

Title: Bipartite Graphs, Quivers and Gauge Theories

Date: May 14, 2013 02:00 PM

URL: <http://pirsa.org/13050020>

Abstract: We discuss how bipartite graphs on Riemann surfaces encapture a wealth of information about the physics of large classes of supersymmetric gauge theories, especially those with quiver structure and arising from the AdS/CFT context. The correspondence between the gauge theory, the underlying algebraic geometry of its space of vacua, the combinatorics of dimers and toric varieties, as well as the number theory of dessin d'enfants becomes particularly intricate under this light.

# Bipartite Graphs, Quivers and Gauge Theories

YANG-HUI HE

Dept of Mathematics, City University, London;  
School of Physics, NanKai University;  
Merton College, University of Oxford

Perimeter Institute, May, 2013



# Bipartite Graphs, Quivers and Gauge Theories

YANG-HUI HE

Dept of Mathematics, City University, London;  
School of Physics, NanKai University;  
Merton College, University of Oxford

Perimeter Institute, May, 2013



# Acknowledgements

- 0511287, B. Feng, YHH, K. Kennaway, C. Vafa
- 0608050, 0701063: S. Benvenuti, B. Feng, A. Hanany, YHH J. Gray, YHH, V. Jejjala, B. Nelson
- 0909.2879: J. Hewlett, YHH
- 1204.1065, 1107.4101, 1104.5490 Hanany, YHH, Jejjala, Pasukonis, Ramgoolam, Rodriguez-Gomez
- 1201.3633 YHH, J. McKay and 1211.1931 YHH, J. McKay J. Read

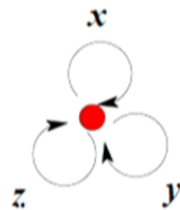
## Two most famous SCFTs in 4-D

- $\mathcal{N} = 4$   $U(N)$  Yang-Mills

- 3 adjoint fields  $X, Y, Z$  with superpotential

$$W = \text{Tr}(X[Y, Z]) = \text{Tr}(XYZ - XZY)$$

- Original AdS/CFT:  $N$  D3-branes transverse to flat  $\mathbb{R}^6$



NODE = Gauge Group

ARROW = Bi-fundamental (Adj) Field

- QUIVER = Finite graph (label =  $\text{rk}(\text{gauge factor})$ ) + relations from SUSY

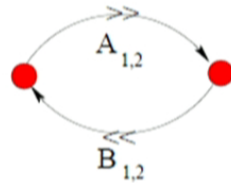
- Matter Content: Nodes + arrows

- Relations (F-Terms):  $D_i W = 0 \rightsquigarrow [X, Y] = [Y, Z] = [X, Z] = 0$

- # gauge factors =  $N_g = 1$ ; # fields =  $N_f = 3$ ; # terms in  $W = N_w = 2$



- Klebanov-Witten  $\mathcal{N} = 1$  “conifold” Theory
  - $SU(N) \times SU(N)$  gauge theory with 4 bi-fundamental fields



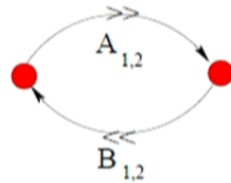
	$SU(N)$	$SU(N)$
$A_{i=1,2}$	$\square$	$\bar{\square}$
$B_{j=1,2}$	$\bar{\square}$	$\square$

$W = \text{Tr}(\epsilon_{il} \epsilon_{jk} A_i B_j A_l B_k)$

QUIVER

- string-theory realization:  $N$  D3-branes transverse to the conifold singularity:
 
$$\{uv = wz\} \subset \mathbb{C}^4$$
- # gauge factors =  $N_g = 2$ ; # fields =  $N_f = 4$ ; # terms in  $W = N_w = 2$
- Observatio Curiosa:  $N_g - N_f + N_w = 0$
- true for almost all known cases in  $AdS_5/CFT_4$

- Klebanov-Witten  $\mathcal{N} = 1$  “conifold” Theory
  - $SU(N) \times SU(N)$  gauge theory with 4 bi-fundamental fields



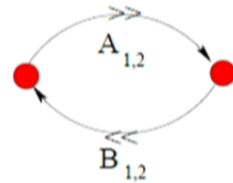
	$SU(N)$	$SU(N)$
$A_{i=1,2}$	$\square$	$\bar{\square}$
$B_{j=1,2}$	$\bar{\square}$	$\square$

$W = \text{Tr}(\epsilon_{il} \epsilon_{jk} A_i B_j A_l B_k)$

QUIVER

- string-theory realization:  $N$  D3-branes transverse to the conifold singularity:
 
$$\{uv = wz\} \subset \mathbb{C}^4$$
- # gauge factors =  $N_g = 2$ ; # fields =  $N_f = 4$ ; # terms in  $W = N_w = 2$
- Observatio Curiosa:  $N_g - N_f + N_w = 0$
- true for almost all known cases in  $AdS_5/CFT_4$

- Klebanov-Witten  $\mathcal{N} = 1$  “conifold” Theory
  - $SU(N) \times SU(N)$  gauge theory with 4 bi-fundamental fields



	$SU(N)$	$SU(N)$
$A_{i=1,2}$	$\square$	$\bar{\square}$
$B_{j=1,2}$	$\bar{\square}$	$\square$

$$W = \text{Tr}(\epsilon_{il} \epsilon_{jk} A_i B_j A_l B_k)$$

QUIVER

- string-theory realization:  $N$  D3-branes transverse to the conifold singularity:
 
$$\{uv = wz\} \subset \mathbb{C}^4$$
- # gauge factors =  $N_g = 2$ ; # fields =  $N_f = 4$ ; # terms in  $W = N_w = 2$
- Observatio Curiosa:  $N_g - N_f + N_w = 0$
- true for almost all known cases in  $AdS_5/CFT_4$



# Gauge Theory Moduli Space & CY Geometry

$$S = \int d^4x \left[ \int d^4\theta \Phi_i^\dagger e^V \Phi_i + \left( \frac{1}{4g^2} \int d^2\theta \text{Tr} \mathcal{W}_\alpha \mathcal{W}^\alpha + \int d^2\theta W(\Phi) + \text{h.c.} \right) \right]$$

$$W = \text{superpotential} \quad V(\phi_i, \bar{\phi}_i) = \sum_i \left| \frac{\partial W}{\partial \phi_i} \right|^2 + \frac{g^2}{4} (\sum_i q_i |\phi_i|^2)^2$$

- VACUUM  $\sim \boxed{V(\phi_i, \bar{\phi}_i) = 0} \Rightarrow \begin{cases} \frac{\partial W}{\partial \phi_i} = 0 & \text{F-TERMS} \\ \sum_i q_i |\phi_i|^2 = 0 & \text{D-TERMS} \end{cases}$
- $\mathcal{M} :=$  vacuum moduli space = space of solutions to F and D-flatness  
= affine (singular) algebraic variety  $\rightsquigarrow$  Quiver Variety
- If  $\mathcal{M}$  CY3, can realize in string theory as D3-brane probing  $\mathcal{M}$ 
  - $\mathcal{M} \rightarrow$  Gauge Theory: Geometrical Engineering (Vafa, Cachazo, et al)
  - Gauge Theory  $\rightarrow \mathcal{M}$ : Forward Algorithm
- for  $N$ -branes, get  $Sym^N \mathcal{M} = \mathcal{M}^N / \Sigma_N$

# Gauge Theory Moduli Space & CY Geometry

$$S = \int d^4x \left[ \int d^4\theta \Phi_i^\dagger e^V \Phi_i + \left( \frac{1}{4g^2} \int d^2\theta \text{Tr} \mathcal{W}_\alpha \mathcal{W}^\alpha + \int d^2\theta W(\Phi) + \text{h.c.} \right) \right]$$

$$W = \text{superpotential} \quad V(\phi_i, \bar{\phi}_i) = \sum_i \left| \frac{\partial W}{\partial \phi_i} \right|^2 + \frac{g^2}{4} (\sum_i q_i |\phi_i|^2)^2$$

- VACUUM  $\sim \boxed{V(\phi_i, \bar{\phi}_i) = 0} \Rightarrow \begin{cases} \frac{\partial W}{\partial \phi_i} = 0 & \text{F-TERMS} \\ \sum_i q_i |\phi_i|^2 = 0 & \text{D-TERMS} \end{cases}$
- $\mathcal{M} :=$  vacuum moduli space = space of solutions to F and D-flatness  
= affine (singular) algebraic variety  $\rightsquigarrow$  Quiver Variety
- If  $\mathcal{M}$  CY3, can realize in string theory as D3-brane probing  $\mathcal{M}$ 
  - $\mathcal{M} \rightarrow$  Gauge Theory: Geometrical Engineering (Vafa, Cachazo, et al)
  - Gauge Theory  $\rightarrow \mathcal{M}$ : Forward Algorithm
- for  $N$ -branes, get  $Sym^N \mathcal{M} = \mathcal{M}^N / \Sigma_N$

# An Computational Approach

- Programme to study the computational algebraic geometry of  $\mathcal{M}$ : w/ Gray, Jejjala, Nelson; Hauenstein, Mehta; Hanany, Mekareeya; Chuang, Hewlett;

①  $n$ -fields: start with polynomial ring  $\mathbb{C}[\phi_1, \dots, \phi_n]$

②  $D =$  set of  $k$  GIO's: a **ring map**  $\mathbb{C}[\phi_1, \dots, \phi_n] \xrightarrow{D} \mathbb{C}[D_1, \dots, D_k]$

③ Now incorporate superpotential: F-flatness

$$\langle f_{i=1, \dots, n} = \frac{\partial W(\phi_i)}{\partial \phi_i} = 0 \rangle \simeq \text{ideal of } \mathbb{C}[\phi_1, \dots, \phi_k]$$

④ Moduli space = image of the ring map

$$\frac{\mathbb{C}[\phi_1, \dots, \phi_n]}{\{F = \langle f_1, \dots, f_n \rangle\}} \xrightarrow{D = \text{GIO}} \mathbb{C}[D_1, \dots, D_k], \quad \mathcal{M} \simeq \text{Im}(D)$$

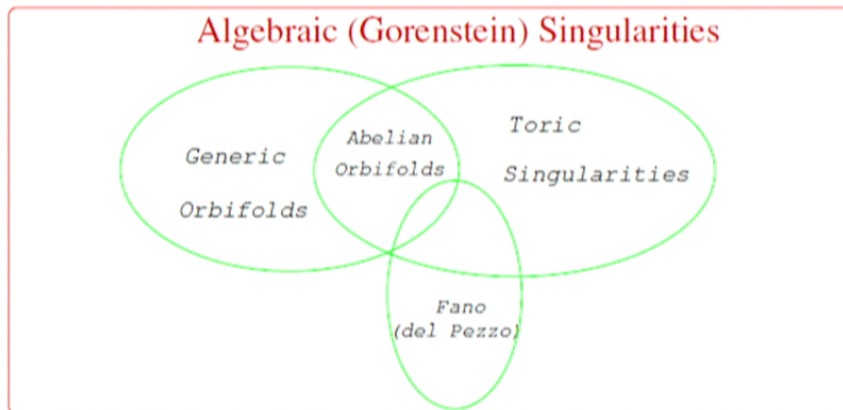
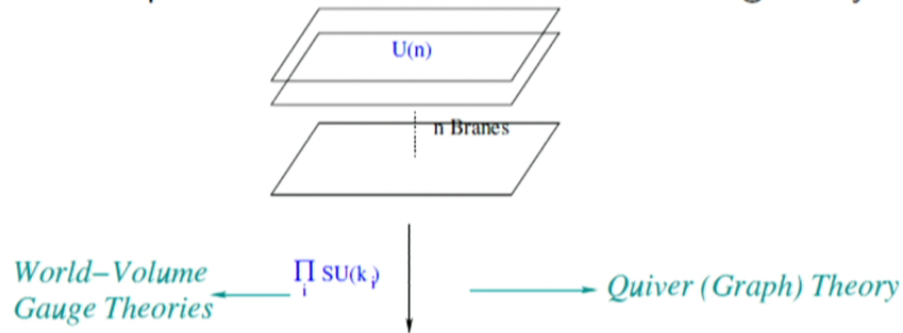
- Image is an ideal of  $\mathbb{C}[D_1, \dots, D_k]$ , i.e.,

$\mathcal{M}$  explicitly realised as an affine variety in  $\mathbb{C}^k$

- Gröbner Bases: implemented in MACAULAY2, Singular...
- Numerical AG: implemented in Bertini...

# Gauge Theory & CY correspondence: Paradigm

Stack of  $n$  parallel branes  $\perp$  a Calabi-Yau singularity  $\mathcal{M}$ :



Orbifolds:

$$\mathbb{C}^3 / (\Gamma \subset SU(3))$$

Toric:

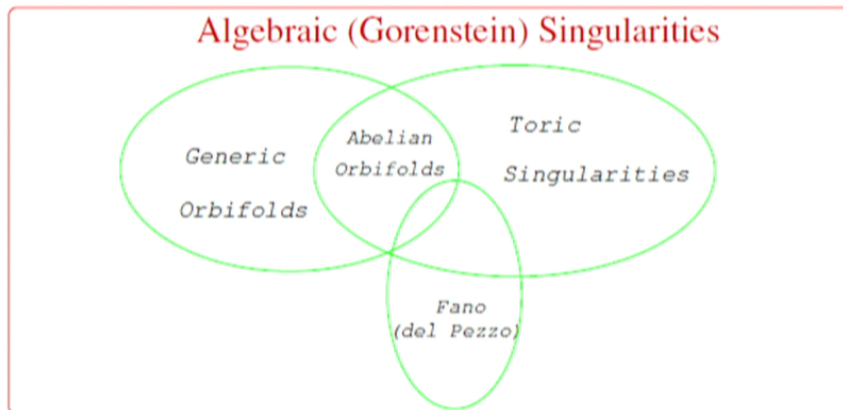
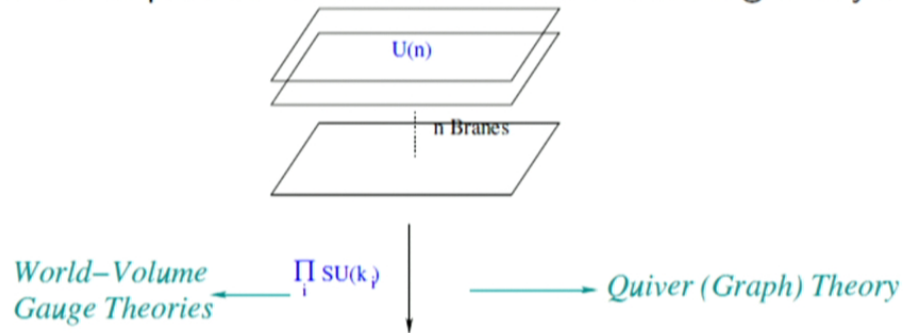
e.g., conifold,  $Y^{p,q}$ ,  
 $L^{p,q} \dots$

Fano (del Pezzo):

e.g.,  $dP_{0,\dots,8}$

# Gauge Theory & CY correspondence: Paradigm

Stack of  $n$  parallel branes  $\perp$  a Calabi-Yau singularity  $\mathcal{M}$ :



Orbifolds:

$$\mathbb{C}^3 / (\Gamma \subset SU(3))$$

Toric:

e.g., conifold,  $Y^{p,q}$ ,  
 $L^{p,q} \dots$

Fano (del Pezzo):

e.g.,  $dP_{0,\dots,8}$

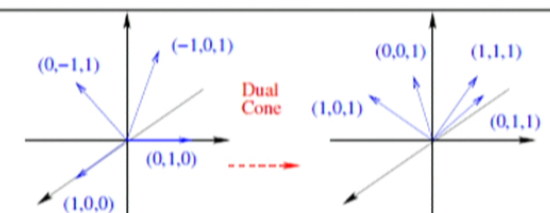
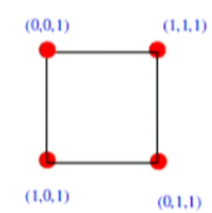
# Toric $AdS_5/CFT_4$ , or Gauge Theories with $\mathcal{M}$ Toric CY3

- By far the largest class known and studied
- World-volume physics and Geometry of  $\mathcal{M} \sim$  Combinatorial Data  $\sim$  (integer cones  $\sigma$  in  $\mathbb{Z}^r$ -lattice); (CY3  $\rightsquigarrow$  planar toric diag)
  - Computing  $\mathcal{M}$ : (Aspinwall, Diaconescu, Douglas, Greene, Morrison, Plesser, et al.) from Witten's GLSM;
  - Feng-Hanany-YHH (Inverse Algorithm toric diag  $\rightsquigarrow$  gauge theory) 0003085
- Explicit Ricci-flat metric known for infinite families  $Y^{p,q}, L^{abc}$  (conifold, special case); [Candelas-de la Ossa, Cvetič, Hanany, Pope, Sparks, Waldram ...]
- Dual Sasaki-Einstein Cone has  $U(1)^3$  isometry
- No known compact CY3 metrics, all known affine CY3 metrics are toric so far.

# Toric $AdS_5/CFT_4$ , or Gauge Theories with $\mathcal{M}$ Toric CY3

- By far the largest class known and studied
- World-volume physics and Geometry of  $\mathcal{M} \sim$  Combinatorial Data  $\sim$  (integer cones  $\sigma$  in  $\mathbb{Z}^r$ -lattice); (CY3  $\rightsquigarrow$  planar toric diag)
  - Computing  $\mathcal{M}$ : (Aspinwall, Diaconescu, Douglas, Greene, Morrison, Plesser, et al.) from Witten's GLSM;
  - Feng-Hanany-YHH (Inverse Algorithm toric diag  $\rightsquigarrow$  gauge theory) 0003085
- Explicit Ricci-flat metric known for infinite families  $Y^{p,q}, L^{abc}$  (conifold, special case); [Candelas-de la Ossa, Cvetič, Hanany, Pope, Sparks, Waldram ...]
- Dual Sasaki-Einstein Cone has  $U(1)^3$  isometry
- No known compact CY3 metrics, all known affine CY3 metrics are toric so far.

## Notation for Affine Toric Variety

Def		Example (Conifold)
Comb.:	<p>Convex Cone <math>\sigma \in \mathbb{Z}^d \rightsquigarrow</math></p> <p>Dual Cone <math>\sigma^\vee \rightsquigarrow X =</math></p> <p><math>\text{Spec}_{Max} \mathbb{C}[S_\sigma = x_i^{\text{gen}(\sigma^\vee) \cap \mathbb{Z}^d}]</math></p> <p>Toric Diagram = <math>S_\sigma</math></p>	 <p><math>S_\sigma = \langle a = z, c = yz, b = xyz, d = xz \rangle</math></p> <p><math>ab = cd</math> in <math>\mathbb{C}^4[a, b, c, d]</math></p>
Symp:	<p>Generalise <math>\mathbb{P}^n</math>:</p> <p>a <math>(\mathbb{C}^*)^{q-d}</math> action on <math>\mathbb{C}_{[x_i]}^q</math></p> <p><math>x_i \mapsto \lambda_a^{Q_i^a=1 \dots q-d} x_i</math> with</p> <p>Relations: <math>\sum_{i=1}^d Q_i^a v_i = 0</math></p> <p>Toric Diagram = <math>v_i</math></p>	 <p><math>Q = [-1, -1, 1, 1]</math></p> <p><math>\mathbb{C}^*</math> on <math>\mathbb{C}^4 \rightsquigarrow</math></p> <p><math>\ker Q = G_t =</math></p> $\begin{pmatrix} 1 & 0 & 0 & 1 \\ 1 & 0 & 1 & 0 \\ 1 & 1 & 1 & 1 \end{pmatrix}$
Comp:	Binomial Ideal $\langle \prod p_i = \prod q_j \rangle$	$ab = cd$ in $\mathbb{C}^4$



# toric gauge theory $\simeq$ dimer/tiling configurations

- $G - E + N_T = 0$  is Euler relation for a torus!  $G$  nodes (vol(toric diag)),  $E$  edges (fields),  $N_T$  faces (loops in the quiver,  $W$ -terms)
- Quiver: loops / F-term relations  $\sim \mathcal{M} = \text{CY3}$  (planar toric diag  $D$ )
  - dual of  $D = (p, q)$ -configuration of branes
  - repetition in  $D =$  multiplicities in homogeneous coord
- Dual Periodic Quiver: (coloured graph dual) black (white) nodes  $\sim W$  terms with  $+$  ( $-$ ) sign; (Hanany, Kennaway, Feng, Franco, YHH, Uranga, Vegh, Wecht)
  - bi-partite: each field appears twice with opposite sign; periodic: torus
  - Dimer Model on a torus (i.e., periodic planar tiling)
  - Kasteleyn matrix = weighted adjacency matrix of a dimer model
  - $P(z, w) = \text{Det}(Kas) = 2\text{-var poly} \rightsquigarrow$  Perfect matchings = coefficients of  $\text{Det}(Kas) =$  multiplicities in GLSM fields = repetitions of the nodes in  $D$



## toric gauge theory $\simeq$ dimer/tiling configurations

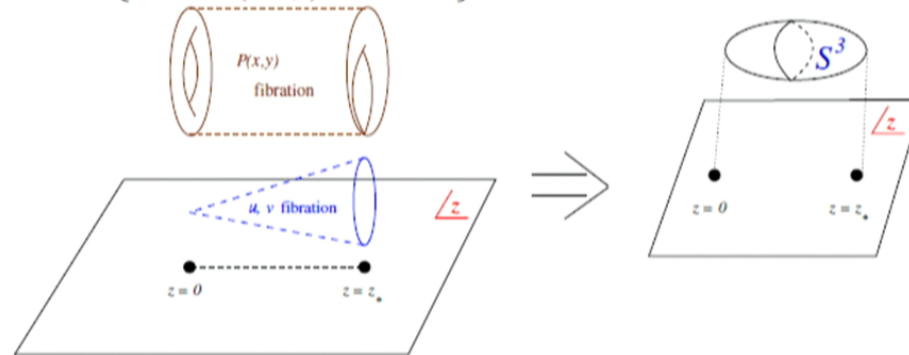
- $G - E + N_T = 0$  is Euler relation for a torus!  $G$  nodes (vol(toric diag)),  $E$  edges (fields),  $N_T$  faces (loops in the quiver,  $W$ -terms)
- Quiver: loops / F-term relations  $\sim \mathcal{M} = \text{CY3}$  (planar toric diag  $D$ )
  - dual of  $D = (p, q)$ -configuration of branes
  - repetition in  $D =$  multiplicities in homogeneous coord
- Dual Periodic Quiver: (coloured graph dual) black (white) nodes  $\sim W$  terms with  $+$  ( $-$ ) sign; (Hanany, Kennaway, Feng, Franco, YHH, Uranga, Vegh, Wecht)
  - bi-partite: each field appears twice with opposite sign; periodic: torus
  - Dimer Model on a torus (i.e., periodic planar tiling)
  - Kasteleyn matrix = weighted adjacency matrix of a dimer model
  - $P(z, w) = \text{Det}(Kas) = 2\text{-var poly} \rightsquigarrow$  Perfect matchings = coefficients of  $\text{Det}(Kas) =$  multiplicities in GLSM fields = repetitions of the nodes in  $D$



# Toric CY3, Mirror Symmetry & Dimers/Tilings

- [Feng-Kennaway-YHH-Vafa] torus  $T^2$  lives in  $T^3$  of mirror symmetry;  
[Heckman-Vafa-Xie-Yamazaki] stringy origins of BFT [Franco]
- Mirror Symmetry:
  - Thrice T-duality: D3-brane on CY3  $\rightsquigarrow$  D6-branes wrapping 3-cycles in mirror

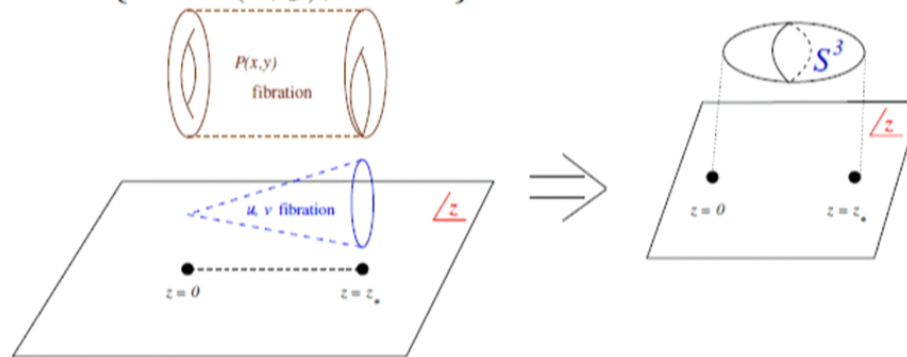
CY3:  $\{z = P(x, y), z = uv\} \subset \mathbb{C}^5$



- Intersection in  $H_3$  of mirror gives quiver:  $a_{ij} = S_j \cap S_j$
- zig-zag paths in  $T^2 \sim (p, q)$ -cycle  $\sim$  winding paths in  $\Sigma$  (show by Pick's)

# Toric CY3, Mirror Symmetry & Dimers/Tilings

- [Feng-Kennaway-YHH-Vafa] torus  $T^2$  lives in  $T^3$  of mirror symmetry;  
[Heckman-Vafa-Xie-Yamazaki] stringy origins of BFT [Franco]
- Mirror Symmetry:
  - Thrice T-duality: D3-brane on CY3  $\rightsquigarrow$  D6-branes wrapping 3-cycles in mirror  
CY3:  $\{z = P(x, y), z = uv\} \subset \mathbb{C}^5$



- Intersection in  $H_3$  of mirror gives quiver:  $a_{ij} = S_j \cap S_j$
- zig-zag paths in  $T^2 \sim (p, q)$ -cycle  $\sim$  winding paths in  $\Sigma$  (show by Pick's)

# Amoebae & Algae

- **Amoeba Projection**  $Log(z, w) \rightarrow (\log |z|, \log |w|)$

$$A = Amoeba(P(z, w) \subset (\mathbb{C}^*)^2) = Log(P) \subset \mathbb{R}^2 \rightsquigarrow$$

skeleton of  $A$  is the  $(p, q)$ -configuration

- $T^2$  of dimer model lives in the  $T^3$  of mirror symmetry
  - $P(z, w) = 0$  describes fiber  $\Sigma$  over  $s = 0$  in mirror CY3
  - $(\cap 3\text{-cycles}) \cap \Sigma$  at a graph  $\Gamma$  on  $T^2 \subset T^3 \rightsquigarrow$  **periodic tiling**
  - **Alga Projection**:  $Arg(z, w) \rightarrow (\arg(z), \arg(w))$

$$Alga(P(z, w) \subset (\mathbb{C}^*)^2) = Arg(P) \subset [0, 2\pi)^2 \rightsquigarrow$$

fundamental region of dimer

# Amoebae & Algae

- **Amoeba Projection**  $Log(z, w) \rightarrow (\log |z|, \log |w|)$

$$A = Amoeba(P(z, w) \subset (\mathbb{C}^*)^2) = Log(P) \subset \mathbb{R}^2 \rightsquigarrow$$

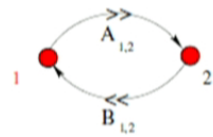
skeleton of  $A$  is the  $(p, q)$ -configuration

- $T^2$  of dimer model lives in the  $T^3$  of mirror symmetry
  - $P(z, w) = 0$  describes fiber  $\Sigma$  over  $s = 0$  in mirror CY3
  - $(\cap 3\text{-cycles}) \cap \Sigma$  at a graph  $\Gamma$  on  $T^2 \subset T^3 \rightsquigarrow$  **periodic tiling**
  - **Alga Projection**:  $Arg(z, w) \rightarrow (\arg(z), \arg(w))$

$$Alga(P(z, w) \subset (\mathbb{C}^*)^2) = Arg(P) \subset [0, 2\pi)^2 \rightsquigarrow$$

fundamental region of dimer

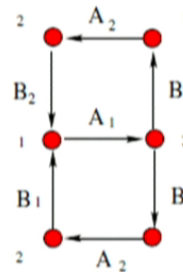
# Quivers, CY3, Tilings and Dimers



Quiver

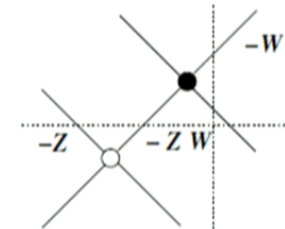
$$W = \text{Tr}(A_1 B_1 A_2 B_2 - A_1 B_2 A_2 B_1)$$

Draw on torus



Periodic Quiver  
W encoded

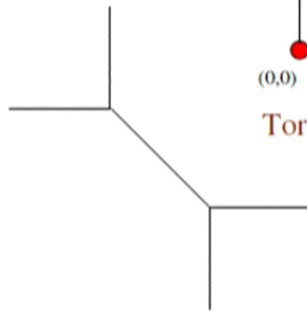
Graph Dual



Dimer Model

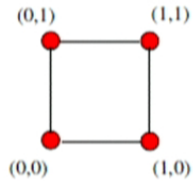
Period Tiling  
of Plane

Geometric  
Engineering



(p,q)-Web

YANG-HUI HE (London/Tianjin/Oxford)



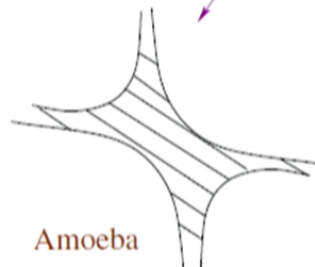
Toric Diagram

Graph Dual

Newton  
Polynomial

Kastelyn Matrix  
 $\det K(Z,W)_{1 \times 1} = 1 - Z - W - WZ$

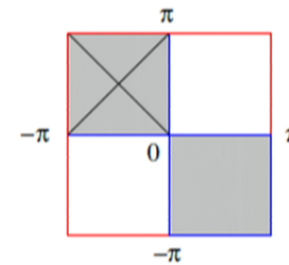
Projection



Amoeba

Bipartita

Projection



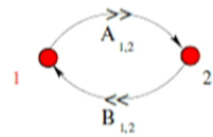
Alga



PI, May, 2013

13 / 29

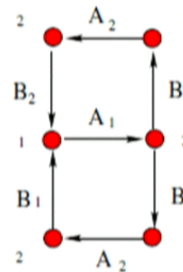
# Quivers, CY3, Tilings and Dimers



Quiver

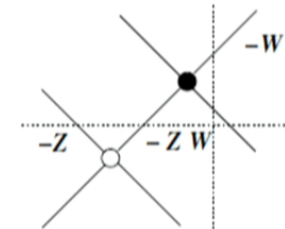
$$W = \text{Tr}(A_1 B_1 A_2 B_2 - A_1 B_2 A_2 B_1)$$

Draw on torus



Periodic Quiver  
W encoded

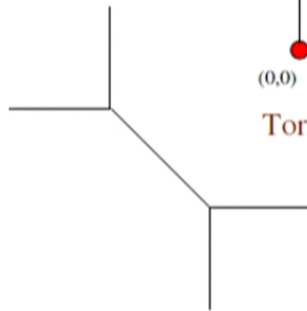
Graph Dual



Dimer Model

Period Tiling  
of Plane

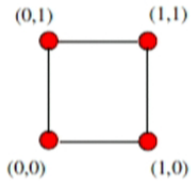
Geometric  
Engineering



(p,q)-Web

YANG-HUI HE (London/Tianjin/Oxford)

Newton  
Polynomial

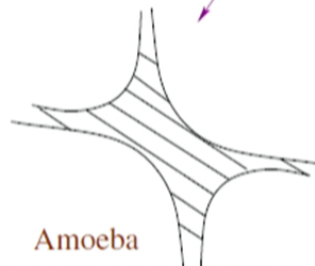


Toric Diagram

Kastelyn Matrix

$$\det K(Z,W)_{1 \times 1} = 1 - Z - W - WZ$$

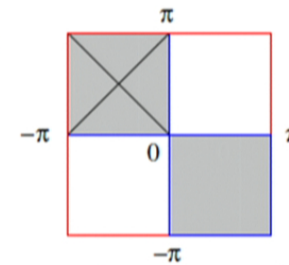
Projection



Amoeba

Bipartita

Projection



Alga



PI, May, 2013

13 / 29



# Seiberg Duality for $\mathcal{N} = 1$

- 2 theories: **Direct Electric** theory:  $N_c$  with  $N_f$  flavours; **Dual Magnetic** theory:  $N_f - N_c$  (take  $\frac{3}{2}N_c \leq N_f \leq 3N_c$ ) with  $N_f$  flavours

	$SU(N_c)$	$SU(N_f)_L$	$SU(N_f)_R$		$SU(N_f - N_c)$	$SU(N_f)_L$	$SU(N_f)_R$
$Q$	$\square$	$\square$	$1$	$q$	$\square$	$\bar{\square}$	$1$
$Q'$	$\bar{\square}$	$1$	$\bar{\square}$	$q'$	$\bar{\square}$	$1$	$\square$
				$M$	$1$	$\square$	$\bar{\square}$

$W = 0$ 
 $W = Mqq'$

- two theories with different gauge symmetries, same global symmetries and different matter content/superpotential
- have the same IR dynamics, i.e., **same vacuum moduli space**

In matching, think of  $M$  as a composite meson made of the quarks in the direct theory.

# Seiberg Duality for $\mathcal{N} = 1$

- 2 theories: **Direct Electric** theory:  $N_c$  with  $N_f$  flavours; **Dual Magnetic** theory:  $N_f - N_c$  (take  $\frac{3}{2}N_c \leq N_f \leq 3N_c$ ) with  $N_f$  flavours

	$SU(N_c)$	$SU(N_f)_L$	$SU(N_f)_R$		$SU(N_f - N_c)$	$SU(N_f)_L$	$SU(N_f)_R$
$Q$	$\square$	$\square$	$1$	$q$	$\square$	$\bar{\square}$	$1$
$Q'$	$\bar{\square}$	$1$	$\bar{\square}$	$q'$	$\bar{\square}$	$1$	$\square$
				$M$	$1$	$\square$	$\bar{\square}$

$W = 0$ 
 $W = Mqq'$

- two theories with different gauge symmetries, same global symmetries and different matter content/superpotential
- have the same IR dynamics, i.e., **same vacuum moduli space**

In matching, think of  $M$  as a composite meson made of the quarks in the direct theory.

# A Quiver Duality from Seiberg Duality

We have quiver labeled by  $(N_c)_i$  and arrows  $a_{ij}$ :

- 1 Pick dualisation node  $i_0$  with  $N_c$ , and define:

$I_{in} :=$  nodes having arrows going into  $i_0$

$I_{out} :=$  nodes having arrow coming from  $i_0$

$I_{no} :=$  nodes unconnected with  $i_0$

- 2  $N_c \rightarrow N_f - N_c$  (where  $N_f = \sum_{i \in I_{in}} a_{i,i_0} = \sum_{i \in I_{out}} a_{i_0,i}$ )

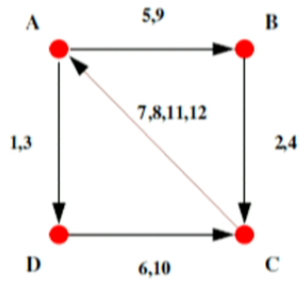
- 3 Reverse arrows going in or out of  $i_0$ , leave  $I_{no}$ , and change affected nodes:

$$a_{ij}^{dual} = \begin{cases} a_{ji} & \text{if either } i, j = i_0 \\ a_{ij} - a_{i_0,i} a_{j,i_0} & \text{if both } A \in I_{out}, B \in I_{in} \end{cases}$$

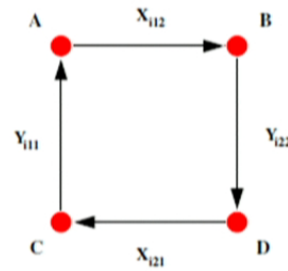
If negative, take it to mean  $-a^{dual}$  arrows from  $j$  to  $i$ .

# Many Examples

F0

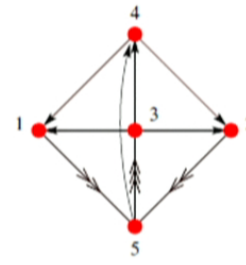


Model I

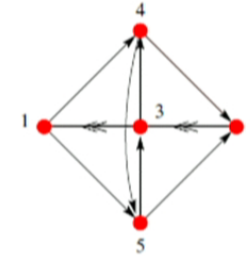


Model II

dP2

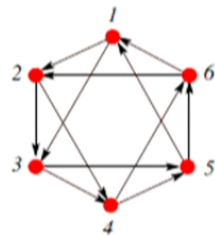


Model I

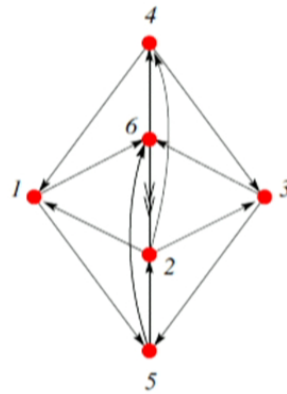


Model II

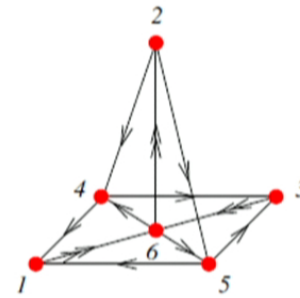
dP3



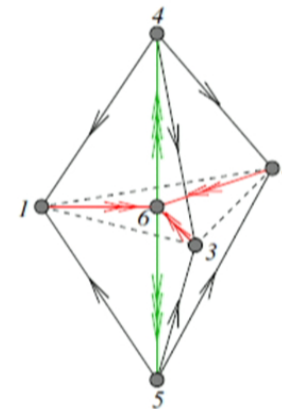
Model I



Model II



Model III

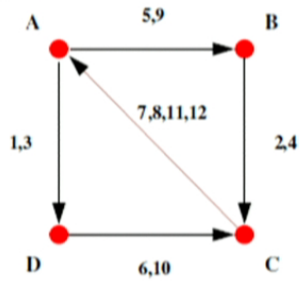


Model IV

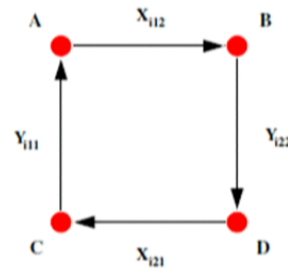


# Many Examples

F0

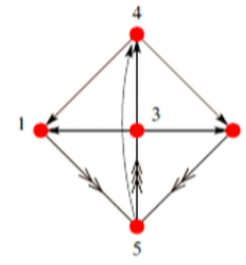


Model I

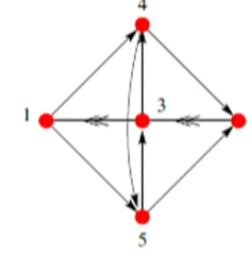


Model II

dP2

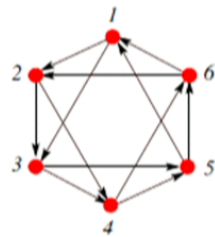


Model I

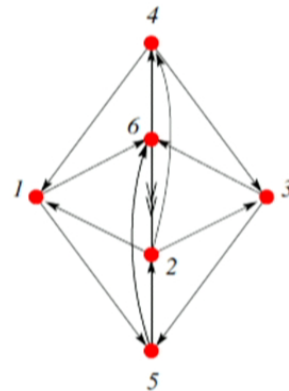


Model II

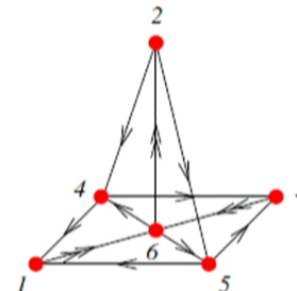
dP3



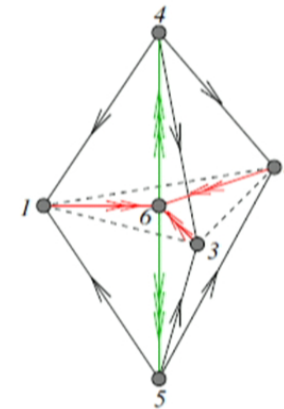
Model I



Model II



Model III

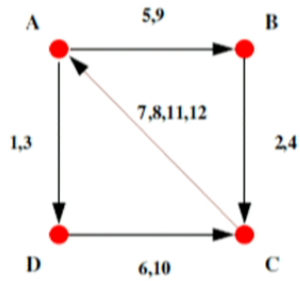


Model IV

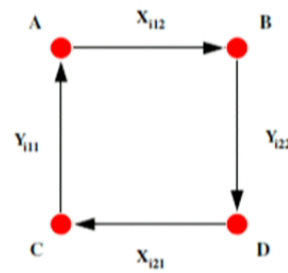


# Many Examples

F0

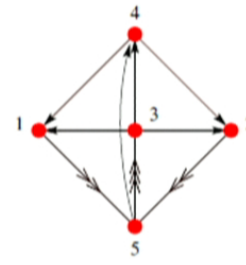


Model I

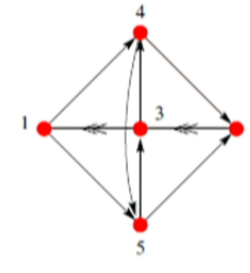


Model II

dP2

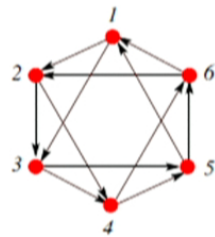


Model I

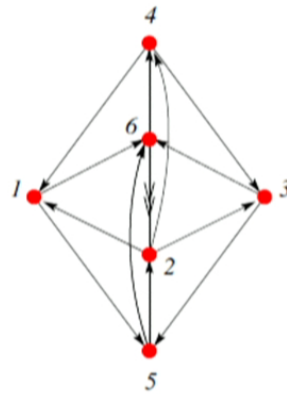


Model II

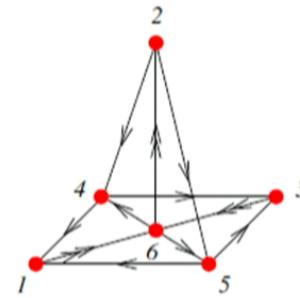
dP3



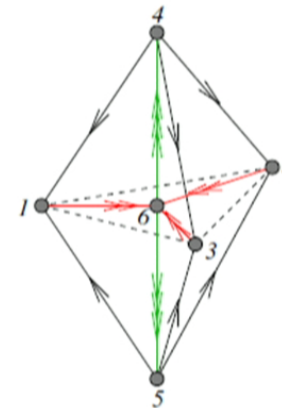
Model I



Model II



Model III



Model IV



# Perspectives on Seiberg Duality

- **Mirror Picture** (Type IIA)

- D6-branes wrapping  $SL-k + 3$  cycles  $S_i$  in the mirror  $Y$
- Quiver = intersection matrix  $A_{ij} = S_i \circ S_j$
- **Picard-Lefschetz** transformations about the  $S_{i_0}$  vanishing cycle:  
$$S_i \rightarrow S_i - (S_i \circ S_{i_0})S_{i_0}$$

- **Derived Category** (Type IIB)

- think of brane as support for coherent sheaf with Mukai vector:  
$$\text{ch}(F_i) := (\text{rk}, c_1, c_2)$$
- Quiver:  $A_{ij} = \chi(F_i, F_j) := \sum_m (-1)^m \dim_{\mathbb{C}} \text{Ext}^m(F_i, F_j)$
- mutation of exceptional collection of  $F_i$

- **Cluster Algebra**

- cluster mutation rules on cluster (matrix) variables
- relation to total positivity and Grassmannian?



# Perspectives on Seiberg Duality

- **Mirror Picture** (Type IIA)

- D6-branes wrapping  $SL-k + 3$  cycles  $S_i$  in the mirror  $Y$
- Quiver = intersection matrix  $A_{ij} = S_i \circ S_j$
- **Picard-Lefschetz** transformations about the  $S_{i_0}$  vanishing cycle:  
$$S_i \rightarrow S_i - (S_i \circ S_{i_0})S_{i_0}$$

- **Derived Category** (Type IIB)

- think of brane as support for coherent sheaf with Mukai vector:  
$$\text{ch}(F_i) := (\text{rk}, c_1, c_2)$$
- Quiver:  $A_{ij} = \chi(F_i, F_j) := \sum_m (-1)^m \dim_{\mathbb{C}} \text{Ext}^m(F_i, F_j)$
- mutation of exceptional collection of  $F_i$

- **Cluster Algebra**

- cluster mutation rules on cluster (matrix) variables
- relation to total positivity and Grassmannian?



# Grothendieck's Dessin d'Enfant

- **Belyĭ Map**: rational map  $\beta : \Sigma \longrightarrow \mathbb{P}^1$  ramified only at  $(0, 1, \infty)$ 
  - **Theorem [Belyĭ]**:  $\beta$  exists  $\Leftrightarrow \Sigma$  can be defined over  $\overline{\mathbb{Q}}$
  - $(\beta, \Sigma)$  Belyĭ Pair
  - **Dessin d'Enfants** =  $\beta^{-1}([0, 1] \in \mathbb{P}^1)$  is a *bi-partite graph* on  $\Sigma$ : label  $\beta^{-1}(0)$  black and  $\beta^{-1}(1)$  white, then  $\beta^{-1}(\infty)$  lives one per face
  - Ramification data: 
$$\left\{ \begin{array}{l} r_0(1), r_0(2), \dots, r_0(B) \\ r_1(1), r_1(2), \dots, r_1(W) \\ r_\infty(1), r_\infty(2), \dots, r_\infty(I) \end{array} \right\}$$
- equivalently, **Permutation Triple**:  $\sigma_B, \sigma_W$  and  $\sigma_B \sigma_W \sigma_\infty = \mathbb{I}$  (encodes how the sheets are permuted at the ramification points; cf. Ramgoolam et al.)
 
$$\sigma_B = (\dots)_{r_0(1)} (\dots)_{r_0(2)} \dots (\dots)_{r_0(B)}, \quad \sigma_W = (\dots)_{r_1(1)} (\dots)_{r_1(2)} \dots (\dots)_{r_1(W)}$$
 Cartographic group:  $\langle \sigma_B, \sigma_W \rangle$

# Grothendieck's Dessin d'Enfant

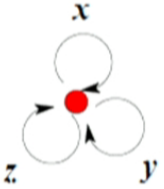
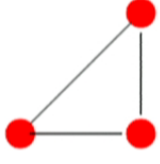
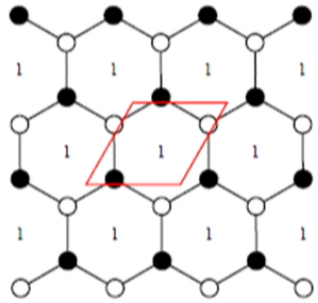
- **Belyĭ Map**: rational map  $\beta : \Sigma \longrightarrow \mathbb{P}^1$  ramified only at  $(0, 1, \infty)$ 
  - **Theorem [Belyĭ]**:  $\beta$  exists  $\Leftrightarrow \Sigma$  can be defined over  $\overline{\mathbb{Q}}$
  - $(\beta, \Sigma)$  Belyĭ Pair
  - **Dessin d'Enfants** =  $\beta^{-1}([0, 1] \in \mathbb{P}^1)$  is a *bi-partite graph* on  $\Sigma$ : label  $\beta^{-1}(0)$  black and  $\beta^{-1}(1)$  white, then  $\beta^{-1}(\infty)$  lives one per face
  - Ramification data:  $\left\{ \begin{array}{l} r_0(1), r_0(2), \dots, r_0(B) \\ r_1(1), r_1(2), \dots, r_1(W) \\ r_\infty(1), r_\infty(2), \dots, r_\infty(I) \end{array} \right\}$
- equivalently, **Permutation Triple**:  $\sigma_B, \sigma_W$  and  $\sigma_B \sigma_W \sigma_\infty = \mathbb{I}$  (encodes how the sheets are permuted at the ramification points; cf. Ramgoolam et al.)
 
$$\sigma_B = (\dots)_{r_0(1)} (\dots)_{r_0(2)} \dots (\dots)_{r_0(B)}, \quad \sigma_W = (\dots)_{r_1(1)} (\dots)_{r_1(2)} \dots (\dots)_{r_1(W)}$$
 Cartographic group:  $\langle \sigma_B, \sigma_W \rangle$

# Grothendieck's Dessin d'Enfant

- **Belyĭ Map**: rational map  $\beta : \Sigma \longrightarrow \mathbb{P}^1$  ramified only at  $(0, 1, \infty)$ 
  - **Theorem [Belyĭ]**:  $\beta$  exists  $\Leftrightarrow \Sigma$  can be defined over  $\overline{\mathbb{Q}}$
  - $(\beta, \Sigma)$  Belyĭ Pair
  - **Dessin d'Enfants** =  $\beta^{-1}([0, 1] \in \mathbb{P}^1)$  is a *bi-partite graph* on  $\Sigma$ : label  $\beta^{-1}(0)$  black and  $\beta^{-1}(1)$  white, then  $\beta^{-1}(\infty)$  lives one per face
  - Ramification data:  $\left\{ \begin{array}{l} r_0(1), r_0(2), \dots, r_0(B) \\ r_1(1), r_1(2), \dots, r_1(W) \\ r_\infty(1), r_\infty(2), \dots, r_\infty(I) \end{array} \right\}$
- equivalently, **Permutation Triple**:  $\sigma_B, \sigma_W$  and  $\sigma_B \sigma_W \sigma_\infty = \mathbb{I}$  (encodes how the sheets are permuted at the ramification points; cf. Ramgoolam et al.)
 
$$\sigma_B = (\dots)_{r_0(1)} (\dots)_{r_0(2)} \dots (\dots)_{r_0(B)}, \quad \sigma_W = (\dots)_{r_1(1)} (\dots)_{r_1(2)} \dots (\dots)_{r_1(W)}$$
 Cartographic group:  $\langle \sigma_B, \sigma_W \rangle$

# Gauge Theories and Dessins

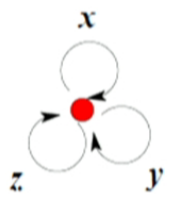
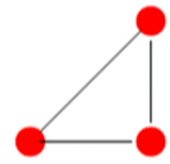
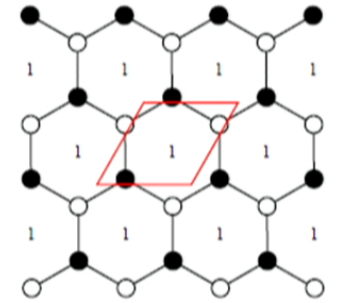
- Matter content and interaction of SUSY gauge theory with toric moduli space is specified by Belyĭ pair
- Our most familiar example of  $\mathcal{N} = 4$  super-Yang-Mills:

Theory	Toric Diag	Belyĭ Pair	Dessin on $T^2$ (dimer)
 <p><math>W = \text{Tr}(X[Y, Z])</math></p>	 <p><math>\mathcal{M} \simeq \mathbb{C}^3</math></p>	$y^2 = x^3 + 1$ $\beta(x, y) = \frac{y+1}{2}$	

- Rmk: Absolute Galois group  $Gal(\overline{\mathbb{Q}}/\mathbb{Q})$  acts *faithfully* over the dessin, even on subsets like dessin on  $\mathbb{P}^1$  or  $T^2$ ...

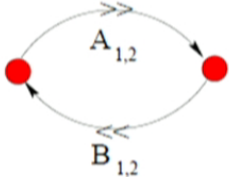
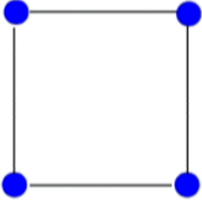
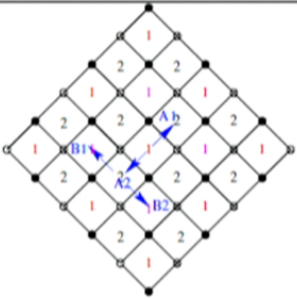
# Gauge Theories and Dessins

- Matter content and interaction of SUSY gauge theory with toric moduli space is specified by Belyi pair
- Our most familiar example of  $\mathcal{N} = 4$  super-Yang-Mills:

Theory	Toric Diag	Belyi Pair	Dessin on $T^2$ (dimer)
 <p><math>W = \text{Tr}(X[Y, Z])</math></p>	 <p><math>\mathcal{M} \simeq \mathbb{C}^3</math></p>	$y^2 = x^3 + 1$ $\beta(x, y) = \frac{y+1}{2}$	

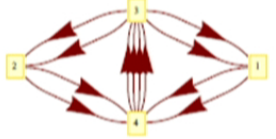
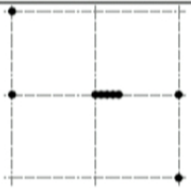
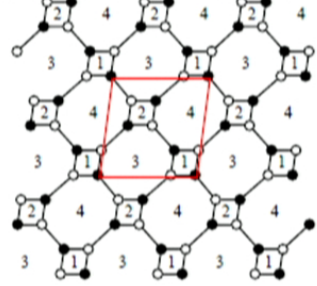
- Rmk: Absolute Galois group  $Gal(\overline{\mathbb{Q}}/\mathbb{Q})$  acts *faithfully* over the dessin, even on subsets like dessin on  $\mathbb{P}^1$  or  $T^2$ ...

- Klebanov-Witten's Conifold Theory

Theory	Toric Diag
 $W = \text{Tr}(\epsilon_{il}\epsilon_{jk}A_iB_jA_lB_k)$	 $\mathcal{M} \simeq \{uv - wz = 0\} \subset \mathbb{C}^4$
Belyi Pair	Dessin on $T^2$ (dimer)
$y^2 = x(x - 1)(x - \frac{1}{2})$ $\beta(x, y) = \frac{x^2}{2x-1}$	

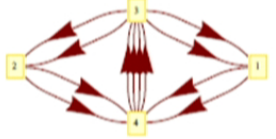
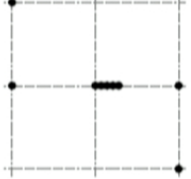
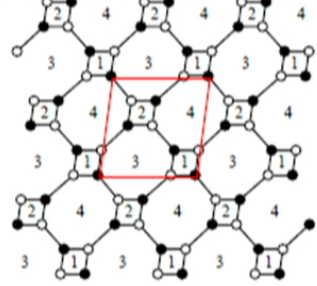
# Plethora of Non-Trivial Examples

e.g., Cone over  $F_0 \simeq \mathbb{P}^1 \times \mathbb{P}^1$  (zeroth Hirzebruch surface);

Theory	Toric Diag
 $w = x_{14}^1 x_{43}^4 x_{31}^2 + x_{14}^2 x_{43}^2 x_{31}^1 + x_{24}^1 x_{43}^1 x_{32}^1 + x_{24}^2 x_{43}^3 x_{32}^2 - x_{43}^1 x_{31}^1 x_{14}^1 - x_{43}^2 x_{32}^2 x_{24}^1 - x_{43}^3 x_{31}^2 x_{14}^2 - x_{43}^4 x_{32}^1 x_{24}^2$	 <p><math>\mathcal{M} \simeq \text{Hirzebruch } 0</math></p>
Belyĭ Pair	Dessin on $T^2$ (dimer)
$y^2 = x - x^3$ $\beta(x, y) = \frac{i(x^2 - \sqrt[3]{-1})^3}{3\sqrt{3}x^2(x^2 - 1)}$	

# Plethora of Non-Trivial Examples

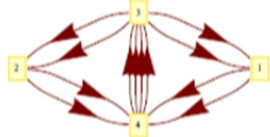
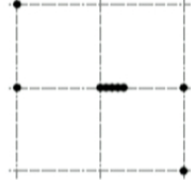
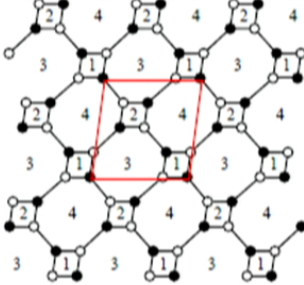
e.g., Cone over  $F_0 \simeq \mathbb{P}^1 \times \mathbb{P}^1$  (zeroth Hirzebruch surface);

Theory	Toric Diag
 $w = x_{14}^1 x_{43}^4 x_{31}^2 + x_{14}^2 x_{43}^2 x_{31}^1 + x_{24}^1 x_{43}^1 x_{32}^1 + x_{24}^2 x_{43}^3 x_{32}^2 - x_{43}^1 x_{31}^1 x_{14}^1 - x_{43}^2 x_{32}^2 x_{24}^1 - x_{43}^3 x_{31}^2 x_{14}^2 - x_{43}^4 x_{32}^1 x_{24}^2$	 <p><math>\mathcal{M} \simeq \text{Hirzebruch } 0</math></p>
Belyĭ Pair	Dessin on $T^2$ (dimer)
$y^2 = x - x^3$ $\beta(x, y) = \frac{i(x^2 - \sqrt[3]{-1})^3}{3\sqrt{3}x^2(x^2 - 1)}$	



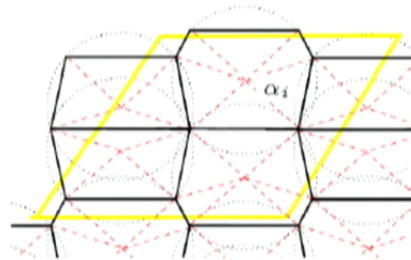
# Plethora of Non-Trivial Examples

e.g., Cone over  $F_0 \simeq \mathbb{P}^1 \times \mathbb{P}^1$  (zeroth Hirzebruch surface);

Theory	Toric Diag
	
$w = x_{14}^1 x_{43}^4 x_{31}^2 + x_{14}^2 x_{43}^2 x_{31}^1 + x_{24}^1 x_{43}^1 x_{32}^1 + x_{24}^2 x_{43}^3 x_{32}^2 - x_{43}^1 x_{31}^1 x_{14}^1 - x_{43}^2 x_{32}^2 x_{24}^1 - x_{43}^3 x_{31}^2 x_{14}^2 - x_{43}^4 x_{32}^1 x_{24}^2$	$\mathcal{M} \simeq \text{Hirzebruch } 0$
Belyĭ Pair	Dessin on $T^2$ (dimer)
$y^2 = x - x^3$ $\beta(x, y) = \frac{i(x^2 - \sqrt[3]{-1})^3}{3\sqrt{3}x^2(x^2 - 1)}$	

# Rigidity

- Dessins are rigid: Belyi pair has no complex structure (frozen at algebraic points in moduli space); in particular elliptic curve has exact  $\tau$
- There is a distinguished dimer in the gauge theory:
  - **NSVZ beta-function**  $\beta(g_a) \sim 3 - \frac{1}{2} \sum_{i \in E_a} (1 - \gamma(X_i))$   
 Conformal dimension =  $\Delta(X_i) = \frac{3}{2}R(X_i) = 1 + \frac{1}{2}\gamma(X_i)$  ( $R(X_i)$  is **R-charge**)
  - Conformality  $\rightsquigarrow \beta(g_a) = 0 \rightsquigarrow \sum_{i \in E_a} (1 - R(X_i)) = 2$  (\*)
  - $W$ -terms have R-charge 2  $\rightsquigarrow \sum_e R(X_e) = 2$  (†)
  - rmk: sum (\*) over nodes + sum (†) over faces  $\rightsquigarrow G - E + N_T = 0$

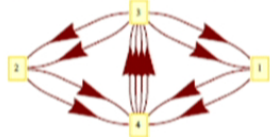
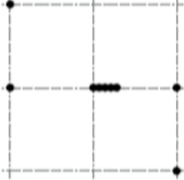
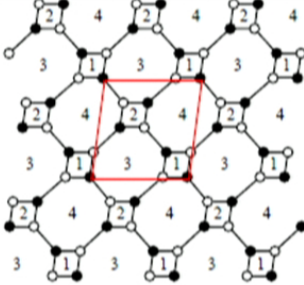


- Set  $R(X_i) = \frac{\alpha_i}{\pi}$
- all parallelgrams are rhombi
- **isoradial embedding**



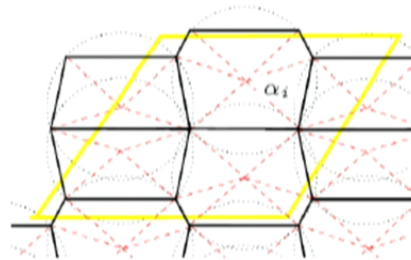
# Plethora of Non-Trivial Examples

e.g., Cone over  $F_0 \simeq \mathbb{P}^1 \times \mathbb{P}^1$  (zeroth Hirzebruch surface);

Theory	Toric Diag
	
$w = x_{14}^1 x_{43}^4 x_{31}^2 + x_{14}^2 x_{43}^2 x_{31}^1 + x_{24}^1 x_{43}^1 x_{32}^1 + x_{24}^2 x_{43}^3 x_{32}^2 - x_{43}^1 x_{31}^1 x_{14}^1 - x_{43}^2 x_{32}^2 x_{24}^1 - x_{43}^3 x_{31}^2 x_{14}^2 - x_{43}^4 x_{32}^1 x_{24}^2$	$\mathcal{M} \simeq \text{Hirzebruch } 0$
Belyĭ Pair	Dessin on $T^2$ (dimer)
$y^2 = x - x^3$ $\beta(x, y) = \frac{i(x^2 - \sqrt[3]{-1})^3}{3\sqrt{3}x^2(x^2 - 1)}$	

# Rigidity

- Dessins are rigid: Belyi pair has no complex structure (frozen at algebraic points in moduli space); in particular elliptic curve has exact  $\tau$
- There is a distinguished dimer in the gauge theory:
  - **NSVZ beta-function**  $\beta(g_a) \sim 3 - \frac{1}{2} \sum_{i \in E_a} (1 - \gamma(X_i))$   
 Conformal dimension =  $\Delta(X_i) = \frac{3}{2}R(X_i) = 1 + \frac{1}{2}\gamma(X_i)$  ( $R(X_i)$  is **R-charge**)
  - Conformality  $\rightsquigarrow \beta(g_a) = 0 \rightsquigarrow \sum_{i \in E_a} (1 - R(X_i)) = 2$  (\*)
  - $W$ -terms have R-charge 2  $\rightsquigarrow \sum_e R(X_e) = 2$  (†)
  - rmk: sum (\*) over nodes + sum (†) over faces  $\rightsquigarrow G - E + N_T = 0$

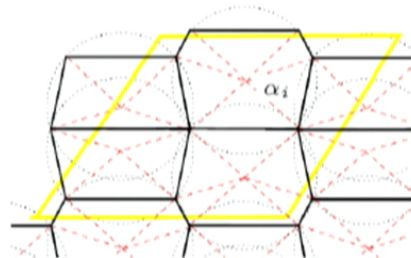


- Set  $R(X_i) = \frac{\alpha_i}{\pi}$
- all parallelgrams are rhombi
- **isoradial embedding**



# Rigidity

- Dessins are rigid: Belyi pair has no complex structure (frozen at algebraic points in moduli space); in particular elliptic curve has exact  $\tau$
- There is a distinguished dimer in the gauge theory:
  - **NSVZ beta-function**  $\beta(g_a) \sim 3 - \frac{1}{2} \sum_{i \in E_a} (1 - \gamma(X_i))$   
 Conformal dimension =  $\Delta(X_i) = \frac{3}{2}R(X_i) = 1 + \frac{1}{2}\gamma(X_i)$  ( $R(X_i)$  is **R-charge**)
  - Conformality  $\rightsquigarrow \beta(g_a) = 0 \rightsquigarrow \sum_{i \in E_a} (1 - R(X_i)) = 2$  (\*)
  - $W$ -terms have R-charge 2  $\rightsquigarrow \sum_e R(X_e) = 2$  (†)
  - rmk: sum (\*) over nodes + sum (†) over faces  $\rightsquigarrow G - E + N_T = 0$



- Set  $R(X_i) = \frac{\alpha_i}{\pi}$
- all parallelgrams are rhombi
- **isoradial embedding**

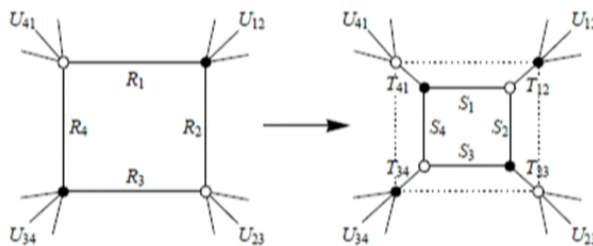


# Square move, Seiberg Duality & Isogeny

- To fix R-charges:  $a$ -maximization or volume  $Z$ -minimization of SE (Intriligator-Wecht, Martelli-Sparks-Yau);

$$a(R) \sim 3 \operatorname{Tr} R^3 - \operatorname{Tr} R \sim \sum_i (R_i - 1)^3 \text{ need to maximize } a(R)$$

- Hanany, YHH, Jejjala, Pasukonis, Ramgoolam, Rodriguez-Gomez
- R-charges and normalized volume of dual geometry are *algebraic numbers*
- Many open puzzles** : e.g. in some cases  $\tau(\text{isoradial}) = \tau(\text{dessin})$ , why? match  $\tau(\text{dimer})$  &  $T^2$  in mirror  $T^3$ -fibration?
- Seiberg Duality = so-called “Urban Renewal”



- $j(\tau)$  of isoradial dimer invariant: elliptic curve for dessins isogenous
- transcendence deg over  $\mathbb{Q}$  inv