

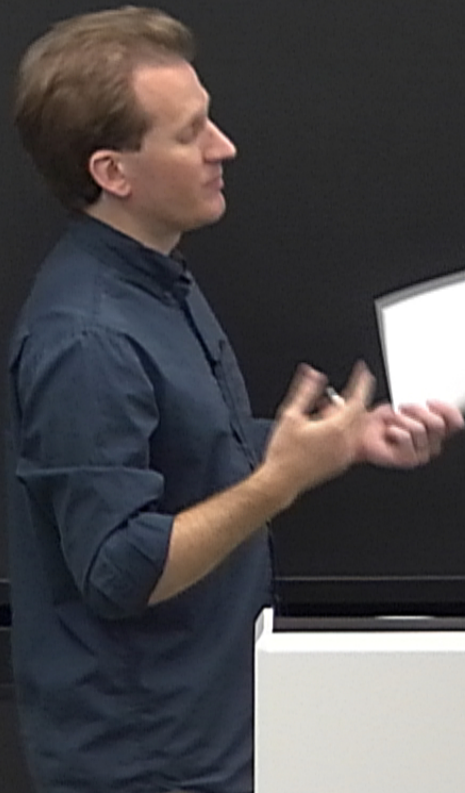
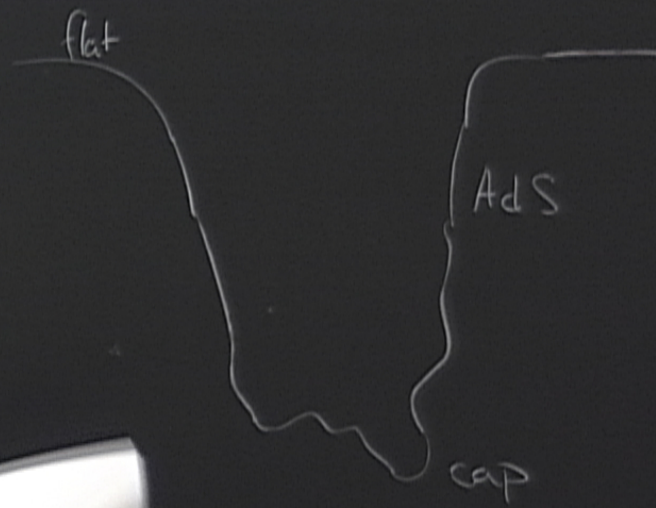
Title: Fuzzballs to Firewalls: A Post-Firewall Review of the Fuzzball Proposal: Lecture 3

Date: Jun 04, 2013 02:00 PM

URL: <http://pirsa.org/13060005>

Abstract: The fuzzball proposal makes a conjecture about the nature of black hole microstates. Now, more than a decade old and including several different philosophies and perspectives, it is especially relevant after the recent firewall argument and ensuing debate. Over three lectures, I plan to start with a very general discussion of the general ideas and motivations, then review the theoretical evidence from string theory, and finally close by discussing open questions, including the fate of a freely falling observer as he/she passes through the black hole horizon.

$\vec{F}(v) \longrightarrow$  SUGRA solution



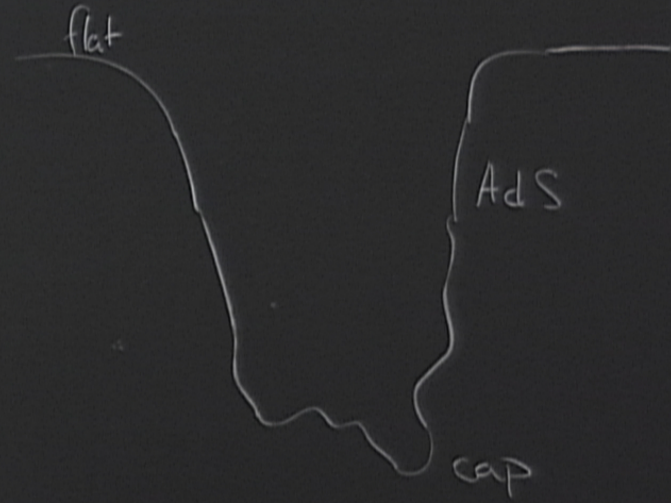


$\vec{F}(v) \longrightarrow$  SUGRA solution

Lunin-Mathur solutions



RR vac

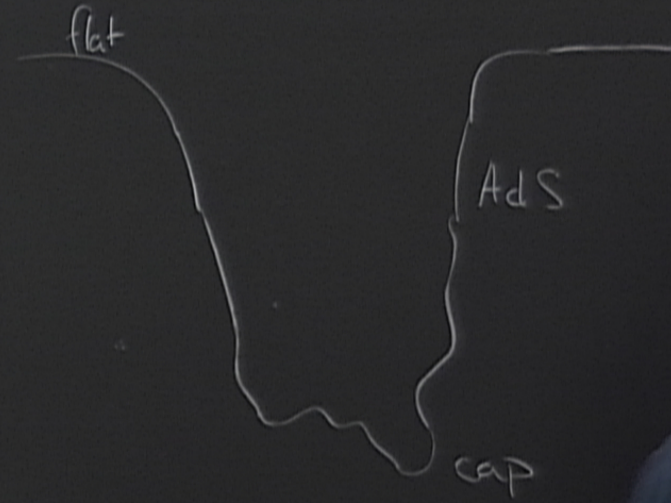


$\vec{F}(v) \longrightarrow$  SUGRA solution

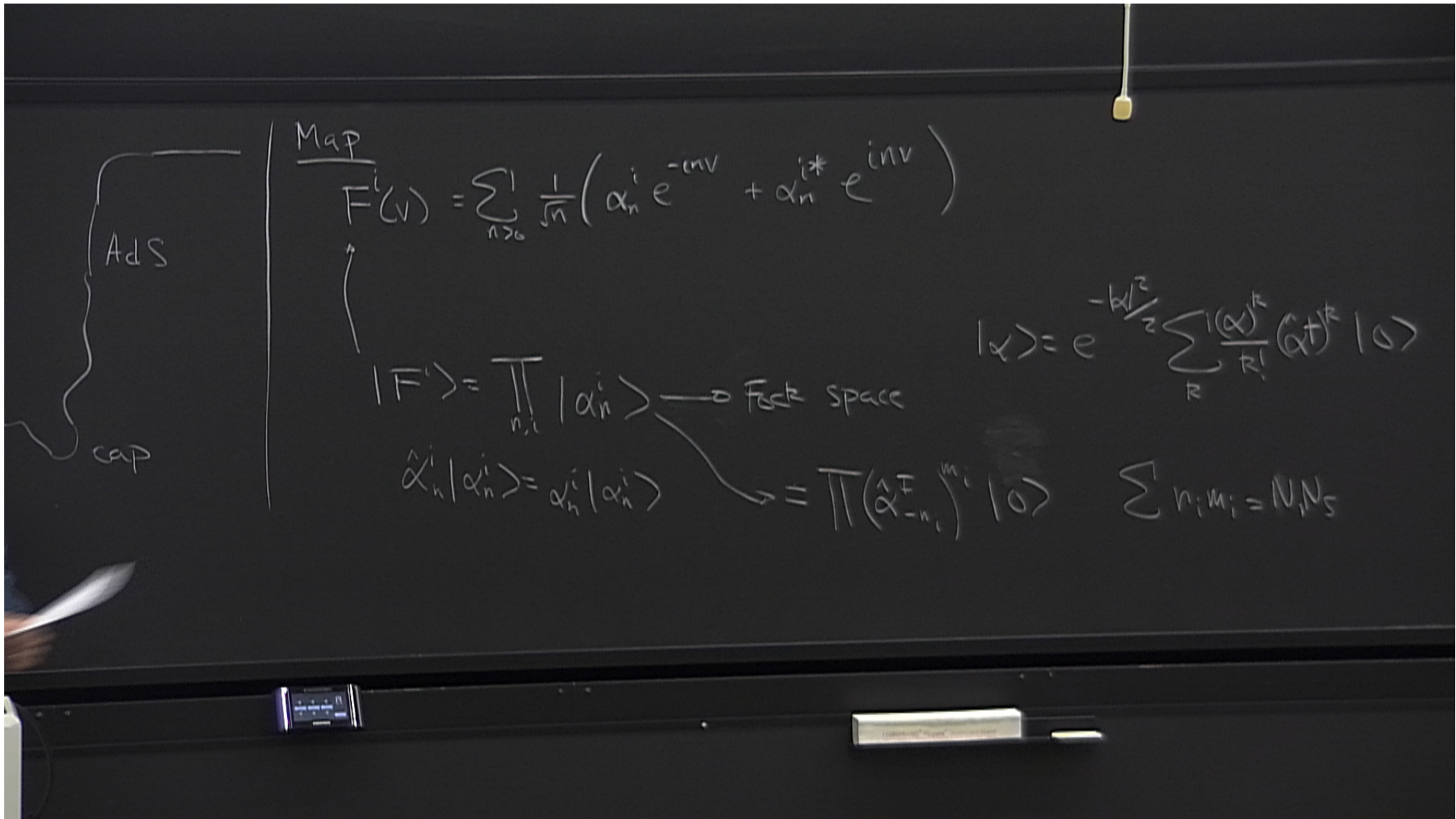
Lunin-Mathur solutions



RR vacuum of  $CFT_2$







AdS

cap

Map

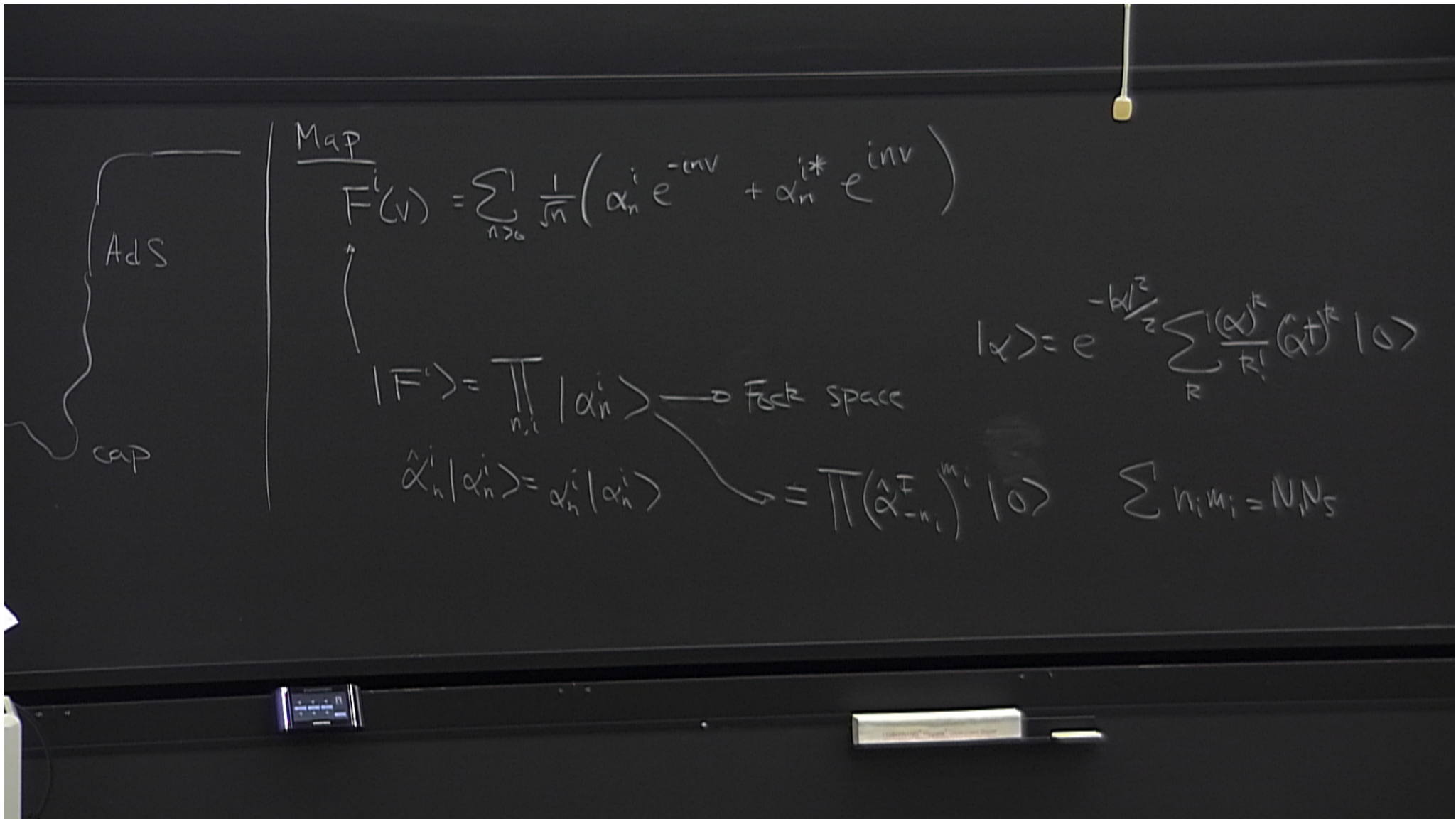
$$F^i(v) = \sum_{n \geq 0} \frac{1}{\sqrt{n}} \left( \alpha_n^i e^{-nv} + \alpha_n^{i*} e^{nv} \right)$$

$$|F^i\rangle = \prod_{n,i} |\alpha_n^i\rangle \rightarrow \text{Fock space}$$

$$\hat{\alpha}_n^i |\alpha_n^i\rangle = \alpha_n^i |\alpha_n^i\rangle \rightarrow = \prod (\hat{\alpha}_{-n_i}^i)^{m_i} |0\rangle$$

$$|\alpha\rangle = e^{-|\alpha|^2/2} \sum_{\mathbf{R}} \frac{(\alpha)^{\mathbf{R}}}{\mathbf{R}!} (\hat{\alpha}^\dagger)^{\mathbf{R}} |0\rangle$$

$$\sum n_i m_i = N_i N_5$$



AdS

cap

Map

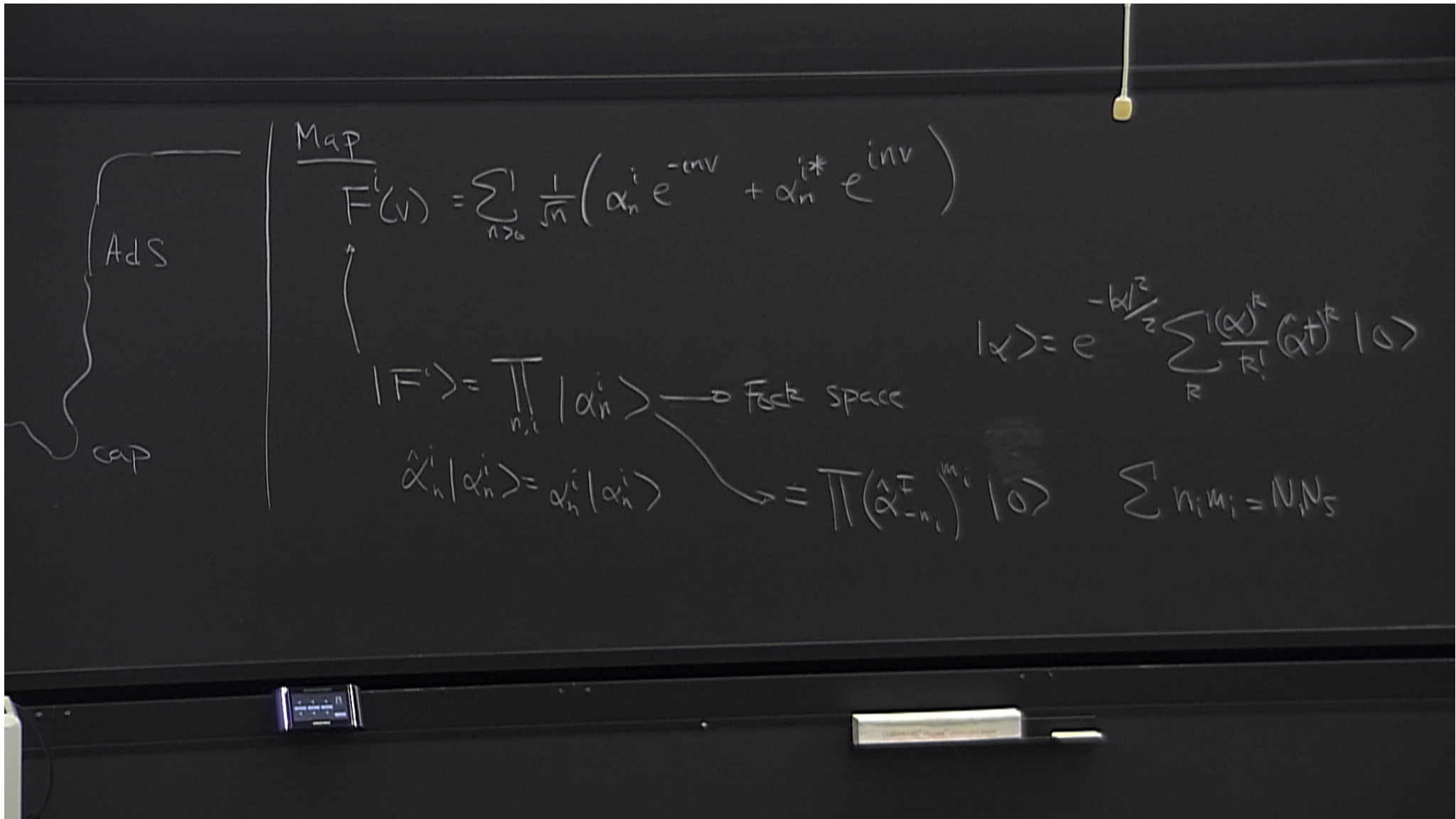
$$F^i(\nu) = \sum_{n \geq 0} \frac{1}{\sqrt{n}} \left( \alpha_n^i e^{-i\nu n} + \alpha_n^{i*} e^{i\nu n} \right)$$

$$|F^i\rangle = \prod_{n,i} |\alpha_n^i\rangle \rightarrow \text{Fock space}$$

$$\hat{\alpha}_n^i |\alpha_n^i\rangle = \alpha_n^i |\alpha_n^i\rangle \rightarrow = \prod (\hat{\alpha}_{-n_i}^i)^{m_i} |0\rangle \quad \sum n_i m_i = N_i N_5$$

$$|\alpha\rangle = e^{-|\alpha|^2/2} \sum_{\mathbf{r}} \frac{1}{\mathbf{r}!} (\hat{\alpha}^\dagger)^{\mathbf{r}} |0\rangle$$





AdS

cap

Map

$$F^i(v) = \sum_{n \geq 0} \frac{1}{\sqrt{n}} \left( \alpha_n^i e^{-nv} + \alpha_n^{i*} e^{nv} \right)$$

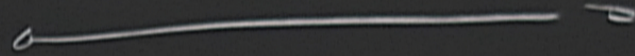
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$$|\alpha\rangle = e^{-|\alpha|^2/2} \sum_R \frac{1}{R!} (\hat{\alpha}^\dagger)^R |0\rangle$$

geometry / AdS

(smooth geometry)

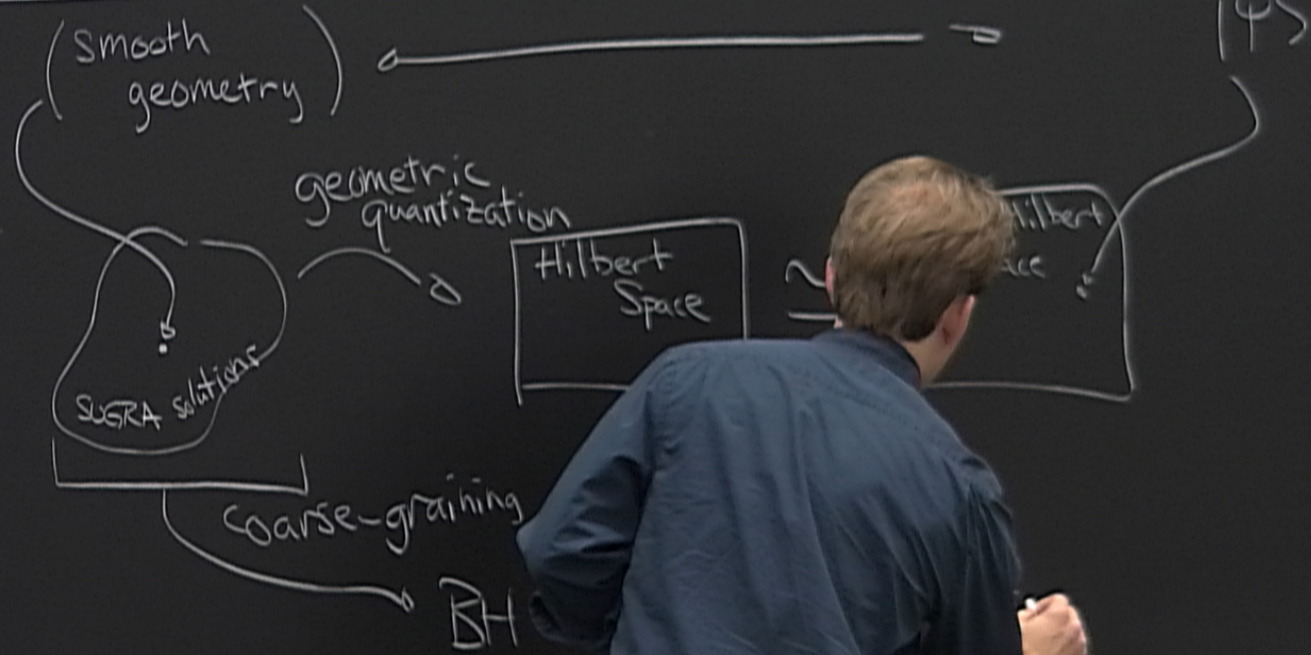


CFT  
|  
QFT



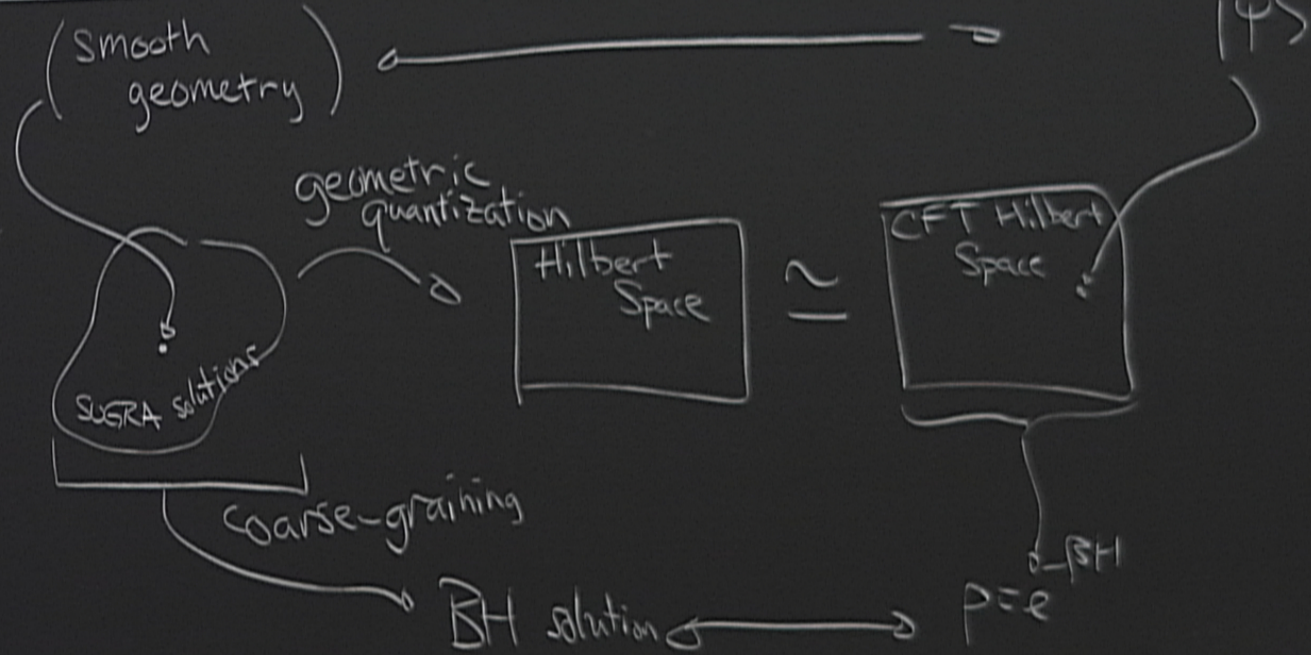


geometry / AdS



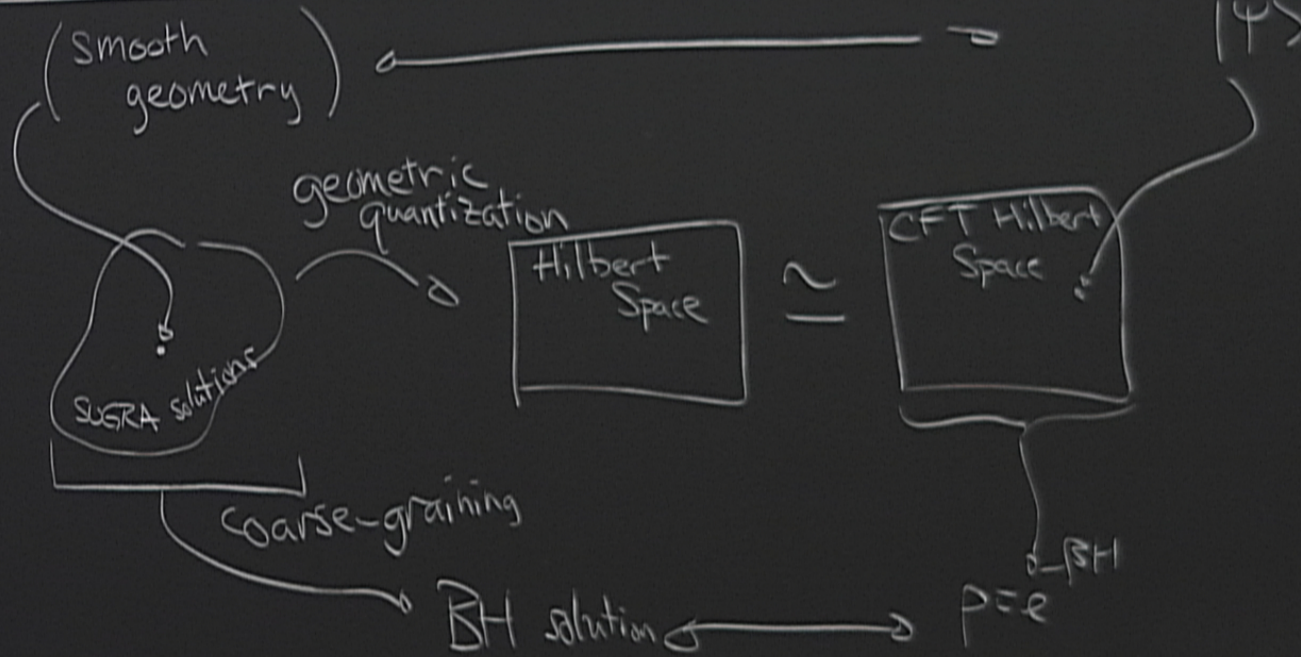


# Geometry / AdS





# Geometry / AdS



3-charge BPS

Bena-Warner

$$ds_{11}^2 = ds_5^2 \frac{(z_2 z_3)}{z_1}$$



3-charge BPS

Bena-Warner

$$ds_{11}^2 = ds_5^2 + \left( \frac{z_2 z_3}{z_1} \right)^{\frac{1}{3}} (dx_5^2 + dx_6^2) + \left( \frac{z_1 z_3}{z_2} \right)^{\frac{1}{3}} (dx_9^2 + dx_{10}^2)$$



3-charge BPS

Bena-Warner

$$ds_{11}^2 = ds_5^2 + \left(\frac{z_2 z_3}{z_1}\right)^{\frac{1}{3}} (dx_5^2 + dx_6^2) + \left(\frac{z_1 z_3}{z_2}\right)^{\frac{1}{3}} (dx_7^2 + dx_8^2) + \left(\frac{z_1 z_2}{z_3}\right)^{\frac{1}{3}} dx_{10}^2$$
$$\rightarrow - (z_1 z_2 z_3)^{\frac{2}{3}} (dt + k)^2 + (z_1 z_2 z_3)^{\frac{1}{3}} dy^2$$

$$C^{(3)} = A^{(1)} \wedge dx_5 \wedge dx_6 + \dots$$



$$\textcircled{H}^{(I)} = dA^{(I)} + d\left(\frac{dt+k}{z}\right)$$

$$\textcircled{H}^{(I)} = *_4 \textcircled{H}^{(I)}$$

$$z^2 z_I$$

$$\textcircled{1} \textcircled{I} = dA^{(1)} + d\left(\frac{dt+k}{z}\right)$$

$$(1) \textcircled{I}^{(I)} = *_{4} \textcircled{I}^{(H)}$$

$$(2) \nabla^2 z_I = \frac{1}{2} C_{IJK} *_{4} (\textcircled{I}^{(5)} \wedge \textcircled{I}^{(10)})$$

$$(3) dk + *_{4} dk = z_I \textcircled{I}^{(H)}$$



$$\textcircled{1} \textcircled{I} = dA^{(1)} + d\left(\frac{dt+k}{z}\right)$$

$$(1) \textcircled{I}^{(I)} = *_{4} \textcircled{I}^{(I)}$$

$$(2) \nabla^2 Z_I = \frac{1}{2} C_{ISK} *_{4} (\textcircled{I}^{(I)} \wedge \textcircled{I}^{(I)})$$

$$(3) dk + *_{4} dk = Z_I \textcircled{I}^{(I)}$$

open question: are there enough SUGRA solutions  
to explain 3-charge BH entropy?

$$\Omega \sim e^{\#N}$$

(10)



open question: are there enough SUGRA solutions  
to explain 3-charge BH entropy?

$$\Omega \sim e^{\#N}$$

(10)

$$\textcircled{1} \textcircled{I} = dA^{(1)} + d\left(\frac{dt+k}{z}\right)$$

open question: a  
to

$$(1) \textcircled{1} \textcircled{I} = *_{4} \textcircled{1} \textcircled{I}$$

$$(2) \nabla^2 Z_I = \frac{1}{2} C_{IJK} *_{4} (\textcircled{1} \textcircled{I} \textcircled{J} \textcircled{K})$$

$$(3) dk + *_{4} dk = Z_I \textcircled{1} \textcircled{I}$$



$$\textcircled{1} \textcircled{I} = dA^{(1)} + d(\dots)$$

$$(1) \textcircled{I} = *_{4}$$

$$(2) \nabla \cdot \dots = \text{ISK} * (\textcircled{5} \wedge \textcircled{70})$$

$$(3) \dots \textcircled{I}$$

open question: are there end  
to explain 3-

$$\Omega \sim e^{\#1}$$

JMRT

family of DLDSP nonextremal geometries

parameter  $(\mu, \bar{n}, \kappa) \in \mathbb{Z}^3$

$$M_{\mu, \bar{n}} \times S^1 \times T^4$$



JMRT

family of DIDSP nonextremal geometries

parameters  $(n, \bar{n}, k) \in \mathbb{Z}^3$

$M_{f,1} \times S^1 \times T^4$

AdS

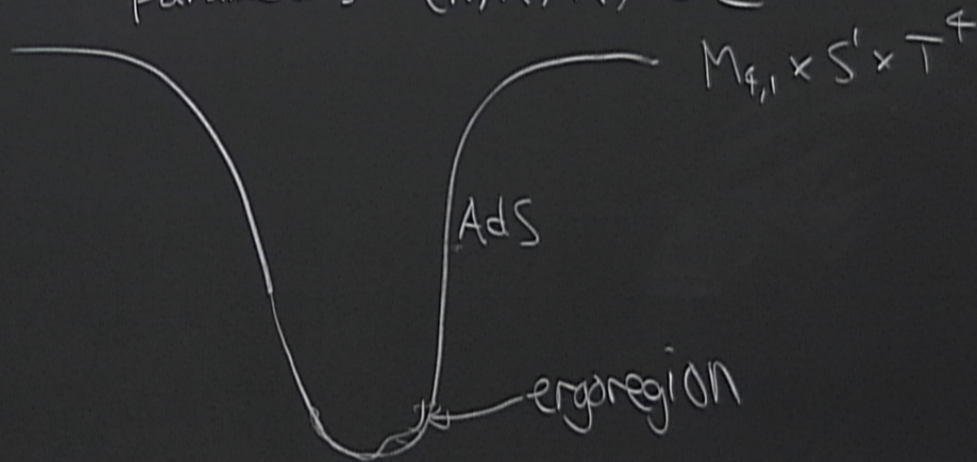
ergoregion



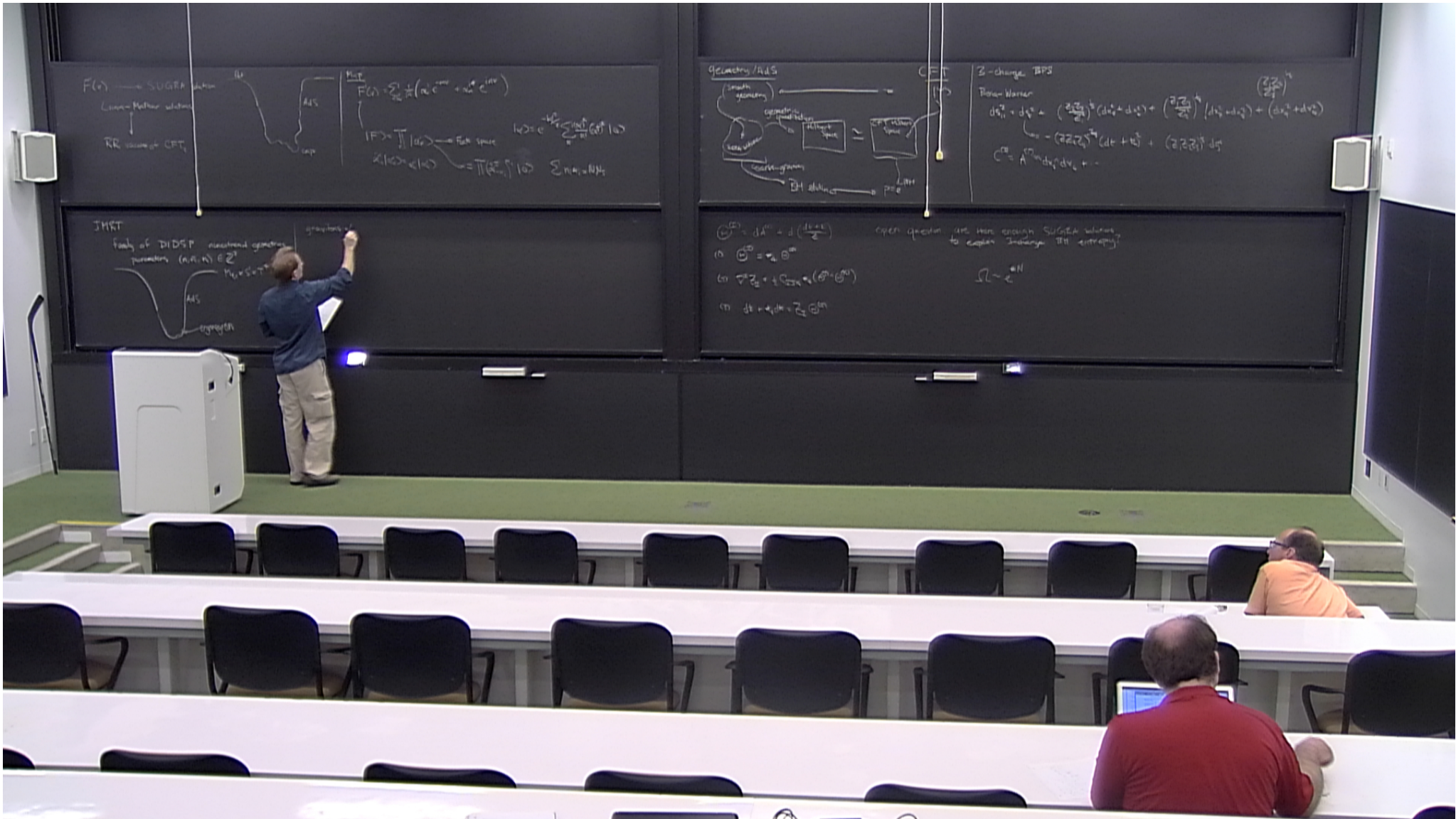
JMRT

family of D1D5P nonextremal geometries

parameters  $(n, \bar{n}, k) \in \mathbb{Z}^3$









gravitons w/ legs on  $T^4$

— minimally coupled scalars

metries

4  
live



gravitons w/ legs on  $T^4$   
— minimally coupled scalars

Solve  $\square\psi = 0$

$$e^{i(\lambda y - \omega t)} R(r) Y(\theta, \varphi, \delta)$$

gravitons w/ legs on  $T^4$   
 — minimally coupled scalars

Solve  $\square \Psi = 0$

$$\Psi = e^{i(\lambda y - \omega t)} R(r) Y(\theta, \varphi, \delta)$$

$$- \omega_{\pm} > 0$$

- classical instability

$$\lambda = \bar{a} - a$$

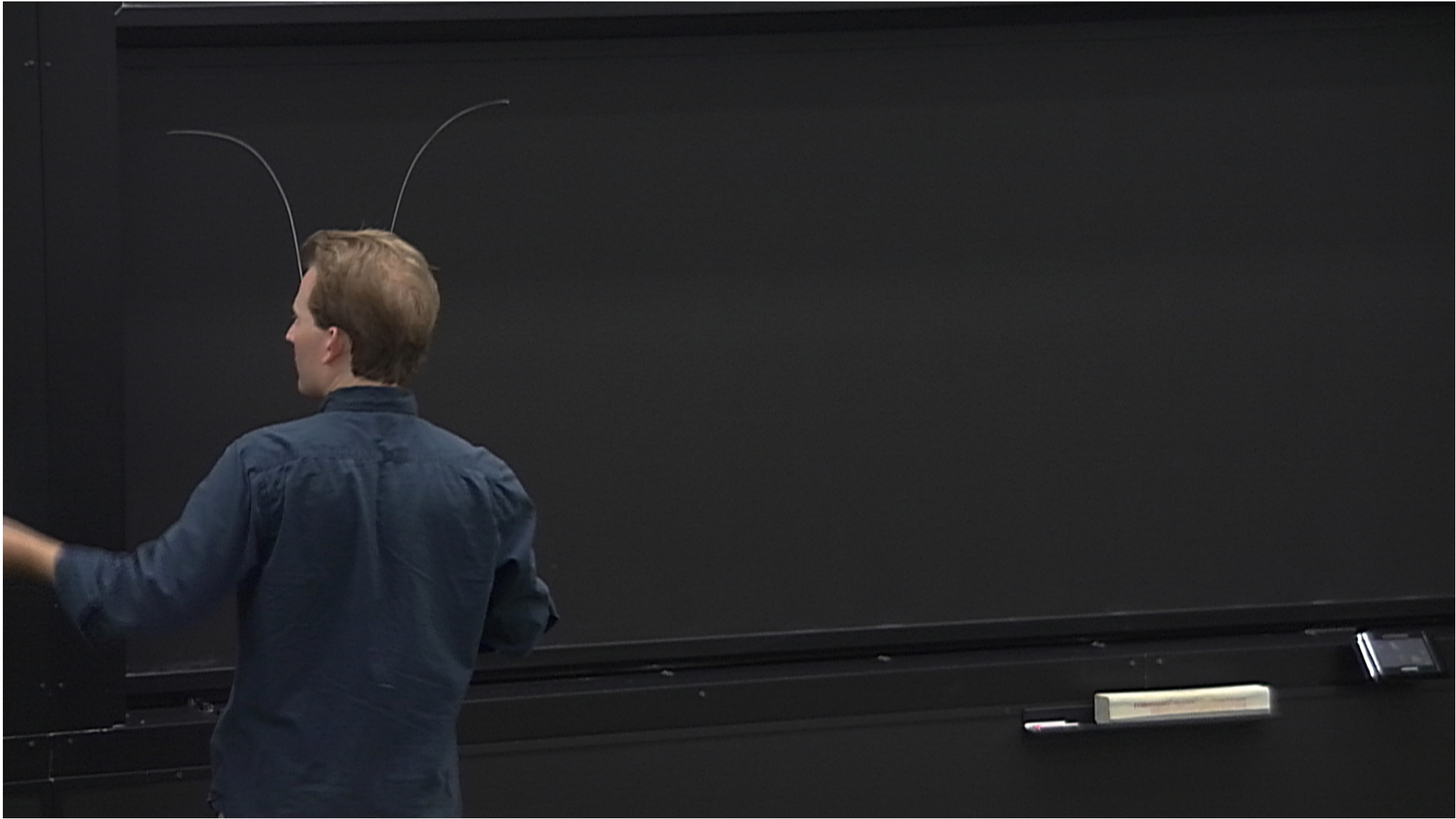
$$\omega_{\pm} = E = a + \bar{a}$$

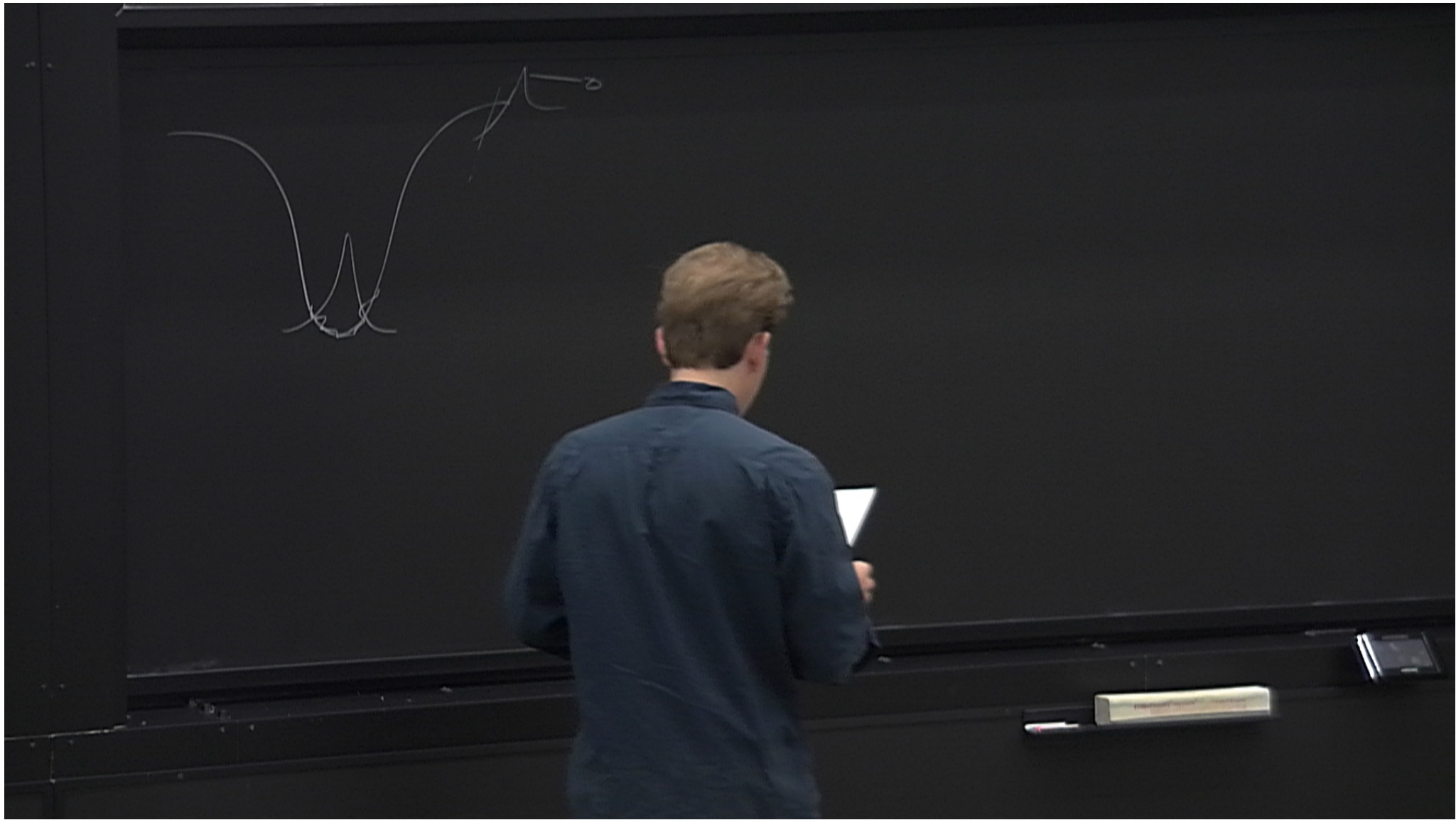
$$\omega_{\pm} = \frac{2\pi}{R^2} \left[ \frac{E^2 - \lambda^2}{4R^2} \frac{g^2 \alpha'^4}{V} n, n_s \right]^{l+1} \binom{N+l+1}{l+1} \binom{\bar{N}+l+1}{l+1}$$

$$a = -\frac{1}{R} \left[ \frac{l}{2} + (n + \frac{1}{2})(m_{\psi} - m_{\phi}) + N + 1 \right]$$

$$\bar{a} = -\frac{1}{R} \left[ \frac{l}{2} + (\bar{n} + \frac{1}{2})(m_{\phi} + m_{\psi}) + \bar{N} + 1 \right]$$

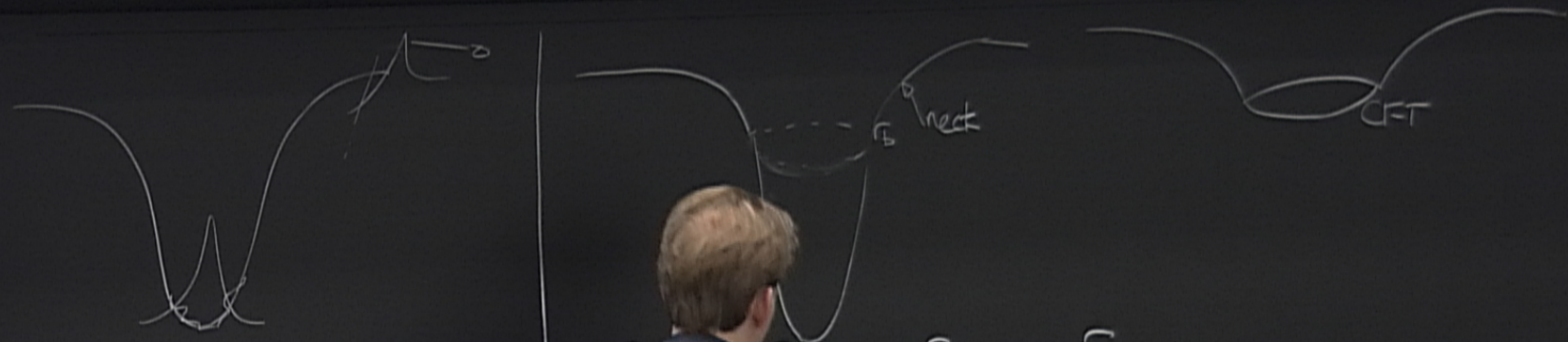




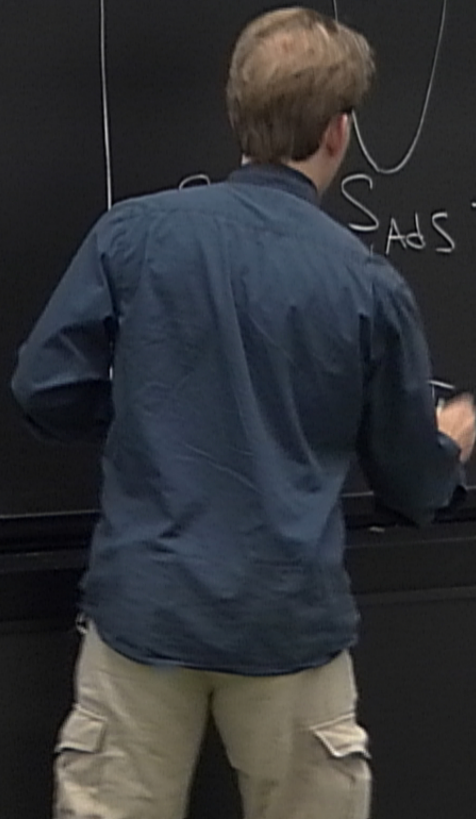




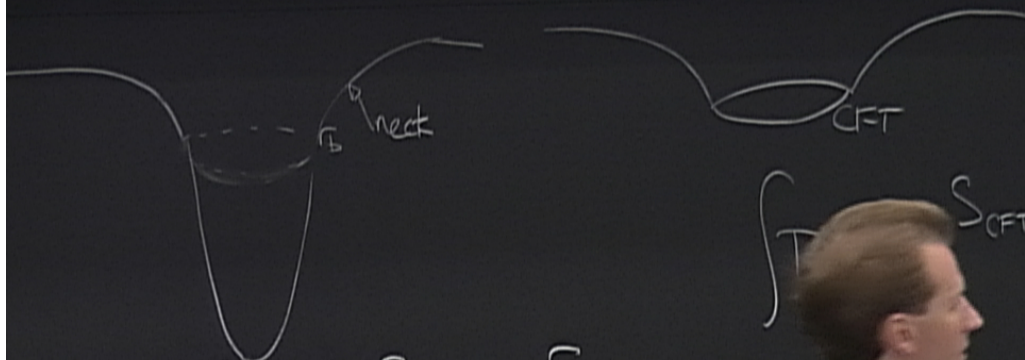
BH solution  $\longleftrightarrow$   $P=e$



$$S_{\text{AdS}} + S_{\text{neck}} + S_{\text{flat}}$$



→ p=0

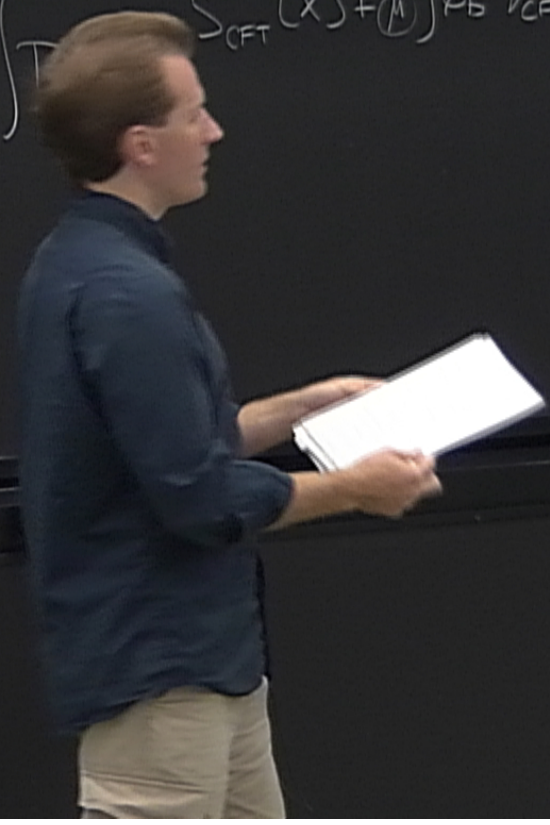


$$S_{tot} = S_{Ads} + S_{neck} + S_{flat}$$

↓

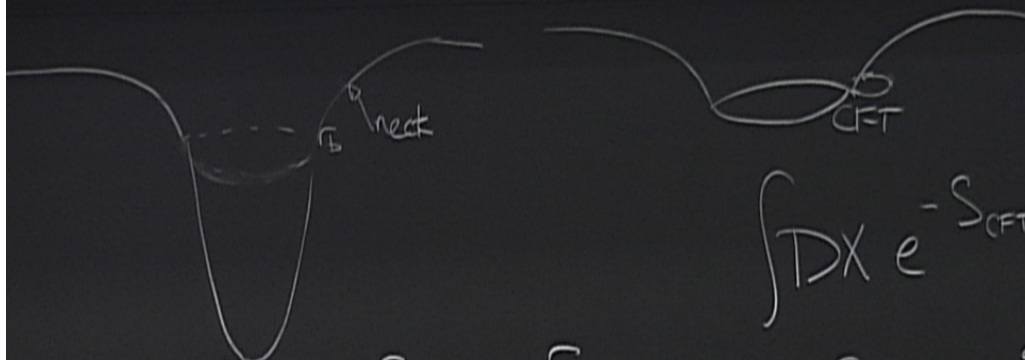
$$S_{CFT} + S_{int} + S_{flat}$$

$$\int \mathcal{D}\phi \int \mathcal{D}\psi \int \mathcal{D}\chi \exp(-S_{CFT}[\chi] + \mathcal{M} \int \phi_b V_{CFT}) = \int \mathcal{D}\phi \int \mathcal{D}\psi \exp(-S_{Ads}[\phi])$$





→ p.e



$$S_{\text{tot}} = S_{\text{AdS}} + S_{\text{neck}} + S_{\text{flat}}$$

$$\downarrow$$

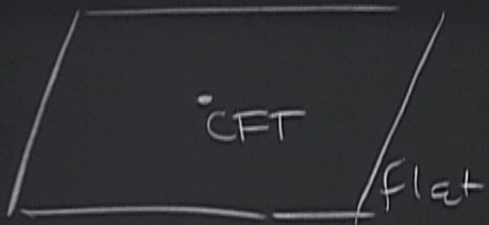
$$S_{\text{CFT}} + S_{\text{int}} + S_{\text{flat}}$$

$$\int DX e^{-S_{\text{CFT}}[X] + (M) \int \phi_b V_{\text{CFT}}} = \int D\phi|_{\phi_b} e^{-S_{\text{AdS}}[\phi]}$$

$$S_{\text{int}} = -c_2 \sum_m \int dy^d [\partial_\mu \phi]^2|_{r=0} V_{2,m}$$

$c_2 = \frac{16\pi^2}{b} \frac{1}{b}$

(flat space field)

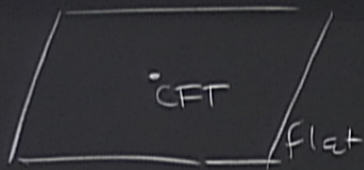


$$\mathcal{A} = \sum_{\text{CFT}} \langle f | \langle F | (S_{\text{int}} | I \rangle_{\text{Flat}} | i \rangle_{\text{CFT}})$$

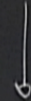


$$\frac{d\mathcal{T}}{dE} = \frac{2\pi}{2^{2\ell+1} \ell!^2} \frac{(Q_1 Q_5)^{\ell+1}}{R^{2\ell+3}} (E^2 - X^2)^{\ell+1} \underbrace{|\langle f | \nu | i \rangle|^2}_{\langle (L_{-1})^N \nu \nu \rangle} \delta_{\nu, \nu_0} \delta(E)$$

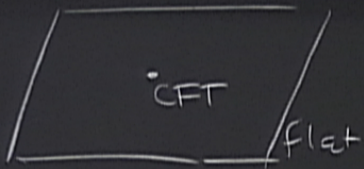




$$A = \sum_{\text{CFT}} f | \langle F | S_{\text{int}} | I \rangle_{\text{flat}} | i \rangle_{\text{CFT}}$$



$$\frac{d\Gamma}{dE} = \frac{2\pi}{2^{2L+1}} \frac{(Q_s)^{2L+1}}{R^{2D+2}} (E^2 - X^2)^{L+1} \underbrace{|\langle f | V | i \rangle|^2}_{\langle L_i^N \rangle} \delta_{\text{Dirac}} \delta(E - E_0)$$

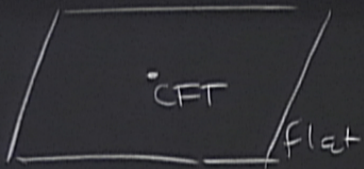


$$A = \sum_{\text{CFT}} f | \langle F | S_{int} | I \rangle_{\text{flat}} | i \rangle_{\text{CFT}}$$



$$\frac{d\Gamma}{dE} = \frac{2\pi}{2^{2L+1}} \frac{(Q, Q_S)^{2L+1}}{R^{2D+3}} (E^2 - X^2)^{2L+1} \underbrace{|\langle f | V | i \rangle|^2}_{\langle L, N \rangle} \delta_{\lambda, \lambda_0} \delta(E - E_0)$$





$$A = \sum_{\text{CFT}} f | \langle F | S_{\text{int}} | I \rangle_{\text{flat}} | i \rangle_{\text{CFT}}$$



$$\frac{d\Gamma}{dE} = \frac{2\pi}{2^{2l+1} l!} \frac{(Q_s \omega_s)^{2l+1}}{R^{2l+3}} (E^2 - X^2)^{2l+1} \underbrace{|\langle f | V | i \rangle|^2}_{\langle L_{-1}^N \rangle \langle L_1^N \rangle} \delta_{\text{Dirac}} \delta(E - E_0)$$



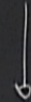
$$A = \sum_{\text{CFT}} \langle f | \langle F | \langle S_{\text{int}} | I \rangle_{\text{Flat}} | i \rangle_{\text{CFT}}$$

$$= \frac{2\pi}{R^{2L+1}} \frac{(Q, Q_S)^{L+1}}{L!} (E^2 - X^2)^{L+1} \underbrace{|\langle f | V | i \rangle|^2}_{\langle L, N \rangle} \delta_{\lambda, \lambda_0} \delta(E - E_0)$$





$$A = \sum_{\text{CFT}} f | \langle F | S_{\text{int}} | I \rangle_{\text{Flat}} | i \rangle_{\text{CFT}}$$



$$\frac{d\Gamma}{dE} = \frac{2\pi}{2^{2l+1} l!^2} \frac{(Q, Q_S)^{2l+1}}{R^{2l+3}} (E^2 - X^2)^{2l+1} \underbrace{|\langle f | V | i \rangle|^2}_{\langle L_{-1}^N \rangle} \delta_{\lambda, \lambda_0} \delta(E - E_0)$$

## Open Questions

- Is SUGRA enough?
- What use are microstates? " BH solution?
- What are the typical microstates like?
- Relationship to Euclidean BH?

4  
scalars

$$(-)Y(0, \psi, \beta)$$

$$\lambda = \bar{a} -$$
$$\omega_R = E$$
$$\omega_I =$$
$$a =$$
$$\bar{a} =$$



## Open Questions

- Is SUGRA enough?
- What use microstates? " BH solution?
- What typical microstates like?
- Rel Euclidean BH?
- Area

4  
scalars

$(-1)Y(0,4,6)$

$\lambda = \bar{a}$   
 $\omega_R = E$   
 $\omega_I =$   
 $a =$   
 $\bar{a} =$