ness of boundary data.

Can we have:

Necessary condition for existence? Degree of non-existence?

Recent works by Riello, Dittrich, Bonzom, as well as previous work on measure factors by Perez and Rovelli suggest we are far from optimal.

Relevance for GFT? Ala Riello? (Ben Geloun, Gurau, Rivasseau, et.al.)

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What equations of motion (if any) do our discrete geometries satisfy?

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