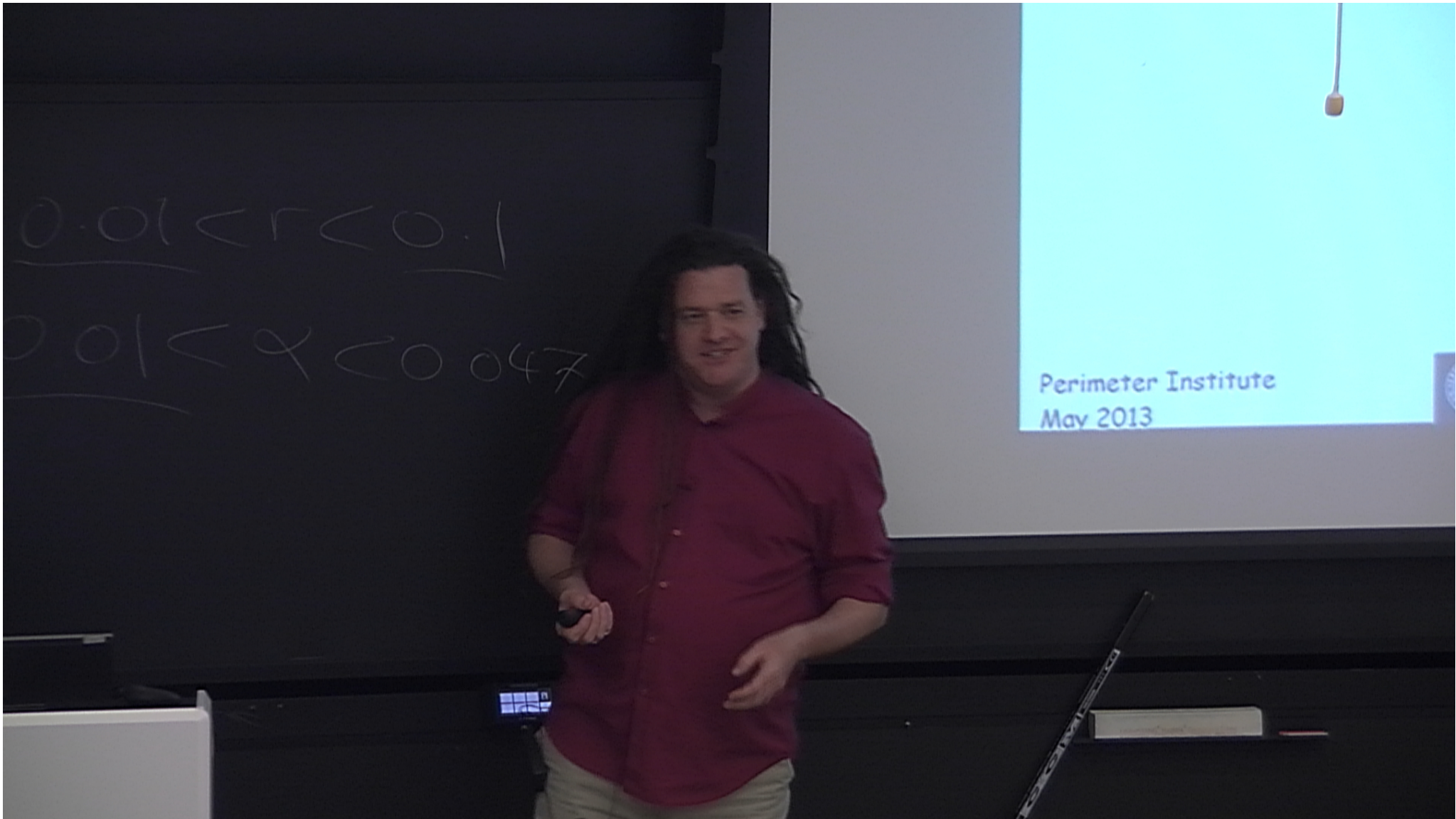


Title: Limitations on the psi-epistemic view of quantum states

Date: May 07, 2013 03:30 PM

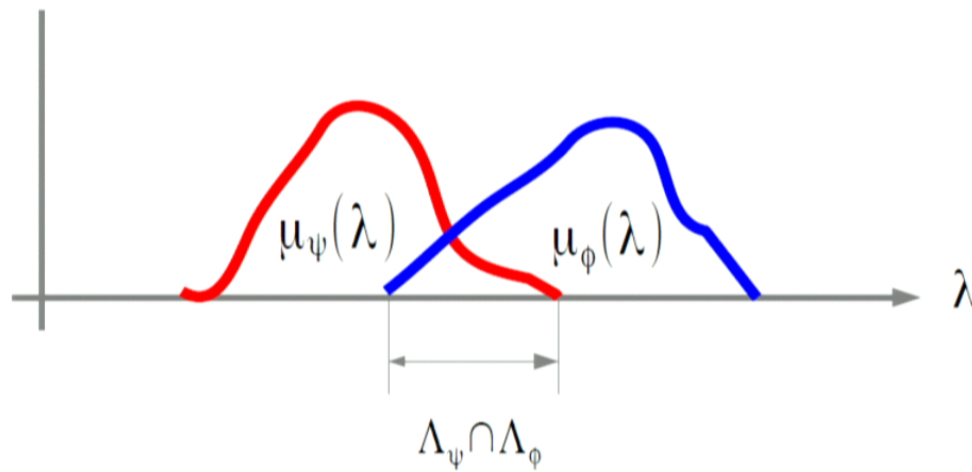
URL: <http://pirsa.org/13050018>

Abstract: <span>The "psi-epistemic" view is that the quantum state does not represent a state of the world, but a state of knowledge about the world. &nbsp;It is motivated, in part, by the observation of qualitative similarities between characteristic properties of non-orthogonal quantum wavefunctions and between overlapping classical probability distributions. &nbsp;It might be suggested that this gives a natural explanation for these properties, which seem puzzling for the alternative "psi-ontic" view. &nbsp;I will examine two such similarities, quantum state overlap and quantum state discrimination, and ask how far can we reproduce the quantitative values given by quantum theory. &nbsp;It will be shown that the psi-epistemic view cannot properly account for the quantitative values, and so must still rely on the same kind of explanations as the "psi-ontic" view.<strongr><br></strongr></span>





# Limitations on the $\Psi$ -epistemic view of quantum states



# Some Recent Activity

MF Pusey, J Barrett, and T Rudolph, *Nature Physics* **8**, 476

MJW Hall, arXiv:1111.6304

PG Lewis, D Jennings, J Barrett, and T Rudolph. arXiv:1201.6554

DJ Miller, arXiv:1202.6465

R Colbeck and R Renner, *Physical Review Letters* **108**, 150402

M Schlosshauer and A Fine, *Physical Review Letters* **108**, 260404

L Hardy, arXiv:1205.1439v2.

**OJE Maroney, arXiv:1207.6906**

MS Leifer, OJE Maroney, *Physical Review Letters* **110**, 120401

MK Patra, S Pironio, S Massar, arXiv:1211.1179

S Aaronson, A Bouland, L Chua and G Lowther, arXiv:1303.2834

... amongst others....

Perimeter Institute  
May 2013



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Limitations on the  $\Psi$ -epistemic  
view of quantum states

- $\Psi$ -epistemic theories
  - What are they? Why are they?
  - PBR and Hardy theorems
- Quantum State Overlap
  - Epistemic vs. Ontic explanation
  - Limitations of the epistemic explanation
- Quantum State Discrimination
  - Epistemic vs. Ontic explanation
  - Limitations of the epistemic explanation
- Where does that leave us?

# Why $\Psi$ -epistemic?

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# Ontological models for Preparations and Measurements

Perform some sequence of physical procedures,  $P$ , that prepares a system to be in some physical state. Label that state:  $\lambda$ .

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Perform some sequence of procedures,  $M$ , on the prepared system, that produces one of a number of distinct outcomes,  $q$ , which occur with a non-trivial relationship to the preparation:  $P(q|M,P)$ .

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An "Ontological Model" seeks to explain this regularity in terms of the "ontic" states  $\lambda$ . The preparation,  $P$ , is associated with a probability distribution over the ontic states  $\mu_P(\lambda)$ , and the measurement with a conditional response probability  $\epsilon(q|\lambda)$ , such that

$$P(q|M,P) = \int d\lambda \mu_P(\lambda) \epsilon_M(q|\lambda)$$



# Ontic Models Framework

Ontic State  $\lambda$       Preparation  $\psi$        $\mu_\psi(\lambda)$       Measurement  $M$ , Outcome,  $\phi$        $\epsilon_M(\phi|\lambda)$

$\lambda$  represents the actual physical state of the world after the preparation. It carries all the factual properties about the system. Any future interactions are based upon  $\lambda$ , and only depend upon  $P$  via the properties of  $\lambda$ .

$$p(\phi|M, \psi) = \int d\lambda \mu_\psi(\lambda) \epsilon_M(\phi|\lambda) = |\langle \phi | \psi \rangle|^2$$

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Wavefunction collapse:       $\lambda \in \{|\psi\rangle\}$        $\mu_\psi(\lambda) = \delta_{|\phi\rangle, |\psi\rangle}$   
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de Broglie Bohm:       $\lambda \in \{(|\psi\rangle, x)\}$        $\mu_\psi(\lambda) = (\delta_{|\psi\rangle, |\psi\rangle}, |\langle \phi | X | \phi \rangle|^2)$   
 $\epsilon_M(\phi|\lambda) \in \{0, 1\}$  contextually

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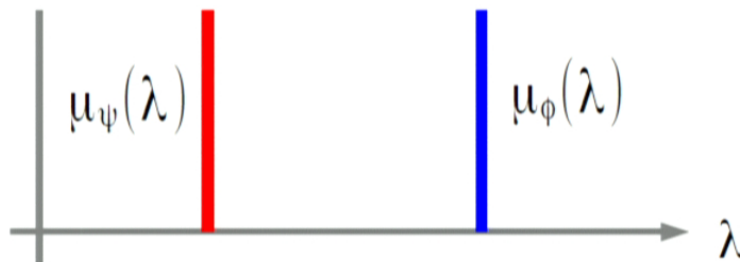
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Must  $\psi$  be part of  $\lambda$ ?



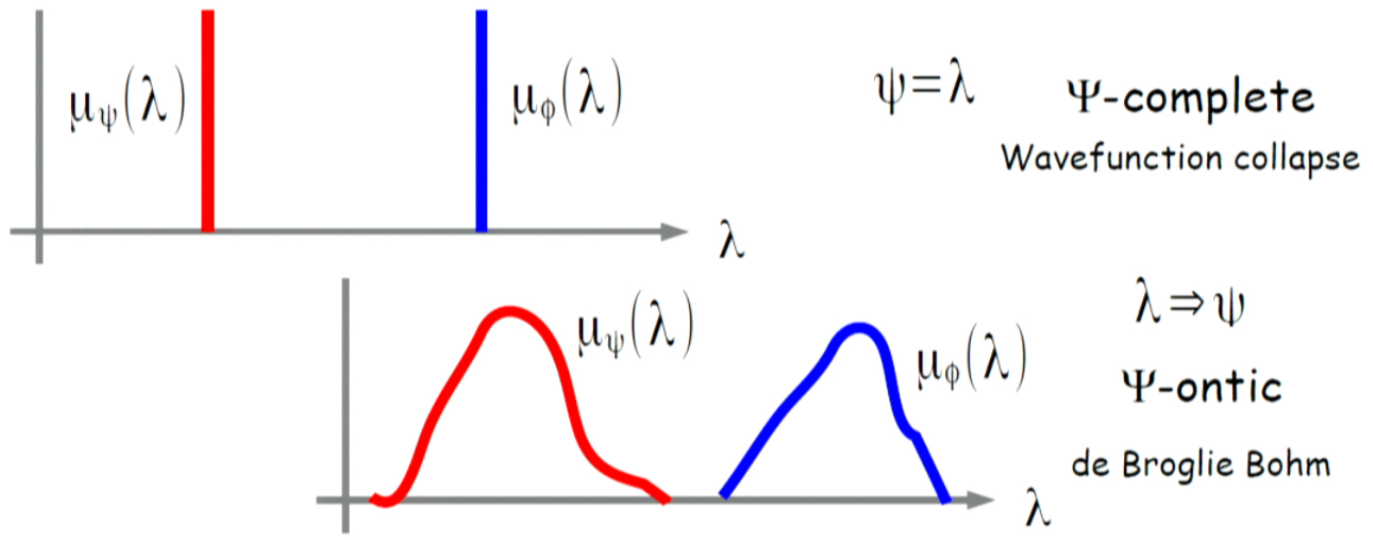
# Is $\psi$ part of $\lambda$ ?



$$\psi = \lambda$$

$\Psi$ -complete  
Wavefunction collapse

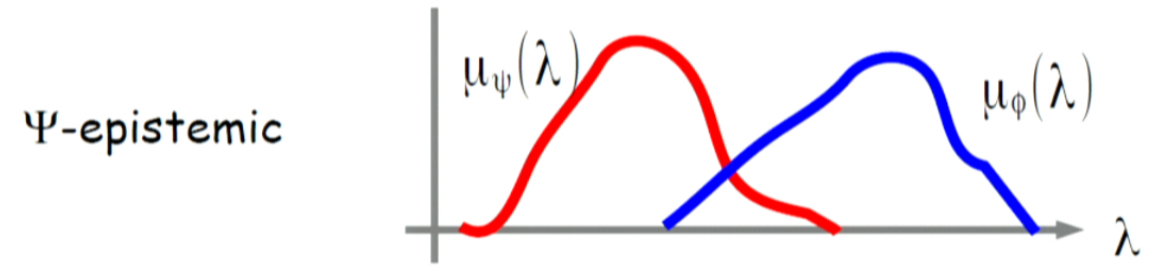
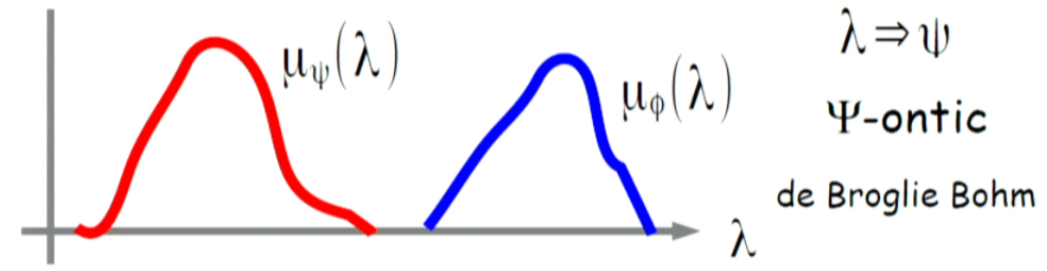
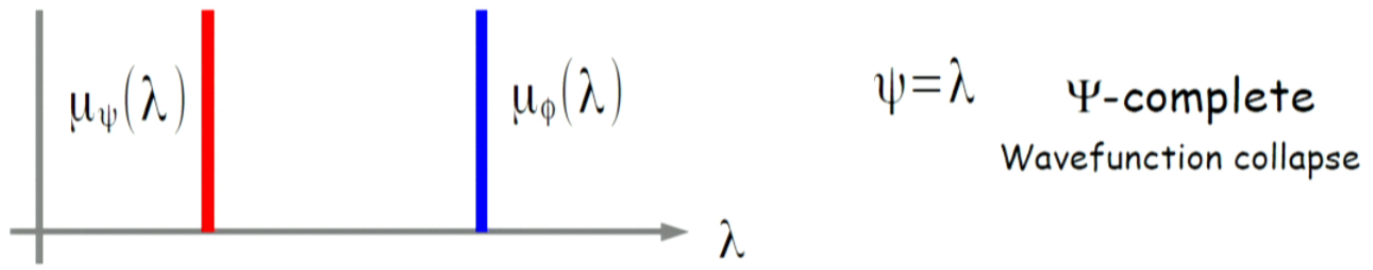
# Is $\psi$ part of $\lambda$ ?



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Limitations on the  $\Psi$ -epistemic view of quantum states

# Is $\psi$ part of $\lambda$ ?



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# Why $\Psi$ -epistemic?

- Qualitative similarities between quantum wavefunctions and classical probability distributions:
  - Indistinguishability of non-orthogonal quantum states
  - Collapse as Bayesian updating
  - Exponential increase in complexity
  - No-cloning of probability distributions
  - Non-unique decomposition of a density matrix
  - et al.
- But can the  $\Psi$ -epistemic point of view account for the (quantitative) quantum properties?



# Some problems

- Hardy's ontic state indifference theorem
  - Operations performed at remote locations must change  $\lambda$  locally
  - A novel form of non-locality for isolated systems.
- PBR's factorisability theorem
  - Preparations of remote product states must be correlated to  $\lambda$  locally
  - A novel form of non-locality for product states.
- If we can accept non-locality for  $\Psi$ -ontic theories, then we cannot dismiss  $\Psi$ -epistemic theories on this basis.
  - Perhaps this is what we should expect from QFT? Even the global vacuum can show Bell Inequality violations in QFT, indicating non-locality of the global ontic state.
  - Local probability distribution may still be stationary under remote operations. Then  $\mu_p(\lambda)$ , at least, can be well defined for the local operations.

# Some (more) problems

- How well does the  $\Psi$ -epistemic view reproduce the *quantitative* properties?
- Indistinguishability of non-orthogonal quantum states
  - Quantum state overlap
    - A test designed to perfectly select a given quantum state, will select a non-orthogonal quantum state with a minimum probability  $a > 0$
  - Quantum state discrimination
    - Two non-orthogonal quantum states can be distinguished in a single test with a maximum probability  $b < 1$

# Quantum State Overlap

Take two preparations,  $\phi$ , and  $\psi$ .

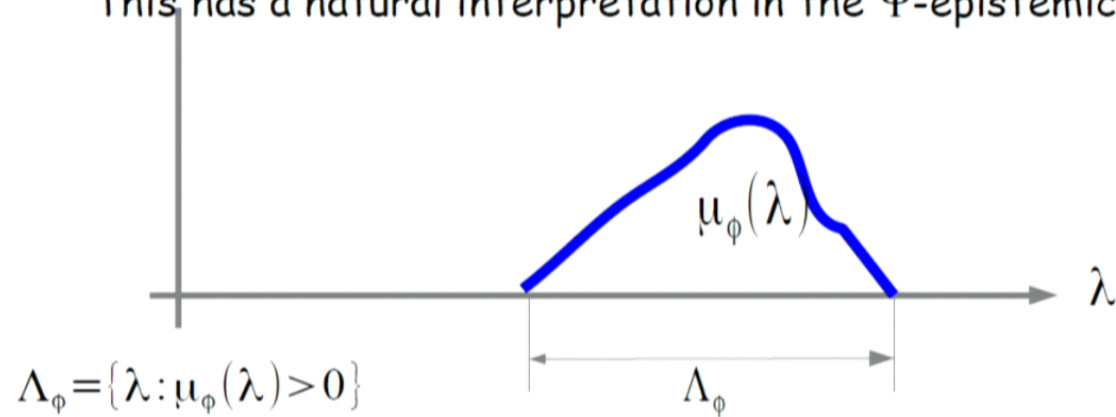
For any measurement for which:  $p(q|M, \phi) = 1$

it turns out  $p(q|M, \psi) \geq a > 0$

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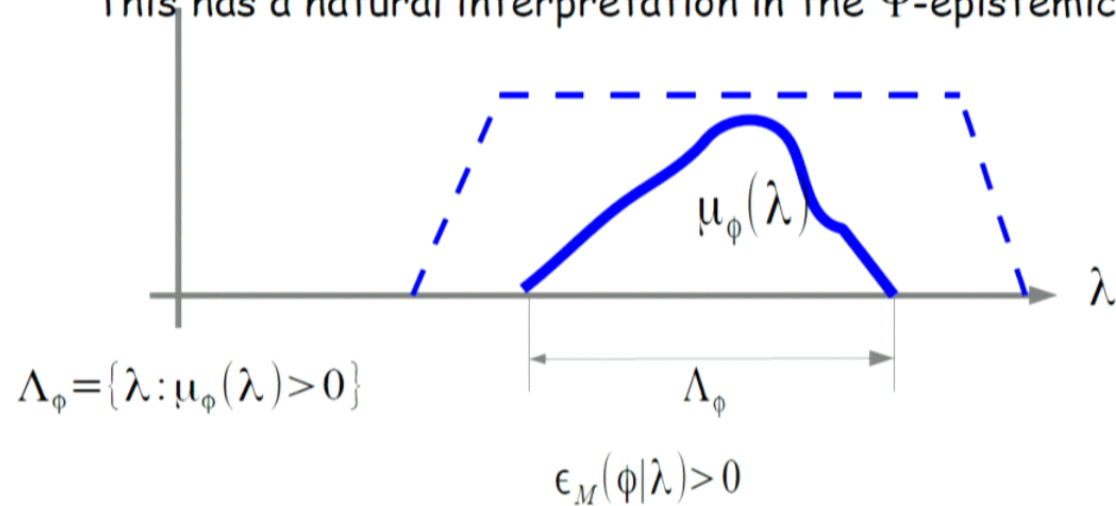
This has a natural interpretation in the  $\Psi$ -epistemic view



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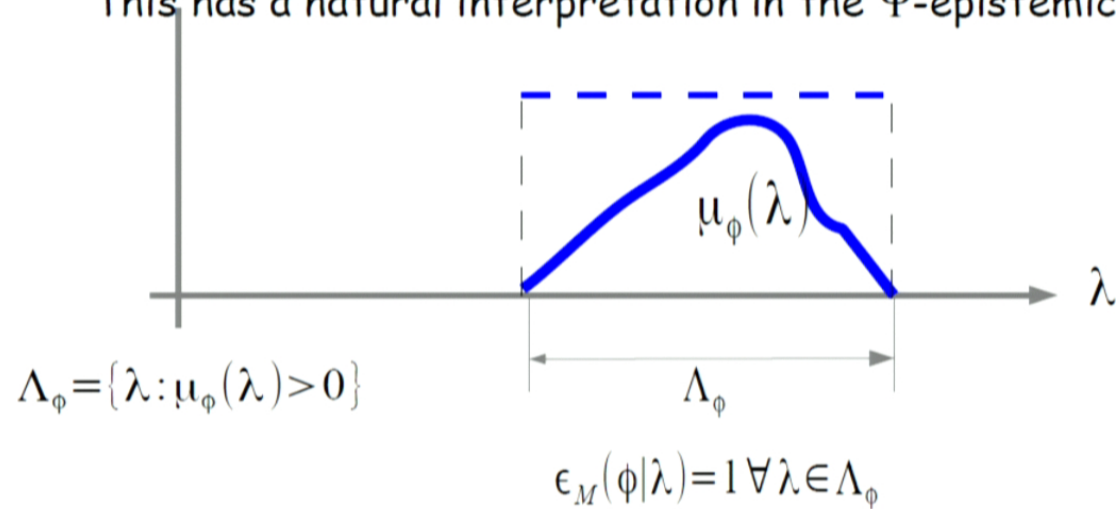
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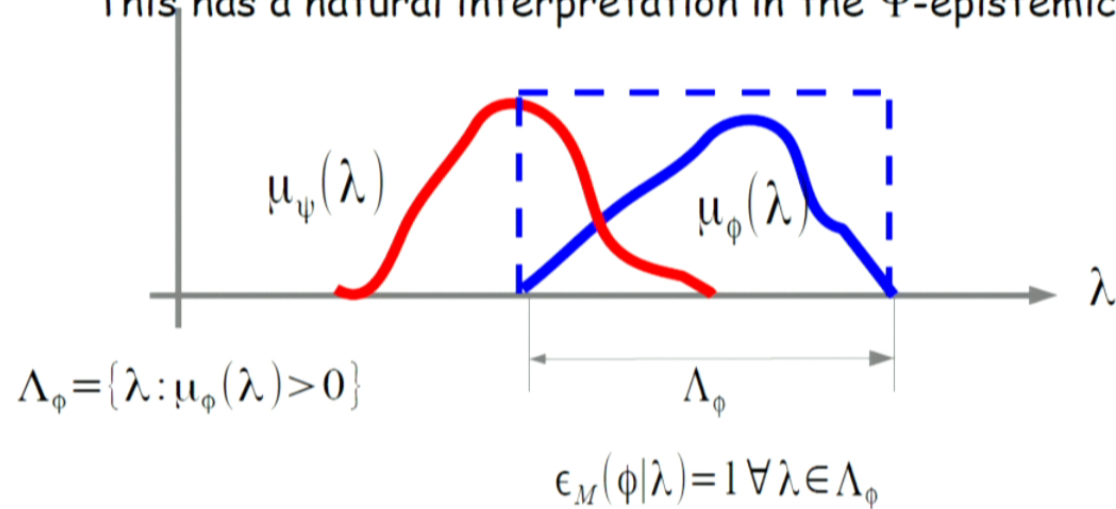




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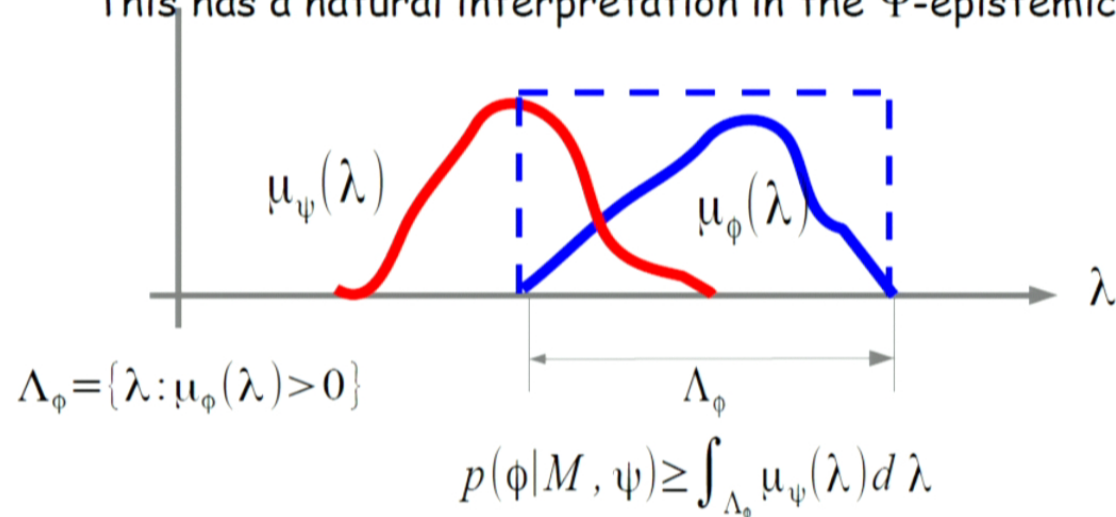
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# Quantum State Overlap

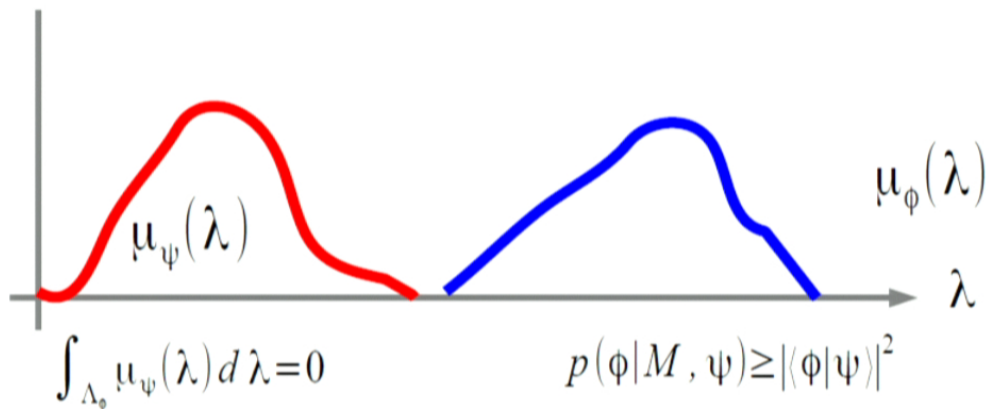
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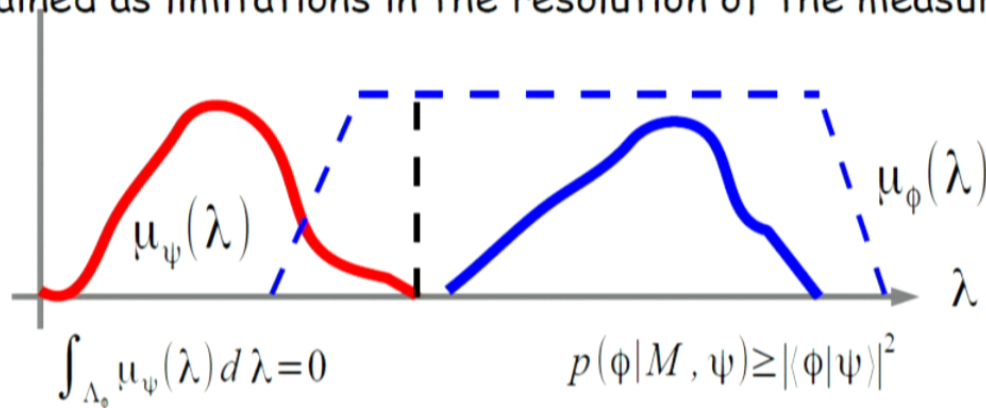
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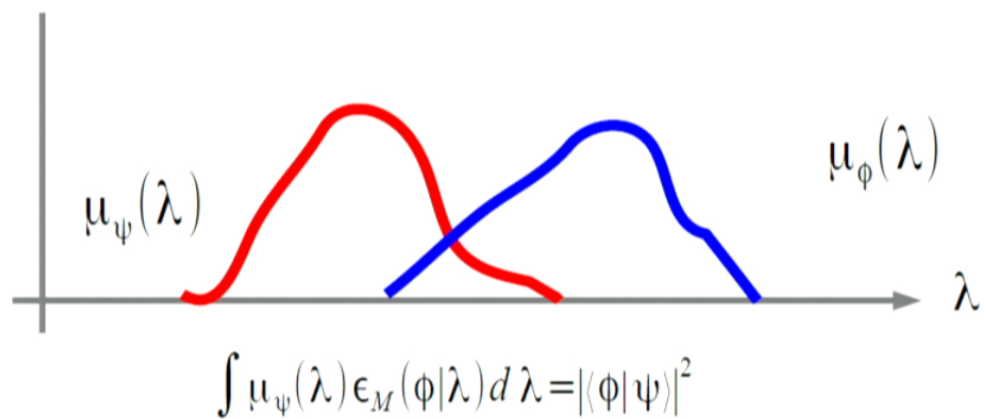
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For the  $\Psi$ -ontic view, the quantum state overlap must be explained as limitations in the resolution of the measurement



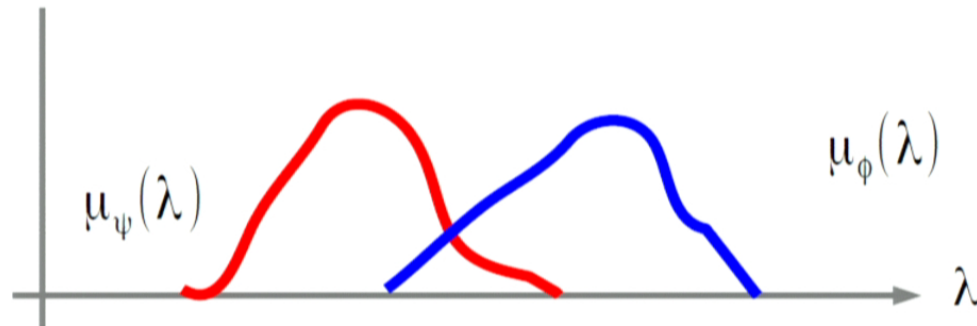
# Maximally $\Psi$ -Epistemic



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# Maximally $\Psi$ -Epistemic



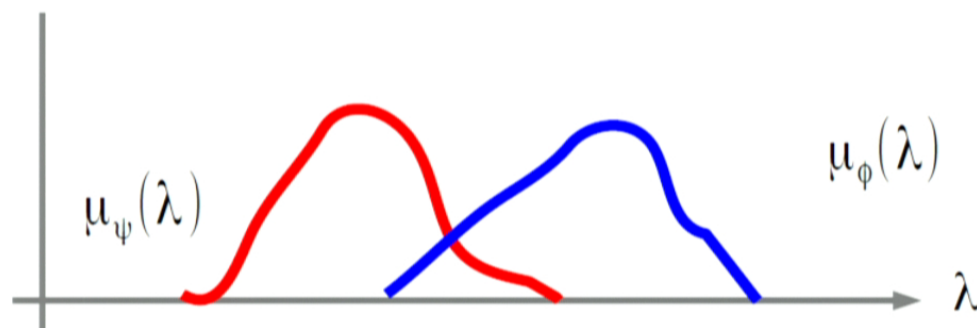
$$\int \mu_\psi(\lambda) \epsilon_M(\phi|\lambda) d\lambda = |\langle \phi | \psi \rangle|^2$$

$$\int_{\Lambda_\phi} \mu_\psi(\lambda) d\lambda = \int_{\Lambda_\phi} \mu_\psi(\lambda) \epsilon_M(\phi|\lambda) d\lambda \leq |\langle \phi | \psi \rangle|^2$$

$$\int_{\Lambda_\phi} \mu_\psi(\lambda) d\lambda = \omega(\phi, \psi) |\langle \phi | \psi \rangle|^2$$



# Maximally $\Psi$ -Epistemic



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$$0 \leq \omega(\phi, \psi) \leq 1$$

$\Psi$ -ontic     $\omega(\phi, \psi) = 0$     Maximally  $\Psi$ -epistemic     $\omega(\phi, \psi) = 1$



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# Quantum Deficiency Theorem

Harrigan and Rudolph (2007)

$$\Lambda_\phi = \{\lambda : \mu_\phi(\lambda) > 0\} \quad \int \mu_\phi(\lambda) \epsilon(\phi|\lambda) d\lambda = 1 \quad \epsilon(\phi|\lambda) = 1, \quad \forall \lambda \in \Lambda_\phi$$

$$\int \mu_\psi(\lambda) \epsilon(\phi|\lambda) d\lambda = \int_{\lambda \in \Lambda_\phi} \mu_\psi(\lambda) d\lambda + \int_{\lambda \notin \Lambda_\phi} \mu_\psi(\lambda) \epsilon(\phi|\lambda) d\lambda = |\langle \psi | \phi \rangle|^2$$

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If maximally  $\Psi$ -epistemic:  $\int_{\lambda \in \Lambda_\phi} \mu_\psi(\lambda) d\lambda = |\langle \psi | \phi \rangle|^2$

$$\int_{\lambda \notin \Lambda_\phi} \mu_\psi(\lambda) \epsilon(\phi|\lambda) d\lambda = 0 \quad \epsilon(\phi|\lambda) = 0, \quad \forall \lambda \notin \Lambda_\phi$$

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$$\epsilon(\phi|\lambda) = \{0, 1\}$$

In a Hilbert space  $d > 2$ , there are no maximally  $\Psi$ -epistemic theories

But between which states and how close can you get?

# How $\Psi$ -epistemic can we get?

$$\int_{\Lambda_\phi} \mu_\psi(\lambda) d\lambda = \omega(\phi, \psi) |\langle \phi | \psi \rangle|^2$$

OJE Maroney, arXiv:1207.6906



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If we assume  $\omega(\phi, \psi)$  is a constant

between all pairs of states:

$$\omega_d \leq \frac{d^2}{2d^2 - 4d + 4} \rightarrow \frac{1}{2} + \frac{1}{d}$$

For any Hilbert space of dimension  $d$ , there must exist pairs of states for which  $\omega(\phi, \psi) \leq \omega_d$



## How $\Psi$ -epistemic can we get?

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We can improve on this!

(Barrett, Cavalcanti, Lal, Maroney, forthcoming)



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# How $\Psi$ -epistemic can we get?

Suppose for some three states  $\phi_a, \phi_b, \phi_c$ , there exists a measurement  $M$ :

	$M$		
	$ q_1\rangle$	$ q_2\rangle$	$ q_3\rangle$
$ \phi_a\rangle$	0	$p(q_2 \phi_a)$	$p(q_3 \phi_a)$
$ \phi_b\rangle$	$p(q_1 \phi_b)$	0	$p(q_3 \phi_b)$
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Caves, Fuchs, Schack (2002): such a basis exists whenever

$$x_{ab} + x_{bc} + x_{ca} \leq 1 \quad (x_{ab} + x_{bc} + x_{ca} - 1) \geq 4x_{ab}x_{bc}x_{ca} \quad x_{ab} = |\langle \phi_a | \phi_b \rangle|^2$$

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If  $\phi_a, \phi_b, \phi_c$ , are drawn from three Mutually Unbiased Bases  $x=1/d$

$$\frac{3}{d} \leq 1 \quad \left(\frac{d-3}{d}\right)^2 \geq \frac{4}{d^3}$$

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If  $\phi_a, \phi_b, \phi_c$ , are drawn from three Mutually Unbiased Bases  $x=1/d$   
 $d \geq 4$

In any prime power dimension Hilbert space there are  $d+1$  such MUB's

# How $\Psi$ -epistemic can we get?

Suppose  $\phi_a, \phi_b, \phi_c$ , are from any three MUBs in a Hilbert space of dimension  $d > 3$ .

There exists a measurement  $M$ :

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# How $\Psi$ -epistemic can we get?

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There exists a measurement  $M$ :

	$M$		
	$ q_1\rangle$	$ q_2\rangle$	$ q_3\rangle$
$ \phi_a\rangle$	0	$p(q_2 \phi_a)$	$p(q_3 \phi_a)$
$ \phi_b\rangle$	$p(q_1 \phi_b)$	0	$p(q_3 \phi_b)$
$ \phi_c\rangle$	$p(q_1 \phi_c)$	$p(q_2 \phi_c)$	0

$$\forall \lambda \in \Lambda_{\phi_a}, \epsilon_M(q_1|\lambda) = 0$$

$$\forall \lambda \in \Lambda_{\phi_b}, \epsilon_M(q_2|\lambda) = 0$$

$$\forall \lambda \in \Lambda_{\phi_c}, \epsilon_M(q_3|\lambda) = 0$$

# How $\Psi$ -epistemic can we get?

Suppose  $\phi_a, \phi_b, \phi_c$ , are from any three MUBs in a Hilbert space of dimension  $d > 3$ .

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$$\forall \lambda \in \Lambda_{\phi_a}, \epsilon_M(q_1|\lambda) = 0$$

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$$\forall \lambda \in \Lambda_{\phi_c}, \epsilon_M(q_3|\lambda) = 0$$

$$\forall \lambda, \sum_q \epsilon_M(q|\lambda) = 1$$

$$\Lambda_{\phi_a} \cap \Lambda_{\phi_b} \cap \Lambda_{\phi_c} = \emptyset$$

## How $\Psi$ -epistemic can we get?

In a  $d$ -dimensional prime power Hilbert space, let  $|\alpha_i\rangle$  be the  $i$ 'th basis state from the  $\alpha$ 'th MUB.

$$\begin{aligned}\Lambda_{\alpha_i} \cap \Lambda_{\alpha_j} &= \emptyset & i \neq j \\ \Lambda_{\alpha_i} \cap \Lambda_{\beta_j} \cap \Lambda_{\gamma_k} &= \emptyset & \alpha \neq \beta \neq \gamma\end{aligned}$$

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Choose any state  $|\alpha_j\rangle$  from the  $\alpha$ 'th basis.

$$1 \geq \int_{\cup_{\beta \neq \alpha, i} \Lambda_{\beta i}} \mu_{\alpha_j}(\lambda) d\lambda = \sum_{\beta \neq \alpha, i} \int_{\Lambda_{\beta i}} \mu_{\alpha_j}(\lambda) d\lambda$$

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$$\int_{\Lambda_{\phi}} \mu_{\psi}(\lambda) d\lambda = \omega(\phi, \psi) |\langle \phi | \psi \rangle|^2$$

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Choose any state  $|\alpha_j\rangle$  from the  $\alpha$ 'th basis.

$$\sum_{\beta \neq \alpha, i} \omega(\beta_i, \alpha_j) |\langle \beta_i | \alpha_j \rangle|^2 \leq 1$$



## How $\Psi$ -epistemic can we get?

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Choose any state  $|\alpha_j\rangle$  from the  $\alpha$ 'th basis.

$$\sum_{\beta \neq \alpha, i} \frac{\omega(\beta i, \alpha j)}{d} \leq 1 \qquad \overline{\omega}_d = \frac{\sum_{\beta \neq \alpha, i} \omega(\beta i, \alpha j)}{d^2} \leq \frac{1}{d}$$

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For any Hilbert space of dimension  $d$ , there must exist pairs of states for which  $\omega(\phi, \psi) \leq 1/d$

# Limitations on $\omega(\psi, \phi)$

- The 'ontic explanation' of a quantum state overlap is in terms of limitations in the resolution of the measurement process.
  - $\Psi$ -ontic theories can account for quantum state overlap entirely in terms of limitations of the measurement process
- The 'epistemic explanation' of a quantum state overlap is in terms of overlaps in the prepared probability distributions.
  - $\Psi$ -epistemic theories cannot account for quantum state overlap entirely in terms of overlaps in probability distributions.
  - In the limit of large Hilbert spaces, the quantum state overlap must still be accounted for entirely due to limitations of the measurement process, some pairs of states.

## How $\Psi$ -epistemic can we get?

In a  $d$ -dimensional prime power Hilbert space, let  $|\alpha_i\rangle$  be the  $i$ 'th basis state from the  $\alpha$ 'th MUB.

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Choose any state  $|\alpha_j\rangle$  from the  $\alpha$ 'th basis.

$$\sum_{\beta \neq \alpha, i} \frac{\omega(\beta_i, \alpha_j)}{d} \leq 1 \qquad \overline{\omega}_d = \frac{\sum_{\beta \neq \alpha, i} \omega(\beta_i, \alpha_j)}{d^2} \leq \frac{1}{d}$$

For any Hilbert space of dimension  $d$ , there must exist pairs of states for which  $\omega(\phi, \psi) \leq 1/d$



# Limitations on $\omega(\psi, \phi)$

- The 'ontic explanation' of a quantum state overlap is in terms of limitations in the resolution of the measurement process.
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  - In the limit of large Hilbert spaces, the quantum state overlap must still be accounted for entirely due to limitations of the measurement process, some pairs of states.



# How $\Psi$ -epistemic?

## Quantum State Discrimination

Perimeter Institute  
May 2013



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FOUNDATION

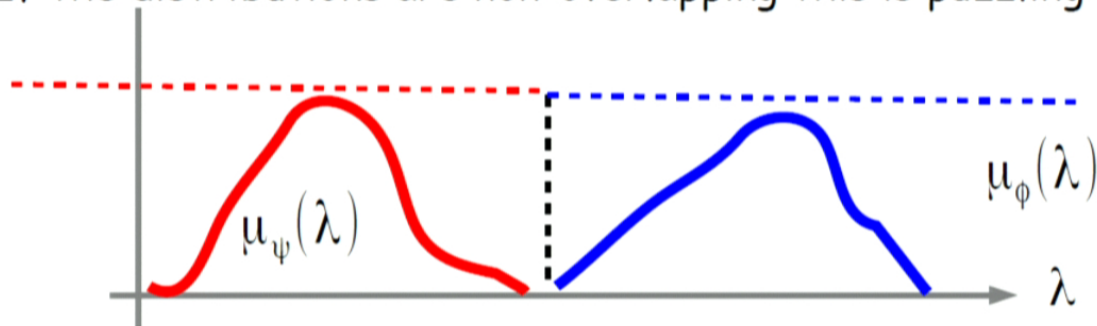
Limitations on the  $\Psi$ -epistemic  
view of quantum states

# Quantum State Discrimination

Take two preparations,  $\phi$ , and  $\psi$ .

Given a system which may have been prepared either way, your best guess as to which cannot succeed better than  $b < 1$

If the distributions are non-overlapping this is puzzling:

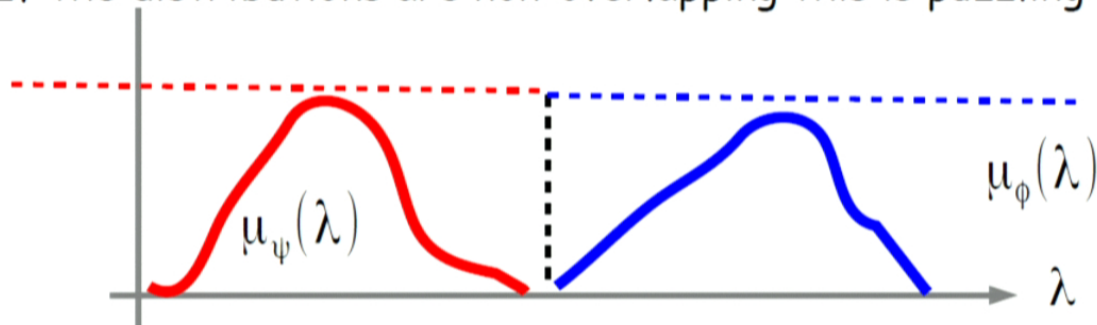


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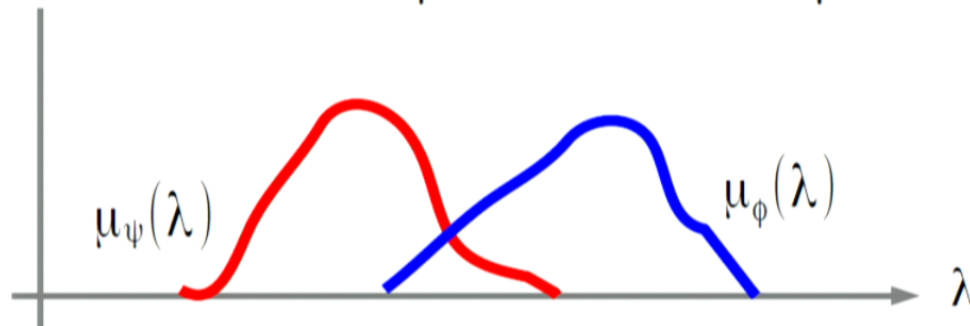


The  $\Psi$ -ontic view must again explain this in terms of limitations of all possible measuring devices.

# Quantum State Discrimination

Take two preparations,  $\phi$ , and  $\psi$ .  
Given a system which may have been prepared either way,  
your best guess as to which cannot succeed better than  $b < 1$

This has a natural interpretation in the  $\Psi$ -epistemic view

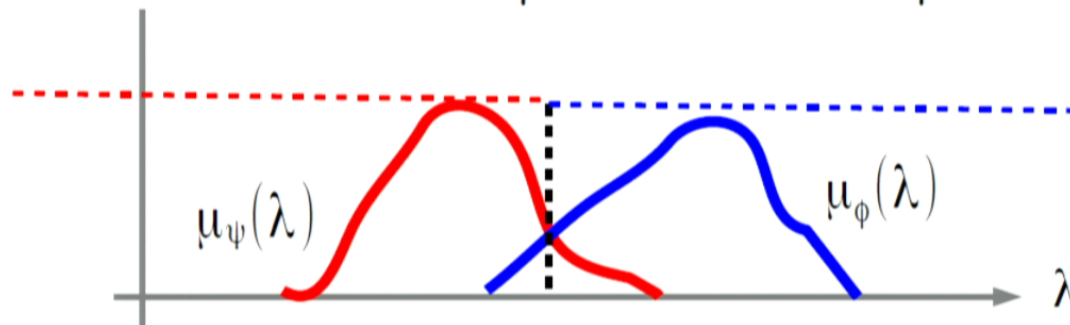


# Quantum State Discrimination

Take two preparations,  $\phi$ , and  $\psi$ .

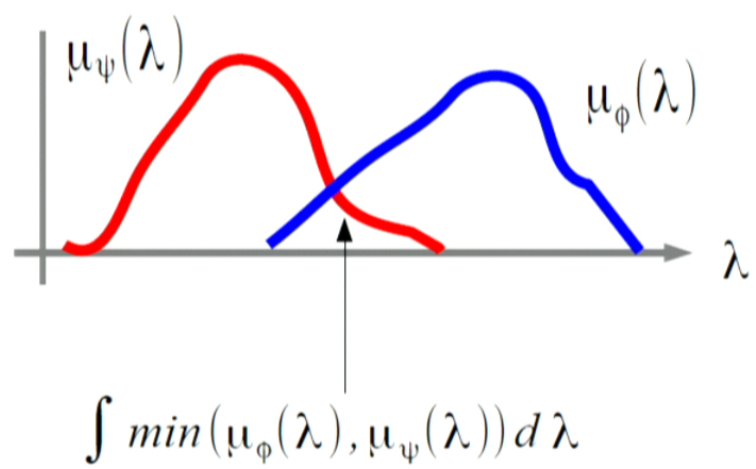
Given a system which may have been prepared either way, your best guess as to which cannot succeed better than  $b < 1$

This has a natural interpretation in the  $\Psi$ -epistemic view



Chance of guessing right:  $1 - \frac{1}{2} \int \min(\mu_\psi(\lambda), \mu_\phi(\lambda)) d\lambda$

# Symmetric overlap



# Symmetric overlap

$$\int \min(\mu_\psi(\lambda), \mu_\phi(\lambda)) \epsilon_M(q|\lambda) d\lambda \leq \int \mu_\psi(\lambda) \epsilon_M(q|\lambda) d\lambda$$
$$\leq \min\left(\int \mu_\psi(\lambda) \epsilon_M(q|\lambda) d\lambda, \int \mu_\phi(\lambda) \epsilon_M(q|\lambda) d\lambda\right)$$

$$\sum_q \epsilon_M(q|\lambda) = 1$$

$$\int \min(\mu_\psi(\lambda), \mu_\phi(\lambda)) d\lambda \leq \sum_q \min(p(q|M, \psi), p(q|M, \phi))$$



# Symmetric overlap

$$\int \min(\mu_\psi(\lambda), \mu_\phi(\lambda)) \epsilon_M(q|\lambda) d\lambda \leq \int \mu_\psi(\lambda) \epsilon_M(q|\lambda) d\lambda$$

$$\leq \min\left(\int \mu_\psi(\lambda) \epsilon_M(q|\lambda) d\lambda, \int \mu_\phi(\lambda) \epsilon_M(q|\lambda) d\lambda\right)$$

$$\sum_q \epsilon_M(q|\lambda) = 1$$

$$\int \min(\mu_\psi(\lambda), \mu_\phi(\lambda)) d\lambda \leq \sum_q \min(p(q|M, \psi), p(q|M, \phi))$$

	$q_1$	$q_2$	$\Sigma$
$\psi$	$p(q_1 \psi)$	$1 - p(q_1 \psi)$	1
$\phi$	$p(q_1 \phi)$	$1 - p(q_1 \phi)$	1
$\min$	$\min(p(q_1 \psi), p(q_1 \phi))$	$\min(1 - p(q_1 \psi), 1 - p(q_1 \phi))$	$1 -  p(q_1 \psi) - p(q_1 \phi) $

# Symmetric overlap

$$\int \min(\mu_\psi(\lambda), \mu_\phi(\lambda)) \epsilon_M(q|\lambda) d\lambda \leq \int \mu_\psi(\lambda) \epsilon_M(q|\lambda) d\lambda$$

$$\leq \min\left(\int \mu_\psi(\lambda) \epsilon_M(q|\lambda) d\lambda, \int \mu_\phi(\lambda) \epsilon_M(q|\lambda) d\lambda\right)$$

$$\sum_q \epsilon_M(q|\lambda) = 1$$

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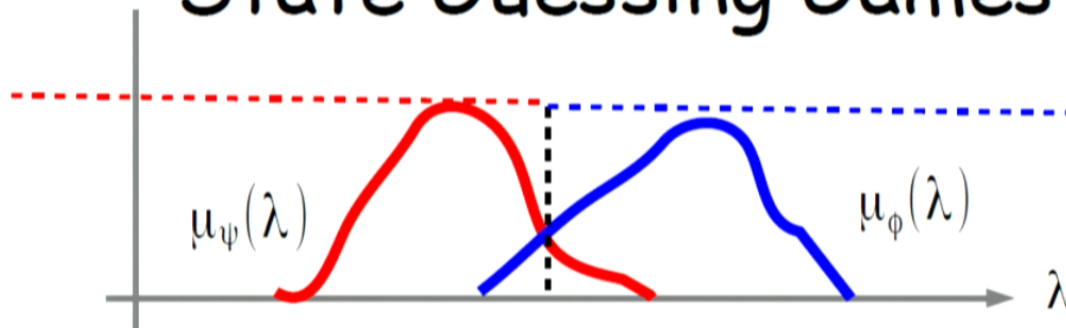
	$q_1$	$q_2$	$\Sigma$
$\psi$	$p(q_1 \psi)$	$1 - p(q_1 \psi)$	1
$\phi$	$p(q_1 \phi)$	$1 - p(q_1 \phi)$	1
$\min$	$\min(p(q_1 \psi), p(q_1 \phi))$	$\min(1 - p(q_1 \psi), 1 - p(q_1 \phi))$	$1 -  p(q_1 \psi) - p(q_1 \phi) $

$$\int \min_i(\mu_{\Phi_i}(\lambda)) \epsilon_M(q|\lambda) d\lambda \leq \int \mu_{\Phi_i}(\lambda) \epsilon_M(q|\lambda) d\lambda$$

$$\leq \min_i\left(\int \mu_{\Phi_i}(\lambda) \epsilon_M(q|\lambda) d\lambda\right)$$

$$\int \min_i(\mu_{\Phi_i}(\lambda)) d\lambda \leq \sum_q \min_i(p(q|M, \Phi_i))$$

# State Guessing Games



Chance of guessing right:  $1 - \frac{1}{2} \int \min(\mu_\psi(\lambda), \mu_\phi(\lambda)) d\lambda$

Working only from the results of a quantum measurement can't improve your chance of guessing right.... So what is the best quantum guess?

# Quantum Best Guess

	$ \psi\rangle$	$ \psi_{\perp}\rangle$
$\psi$	1	0
$\phi$	$ \langle\psi \phi\rangle ^2$	$1- \langle\psi \phi\rangle ^2$
<i>max</i>	1	$1- \langle\psi \phi\rangle ^2$

$$\frac{1}{2} \sum_q \max_i (p(q|M, \psi_i))$$

# Quantum Best Guess

	$ \psi\rangle$	$ \psi^\perp\rangle$
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<i>max</i>	1	$1- \langle\psi \phi\rangle ^2$

$$\frac{1}{2} \sum_q \max_i (p(q|M, \psi_i)) = 1 - \frac{1}{2} |\langle\phi|\psi\rangle|^2$$

$$|u\rangle = \frac{\cos(\alpha/2)|\phi\rangle - \sin(\alpha/2)|\psi\rangle}{\cos\alpha}$$

$$|v\rangle = \frac{\cos(\alpha/2)|\psi\rangle - \sin(\alpha/2)|\phi\rangle}{\cos\alpha}$$

$$\sin\alpha = |\langle\psi|\phi\rangle|$$

	$ u\rangle$	$ v\rangle$
$\psi$	$\cos^2(\alpha/2)$	$\sin^2(\alpha/2)$
$\phi$	$\sin^2(\alpha/2)$	$\cos^2(\alpha/2)$

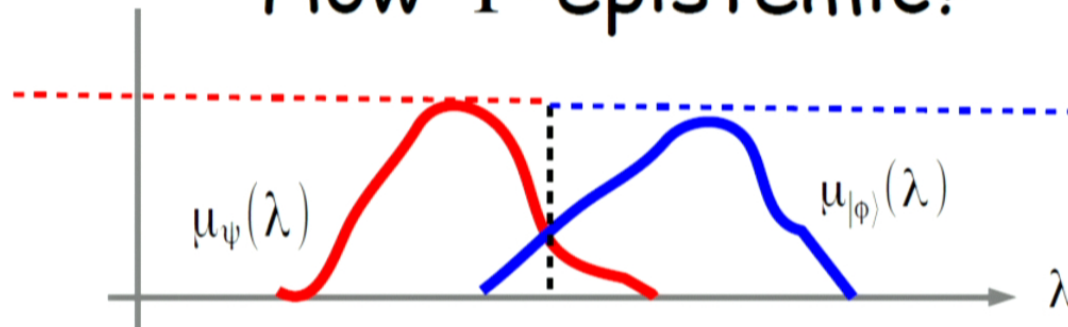
$$\max(2 \sin^2(\alpha/2), 2 \cos^2(\alpha/2)) = 1 + |\cos^2(\alpha/2) - \sin^2(\alpha/2)| = 1 + \sqrt{1 - |\langle\psi|\phi\rangle|^2}$$



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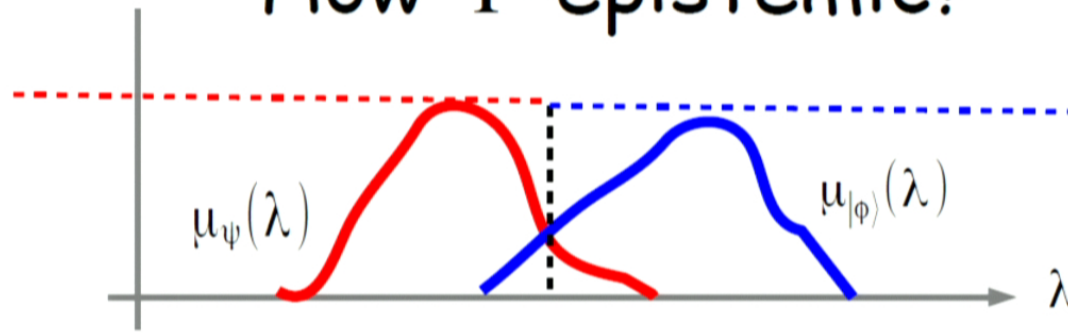
Limitations on the  $\Psi$ -epistemic  
view of quantum states

# How $\Psi$ -epistemic?



$$1 - \frac{1}{2} \int \min(\mu_{|\psi\rangle}(\lambda), \mu_{|\phi\rangle}(\lambda)) d\lambda \geq \frac{1}{2} (1 + \sqrt{1 - |\langle \psi | \phi \rangle|^2})$$

# How $\Psi$ -epistemic?



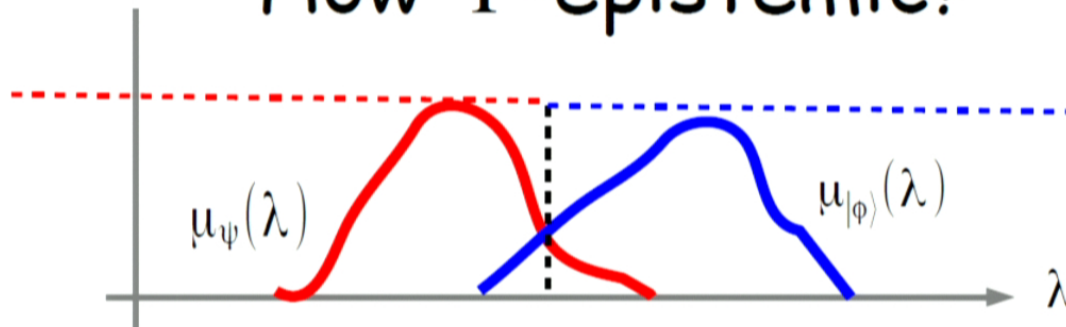
$$1 - \frac{1}{2} \int \min(\mu_\psi(\lambda), \mu_\phi(\lambda)) d\lambda \geq \frac{1}{2} (1 + \sqrt{1 - |\langle \psi | \phi \rangle|^2})$$

$$\int \min(\mu_\psi(\lambda), \mu_\phi(\lambda)) d\lambda \leq (1 - \sqrt{1 - |\langle \psi | \phi \rangle|^2}) < |\langle \psi | \phi \rangle|^2$$

$$\int \min(\mu_\psi(\lambda), \mu_\phi(\lambda)) d\lambda = \varpi(\psi, \phi) (1 - \sqrt{1 - |\langle \psi | \phi \rangle|^2})$$



# How $\Psi$ -epistemic?



$$1 - \frac{1}{2} \int \min(\mu_\psi(\lambda), \mu_\phi(\lambda)) d\lambda \geq \frac{1}{2} (1 + \sqrt{1 - |\langle \psi | \phi \rangle|^2})$$

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$$\int \min(\mu_\psi(\lambda), \mu_\phi(\lambda)) d\lambda = \varpi(\psi, \phi) (1 - \sqrt{1 - |\langle \psi | \phi \rangle|^2})$$

$$0 \leq \varpi(\phi, \psi) \leq 1$$

$\Psi$ -ontic  $\varpi(\phi, \psi) = 0$       Maximally  $\Psi$ -epistemic  $\varpi(\phi, \psi) = 1$



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Limitations on the  $\Psi$ -epistemic  
view of quantum states

# Symmetric overlap

$$\int \min(\mu_\psi(\lambda), \mu_\phi(\lambda)) d\lambda \leq \sum_q \min(p(q|M, \psi), p(q|M, \phi))$$

	$ \psi\rangle$	$ \psi^\perp\rangle$
$\psi$	1	0
$\phi$	$ \langle\psi \phi\rangle ^2$	$1- \langle\psi \phi\rangle ^2$
<i>min</i>	$ \langle\psi \phi\rangle ^2$	0

$$\int \min(\mu_\psi(\lambda), \mu_\phi(\lambda)) d\lambda \leq \int_{\Lambda_\phi} \mu_\psi(\lambda) d\lambda \leq |\langle\psi|\phi\rangle|^2$$

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$$|u\rangle = \frac{\cos(\alpha/2)|\phi\rangle - \sin(\alpha/2)|\psi\rangle}{\cos\alpha}$$

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$$\min(2\sin^2(\alpha/2), 2\cos^2(\alpha/2)) = 1 - |\cos^2(\alpha/2) - \sin^2(\alpha/2)| = 1 - \sqrt{1 - |\langle\psi|\phi\rangle|^2}$$

$$\int \min(\mu_\psi(\lambda), \mu_\phi(\lambda)) d\lambda \leq 1 - \sqrt{1 - |\langle\psi|\phi\rangle|^2}$$

# How $\Psi$ -epistemic?

Suppose  $\phi_a, \phi_b, \phi_c$ , are from any three MUBs in a Hilbert space of dimension  $d > 3$ .

There exists a measurement  $M$ :

	$M$		
	$ q_1\rangle$	$ q_2\rangle$	$ q_3\rangle$
$ \phi_a\rangle$	0	$p(q_2 \phi_a)$	$p(q_3 \phi_a)$
$ \phi_b\rangle$	$p(q_1 \phi_b)$	0	$p(q_3 \phi_b)$
$ \phi_c\rangle$	$p(q_1 \phi_c)$	$p(q_2 \phi_c)$	0

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$ \phi_c\rangle$	$p(q_1 \phi_c)$	$p(q_2 \phi_c)$	0

$$\int \min_i(\mu_{\psi_i}(\lambda)) d\lambda \leq \sum_q \min_i(p(q|M, \psi_i))$$

$$\int \min(\mu_{\alpha i}(\lambda), \mu_{\beta j}(\lambda), \mu_{\gamma k}(\lambda)) d\lambda = 0 \quad \alpha \neq \beta \neq \gamma$$

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$$\int \min(\mu_{\alpha i}(\lambda), \mu_{\beta j}(\lambda), \mu_{\gamma k}(\lambda)) d\lambda = 0 \quad \alpha \neq \beta \neq \gamma$$

From orthogonality:  $\int \min(\mu_{\alpha i}(\lambda), \mu_{\alpha j}(\lambda)) d\lambda = 0 \quad i \neq j$

# How $\Psi$ -epistemic?

A general result: the inclusion-exclusion principle....

$$v\left(\bigcup_i A_i\right) = \sum_i v(A_i) - \sum_{i < j} v(A_i \cap A_j) + \sum_{i < j < k} v(A_i \cap A_j \cap A_k) \dots$$



# How $\Psi$ -epistemic?

A general result: the inclusion-exclusion principle....

$$v(\cup_i A_i) = \sum_i v(A_i) - \sum_{i < j} v(A_i \cap A_j) + \sum_{i < j < k} v(A_i \cap A_j \cap A_k) \dots$$

and the Bonferroni inequality:

$$v(\cup_i A_i) \geq \sum_i v(A_i) - \sum_{i < j} v(A_i \cap A_j)$$

# How $\Psi$ -epistemic?

A general result: the inclusion-exclusion principle....

$$v(\cup_i A_i) = \sum_i v(A_i) - \sum_{i < j} v(A_i \cap A_j) + \sum_{i < j < k} v(A_i \cap A_j \cap A_k) \dots$$

and the Bonferroni inequality:

$$v(\cup_i A_i) \geq \sum_i v(A_i) - \sum_{i < j} v(A_i \cap A_j)$$

Substitute:

$$A_i = \Phi_{\alpha i} \cap \Phi_{\beta j} \quad \begin{aligned} v(A \cap B) &= \int \min(\mu_A(\lambda), \mu_B(\lambda)) d\lambda \\ v(A \cup B) &= \int \max(\mu_A(\lambda), \mu_B(\lambda)) d\lambda \end{aligned}$$

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## How $\Psi$ -epistemic can we get?

In a  $d$ -dimensional prime power Hilbert space, let  $|\beta_j\rangle$  be the  $j$ 'th basis state from the  $\beta$ 'th MUB.

Choose any state  $|\alpha_j\rangle$  from the  $\alpha$ 'th basis.

$$\sum_{\beta \neq \alpha, j} \omega(\phi_{\alpha_i}, \phi_{\beta_j}) \leq \frac{1}{1 - \sqrt{1 - 1/d}} \rightarrow 2d$$

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$$\overline{\omega}_d = \frac{\sum_{\beta \neq \alpha, j} \overline{\omega}(\phi_{\alpha i}, \phi_{\beta j})}{d^2} \leq \frac{1}{d^2 - d\sqrt{d(d-1)}} \rightarrow \frac{2}{d}$$

For any Hilbert space of dimension  $\geq d$ , there must exist pairs of states for which  $\overline{\omega}(\phi, \psi) \leq \overline{\omega}_d$

# Limitations on $\omega(\psi, \phi)$

- As in the case of the quantum state overlap,  $\Psi$ -ontic theories account for the inability to perfectly discriminate quantum states entirely in terms of limitations of the measurement process.
- $\Psi$ -epistemic theories cannot account for quantum state overlap entirely in terms of overlaps in probability distributions.
  - In the limit of large Hilbert spaces, the indistinguishability must still be entirely due to limitations of the measurement process for some pairs of states.

# Noise Tolerance

Perimeter Institute  
May 2013



JOHN TEMPLETON  
FOUNDATION

Limitations on the  $\Psi$ -epistemic  
view of quantum states

# Finite precision loophole?

	$M$		
	$ q_1\rangle$	$ q_2\rangle$	$ q_3\rangle$
$ \Phi_a\rangle$	$\epsilon$	$p(q_2 \Phi_a)$	$p(q_3 \Phi_a)$
$ \Phi_b\rangle$	$p(q_1 \Phi_b)$	$\epsilon$	$p(q_3 \Phi_a)$
$ \Phi_c\rangle$	$p(q_1 \Phi_c)$	$p(q_2 \Phi_c)$	$\epsilon$

$$\Lambda_a \cap \Lambda_b \cap \Lambda_c \neq \emptyset$$

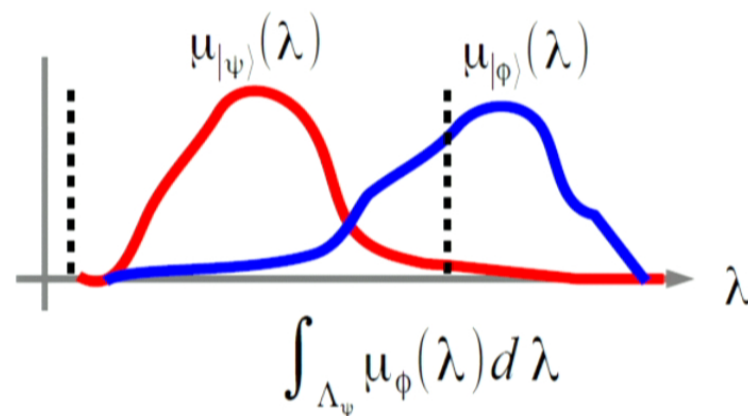


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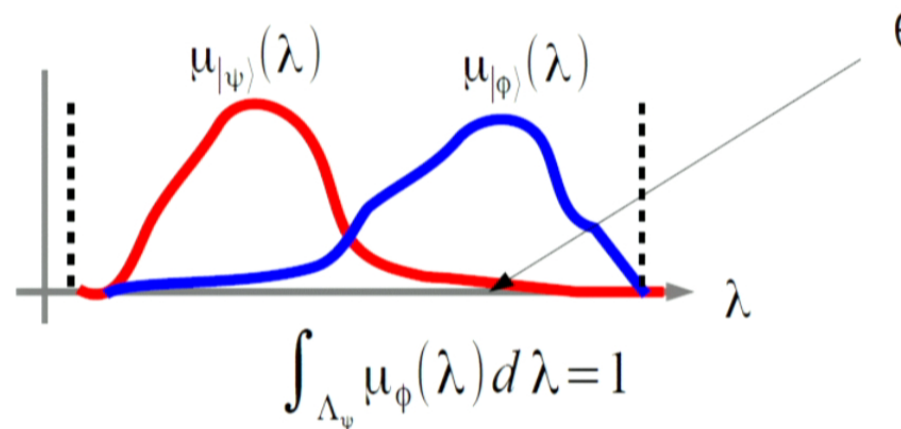
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# Quantum State Overlap



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For any Hilbert space of dimension  $\geq d$ , there must exist pairs of states for which  $\overline{\omega}(\phi, \psi) \leq \overline{\omega}_d$



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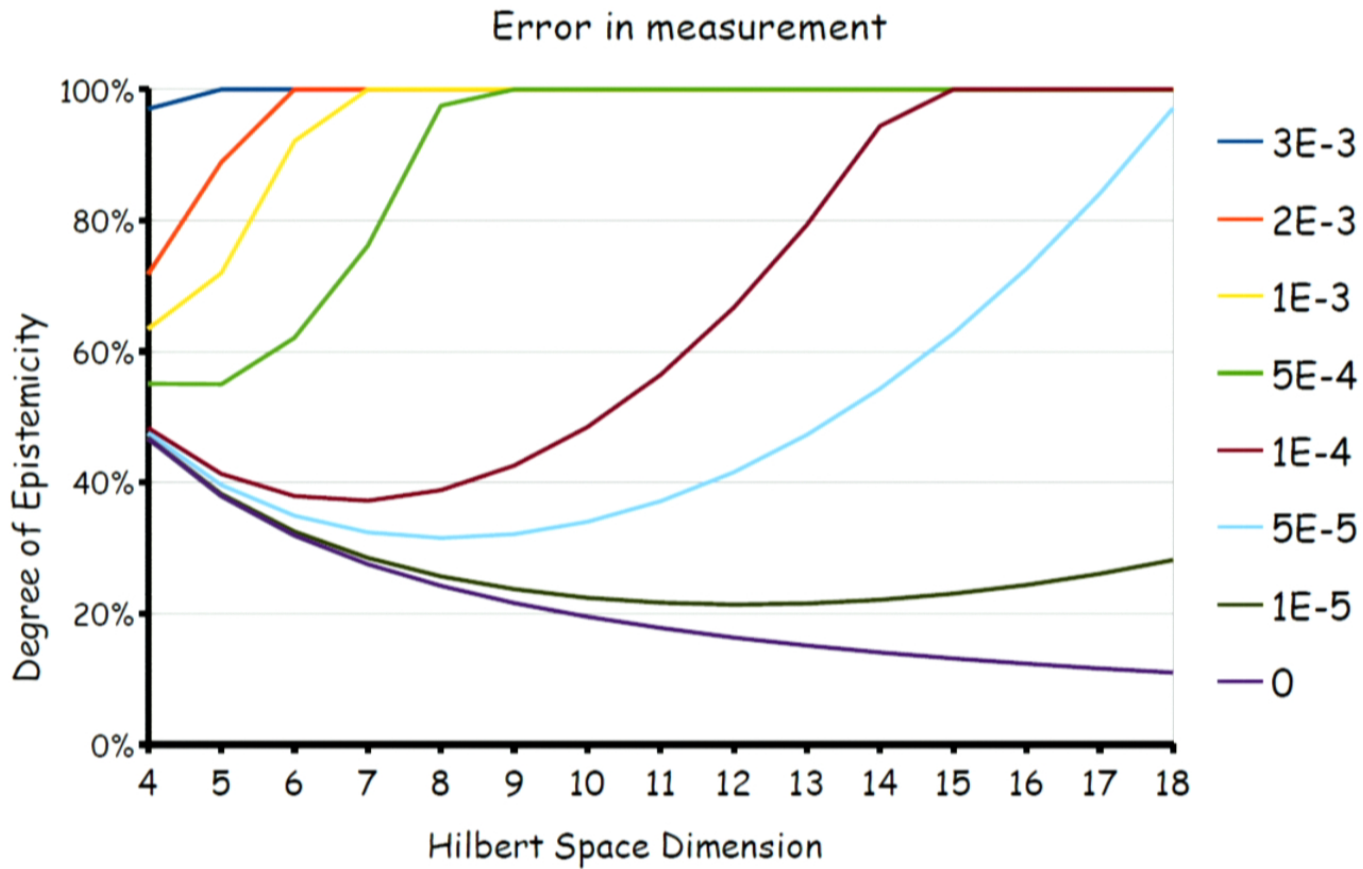
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**There is no finite precision loophole.**

The precision required is high (a minimum fidelity of  $\epsilon > 99.7\%$  for  $d=4$ ) but is not beyond the bounds of possibility with the best ion traps.

# Outlook

- Qualitatively quantum state overlap and quantum state discrimination are accounted for by:
  - Overlapping probability distributions in the epistemic view
  - Resolution of measurement process in the ontic view
- Quantitatively neither quantum state overlap nor quantum state discrimination can be accounted for entirely in terms of overlapping distributions.
  - The  $\psi$ -epistemic viewpoint must also rely on resolution of the measurement process.
- Insofar as there is a puzzle for the  $\psi$ -ontic point of view, there is a puzzle for the  $\psi$ -epistemic point of view.

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# Outlook

- In very high dimension Hilbert spaces,  $\omega, \bar{\omega} \rightarrow 0$ .
  - One of the principal responses to the Hardy and PBR theorems would be that we cannot isolate separate Hilbert spaces. So there is only One Big Hilbert space - of very high dimension!
- Are  $\omega(\psi, \phi)$  and  $\bar{\omega}(\psi, \phi)$  constant between all pairs of states? Unlikely.
  - The average value  $\langle \omega \rangle = \int \omega(\psi, \phi) d\psi d\phi$  may be higher.
  - One possibility that remains: maximally epistemic with respect to a given basis.
- The only constructive models of  $\psi$ -epistemic theories (LJBR and ABCL) seem to have values of  $\omega, \bar{\omega}$  lower than this.
  - Either a better model is possible, or a better limitation!