

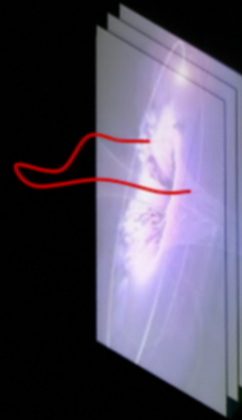
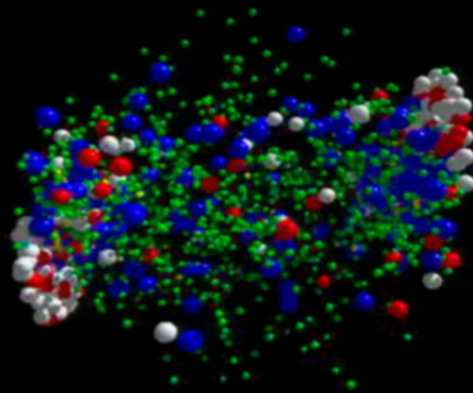
Title: Strong Coupling Isotropization Simplified

Date: Jun 04, 2012 09:00 AM

URL: <http://pirsa.org/12060013>

Abstract: We study the isotropization of a homogeneous, strongly coupled, non-Abelian plasma by means of its gravity dual. We compare the time evolution of a large number of initially anisotropic states as determined, on the one hand, by the full non-linear Einstein's equations and, on the other, by the Einstein's equations linearized around the final equilibrium state. The linear approximation works remarkably well even for states that exhibit large anisotropies. For example, it predicts with a 20% accuracy the isotropization time, which is of order $1/T$, with T the final equilibrium temperature. We comment on possible extensions to less symmetric situations.

Strong Coupling Isotropization Simplified

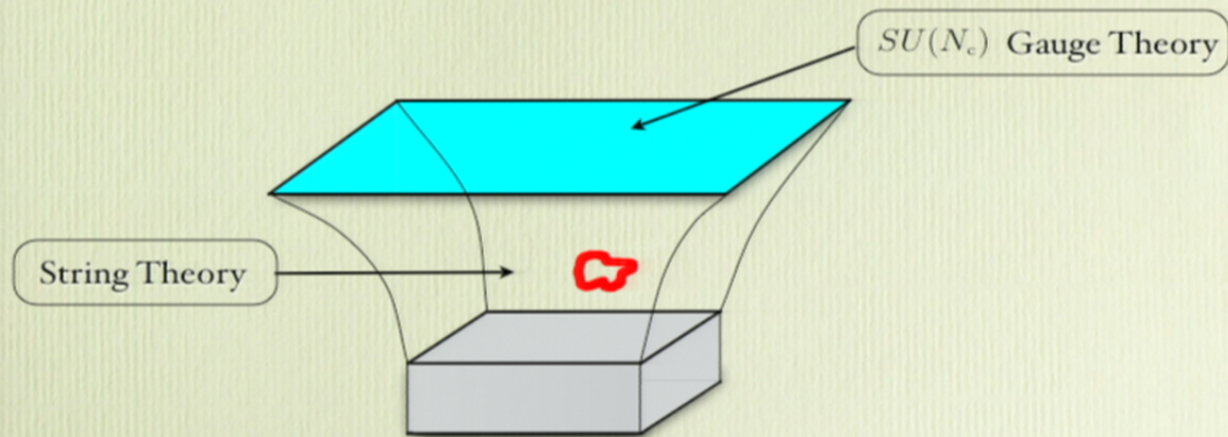


David Mateos
ICREA & University of Barcelona

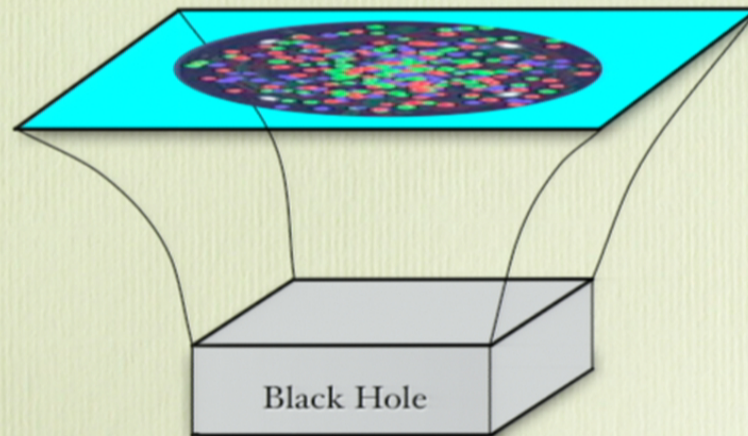
Work with M.Heller, W. van der Schee, M. Spalinski & D. Trancanelli

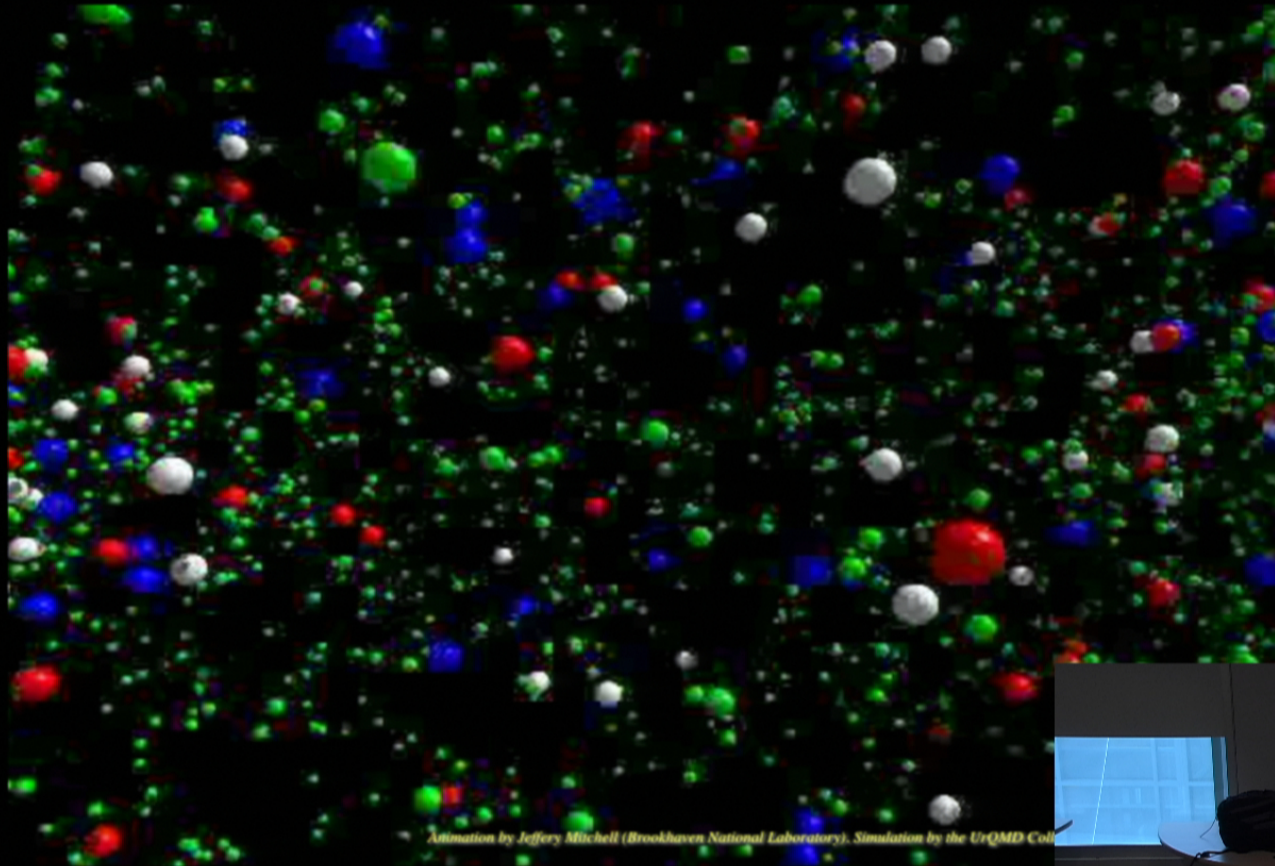
The gauge/string duality

Maldacena '97



Deconfinement (QGP) = Black Hole



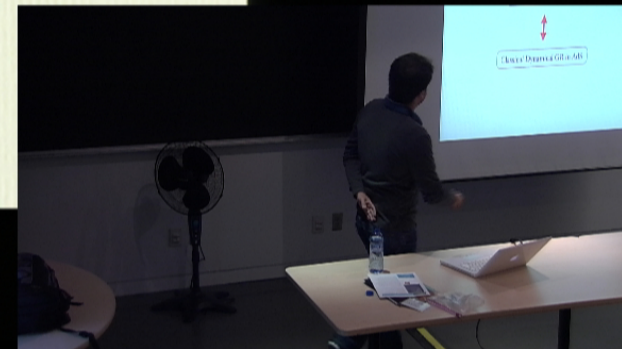


Out of equilibrium

Out-of-equilibrium QFT



Classical Dynamical GR in AdS



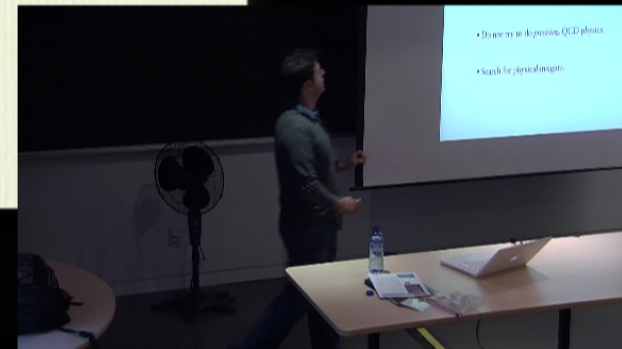
Remember

- QCD dual is beyond supergravity.



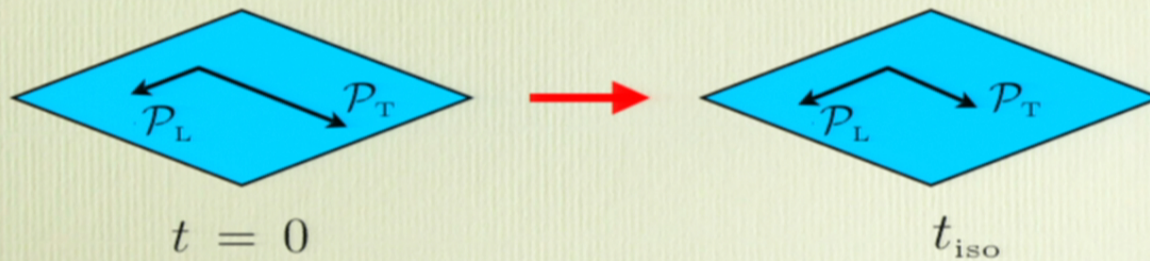
Remember

- QCD dual is beyond supergravity.
- Do not try to do *precision* QCD physics.
- Search for physical insights.



In the context of HIC

- Fast isotropization of the QGP (~ 1 fm/c) remains outstanding challenge.
- Consider simplest possible set-up in AdS/CFT:
Isotropization of homogeneous 4D CFT plasma
(e.g. $N=4$ SYM plasma).



In the context of HIC

- Fast isotropization of the QGP (~ 1 fm/c) remains outstanding challenge.
- Consider simplest possible set-up in AdS/CFT:
Isotropization of homogeneous 4D CFT plasma
(e.g. $N=4$ SYM plasma).
- Homogeneity: $\partial_\mu T^{\mu\nu} = 0 \rightarrow \partial_t T^{00} = 0 \rightarrow \mathcal{E}_i = \mathcal{E}_f$

In the context of HIC

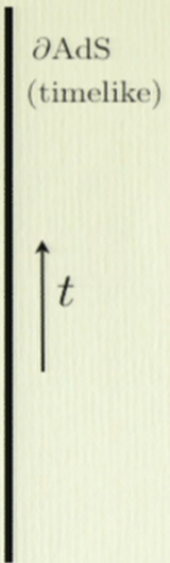
- Fast isotropization of the QGP (~ 1 fm/c) remains outstanding challenge.
- Consider simplest possible set-up in AdS/CFT:
Isotropization of homogeneous 4D CFT plasma
(e.g. $N=4$ SYM plasma).
- Homogeneity: $\partial_\mu T^{\mu\nu} = 0 \rightarrow \partial_t T^{00} = 0 \rightarrow \mathcal{E}_i = \mathcal{E}_f$

In the context of HIC

- Fast isotropization of the QGP (~ 1 fm/c) remains outstanding challenge.
- Consider simplest possible set-up in AdS/CFT:
Isotropization of homogeneous 4D CFT plasma
(e.g. $N=4$ SYM plasma).
- Homogeneity: $\partial_\mu T^{\mu\nu} = 0 \rightarrow \partial_t T^{00} = 0 \rightarrow \mathcal{E}_i = \mathcal{E}_f$
- No hydrodynamics: $\omega \rightarrow 0$ as $q \rightarrow 0$
- Only “quasi-normal modes” (QNM).



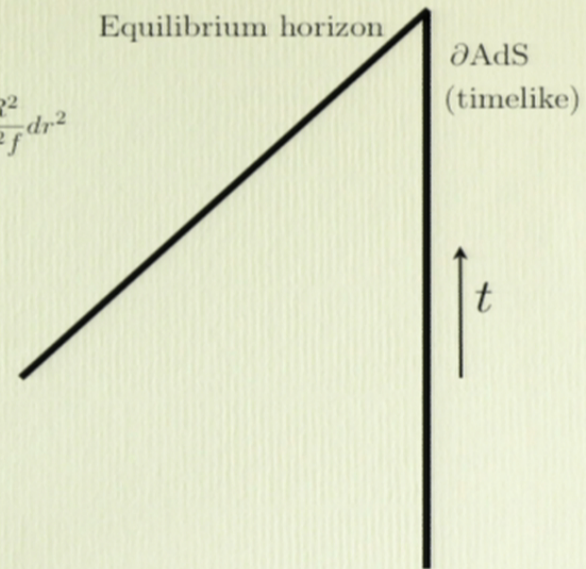
Posing the problem: Causal Structure



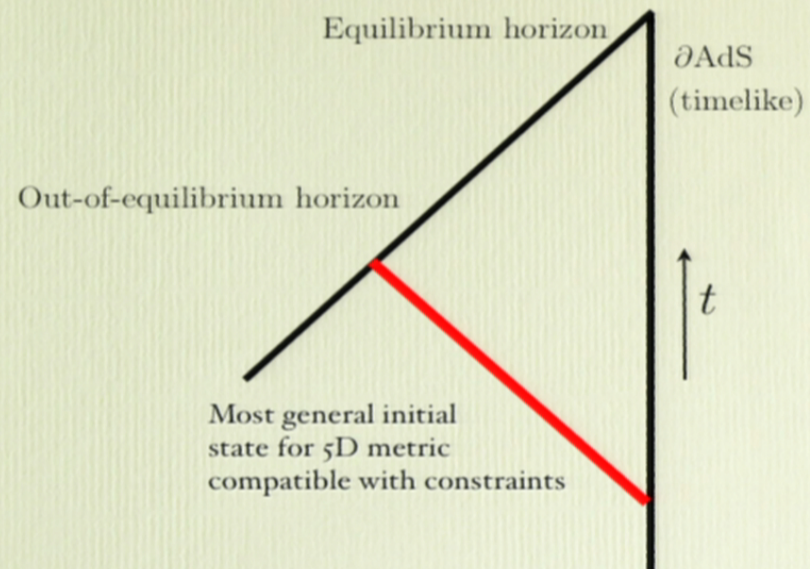
Posing the problem: Causal Structure

$$ds^2 = \frac{r^2}{R^2} (-f dt^2 + dx_1^2 + dx_2^2 + dx_3^2) + \frac{R^2}{r^2 f} dr^2$$

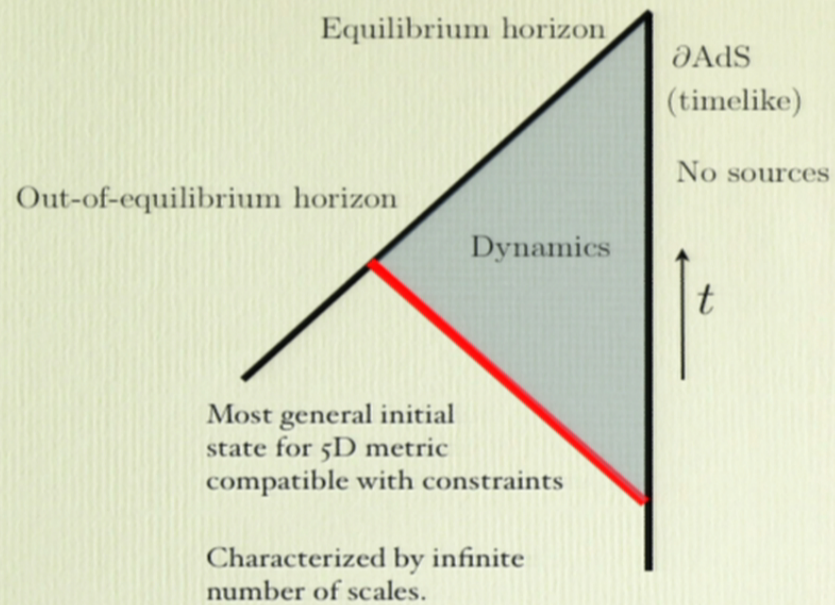
$$f(r) = 1 - \frac{r_0^4}{r^4}$$



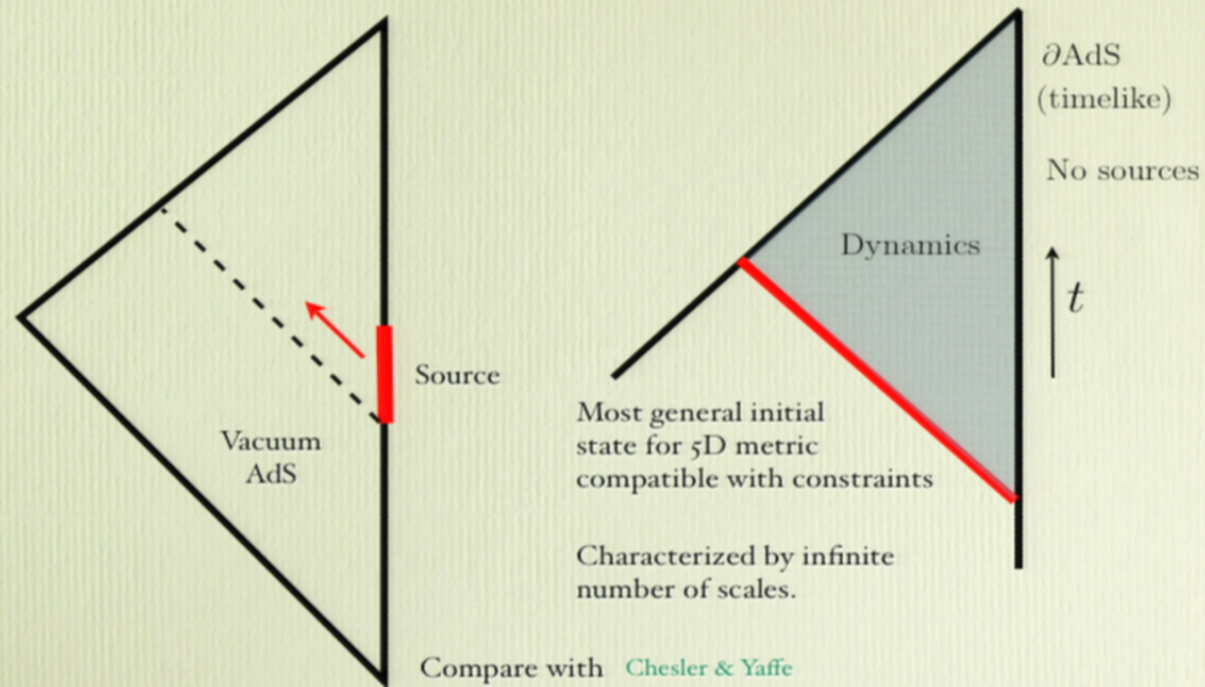
Posing the problem: Causal Structure



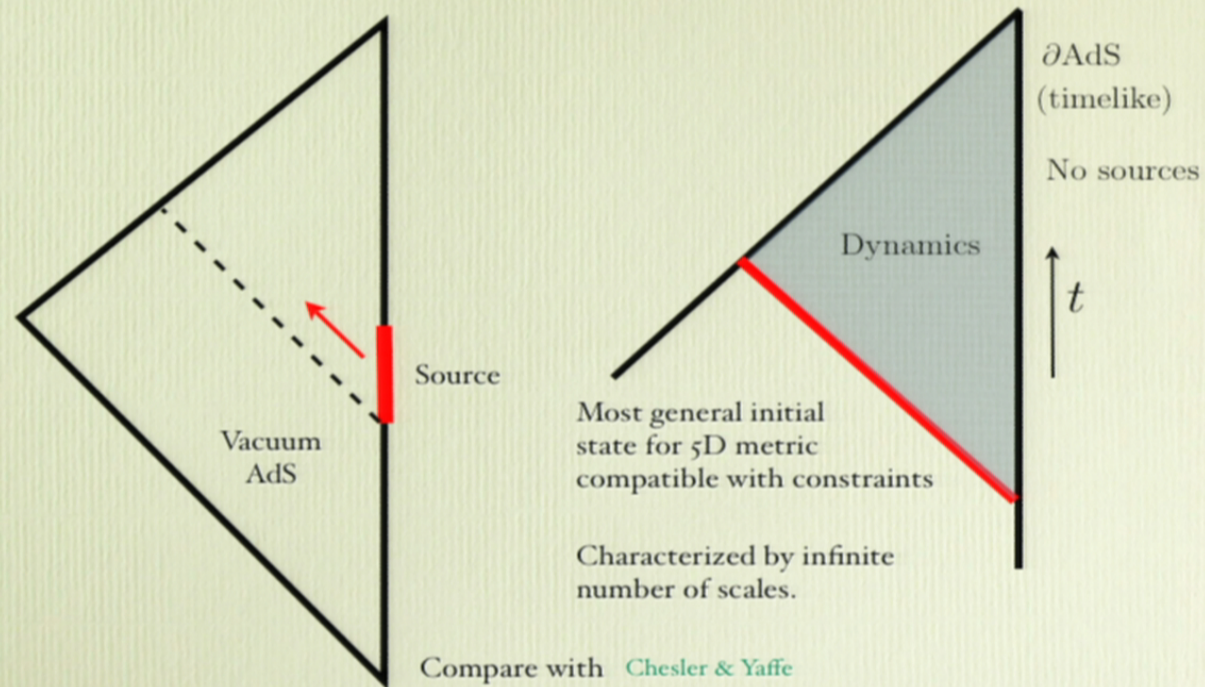
Posing the problem: Causal Structure



Posing the problem: Causal Structure



Posing the problem: Causal Structure

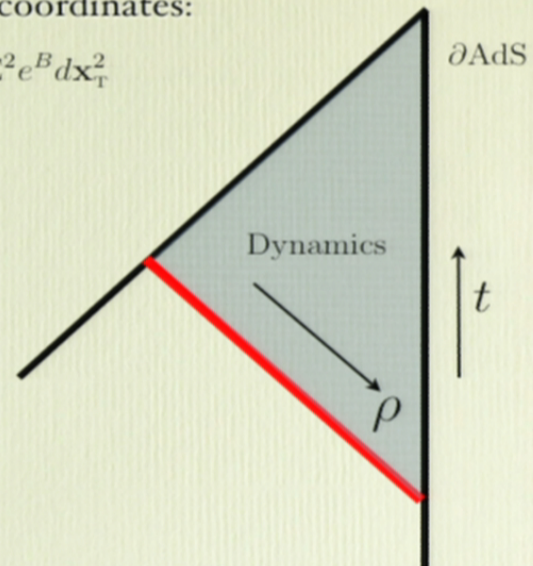


Posing the problem

- Generalized Eddington-Finkelstein coordinates:

$$ds^2 = 2dtd\rho - A dt^2 + \Sigma^2 e^{-2B} dx_L^2 + \Sigma^2 e^B d\mathbf{x}_T^2$$

A, Σ, B functions of t, ρ only.



Posing the problem

- Generalized Eddington-Finkelstein coordinates:

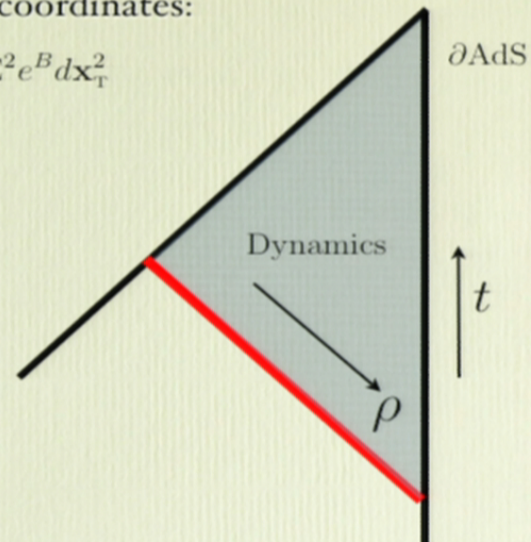
$$ds^2 = 2dtd\rho - Adt^2 + \Sigma^2 e^{-2B} dx_L^2 + \Sigma^2 e^B d\mathbf{x}_T^2$$

A, Σ, B functions of t, ρ only.

- In equilibrium:

$$A = \rho^2(1 - \rho_h^4/\rho^4), \quad \Sigma = \rho$$

$$B = 0, \quad \rho_h = \pi T$$



Posing the problem

- Generalized Eddington-Finkelstein coordinates:

$$ds^2 = 2dtd\rho - Adt^2 + \Sigma^2 e^{-2B} dx_L^2 + \Sigma^2 e^B d\mathbf{x}_T^2$$

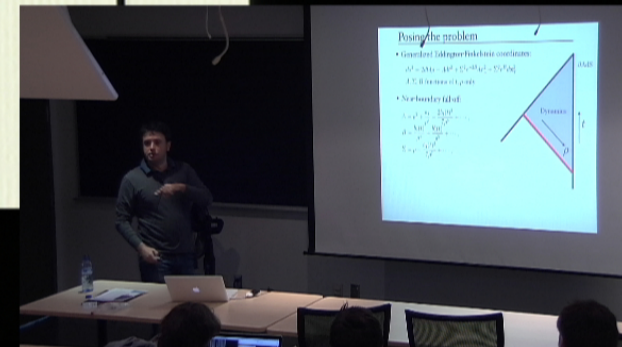
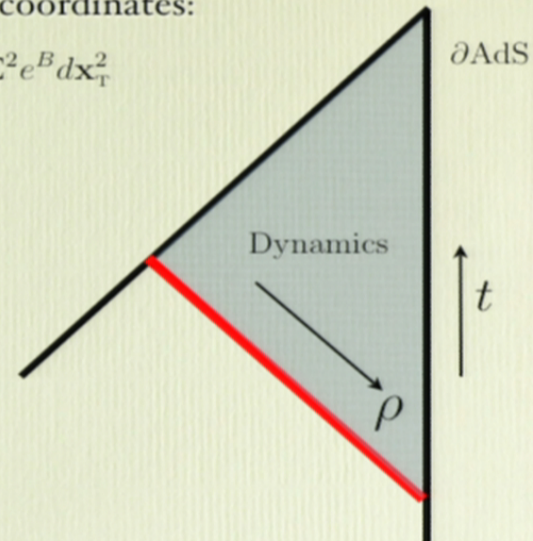
A, Σ, B functions of t, ρ only.

- Near-boundary fall-off:

$$A = \rho^2 + \frac{a_4}{\rho^2} - \frac{2b_4(t)^2}{7\rho^6} + \dots,$$

$$B = \frac{b_4(t)}{\rho^4} + \frac{b'_4(t)}{\rho^5} + \dots,$$

$$\Sigma = \rho - \frac{b_4(t)^2}{7\rho^7} + \dots,$$



Posing the problem

- Generalized Eddington-Finkelstein coordinates:

$$ds^2 = 2dtd\rho - Adt^2 + \Sigma^2 e^{-2B} dx_L^2 + \Sigma^2 e^B d\mathbf{x}_T^2$$

A, Σ, B functions of t, ρ only.

- Near-boundary fall-off:

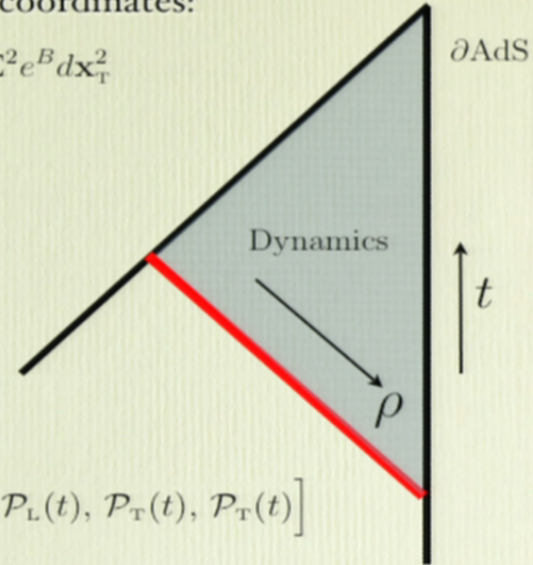
$$A = \rho^2 + \frac{a_4}{\rho^2} - \frac{2b_4(t)^2}{7\rho^6} + \dots,$$

$$B = \frac{b_4(t)}{\rho^4} + \frac{b'_4(t)}{\rho^5} + \dots,$$

$$\Sigma = \rho - \frac{b_4(t)^2}{7\rho^7} + \dots,$$

- Determines: $\langle T_{\mu\nu} \rangle = \frac{N_c^2}{2\pi^2} \text{diag}[\mathcal{E}, \mathcal{P}_L(t), \mathcal{P}_T(t), \mathcal{P}_T(t)]$

$$\mathcal{E} = -3a_4/4 \quad \text{and} \quad \Delta\mathcal{P}(t) = 3b_4(t)$$



Posing the problem

- Generalized Eddington-Finkelstein coordinates:

$$ds^2 = 2dtd\rho - Adt^2 + \Sigma^2 e^{-2B} dx_L^2 + \Sigma^2 e^B d\mathbf{x}_T^2$$

A, Σ, B functions of t, ρ only.

- Near-boundary fall-off:

$$A = \rho^2 + \frac{a_4}{\rho^2} - \frac{2b_4(t)^2}{7\rho^6} + \dots,$$

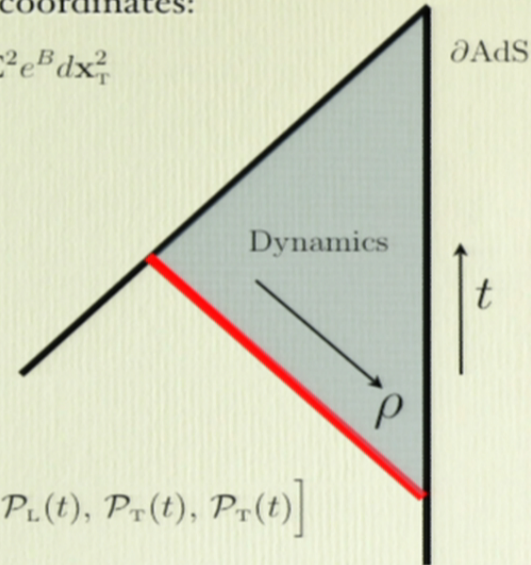
$$B = \frac{b_4(t)}{\rho^4} + \frac{b'_4(t)}{\rho^5} + \dots,$$

$$\Sigma = \rho - \frac{b_4(t)^2}{7\rho^7} + \dots,$$

- Determines: $\langle T_{\mu\nu} \rangle = \frac{N_c^2}{2\pi^2} \text{diag}[\mathcal{E}, \mathcal{P}_L(t), \mathcal{P}_T(t), \mathcal{P}_T(t)]$

$$\mathcal{E} = -3a_4/4 \quad \text{and} \quad \Delta\mathcal{P}(t) = 3b_4(t)$$

- In particular, B determines $\Delta\mathcal{P} = \mathcal{P}_T - \mathcal{P}_L$.



Posing the problem

- Generalized Eddington-Finkelstein coordinates:

$$ds^2 = 2dtd\rho - Adt^2 + \Sigma^2 e^{-2B} dx_L^2 + \Sigma^2 e^B d\mathbf{x}_T^2$$

A, Σ, B functions of t, ρ only.

- Einstein's equations:

$$0 = \Sigma (\dot{\Sigma})' + 2\Sigma' \dot{\Sigma} - 2\Sigma^2,$$

$$0 = \Sigma (\dot{B})' + \frac{3}{2}(\Sigma' \dot{B} + B' \dot{\Sigma}),$$

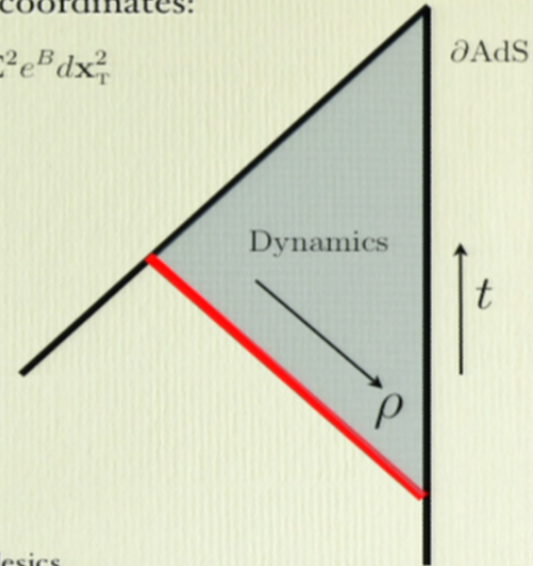
$$0 = A'' + 3B' \dot{B} - 12\Sigma' \dot{\Sigma}/\Sigma^2 + 4,$$

$$0 = \ddot{\Sigma} + \frac{1}{2}(\dot{B}^2 \Sigma - A' \dot{\Sigma}),$$

$$0 = \Sigma'' + \frac{1}{2}B'^2 \Sigma,$$

Derivatives along *ingoing* and *outgoing* null geodesics.

$$h' \equiv \partial_r h \quad \dot{h} \equiv \partial_t h + \frac{1}{2}A \partial_r h$$



Posing the problem

- Generalized Eddington-Finkelstein coordinates:

$$ds^2 = 2dt d\rho - A dt^2 + \Sigma^2 e^{-2B} dx_L^2 + \Sigma^2 e^B d\mathbf{x}_T^2$$

A, Σ, B functions of t, ρ only.

- Einstein's equations:

$$0 = \Sigma (\dot{\Sigma})' + 2\Sigma' \dot{\Sigma} - 2\Sigma^2,$$

$$0 = \Sigma (\dot{B})' + \frac{3}{2}(\Sigma' \dot{B} + B' \dot{\Sigma}),$$

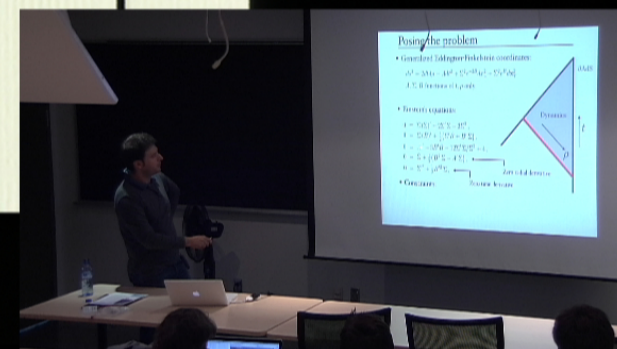
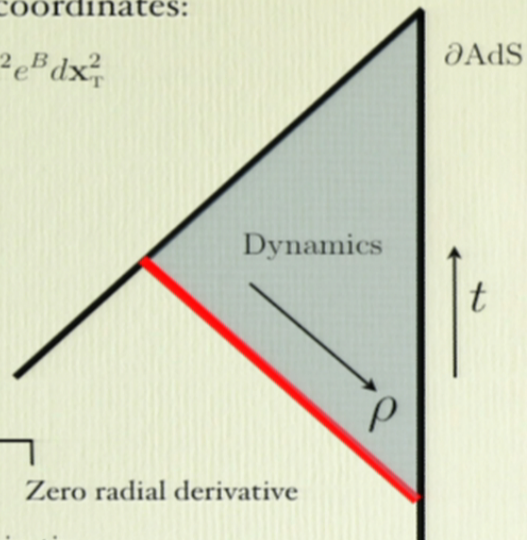
$$0 = A'' + 3B' \dot{B} - 12\Sigma' \dot{\Sigma} / \Sigma^2 + 4,$$

$$0 = \ddot{\Sigma} + \frac{1}{2}(\dot{B}^2 \Sigma - A' \dot{\Sigma}),$$

$$0 = \Sigma'' + \frac{1}{2}B'^2 \Sigma,$$

- Constraints:

Zero time derivative



Posing the problem

- Generalized Eddington-Finkelstein coordinates:

$$ds^2 = 2dt d\rho - A dt^2 + \Sigma^2 e^{-2B} dx_L^2 + \Sigma^2 e^B d\mathbf{x}_T^2$$

A, Σ, B functions of t, ρ only.

- Einstein's equations:

$$0 = \Sigma (\dot{\Sigma})' + 2\Sigma' \dot{\Sigma} - 2\Sigma^2,$$

$$0 = \Sigma (\dot{B})' + \frac{3}{2}(\Sigma' \dot{B} + B' \dot{\Sigma}),$$

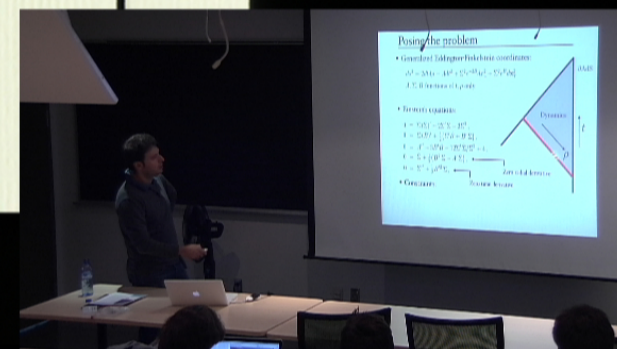
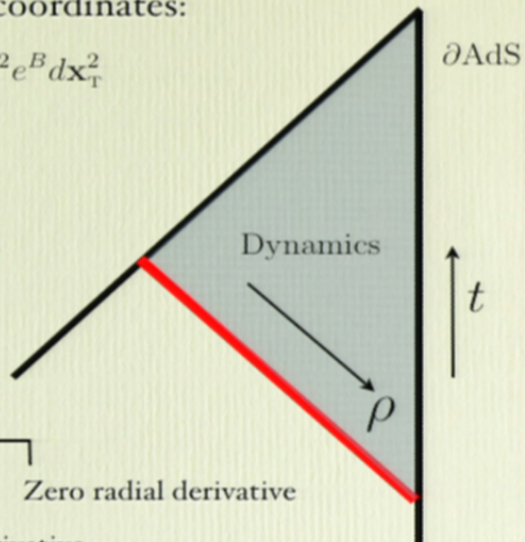
$$0 = A'' + 3B' \dot{B} - 12\Sigma' \dot{\Sigma} / \Sigma^2 + 4,$$

$$0 = \ddot{\Sigma} + \frac{1}{2}(\dot{B}^2 \Sigma - A' \dot{\Sigma}),$$

$$0 = \Sigma'' + \frac{1}{2}B'^2 \Sigma,$$

- Constraints:

Zero time derivative



Posing the problem

- Generalized Eddington-Finkelstein coordinates:

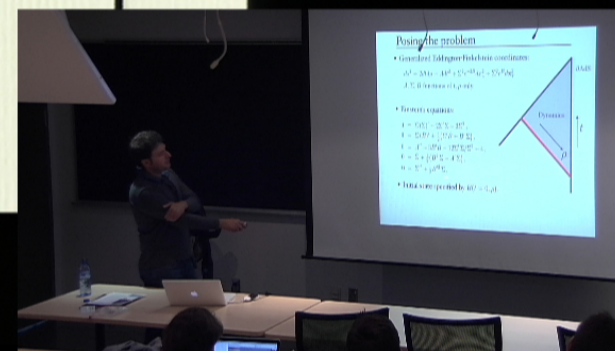
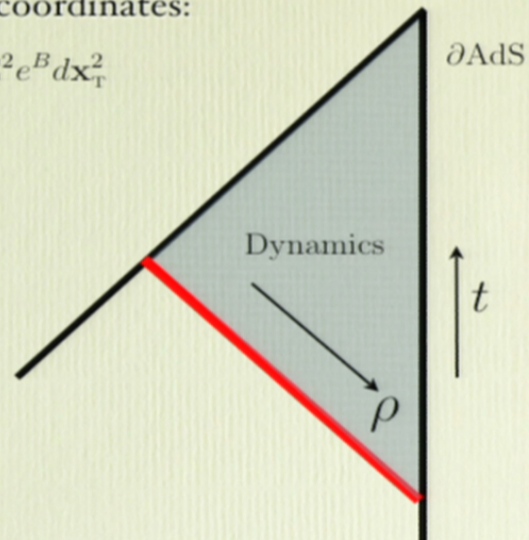
$$ds^2 = 2dt d\rho - A dt^2 + \Sigma^2 e^{-2B} dx_L^2 + \Sigma^2 e^B d\mathbf{x}_T^2$$

A, Σ, B functions of t, ρ only.

- Einstein's equations:

$$\begin{aligned} 0 &= \Sigma (\dot{\Sigma})' + 2\Sigma' \dot{\Sigma} - 2\Sigma^2, \\ 0 &= \Sigma (\dot{B})' + \frac{3}{2}(\Sigma' \dot{B} + B' \dot{\Sigma}), \\ 0 &= A'' + 3B' \dot{B} - 12\Sigma' \dot{\Sigma} / \Sigma^2 + 4, \\ 0 &= \ddot{\Sigma} + \frac{1}{2}(\dot{B}^2 \Sigma - A' \dot{\Sigma}), \\ 0 &= \Sigma'' + \frac{1}{2}B'^2 \Sigma, \end{aligned}$$

- Initial state specified by $B(t = 0, \rho)$.



Posing the problem

- Generalized Eddington-Finkelstein coordinates:

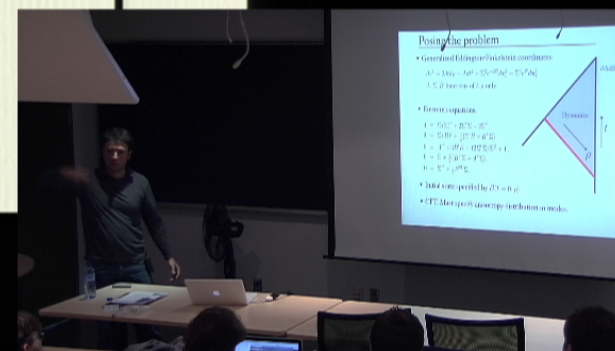
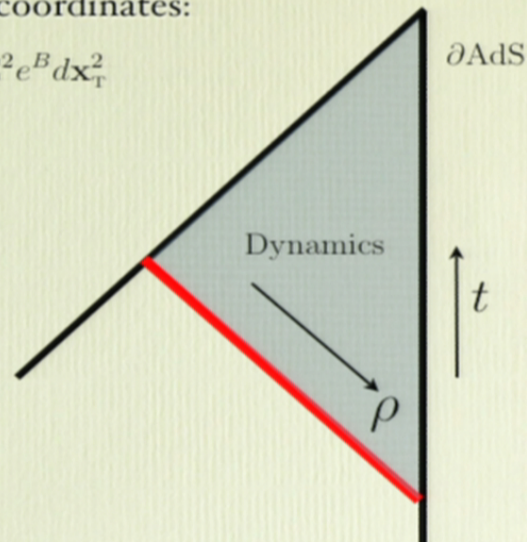
$$ds^2 = 2dt d\rho - A dt^2 + \Sigma^2 e^{-2B} dx_L^2 + \Sigma^2 e^B d\mathbf{x}_T^2$$

A, Σ, B functions of t, ρ only.

- Einstein's equations:

$$\begin{aligned} 0 &= \Sigma (\dot{\Sigma})' + 2\Sigma' \dot{\Sigma} - 2\Sigma^2, \\ 0 &= \Sigma (\dot{B})' + \frac{3}{2}(\Sigma' \dot{B} + B' \dot{\Sigma}), \\ 0 &= A'' + 3B' \dot{B} - 12\Sigma' \dot{\Sigma} / \Sigma^2 + 4, \\ 0 &= \ddot{\Sigma} + \frac{1}{2}(\dot{B}^2 \Sigma - A' \dot{\Sigma}), \\ 0 &= \Sigma'' + \frac{1}{2}B'^2 \Sigma, \end{aligned}$$

- Initial state specified by $B(t = 0, \rho)$.
- CFT: Must specify anisotropy distribution in modes.



Posing the problem

- Generalized Eddington-Finkelstein coordinates:

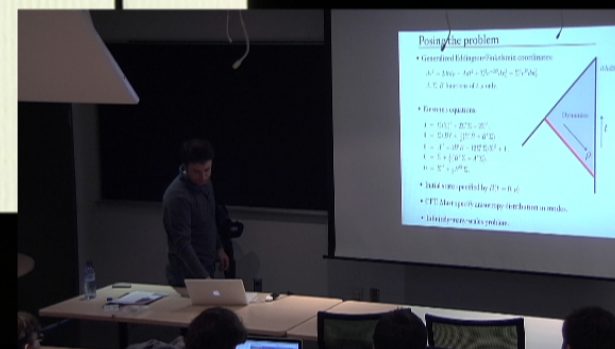
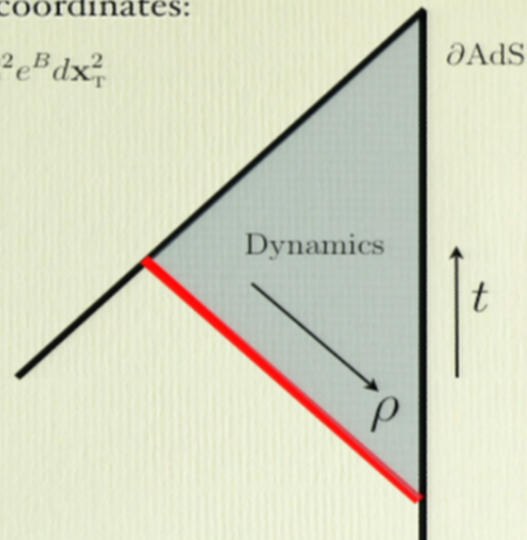
$$ds^2 = 2dt d\rho - A dt^2 + \Sigma^2 e^{-2B} dx_L^2 + \Sigma^2 e^B d\mathbf{x}_T^2$$

A, Σ, B functions of t, ρ only.

- Einstein's equations:

$$\begin{aligned} 0 &= \Sigma (\dot{\Sigma})' + 2\Sigma' \dot{\Sigma} - 2\Sigma^2, \\ 0 &= \Sigma (\dot{B})' + \frac{3}{2}(\Sigma' \dot{B} + B' \dot{\Sigma}), \\ 0 &= A'' + 3B' \dot{B} - 12\Sigma' \dot{\Sigma} / \Sigma^2 + 4, \\ 0 &= \ddot{\Sigma} + \frac{1}{2}(\dot{B}^2 \Sigma - A' \dot{\Sigma}), \\ 0 &= \Sigma'' + \frac{1}{2}B'^2 \Sigma, \end{aligned}$$

- Initial state specified by $B(t = 0, \rho)$.
- CFT: Must specify anisotropy distribution in modes.
- Infinitely-many-scales problem.



Time evolution

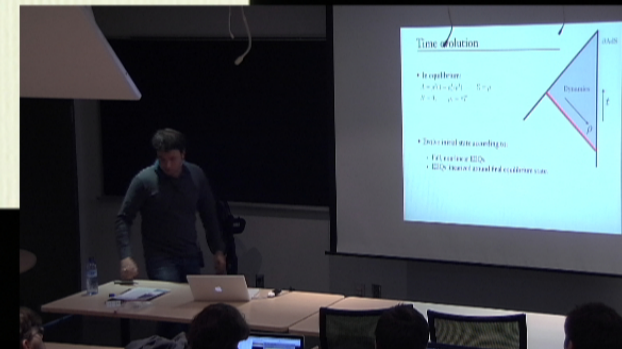
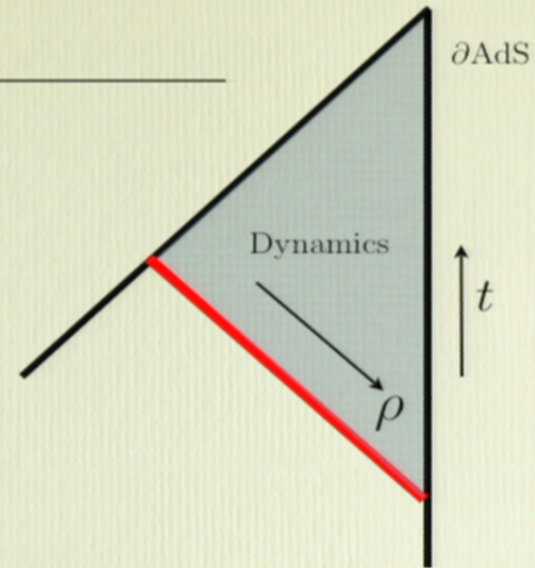
- In equilibrium:

$$A = \rho^2(1 - \rho_h^4/\rho^4), \quad \Sigma = \rho$$

$$B = 0, \quad \rho_h = \pi T$$

- Evolve initial state according to:

- Full, non-linear EEQs.
- EEQs linearized around final equilibrium state.



Motivation: Close-limit Approximation

Price & Pullin '94

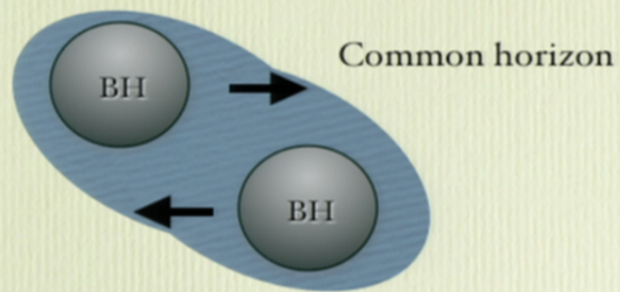
- BH collision in asymptotically flat 4D general relativity:



Motivation: Close-limit Approximation

Price & Pullin '94

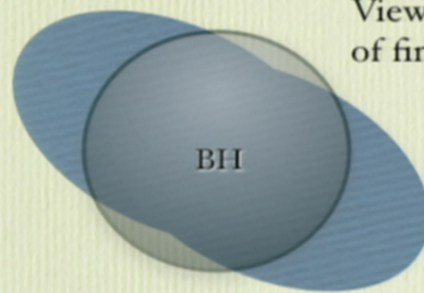
- BH collision in asymptotically flat 4D general relativity:



Motivation: Close-limit Approximation

Price & Pullin '94

- BH collision in asymptotically flat 4D general relativity:



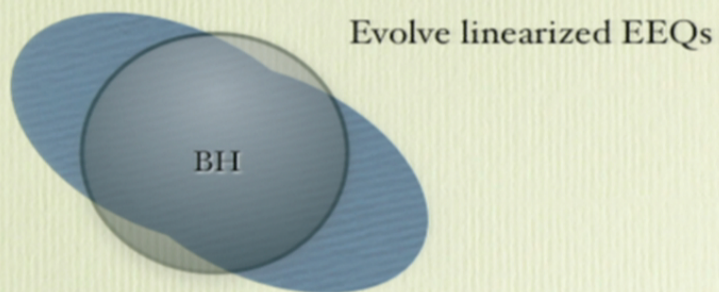
View as perturbation
of final BH



Motivation: Close-limit Approximation

Price & Pullin '94

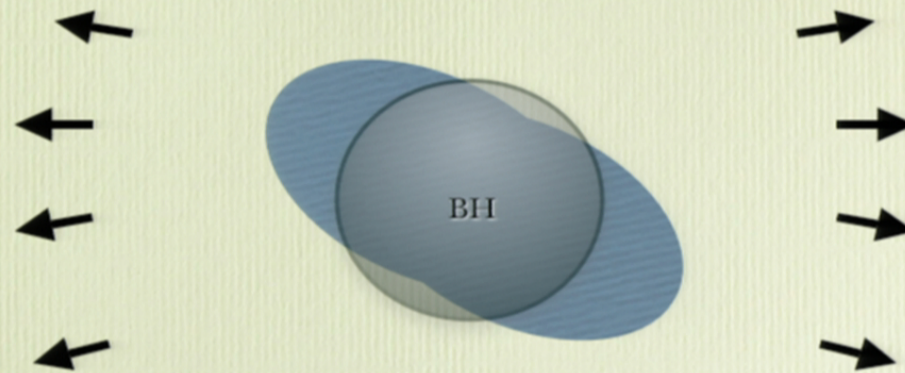
- BH collision in asymptotically flat 4D general relativity:



Motivation: Close-limit Approximation

Price & Pullin '94

- BH collision in asymptotically flat 4D general relativity:

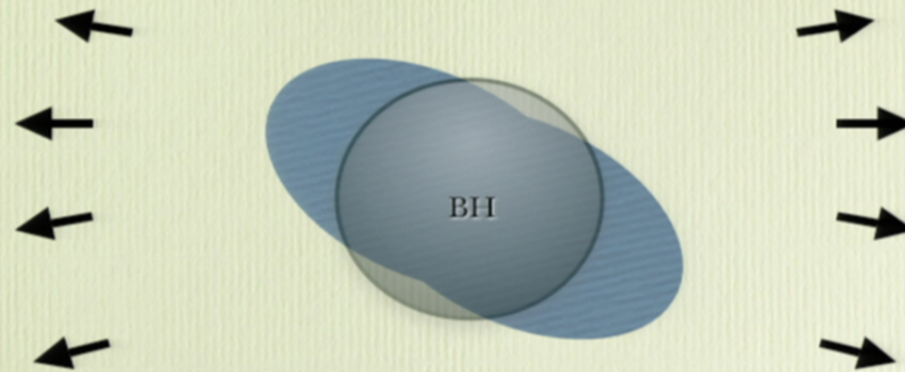


- Wave-form at infinity accurately reproduced (but perhaps non-asymptotic properties would not be).

Motivation: Close-limit Approximation

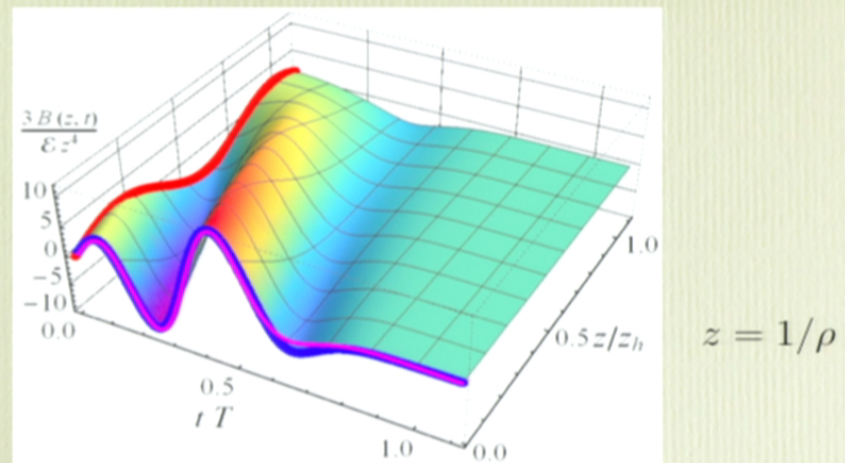
Price & Pullin '94

- BH collision in asymptotically flat 4D general relativity:

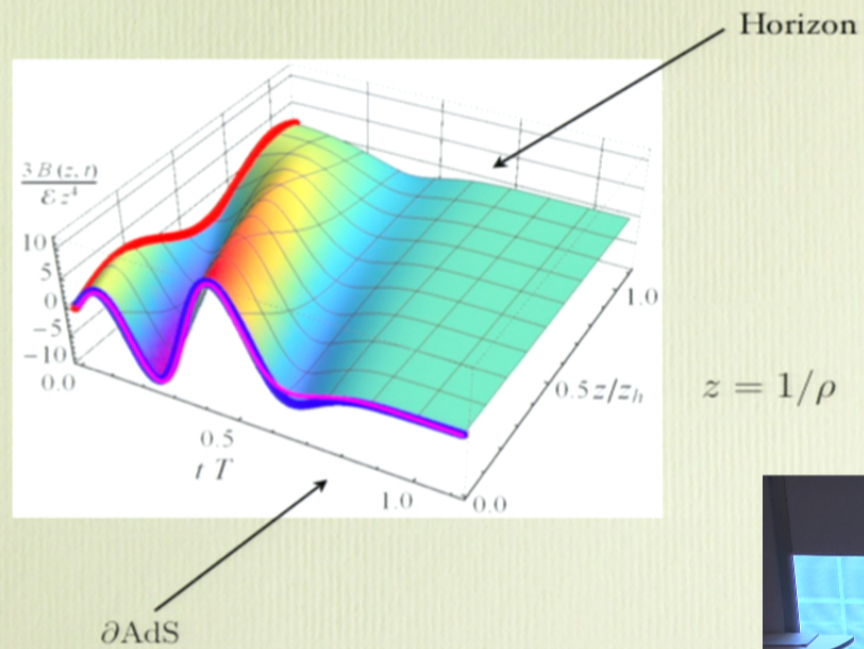


- Wave-form at infinity accurately reproduced (but perhaps non-asymptotic properties would not be).
- Analog in AdS: Boundary stress-tensor.

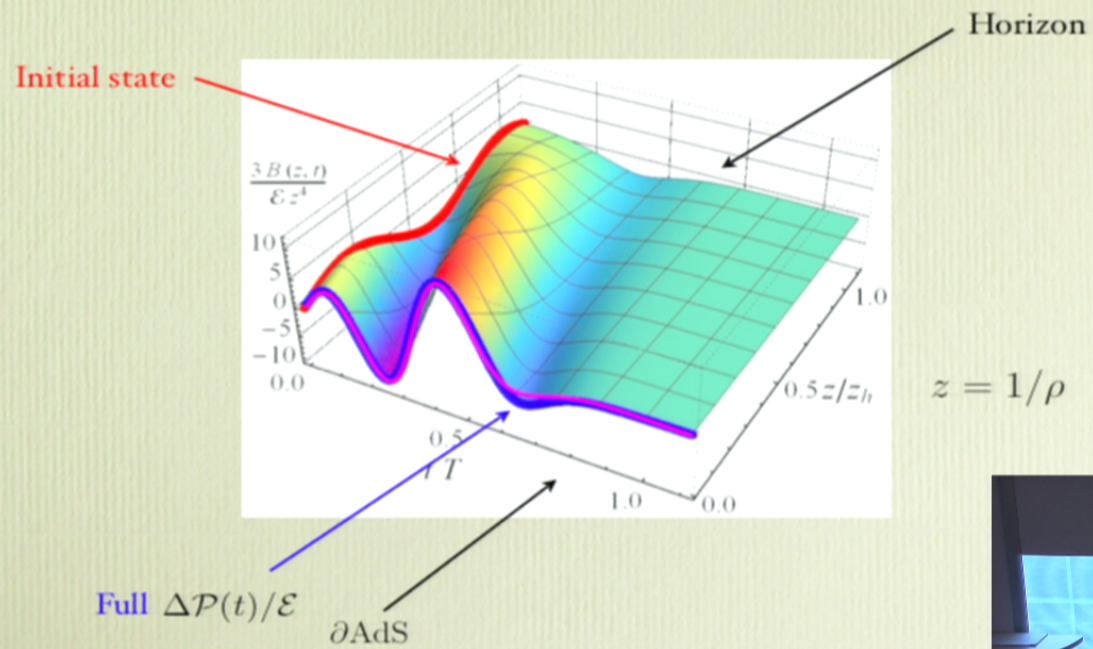
Results



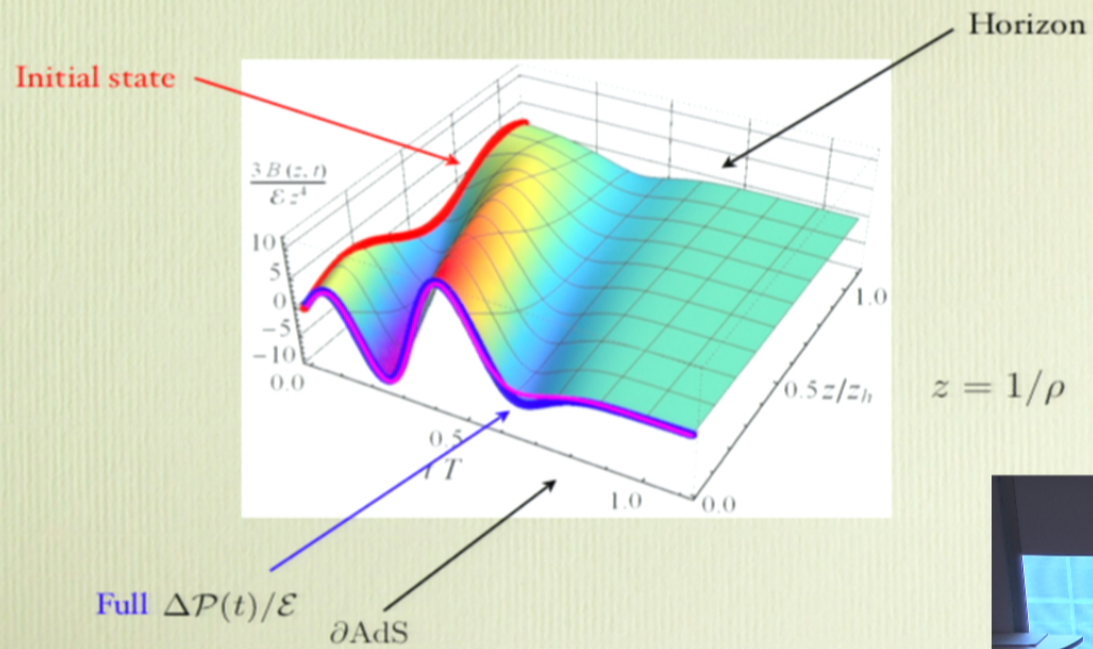
Results



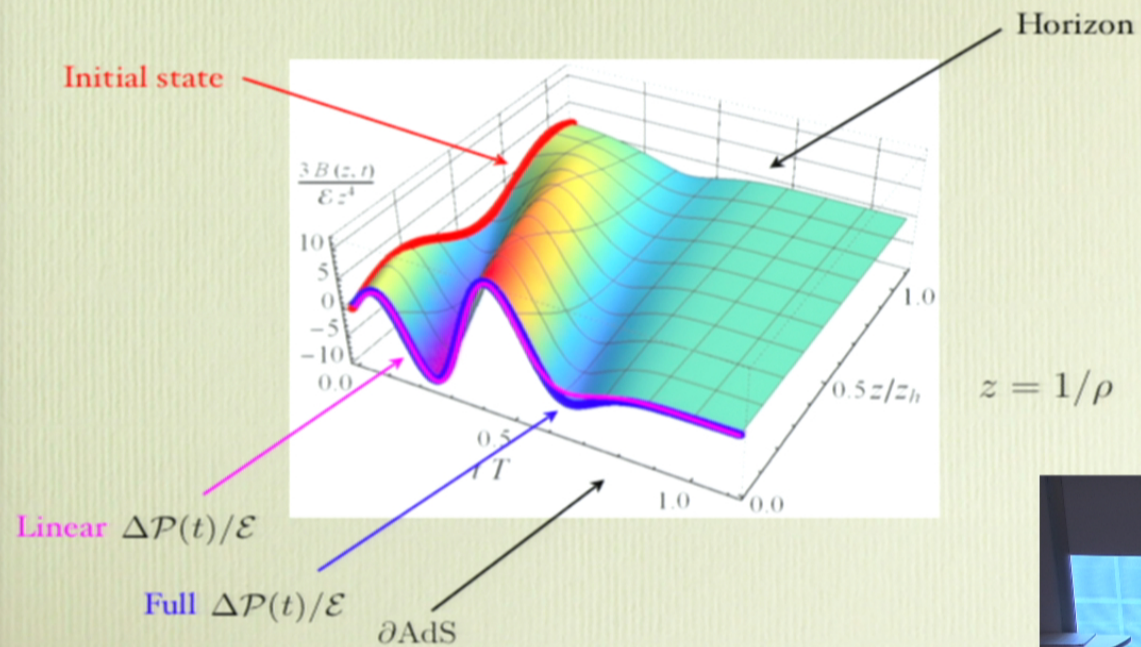
Results



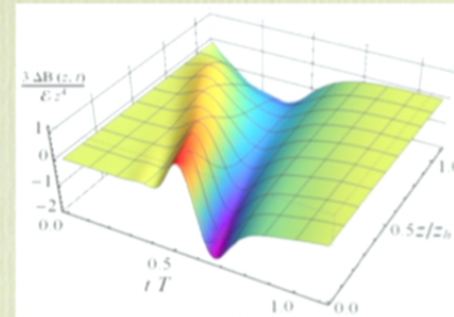
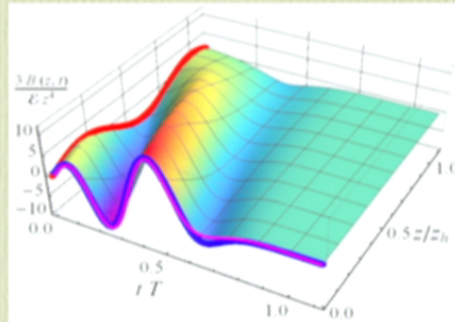
Results



Results



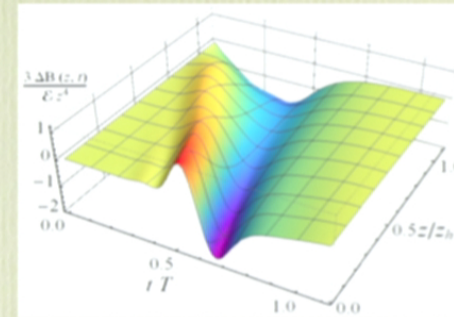
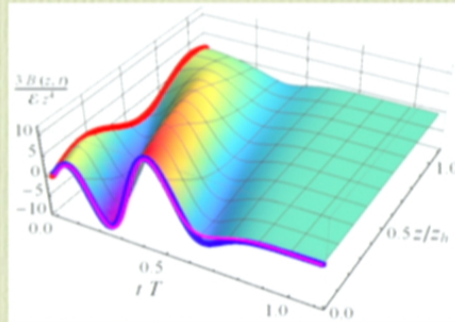
Results



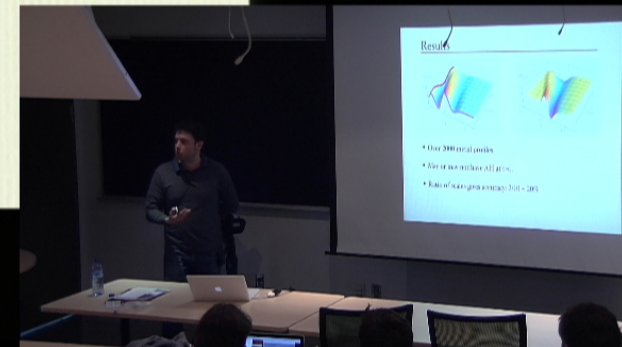
- Over 2000 initial profiles.



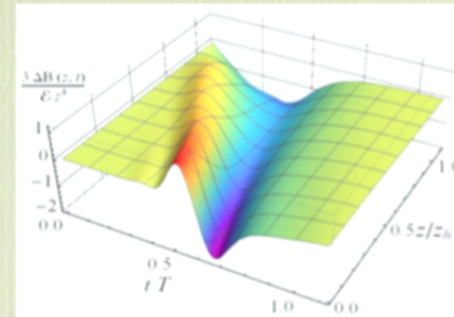
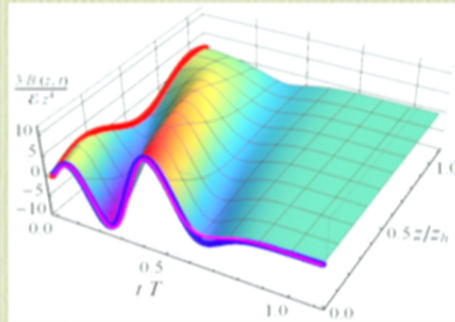
Results



- Over 2000 initial profiles.
- May or may not have AH at $t=0$.
- Ratio of scales gives accuracy: $2/10 \sim 20\%$



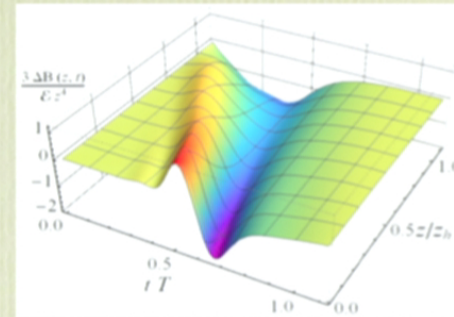
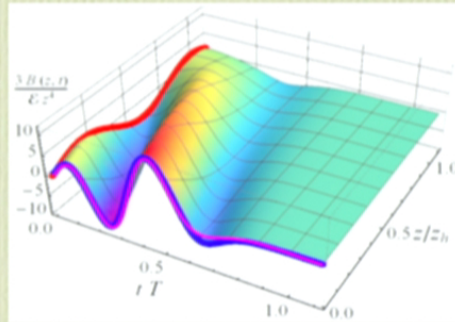
Results



- Over 2000 initial profiles.
- May or may not have AH at $t=0$.
- Ratio of scales gives accuracy: $2/10 \sim 20\%$



Results

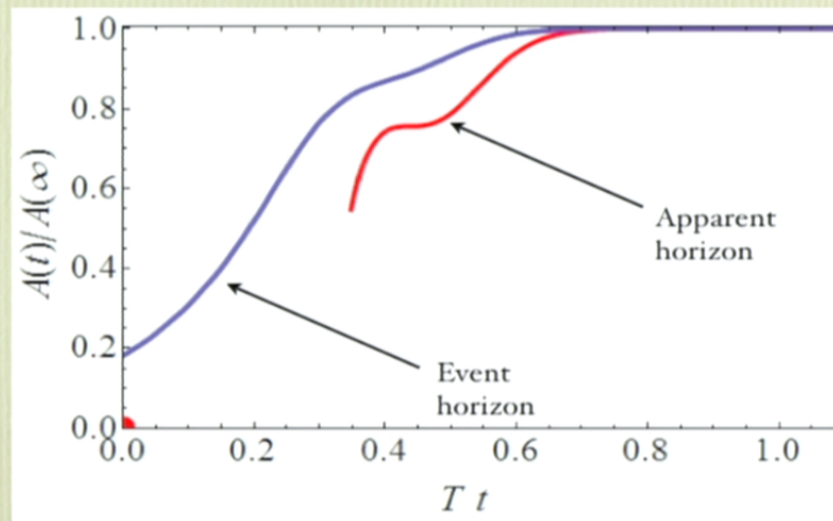


- Over 2000 initial profiles.
- May or may not have AH at $t=0$.
- Ratio of scales gives accuracy: $2/10 \sim 20\%$
- $\Delta \mathcal{P}(t)/\mathcal{E} \sim 10$ implies far from equilibrium.



Results

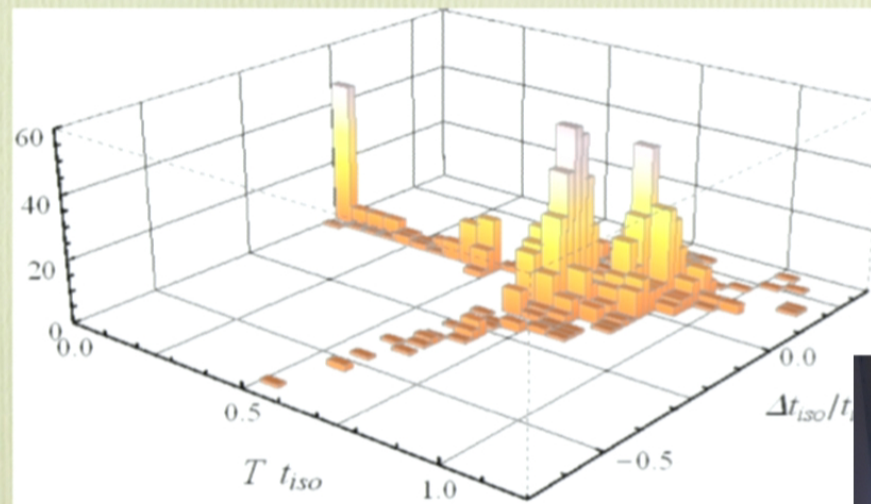
- “Entropy” increases during isotropization.



Results

- Isotropization time of order $1/T$ predicted by LA within 20%.

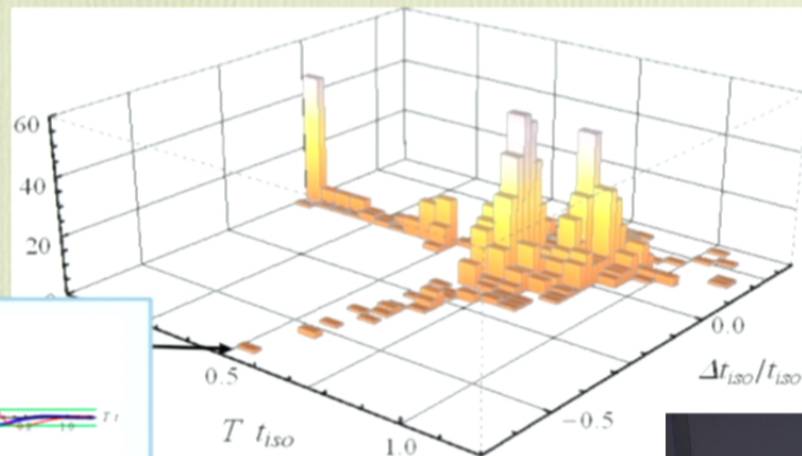
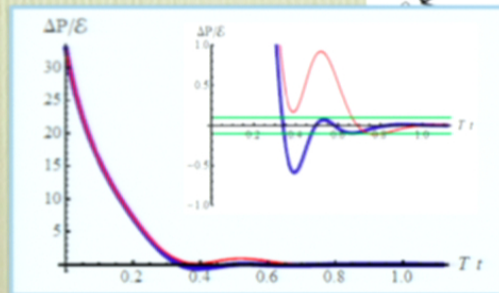
$$\Delta \mathcal{P}(t)/\mathcal{E} \leq 0.1$$

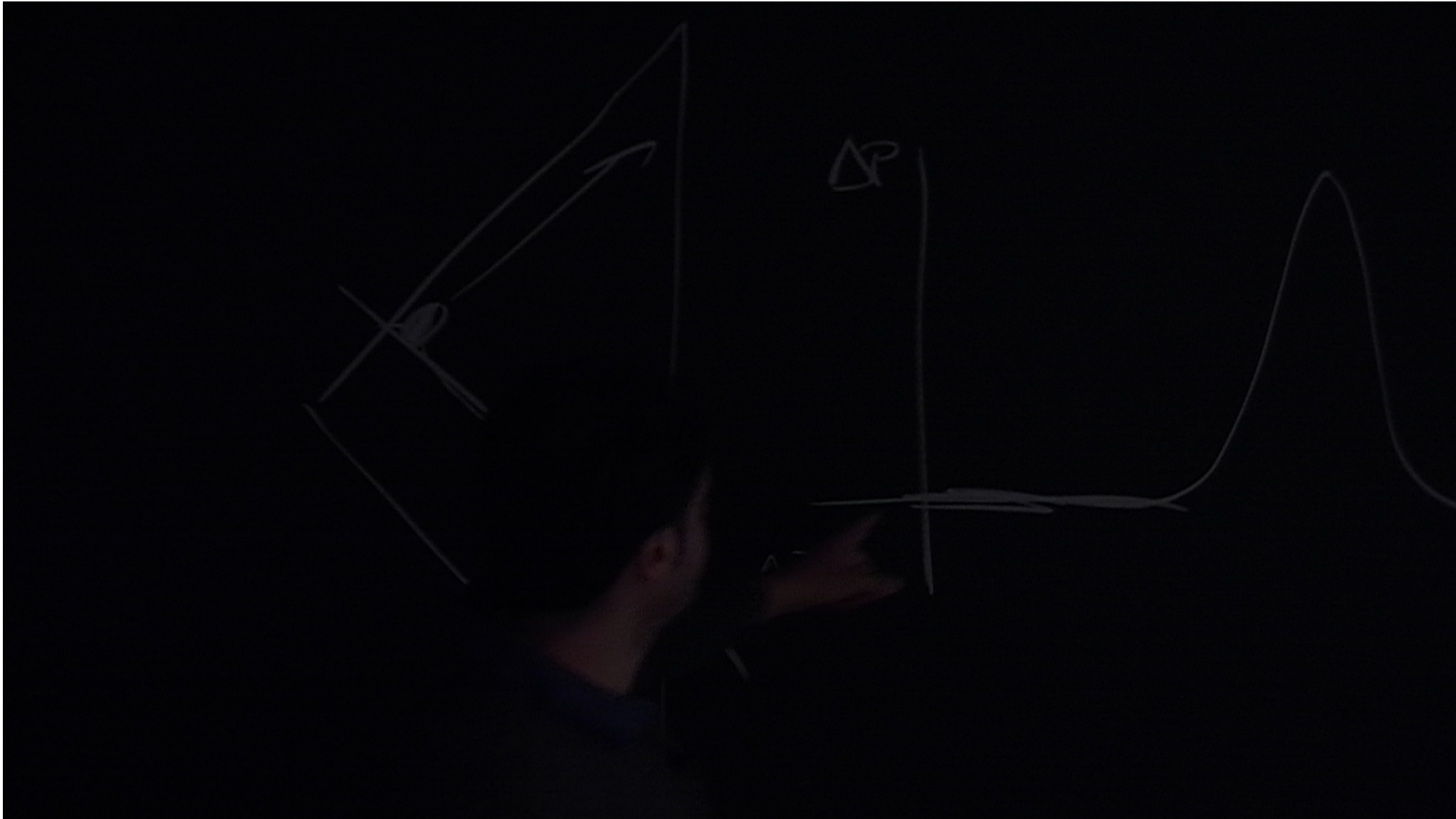


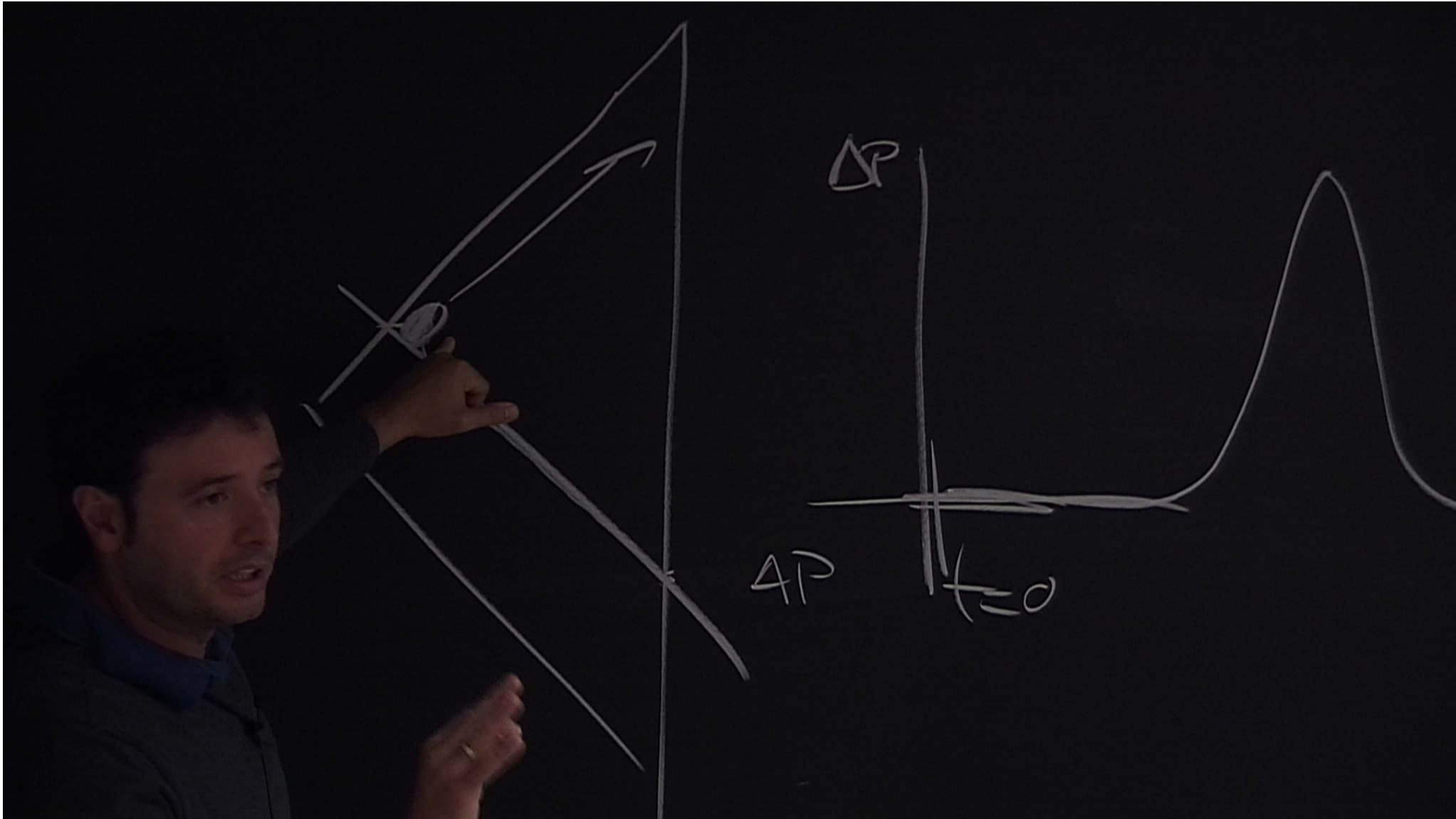
Results

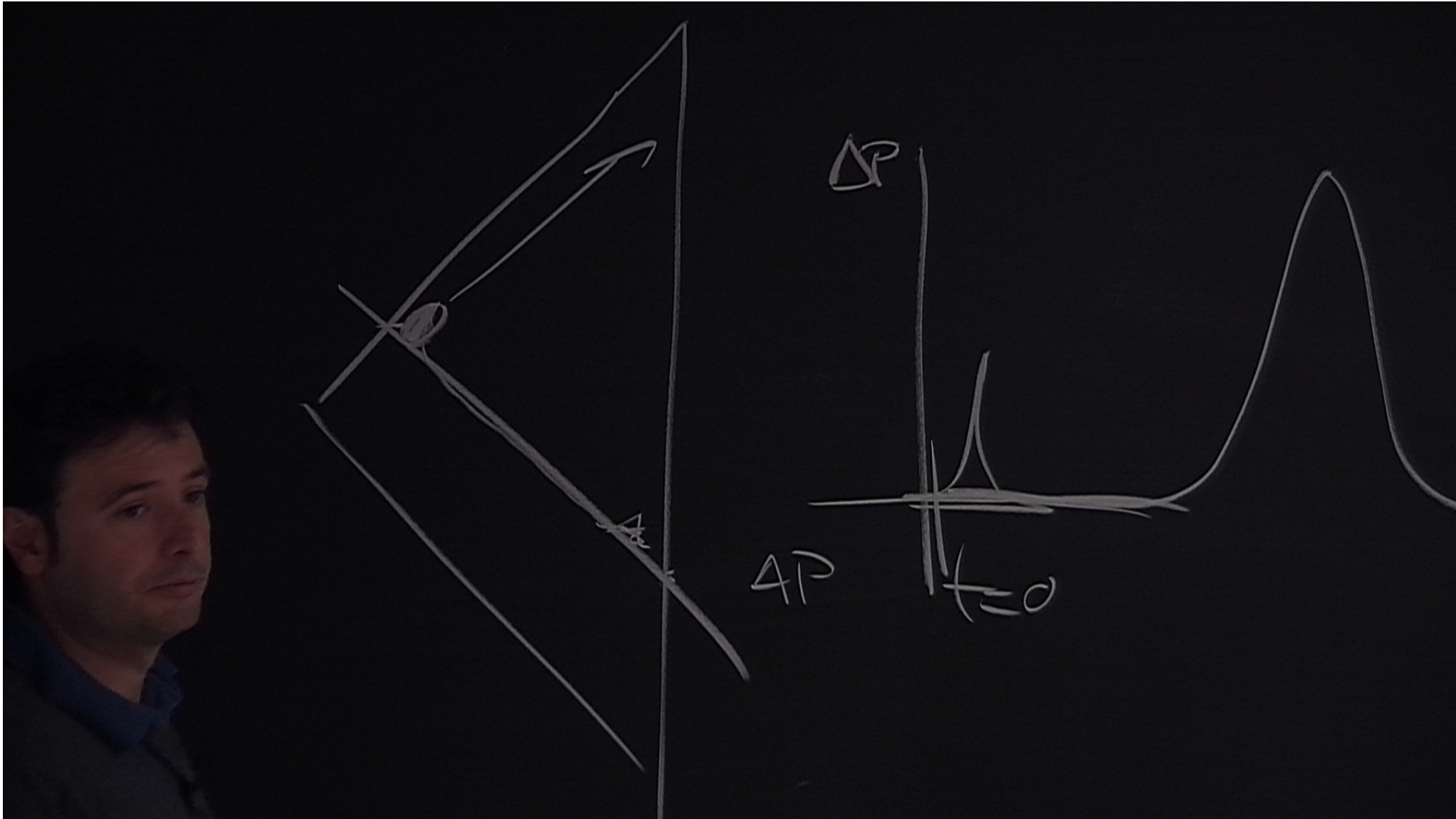
- Isotropization time of order $1/T$ predicted by LA within 20%.

$$\Delta P(t)/\mathcal{E} \leq 0.1$$



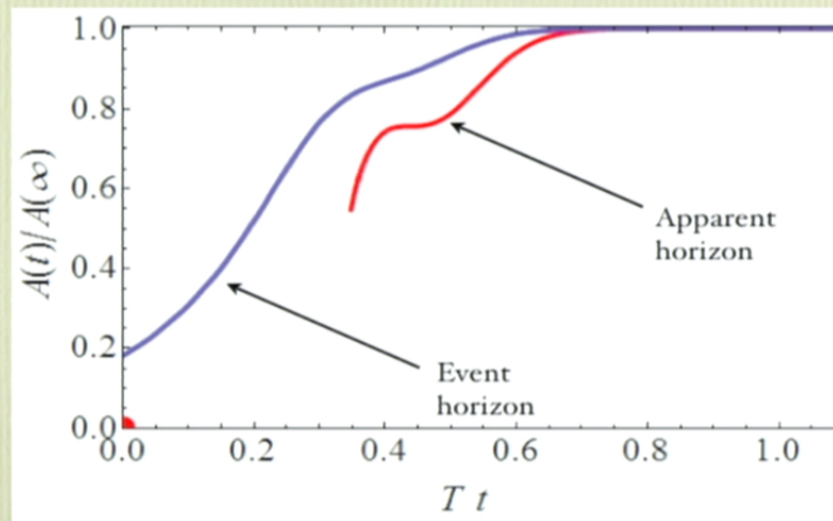




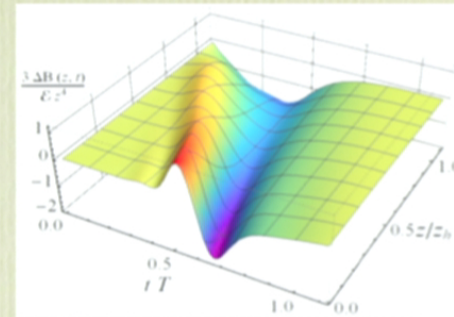
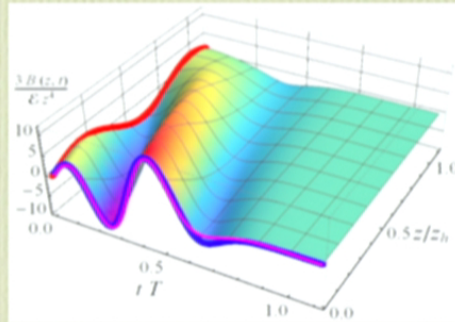


Results

- “Entropy” increases during isotropization.



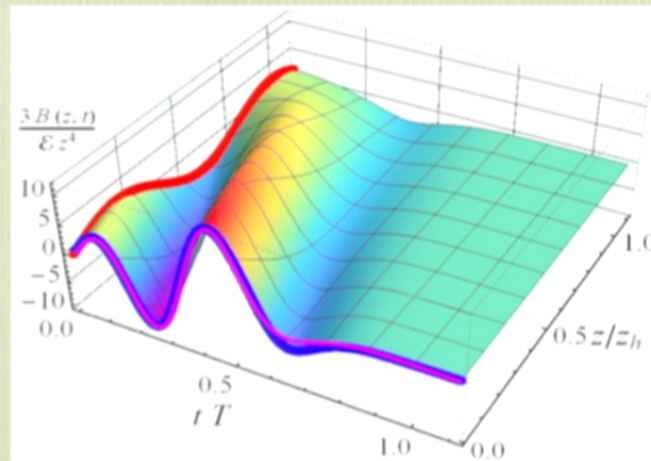
Results



- Over 2000 initial profiles.

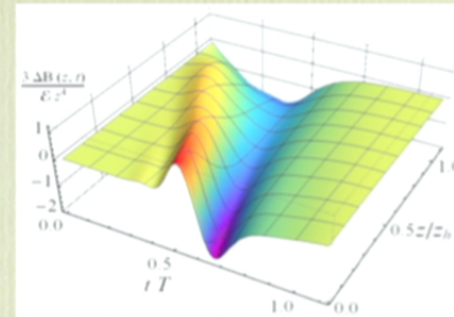
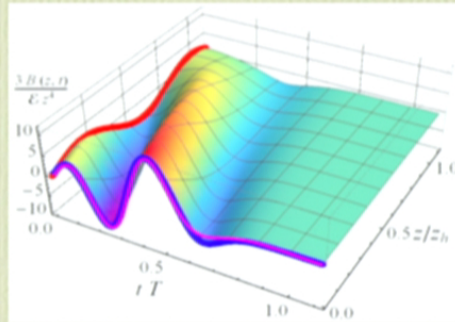


Results



$$z = 1/\rho$$

Results



- Over 2000 initial profiles.



Expand in QNMs

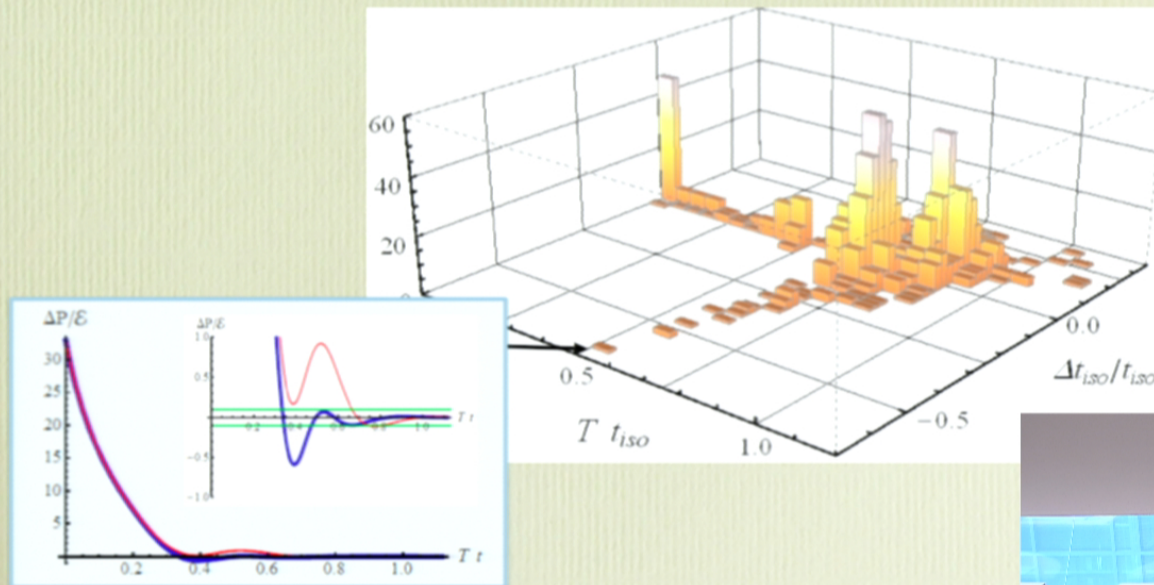
$$\delta B(t, r) = \sum_i c_i b_i(r) e^{-i\omega_i t}$$



Results

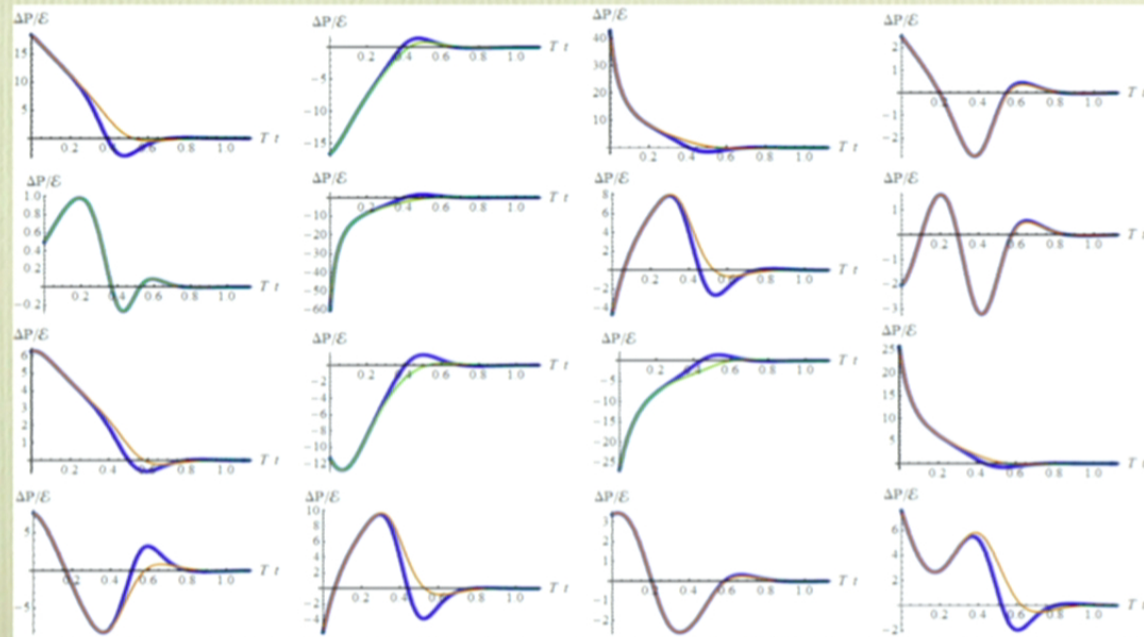
- Isotropization time of order $1/T$ predicted by LA within 20%.

$$\Delta P(t)/\mathcal{E} \leq 0.1$$



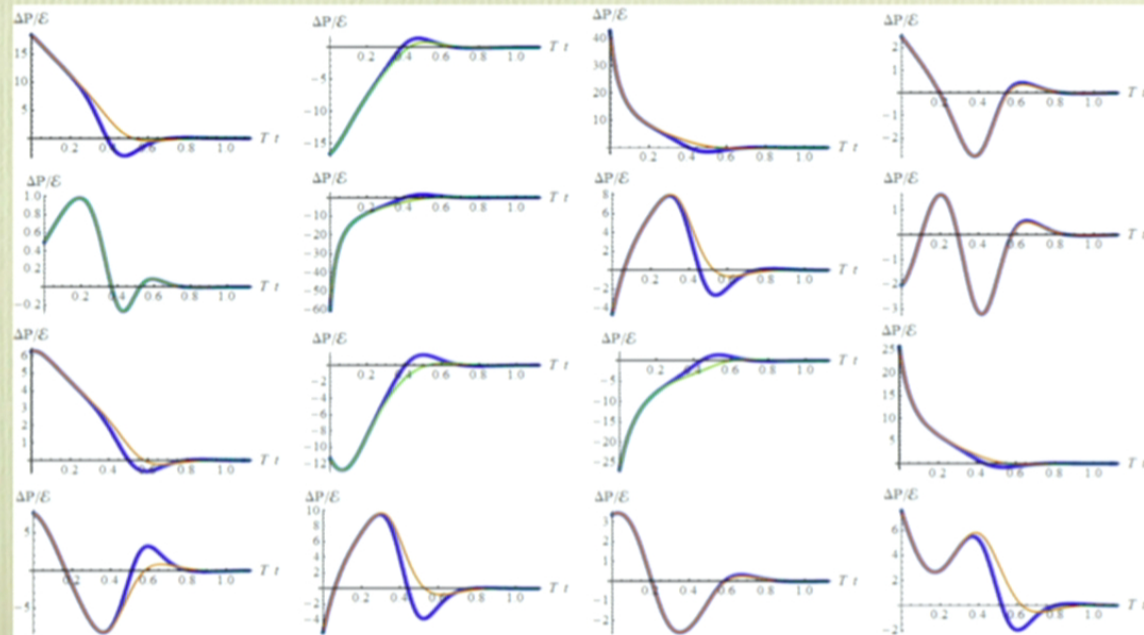
Expand in QNMs (Full / Linear / QNM)

$$\delta B(t, r) = \sum_i c_i b_i(r) e^{-i\omega_i t}$$



Expand in QNMs (Full / Linear / QNM)

$$\delta B(t, r) = \sum_i c_i b_i(r) e^{-i\omega_i t}$$



Discussion

- Gauge: Small perturbations around equilibrium plasma
→ Linear response theory

Discussion

- Gauge: Small perturbations around equilibrium plasma
→ Linear response theory
- Gravity: Small perturbations around equilibrium black hole
→ Linearized Einstein's equations
- In both cases expect linear approx. if $\Delta\mathcal{P}/\mathcal{E} \ll 1$.



Discussion

- Gauge: Small perturbations around equilibrium plasma
→ Linear response theory
- Gravity: Small perturbations around equilibrium black hole
→ Linearized Einstein's equations
- In both cases expect linear approx. if $\Delta\mathcal{P}/\mathcal{E} \ll 1$.
- In Fourier space:
 - HDMs: $\omega \rightarrow 0$ as $q \rightarrow 0$
 - QNMs: $\omega(0) \neq 0$



Discussion

- Gauge: Small perturbations around equilibrium plasma
→ Linear response theory
- Gravity: Small perturbations around equilibrium black hole
→ Linearized Einstein's equations
- In both cases expect linear approx. if $\Delta\mathcal{P}/\mathcal{E} \ll 1$.
- In Fourier space:
 - HDMs: $\omega \rightarrow 0$ as $q \rightarrow 0$
 - QNMs: $\omega(0) \neq 0$
- We have studied far-from-equilibrium dynamics of QNMs.

Discussion

- For small perturbations:

QNMs relax *linearly* and *independently*, with $t_i^{\text{relax}} \sim 1/\text{Im } \omega_i$.

Discussion

- For small perturbations:
QNMs relax *linearly* and *independently*, with $t_i^{\text{relax}} \sim 1/\text{Im } \omega_i$.
- Extend to not-so-small perturbations by adding interactions.

Discussion

- For small perturbations:
QNMs relax *linearly* and *independently*, with $t_i^{\text{relax}} \sim 1/\text{Im } \omega_i$.
- Extend to not-so-small perturbations by adding interactions.
- Expected to break down for $\Delta\mathcal{P}/\mathcal{E} \sim 1$ but it does not.

Discussion

- For small perturbations:
QNMs relax *linearly* and *independently*, with $t_i^{\text{relax}} \sim 1/\text{Im } \omega_i$.
- Extend to not-so-small perturbations by adding interactions.
- Expected to break down for $\Delta\mathcal{P}/\mathcal{E} \sim 1$ but it does not.
- Relaxation still characterized by few frequencies.

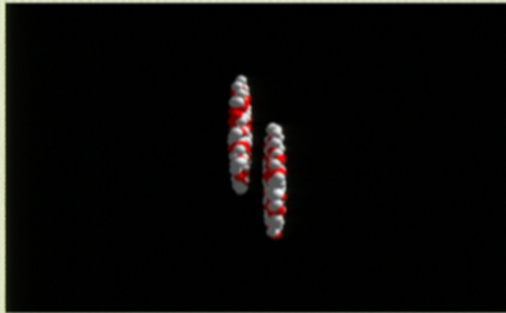
Discussion

- For small perturbations:
QNMs relax *linearly* and *independently*, with $t_i^{\text{relax}} \sim 1/\text{Im } \omega_i$.
- Extend to not-so-small perturbations by adding interactions.
- Expected to break down for $\Delta\mathcal{P}/\mathcal{E} \sim 1$ but it does not.
- Relaxation still characterized by few frequencies.
- Linear approx. valid for stress tensor 1-point function;
other observables probably not well captured.

Discussion

- For small perturbations:
QNMs relax *linearly* and *independently*, with $t_i^{\text{relax}} \sim 1/\text{Im } \omega_i$.
- Extend to not-so-small perturbations by adding interactions.
- Expected to break down for $\Delta\mathcal{P}/\mathcal{E} \sim 1$ but it does not.
- Relaxation still characterized by few frequencies.
- Linear approx. valid for stress tensor 1-point function;
other observables probably not well captured.
- Next step: Include hydrodynamics (boost-invariant case).
- Preliminary results indicate it works.

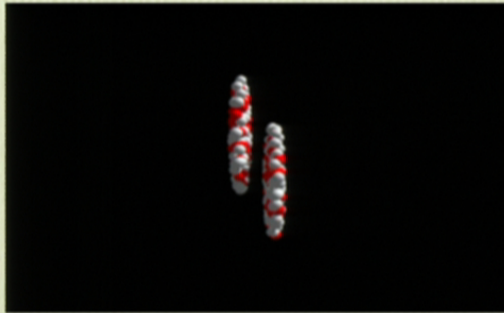
Potential implication



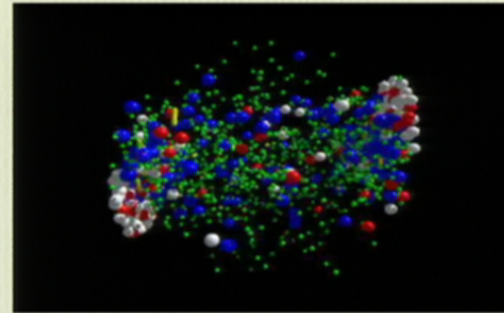
Initial state



Potential implication



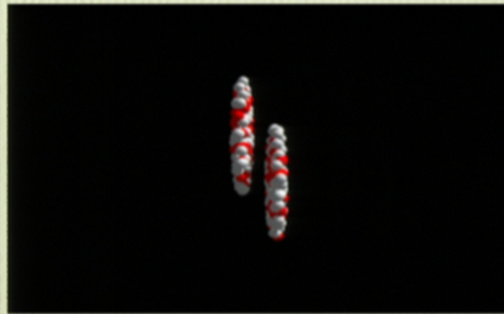
Initial state



Hydro becomes applicable

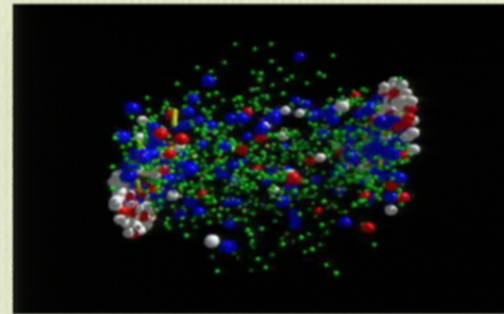


Potential implication



Initial state

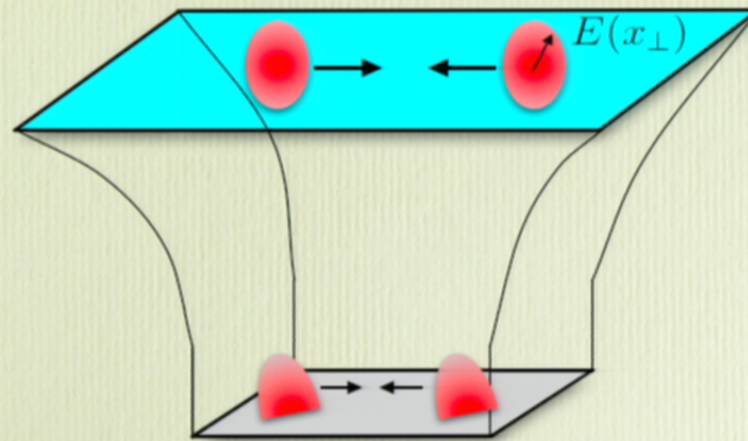
Map?
→



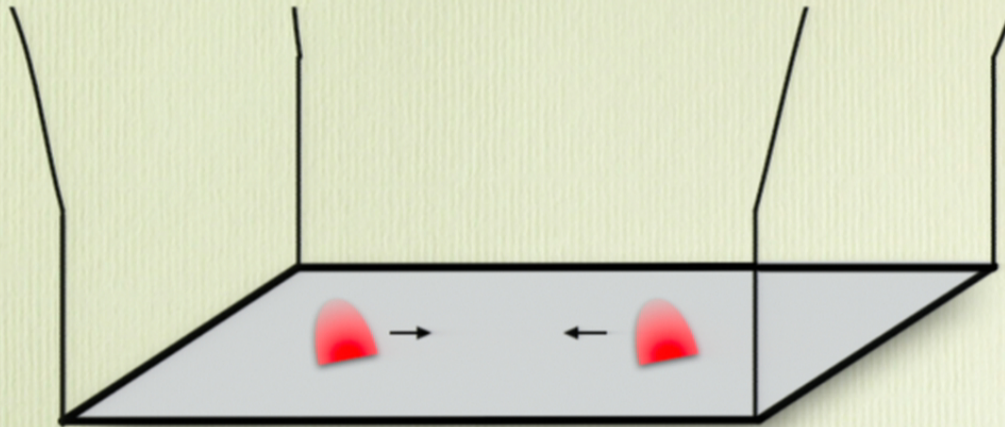
Hydro becomes applicable



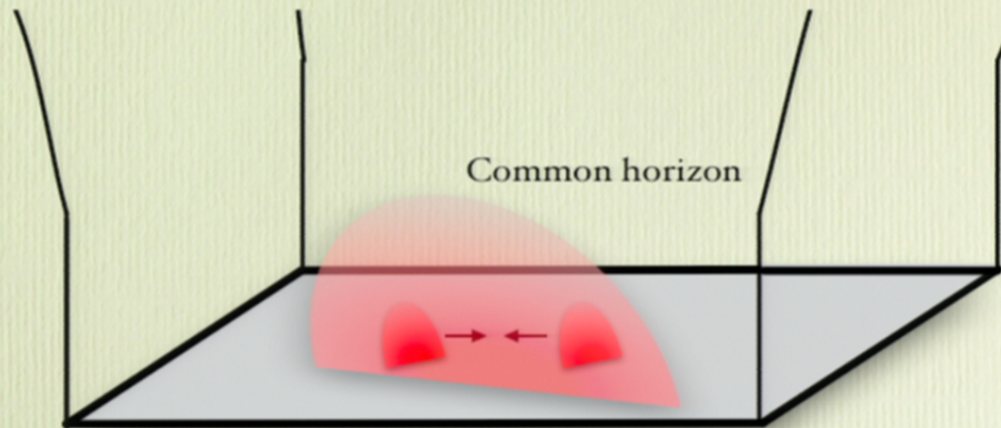
Potential implication



Potential implication



Potential implication



Full non-linearity of gravity encoded in the initial horizon.

Thank you.