

Title: Boosted Multijet Resonances and New Color-Flow Variables

Date: Jan 15, 2013 01:00 PM

URL: <http://pirsa.org/13010104>

Abstract:

Boosted Multijet Resonances and New Color-Flow Variables.

D. Curtin, R. Essig, BS, arXiv:1210.5523

Brian Shuve

PI Particle Physics Seminar

January 15, 2013

Boosted Multijet Resonances and New Color-Flow Variables.

D. Curtin, R. Essig, BS, arXiv:1210.5523

Brian Shuve

PI Particle Physics Seminar

January 15, 2013

Motivation

- Expect first evidence of new physics to appear in strong sector
- Channels with lowest backgrounds:
 - ▶ Jets + missing energy
 - ▶ Jets + leptons
 - ▶ Jets + photons

I

Motivation

- Expect first evidence of new physics to appear in strong sector
- Channels with lowest backgrounds:
 - ▶ Jets + missing energy
 - ▶ Jets + leptons
 - ▶ Jets + photons
- In many BSM theories (ex. conventional SUSY), we can get such signatures because:
 - ▶ Stable neutral particle gives missing energy
 - ▶ Tops in cascade decays can give leptons and missing energy

Motivation

- Expect first evidence of new physics to appear in strong sector
- Channels with lowest backgrounds:
 - ▶ Jets + missing energy
 - ▶ Jets + leptons
 - ▶ Jets + photons
- In many BSM theories (ex. conventional SUSY), we can get such signatures because:
 - ▶ Stable neutral particle gives missing energy
 - ▶ Tops in cascade decays can give leptons and missing energy
- As a result, bounds on such theories are creeping ever upwards:
 - ▶ Gluino bounds in 1-1.5 TeV range
 - ▶ Squark bounds in 1 TeV range
 - ▶ Stop/sbottom bounds in 500 GeV range (except for top mass window)

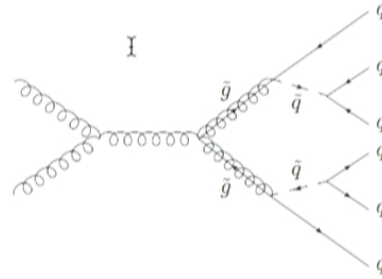
Motivation

- Bounds are considerably weaker in all-hadronic final states because of large QCD backgrounds

I

Motivation

- Bounds are considerably weaker in all-hadronic final states because of large QCD backgrounds
- As a benchmark, we consider a simplified model with gluinos of mass $m_{\tilde{g}}$ decaying to three light-flavor quarks through an R -parity-violating coupling



Boosted final states

- If gluinos are boosted, then final state jets are collimated
 - ▶ Use jet substructure techniques to separate signal from QCD background

I

Boosted final states

- If gluinos are boosted, then final state jets are collimated
 - ▶ Use jet substructure techniques to separate signal from QCD background
- Boosted object tagging is conventionally used when the multijet resonances are produced in a decay chain at some very high momentum
 - ▶ For example, Z' production of highly boosted hadronic tops

I

Boosted final states

- If gluinos are boosted, then final state jets are collimated
 - ▶ Use jet substructure techniques to separate signal from QCD background
- Boosted object tagging is conventionally used when the multijet resonances are produced in a decay chain at some very high momentum
 - ▶ For example, Z' production of highly boosted hadronic tops
- We argue that, even for hadronic resonances produced predominantly at rest, it can be possible to tag and study the boosted fraction
 - ▶ Small number of events, but aggressive substructure cuts can yield high-purity signal samples
 - ▶ Look for resonance in jet mass through shape analysis
 - ▶ Can be competitive with and is complementary to resolved analyses

Boosted final states

- If gluinos are boosted, then final state jets are collimated
 - ▶ Use jet substructure techniques to separate signal from QCD background
- Boosted object tagging is conventionally used when the multijet resonances are produced in a decay chain at some very high momentum
 - ▶ For example, Z' production of highly boosted hadronic tops
- We argue that, even for hadronic resonances produced predominantly at rest, it can be possible to tag and study the boosted fraction
 - ▶ Small number of events, but aggressive substructure cuts can yield high-purity signal samples
 - ▶ Look for resonance in jet mass through shape analysis
 - ▶ Can be competitive with and is complementary to resolved analyses
- We also introduce color-flow observables that can further enhance S/B
 - ▶ Excellent signal/background discrimination possible in highly boosted sample

Outline

- Gluinos and R -parity violation (RPV)
- Existing searches for RPV gluinos
- Boosted gluino signatures
 - ▶ Color-flow observables
- Monte Carlo study
 - ▶ Heavy gluino search
 - ▶ Top-mass gluino search

Gluginos and RPV

- R -parity is a symmetry that removes baryon- and lepton-number-violating operators from SUSY theories
 - ▶ Motivated by stringent bounds on proton decay, neutrino masses, etc.
 - ▶ Generally leads to jets + MET in colliders, which is strongly constrained

I

Gluginos and RPV

- R -parity is a symmetry that removes baryon- and lepton-number-violating operators from SUSY theories
 - ▶ Motivated by stringent bounds on proton decay, neutrino masses, etc.
 - ▶ Generally leads to jets + MET in colliders, which is strongly constrained
- As bounds on R -parity-conserving SUSY get stronger, attention is turning to RPV, which is much less constrained at colliders
- Constraints on proton decay, neutrino masses loosen considerably if we preserve lepton number, only violate baryon number (B -violating RPV):

$$W_{\mathcal{B}} = \frac{1}{2} \lambda''_{ijk} \epsilon^{\alpha\beta\gamma} \bar{u}_\alpha^i \bar{d}_\beta^j \bar{d}_\gamma^k$$

- ▶ Squarks can decay through this operator: $\tilde{q} \rightarrow qq$
- ▶ If $m_{\tilde{g}} < m_{\tilde{q}}$, then $\tilde{g} \rightarrow qqq$ through an off-shell squark

Gluginos and RPV

- R -parity is a symmetry that removes baryon- and lepton-number-violating operators from SUSY theories
 - ▶ Motivated by stringent bounds on proton decay, neutrino masses, etc.
 - ▶ Generally leads to jets + MET in colliders, which is strongly constrained
- As bounds on R -parity-conserving SUSY get stronger, attention is turning to RPV, which is much less constrained at colliders
- Constraints on proton decay, neutrino masses loosen considerably if we preserve lepton number, only violate baryon number (B -violating RPV):

$$W_{\mathcal{B}} = \frac{1}{2} \lambda''_{ijk} \epsilon^{\alpha\beta\gamma} \bar{u}_\alpha^i \bar{d}_\beta^j \bar{d}_\gamma^k$$

- ▶ Squarks can decay through this operator: $\tilde{q} \rightarrow qq$
- ▶ If $m_{\tilde{g}} < m_{\tilde{q}}$, then $\tilde{g} \rightarrow qqq$ through an off-shell squark

Gluginos and RPV

Implications for gluino decays:

- Gluino width (if $m_{\tilde{g}} \ll m_{\tilde{q}}$):

$$\Gamma(\tilde{g} \rightarrow qqq) \sim \frac{\alpha_s |\lambda''|^2 m_{\tilde{g}}^5}{384 \pi^2 m_{\tilde{q}}^4}$$

Gluginos and RPV

Implications for gluino decays:

- Gluino width (if $m_{\tilde{g}} \ll m_{\tilde{q}}$):

$$\Gamma(\tilde{g} \rightarrow qqq) \sim \frac{\alpha_s |\lambda''|^2 m_{\tilde{g}}^5}{384\pi^2 m_{\tilde{q}}^4}$$

- $\Gamma_{\tilde{g}} < \Lambda_{\text{QCD}}$ whenever $m_{\tilde{q}} > m_{\tilde{g}}$, even if $\lambda'' \sim 1$
 - ▶ Gluino hadronizes before decay

Gluginos and RPV

Implications for gluino decays:

- Gluino width (if $m_{\tilde{g}} \ll m_{\tilde{q}}$):

$$\Gamma(\tilde{g} \rightarrow qqq) \sim \frac{\alpha_s |\lambda''|^2 m_{\tilde{g}}^5}{384\pi^2 m_{\tilde{q}}^4}$$

- $\Gamma_{\tilde{g}} < \Lambda_{\text{QCD}}$ whenever $m_{\tilde{q}} > m_{\tilde{g}}$, even if $\lambda'' \sim 1$
 - ▶ Gluino hadronizes before decay
- No displaced vertex if $\lambda'' \gtrsim 10^{-3}$ (depends on squark mass)

Existing searches for RPV gluinos

Resonance searches:

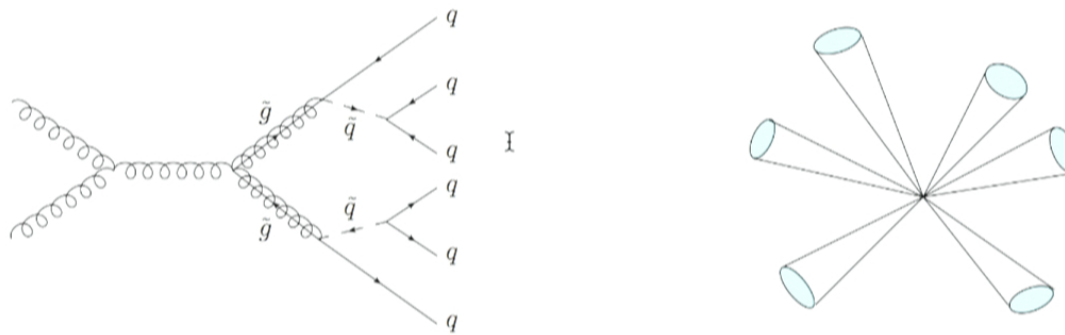
- We can reconstruct the gluino mass in **triplets** of jets

I

Existing searches for RPV gluinos

Resonance searches:

- We can reconstruct the gluino mass in **triplets** of jets



- But which jets do we combine? There are $\binom{6}{3} = 20$ combinations!
 - ▶ Best to consider ensemble of all combinations, although this results in large combinatoric background

Existing searches for RPV gluinos

Resonance searches:

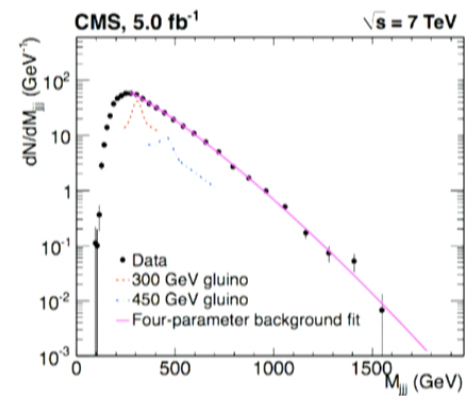
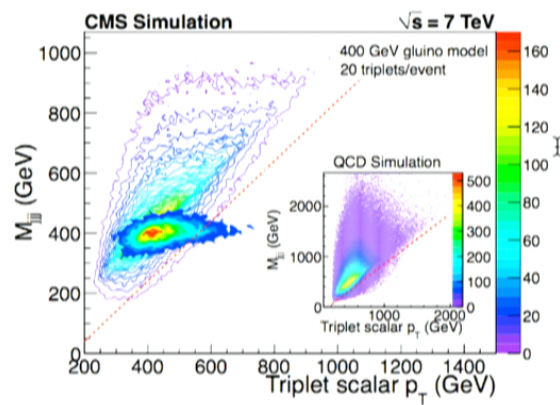
- Backgrounds are reduced using the fact that
 - ▶ $m_{jjj} \propto \sum p_{T i}$ for QCD/combinatoric background
 - ▶ $m_{jjj} \propto m_{\text{res}}$ for signal

I

Existing searches for RPV gluinos

Resonance searches:

- Backgrounds are reduced using the fact that
 - ▶ $m_{jjj} \propto \sum p_{T i}$ for QCD/combinatoric background
 - ▶ $m_{jjj} \propto m_{\text{res}}$ for signal



(CMS analysis, arXiv:1208.2931)

Existing searches for RPV gluinos

Counting experiments:

- ATLAS has a recent search that looks for a high multiplicity of hard jets (arXiv:1210.4813)
 - ▶ Require 6 anti- k_T , $R = 0.4$ jets with $p_T > 80 - 160$ GeV, depending on gluino mass
 - ▶ Exclude RPV gluinos up to 666 GeV
 - ▶ No sensitivity to underlying mass scale of new physics, except indirectly through p_T cut

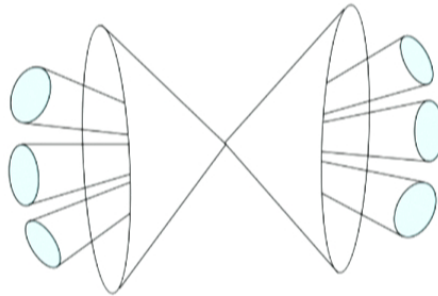
Existing searches for RPV gluinos

Counting experiments:

- ATLAS has a recent search that looks for a high multiplicity of hard jets (arXiv:1210.4813)
 - ▶ Require 6 anti- k_T , $R = 0.4$ jets with $p_T > 80 - 160$ GeV, depending on gluino mass
 - ▶ Exclude RPV gluinos up to 666 GeV
 - ▶ No sensitivity to underlying mass scale of new physics, except indirectly through p_T cut

Boosted searches

- For six-jet resonances, combinatoric background is a major problem
- If heavy particle is **boosted**, decay products are collimated
 - ▶ Removes ambiguity of choosing jet triplets
 - ▶ All resonance decay products can be clustered into one fat jet that reconstructs its mass



Existing searches for RPV gluinos

Boosted search:

- First experimental search for exotic boosted colored resonances released in October!
 - ▶ Counting experiment (arXiv:1210.4813)
 - ▶ Look for two fat jets ($R = 1$), ≥ 4 thin jets ($R = 0.4$), and place cuts on jet mass and a variable called N -subjettiness
 - ▶ Exclude gluino masses up to 255 GeV, closing the allowed mass window with $m_{\tilde{g}} \sim m_t$ left by the CDF and CMS searches

Existing searches for RPV gluinos

Theoretical study:

- Boosted gluino search first considered by Raklev, Salam, Wacker (arXiv:1005.1229)
 - ▶ Analysis done for the Tevatron
 - ▶ Require two hard ($p_T > 350$ GeV) jets that are comparable in mass and exhibit some substructure (cut on merging scale y_1)
 - ▶ Find that the Tevatron exclusion reach could be as high as 280 GeV for RPV gluinos (much better than 145 GeV bound from six-jet analysis!)

RPV gluino signatures

- There are two properties of gluino decay that we exploit:
 - 1 Kinematics of the hard final-state partons
 - ▶ Two “fat” jets reconstructing the same mass
 - ▶ Each fat jet has three hard centres of radiation corresponding to the three-quark final states of gluino decay

RPV gluino signatures

- We use several substructure variables that are sensitive to these properties

1 Kinematics of the hard final-state partons

- ▶ Two “fat” jets reconstructing the same mass

- ★ **Jet mass symmetry**

$$s_{\text{m}} \equiv \frac{|m_1 - m_2|}{(m_1 + m_2)/2}$$

- ▶ Each fat jet has three hard centres of radiation corresponding to the three-quark final states of gluino decay

- ★ **N -subjettiness** (defined momentarily)

- ▶ The p_{T} of the three hardest subjects should be comparable

- ★ **Subject hierarchy**

$$h_{31} \equiv \frac{p_{\text{T}3}}{p_{\text{T}1}}$$

Variable: N -subjettiness

- N -subjettiness (τ_N) is a variable that characterizes whether a jet looks like it has N subjects
 - ▶ $\tau_N \ll 1$ if a jet is well-described by N subjects
 - ▶ $\tau_N \sim 1$ if a jet is **not** well-described by N subjects (i.e. there is a lot of radiation that is not aligned with any of the subjects)
- N -subjettiness: given a set of N axes,

$$\tau_N^{(\beta)} = \frac{1}{d_0} \sum_{i \in \text{jet}} p_{T i} \min[(\Delta R_{i1})^\beta, \dots, (\Delta R_{iN})^\beta]$$
$$d_0 = \sum_{i \in \text{jet}} p_{T i} R_0^\beta$$

- ▶ R_0 is the radius used to cluster the fat jet
- ▶ ΔR_{ij} is the distance in the η - ϕ plane between particle i and axis j

Variable: N -subjettiness

- N -subjettiness (τ_N) is a variable that characterizes whether a jet looks like it has N subjects
 - ▶ $\tau_N \ll 1$ if a jet is well-described by N subjects
 - ▶ $\tau_N \sim 1$ if a jet is **not** well-described by N subjects (i.e. there is a lot of radiation that is not aligned with any of the subjects)
- N -subjettiness: given a set of N axes,

$$\tau_N^{(\beta)} = \frac{1}{d_0} \sum_{i \in \text{jet}} p_{T i} \min[(\Delta R_{i1})^\beta, \dots, (\Delta R_{iN})^\beta]$$

$$d_0 = \sum_{i \in \text{jet}} p_{T i} R_0^\beta$$

- ▶ R_0 is the radius used to cluster the fat jet
- ▶ ΔR_{ij} is the distance in the η - ϕ plane between particle i and axis j

Variable: N -subjettiness

I

$$\tau_N^{(\beta)} = \frac{1}{d_0} \sum_{i \in \text{jet}} p_{T i} \min[(\Delta R_{i1})^\beta, \dots, (\Delta R_{iN})^\beta]$$

- How do we figure out which axes to use?
 - ▶ Because axes \sim subjects, we can use jet clustering algorithm to re-cluster jet constituents into N subjects (ATLAS uses k_T)

Variable: N -subjettiness

I

$$\tau_N^{(\beta)} = \frac{1}{d_0} \sum_{i \in \text{jet}} p_{T i} \min[(\Delta R_{i1})^\beta, \dots, (\Delta R_{iN})^\beta]$$

- How do we figure out which axes to use?
 - ▶ Because axes \sim subjects, we can use jet clustering algorithm to re-cluster jet constituents into N subjects (ATLAS uses k_T)
 - ▶ Alternatively, there is a **unique** choice of axes that minimizes τ_N and represents the “best” choice of axes
- We use the τ_N -minimization technique outlined in arXiv:1108.2701 (Thaler, van Tilburg)
 - ▶ Get substantial improvement over k_T -clustering (need to get experimentalists to switch too!)

Variable: N -subjettiness

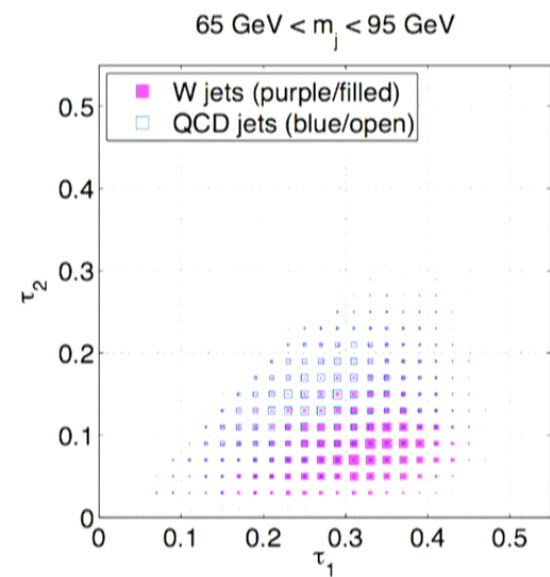
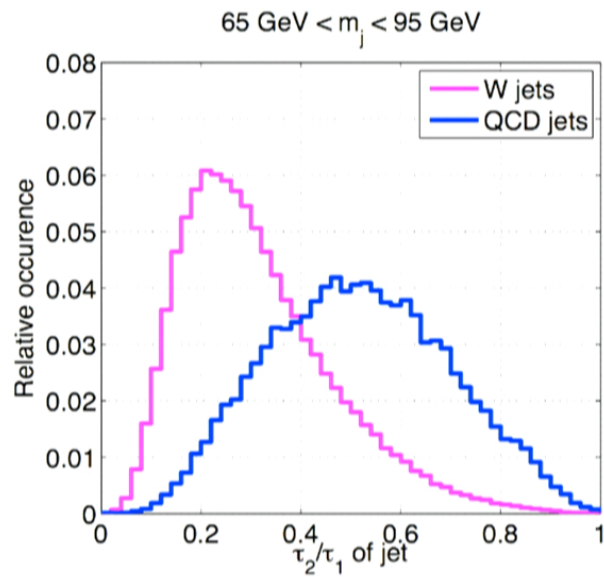
I

- It turns out that it's better to look at ratios of N -subjettiness
($\tau_{NN-1} \equiv \tau_N / \tau_{N-1}$)
 - ▶ A sudden drop in τ_N indicates that N -subjettiness has locked onto the correct number of jets



- ▶ Crosses: 1-subjettiness axes, circles: 2-subjettiness axes
- ▶ Left: schematic of signal event (hadronic W decay), right: schematic of QCD fat jet

Variable: N -subjettiness



(Thaler, Van Tilburg, arXiv:1011.2268)

- We use τ_{32} as a discriminant to isolate boosted gluino events

Color flow in gluino decays

I

- Quarks arising from color-singlet particle decay form a color dipole

Color flow in gluino decays

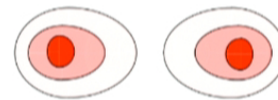
I

- Quarks arising from color-singlet particle decay form a color dipole
 - ▶ Radiation tends to be concentrated between them

Color flow in gluino decays

I

- Quarks arising from color-singlet particle decay form a color dipole
 - ▶ Radiation tends to be concentrated between them
 - ▶ 'Color string' picture



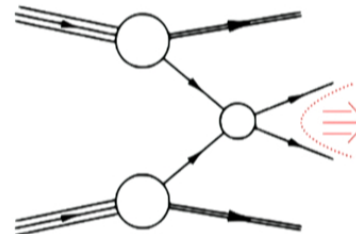
Color flow in gluino decays

I

- Quarks arising from color-singlet particle decay form a color dipole
 - ▶ Radiation tends to be concentrated between them
 - ▶ 'Color string' picture



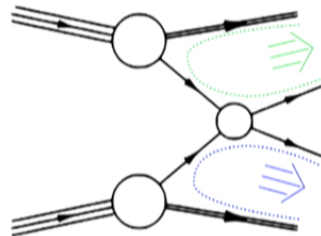
- This is the color configuration of quarks from hadronic W or Higgs decays



Color flow in gluino decays

I

- When quarks contained with protons collide, there remain colored **beam remnants**
 - ▶ Final state quarks can also be color-connected with the beam remnant



Color flow in gluino decays

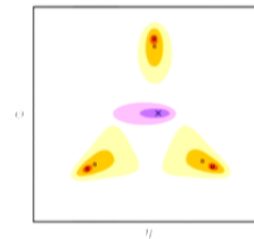
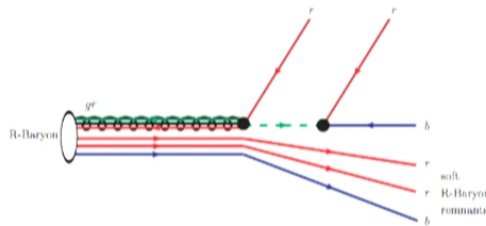
I

- For RPV-decaying gluino LSPs, the decay happens after hadronization
 - ▶ Gluino decay results in 3 hard quarks and a hadron remnant (analogous to the beam remnant)

Color flow in gluino decays

I

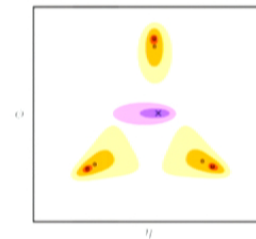
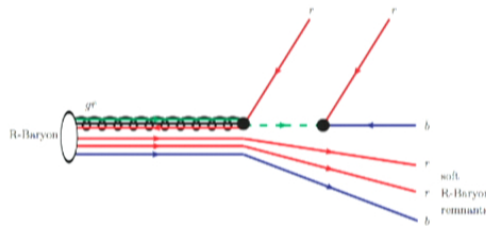
- For RPV-decaying gluino LSPs, the decay happens after hadronization
 - ▶ Gluino decay results in 3 hard quarks and a hadron remnant (analogous to the beam remnant)
 - ▶ As before, this leads to color connection between the gluino decay products and the hadron remnant (as well as between the gluino decay products)



Color flow in gluino decays

I

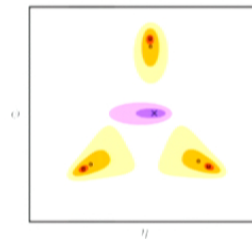
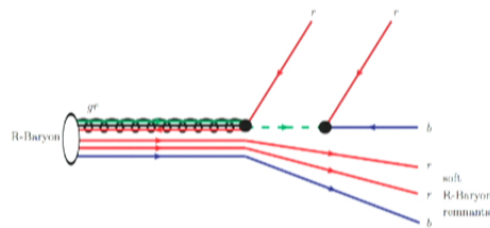
- For RPV-decaying gluino LSPs, the decay happens after hadronization
 - ▶ Gluino decay results in 3 hard quarks and a hadron remnant (analogous to the beam remnant)
 - ▶ As before, this leads to color connection between the gluino decay products and the hadron remnant (as well as between the gluino decay products)



Color flow in gluino decays

I

- For RPV-decaying gluino LSPs, the decay happens after hadronization
 - ▶ Gluino decay results in 3 hard quarks and a hadron remnant (analogous to the beam remnant)
 - ▶ As before, this leads to color connection between the gluino decay products and the hadron remnant (as well as between the gluino decay products)



Variable: Pull

I

- In general, radiation from color flow tends to make jets lop-sided instead of round

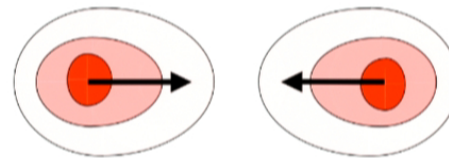
Variable: Pull

I

- In general, radiation from color flow tends to make jets lop-sided instead of round
- Gallicchio and Schwartz (arXiv:1001.5027) developed an observable to quantify this: the **pull vector**

$$\vec{t} = \sum_{i \in \text{jet}} \frac{p_{T i}}{p_{T \text{jet}}} |\vec{r}_i|^2 \hat{r}_i$$

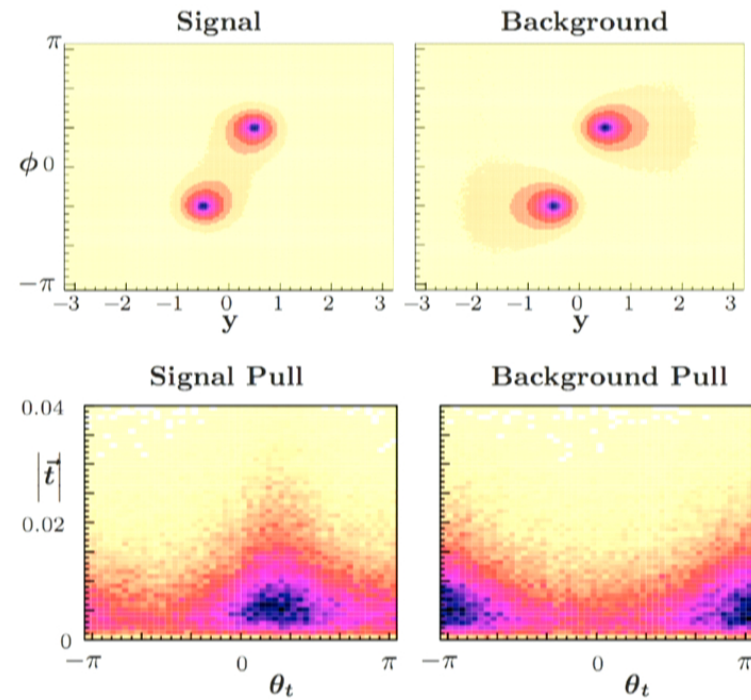
- ▶ \vec{r}_i is a vector from i to the jet center



- Pull of a jet points in the direction of the **other** pole of the color dipole

Variable: Pull

- Original example: Zh production, with $h \rightarrow b\bar{b}$, $Z \rightarrow \ell^+\ell^-$
 - ▶ Look at pulls of two b jets

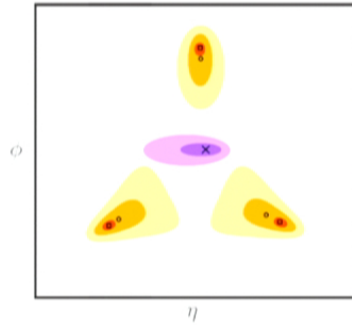


(Gallicchio, Schwartz, arXiv:1001.5027)

New variable: Radial pull

I

- In our signal:



- Subject pulls tend to point **in** toward the fat jet center
- Background tends to have more color connection with beam; some pulls point outward

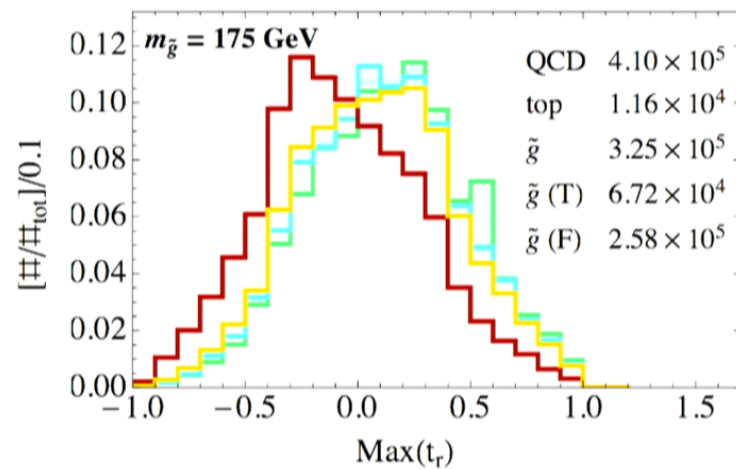
- We propose a new variable called **radial pull**:

$$t_r = -\frac{1}{N} \sum_{i=1}^N \hat{t}_i \cdot \hat{r}_i$$

- ▶ \hat{t}_i is the subject pull
- ▶ \hat{r}_i is the unit vector pointing from the subject to the fat jet center

New variable: Radial pull

- Radial pull distribution:



- Cuts:

- ▶ Pass trigger (6 jets, $p_T > 60$ GeV)
- ▶ 2 fat jets (anti- k_T , $R = 1.5$) with $p_T > 200$ GeV
- ▶ $\tau_{32} < 0.7$

Event generation

I

We consider two ranges of gluino masses:

- “Heavy” gluinos: $m_{\tilde{g}} = 450 - 1000$ GeV
 - ▶ To determine mass reach of boosted searches

Event generation

I

We consider two ranges of gluino masses:

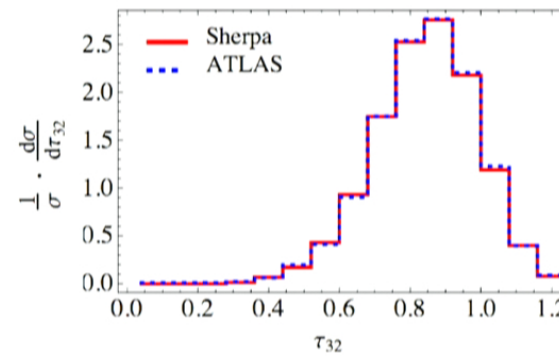
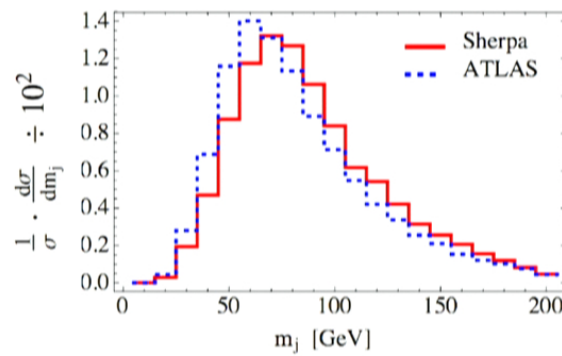
- “Heavy” gluinos: $m_{\tilde{g}} = 450 - 1000$ GeV
 - ▶ To determine mass reach of boosted searches
- “Top-mass”: $m_{\tilde{g}} \sim m_t$
 - ▶ Sample of highly boosted gluinos to demonstrate efficacy of jet substructure and color-flow observables

Event generation

- Analysis done using FastJet 3.0.2 with the N -subjettiness plug-in
- We validated our backgrounds with unfolded data from ATLAS substructure analyses at 7 TeV, 35 pb^{-1} (arXiv:1203.4606)
 - ▶ Extracted K -factor of ≈ 2 for background normalization

Event generation

- Analysis done using FastJet 3.0.2 with the N -subjettiness plug-in
- We validated our backgrounds with unfolded data from ATLAS substructure analyses at 7 TeV, 35 pb^{-1} (arXiv:1203.4606)
 - ▶ Extracted K -factor of ≈ 2 for background normalization
- For jets in 300-400 GeV bin: (similar results for other p_T)



Pile-up

- We do not explicitly include pile-up in our analysis
- We checked that applying jet grooming algorithms did not affect total efficiencies of kinematic cuts
 - ▶ Tried jet trimming with $R = 0.4$, $f_{\text{cut}} = 0.02 - 0.05$ (depending on jet p_T)
 - ▶ Found that distributions shifted, but could choose value for cuts that give comparable signal/background efficiencies

Pile-up

- We do not explicitly include pile-up in our analysis
- We checked that applying jet grooming algorithms did not affect total efficiencies of kinematic cuts
 - ▶ Tried jet trimming with $R = 0.4$, $f_{\text{cut}} = 0.02 - 0.05$ (depending on jet p_T)
 - ▶ Found that distributions shifted, but could choose value for cuts that give comparable signal/background efficiencies
- Color-flow variables depend on soft radiation concentrated around subjet centers
 - ▶ Can be mitigated by using only particles within $\Delta R = 0.4$ of subjet center to minimize effects of pile-up

Pile-up

- We do not explicitly include pile-up in our analysis
- We checked that applying jet grooming algorithms did not affect total efficiencies of kinematic cuts
 - ▶ Tried jet trimming with $R = 0.4$, $f_{\text{cut}} = 0.02 - 0.05$ (depending on jet p_T)
 - ▶ Found that distributions shifted, but could choose value for cuts that give comparable signal/background efficiencies
- Color-flow variables depend on soft radiation concentrated around subjet centers
 - ▶ Can be mitigated by using only particles within $\Delta R = 0.4$ of subjet center to minimize effects of pile-up
- New paper (Soyez *et al.*, arXiv:1211.2811) proposes a subtraction method that works very well for jet shapes
 - ▶ Could further mitigate effects of pile-up!

Cut flow for 650 GeV gluino

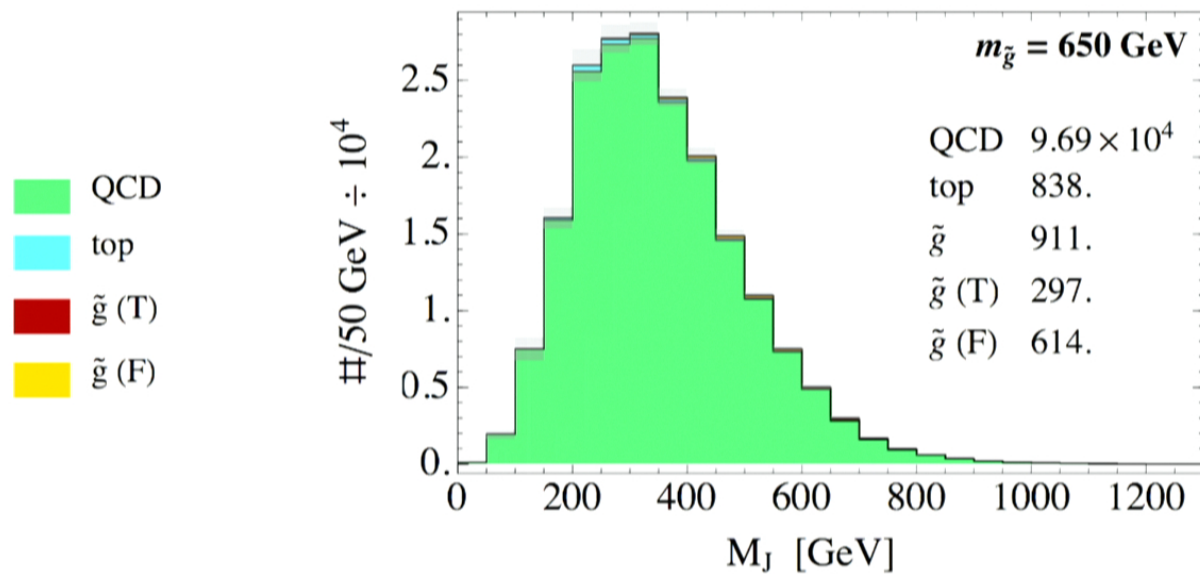
- 8 TeV, 20 fb⁻¹
- Initial sample:
 - ▶ Fat jet $p_T > 600$ GeV
 - ▶ $\tau_{32} < 0.7$

I

Cut flow for 650 GeV gluino

- 8 TeV, 20 fb⁻¹
- Initial sample:
 - ▶ Fat jet $p_T > 600$ GeV
 - ▶ $\tau_{32} < 0.7$

I



Cut flow for 650 GeV gluino

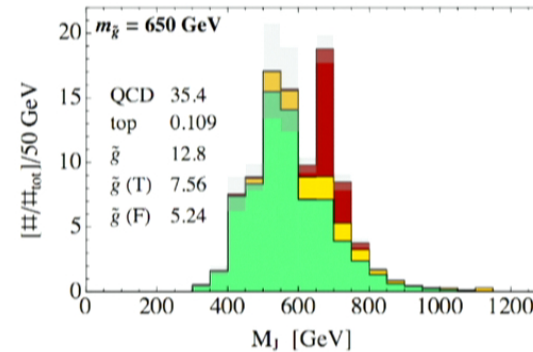
■ QCD

■ top

■ \tilde{g} (T)

■ \tilde{g} (F)

- The “kinematic” cuts give the distribution at **right**
- Adding a soft cut on axis contraction cleans up distribution (see **below**)



Shape analysis results

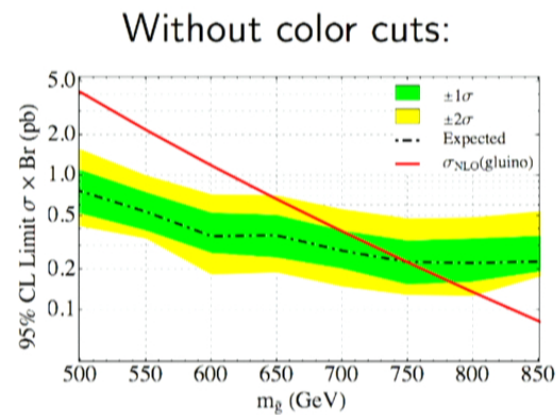
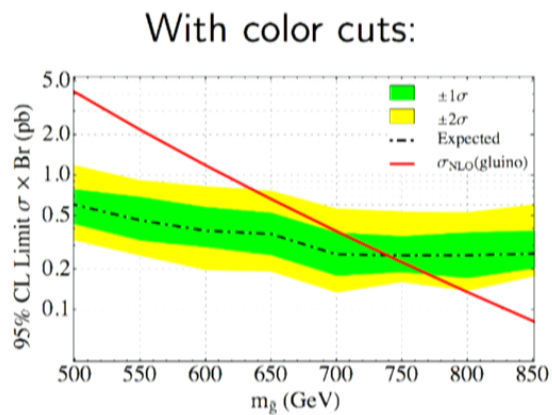
- Perform shape analysis for range of gluino masses 450 – 1000 GeV to determine exclusion reach

I

Shape analysis results

- Perform shape analysis for range of gluino masses 450 – 1000 GeV to determine exclusion reach
- At 20 fb^{-1} :

I

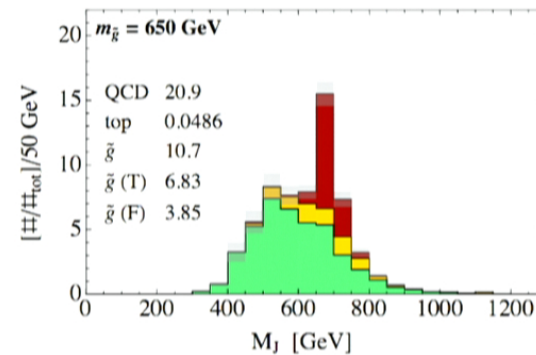
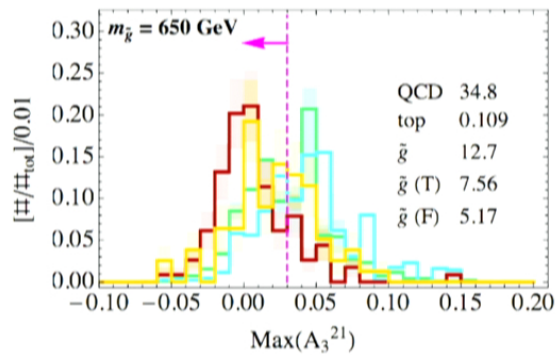
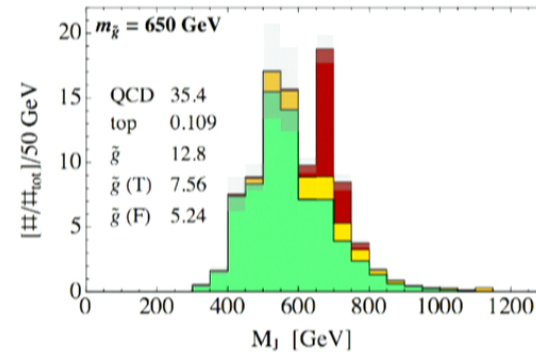


- Expected exclusion is $\sim 750 \text{ GeV}$ ($\sim 650 \text{ GeV}$) at 20 fb^{-1} (5 fb^{-1})

Cut flow for 650 GeV gluino

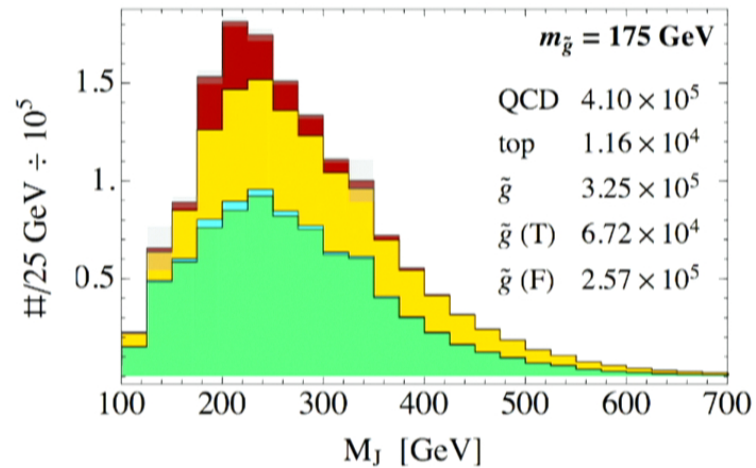


- The “kinematic” cuts give the distribution at **right**
- Adding a soft cut on axis contraction cleans up distribution (see **below**)



Analysis details

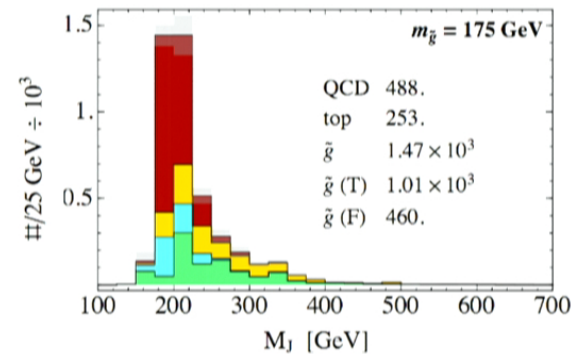
- Much higher cross section and boosted fraction, can afford harsher cuts
 - ▶ For example, a cut on radial pull has signal efficiency of 12%, but combinatoric background/top mistag rate of 2% and QCD mistag rate of 0.2%!
- Initial sample:
 - ▶ Fat jet $p_T > 200$ GeV
 - ▶ $\tau_{32} < 0.7$



Cut flow for top-mass GeV gluino



- The “kinematic” cuts give the distribution at **right** ($S/B \sim 2$)
- Adding a harsh cut on radial pull leaves almost no background!



Analysis details

I

- Compare with recent ATLAS results for $m_{\tilde{g}} = 175$ GeV
 - ▶ ATLAS excludes $\text{Br}(\tilde{g} \rightarrow qqq) < 0.25$ at 95% C.L. with 7 TeV, 5 fb^{-1}
 - ▶ Our analysis would exclude $\text{Br}(\tilde{g} \rightarrow qqq) < 0.15$ at 95% C.L. with 8 TeV, 5 fb^{-1} and $\text{Br}(\tilde{g} \rightarrow qqq) < 0.05$ at 20 fb^{-1}
- We show that harsh cuts on a combination of substructure and color-flow observables can give an extremely pure sample of gluinos
 - ▶ Could be useful if evidence for new physics found in hadronic final states!

Conclusions

- Boosted searches provide a complementary way to study all-hadronic final states
 - ▶ Independent systematics from resolved analyses and other resonance searches
 - ▶ Higher S/B
 - ▶ Can be competitive with resolved searches, even at higher masses
 - ▶ Good way to tag new-physics events
- Color-flow observables can be powerful!
 - ▶ First study of color-flow pattern from R -hadron decay
 - ▶ Introduced new observables
 - ★ Radial pull, generalizing pull-based dipole taggers
 - ★ Axis contraction, based on N -subjettiness
- Future work: explore other applications for these methods
 - ▶ $\tilde{g} \rightarrow q\bar{q}\tilde{\chi}_0, \tilde{\chi}_0 \rightarrow udd$
 - ▶ $h \rightarrow gg?$