

Title: Pinning of Fermionic Occupation Numbers

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Abstract: The problem of determining and describing the family of 1-particle reduced density operators (1-RDO) arising from N-fermion pure states (via partial trace) is known as the fermionic quantum marginal problem. We present its solution, a multitude of constraints on the eigenvalues of the 1-RDO, generalizing the Pauli exclusion principle. To explore the relevance of these constraints we study an analytically solvable model of N fermions in a harmonic potential and determine the spectral 'trajectory' corresponding to the ground state as function of the fermion-fermion interaction strength. Intriguingly, we find that the occupation numbers are almost, but not exactly, pinned to the boundary of the allowed region (quasi-pinned). Our findings suggest a generalization of the Hartree-Fock approximation.
see also: <http://arxiv.org/abs/1210.5531>

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Credits: M.Christandl & D.Gross



Physical Review Letters

arXiv:1210.5531 [quant-ph]

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Outline

- 1) Motivation
- 2) Generalized Pauli Constraints
- 3) Physical Model
- 4) Pinning Analysis
- 5) Physical Relevance of Pinning



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'no two identical fermions in
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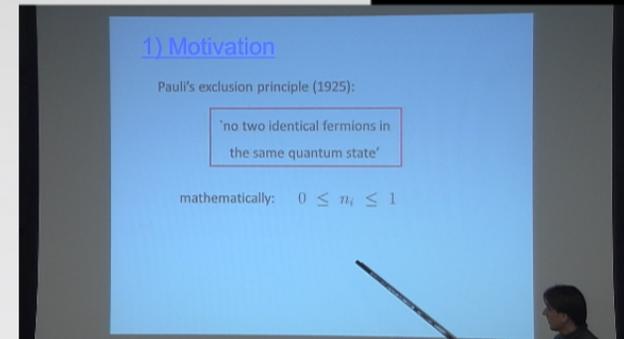


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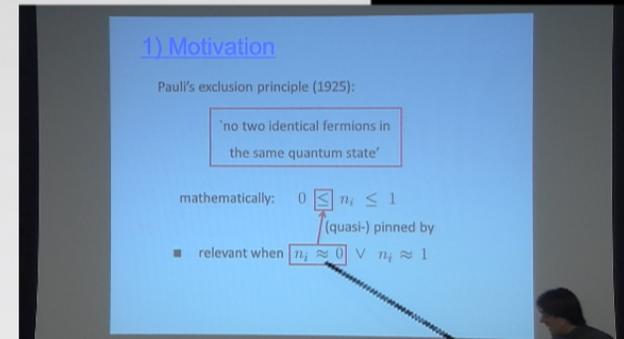
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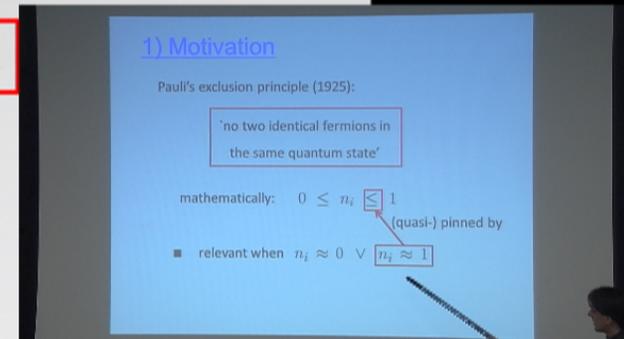
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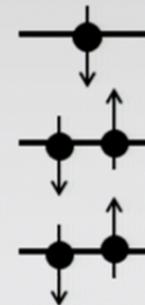
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- → Aufbau principle for atoms



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mathematical objects ?

$$\text{N-fermion states } |\Psi_N\rangle \rightarrow \rho_N \equiv |\Psi_N\rangle\langle\Psi_N|$$

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1-particle reduced
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translate antisymmetry of $|\Psi_N\rangle$
to 1-particle picture

2) Generalized Pauli Constraints

$$\{|\Psi_N\rangle\} \Rightarrow \{\rho_1\}$$

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A: ■ unitary equivalence:

$$\rho_1 \text{ arises from } |\Psi_N\rangle \quad \Rightarrow$$

$$U \rho_1 U^\dagger \text{ can arise from } U^{\otimes N} |\Psi_N\rangle$$

→ only natural occupation numbers $\vec{\lambda}$ relevant

■ depends on N & d

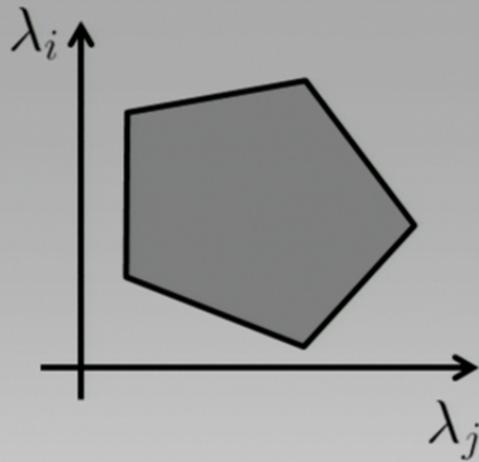
$$N \text{ fermions \& } \dim(\mathcal{H}_1) = d$$

■ $\{\vec{\lambda}\} = \mathcal{P}_{N,d} \subset \mathbb{R}^d$

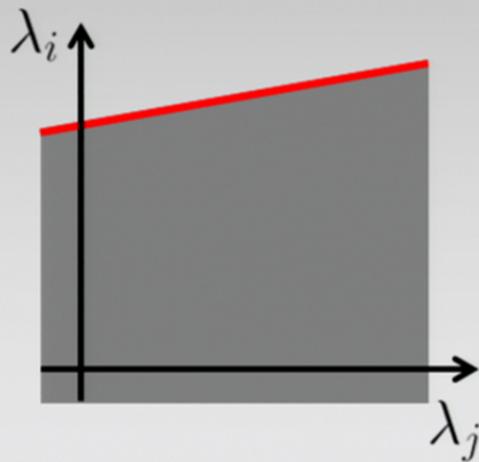
Polytope

[A.Klyachko., CMP 282, p287-322, 2008]

[A.Klyachko, J.Phys 36, p72-86, 2006]



polytope
 =
 intersection of
 finitely many half
 spaces



facet: $c_1 \lambda_1 + \dots + c_d \lambda_d = \kappa$

half
 space: $c_1 \lambda_1 + \dots + c_d \lambda_d \leq \kappa$

Example: $N = 3$ & $d = 7$

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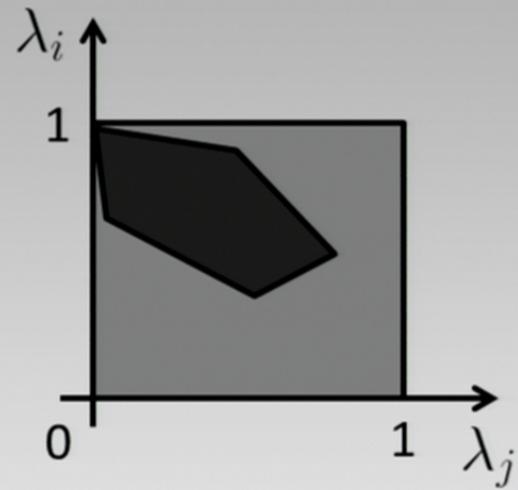
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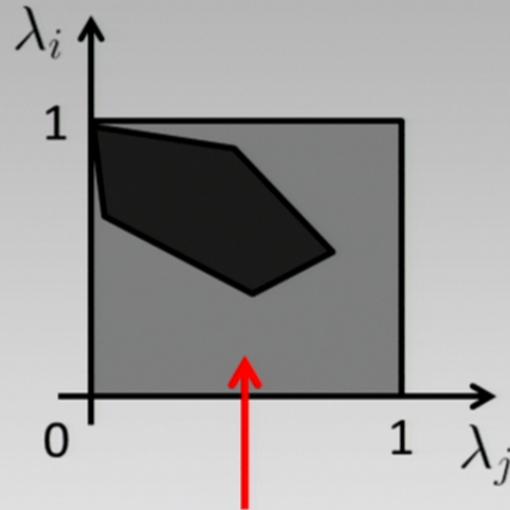
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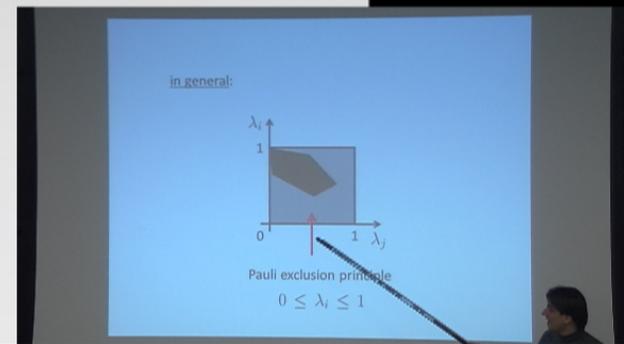


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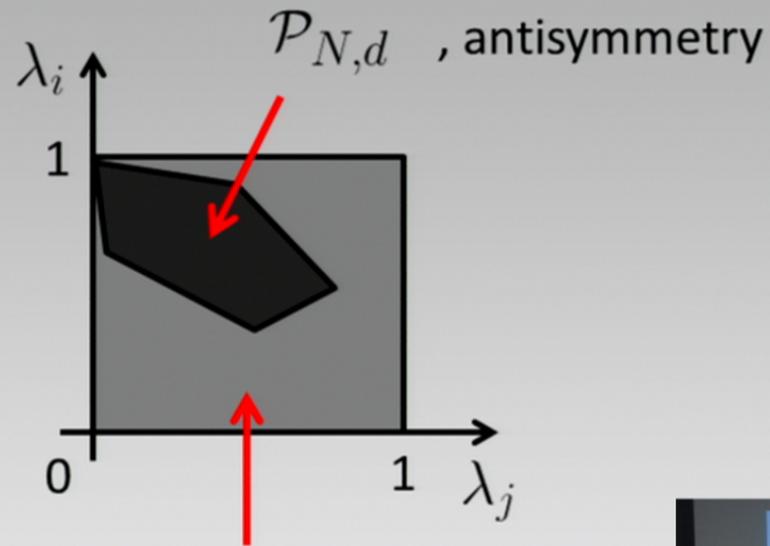


Pauli exclusion principle

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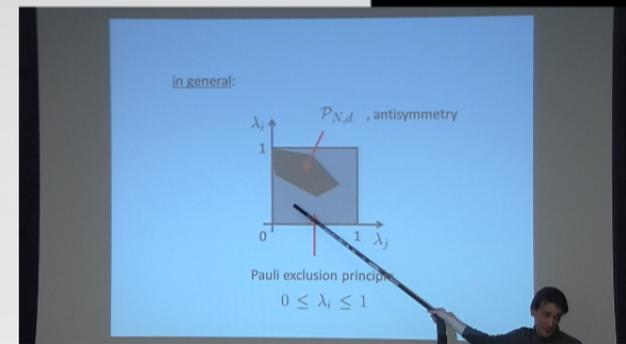


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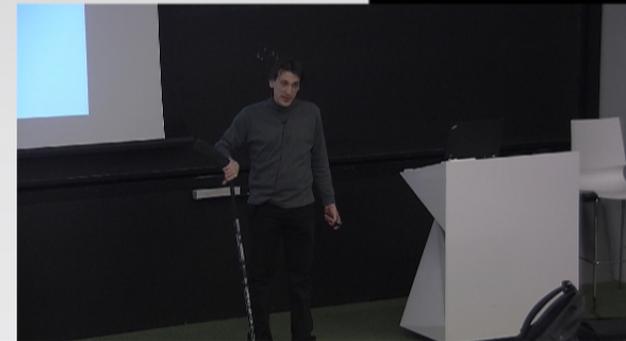
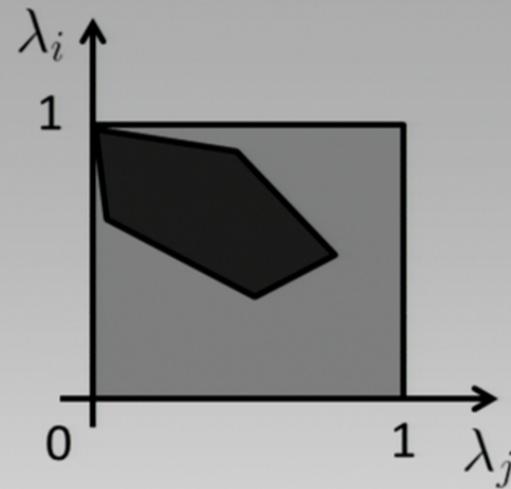


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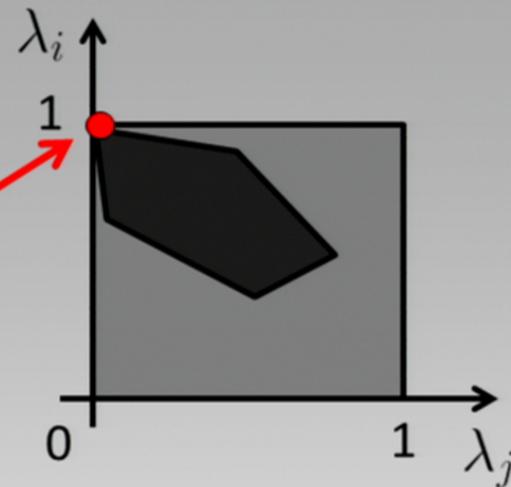
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Position of relevant states
(e.g. ground state) ?



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Slater determinants

$$|\Psi_N\rangle = |1, 2, \dots, N\rangle$$



$$\rho_1 = |1\rangle\langle 1| + \dots + |N\rangle\langle N|$$



$$\vec{\lambda} = \text{spec}(\rho_1) = (\underbrace{1, \dots, 1}_N, 0, 0, \dots)$$

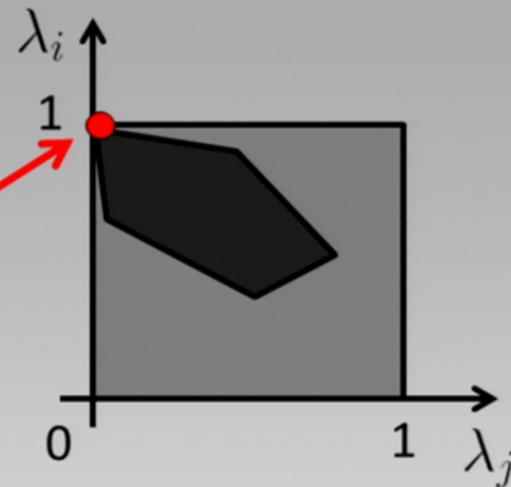
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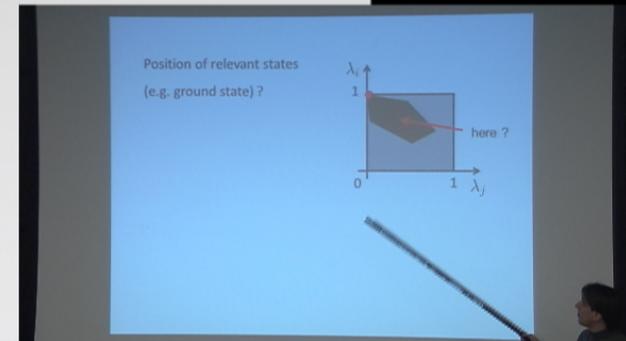
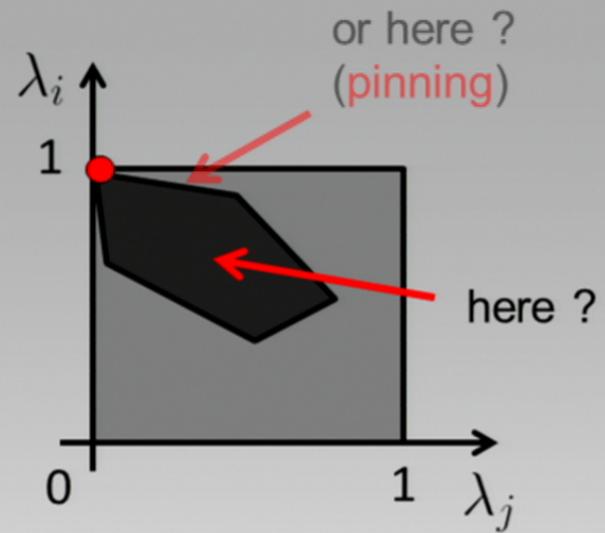


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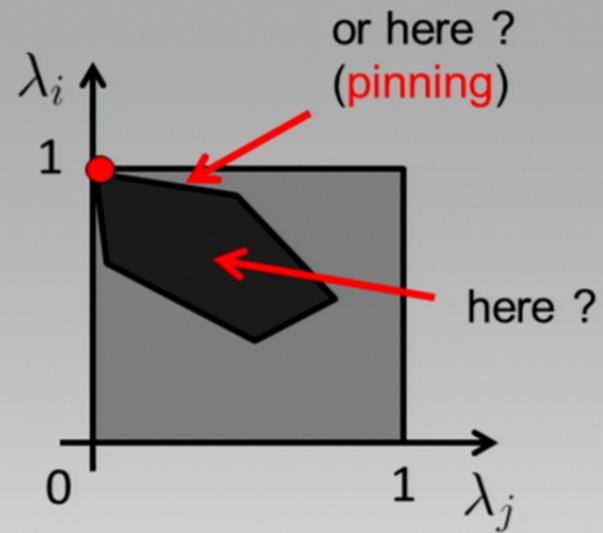
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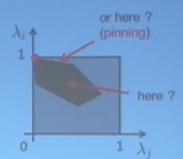
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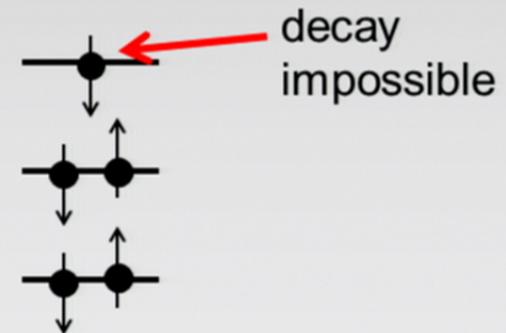
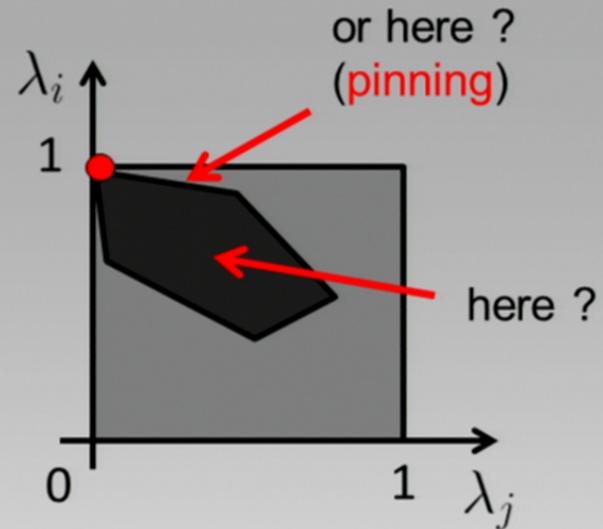
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kinematical constraints

generalization of:



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$$\rho_1 = N \text{Tr}_{N-1}[\rho_N]$$



$$\vec{\lambda} = \text{spec}(\rho_1)$$

Hamiltonian:

$$H = \sum_{i=1}^N \left(\frac{p_i^2}{2m} + \frac{1}{2} m \omega^2 x_i^2 \right) + \frac{1}{2} D \sum_{i,j=1}^N (x_i - x_j)^2$$

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→ restrict to $\wedge^N[L^2(\mathbb{R})] \subseteq L^2(\mathbb{R})^{\otimes N}$

ground state: [Z.Wang et al., arXiv 1108.1607, 2011]

$$\Psi_N(\vec{x}) = \mathcal{N} \left(\prod_{1 \leq i < j \leq N} (x_i - x_j) \right) \cdot \exp \left[-\frac{1}{2N} \left(\frac{1}{l^2} - \frac{1}{\tilde{l}^2} \right) (x_1 + \dots + x_N)^2 - \frac{1}{2\tilde{l}^2} \vec{x}^2 \right]$$

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from now on : $N = 3$

properties of $\vec{\lambda}$:

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- non-trivial duality

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$$\vec{\lambda}(\tilde{l}/l) = \vec{\lambda}(l/\tilde{l})$$

$$e^{-\delta} := \tilde{l}/l$$

$$\longrightarrow \vec{\lambda}(\delta) = \vec{\lambda}(-\delta)$$

$$\longrightarrow \text{weak-interacting} \hat{=} |\delta| \ll 1$$

- 'Boltzmann distribution law':

$$\lambda_k \sim e^{-\gamma k} \quad k \gg 1$$

- $1 - \lambda_1 = \frac{40}{729} \delta^6 + \dots$
 $1 - \lambda_2 = \frac{2}{9} \delta^4 + \dots$
 $1 - \lambda_3 = \frac{2}{9} \delta^4 + \dots$

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| $1 - \lambda_3 = \frac{2}{9}\delta^4 + \dots$ | $\lambda_7 = \frac{80}{2187}\delta^8 + \dots$ |
| $\lambda_4 = \frac{2}{9}\delta^4 + \dots$ | $\lambda_8 = O(\delta^{10})$ |

- hierarchy:

$$\lambda_k = c_k \delta^{2k-6} + O(\delta^{2k-4}), \forall k \geq 5$$

4) Pinning Analysis

$$\dim(L^2(\mathbb{R})) = \infty \quad \Rightarrow \quad \vec{\lambda} \in \mathcal{P}_{3,\infty}$$

4) Pinning Analysis

too difficult/
not known yet

$$\dim(L^2(\mathbb{R})) = \infty \quad \Rightarrow \quad \vec{\lambda} \in \mathcal{P}_{3,\infty}$$

instead: check $\vec{\lambda}_d := (\lambda_1, \dots, \lambda_d)$ w.r.t $\mathcal{P}_{3,d}$

(quasi-)pinning of $\mathcal{P}_{3,d}$ - constraint

4) Pinning Analysis

too difficult/
not known yet

$$\dim(L^2(\mathbb{R})) = \infty \quad \Rightarrow \quad \vec{\lambda} \in \mathcal{P}_{3,\infty}$$

instead: check $\vec{\lambda}_d := (\lambda_1, \dots, \lambda_d)$ w.r.t $\mathcal{P}_{3,d}$

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\Leftrightarrow

(quasi-)pinning of $\mathcal{P}_{3,\infty}$ - constraint $\quad \text{mod } O(\lambda_{d+1})$

■ $\wedge^3[\mathcal{H}_6]$

[Borland&Dennis, J.Phys. B, 5,1, 1972]

[Ruskai, Phys. Rev. A, 40,45, 2007]

$$\lambda_1 + \lambda_6 = \lambda_2 + \lambda_5 = \lambda_3 + \lambda_4 = 1$$

$$D^{(6)} := \lambda_5 + \lambda_6 - \lambda_4 \geq 0$$

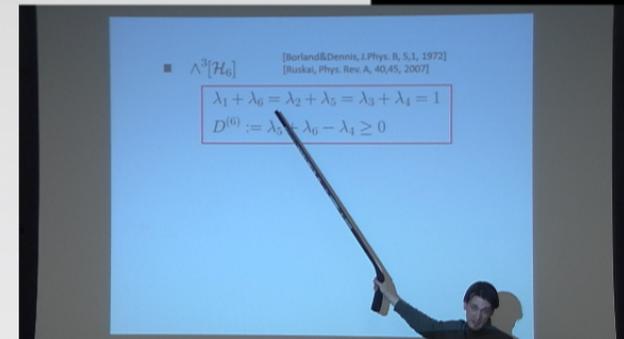
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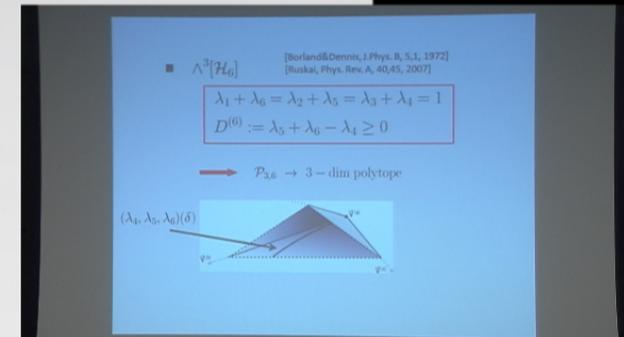
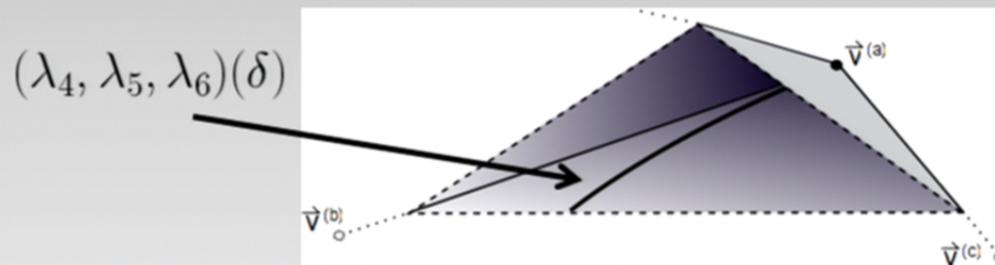
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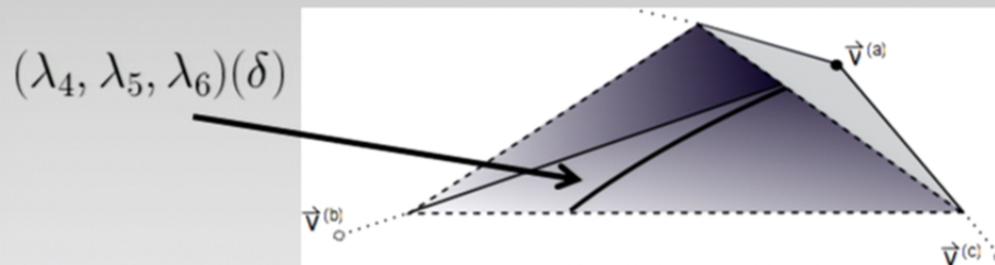
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$(\lambda_4, \lambda_5, \lambda_6)(\delta)$

$\vec{v}^{(a)}$

$\vec{v}^{(b)}$

$\vec{v}^{(c)}$

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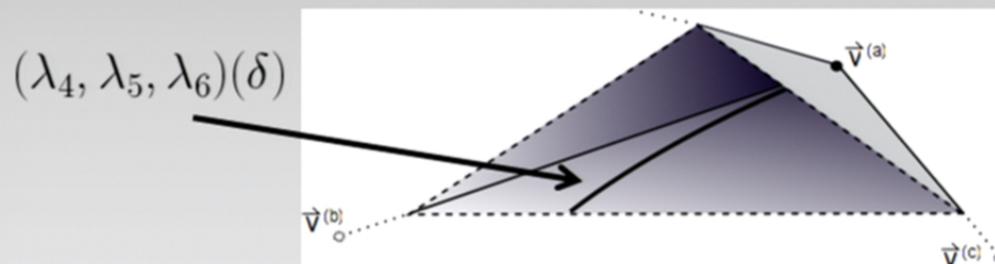
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$$0 \leq D^{(6)}(\delta) = \frac{4510}{59049} \delta^8 + O(\delta^{10})$$

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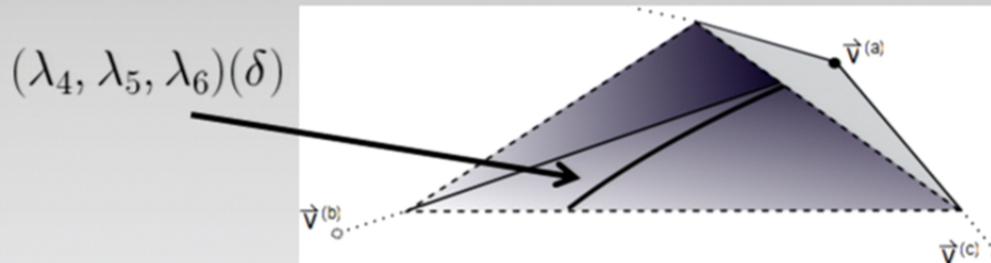
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relevant as long
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→ lower bound on
pinning order

$$\blacksquare \quad \wedge^3[\mathcal{H}_7]$$

$$D_1^{(7)} = \frac{20}{2187} \delta^8 + O(\delta^{10})$$

$$D_2^{(7)} = \frac{10}{243} \delta^8 + O(\delta^{10})$$

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quasi-pinning

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Physical Relevance of Pinning ?

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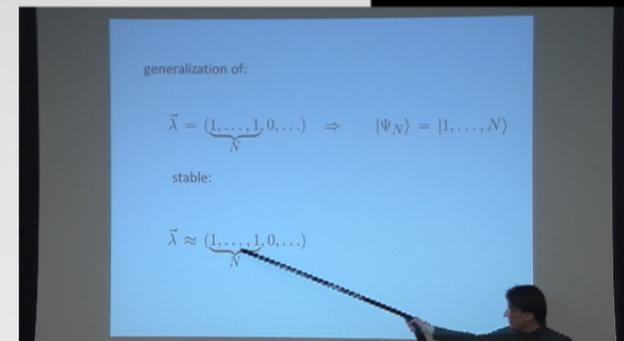


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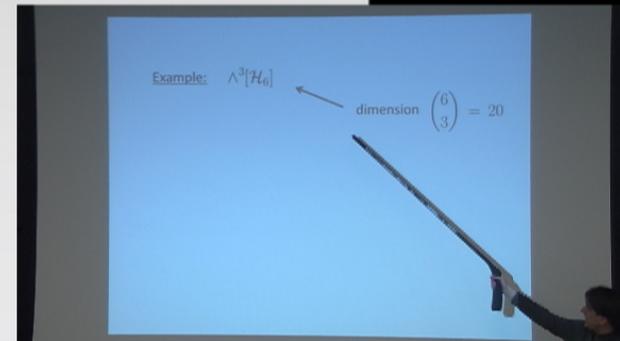
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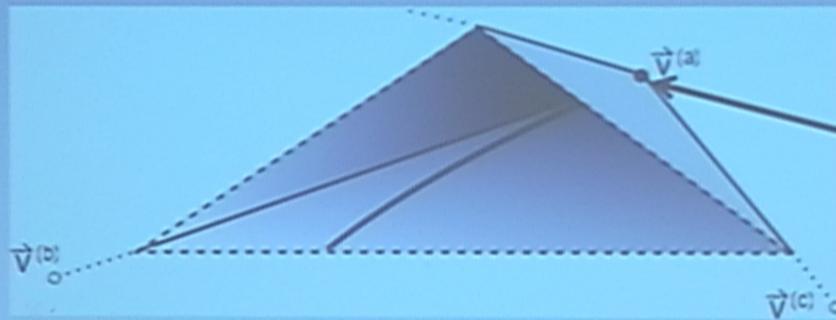
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approximate unknown ground state $|\Psi_3\rangle$



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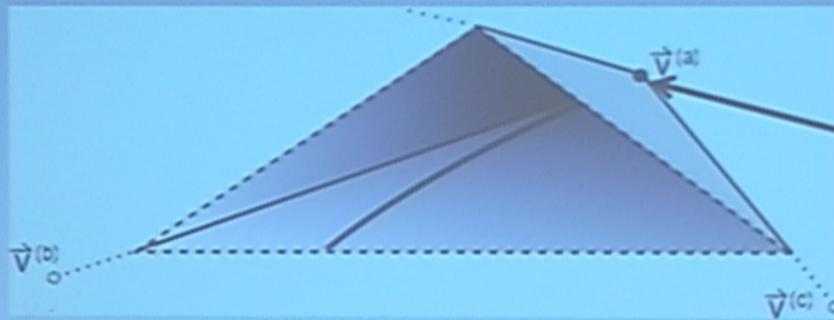
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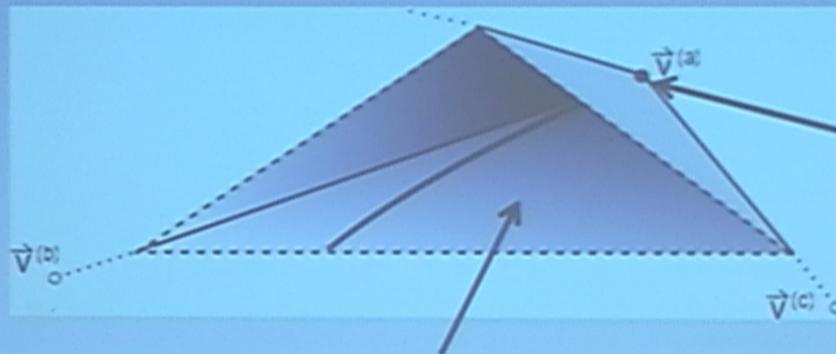
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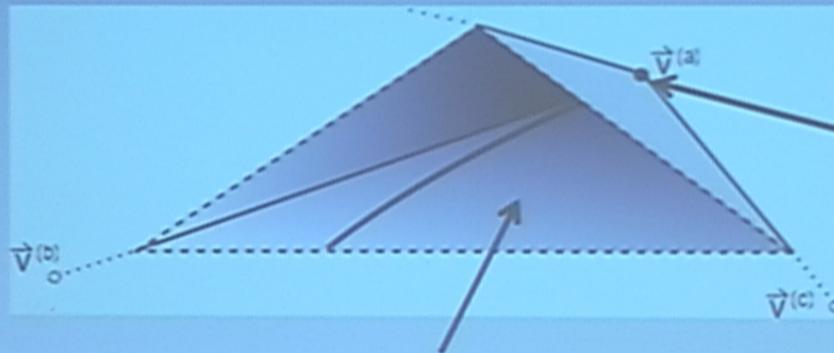
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Conclusions

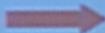
- antisymmetry of $|\Psi_N\rangle$
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Generalized
Pauli constraints



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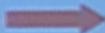


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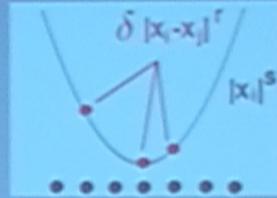


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- Fermionic Ground States simpler than appreciated (?)

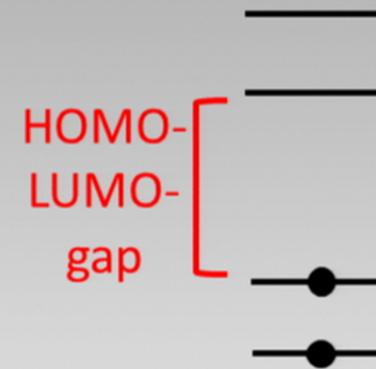
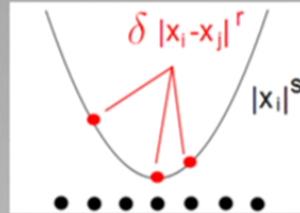
Outline

- generic for:
- Hubbard model
- Quantum Chemistry: Atoms

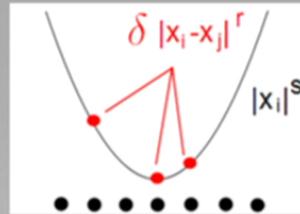


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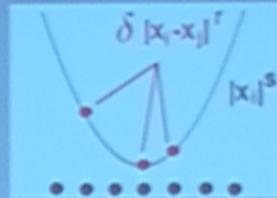


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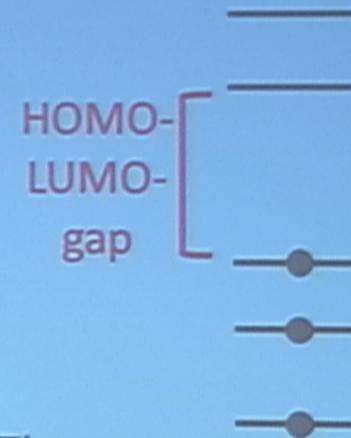
HOMO-
LUMO-
gap



Outline



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- Antisymmetry ⚡ Energy Minimization



- 'Boltzmann distribution law':

$$\lambda_k \sim e^{-\gamma k} \quad k \gg 1$$

- | | |
|--|---|
| $1 - \lambda_1 = \frac{40}{729}\delta^5 + \dots$ | $\lambda_5 = \frac{2}{9}\delta^4 + \dots$ |
| $1 - \lambda_2 = \frac{2}{9}\delta^4 + \dots$ | $\lambda_6 = \frac{40}{729}\delta^5 + \dots$ |
| $1 - \lambda_3 = \frac{2}{9}\delta^4 + \dots$ | $\lambda_7 = \frac{80}{2187}\delta^8 + \dots$ |
| $\lambda_4 = \frac{2}{9}\delta^4 + \dots$ | $\lambda_8 = O(\delta^{10})$ |

hierarchy:

$$= c_k \delta^{2k-6} + O(\delta^{2k-4}), \forall k \geq 5$$

$$\Delta^N [\mathbb{C}^d] \cong \Lambda^{d-N} [\mathbb{C}^d]$$

$$a_i + a_j$$

$$N=2$$

$$d$$

$$\vec{\lambda}$$

$$\lambda_1 = \lambda_2$$

$$\lambda_3 = \lambda_4$$

$$N=3$$