

Title: Entanglement and the Fermi surface

Date: Nov 27, 2012 03:30 PM

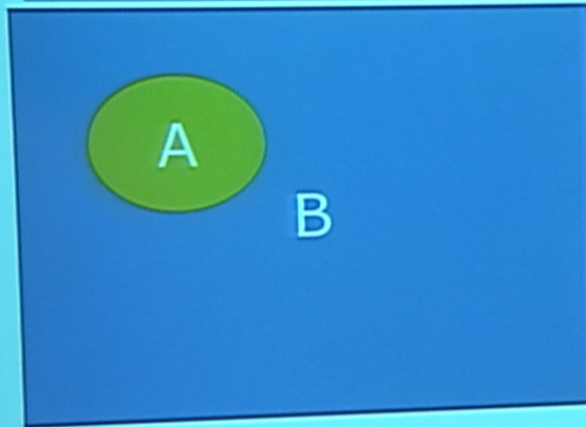
URL: <http://www.pirsa.org/12110082>

Abstract: In this talk I will describe my work characterizing quantum entanglement in systems with a Fermi surface. This class includes everything from Fermi liquids to exotic spin liquids in frustrated magnets and perhaps even holographic systems. I review my original scaling argument and then describe in detail a number of new precise results on entanglement in Fermi liquids. I will also discuss recent quantum Monte Carlo calculations of Renyi entropies and will argue that we now have a rather complete agreement between theory and numerics for Fermi liquid entanglement. I will also discuss universal crossovers between thermal and entanglement entropy and a class of solvable interacting models where we can prove the universality of the Widom formula for Fermi surface entanglement. If there is time, I will comment on several other topics including fluctuations of conserved quantities and connections to holography.

Acknowledgements

- Collaborators: Liza Huijse, Subir Sachdev, Senthil, Jeremy McMinis, Norm Tubman
- Funding from the Simons Foundation

Entanglement



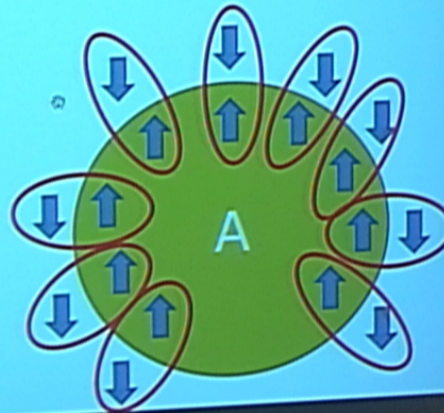
$$\begin{pmatrix} \downarrow \\ \uparrow \end{pmatrix} = |\uparrow_A \downarrow_B\rangle - |\downarrow_A \uparrow_B\rangle$$

$$S_n(A) = S_n(B) = \ln 2$$

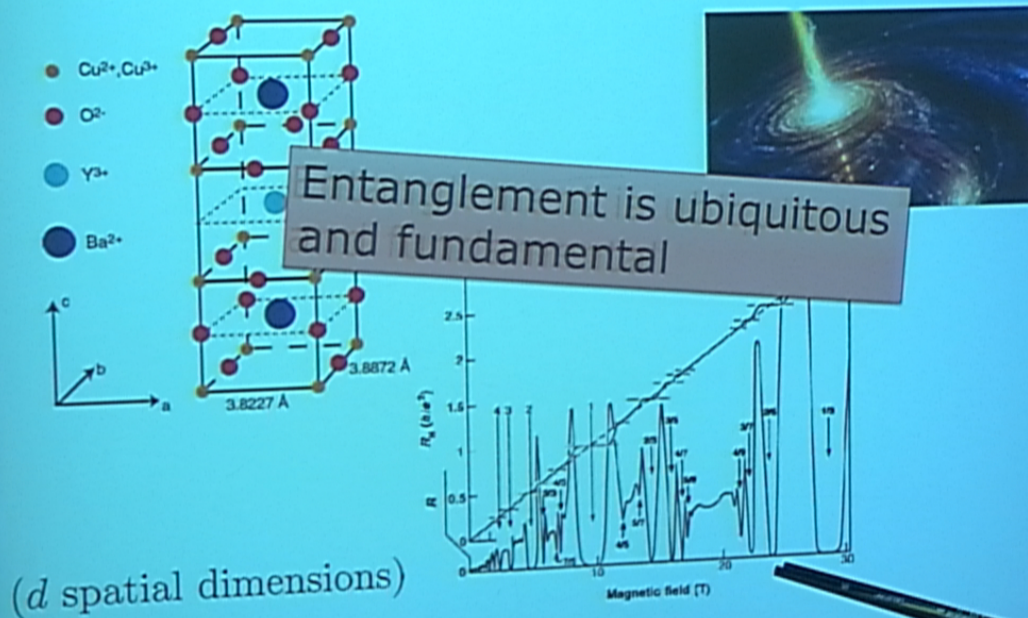
Renyi entropy:

$$S_n = \frac{1}{1-n} \ln (\text{tr}(\rho_A^n))$$

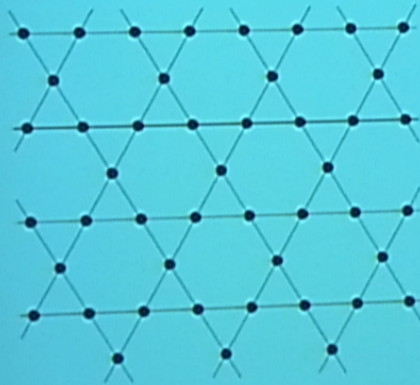
$$S(A) = \lim_{n \rightarrow 1} S_n(A)$$



Quantum Matter

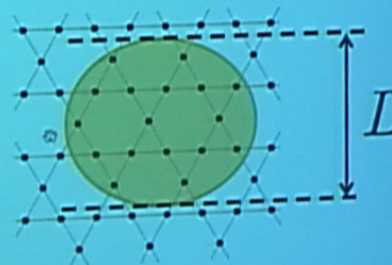


Example: Kagome



$$H = J \sum_{r,r'} \vec{S}_r \cdot \vec{S}_{r'}$$

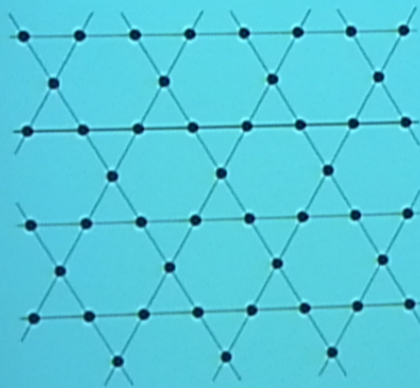
Spin liquid?



Topological entanglement entropy: $S = \alpha L - \beta$

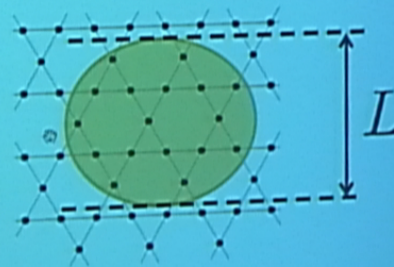
(TEE Kitaev-Preskill, Levin-Wen;
DMRG Yan-White-Huse, Jiang-Wang-Balents, ...)

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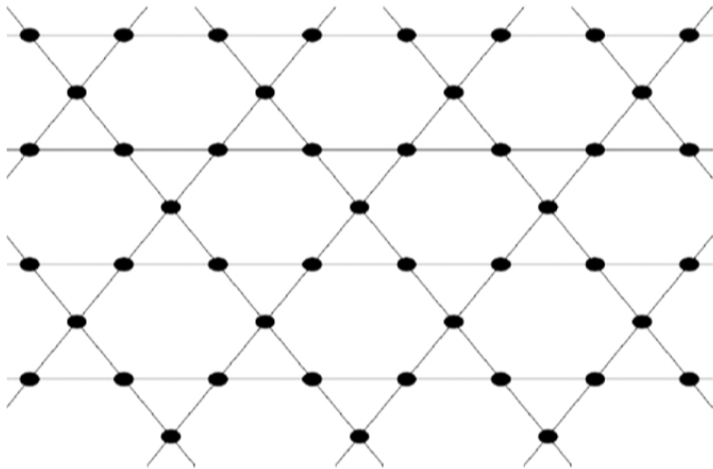
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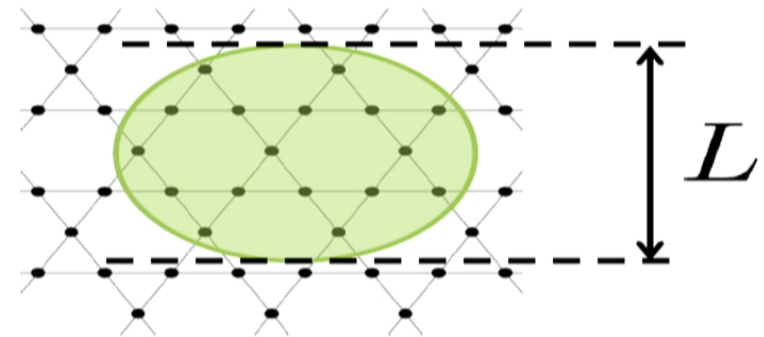
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Topological entanglement entropy: $S = \alpha L - \beta$
*(TEE Kitaev-Preskill, Levin-Wen;
 DMRG Yan-White-Huse, Jiang-Wang-Balents, ...)* $\beta = \ln(2)$

Measuring entanglement

- Possible, in principle, to measure integer Renyi entropy S_n , $n > 1$
- Fine print: requires a delicate quantum coherent process involving multiple copies of the system, signal is exponentially small in S_n
- **IMPORTANT: classical numerical algorithms also possible, “swap trick”**
(QMC Kallin-Isakov-Inglis-Melko-Hastings-...)

Why Fermi liquids?

- Interesting and ubiquitous phase of matter
- Highly entangled, tells us about the limits of entanglement based simulation methods
- Interacting yet tractable
- Contact with more exotic systems, e.g. organic salts and holographic systems?
- Wonderful test bed!

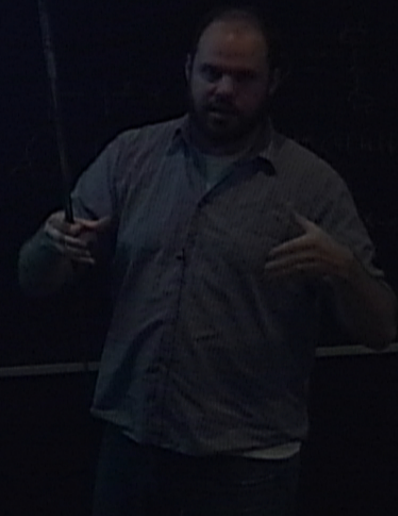
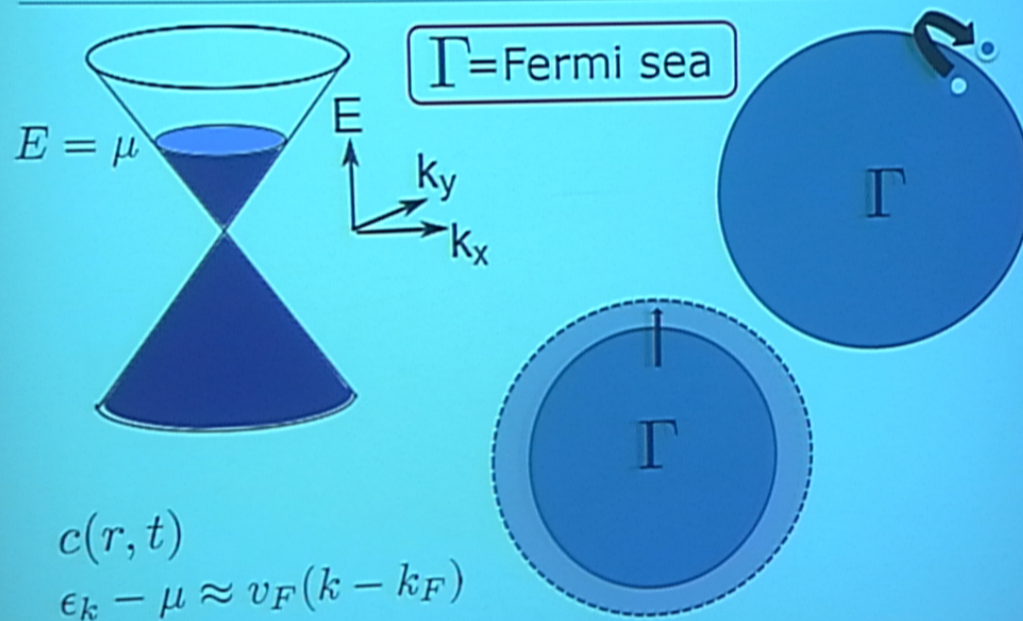
Main question

Entanglement structure of a
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Fermi surfaces



Widom formula

- Summing then gives the Widom formula

$$S_n(A) = \left(1 + \frac{1}{n}\right) \frac{1}{24} \frac{1}{(2\pi)^{d-1}} \int_{\partial A} \int_{\partial \Gamma} |n_x \cdot n_k| \ln(L)$$

- Conjectured ($n=1$) by Klich and Gioev, subsequently justified using one-dimensional mode counting (Klich-Gioev '06, BGS '09)
- Later additions include Renyi index, non-convex regions, thermal crossovers, mutual information, interactions, ...

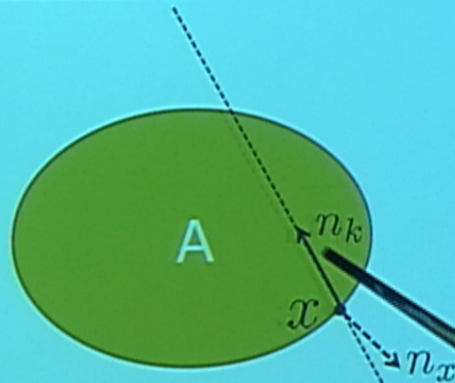
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(BGS '10)

Effective length

$$x \in \partial A$$

$$k \in \partial \Gamma$$



$$L_{\text{eff}}(x, k) = \text{length}(\{x + n_k s \mid s \in (-\infty, \infty)\} \cap A)$$

Important later for thermal crossovers and
subleading terms

(BGS '10)

Widom formula

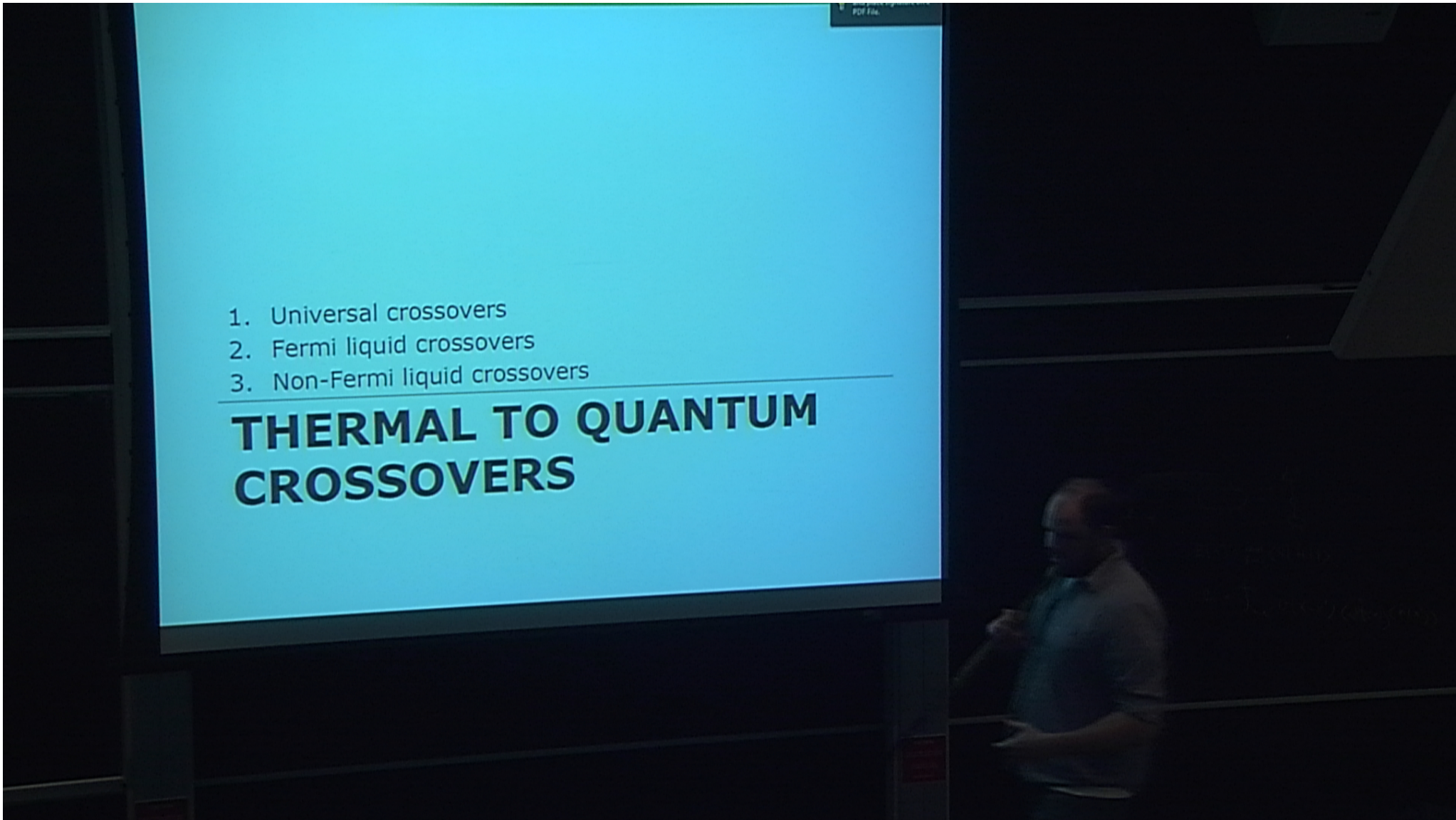
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(*BGS '11*)

- 
1. Universal crossovers
 2. Fermi liquid crossovers
 3. Non-Fermi liquid crossovers

THERMAL TO QUANTUM CROSSOVERS

Entropy crossovers, $d=1$

- Interval, length L , temperature T :

$$S(L, T) = \frac{c}{3} \ln \left(\frac{v}{\pi T \epsilon} \sinh \left(\frac{\pi T L}{v} \right) \right)$$

- Matches proposed crossover form up to "area law respecting" terms

$$S(A, T) \sim T^\phi f(T^{1/z} L)$$

- Similar crossover for $z \neq 1$ leads to conjecture: "reasonable" gapless 1d systems cannot be more than $\log(L)$ entangled (Senthil-BGS '11)

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Fermi liquid entanglement

- Quasiparticles: renormalize velocity

$$v_F \rightarrow v_F^*$$

$$S(A, T) = \frac{1}{12} \frac{1}{(2\pi)^{d-1}} \int_{\partial A} \int_{FS} |n_x \cdot n_k| \ln \left(\frac{v_F^*}{\pi T c} \sinh \left(\frac{\pi T L_{\text{ent}}(x, k)}{v_F^*} \right) \right)$$

- Correctly reproduces thermal entropy, supports claim of universality for Fermi liquids (BGS '10)

Field theory interpretation

$$d = 1 \quad c(x, t) = e^{ik_F x} \psi_R(x, t) + e^{-ik_F x} \psi_L(x, t)$$

Replica trick
 $\text{tr}(\rho_A^n) = e^{-(n-1)S_n}$

(Cardy-Calabrese '10,
BGS-McMinis-Tubman '12)

$$c^\dagger c \sim e^{i2k_F x} \psi_R \psi_L^\dagger + \dots$$

Branch points (impurities)
"bind" relevant operators

$$w = \left(\frac{z - z_1}{z - z_2} \right)^{1/n}$$

Conformal
transformation

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Ground state/correlations

- Ground state is simple in decoupled variables $|G\rangle = |f\rangle|\tilde{\tau}\rangle$
- BUT, f and $\tilde{\tau}$ are non-locally related to c and σ
- Nevertheless, many correlators are simple

$$\langle c(r, t) c^\dagger(0, 0) \rangle = \langle f(r, t) f^\dagger(0, 0) \rangle \langle \tau^x(r, t) \tau^x(0, 0) \rangle$$
$$\rightarrow Z = \langle \tau^x \rangle^2$$