

Title: Numerical Holographic Striped Phases

Date: Jun 08, 2012 03:15 PM

URL: <http://pirsa.org/12060029>

Abstract: TBA



a place of mind



Numerical Holographic Striped Phases

Exploring AdS-CFT Dualities in Dynamical Settings

Jared Stang

June 8th, 2012

with M. Rozali, D. Smyth, and E. Sorkin

Motivation and goal I

1. AdS/CMT: stripes in high- T_c superconductors
2. Understanding strongly coupled physics

Goal:

Gravity dual of a 2+1 system with stripes

Motivation and goal I

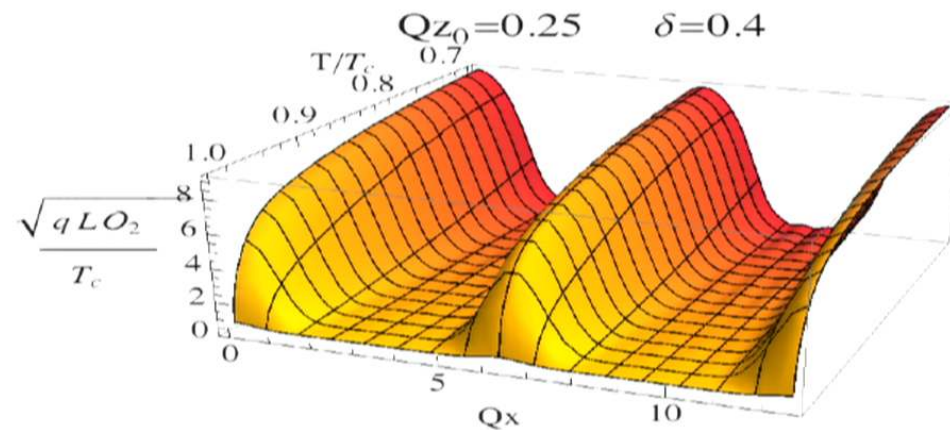
1. AdS/CMT: stripes in high- T_c superconductors
2. Understanding strongly coupled physics

Goal:

Gravity dual of a 2+1 system with stripes

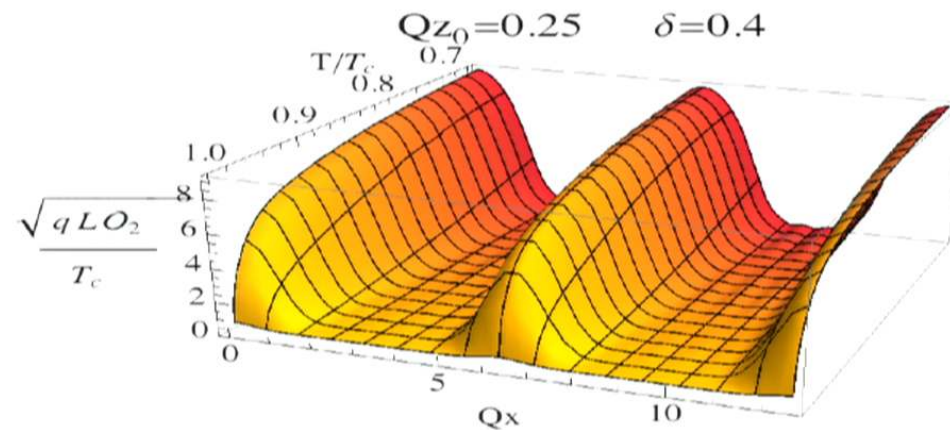
Background I

- One method: source stripes in the Hartnoll-Herzog-Horowitz superconductor Flauger, Pajer and Papanikolaou, 1010.1775



Background I

- One method: source stripes in the Hartnoll-Herzog-Horowitz superconductor Flauger, Pajer and Papanikolaou, 1010.1775



Background II

- Instabilities with Chern-Simons terms
 - Holographic QCD
Domokos and Harvey, 0704.1604
 - Maxwell-Chern-Simons in $4+1$
Nakamura, Ooguri and Park, 0911.0679
 - D7 probe brane in D3 background
Bergman *et al.*, 1106.3883
 - Einstein-Maxwell-axion in $3+1$
Donos and Gauntlett, 1106.2004

Motivation and goal II

1. AdS/CMT: stripes
2. Understanding strongly coupled physics
3. How do the stripes back-react on the geometry?

Goal:

Gravity dual of a $2+1$ system with stripes

Motivation and goal II

1. AdS/CMT: stripes
2. Understanding strongly coupled physics
3. How do the stripes back-react on the geometry?

Goal:

Back-reacted gravity dual of $2+1$ system with **spontaneous stripes**

The model

- AdS gravity, neutral pseudo-scalar, and gauge field:

Donos and Gauntlett, 1106.2004

$$\mathcal{L} = \frac{1}{2} \sqrt{-g} (R + \Lambda) - \frac{1}{2} \sqrt{-g} (\partial^\mu \psi \partial_\mu \psi + m^2 \psi^2) - \frac{1}{4} \sqrt{-g} F^{\mu\nu} F_{\mu\nu} - \frac{c_1}{8} \epsilon^{\mu\nu\rho\sigma} \psi F_{\mu\nu} F_{\rho\sigma}$$

- Homogeneous solution is AdS-Reissner-Nordstrom:

$$ds_{RN}^2 = -2r^2 f_{RN} dt^2 + \frac{dr^2}{2r^2 f_{RN}} + 2r^2 (dx^2 + dy^2)$$

$$A = \left(1 - \frac{r_+}{r}\right) dt, \quad \psi = 0 \quad f_{RN} = 1 - \left(1 + \frac{1}{4r_+^2}\right) \left(\frac{r_+}{r}\right)^3 + \frac{1}{4r_+^2} \left(\frac{r_+}{r}\right)^4$$

The model

- AdS gravity, neutral pseudo-scalar, and gauge field:

Donos and Gauntlett, 1106.2004

$$\mathcal{L} = \frac{1}{2} \sqrt{-g} (R + \Lambda) - \frac{1}{2} \sqrt{-g} (\partial^\mu \psi \partial_\mu \psi + m^2 \psi^2) - \frac{1}{4} \sqrt{-g} F^{\mu\nu} F_{\mu\nu} - \frac{c_1}{8} \epsilon^{\mu\nu\rho\sigma} \psi F_{\mu\nu} F_{\rho\sigma}$$

- Homogeneous solution is AdS-Reissner-Nordstrom:

$$ds_{RN}^2 = -2r^2 f_{RN} dt^2 + \frac{dr^2}{2r^2 f_{RN}} + 2r^2(dx^2 + dy^2)$$

$$A = \left(1 - \frac{r_+}{r}\right) dt, \quad \psi = 0 \quad f_{RN} = 1 - \left(1 + \frac{1}{4r_+^2}\right) \left(\frac{r_+}{r}\right)^3 + \frac{1}{4r_+^2} \left(\frac{r_+}{r}\right)^4$$

The linearised system

Donos and Gauntlett, 1106.2004

- Consider this fluctuation and look for normalisable modes:

$$\delta g_{ty} = \lambda r(r - r_+)w(r) \sin(kx)$$

$$\delta A_y = \lambda a(r) \sin(kx)$$

$$\delta \psi = \lambda \phi(r) \cos(kx)$$

- Asymptotic expansions:

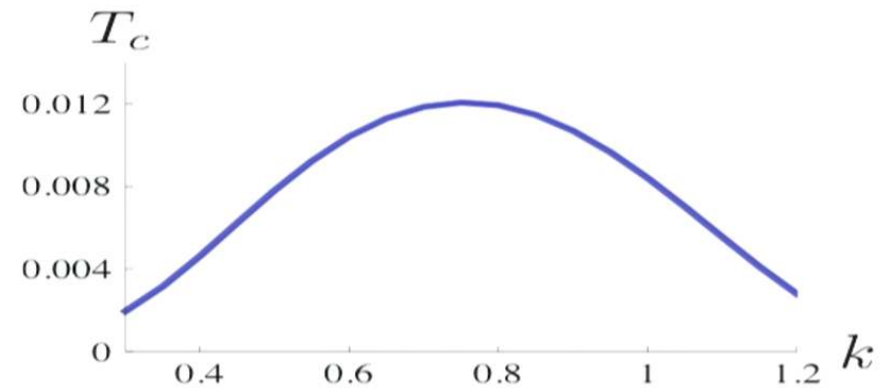
$$w = w_0 + \dots + \frac{w_3}{r^3} + \dots$$

$$a = a_0 + \dots + \frac{a_1}{r} + \dots$$

$$\phi = \frac{\phi_1}{r} + \dots + \frac{\phi_2}{r^2} + \dots$$

- At the next order:

$$\delta Z = \lambda^2(Z^{(0)}(r) + Z^{(1)}(r) \cos(2kx)), \quad Z = \{g_{tt}, g_{xx}, g_{yy}, A_t\}$$



The linearised system

Donos and Gauntlett, 1106.2004

- Consider this fluctuation and look for normalisable modes:

$$\delta g_{ty} = \lambda r(r - r_+)w(r) \sin(kx)$$

$$\delta A_y = \lambda a(r) \sin(kx)$$

$$\delta \psi = \lambda \phi(r) \cos(kx)$$

- Asymptotic expansions:

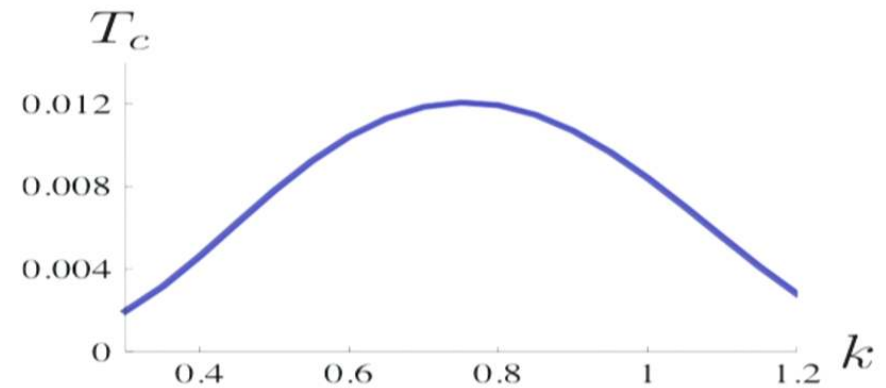
$$w = w_0 + \dots + \frac{w_3}{r^3} + \dots$$

$$a = a_0 + \dots + \frac{a_1}{r} + \dots$$

$$\phi = \frac{\phi_1}{r} + \dots + \frac{\phi_2}{r^2} + \dots$$

- At the next order:

$$\delta Z = \lambda^2(Z^{(0)}(r) + Z^{(1)}(r) \cos(2kx)), \quad Z = \{g_{tt}, g_{xx}, g_{yy}, A_t\}$$



The linearised system

Donos and Gauntlett, 1106.2004

- Consider this fluctuation and look for normalisable modes:

$$\delta g_{ty} = \lambda r(r - r_+) w(r) \sin(kx)$$

$$\delta A_y = \lambda a(r) \sin(kx)$$

$$\delta \psi = \lambda \phi(r) \cos(kx)$$

- Asymptotic expansions:

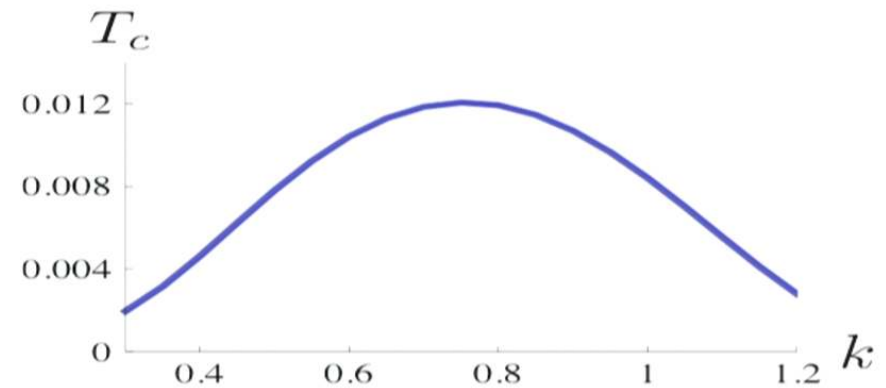
$$w = w_0 + \dots + \frac{w_3}{r^3} + \dots$$

$$a = a_0 + \dots + \frac{a_1}{r} + \dots$$

$$\phi = \frac{\phi_1}{r} + \dots + \frac{\phi_2}{r^2} + \dots$$

- At the next order:

$$\delta Z = \lambda^2 (Z^{(0)}(r) + Z^{(1)}(r) \cos(2kx)), \quad Z = \{g_{tt}, g_{xx}, g_{yy}, A_t\}$$



The linearised system

Donos and Gauntlett, 1106.2004

- Consider this fluctuation and look for normalisable modes:

$$\delta g_{ty} = \lambda r(r - r_+)w(r) \sin(kx)$$

$$\delta A_y = \lambda a(r) \sin(kx)$$

$$\delta \psi = \lambda \phi(r) \cos(kx)$$

- Asymptotic expansions:

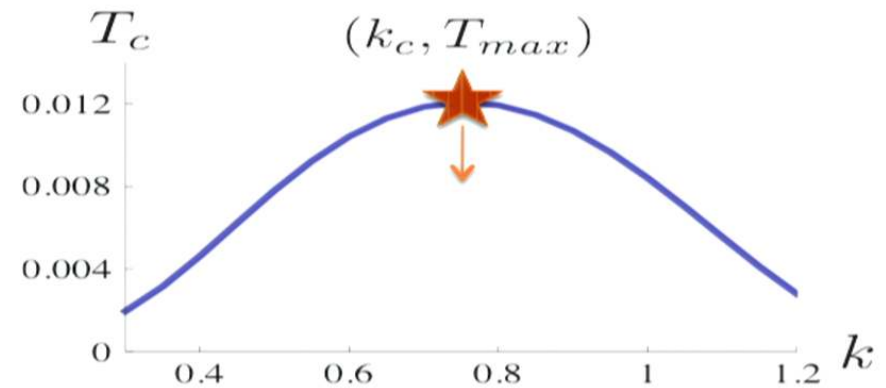
$$w = \cancel{w_0} + \dots + \frac{w_3}{r^3} + \dots$$

$$a = \cancel{a_0} + \dots + \frac{a_1}{r} + \dots$$

$$\phi = \cancel{\frac{\phi_1}{r}} + \dots + \frac{\phi_2}{r^2} + \dots$$

- At the next order:

$$\delta Z = \lambda^2 (Z^{(0)}(r) + Z^{(1)}(r) \cos(2kx)), \quad Z = \{g_{tt}, g_{xx}, g_{yy}, A_t\}$$



The linearised system

Donos and Gauntlett, 1106.2004

- Consider this fluctuation and look for normalisable modes:

$$\delta g_{ty} = \lambda r(r - r_+)w(r) \sin(kx)$$

$$\delta A_y = \lambda a(r) \sin(kx)$$

$$\delta \psi = \lambda \phi(r) \cos(kx)$$

- Asymptotic expansions:

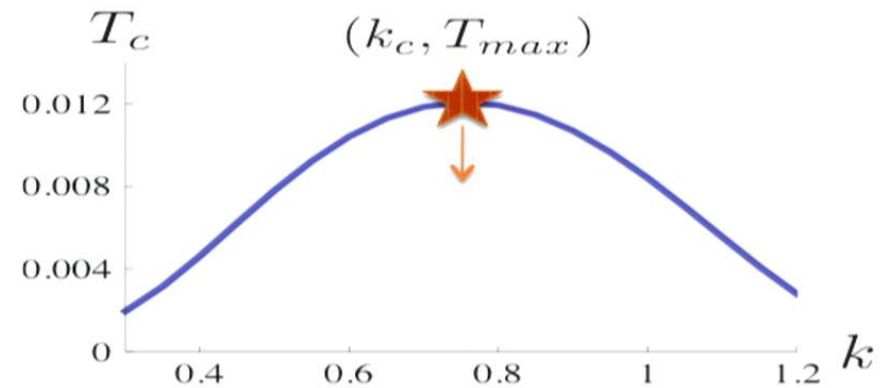
$$w = \cancel{w_0} + \dots + \frac{w_3}{r^3} + \dots$$

$$a = \cancel{a_0} + \dots + \frac{a_1}{r} + \dots$$

$$\phi = \cancel{\frac{\phi_1}{r}} + \dots + \frac{\phi_2}{r^2} + \dots$$

- At the next order:

$$\delta Z = \lambda^2 (Z^{(0)}(r) + Z^{(1)}(r) \cos(2kx)), \quad Z = \{g_{tt}, g_{xx}, g_{yy}, A_t\}$$



The linearised system

Donos and Gauntlett, 1106.2004

- Consider this fluctuation and look for normalisable modes:

$$\delta g_{ty} = \lambda r(r - r_+)w(r) \sin(kx)$$

$$\delta A_y = \lambda a(r) \sin(kx)$$

$$\delta \psi = \lambda \phi(r) \cos(kx)$$

- Asymptotic expansions:

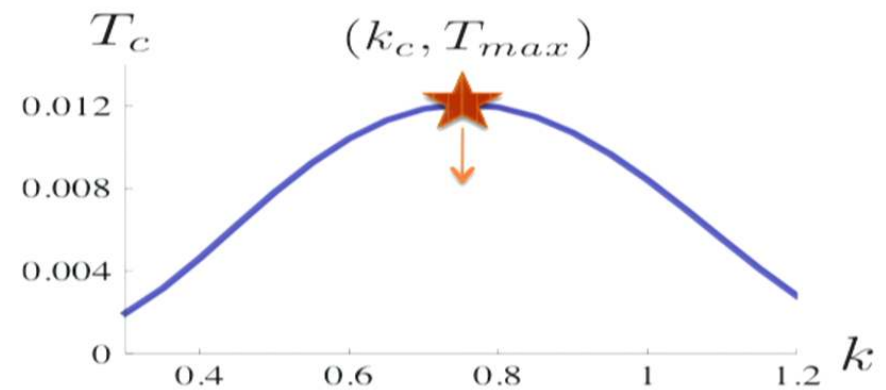
$$w = \cancel{w_0} + \dots + \frac{w_3}{r^3} + \dots$$

$$a = \cancel{a_0} + \dots + \frac{a_1}{r} + \dots$$

$$\phi = \cancel{\frac{\phi_1}{r}} + \dots + \frac{\phi_2}{r^2} + \dots$$

- At the next order:

$$\delta Z = \lambda^2(Z^{(0)}(r) + Z^{(1)}(r) \cos(2kx)), \quad Z = \{g_{tt}, g_{xx}, g_{yy}, A_t\}$$



Non-linear system: Ansatz

$$ds^2 = -2r^2 f_{RN} e^{2A(r,x)} dt^2 + e^{2B(r,x)} \left(\frac{dr^2}{2r^2 f_{RN}} + 2r^2 dx^2 \right) + 2r^2 e^{2C(r,x)} (dy - W(r,x)dt)^2$$

$$A = A_t(r,x)dt + A_y(r,x)dy, \quad \psi = \psi(r,x)$$

$$T = T_{max} e^{A-B} \Big|_{r=r_+}$$

$$x = \pi/k_c$$

7 fields:

$A, B, C, A_t, W, A_y, \psi$

$$x = 0$$

$$r = r_+$$

$$r = r_{max}$$

Non-linear system: Ansatz

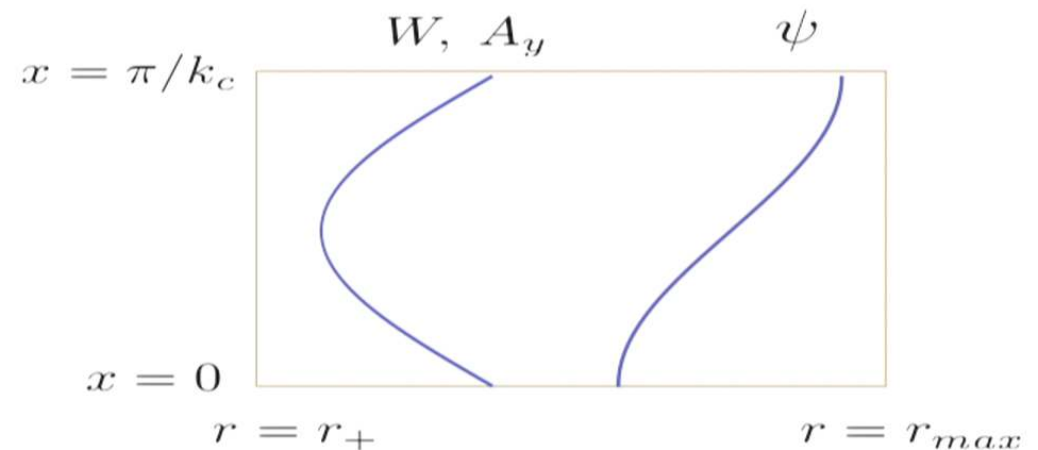
$$ds^2 = -2r^2 f_{RN} e^{2A(r,x)} dt^2 + e^{2B(r,x)} \left(\frac{dr^2}{2r^2 f_{RN}} + 2r^2 dx^2 \right) + 2r^2 e^{2C(r,x)} (dy - W(r,x)dt)^2$$

$$A = A_t(r,x)dt + A_y(r,x)dy, \quad \psi = \psi(r,x)$$

$$T = T_{max} e^{A-B} |_{r=r_+}$$

7 fields:

$A, B, C, A_t, W, A_y, \psi$



Non-linear system: Equations of motion

$$\left. \begin{aligned} G_t^t &= T_t^t \\ G_y^t &= T_y^t \\ G_y^t &= T_y^t \\ G_r^r + G_x^x &= T_r^r + T_x^x \end{aligned} \right\} \begin{array}{l} 2^{\text{nd}} \text{ order, elliptic E.O.M. for} \\ A, B, C, W \end{array}$$

$$\left. \begin{aligned} G_x^r &= T_x^r \\ G_r^r - G_x^x &= T_r^r - T_x^x \end{aligned} \right\} \text{constraint equations}$$

E.O.M. + boundary conditions \implies constraints are satisfied

Wiseman, hep-th/0209051

Non-linear system: Boundary conditions

1. $r = r_{max}$: falloff conditions

$$A = A_0(x) + \dots + \frac{A_3(x)}{r^3} + \dots$$

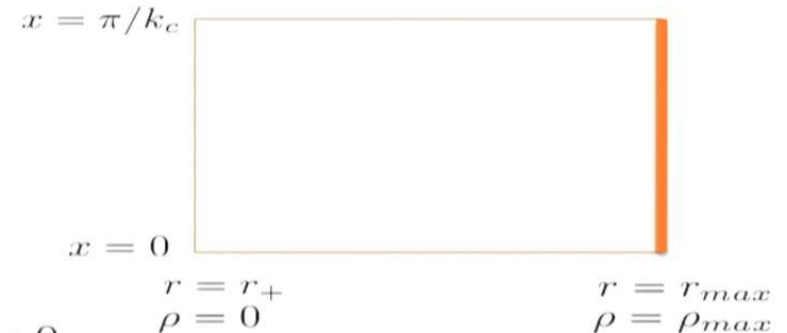
2. $r = r_+$:

- regularity conditions

With $\rho = \sqrt{r^2 - r_+^2}$, $\partial_\rho C|_{\rho=0} = 0$

- constraints

$$\partial_x(A - B)|_{\rho=0} = 0 \Rightarrow A|_{\rho=0} = B|_{\rho=0} - d_0 \quad (T = T_{max}e^{-d_0})$$



3. $x = 0, x = \pi/k_c$: Neumann or Dirichlet conditions

Solve: Finite difference, use Gauss-Siedel relaxation

Non-linear system: Boundary conditions

1. $r = r_{max}$: falloff conditions

$$A = A_0(x) + \dots + \frac{A_3(x)}{r^3} + \dots$$

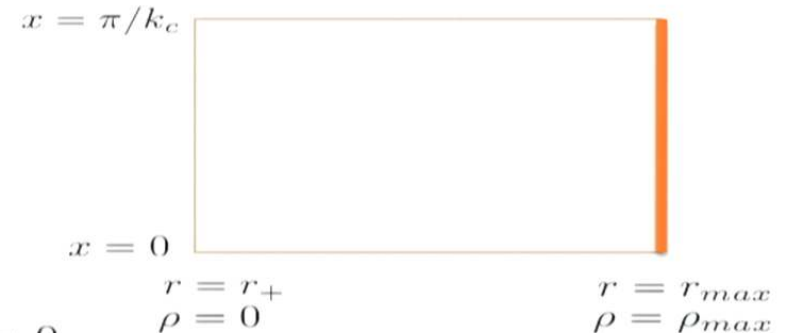
2. $r = r_+$:

- regularity conditions

With $\rho = \sqrt{r^2 - r_+^2}$, $\partial_\rho C|_{\rho=0} = 0$

- constraints

$$\partial_x(A - B)|_{\rho=0} = 0 \Rightarrow A|_{\rho=0} = B|_{\rho=0} - d_0 \quad (T = T_{max}e^{-d_0})$$



3. $x = 0, x = \pi/k_c$: Neumann or Dirichlet conditions

Solve: Finite difference, use Gauss-Siedel relaxation

Non-linear system: Boundary conditions

1. $r = r_{max}$: falloff conditions

$$A = A_0(x) + \dots + \frac{A_3(x)}{r^3} + \dots$$

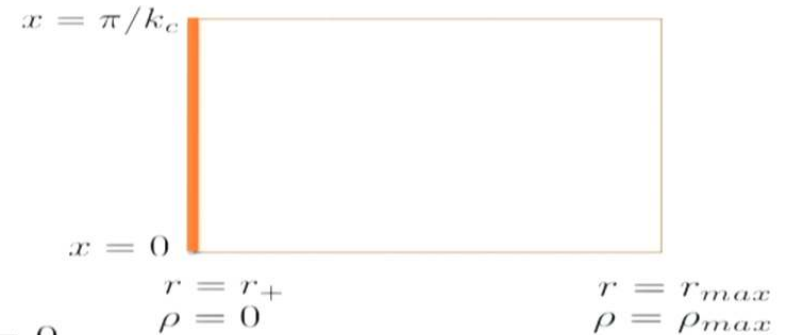
2. $r = r_+$:

- regularity conditions

With $\rho = \sqrt{r^2 - r_+^2}$, $\partial_\rho C|_{\rho=0} = 0$

- constraints

$$\partial_x(A - B)|_{\rho=0} = 0 \Rightarrow A|_{\rho=0} = B|_{\rho=0} - d_0 \quad (T = T_{max}e^{-d_0})$$



3. $x = 0, x = \pi/k_c$: Neumann or Dirichlet conditions

Solve: Finite difference, use Gauss-Siedel relaxation

Non-linear system: Ansatz

$$ds^2 = -2r^2 f_{RN} e^{2A(r,x)} dt^2 + e^{2B(r,x)} \left(\frac{dr^2}{2r^2 f_{RN}} + 2r^2 dx^2 \right) + 2r^2 e^{2C(r,x)} (dy - W(r,x)dt)^2$$

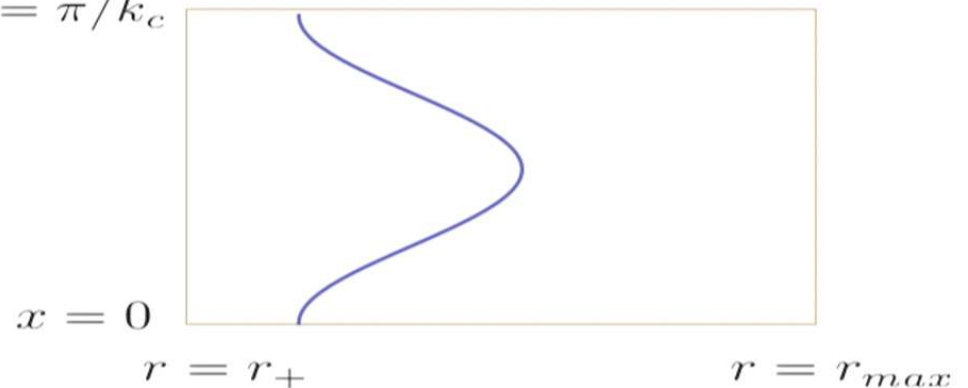
$$A = A_t(r,x)dt + A_y(r,x)dy, \quad \psi = \psi(r,x)$$

$$T = T_{max} e^{A-B} |_{r=r_+}$$

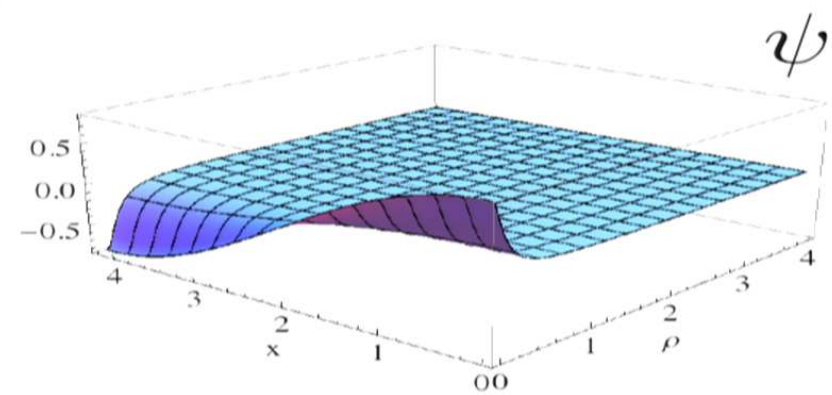
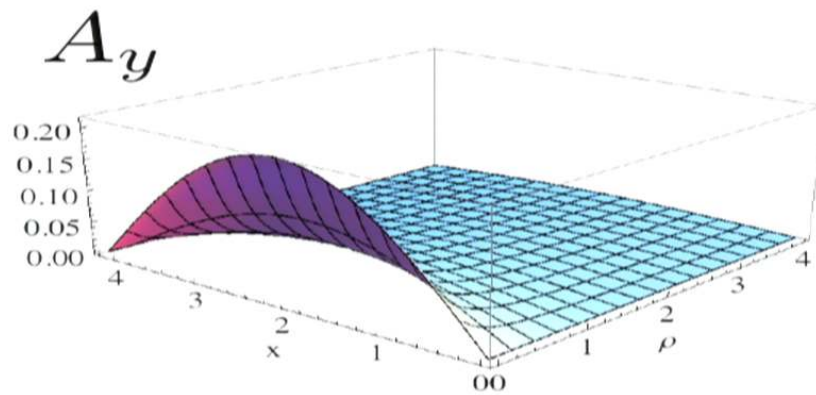
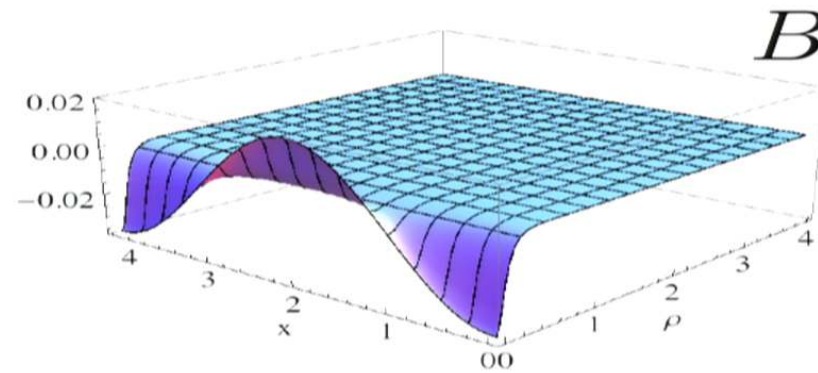
$$x = \pi/k_c$$

7 fields:

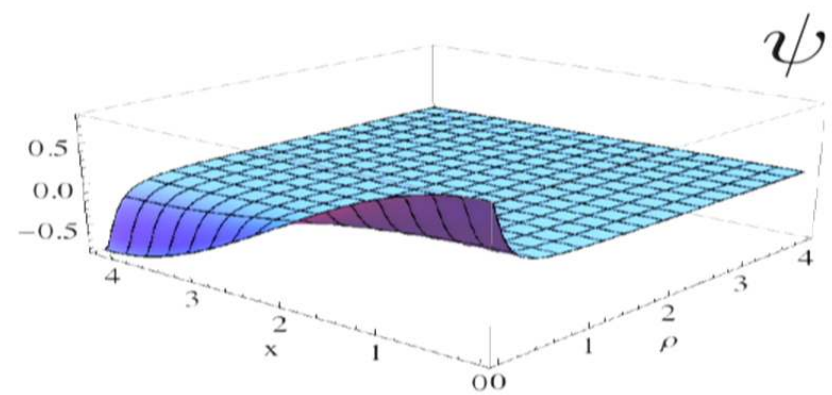
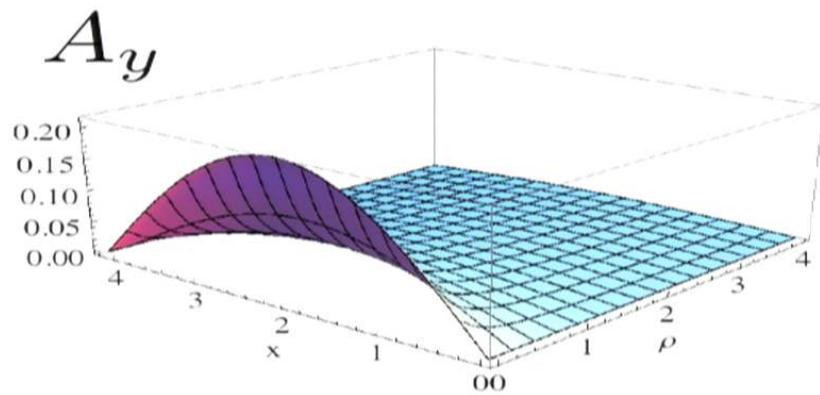
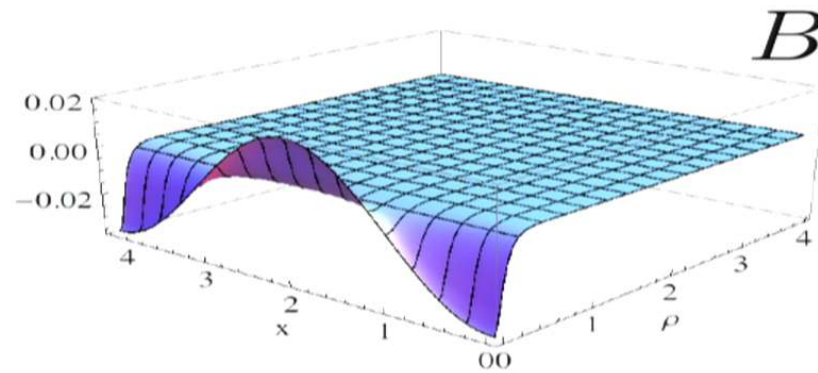
$A, B, C, A_t, W, A_y, \psi$



Preliminary results



Preliminary results



Further questions

1. What are the properties of the field theory?
2. Is this the dominant phase?
3. Zero temperature limit of this phase?