

Title: Condensed Matter - Lecture 14

Date: Oct 26, 2012 10:30 AM

URL: <http://pirsa.org/12100039>

Abstract:

$$\omega = JS_z$$

$\frac{1}{2} \hbar$



$$\text{FM: } \omega = JS_z (1 - \cos ka)$$

$$\text{AFM: } \omega = JS_z \sqrt{1 - \cos ka}$$

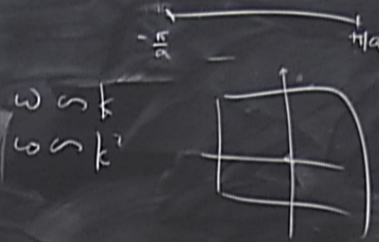




$$\text{FM: } \omega = JS_z (1 - \cos ka)$$

$$\text{AFM: } \omega = JS_z \sqrt{1 - \cos ka}$$

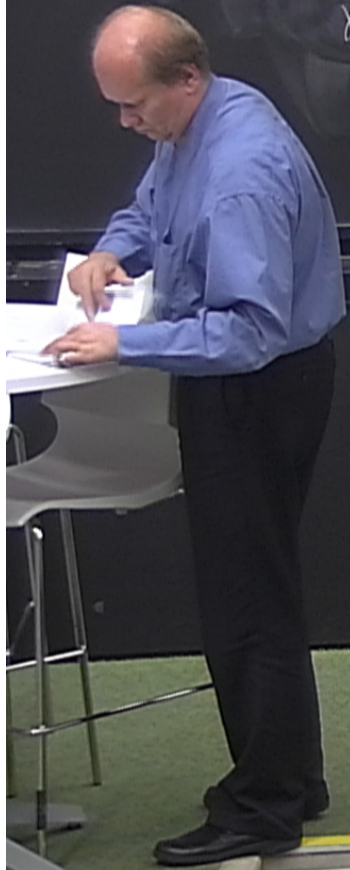
$$\gamma_k =$$



$$\omega = \sqrt{\frac{2C}{m}} \sqrt{1 - \cos ka}$$

$$k \ll \omega \ll k$$

$$v = \frac{\partial \omega}{\partial k} = \sqrt{\frac{2C}{m}} \frac{\sin ka \cdot a}{2\sqrt{1 - \cos ka}}$$



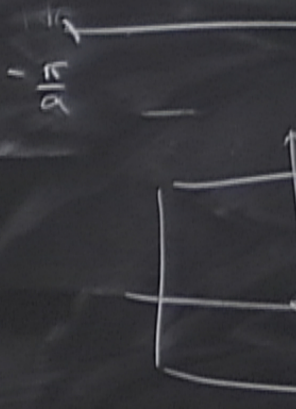


$$\text{FM: } \omega = JS_z (1 - \cos ka)$$

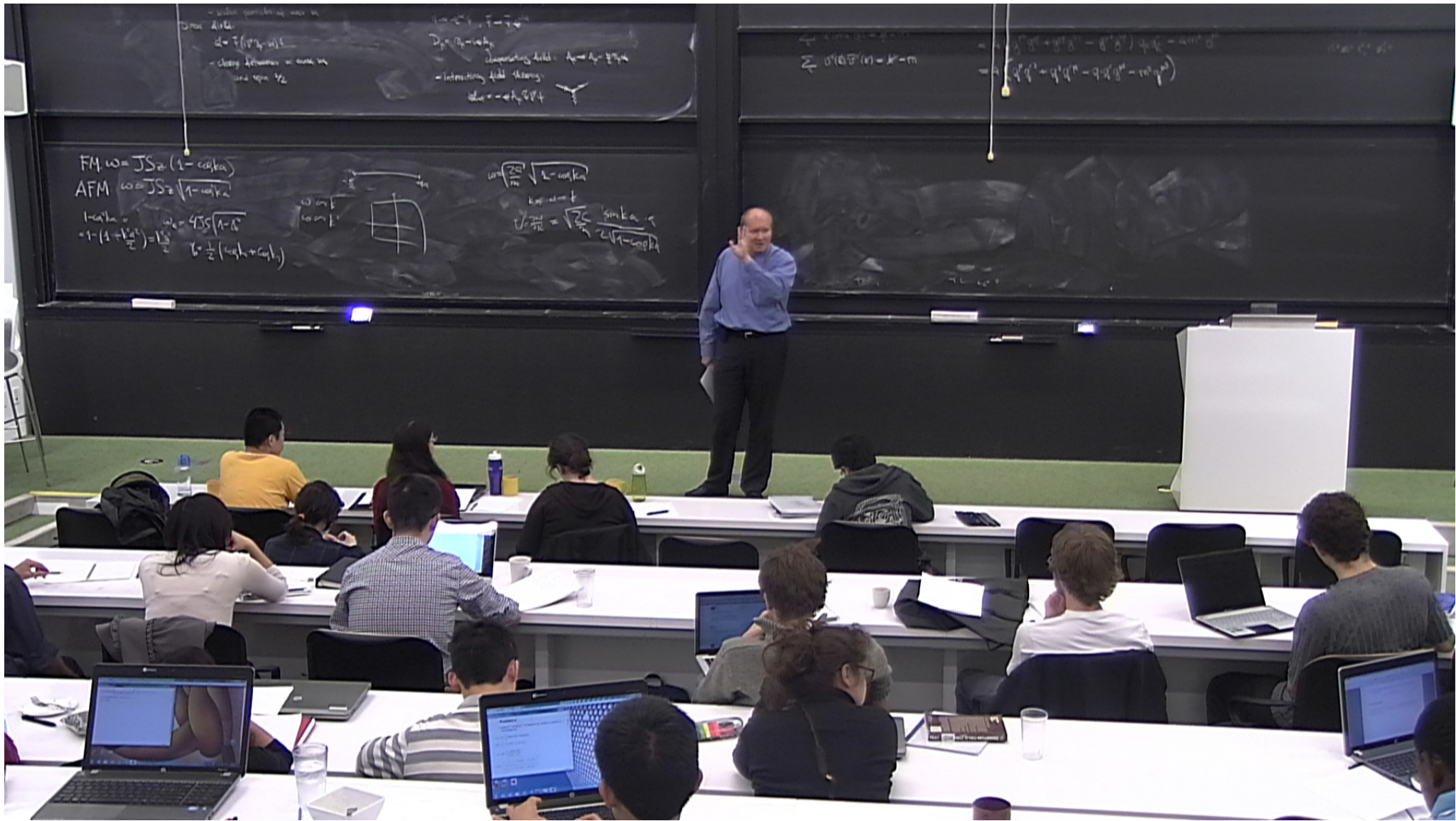
$$\text{AFM: } \omega = JS_z \sqrt{1 - \cos^2 ka}$$

$$\begin{aligned} 1 - \cos^2 ka &= \omega_k = 4JS \sqrt{1 - \gamma_k^2} \\ &= 1 - \left(1 - \frac{k^2 a^2}{2}\right) = \frac{k^2 a^2}{2} \quad \gamma_k = \frac{1}{2} (\cos k_x + \cos k_y) \end{aligned}$$

$\omega \propto k$   
 $\omega \propto k^2$









$$i\hbar \frac{\partial \varphi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \varphi - \mu \varphi + g |\varphi|^2 \varphi$$





$$i\hbar \frac{\partial \varphi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \varphi - \mu \varphi + g |\varphi|^2 \varphi$$

$$\varphi_0 = \sqrt{\rho_0} e^{i\phi}$$

$$\varphi(\vec{x}, t) = \varphi_0 + u e^{i(\omega t - \vec{k} \cdot \vec{x})}$$

$(\omega t - \vec{k} \cdot \vec{x})$



$$i\hbar \frac{\partial \varphi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \varphi - \mu \varphi + g |\varphi|^2 \varphi$$

$$\varphi_0 = \sqrt{\rho_0} e^{i\phi}$$

$$\varphi(\vec{x}, t) = \varphi_0 + u e^{\frac{i(\omega t - \vec{k} \cdot \vec{x})}{\hbar}} + u^* e^{-\frac{i(\omega t - \vec{k} \cdot \vec{x})}{\hbar}}$$

$$\omega(\vec{k}) = \sqrt{\frac{\hbar^2 k^2}{2m} \left( \frac{\hbar^2 k^2}{2m} + 2g\rho_0 \right)} \quad \frac{\hbar^2}{2m} = 1$$

For small  $u$   $\omega(\vec{k})$



$$i\hbar \frac{\partial \varphi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \varphi - \mu \varphi + g |\varphi|^2 \varphi$$

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For small  $k$ ,  $\omega(\vec{k}) \propto |\vec{k}|$  phonons

For large  $k$ ,  $\omega(\vec{k}) \approx \frac{\hbar k^2}{2m}$



$$i\hbar \frac{\partial \varphi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \varphi - \mu \varphi + g |\varphi|^2 \varphi$$

$$\varphi = \sqrt{\rho_0} e^{i\phi}$$

$$\varphi(\vec{x}, t) = \sqrt{\rho_0} e^{i(\omega t - \vec{k} \cdot \vec{x})} + u^* e^{-i(\omega t - \vec{k} \cdot \vec{x})}$$

$$\omega(\vec{k}) = \sqrt{\frac{\hbar^2 k^2}{2m} \left( \frac{\hbar^2 k^2}{2m} + 2g\rho_0 \right)} \quad \hbar=1$$

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$$i\hbar \frac{\partial \varphi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \varphi - \mu \varphi + g |\varphi|^2 \varphi$$

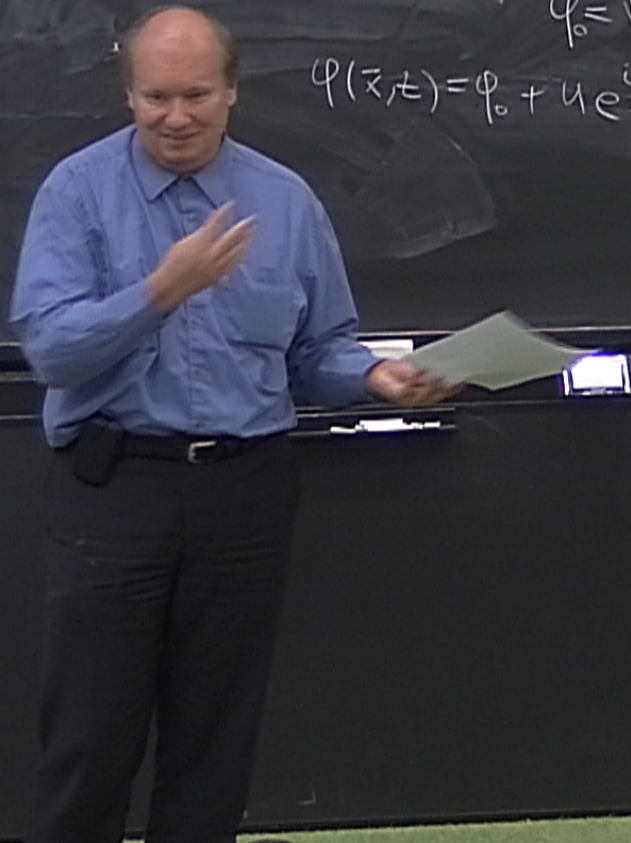
$$\varphi_0 = \sqrt{\rho_0} e^{i\phi}$$

$$\varphi(\vec{x}, t) = \varphi_0 + u e^{\frac{i(\omega t - \vec{k} \cdot \vec{x})}{\hbar}} + u^* e^{-\frac{i(\omega t - \vec{k} \cdot \vec{x})}{\hbar}}$$

$$\omega(\vec{k}) = \sqrt{\frac{\hbar^2 k^2}{2m} \left( \frac{\hbar^2 k^2}{2m} + 2g\rho_0 \right)} \quad \frac{\hbar^2}{2m} = 1$$

For small  $k$ ,  $\omega(\vec{k}) \propto |\vec{k}|$  phonons

For large  $k$ ,  $\omega(\vec{k}) \approx \frac{\hbar k^2}{2m}$





$$i\hbar \frac{\partial \varphi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \varphi - \mu \varphi + g |\varphi|^2 \varphi$$

$$\varphi_0 = \sqrt{\rho_0} e^{i\phi}$$

$$\varphi(\vec{x}, t) = \varphi_0 + u e^{i(\omega t - \vec{k} \cdot \vec{x})} + u^* e^{-i(\omega t - \vec{k} \cdot \vec{x})}$$

$$\omega(\vec{k}) = \sqrt{\frac{\hbar^2 k^2}{2m} \left( \frac{\hbar^2 k^2}{2m} + 2g\rho_0 \right)} \quad \frac{\hbar^2}{2m} = 1$$

For  $k$ ,  $\omega(\vec{k}) \propto |\vec{k}|$  phonons

For  $k$ ,  $\omega(\vec{k}) \approx \frac{\hbar^2 k^2}{2m}$



$$i\hbar \frac{\partial \varphi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \varphi - \mu \varphi + g |\varphi|^2 \varphi$$

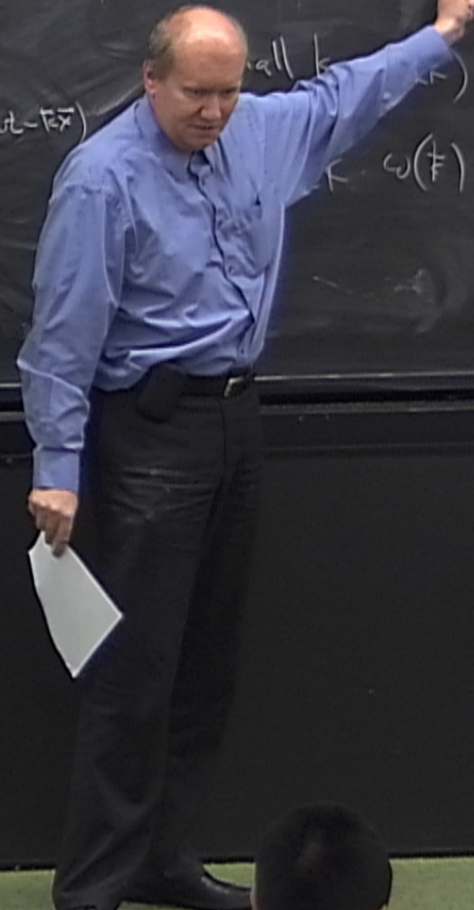
$$\varphi_0 = \sqrt{\rho_0} e^{i\phi}$$

$$\varphi(\vec{x}, t) = \varphi_0 + u e^{i(\omega t - \vec{k} \cdot \vec{x})} + u^* e^{-i(\omega t - \vec{k} \cdot \vec{x})}$$

$$\omega(\vec{k}) = \sqrt{\frac{\hbar^2 k^2}{2m} \left( \frac{\hbar^2 k^2}{2m} + 2g\rho_0 \right)} \quad \frac{\hbar^2}{2m} = 1$$

all  $k$   $\omega(\vec{k}) \propto |\vec{k}|$  phonons

$$\omega(\vec{k}) \approx \frac{\hbar k^2}{2m}$$





$$i\hbar \frac{\partial \varphi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \varphi - \mu \varphi + g |\varphi|^2 \varphi$$

$$\varphi_0 = e^{i\phi}$$

$$\varphi(\vec{x}, t) = \varphi_0 + u e^{-i\omega t - i\vec{k}\vec{x}} + u^* e^{-i(\omega t - \vec{k}\vec{x})}$$

$$\omega(\vec{k}) = \sqrt{\frac{\hbar^2 k^2}{2m} \left( \frac{\hbar^2 k^2}{2m} + 2g\varphi_0 \right)} \quad \frac{\hbar^2 k^2}{2m} \ll 2g\varphi_0$$

For small  $k$ ,  $\omega(\vec{k}) \propto |\vec{k}|$  phonons

For large  $k$ ,  $\omega(\vec{k}) \approx \frac{\hbar^2 k^2}{2m}$

$$\frac{\hbar^2 k^2}{2m} \gg 2g\varphi_0$$



$$i\hbar \frac{\partial \varphi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \varphi - \mu \varphi + g |\varphi|^2 \varphi$$

$$\varphi_0 = \sqrt{\frac{\mu}{g}} e^{i\phi}$$

$$\varphi(\vec{x}, t) = \varphi_0 + \sum_k a_k e^{-i(\omega_k t - \vec{k} \cdot \vec{x})}$$

$$\omega(k) = \sqrt{\frac{\hbar^2 k^2}{2m} \left( \frac{\hbar^2 k^2}{2m} + 2g\varphi_0 \right)} \quad \frac{\hbar^2}{2m} = 1$$

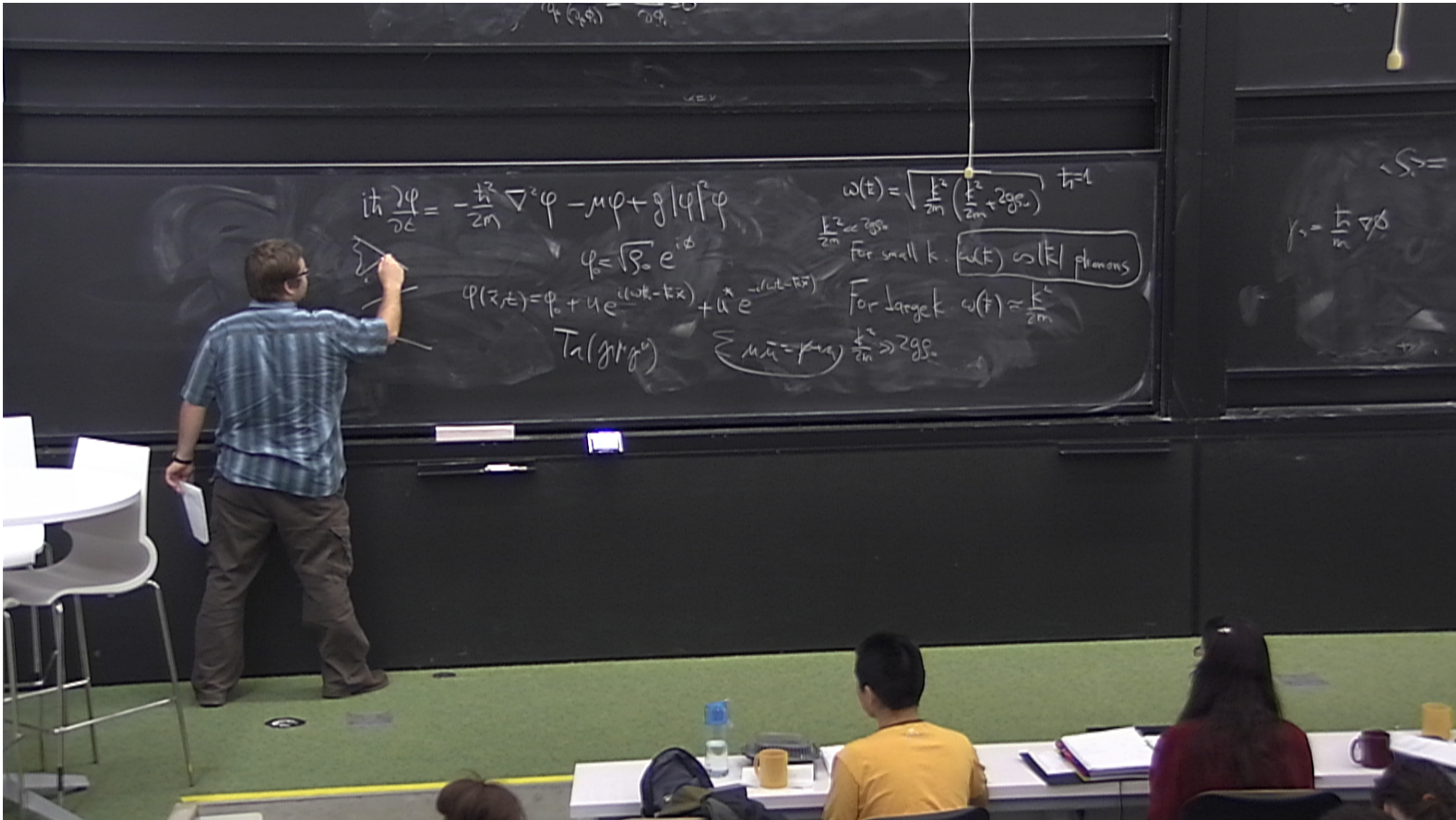
$$\frac{\hbar^2 k^2}{2m} \ll 2g\varphi_0$$

For small  $k$ .  $\omega(k) \propto |k|$  phonons

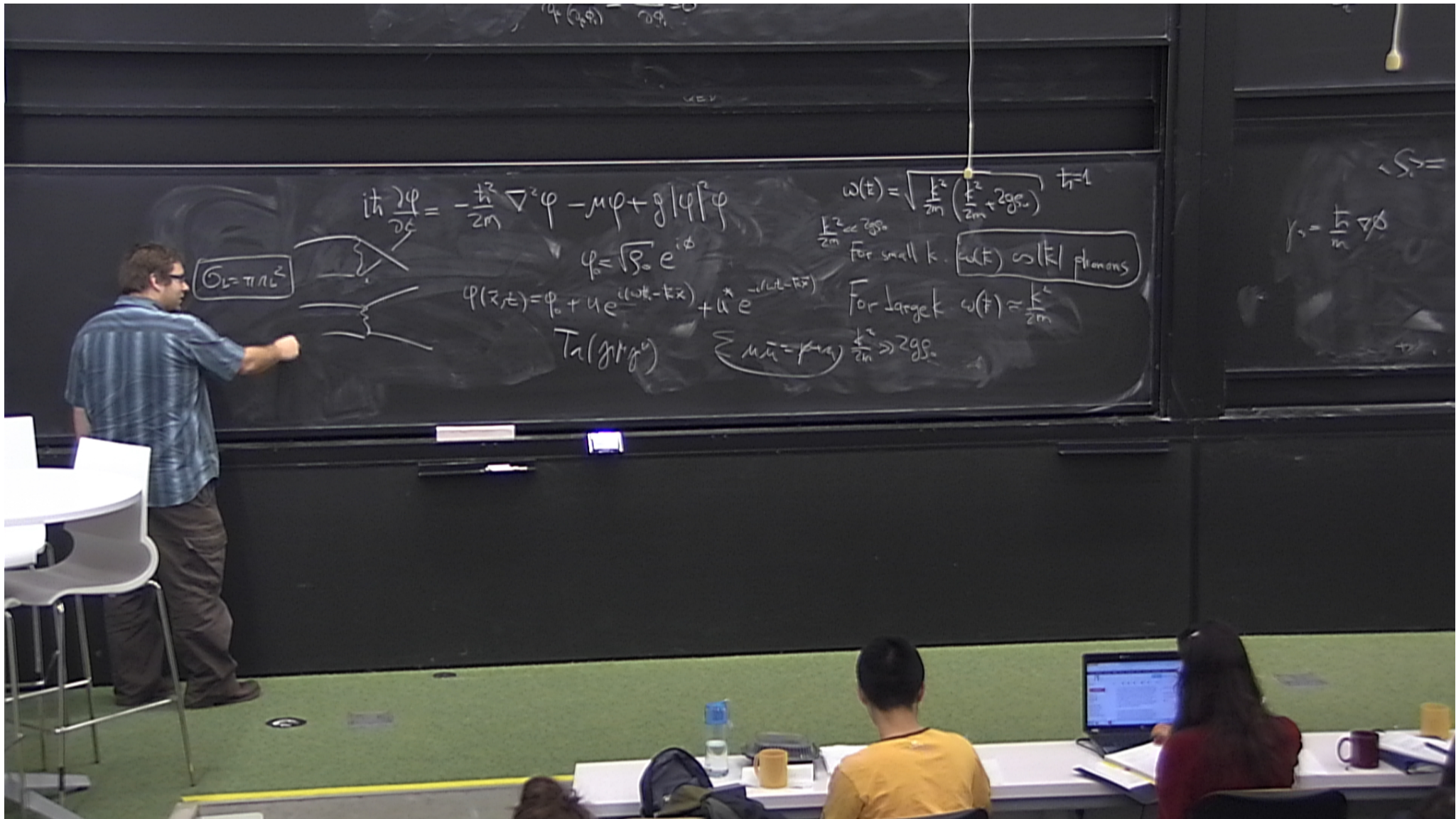
For large  $k$ .  $\omega(k) \approx \frac{\hbar^2 k^2}{2m}$

$$\frac{\hbar^2 k^2}{2m} \gg 2g\varphi_0$$





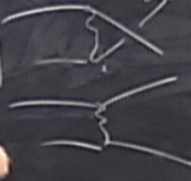




$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \psi - m\phi + g|\psi|^2 \psi$$

$$\omega(k) = \sqrt{\frac{\hbar^2 k^2}{2m} \left( \frac{\hbar^2 k^2}{2m} + 2g\psi_0 \right)} \quad \hbar^{-1}$$

$$\Omega = \pi \Lambda^2$$



$$\psi_0 = \sqrt{\psi_0} e^{i\phi}$$

$$\psi(z,t) = \psi_0 + u e^{i(\omega t - kx)} + u^* e^{-i(\omega t - kx)}$$

For small  $k$ ,  $\omega(k) \approx \omega(k=0) + v k$  phonons

For large  $k$ ,  $\omega(k) \approx \frac{\hbar^2 k^2}{2m}$

$$T_A(g\mu\psi_0)$$

$$\sum u u^* = \mu \psi_0 \quad \frac{\hbar^2 k^2}{2m} \gg 2g\psi_0$$

$$\psi_0 = \frac{\hbar^2}{m} \psi_0$$



$$i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \psi - \mu \psi + g |\psi|^2 \psi$$

$$\psi_0 = \sqrt{\rho_0} e^{i\phi}$$

$$\psi(z,t) = \psi_0 + u e^{i(\omega t - kx)} + u^* e^{-i(\omega t - kx)}$$

$$T_0(g\mu\rho_0)$$

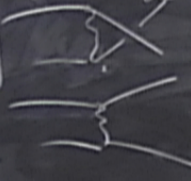
$$\sum u u^* = \mu \rho_0 \quad \frac{\hbar^2 k^2}{2m} \gg 2g\rho_0$$

$$\omega(t) = \sqrt{\frac{\hbar^2 k^2}{2m} \left( \frac{\hbar^2 k^2}{2m} + 2g\rho_0 \right)} \quad \hbar^{-1}$$

For small  $k$ :  $\omega(t) \approx \omega(k)$  phonons

For large  $k$ :  $\omega(t) \approx \frac{\hbar^2 k^2}{2m}$

$$\Theta = \pi \hbar^2$$

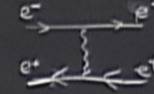


$$\psi = \frac{\hbar}{m} \nabla \phi$$



$$E \approx m_\mu:$$

$$\sigma \sim \frac{\alpha^2}{m_\mu^2} \sqrt{1 - \frac{m_\mu}{m_e}}$$



QFT

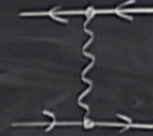
$$[a_p, a_k^\dagger] = (2\pi)^3 \delta(\vec{p} - \vec{k})$$

Bosons/Fermions

Fock space:

$$|\vec{p}_1 \dots \vec{p}_n\rangle = a_{p_1}^\dagger \dots a_{p_n}^\dagger |0\rangle$$

Forces  $\Leftrightarrow$

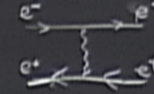


$$S =$$



$$E \approx m_\mu:$$

$$\sigma \sim \frac{\alpha^2}{v} \sqrt{1 - \frac{m_\mu}{E}}$$



QFT

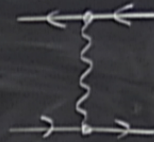
$$[a_p, a_k^\dagger] = (2\pi)^3 \delta(\vec{p} - \vec{k})$$

bosons/fermions

Fock space.

$$|\vec{p}_1 \dots \vec{p}_n\rangle = a_{p_1}^\dagger \dots a_{p_n}^\dagger |0\rangle$$

Forces  $\Leftrightarrow$



$$S = \int d^4x \mathcal{L}(\partial_\mu \phi_i, \phi_i)$$

↑ Lagrangian density.



$$E \approx m_\mu:$$

$$\sigma \sim \frac{\alpha^2}{v} \sqrt{1 - \frac{m_\mu}{E}}$$



QFT

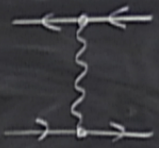
$$[a_p, a_k^\dagger] = (2\pi)^3 \delta(\vec{p} - \vec{k})$$

bosons/fermions

Fock space:

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Forces  $\Leftrightarrow$



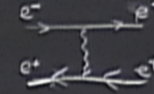
$$S = \int d^4x \mathcal{L}(\partial_\mu \phi_i, \phi_i)$$

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QFT

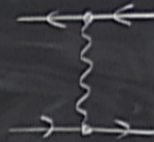
$$[a_p, a_k^\dagger] = (2\pi)^3 \delta(\vec{p} - \vec{k})$$

Bosons / fermions

Fock space

$$|\vec{p}_1 \dots \vec{p}_n\rangle \quad |a_{p_n}^\dagger |0\rangle$$

Forces  $\Leftrightarrow$



$$S = \int d^4x \mathcal{L}(\partial_\mu \phi_i, \phi_i)$$

Lagrangian density

$$\phi_i(\vec{x}, t) \equiv \phi_{i,\vec{x}}(t)$$



QFT

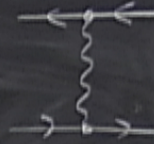
$$[a_p, a_k^\dagger] = (2\pi)^3 \delta(\vec{p}-\vec{k})$$

bosons/fermions

Fock space.

$$|\vec{p}_1 \dots \vec{p}_n\rangle = a_{p_1}^\dagger \dots a_{p_n}^\dagger |0\rangle$$

Forces  $\Leftrightarrow$



$$S = \int d^4x \mathcal{L}(\partial_\mu \phi_i, \phi_i)$$

Lagrangian density.

$$\phi_i(\vec{x}, t) \equiv \phi_{i,\vec{x}}(t) \leftarrow \text{basic dynamical variables.}$$





QFT

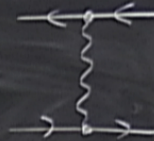
$$[a_p, a_k^\dagger] = (2\pi)^3 \delta(\vec{p} - \vec{k})$$

bosons/fermions

Fock space:

$$|\vec{p}_1 \dots \vec{p}_n\rangle = a_{p_1}^\dagger \dots a_{p_n}^\dagger |0\rangle$$

Forces  $\Leftrightarrow$



$$S = \int d^4x \mathcal{L}(\partial_\mu \phi_i, \phi_i)$$

Lagrangian density.

$\phi_i(\vec{x}, t) \equiv \phi_{i\vec{x}}(t)$  ← basic dynamical variables.

of motion:

$$\partial_\mu \frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi_i)} - \frac{\partial \mathcal{L}}{\partial \phi_i} = 0$$



QFT

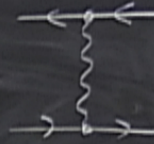
$$[a_p, a_k^\dagger] = (2\pi)^3 \delta(\vec{p}-\vec{k})$$

bosons/fermions

Fock space:

$$|\vec{p}_1 \dots \vec{p}_n\rangle = a_{p_1}^\dagger \dots a_{p_n}^\dagger |0\rangle$$

Forces  $\Leftrightarrow$



$$S = \int d^4x \mathcal{L}(\partial_\mu \phi_i, \phi_i)$$

Lagrangian density.

$\phi_i(\vec{x}, t) \equiv \phi_{i,\vec{x}}(t)$  ← basic dynamical variables.

eq of motion:

$$\partial_\mu \frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi_i)} - \frac{\partial \mathcal{L}}{\partial \phi_i} = 0$$



$$\sigma \sim \frac{\alpha^2}{m_e^2} \sqrt{1 - \frac{m_e}{m_\mu}}$$

QFT

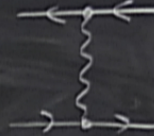
$$[a_p, a_k^\dagger] = (2\pi)^3 \delta(\vec{p} - \vec{k})$$

Bosons/Fermions

Fock space:

$$|\vec{p}_1 \dots \vec{p}_n\rangle = a_{p_1}^\dagger \dots a_{p_n}^\dagger |0\rangle$$

Forces  $\longleftrightarrow$



$$S = \int d^4x \mathcal{L}(\partial_\mu \phi_i, \phi_i)$$

Lagrangian

$$\phi_i(\vec{x}, t) \equiv \phi_{i\vec{x}}(t)$$

variables.

Eq. of motion:

$$\partial_\mu \left( \frac{\partial \mathcal{L}}{\partial (\partial_\mu \phi_i)} \right) - \frac{\partial \mathcal{L}}{\partial \phi_i} = 0$$



$$\sigma \sim \frac{\alpha}{m_p^2} \frac{1}{\sqrt{s}} - \frac{m_p}{m_p^2}$$

QFT

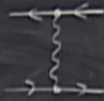
$$[a_p, a_k^\dagger] = (2\pi)^3 \delta(\vec{p} - \vec{k})$$

Bosons/Fermions

Fock space:

$$|\vec{p}_1, \dots, \vec{p}_n\rangle = a_{p_1}^\dagger \dots a_{p_n}^\dagger |0\rangle$$

Forces  $\leftrightarrow$



$$\mathcal{S} = \int d^4x \mathcal{L}(q_i, \phi_i, \psi_i)$$

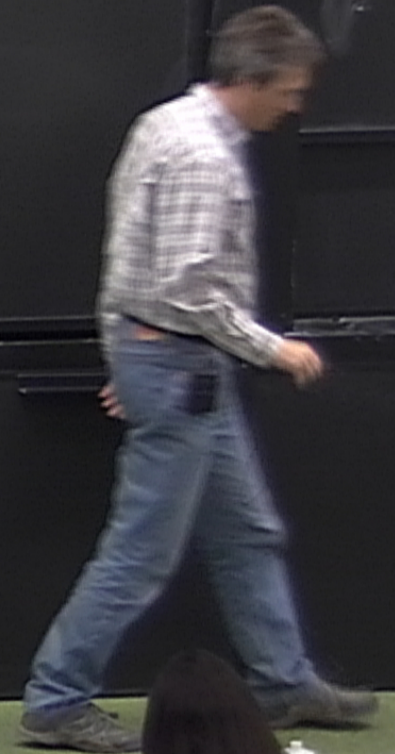
$\mathcal{L}$  Lagrangian density

$\phi_i(x,t) \equiv \phi_{i,\alpha}(x) \leftarrow$  basic dynamical variables.

Eq of motion:

$$q \frac{\partial \mathcal{L}}{\partial q} - \frac{\partial \mathcal{L}}{\partial \dot{q}} = 0$$

Symmetries  $\leftrightarrow$  Conservation laws.





$$\sigma \sim \frac{\alpha}{m_p^2} \frac{1}{\sqrt{s}} - \frac{m_p}{m_p^2}$$

QFT

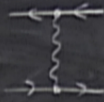
$$[a_p, a_k^\dagger] = (2\pi)^3 \delta(\vec{p}-\vec{k})$$

Bosons/Fermions

Fock space:

$$|\vec{p}_1, \dots, \vec{p}_n\rangle = a_{p_1}^\dagger \dots a_{p_n}^\dagger |0\rangle$$

Forces  $\leftrightarrow$



Symmetries  $\leftrightarrow$  Conservation laws.

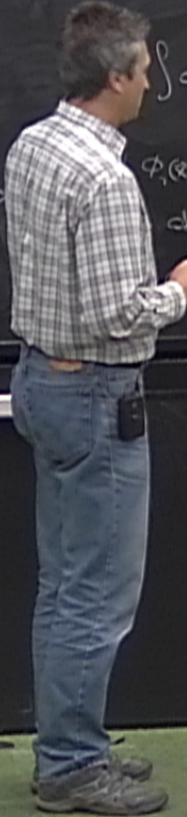
$$\int d^4x \mathcal{L}(\eta_i, \phi_i, \psi_i)$$

↑ Lagrangian density

$\phi_i(x,t) \equiv \phi_{i,\alpha}(x)$  ← basic dynamical variables.

of motion:

$$\frac{\partial \mathcal{L}}{\partial \phi_i} - \frac{\partial \mathcal{L}}{\partial \phi_i} = 0$$





$$\sigma \sim \frac{\alpha}{m_p^2} \frac{1}{\sqrt{s}} - \frac{m_p^2}{m_p^2}$$

QFT

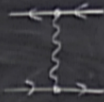
$$[a_p, a_k^\dagger] = (2\pi)^3 \delta(\vec{p} - \vec{k})$$

Bosons/Fermions

Fock space:

$$|\vec{p}_1, \dots, \vec{p}_n\rangle = a_{p_1}^\dagger \dots a_{p_n}^\dagger |0\rangle$$

Forces  $\leftrightarrow$



$$\mathcal{S} = \int d^4x \mathcal{L}(q, \phi_1, \phi_2)$$

Lagrangian

$$\phi_1(x,t) \equiv \phi_2(x,t)$$

Eq. of motion:

$$q \frac{\partial \mathcal{L}}{\partial q} - \frac{\partial \mathcal{L}}{\partial \dot{q}}$$

Symmetries  $\leftrightarrow$  Conservation laws

Use



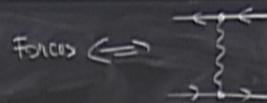
QFT

$$[a_p, a_k^\dagger] = (2\pi)^3 \delta(\vec{p}-\vec{k})$$

↑  
Bosons/fermions

Fock space.

$$|\vec{p}_1 \dots \vec{p}_n\rangle = a_{p_1}^\dagger \dots a_{p_n}^\dagger |0\rangle$$



$$S = \int d^4x \mathcal{L}(q_\mu, \phi_i, \psi_j)$$

↑ Lagrangian density

$\phi_i(x,t) \equiv \phi_i(x) \leftarrow$  basic dynamical variables.

Eqs of motion:

$$q_\mu \left( \frac{\partial \mathcal{L}}{\partial q_\mu} - \frac{\partial \mathcal{L}}{\partial \dot{q}_\mu} \right) = 0$$

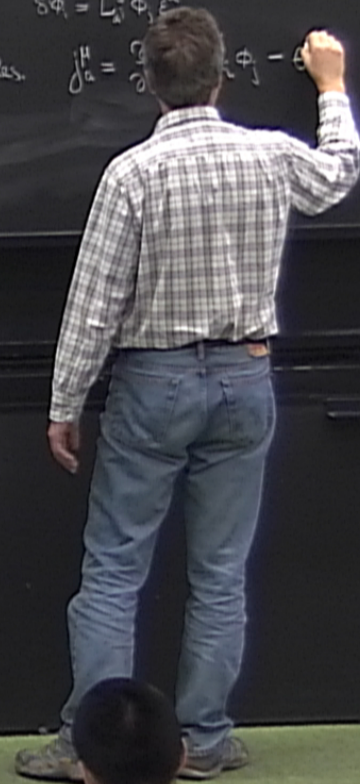
Symmetries  $\leftrightarrow$  Conservation laws.

Noether theorem:

$$\delta x^\mu = \theta^\mu_\nu x^\nu$$

$$\delta \phi = L^\mu_\nu \phi_j \theta^\nu$$

$$j^\mu_\alpha = \frac{\partial \mathcal{L}}{\partial \dot{\phi}_i} \delta \phi_i - \mathcal{L} \delta x^\mu$$





$$\cdot \lambda \neq \infty, v \neq 0$$

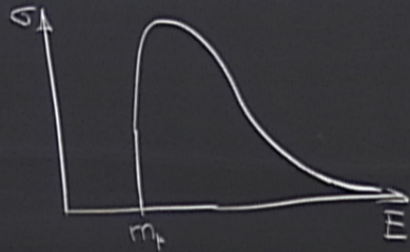
$$g = \frac{1}{m_R^2} + \left( \frac{E}{m_R} \right)$$

$$E \gg \frac{m_R}{2} :$$

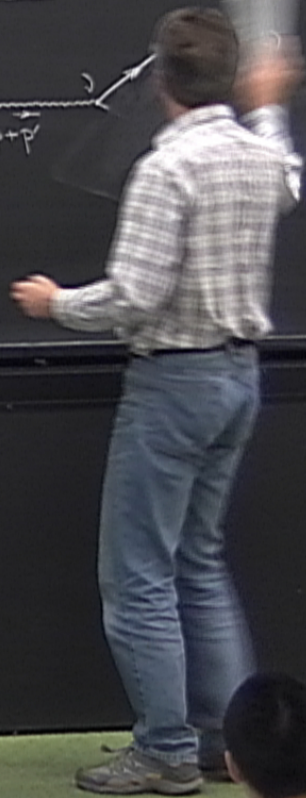
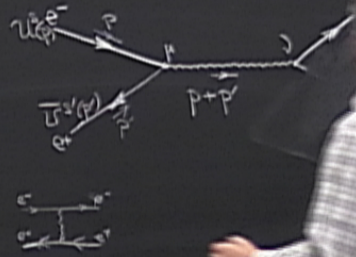
$$g \approx \frac{1}{m_R^2}$$

$$E \gg \frac{m_R}{2} :$$

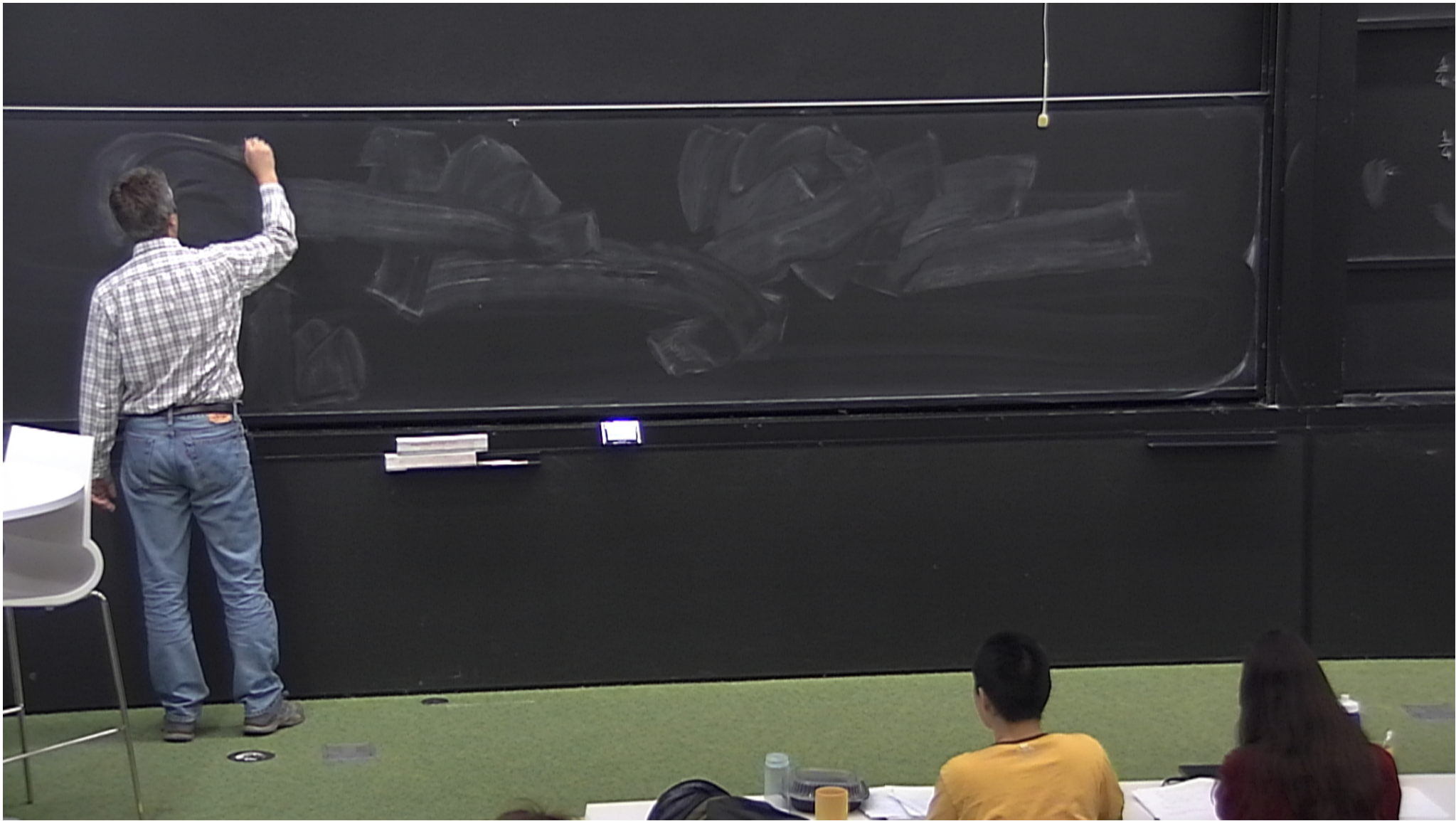
$$g \approx \frac{1}{m_R^2} \sqrt{\frac{E}{m_R}}$$



QED calculation









Klein-Gordon field:

$$\mathcal{L} = \frac{1}{2} (\partial_\mu \varphi)^2 - \frac{1}{2} m^2 \varphi^2$$

- scalar



Klein-Gordon field.

$$\mathcal{L} = \frac{1}{2} (\partial_\mu \phi)^2 - \frac{1}{2} m^2 \phi^2$$

- scalar particles of mass  $m$

Dirac field.

$$\mathcal{L} = \bar{\psi} (i \gamma^\mu \partial_\mu - m) \psi$$

- charge



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fermions



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- charged fermions  $w$ : mass  $m$   
and spin  $\frac{1}{2}$

QED

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \bar{\psi} (i \gamma^\mu D_\mu - m) \psi$$



Klein-Gordon field.

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- scalar  $\phi$   $\rightarrow$  mass  $m$

Dirac field.

$$\mathcal{L} = \bar{\psi} (i \not{\partial} - m) \psi$$

- chiral  $\psi$   $\rightarrow$  mass  $m$

QED

$$\mathcal{L} = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \bar{\psi} (i \not{D} - m) \psi$$

$$\psi \rightarrow e^{i\alpha} \psi, \quad \bar{\psi} \rightarrow \bar{\psi} e^{-i\alpha}$$



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$$D_\mu = \partial_\mu - ieA_\mu$$



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compensating field:  $A_\mu \rightarrow A_\mu - \frac{1}{q} \partial_\mu \alpha$



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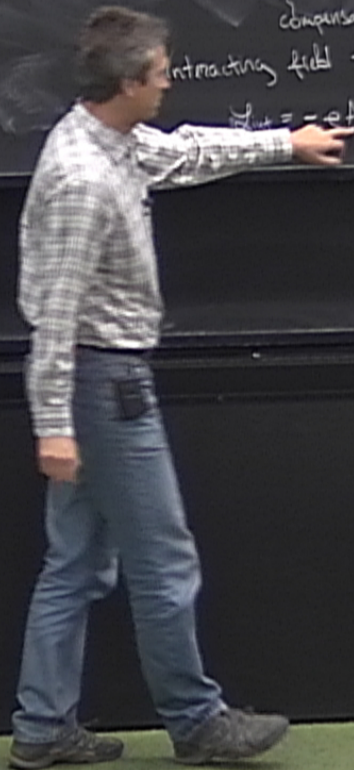
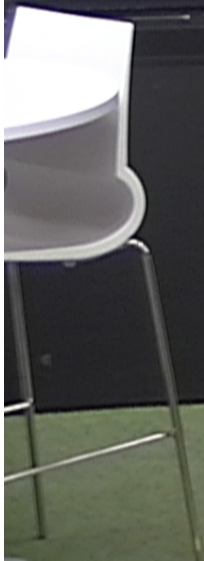
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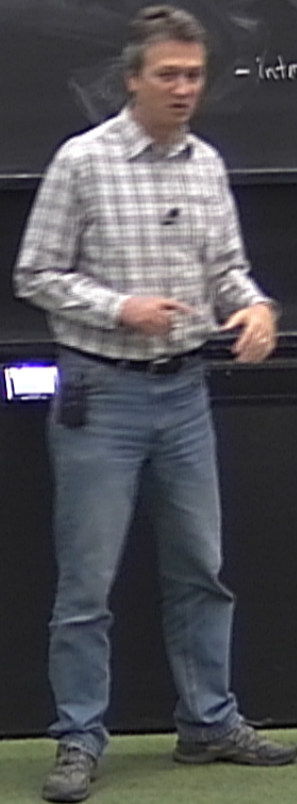
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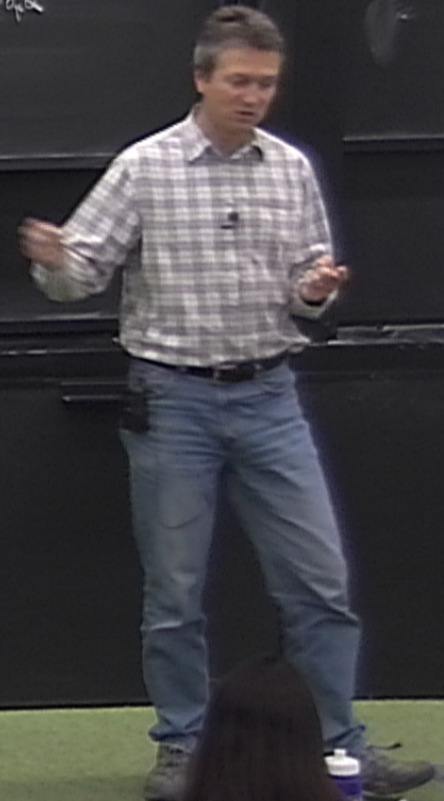
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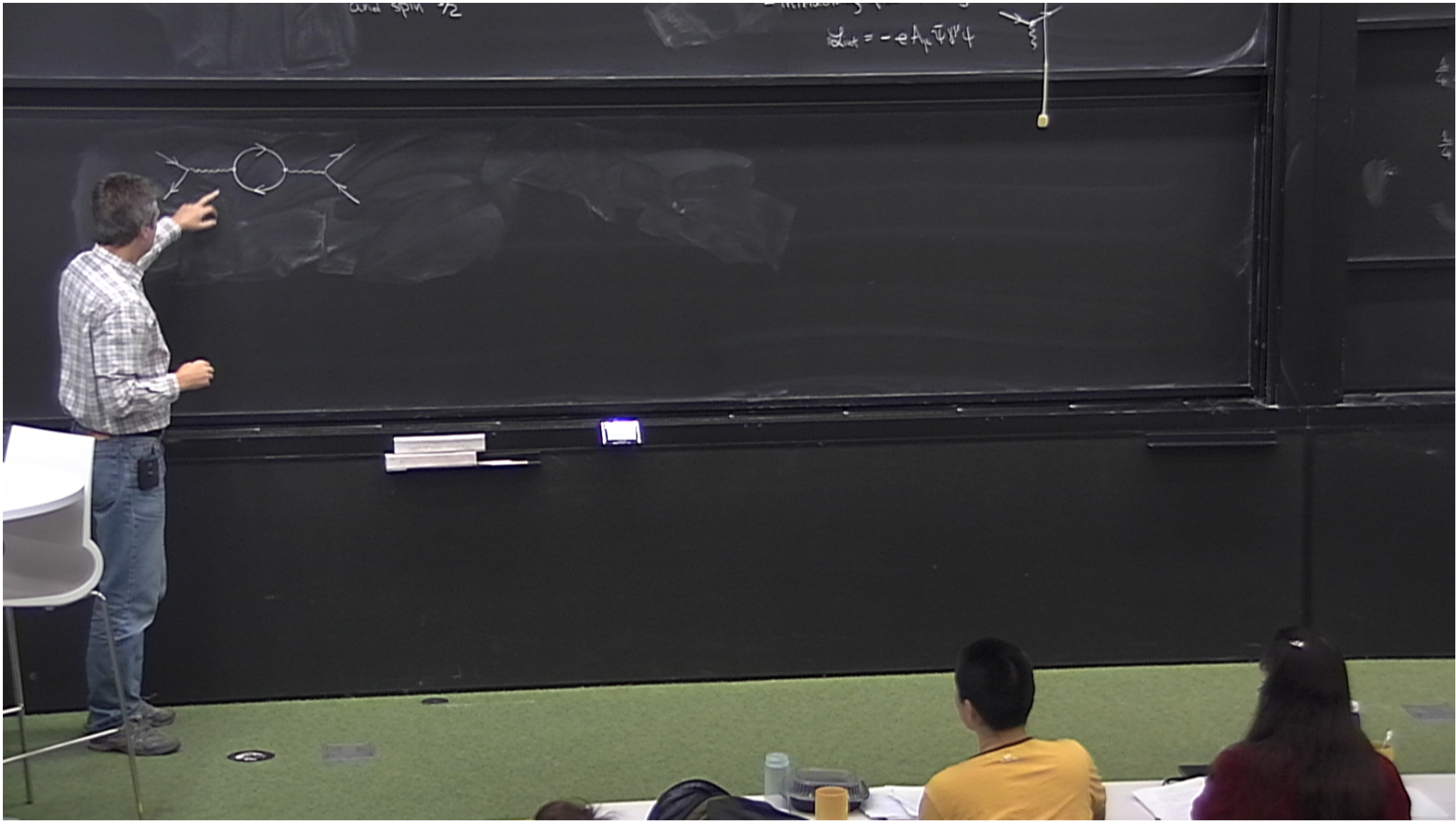
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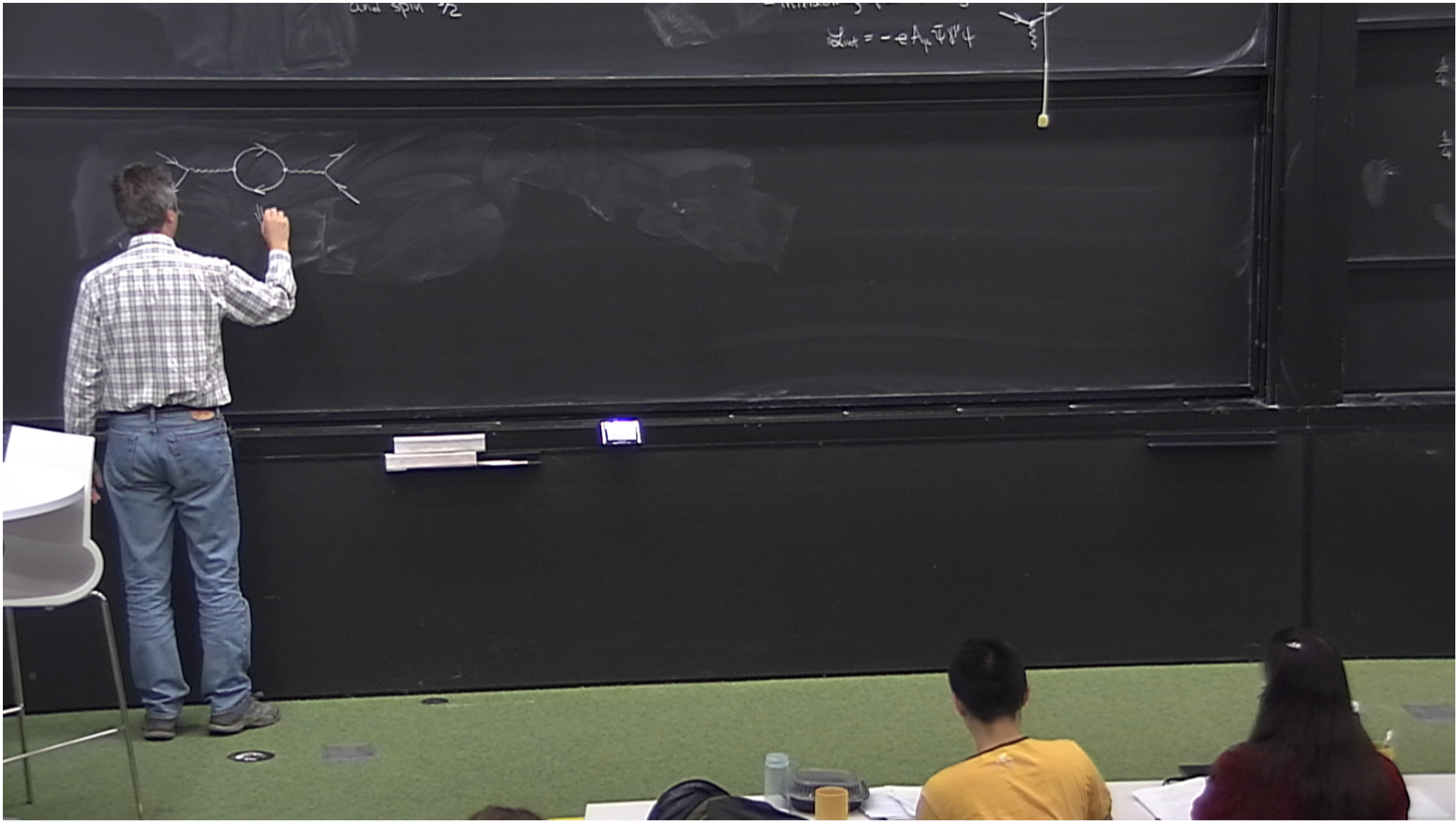
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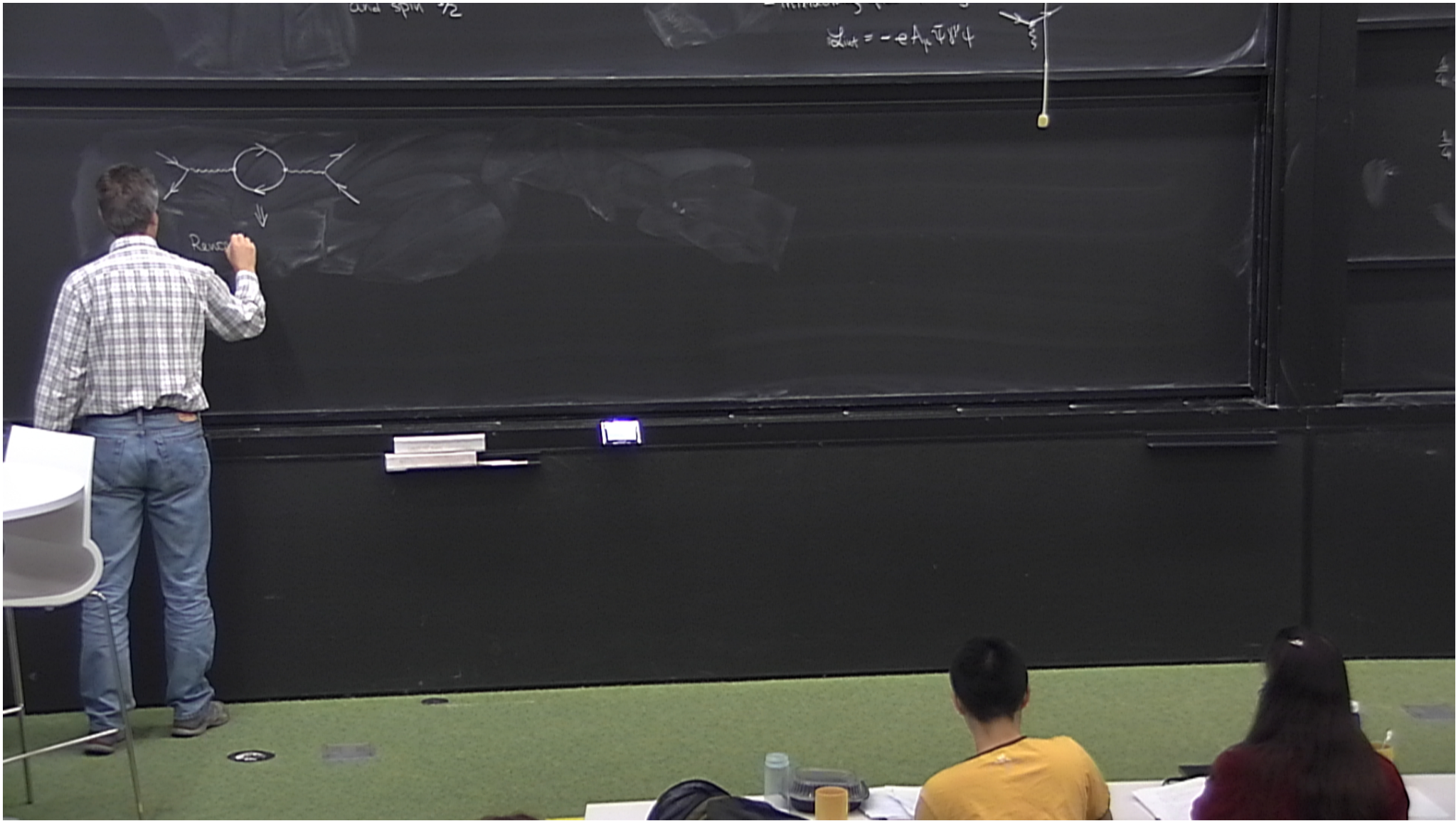








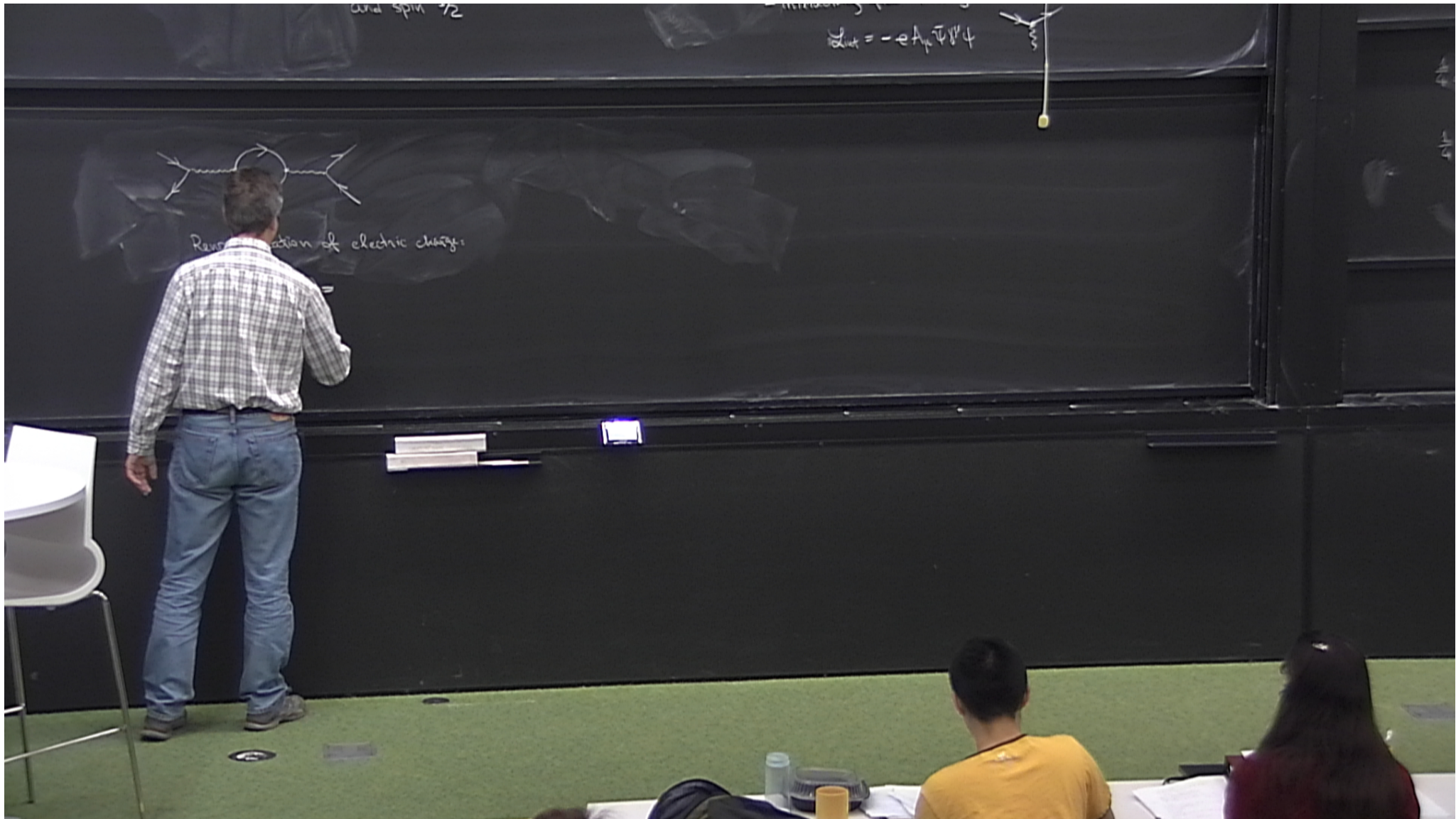




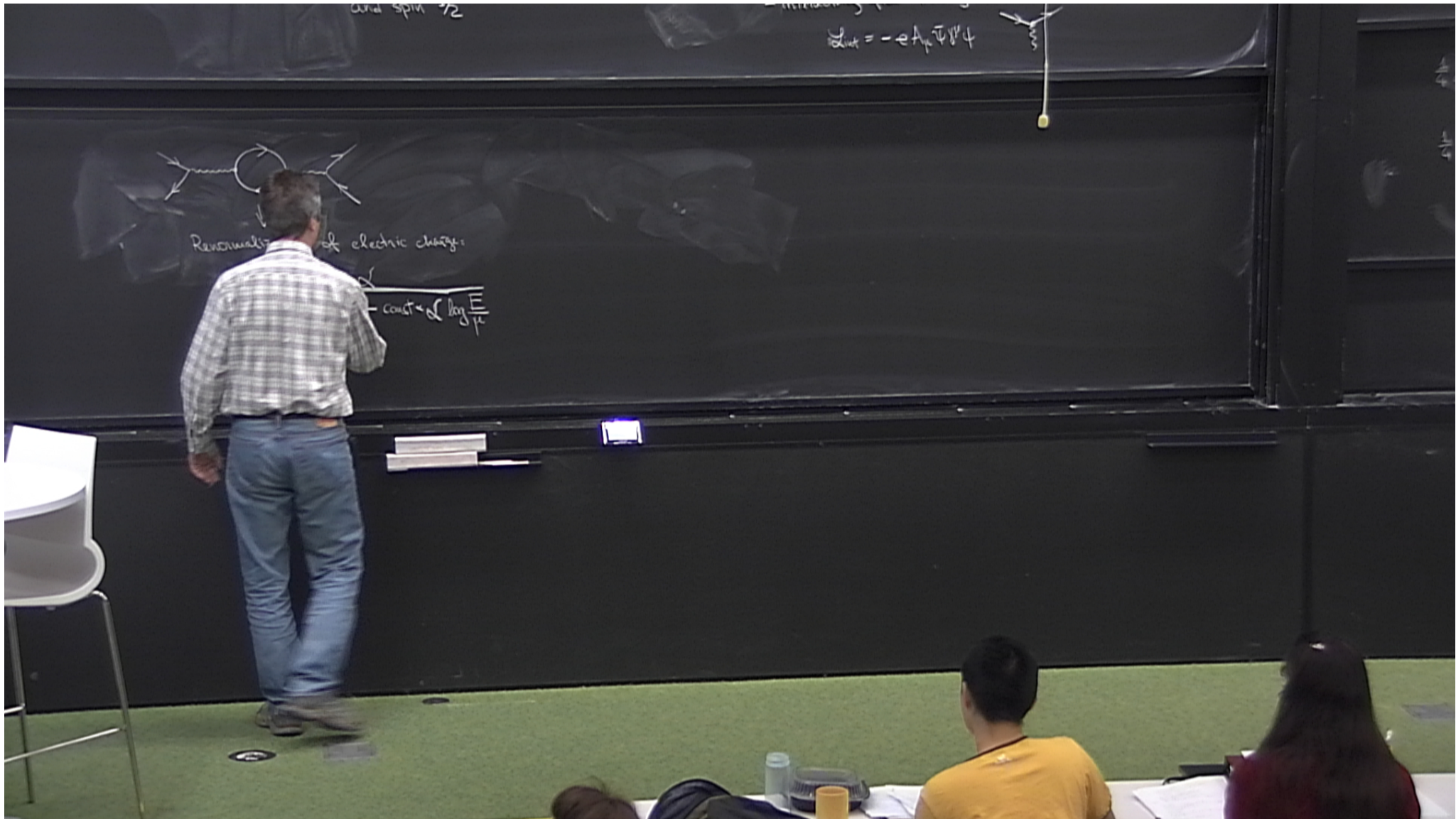














and spin 1/2

$$\mathcal{L}_{int} = -e A_\mu \bar{\psi} \gamma^\mu \psi$$



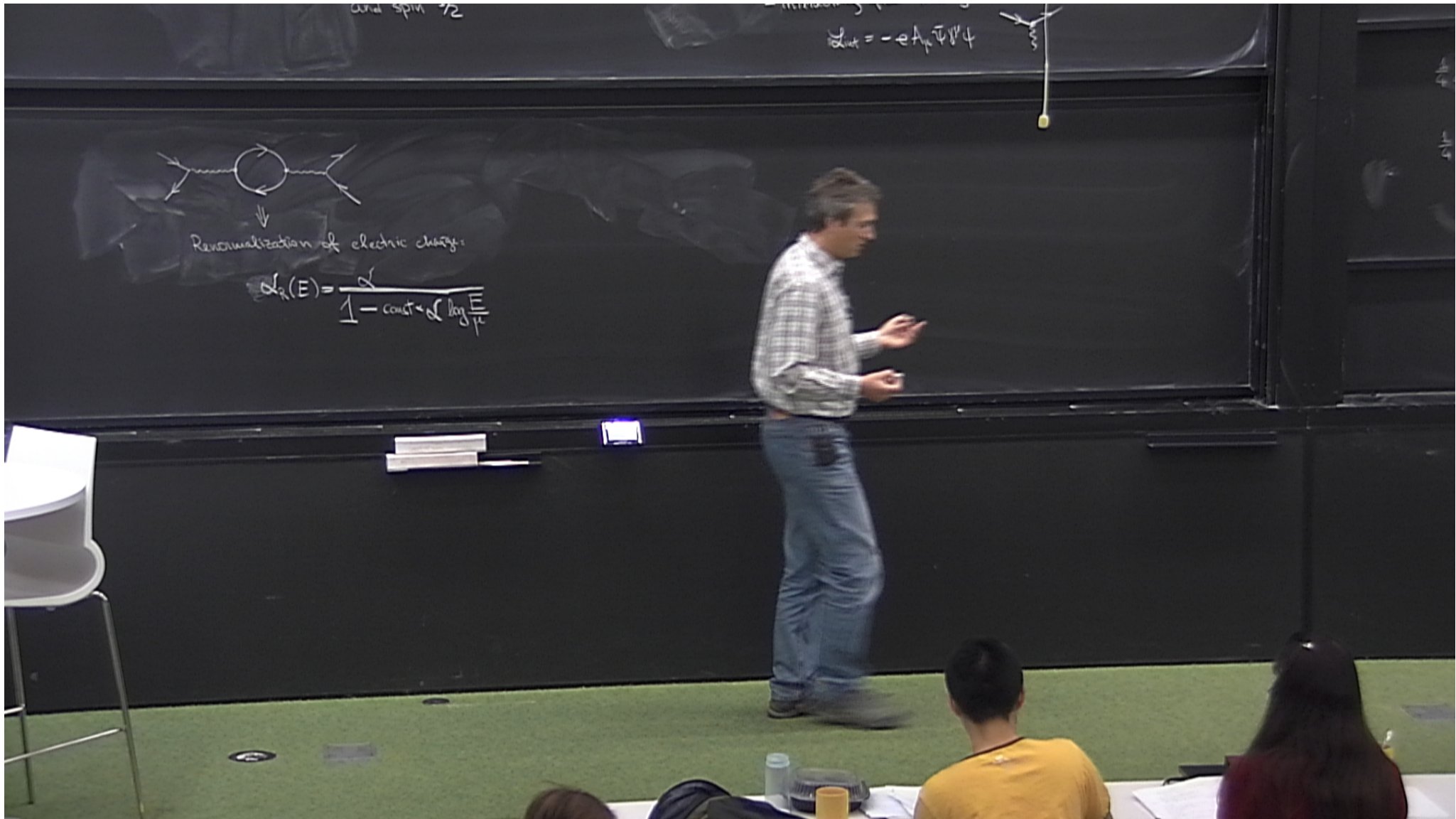
Renormalization of electric charge:

$$\alpha_s(E) = \frac{\alpha_s}{1 - \text{const} \cdot \alpha_s \log \frac{E}{\mu}}$$

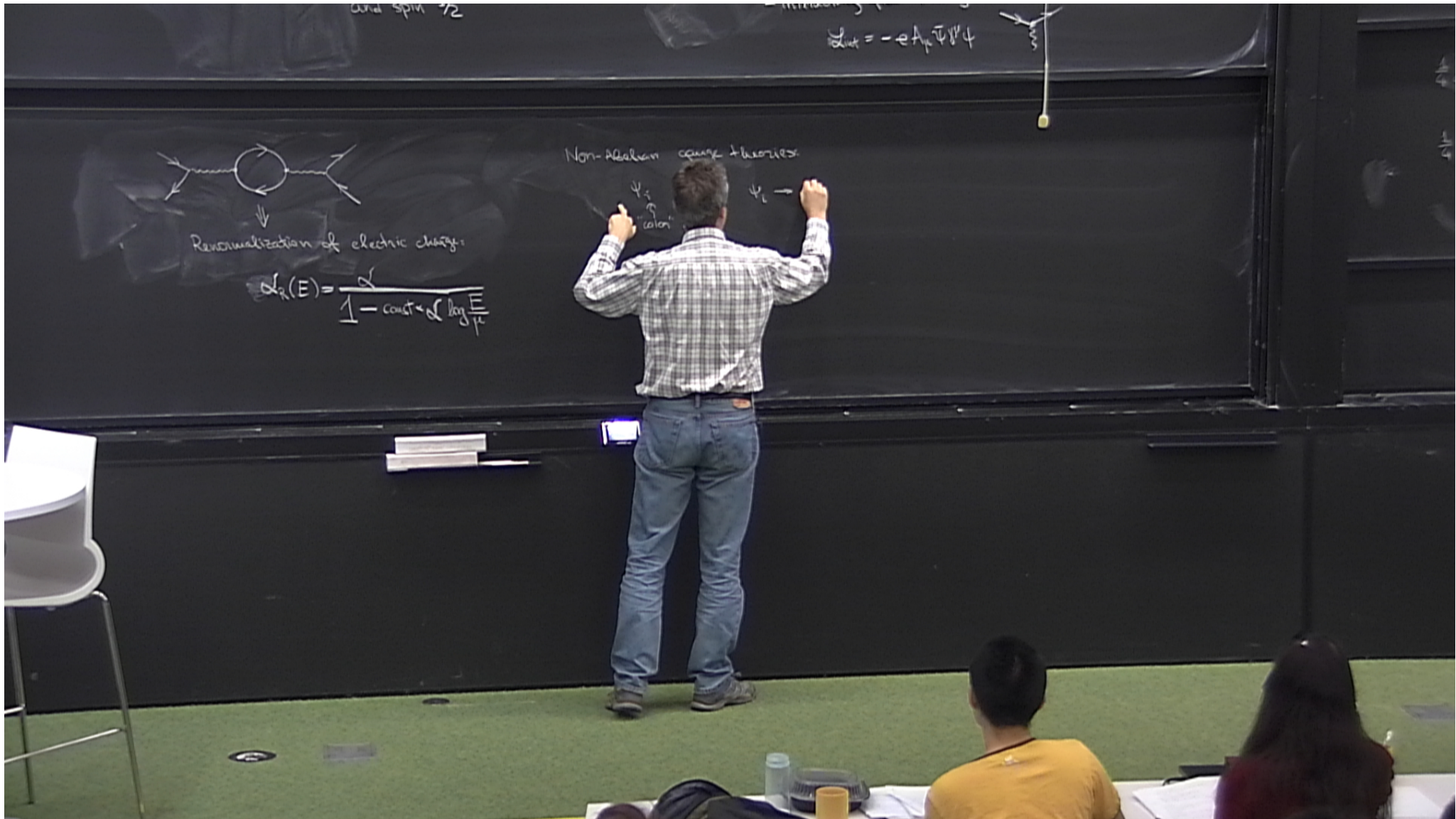
















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Renormalization of electric charge:

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Non-Abelian gauge theories

$$\psi_i \rightarrow U_i \psi_i$$

unitary matrix,  $\det U = 1$   
 $SU(N)$





Renormalization of electric charge:

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Non-Abelian gauge theories:

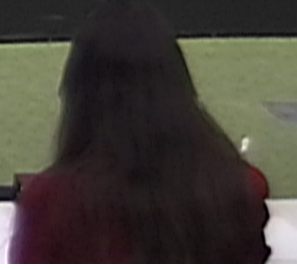
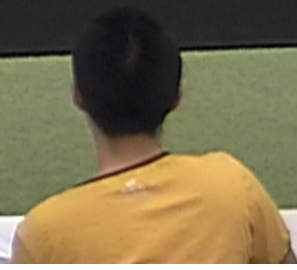
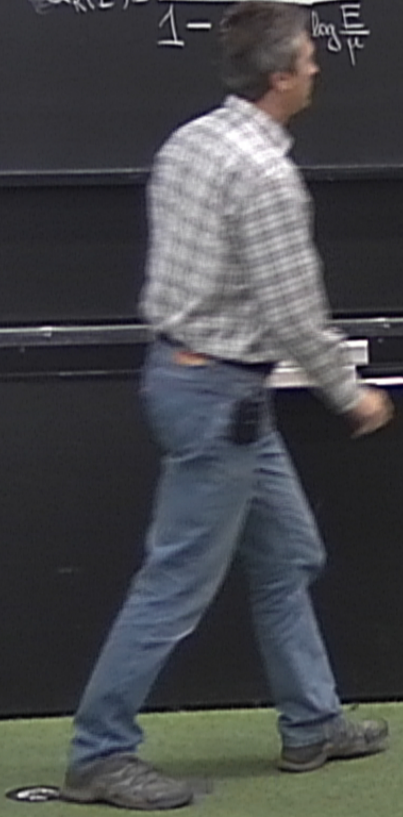
$\psi_i$   
color index

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unitary matrix,  $\det U = 1$   
SU(N)

$$D_{\mu}^{\dagger} = \partial_{\mu}^{\dagger} - A_{\mu}^{\dagger} \cdot T_{\dagger}$$

NxN Hermitian traceless matrix.







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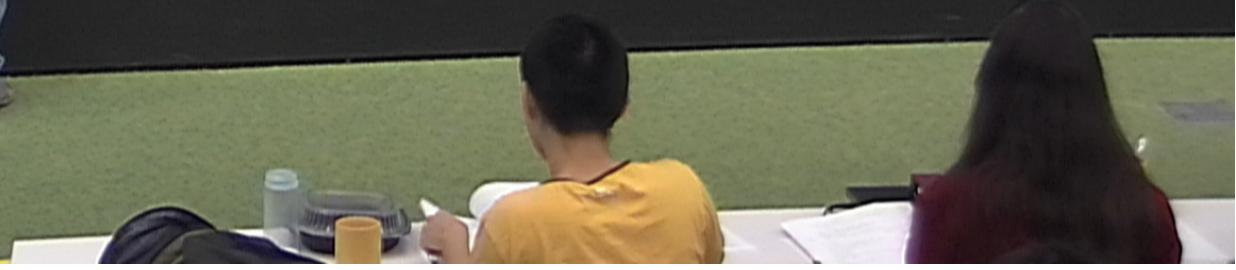
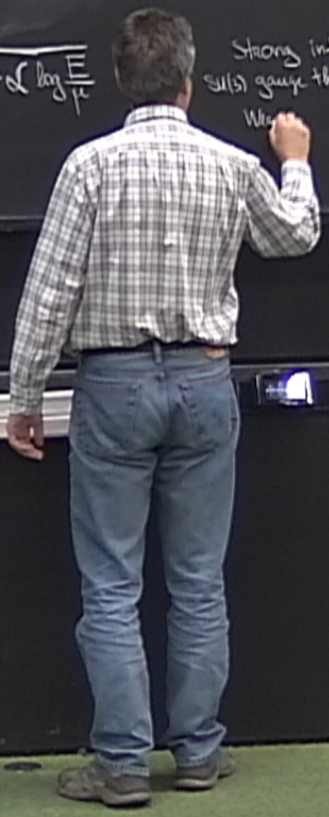
↑  
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$$D_{ij}^k = g_{ij}^k - A_{ij}^k \psi_j$$

$N \times N$  Hermitian traceless matrix.

Strong interactions  
 $SU(3)$  gauge theory: QCD  
Weyl







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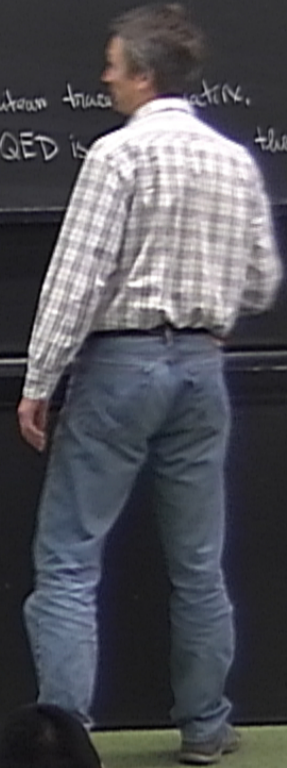
$SU(N)$

$$D_{ij}^k = g A_{ij}^k - A_{ij}^l A_{lk}^k$$

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QED is  $U(1)$  theory

Strong interactions  
 $SU(3)$  gauge theory: QCD  
Weak interactions  
 $SU(2) \times U(1)$





Renormalization of electric charge:

$$\alpha_2(E) = \frac{\alpha}{1 - \text{const} \cdot \alpha \log \frac{E}{\mu}}$$

Strong interactions  
SU(3) gauge theory: QCD  
Weak interactions  
→  
broken

$\psi_i$   
↑  
"color" index

$\psi_i \rightarrow U_i^j \psi_j$   
↑  
unitary matrix,  $\det U = 1$   
SU(N)

$$D_{ij}^a = g A_{ij}^a - A_{ij}^a \gamma_5$$

NxN Hermitian traceless matrix.

QED is a U(1) gauge theory



Renormalization of electric charge:

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QED is a  $U(1)$  gauge theory

Strong interactions  
 $SU(3)$  gauge theory: QCD

Weak interactions  
 $SU(2) \times U(1) \rightarrow U(1)$   
broken  
QED