

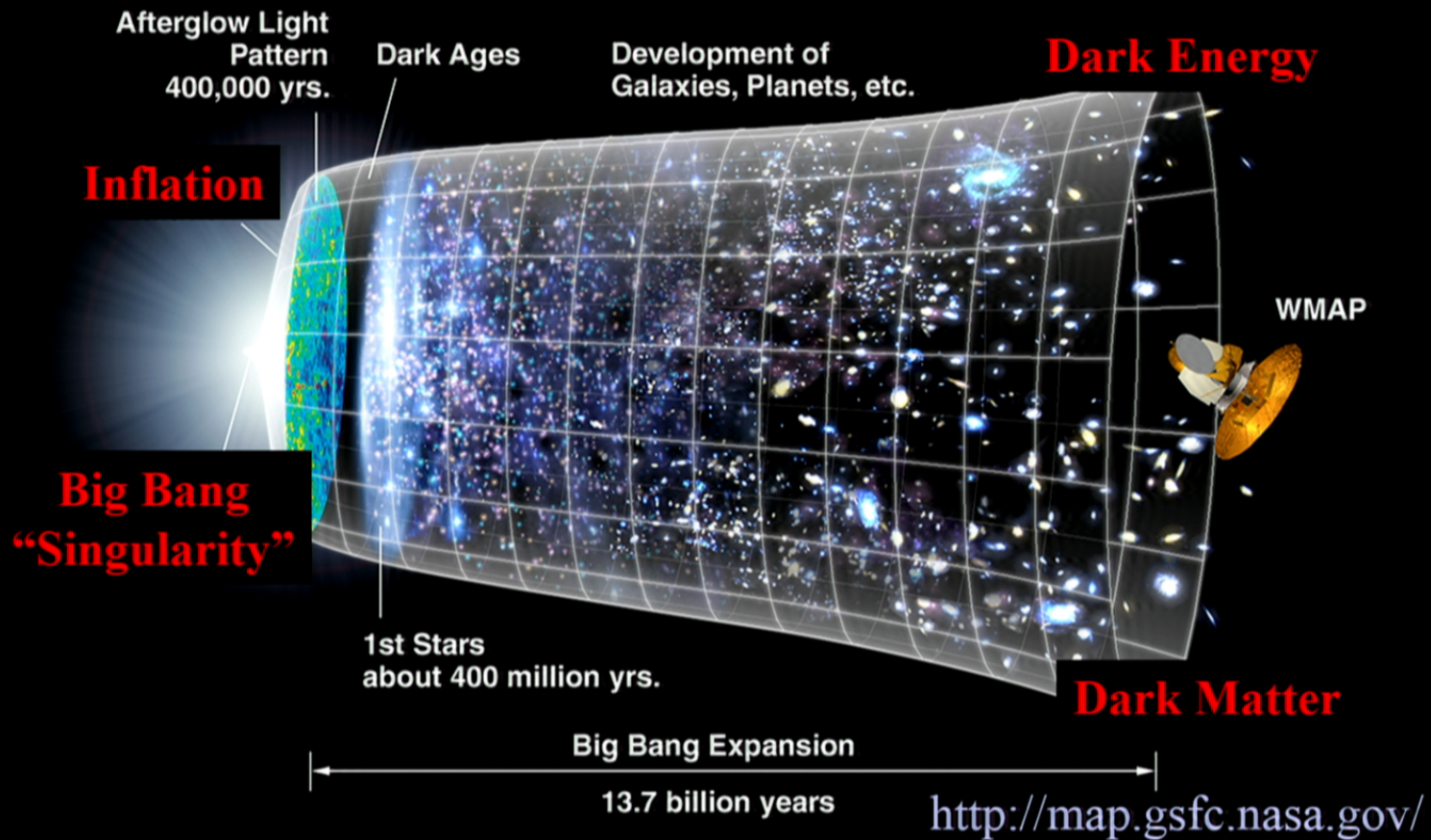
Title: Massive gravity and cosmology

Date: Jul 08, 2013 11:40 AM

URL: <http://pirsa.org/13070002>

Abstract:

Why alternative gravity theories?



Three conditions for good alternative theories of gravity (my personal viewpoint)

1. Theoretically consistent
e.g. no ghost instability
2. Experimentally viable
solar system / table top experiments
3. Predictable
e.g. protected by symmetry

Some examples

- I. Ghost condensation
IR modification of gravity
motivation: dark energy/matter
- II. Nonlinear massive gravity
IR modification of gravity
motivation: “Can graviton have mass?”
- III. Horava-Lifshitz gravity
UV modification of gravity
motivation: quantum gravity
- IV. Superstring theory
UV modification of gravity
motivation: quantum gravity, unified theory

A motivation for IR modification

- Gravity at long distances
 - Flattening galaxy rotation curves
 - extra gravity
 - Dimming supernovae
 - accelerating universe



A motivation for IR modification

- **Gravity at long distances**
Flattening galaxy rotation curves
extra gravity
Dimming supernovae
accelerating universe
- Usual explanation: **new forms of matter (DARK MATTER) and energy (DARK ENERGY).**

Dark component in the solar system?

Precession of perihelion
observed in 1800's...



Dark component in the solar system?

Precession of perihelion
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which people tried to
explain with a “dark
planet”, Vulcan,



But the right answer wasn't “dark planet”, it was
“change gravity” from Newton to GR.

Can we change gravity in IR?

➤ Change Theory?

Massive gravity

Fierz-Pauli 1939

DGP model

Dvali-Gabadadze-Porrati 2000



Can we change gravity in IR?

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➤ Change State?

Higgs phase of gravity

The simplest: Ghost condensation

Arkani-Hamed, Cheng, Luty and Mukohyama, JHEP 0405:074,2004.

Linear massive gravity (Fierz-Pauli 1939)

- Simple question: Can spin-2 field have mass?
- $L = L_{EH}[h] + m_g^2 [\eta^{\mu\rho}\eta^{\nu\sigma} h_{\mu\nu} h_{\rho\sigma} - (\eta^{\mu\nu} h_{\mu\nu})^2]$
 $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$
- Unique linear theory without ghosts
- Broken diffeomorphism
 - no momentum constraint
 - 5 d.o.f. (2 tensor + 2 vector + 1 scalar)

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Massive gravity: history

Simple question: Can graviton have mass?


May lead to acceleration without dark energy

Yes?

No?

Fierz-Pauli theory (1939)

Unique linear theory
without instabilities
(ghosts)

A grey seesaw is shown on a triangular fulcrum. The left side of the seesaw is higher and has a blue box with white text on it. The right side of the seesaw is lower and is empty.

vDVZ vs Vainshtein

- van Dam-Veltman-Zhakharov (1970)
Massless limit \neq Massless theory = GR
5 d.o.f remain \rightarrow PPN parameter $\gamma = \frac{1}{2} \neq 1$
- Vainshtein (1972)
Linear theory breaks down in the limit.
Nonlinear analysis shows continuity and GR is recovered @ $r < r_V = (r_g/m_g^4)^{1/5}$.
Continuity is not uniform w.r.t. distance.

vDVZ vs Vainshtein

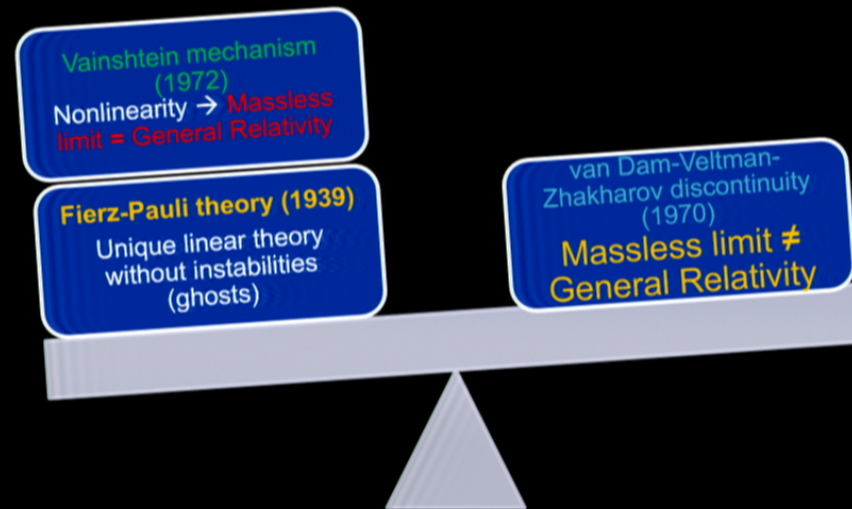
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Naïve nonlinear theory and BD ghost

- FP theory with $\eta^{\mu\nu} \rightarrow g^{\mu\nu}$
$$L = L_{EH}[h] + m_g^2 [g^{\mu\rho} g^{\nu\sigma} h_{\mu\nu} h_{\rho\sigma} - (g^{\mu\nu} h_{\mu\nu})^2]$$
$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$$
- Vainshtein effect (1972)
- Boulware-Deser ghost (1972)
No Hamiltonian constraint @ nonlinear level
 \rightarrow 6 d.o.f. = 5 d.o.f. of massive spin-2 + 1 ghost

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Vainshtein mechanism
(1972)
Nonlinearity \rightarrow Massless
limit = General Relativity

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Nonlinear massive gravity

de Rham, Gabadadze 2010

de Rham, Gabadadze & Tolley 2010

- First example of fully nonlinear massive gravity without BD ghost since 1972!
- Purely classical
- Properties of 5 d.o.f. depend on background
- **4 scalar fields ϕ^a ($a=0,1,2,3$)**
- **Poincare symmetry in the field space:**

$$\phi^a \rightarrow \phi^a + c^a, \quad \phi^a \rightarrow \Lambda_b^a \phi^b$$



$$f_{\mu\nu} \equiv \eta_{ab} \partial_\mu \phi^a \partial_\nu \phi^b$$

Pullback of
Minkowski metric in field space
to spacetime

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Systematic resummation

de Rham, Gabadadze & Tolley 2010

$$I_{mass}[g_{\mu\nu}, f_{\mu\nu}] = M_{Pl}^2 m_g^2 \int d^4x \sqrt{-g} (\mathcal{L}_2 + \alpha_3 \mathcal{L}_3 + \alpha_4 \mathcal{L}_4)$$

$$f_{\mu\nu} \equiv \eta_{ab} \partial_\mu \phi^a \partial_\nu \phi^b$$

$$\mathcal{K}_\nu^\mu = \delta_\nu^\mu - \left(\sqrt{g^{-1} f} \right)_\nu^\mu$$

$$\mathcal{L}_2 = \frac{1}{2} ([\mathcal{K}]^2 - [\mathcal{K}^2])$$

$$\mathcal{L}_3 = \frac{1}{6} ([\mathcal{K}]^3 - 3 [\mathcal{K}] [\mathcal{K}^2] + 2 [\mathcal{K}^3]) \quad [\mathcal{A}] \equiv Tr \mathcal{A}$$

$$\mathcal{L}_4 = \frac{1}{24} ([\mathcal{K}]^4 - 6 [\mathcal{K}]^2 [\mathcal{K}^2] + 3 [\mathcal{K}^2]^2 + 8 [\mathcal{K}] [\mathcal{K}^3] - 6 [\mathcal{K}^4])$$

No helicity-0 ghost, i.e. no BD ghost, in decoupling limit

$$\mathcal{K}_{\mu\nu} = \partial_\mu \partial_\nu \pi \quad \longrightarrow \quad \mathcal{L}_{2,3,4} = (\text{total derivative})$$

No BD ghost away from decoupling limit (Hassan&Rosen)

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Massless limit \neq General Relativity

No FLRW universe?

D'Amico, de Rham, Dubovsky, Gabadadze, Pirtshalava, Tolley (2011)

- Flat FLRW ansatz in “Unitary gauge”

$$g_{\mu\nu} dx^\mu dx^\nu = -N^2(t) dt^2 + a^2(t) (dx^2 + dy^2 + dz^2)$$

$$\phi^a = x^a \quad \rightarrow \quad f_{\mu\nu} = \eta_{\mu\nu}$$

- Bianchi “identity” $\rightarrow a(t) = \text{const.}$

$$\text{c.f.} \quad \nabla^\mu \left(\frac{2}{\sqrt{-g}} \frac{\delta I}{\delta g^{\mu\nu}} \right) = \frac{1}{\sqrt{-g}} \frac{\delta I_g}{\delta \phi^a} \partial_\nu \phi^a$$

\rightarrow no non-trivial flat FLRW cosmology

- “Our conclusions on the absence of the homogeneous and isotropic solutions do not change if we allow for a more general maximally symmetric 3-space”

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Consistent Theory
found in 2010 but

No Viable Cosmology?

de Sitter problem
"ghost" (1976)
First example of nonlinear
massive gravity without
BD ghost since 1976

de Rham, Gabriellini,
Ghosh (2010)

Nonlinearity \rightarrow Massless
limit \neq General Relativity

Fierz-Pauli theory (1939)
Unique linear theory
without instabilities
(ghosts)

D'Ambrosio et al. (2000)
No exact de Sitter
FRW (homogeneous
isotropic) universe!

Boulware-Deser ghost
(1972)

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 \rightarrow instability (ghost)

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Massless limit \neq
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Open FLRW solutions

Gumrukcuoglu, Lin, Mukohyama, arXiv: 1109.3845 [hep-th]

- $f_{\mu\nu}$ spontaneously breaks diffeo.
- Both $g_{\mu\nu}$ and $f_{\mu\nu}$ must respect FLRW symmetry
- Need FLRW coordinates of Minkowski $f_{\mu\nu}$

- No closed FLRW chart

$$\phi^0 = f(t)\sqrt{1 + |K|(x^2 + y^2 + z^2)},$$

$$\phi^1 = \sqrt{|K|}f(t)x,$$

$$\phi^2 = \sqrt{|K|}f(t)y,$$

$$\phi^3 = \sqrt{|K|}f(t)z.$$

- Open FLRW ansatz

$$f_{\mu\nu}dx^\mu dx^\nu = -(\dot{f}(t))^2 dt^2 + |K| (f(t))^2 \Omega_{ij}(x^k) dx^i dx^j$$

$$g_{\mu\nu}dx^\mu dx^\nu = -N(t)^2 dt^2 + a(t)^2 \Omega_{ij} dx^i dx^j,$$

$$\Omega_{ij} dx^i dx^j = dx^2 + dy^2 + dz^2 - \frac{|K|(xdx + ydy + zdz)^2}{1 + |K|(x^2 + y^2 + z^2)},$$

Open FLRW solutions

Gumrukcuoglu, Lin, Mukohyama, arXiv: 1109.3845 [hep-th]

- EOM for ϕ^a ($a=0,1,2,3$)

$$(\dot{a} - \sqrt{|K|N}) \left[\left(3 - \frac{2\sqrt{|K|f}}{a} \right) + \alpha_3 \left(3 - \frac{\sqrt{|K|f}}{a} \right) \left(1 - \frac{\sqrt{|K|f}}{a} \right) + \alpha_4 \left(1 - \frac{\sqrt{|K|f}}{a} \right)^2 \right] = 0$$

- The first sol $\dot{a} = \sqrt{|K|N}$ implies $g_{\mu\nu}$ is Minkowski
 → we consider other solutions

$$f = \frac{a}{\sqrt{|K|}} X_{\pm}, \quad X_{\pm} \equiv \frac{1 + 2\alpha_3 + \alpha_4 \pm \sqrt{1 + \alpha_3 + \alpha_3^2 - \alpha_4}}{\alpha_3 + \alpha_4}$$

- Latter solutions do not exist if $K=0$
- Metric EOM → self-acceleration

$$3H^2 + \frac{3K}{a^2} = \Lambda_{\pm} + \frac{1}{M_{Pl}^2} \rho$$

$$\Lambda_{\pm} \equiv -\frac{m_g^2}{(\alpha_3 + \alpha_4)^2} \left[(1 + \alpha_3) (2 + \alpha_3 + 2\alpha_3^2 - 3\alpha_4) \pm 2 (1 + \alpha_3 + \alpha_3^2 - \alpha_4)^{3/2} \right]$$

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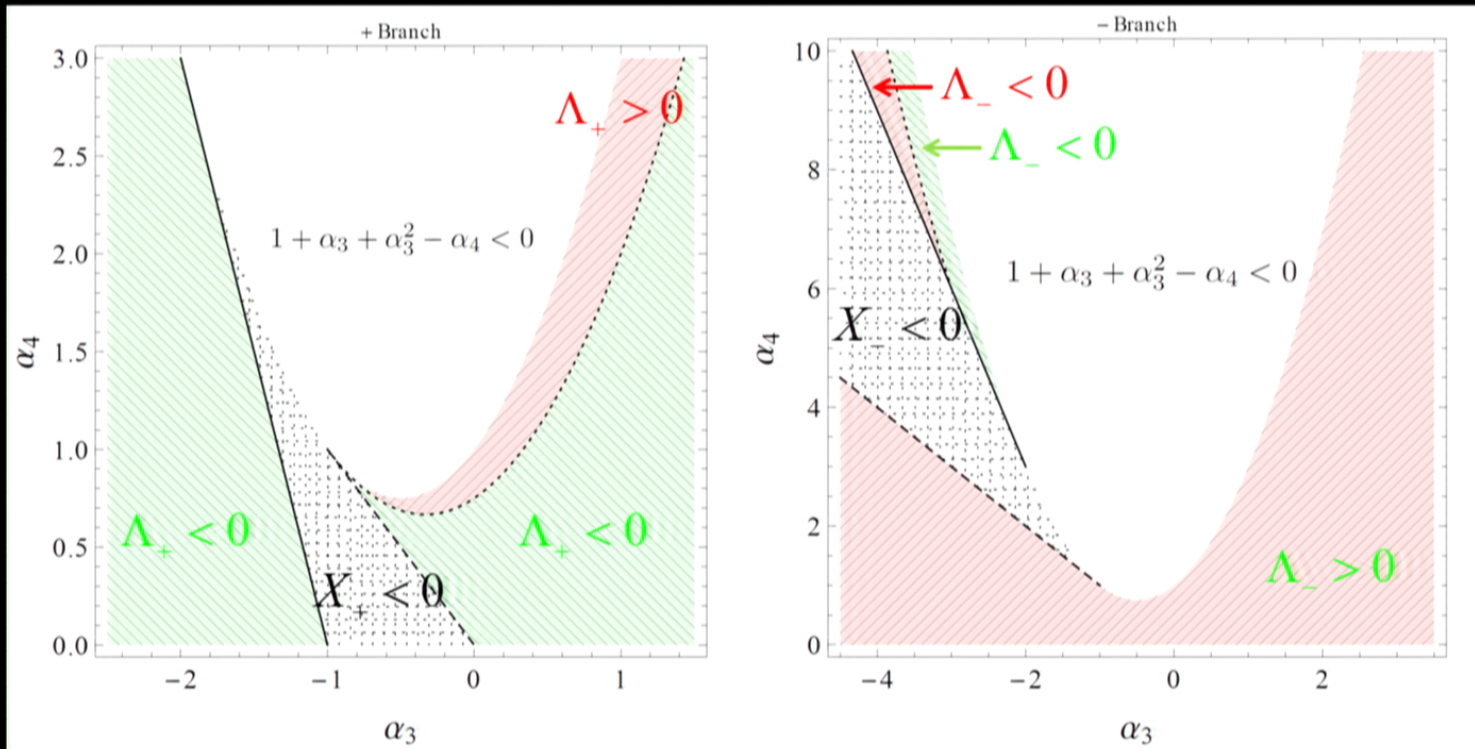
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Our recent contributions

Cosmological solutions of nonlinear massive gravity

Good?

Bad?

More general fiducial
metric $f_{\mu\nu}$
closed/flat/open FRW
universes allowed
GLM (2011b)

Open universes with self-
acceleration
GLM (2011a)

D'Amico, et.al. (2011)
Non-existence of flat
FRW (homogeneous
isotropic) universe!

GLM = Gumrukcuoglu-Lin-Mukohyama
DGM = DeFelice-Gumrukcuoglu-Mukohyama

Summary so far

- Nonlinear massive gravity
free from BD ghost
- FLRW background
No closed/flat universe
Open universes with self-acceleration!
- More general fiducial metric $f_{\mu\nu}$
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Friedmann eq does not depend on $f_{\mu\nu}$
- Cosmological linear perturbations
Scalar/vector sectors \rightarrow same as in GR
Tensor sector \rightarrow time-dependent mass

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Nonlinear instability

DeFelice, Gumrukcuoglu, Mukohyama, arXiv: 1206.2080 [hep-th]

- de Sitter or FLRW fiducial metric
- Pure gravity + bare cc \rightarrow FLRW sol = de Sitter
- Bianchi I universe with axisymmetry + linear perturbation (without decoupling limit)
- Small anisotropy expansion of Bianchi I + linear perturbation
 \rightarrow nonlinear perturbation around flat FLRW
- **Odd-sector:**
1 healthy mode + 1 healthy or ghosty mode
- **Even-sector:**
2 healthy modes + 1 ghosty mode
- This is not BD ghost nor Higuchi ghost.

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New class of cosmological solution

Gumrukcuoglu, Lin, Mukohyama, arXiv: 1206.2723 [hep-th]
+ De Felice, arXiv: 1303.4154 [hep-th]

- Healthy regions with (relatively) large anisotropy
- Are there attractors in healthy region?
- Classification of fixed points
- Local stability analysis
- Global stability analysis

At attractors, physical metric is isotropic but fiducial metric is anisotropic.

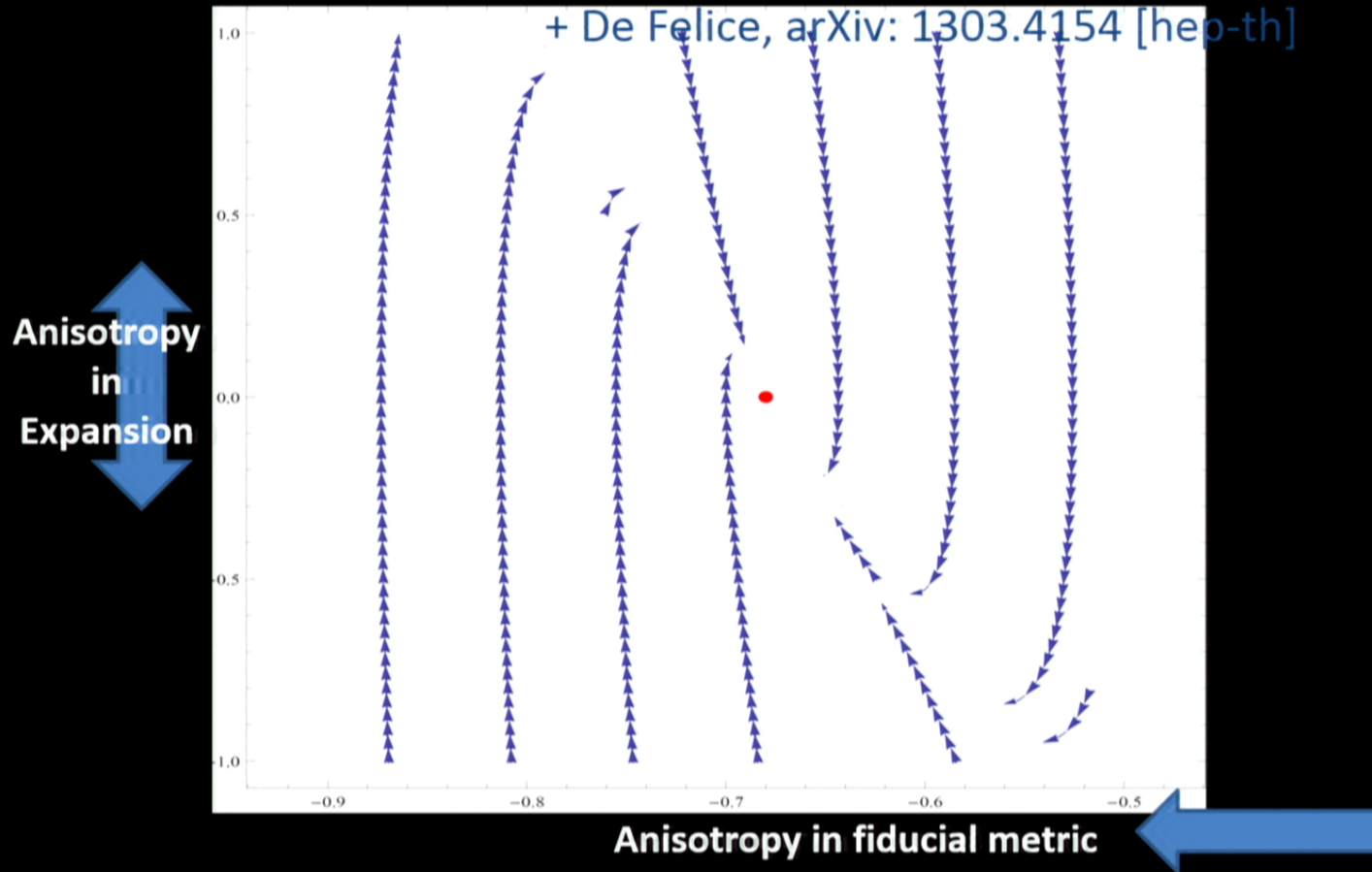
→ Anisotropic FLRW universe!

statistical anisotropy expected
(suppressed by small m_g^2)

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Our recent contributions

Cosmological solutions of nonlinear massive gravity

Anisotropic FRW:

Statistical anisotropy

(suppressed by fineness of graviton mass)

with

isotropic expansion

NEW Stable Solution:
GLM (2011)

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Linearized massive gravity
in FRW (non-homogeneous
isotropic) universe!

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Quasidilaton D'Amico, Gabadadze, Hui, Pirtskhalava, 2012

- New nonlinear instability [DeFelice, Gumrukcuoglu, Mukohyama 2012] → (i) new backgrounds, or (ii) extended theories

- Quasidilaton: scalar σ with global symmetry:

$$\sigma \rightarrow \sigma + \sigma_0 \quad \phi^a \rightarrow e^{-\sigma_0/M_{\text{Pl}}} \phi^a$$

- Action

$$S = \frac{M_{\text{Pl}}^2}{2} \int d^4x \sqrt{-g} \left[R - 2\Lambda - \frac{\omega}{M_{\text{Pl}}^2} \partial_\mu \sigma \partial^\mu \sigma + 2m_g^2 (\mathcal{L}_2 + \alpha_3 \mathcal{L}_3 + \alpha_4 \mathcal{L}_4) \right]$$

$$\mathcal{K}^\mu{}_\nu = \delta^\mu{}_\nu - e^{\sigma/M_{\text{Pl}}} \left(\sqrt{g^{-1} f} \right)^\mu{}_\nu \quad f_{\mu\nu} = \eta_{ab} \partial_\mu \phi^a \partial_\nu \phi^b$$

- **Scaling solution = self-accelerating de Sitter**
($H = \text{const} > 0$ with $\Lambda = 0$)

Stable extension of quasidilaton

arXiv: 1306.5502 [hep-th] /w A. De Felice

- Self-accelerating solution in the original quasidilaton theory has ghost instability

[Gumrukcuoglu, Hinterbichler, Lin, Mukohyama, Trodden 2013; D'Amico, Gabadadze, Hui, Pirtskhalava 2013]

- Simple extension: $f_{\mu\nu} \rightarrow \tilde{f}_{\mu\nu}$

$$\tilde{f}_{\mu\nu} \equiv f_{\mu\nu} - \frac{\alpha_\sigma}{M_{\text{Pl}}^2 m_g^2} e^{-2\sigma/M_{\text{Pl}}} \partial_\mu \sigma \partial_\nu \sigma$$

- Self-accelerating solution is stable if

$$0 < \omega < 6$$

$$X^2 < \frac{\alpha_\sigma H^2}{m_g^2} < r^2 X^2$$

$$X \equiv \frac{e^{\bar{\sigma}/M_{\text{Pl}}}}{a}$$

$$M_{\text{GW}}^2 \equiv \frac{(r-1)X^3 m_g^2}{X-1} + \frac{\omega H^2 (rX + r - 2)}{(X-1)(r-1)} > 0$$

$$r \equiv \frac{n}{N} a$$

Stable extension of quasidilaton

arXiv: 1306.5502 [hep-th] /w A. De Felice

- Self-accelerating solution in the original quasidilaton theory has ghost instability

[Gumrukcuoglu, Hinterbichler, Lin, Mukohyama, Trodden 2013; D'Amico, Gabadadze, Hui, Pirtskhalava 2013]

- Simple extension: $f_{\mu\nu} \rightarrow \tilde{f}_{\mu\nu}$

$$\tilde{f}_{\mu\nu} \equiv f_{\mu\nu} - \frac{\alpha_\sigma}{M_{\text{Pl}}^2 m_g^2} e^{-2\sigma/M_{\text{Pl}}} \partial_\mu \sigma \partial_\nu \sigma$$

- Self-accelerating solution is stable if

$$0 < \omega < 6$$

$$X^2 < \frac{\alpha_\sigma H^2}{m_g^2} < r^2 X^2$$

$$X \equiv \frac{e^{\bar{\sigma}/M_{\text{Pl}}}}{a}$$

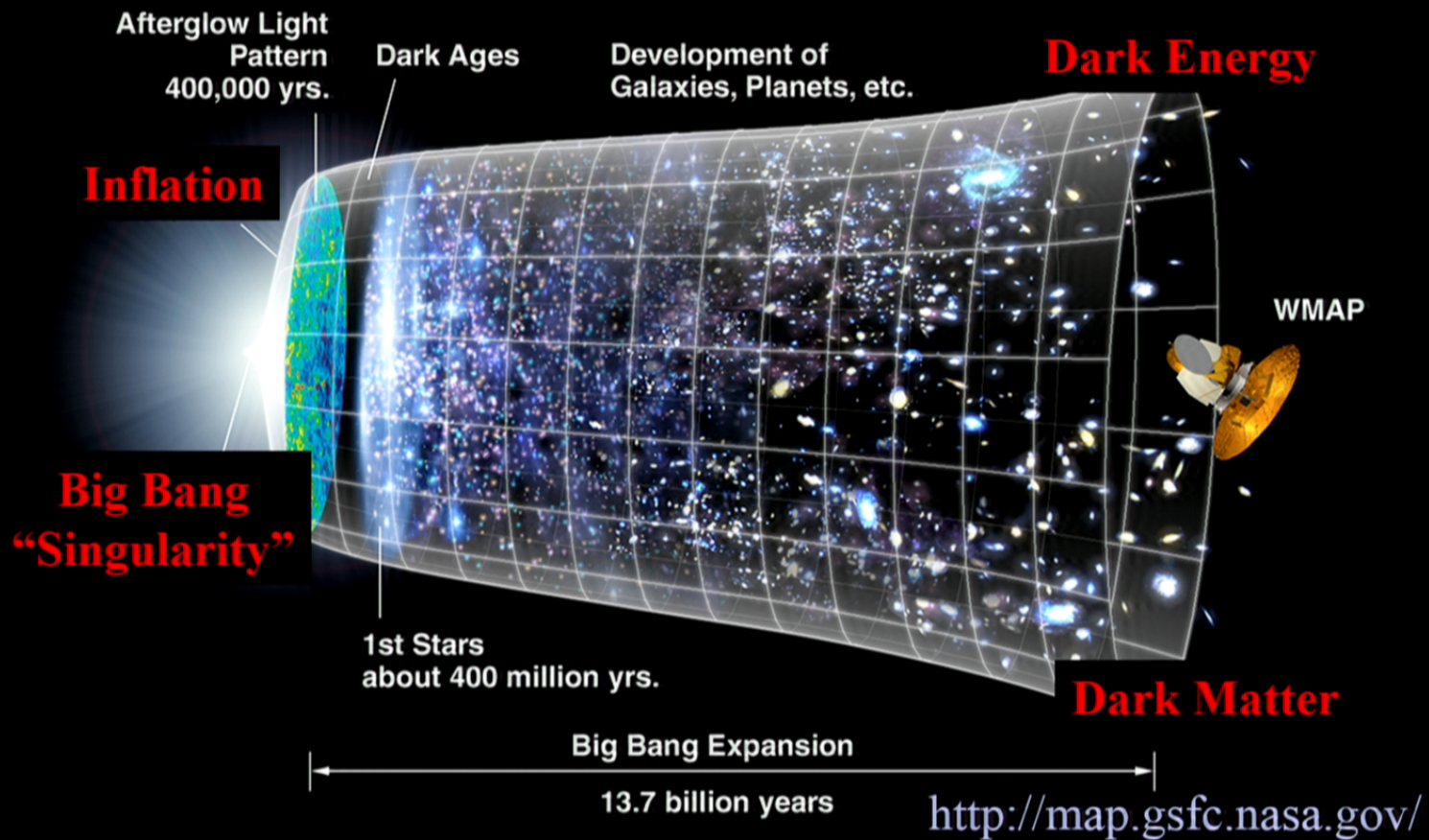
$$M_{\text{GW}}^2 \equiv \frac{(r-1)X^3 m_g^2}{X-1} + \frac{\omega H^2 (rX + r - 2)}{(X-1)(r-1)} > 0$$

$$r \equiv \frac{n}{N} a$$

Summary

- Nonlinear massive gravity
free from BD ghost
- FLRW background
No closed/flat universe
Open universes with self-acceleration!
- More general fiducial metric $f_{\mu\nu}$
closed/flat/open FLRW universes allowed
Friedmann eq does not depend on $f_{\mu\nu}$
- Cosmological linear perturbations
Scalar/vector sectors \rightarrow same as in GR
Tensor sector \rightarrow time-dependent mass
- All homogeneous and isotropic FLRW solutions have ghost
- New class of cosmological solution:
anisotropic FLRW \rightarrow statistical anisotropy
(suppressed by small m_g^2)
Analogue of Ghost Condensate!
- Extended quasidilaton: stable self-accelerating FLRW

Why alternative gravity theories?



General fiducial metric

Appendix of Gumrukcuoglu, Lin, Mukohyama, arXiv: 1111.4107 [hep-th]

- **Poincare symmetry in the field space**
→ $f_{\mu\nu} = (\text{Minkowski})_{ab} \partial_\mu \phi^a \partial_\nu \phi^b$
- **de Sitter symmetry in the field space**
→ $f_{\mu\nu} = (\text{deSitter})_{ab} \partial_\mu \phi^a \partial_\nu \phi^b$
- **FRW symmetry in the field space**
→ $f_{\mu\nu} = (\text{FLRW})_{ab} \partial_\mu \phi^a \partial_\nu \phi^b$

**Flat/closed/open FLRW cosmology allowed
if “fiducial metric” $f_{\mu\nu}$ is de Sitter (or FRW)**

→ Friedmann equation with the same effective cc

$$3H^2 + \frac{3K}{a^2} = \Lambda_\pm + \frac{1}{M_{Pl}^2} \rho$$

$$\Lambda_\pm \equiv -\frac{m_g^2}{(\alpha_3 + \alpha_4)^2} \left[(1 + \alpha_3) (2 + \alpha_3 + 2\alpha_3^2 - 3\alpha_4) \pm 2 (1 + \alpha_3 + \alpha_3^2 - \alpha_4)^{3/2} \right]$$