

Title: Quantification and Evolution of Quantum Correlations

Date: Nov 28, 2012 04:00 PM

URL: <http://www.pirsa.org/12110096>

Abstract:

Quantification and Evolution of Quantum Correlations

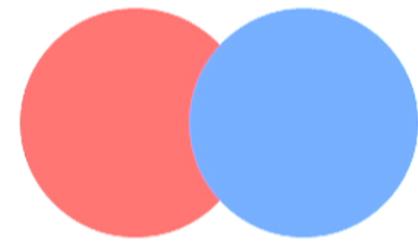


Asma Al-Qasimi
Department of Physics
University of Toronto

Perimeter Institute
for Theoretical Physics
Wednesday, November 28, 2012
4:00pm-5:00pm
The Time Room

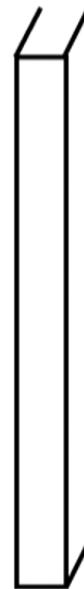
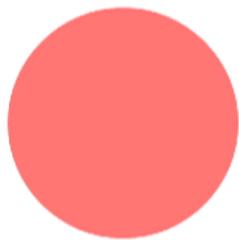
Intro + Outline

- Correlations, Quantum Correlations
- Why study Quantum Correlations?
- Different Quantification Methods:
 - 1) Information-Based
 - 2) Local Operation
 - 3) Measurement-Based
 - 4) Distance Measure
 - 5) Geometric Representation
- Evolution of Quantum Correlations
 - 1) Entanglement Sudden Death in Two Qubits
 - 2) Entanglement Sudden Death in Two Harmonic Oscillators
 - 3) Correlations in Atomic Spectroscopy

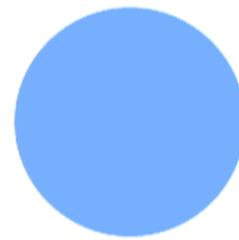




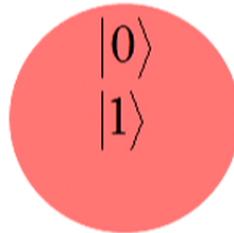
Alice



Bob

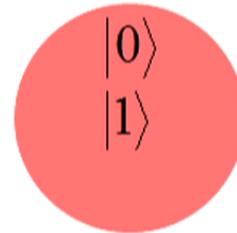


State Representation of Quantum Systems



two-level
quantum
system

State Representation of Quantum Systems

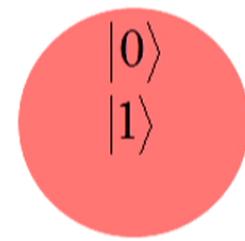


two-level
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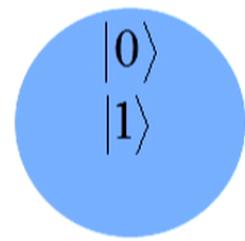
Pure State

$$|\psi\rangle = A|0\rangle + B|1\rangle$$

$$\rho = |\psi\rangle\langle\psi| = \begin{pmatrix} |A|^2 & AB^* \\ BA^* & |B|^2 \end{pmatrix}$$



$|0\rangle$
 $|1\rangle$



$|0\rangle$
 $|1\rangle$

System C = System A + System B

$$|\Psi\rangle_C \neq |\psi\rangle_A \otimes |\psi\rangle_B$$

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$$|\Psi\rangle_C \neq |\psi\rangle_A \otimes |\psi\rangle_B$$

$$|\Psi\rangle = \frac{1}{\sqrt{2}}(|01\rangle - |10\rangle)$$

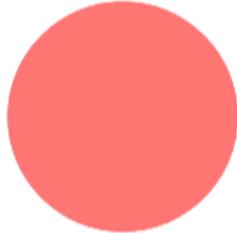
Bell State

Maximally Entangled
Pure State



Alice

Bob

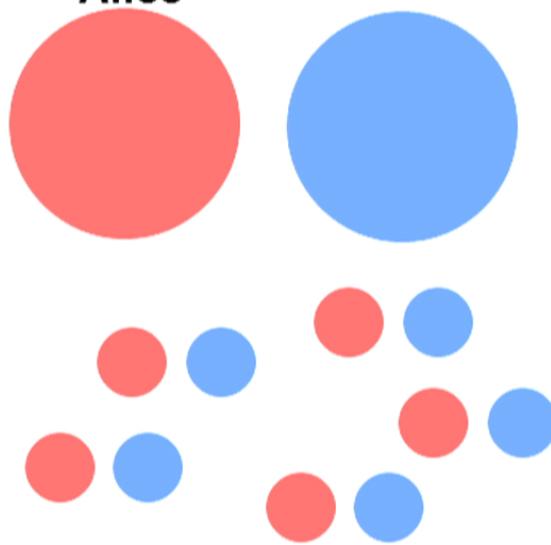


Bennett *et. al.*, PRA, **53** (4), 2046 (1996).



Alice

Bob



How much entanglement is there
between the two qubits?

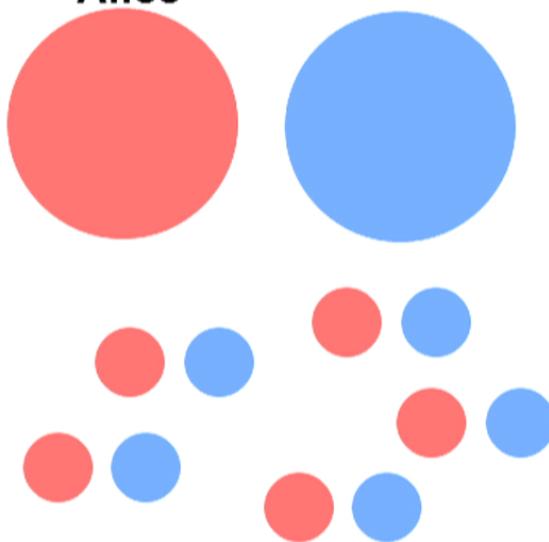
Look at m realizations of the pair
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Bennett *et. al.*, PRA, **53** (4), 2046 (1996).



Alice

Bob



How much entanglement is there
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Look at m realizations of the pair
of qubits, and ask :

By performing local operations and classically communicating
with each, how many Bell states (n) can Alice and Bob create?

Bennett et. al., PRA, 53 (4), 2046 (1996).

How much information is required to create the state of interest?

W. K. Wootters, Phys. Rev. Lett. **80** (10), 2245 (1998).

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Ebits!!

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Entanglement of Formation = von Neumann entropy of the reduced density matrix of ρ

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$$E(\rho) = S(\rho_{red}^A) = S(\rho_{red}^B)$$

straightforward
for pure states

not as straightforward
for mixed states

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not as straightforward
for mixed states

$$\rho = \sum_i p_i |\psi_i\rangle\langle\psi_i|$$

Monotonic Function
of **Concurrence**

$$E(\rho) = \min \sum_i p_i E(\psi_i)$$

W. K. Wootters, Phys. Rev. Lett. **80** (10), 2245 (1998).

Concurrence is a very useful tool to study entanglement

Asma Al-Qasimi and Daniel F. V. James *Phys Rev A* **77**, 012117 (2008).

Concurrence is a very useful tool to study entanglement

Example: Studies of Entanglement Sudden Death (ESD)

Asma Al-Qasimi and Daniel F. V. James *Phys Rev A* **77**, 012117 (2008).

Mathematical Description of System Investigated and its Dynamics

X-states

$$\hat{\rho}(t) = \begin{pmatrix} a(t) & 0 & 0 & w(t) \\ 0 & b(t) & z(t) & 0 \\ 0 & z^*(t) & c(t) & 0 \\ w^*(t) & 0 & 0 & d(t) \end{pmatrix}$$

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Master Equation

$$\frac{\partial \hat{\rho}}{\partial t} = \frac{1}{i\hbar} [\hat{H}, \hat{\rho}] + L_1[\hat{\rho}] + L_2[\hat{\rho}]$$

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Liouvillian of the two
qubits

$$L_i[\hat{\rho}] = \frac{(\bar{n}+1)\Gamma}{2} \left[[\hat{\sigma}_-, \hat{\rho}\hat{\sigma}_+] + [\hat{\sigma}_-^i \hat{\rho}, \hat{\sigma}_+^i] \right] + \frac{\bar{n}\Gamma}{2} \left[[\hat{\sigma}_+, \hat{\rho}\hat{\sigma}_-^i] + [\hat{\sigma}_+^i \hat{\rho}, \hat{\sigma}_-^i] \right]$$

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$$L_i[\hat{\rho}] = \frac{(\bar{n}+1)\Gamma}{2} [(\hat{\sigma}_-, \hat{\rho} \hat{\sigma}_+) + (\hat{\sigma}_- \hat{\rho}, \hat{\sigma}_+)] + \frac{\bar{n}\Gamma}{2} [(\hat{\sigma}_+, \hat{\rho} \hat{\sigma}_-) + (\hat{\sigma}_+ \hat{\rho}, \hat{\sigma}_-)]$$

Depopulation due to stimulated and spontaneous emissions

How to Calculate Concurrence

For mixed density matrix $\hat{\rho}$,
Find eigenvalues, $\{\lambda_i\}$, of the R matrix,

where $R = \hat{\rho} \Lambda \hat{\rho}^* \Lambda$ and $\Lambda = \hat{\sigma}_y \otimes \hat{\sigma}_y$.

$$C = \max(\sqrt{\lambda_1} - \sqrt{\lambda_2} - \sqrt{\lambda_3} - \sqrt{\lambda_4}, 0),$$

$$\lambda_1 > \lambda_2 > \lambda_3 > \lambda_4$$

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$$\lambda_1 > \lambda_2 > \lambda_3 > \lambda_4$$

$C = 0$ no entanglement

$C = 1$ maximally entangled

$$\text{Find } \hat{\rho}(t) = \begin{pmatrix} a(t) & 0 & 0 & w(t) \\ 0 & b(t) & z(t) & 0 \\ 0 & z^*(t) & c(t) & 0 \\ w^*(t) & 0 & 0 & d(t) \end{pmatrix}$$

$$C = 2 \max \left\{ 0, |z(t)| - \sqrt{a(t)d(t)}, |w(t)| - \sqrt{b(t)c(t)} \right\}$$

$C = 0$ equivalent to solving two quartic equations in X .

$$X = e^{-(2\bar{n}+1)\Gamma t}$$

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$C = 0$ equivalent to solving two quartic equations in X .

$$X = e^{-(2\bar{n}+1)\Gamma t} \quad \begin{aligned} t = 0 &\Rightarrow X = 1 \\ t = \infty &\Rightarrow X = 0 \end{aligned}$$

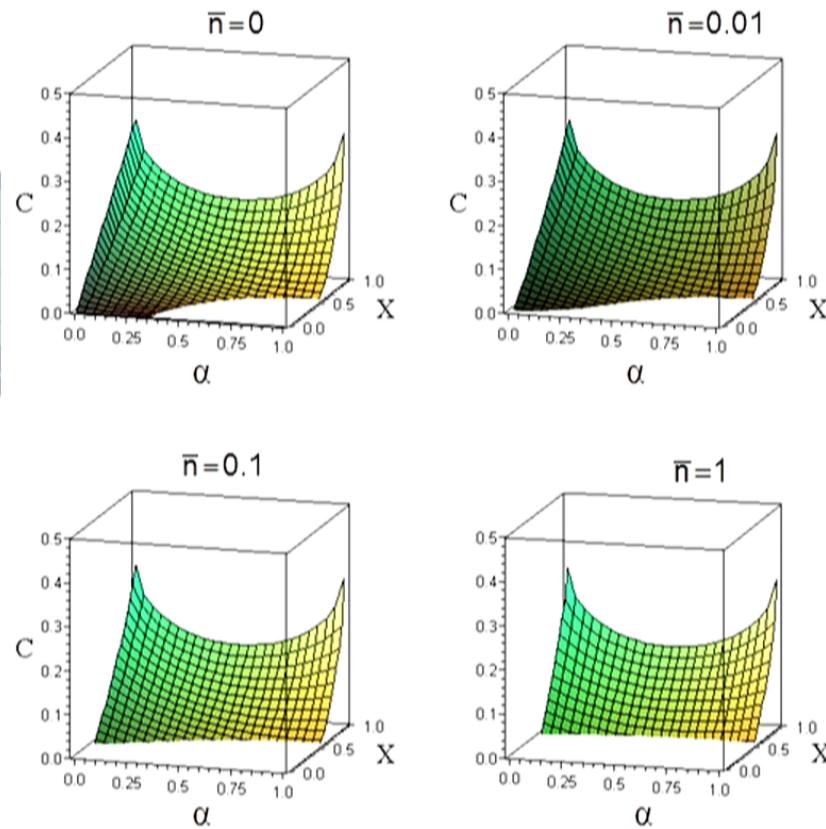
$$\hat{\rho}(0) = \frac{1}{3} \begin{pmatrix} \alpha & 0 & 0 & 0 \\ 0 & 1 & 1 & 0 \\ 0 & 1 & 1 & 0 \\ 0 & 0 & 0 & 1-\alpha \end{pmatrix}$$

T. Yu and J. H. Eberly.
Phys. Rev. Lett **93**,
140404 (2004).

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 showed that, at T=0 there is no sudden death for $0 \leq \alpha \leq 1/3$ only.



Partial Transpose Criterion for Separability

ρ = density matrix of total system

transpose over one system:

$$\sigma = \sum_A w_A (\rho'_A)^T \otimes \rho''_A$$
$$\sigma_{m\mu,n\nu} \equiv \rho_{n\mu,m\nu}$$

A. Peres, *Phys. Rev. Lett.* **77**, No. 8, 1413-1415 (1996).

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σ has at least one negative eigenvalue \rightarrow System entangled
Otherwise \rightarrow System either entangled or separable

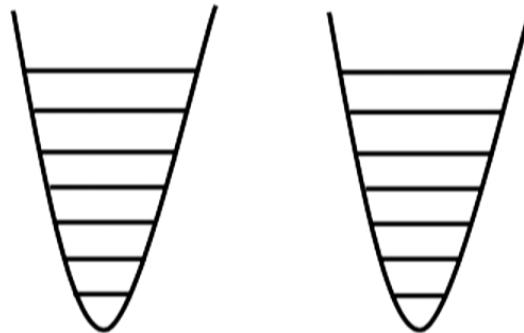
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The weakness in the separability criterion does not prevent it from being useful in some studies.

For example, in the study of the phenomena of Entanglement sudden death (ESD), the nonexistence of ESD can be proven with certainty, but not its existence.

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(General) N-photon-two-mode state undergoing dephasing

$$\begin{aligned}\rho(t) = & \sum_{k=0}^N \rho_{kk}(0) |N-k, k\rangle\langle N-k, k| \\ & + \sum_{\substack{k, m=0 \\ k \neq m}}^N \rho_{km}(0) e^{-\frac{1}{2}(k-m)^2(\Gamma_1 + \Gamma_2)t} |N-k, k\rangle\langle N-m, m|\end{aligned}$$

For σ_{PT} :

$$|\psi\rangle = \frac{1}{\sqrt{2}} \left\{ e^{i\theta} |N-k, m\rangle - |N-m, k\rangle \right\}$$
$$\theta = \text{Arg}(\rho_{km})$$

is an eigenvector with eigenvalue:

$$-\rho_{km}(t) = -\rho_{km}(0) e^{-\frac{1}{2}(k-m)^2(\Gamma_1 + \Gamma_2)t}$$

Al-Qasimi and James , *Optics Letters* **34**, 268-270 (2009).

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Negative!

For a general two-mode-N-photon state undergoing pure dephasing, there is no sudden death of entanglement.

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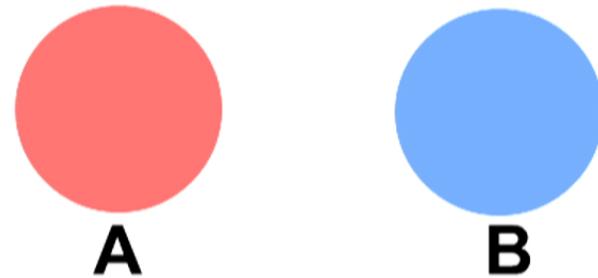
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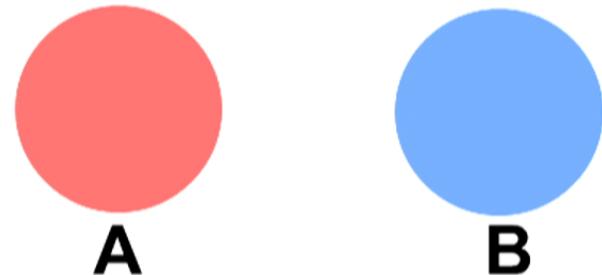
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Quantum Discord: Measurement of Quantum Systems Disturbs Them



H. Ollivier and W. H. Zurek, *Phys. Rev. Lett.* **88**, 017901 (2001).

Quantum Discord: Measurement of Quantum Systems Disturbs Them



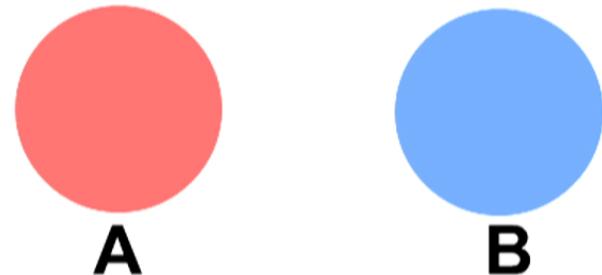
$$I(\rho) = S(\rho^A) + S(\rho^B) - S(\rho)$$

ρ^i = reduced density matrix of system i

Perform a measurement on system B only.

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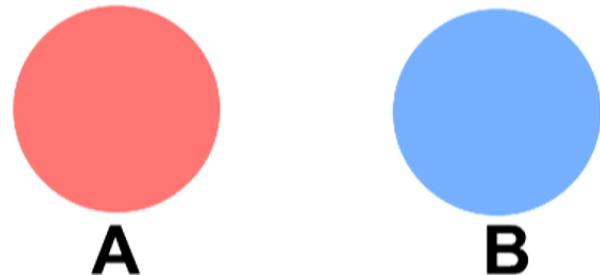
Calculate $I(\rho)$ before and after the measurement and compare.



$$I(\rho)$$

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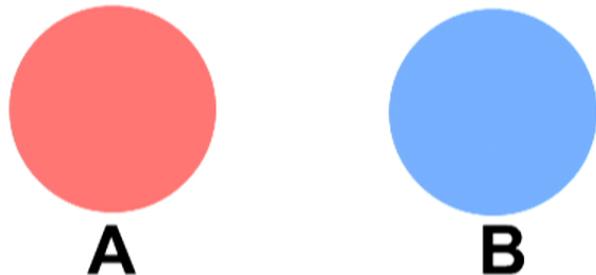
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$$Q(\rho) = I(\rho) - C(\rho)$$

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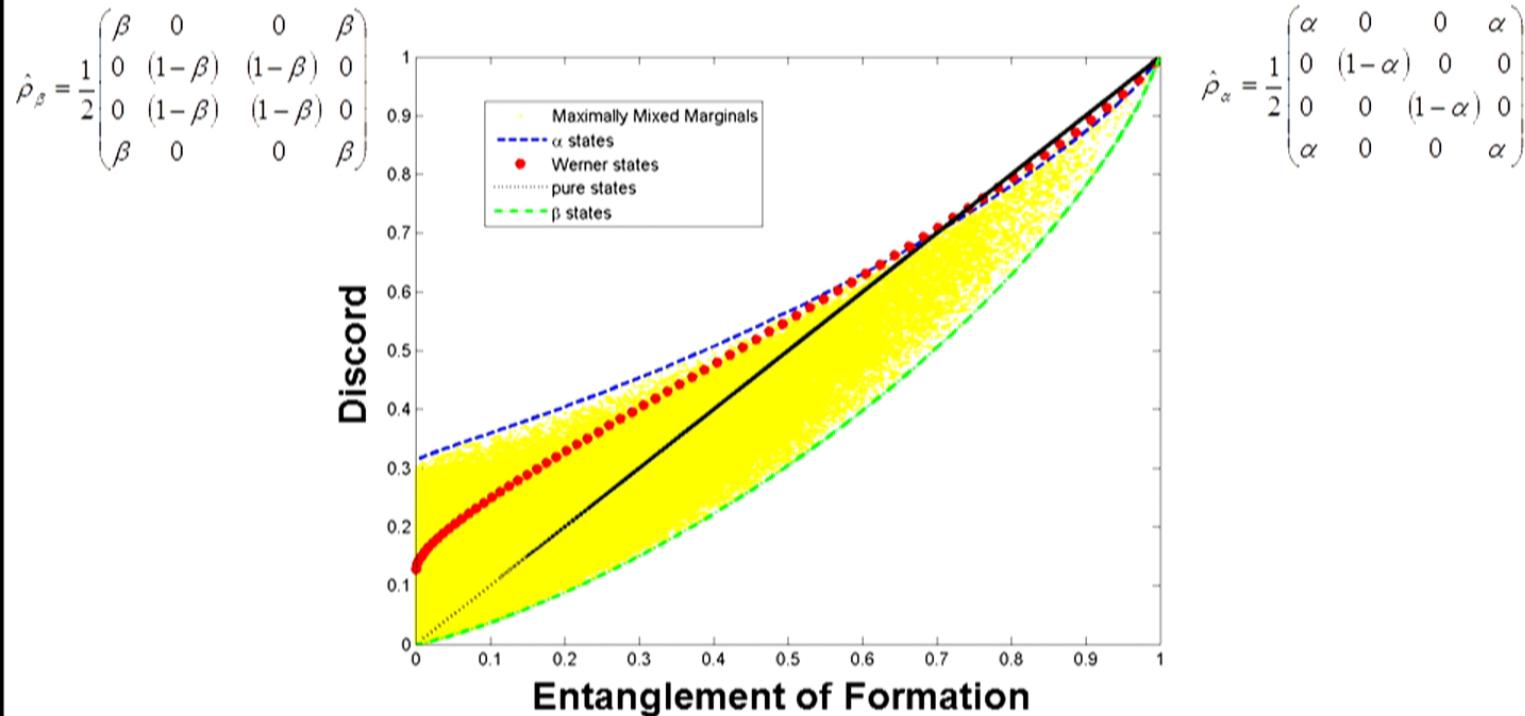
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To obtain final form of $C(\rho)$, maximize over all measurements

H. Ollivier and W. H. Zurek, *Phys. Rev. Lett.* **88**, 017901 (2001).

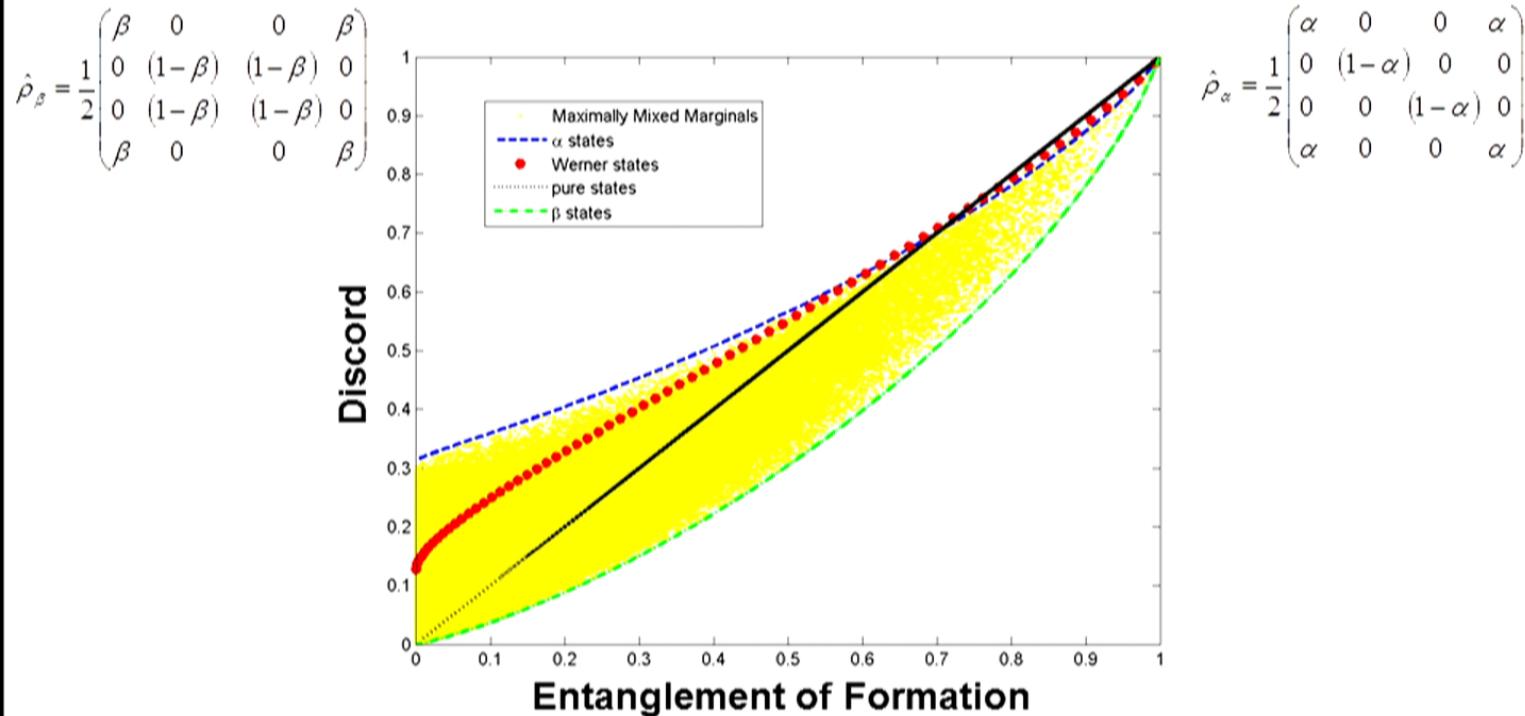
Discord versus Entanglement



$$\rho_{Werner} = \frac{1-z}{4} I + z |\psi\rangle\langle\psi| \quad |\psi\rangle = \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle)$$

Al-Qasimi and James, *Physical Review A*, **83**, 032101 (2011).

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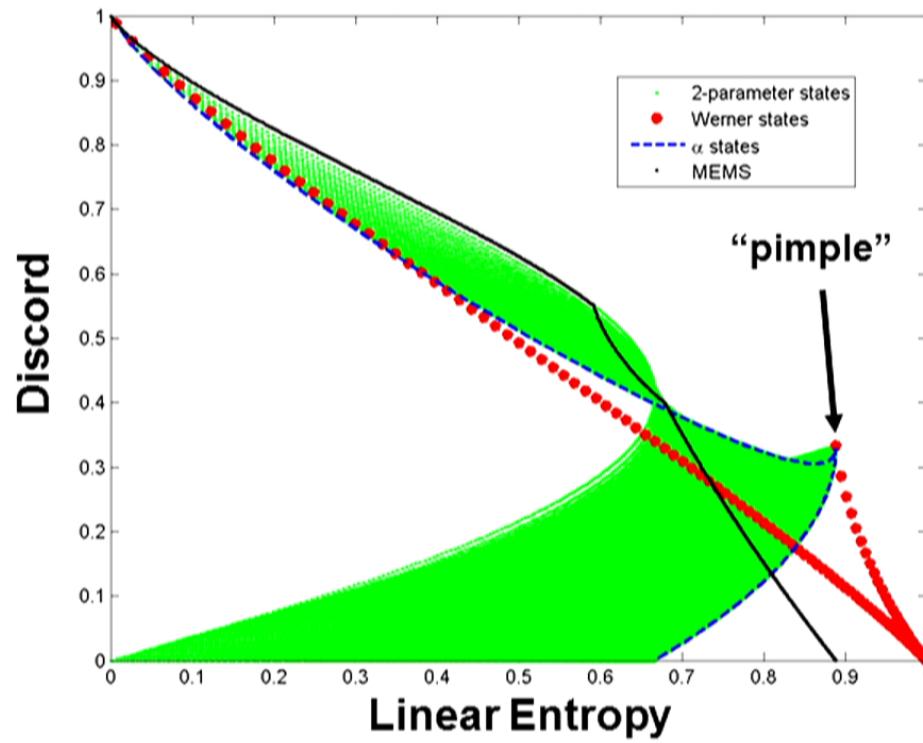


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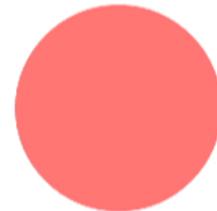
Al-Qasimi and James, *Physical Review A*, **83**, 032101 (2011).

Discord versus Entropy

$$\hat{\rho}_{a,b} = \frac{1}{2} \begin{pmatrix} a & 0 & 0 & a \\ 0 & (1-a-b) & 0 & 0 \\ 0 & 0 & (1-a+b) & 0 \\ a & 0 & 0 & a \end{pmatrix}$$



Spectroscopy Using Atoms

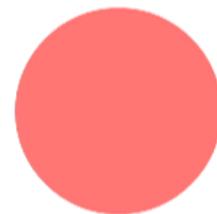


$$|\psi\rangle = \frac{1}{\sqrt{2}}(|g\rangle + e^{i\phi}|e\rangle)$$

In a Ramsey Experiment, by measuring the population of atoms in a certain state, ϕ can be determined.

* Chwalla *et al.*, Appl. Phys. B **89**, 483 (2007).

Spectroscopy Using Atoms

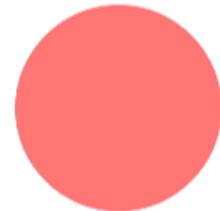


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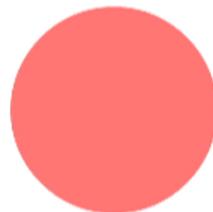
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Correlations in Quantum-Enhanced Atomic Spectroscopy*

2 ions in a trap with *dephasing* going on.

* Al-Qasimi, James, Blatt, and White (In Preparation)

Correlations in Quantum-Enhanced Atomic Spectroscopy*

2 ions in a trap with ***dephasing*** going on.

Initial State

$$|\psi\rangle = A|gg\rangle + B|ge\rangle + C|eg\rangle + D|ee\rangle$$

Evolves into the form:

$$\rho = \begin{pmatrix} a & 0 & 0 & 0 \\ 0 & b & z & 0 \\ 0 & z^* & c & 0 \\ 0 & 0 & 0 & d \end{pmatrix}$$

* Al-Qasimi, James, Blatt, and White (In Preparation)

Parity

$$P = 2 \operatorname{Re}\{z\}$$

how well the relative phase
is being measured

* Quesada, Al-Qasimi, and James, *J. Mod. Opt.*, **59** (15), 1322 (2012).

Parity

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Entanglement

$$C = 2 \max \left\{ 0, |z| - \sqrt{ad} \right\}$$

* Quesada, Al-Qasimi, and James, *J. Mod. Opt.*, **59** (15), 1322 (2012).

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Discord*

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$$N_1 = H \left(\frac{1}{2} + \frac{1}{2} \sqrt{(a-d+b-c)^2 + 4z^2} \right)$$

$$\begin{aligned} H(y) &= -y \log_2 y \\ &\quad -(1-y) \log_2 (1-y) \end{aligned}$$

Binary Entropy Function

* Quesada, Al-Qasimi, and James, *J. Mod. Opt.*, **59** (15), 1322 (2012).

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Binary Entropy Function

$$N_2 = -a \log_2 \left[\frac{a}{a+c} \right] - b \log_2 \left[\frac{b}{b+d} \right]$$

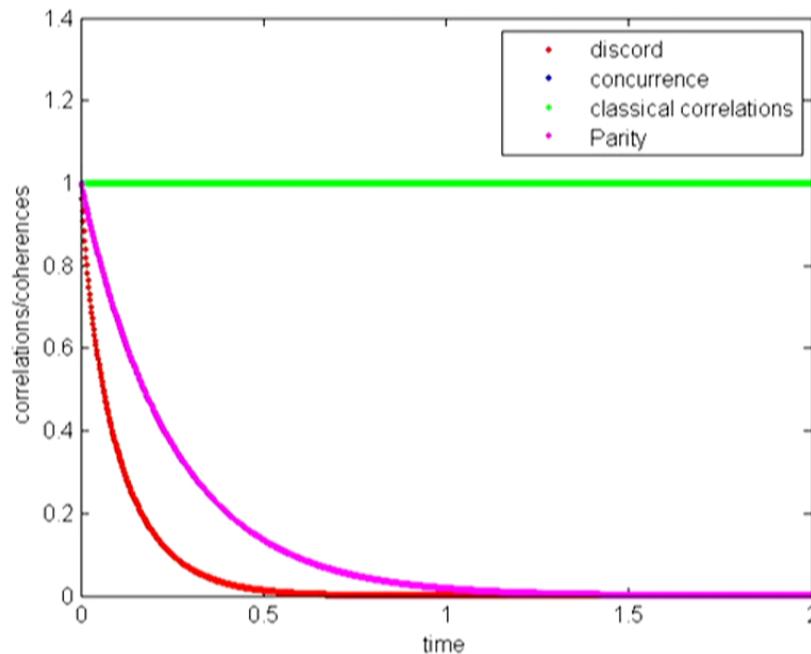
$$-c \log_2 \left[\frac{c}{a+c} \right] - d \log_2 \left[\frac{d}{b+d} \right]$$

* Quesada, Al-Qasimi, and James, *J. Mod. Opt.*, **59** (15), 1322 (2012).

Examples

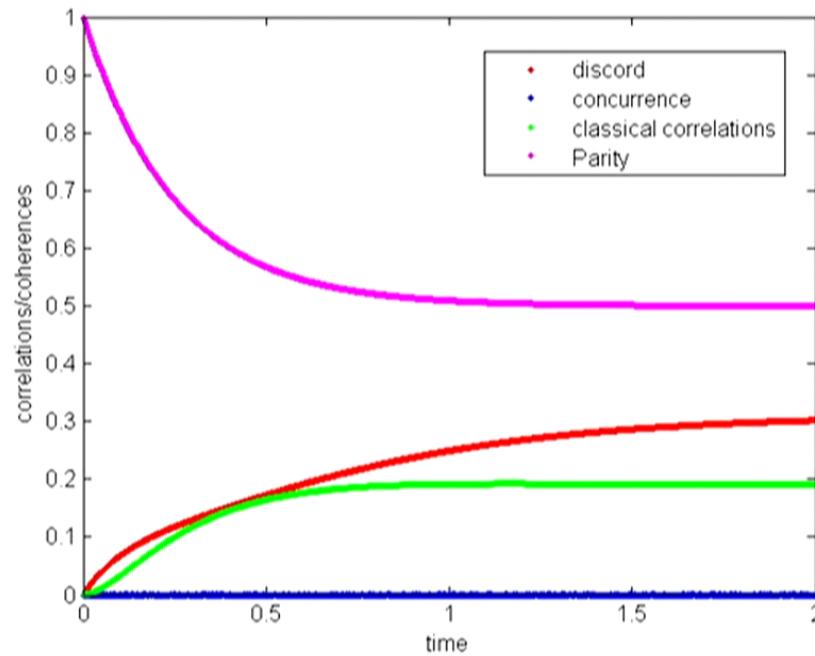
Initial state is a Bell state

$$|\psi\rangle = \frac{1}{\sqrt{2}}(|00\rangle + |11\rangle)$$

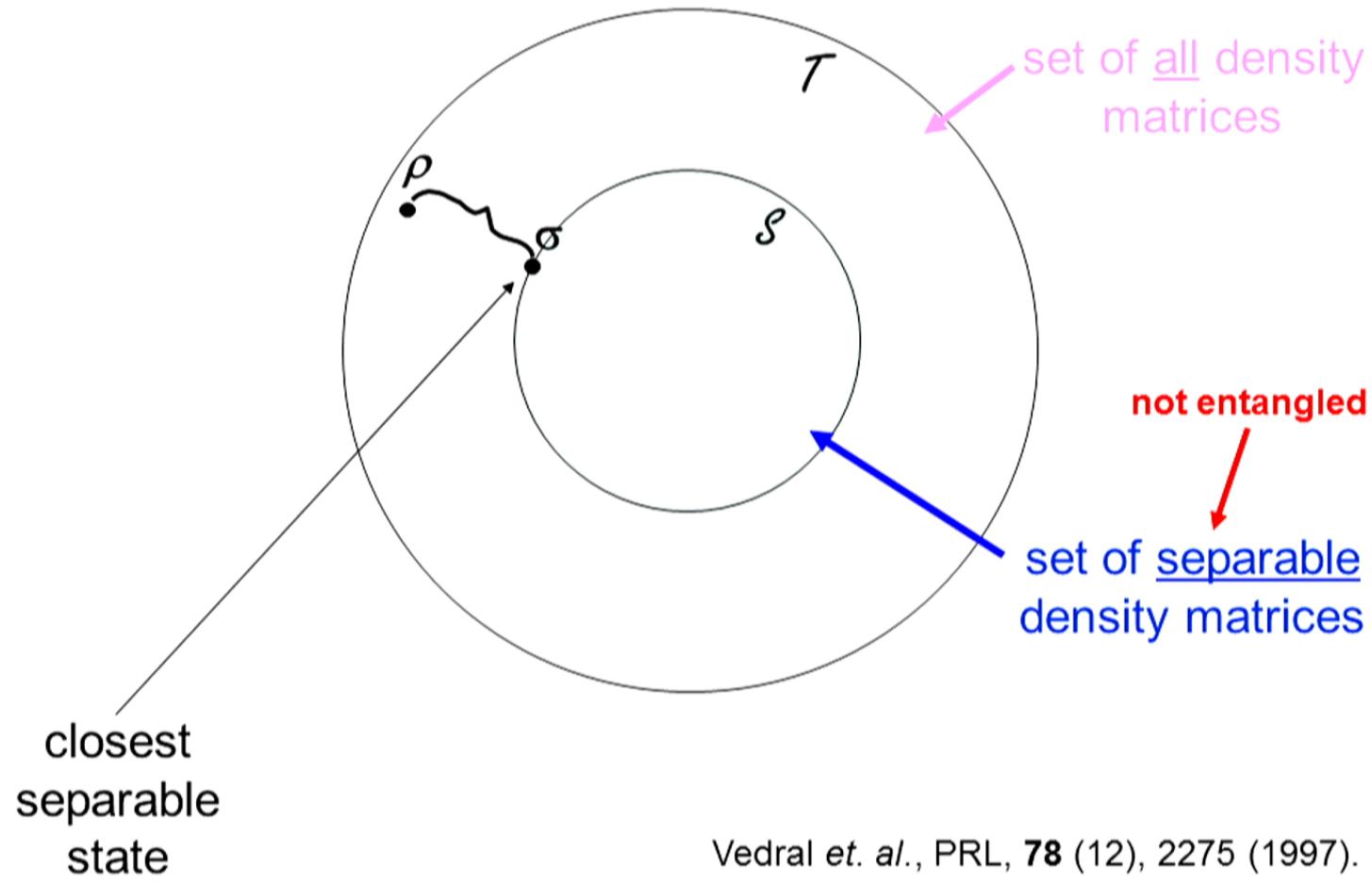


Initial state is a separable state

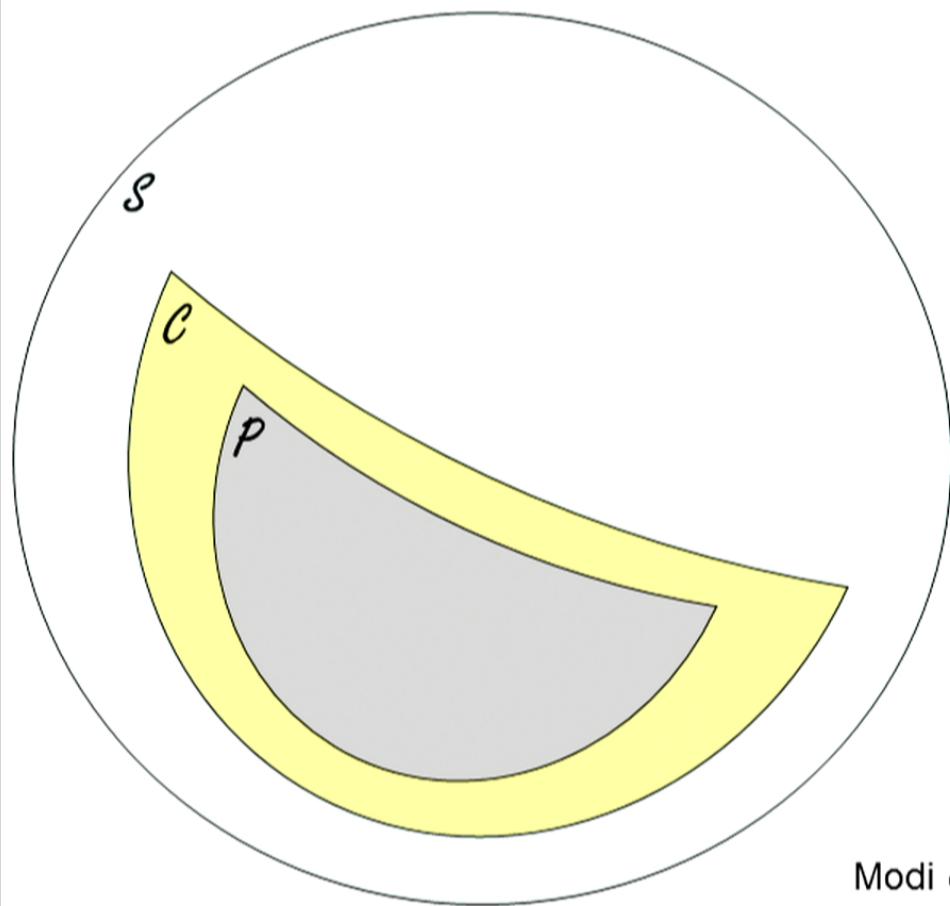
$$|\psi\rangle = \frac{1}{2}(|00\rangle + |01\rangle + |10\rangle + |11\rangle)$$



Using ‘Distance’ to Quantify Correlations

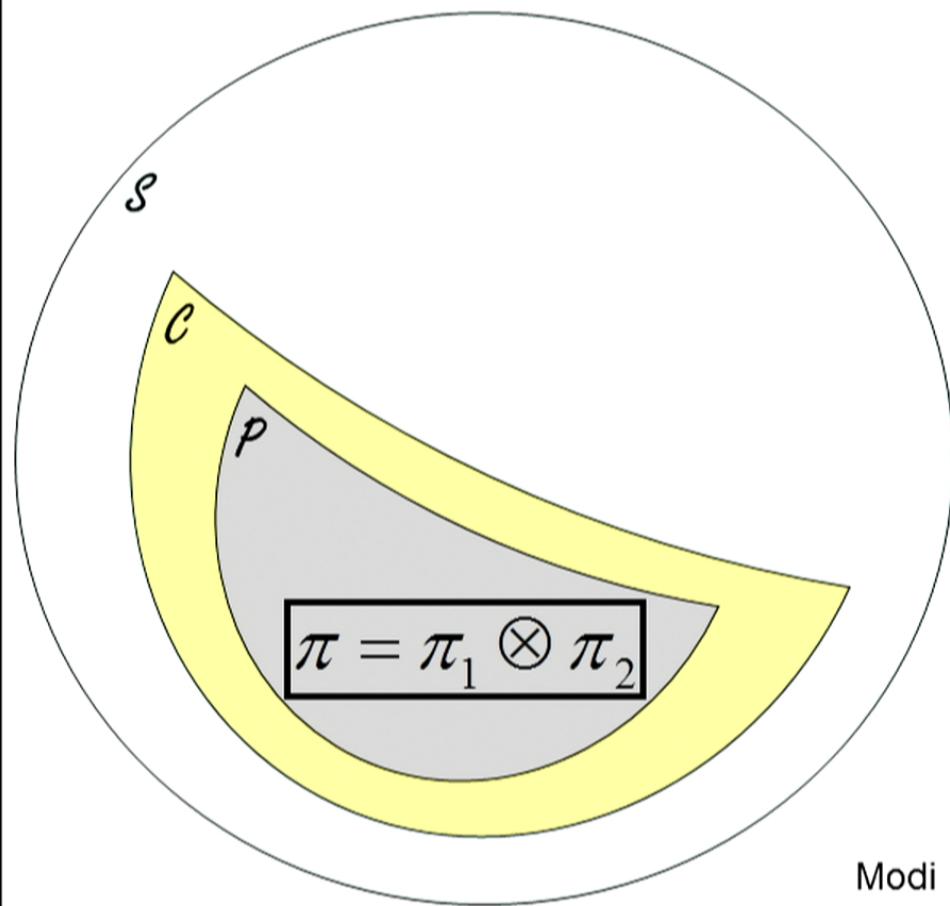


Hierarchy of Separable States



Modi *et. al.*, PRL, **104**, 080501 (2010).

Hierarchy of Separable States

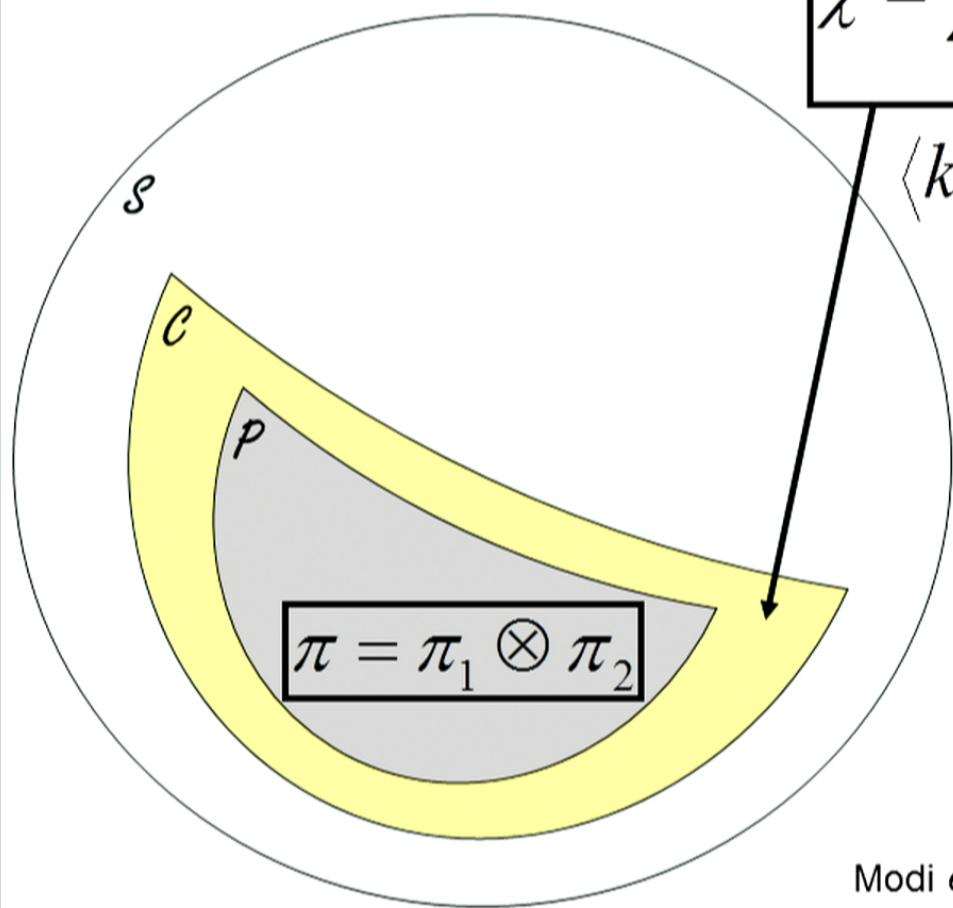


Modi et. al., PRL, 104, 080501 (2010).

Hierarchy of Separable States

$$\chi = \sum_{k_n} p_{k_1 k_2} |k_1 k_2\rangle\langle k_1 k_2|$$

$$\langle k_1 k_2 | k'_1 k'_2 \rangle = \delta_{k_1 k'_1} \delta_{k_2 k'_2}$$



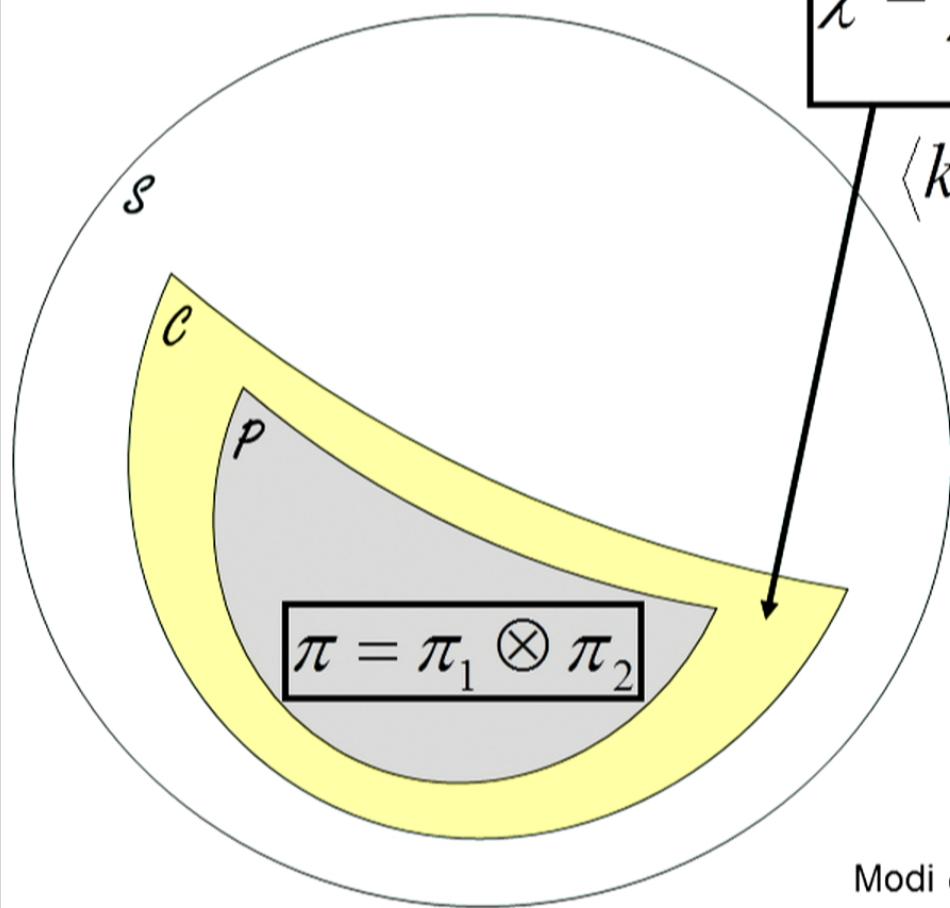
$$\pi = \pi_1 \otimes \pi_2$$

Modi et. al., PRL, 104, 080501 (2010).

Hierarchy of Separable States

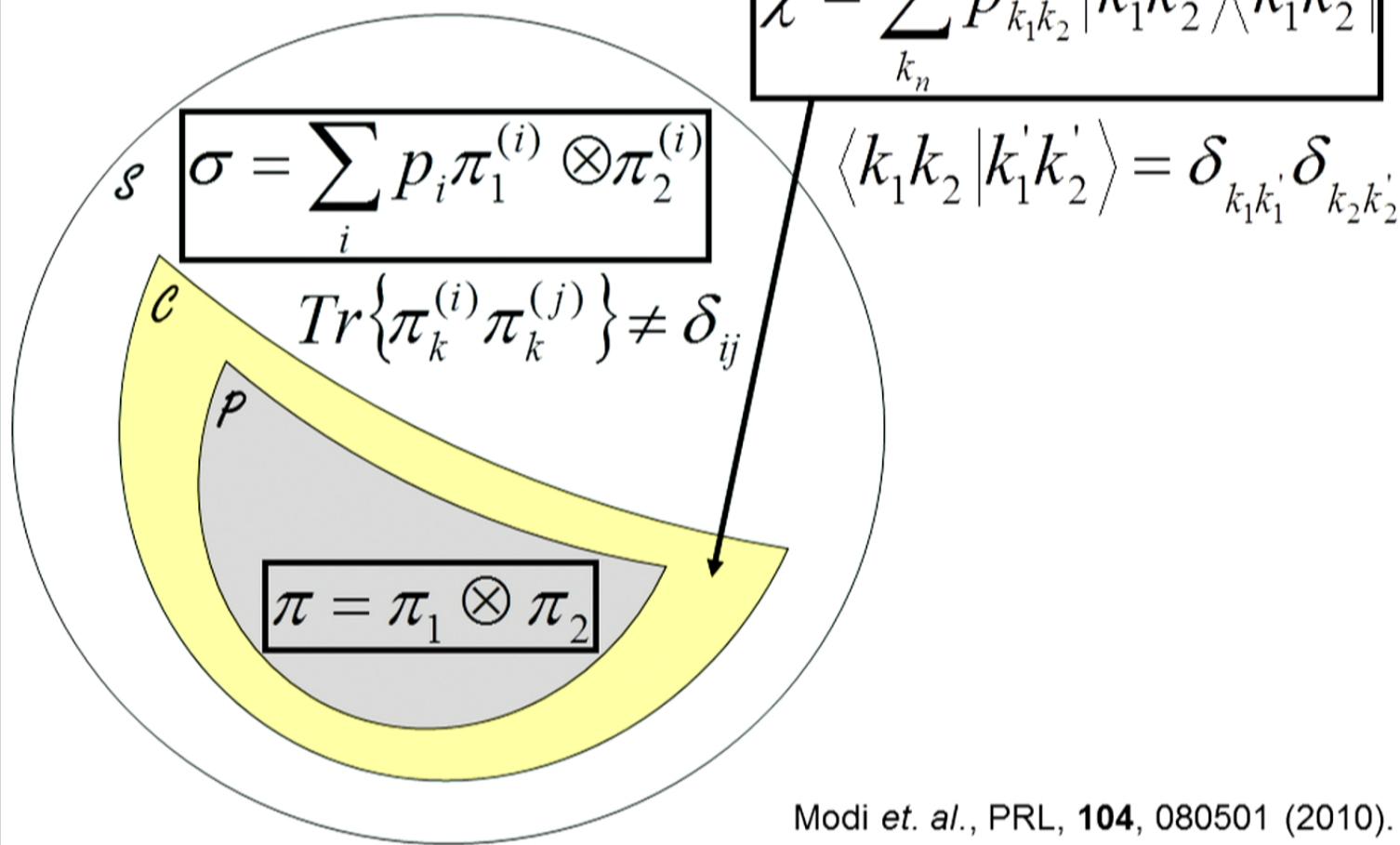
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Modi et. al., PRL, 104, 080501 (2010).

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Modi et. al., PRL, 104, 080501 (2010).

Given density matrix ρ

Modi *et. al.*, PRL, **104**, 080501 (2010).

Given density matrix ρ

Entanglement → shortest distance between ρ and a
separable state (σ')

Discord → shortest distance between ρ and a **classical** state (χ')

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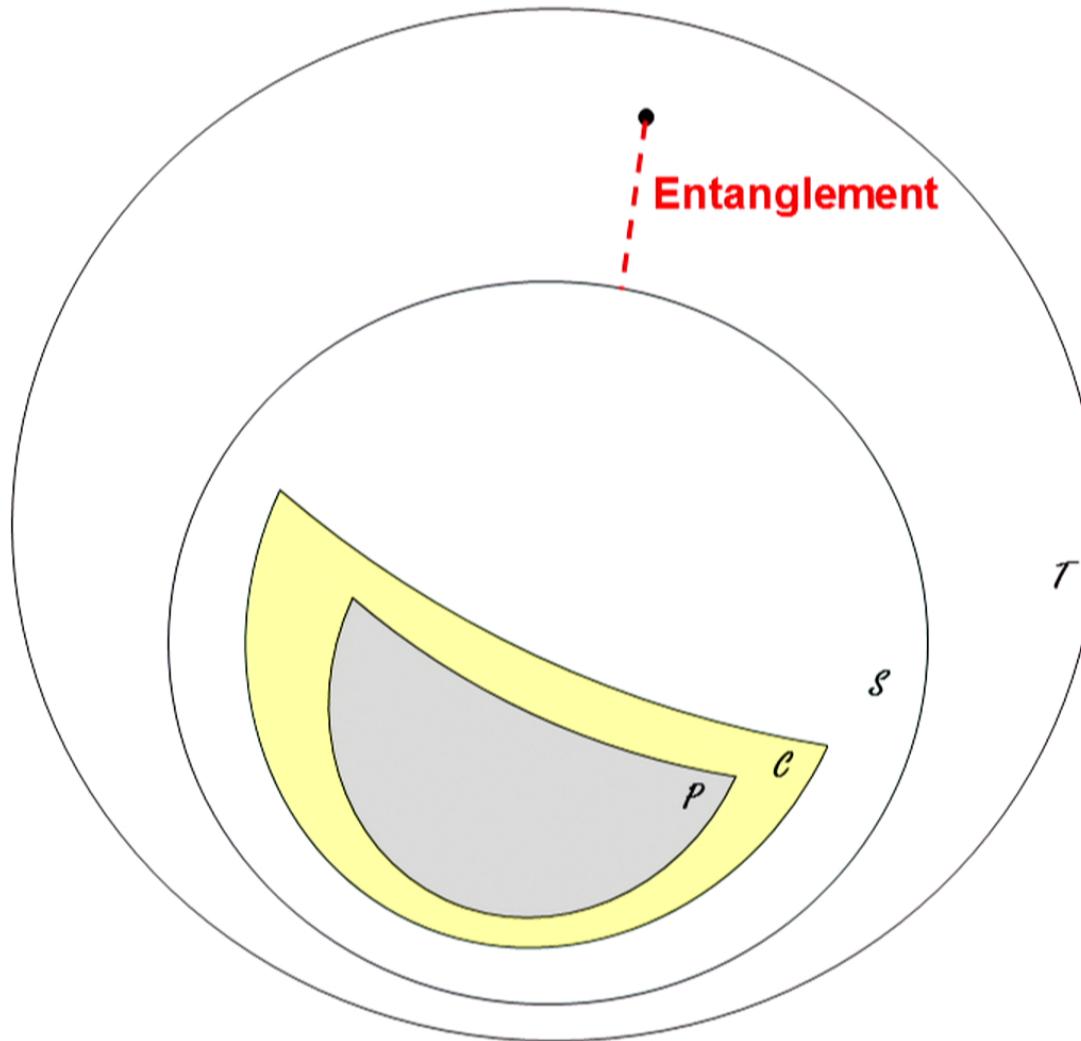
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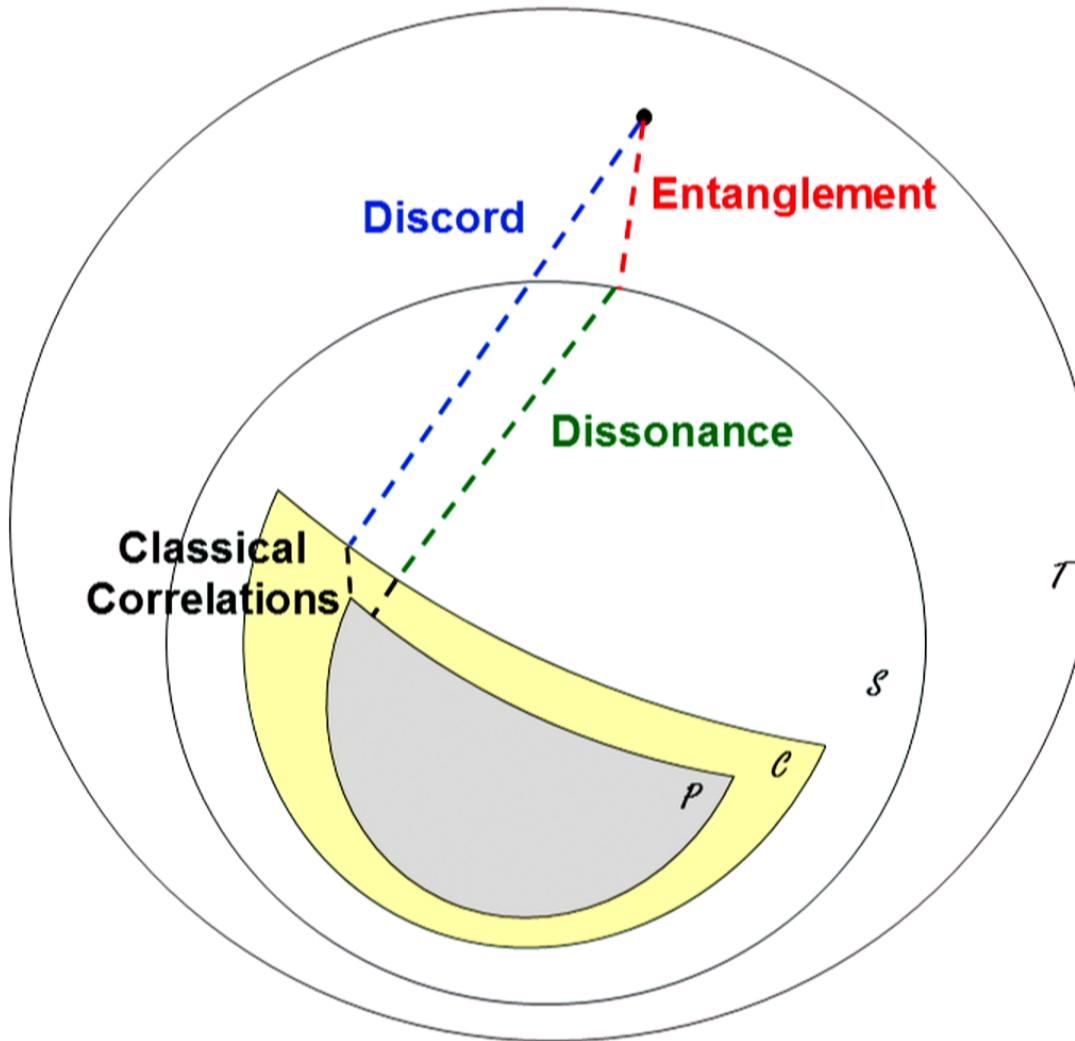
Entanglement → shortest distance between ρ and a
separable state (σ')

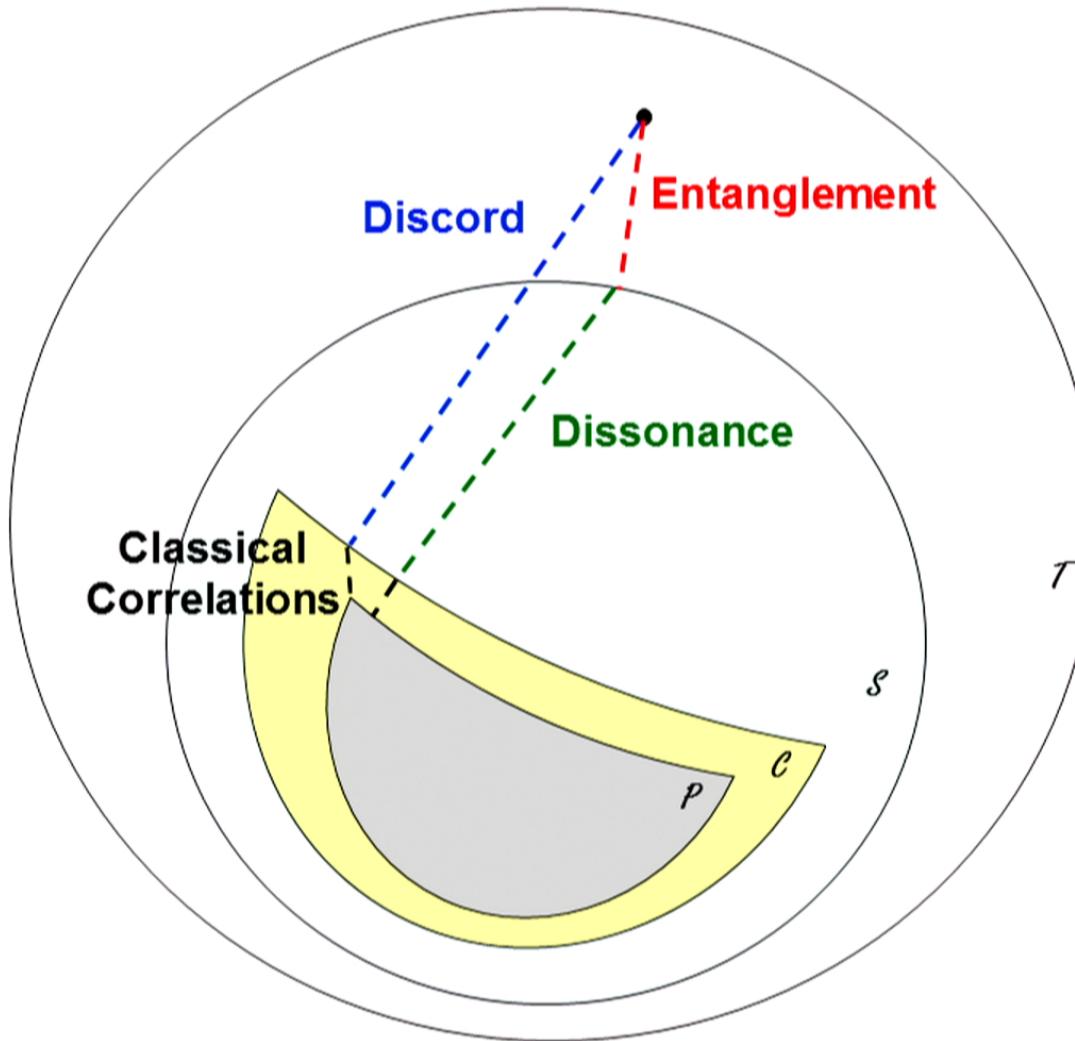
Discord → shortest distance between ρ and a **classical** state (χ')

Dissonance → shortest distance between σ' and a
classical state (χ'')

Modi *et. al.*, PRL, **104**, 080501 (2010).







Measuring the Distance

$$D(\rho \parallel \sigma)$$

Quantity should be zero when it is known that the correlation is absent.

Quantity should be invariant under local unitary operations.

Quantity cannot increase under
local general measurements + classical communications

$$S(\rho \parallel \sigma) = \text{Trace} \left\{ \rho \ln \frac{\rho}{\sigma} \right\}$$

**Von Neumann
Relative Entropy**

Vedral et. al., PRL, **78** (12), 2275 (1997).

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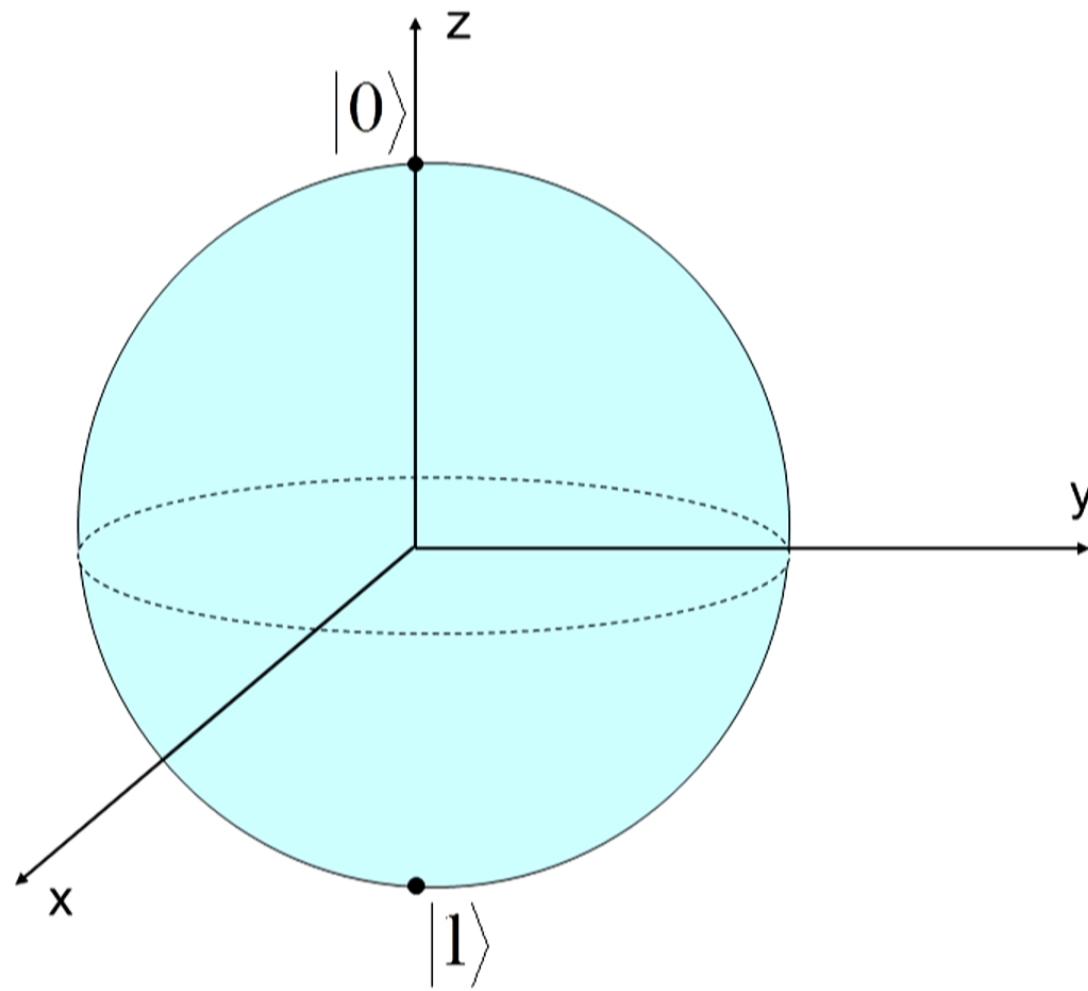
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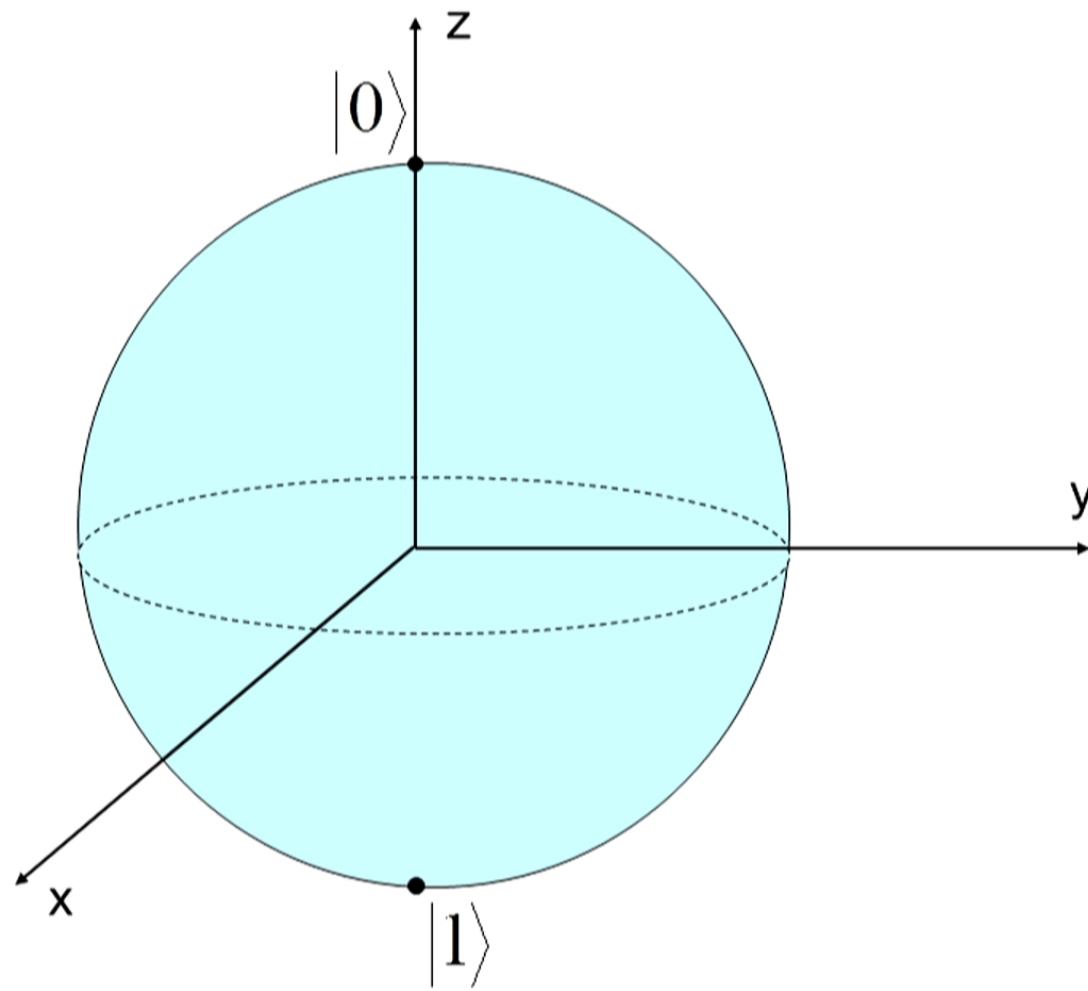
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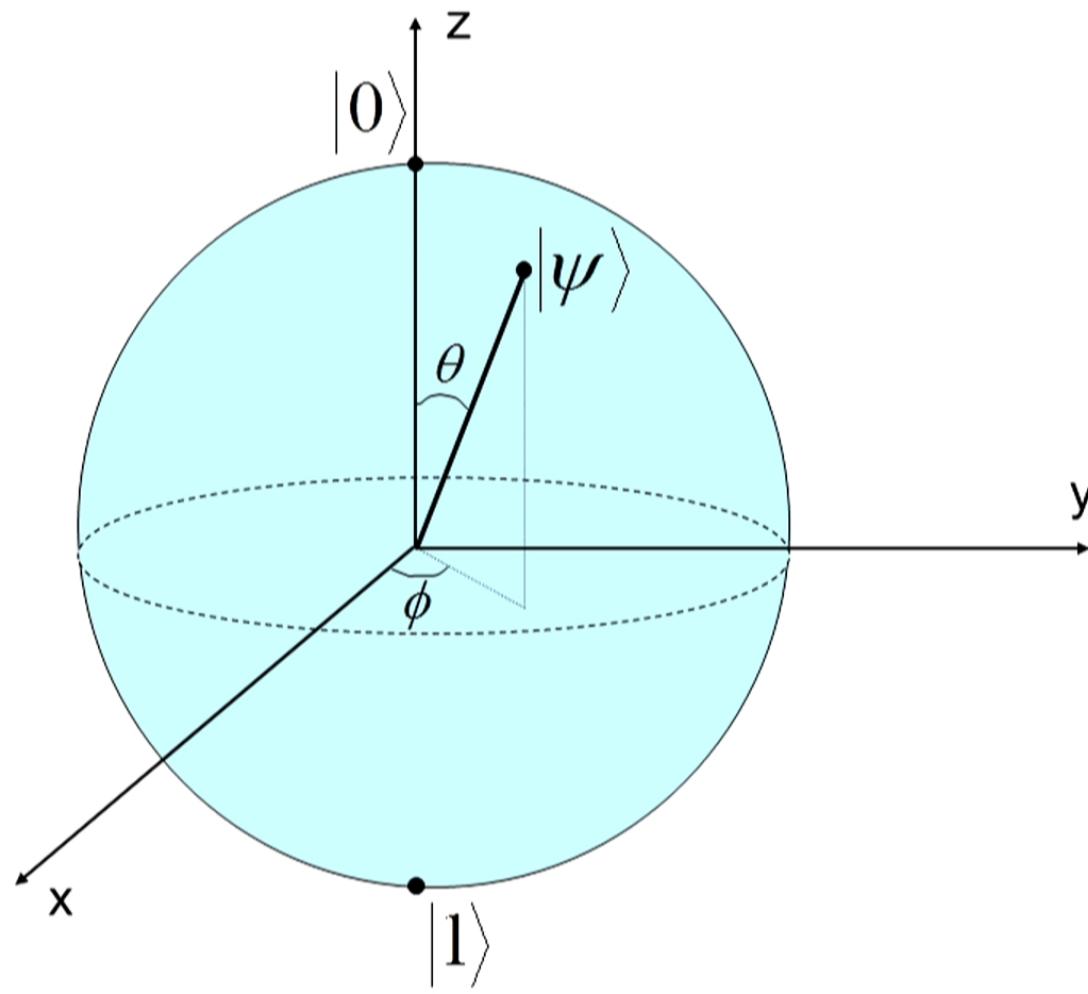
Bloch Sphere Representation of a Qubit



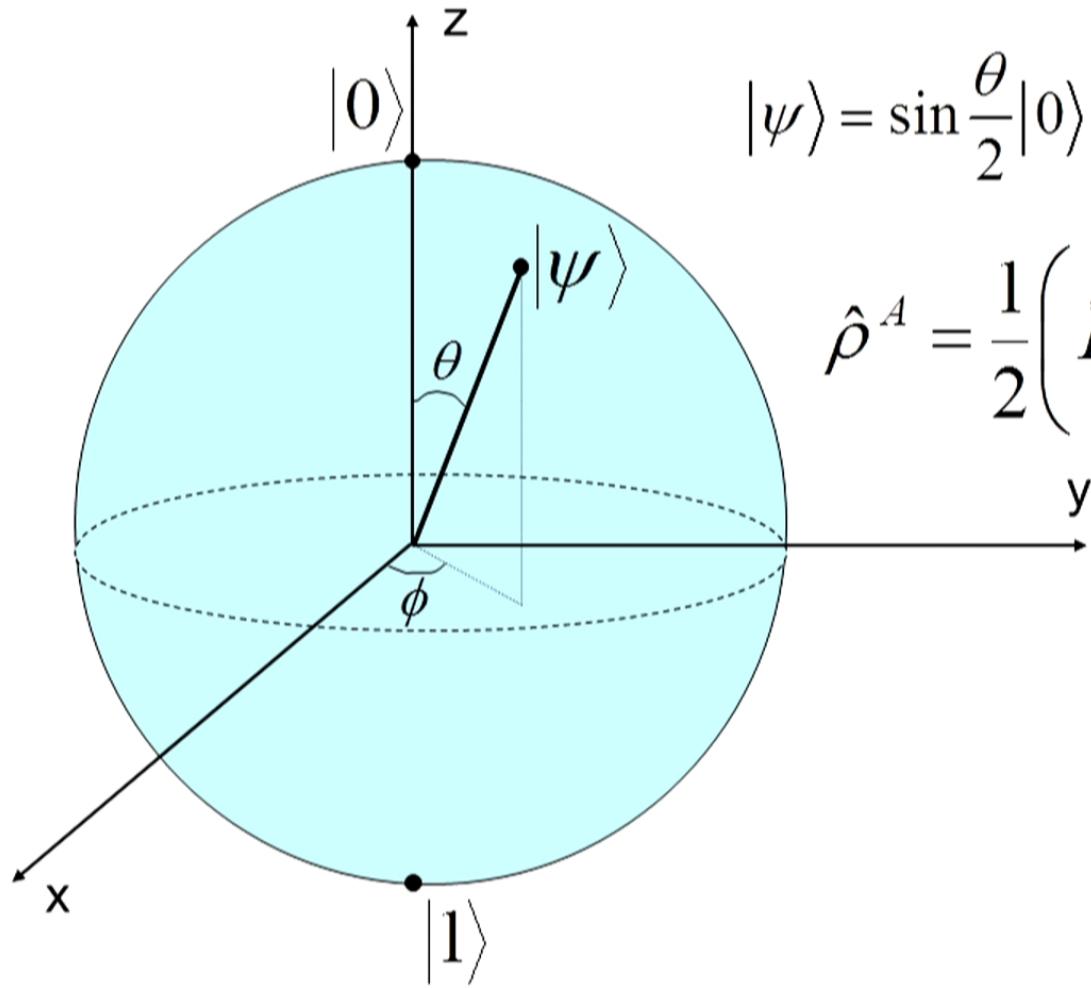
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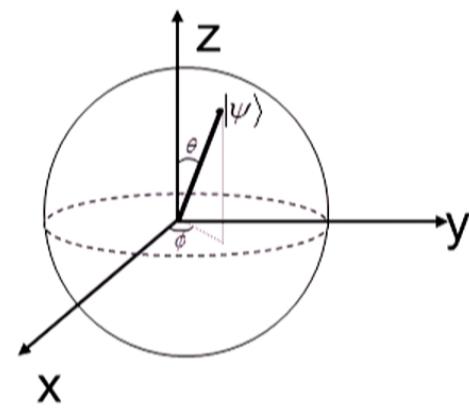
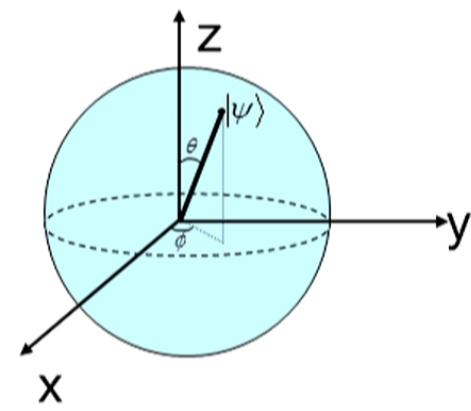


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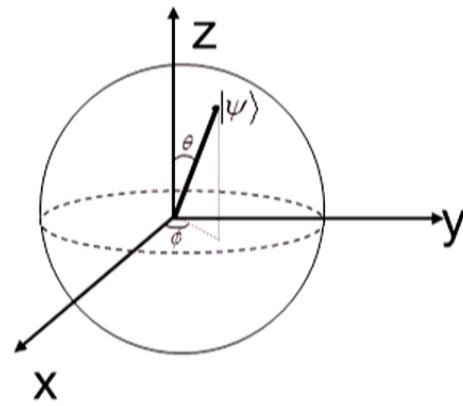
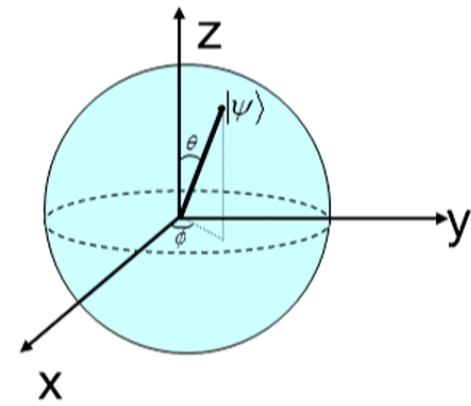


$$|\psi\rangle = \sin \frac{\theta}{2} |0\rangle + e^{i\phi} \cos \frac{\theta}{2} |1\rangle$$

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$$C_{ij} = \langle (\hat{\sigma}_i - A_i \hat{I}) \otimes (\hat{\sigma}_j - B_j \hat{I}) \rangle$$

Cross-Correlation Tensor

$$\begin{aligned}\hat{\rho}^{AB} = & \frac{1}{2} \left(\hat{I} + \sum_{i=1}^3 A_i \hat{\sigma}_i \right) \otimes \frac{1}{2} \left(\hat{I} + \sum_{j=1}^3 B_j \hat{\sigma}_j \right) \\ & + \frac{1}{4} \sum_{i,j=1}^3 C_{ij} \hat{\sigma}_i \otimes \hat{\sigma}_j\end{aligned}$$

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Determinant of C_{ij} tells us about the existence/absence of quantum correlations

- 1) When Entanglement > 0 , $\text{Det}(C_{ij}) < 0$

Al-Qasimi and James (In Preparation)

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- 2) When $\text{Det}(C_{ij}) < 0$, Discord > 0

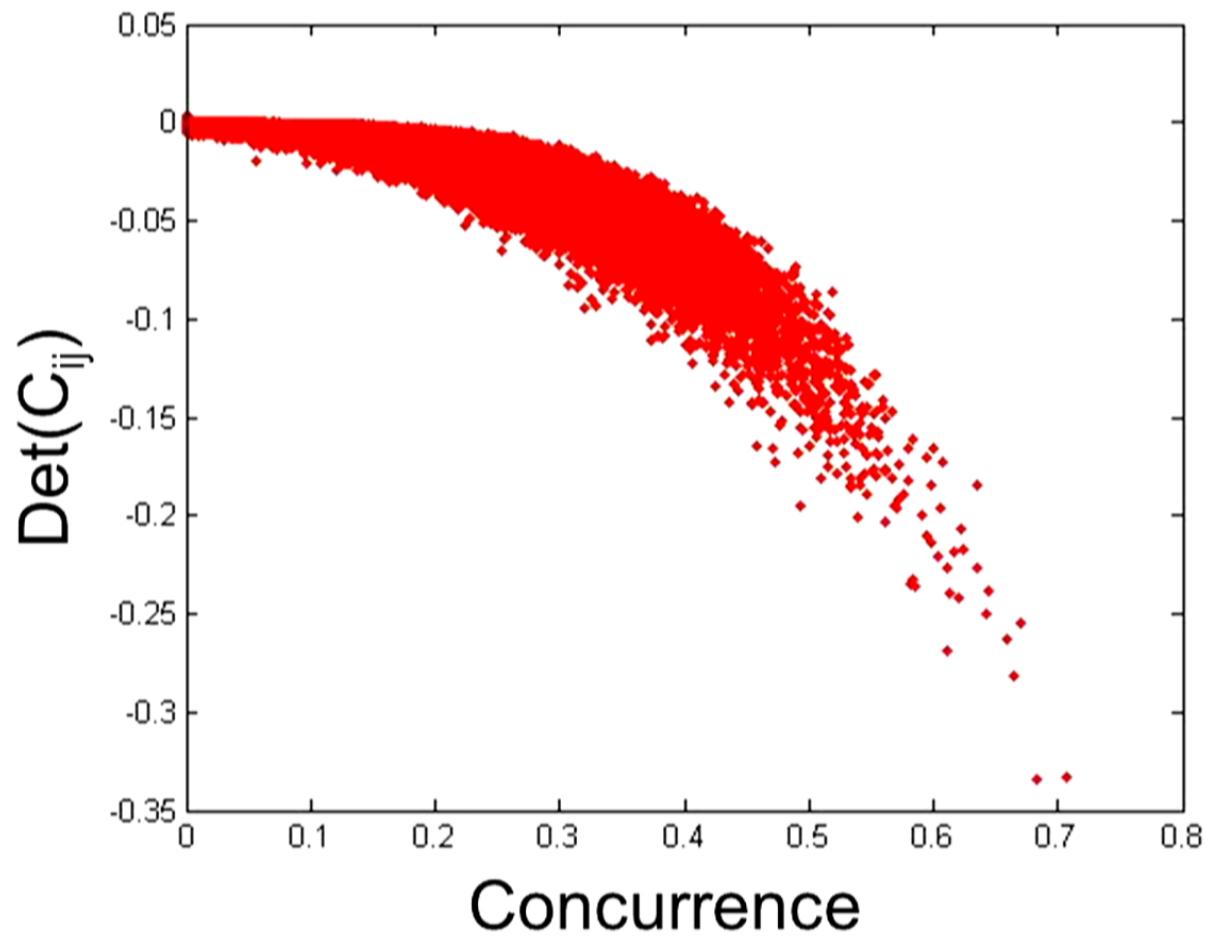
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States with $\text{Det}(C_{ij}) < 0 \rightarrow \text{Sinister States}$

Al-Qasimi and James (In Preparation)



Why Sinister?

$$C_{ij} = \sum_{k=1}^3 s_k a_i^{(k)} b_j^{(k)}$$

Singular Value
Decomposition

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If one of the vector sets is LH (and the other RH), $\text{Det}(C_{ij}) < 0$.

Sinister = Latin for *Left*

Conclusions

- Quantify Quantum Correlations (QC) by the amount of information required to create the state.
- Detect QC by performing a local operation
- Quantify QC by the fact that measurements disturb quantum systems.
- Quantify correlations by the distance between the state of interest and the closest state that lacks the quantity of interest.
- Quantify QC by the determinant of the Cross-Correlation Matrix
- Concurrence and Partial Transpose Criterion for Separability to study ESD
- Correlations in Atomic Spectroscopy