Title: Doubly General Relativity

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Abstract:

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Doubly General Relativity Conformal Nature of the Universe

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Machian Picture

Two Machian Principles

- relativity of clocks: local time reparametrization invariance
 - ⇒ local Hamilton constraints
- relativity of rods: local spatial conformal invariance
 - ⇒ local spatial conformal constraints

Generically, second class constraint system

⇒ can not simultaneously realized as phase space symmetries.

On BRST-extended phase space:

Can simultaneously be realized as nilpotent transformations if on-shell Hamiltonian is doubly invariant.

- ⇒ hidden BRST-symmetry in gravity
- ⇒ Doubly General Relativity due to Shape Dynamics

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Heuristic

Constrained Hamiltonian path integral

```
Z = \int Dq Dp \delta[\chi] \delta[\sigma] |\{\chi, \sigma\}| \exp(i \int dt p.\dot{q})
= \int Dq Dp D\eta DP \exp(i \int dt(p.\dot{q} + P.\dot{\eta} - \{\Omega, \Psi\}))
with BRST-generator \Omega, gauge-fixing \Psi.
```

Wish list

- **1** both Machian invariances implemented as invariances under $s_1 = \{\Omega, .\}$ and $s_2 = \{\Psi, .\}$
- ② locality of $L = p.\dot{q} + P.\dot{\eta} \{\Omega, \Psi\}$
- linearly realized symmetries (not yet)
- intuitive interpretation



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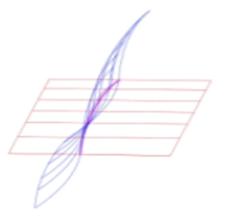
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Duality between ADM and Shape Dynamics



ADM Gravity

$$\begin{split} S(N) &= \int N \left(\frac{G(\pi, \pi)}{\sqrt{|g|}} - (R - 2\Lambda) \sqrt{|g|} \right) \\ &\quad H(v) = \int \pi^{ab} \mathcal{L}_v g_{ab} \end{split}$$



Shape Dynamics

$$\begin{array}{c} H_{SD} = V - V_o \\ Q(\rho) = \int (\pi - \langle \pi \rangle \sqrt{|g|}) \\ H(v) = \int \pi^{ab} \mathcal{L}_v g_{ab} \end{array}$$



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BRST-Formalism (quick and dirty)

Abelian constraints χ_{α}

- BRST-generator $\Omega = \eta^{\alpha} \chi_{\alpha}$ satisfies $\{\Omega, \Omega\} = 0$ (nontrivial: $gh(\Omega) = 1$)
- \Rightarrow defines graded differential $s: f \to \{\Omega, f\}$, i.e. $s^2 = 0$
- Observables as cohomology of s at gh(.) = 0:
 - gauge invariance: $\{\Omega, f(p,q)\} = 0 \Rightarrow f$ (strong observable)
 - equivalence: $\tilde{f} = f + \{\Omega, \Psi\} = f + \sigma^{\alpha} \chi_{\alpha} + \mathcal{O}(\eta)$ (weak observable) for gauge fixing $\Psi = \sigma^{\alpha} P_{\alpha} + \mathcal{O}(\eta)$ with $gh(\Psi) = -1$
- always strong equations on extended phase space
- gauge fixed Hamiltonian $H_{BRS} = H_o + \{\Omega, \Psi\}$ when $\{H_o, \Omega\} = 0$.

Nonabelian constraints $\tilde{\chi}_{\alpha} = M_{\alpha}^{\beta} \chi_{\beta}$

apply canonical transform $\exp(\{\eta_{\alpha}c_{\beta}^{\alpha}P^{\beta},.\})$ to Abelian case $\tilde{\Omega} = \eta^{\alpha}M_{\alpha}^{\beta}\chi_{\beta} + \mathcal{O}(\eta^{2})$ defines \tilde{s} , cohomology same as of s at gh(.) = 0.

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From Symmetry Trading to Symmetry Doubling

Symmetry Trading requires

two first class surfaces (original and equivalent gauge symmetry) that gauge fix one another

BRST-gauge-fixing

- Ω is nilpotent because orig. system is first class
- Ψ can be chosen nilpotent because equiv. system is first class
- if H_o (on shell) Poisson commutes with Ω and Ψ then gauge fixed

$$H_{BRS} = H_o + \{\Omega, \Psi\}$$

is annihilated by both s_{Ω} and s_{Ψ}

Symmetry Doubling:

Canonical action $S = \int dt (p_i \dot{q}^i + P_\alpha \dot{\eta}^\alpha - H_{BRS})$ is invariant under two BRST-transformations (up to a boundary term).

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Different from Anti-BRST

Abelian constraints: $\chi_{\alpha} \approx 0$

- double constraints $\chi_{\alpha}^{1} = \chi_{\alpha}^{2} = \chi_{\alpha}$
- reducibility condition $\chi_{\alpha}^{1} \chi_{\alpha}^{2} = 0$
- \Rightarrow BRST operator $\Omega = \eta_1^{\alpha} \chi_{\alpha}^1 + \eta_2^{\alpha} \chi_{\alpha}^2 + \lambda^{\alpha} (P_{\alpha}^1 P_{\alpha}^2)$

bi-degree expansion yields two commuting generators

$$\Omega = \eta_1^{\alpha} \chi_{\alpha}^1 + \lambda^{\alpha} P_{\alpha}^2$$
 and $\bar{\Omega} = \eta_2^{\alpha} \chi_{\alpha}^2 - \lambda^{\alpha} P_{\alpha}^1$.

Superalgebras

Anti-BRST:
$$\{\Omega,\Omega\}=0=\{\bar{\Omega},\bar{\Omega}\},\{\Omega,H\}=0=\{\bar{\Omega},H\}$$

$$\{\Omega,\bar{\Omega}\}=0.$$

Symmetry Doubling:
$$\{\Omega, \Omega\} = 0 = \{\Psi, \Psi\}, \{\Omega, H\} = 0 = \{\Psi, H\}$$

$$\{\Omega, \Psi\} = H$$



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Construction of Doubly General Relativity (I)

Extending Shape Dynamics

- fixed CMC condition $Q(x) = \pi(x) + \lambda \sqrt{|g|}$
- conformal spatial harmonic gauge $F^k(x) = (g^{ab}\delta^k_c + \frac{1}{3}g^{ak}\delta^b_c)e^c_{\alpha}(\nabla_a \hat{\nabla}_a)e^{\alpha}_b$
- First class system: {Q(x), Q(y)} = 0 = {Fⁱ(x), F^j(y)}
 as well as {Q(x), Fⁱ(y)} = Fⁱ(y)δ(x, y)

Interpretation as "local conformal system"

Q generates spatial dilatations and Poisson brackets resemble C(3) at each point

Gauge fixing ADM

- gauge fixing operator is elliptic and invertible in a region R
- out side R: meager set with finite dimensional kernel

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Construction of Doubly General Relativity (II)

BRST-charges

$$\Omega_{ADM} = \int d^3x \left(\eta S + \eta^a g_{ac} \pi^{cd}_{;d} + \eta^b \eta^a_{,b} P_a + \frac{1}{2} \eta^a \eta_{,a} P + \eta \eta_{,c} P_b g^{bc} \right)$$
 $\Omega_{ESD} = \int d^3x \left(P \frac{\pi}{\sqrt{g}} + P_a F^a + \frac{1}{2} \frac{P}{\sqrt{g}} P_a \eta^a \right)$

Gauge-fixed gravity action

 $S_{gf} = \int dt (\dot{z}_A P(z_A) - \{\Omega_{ADM}, \Omega_{ESD}\})$ is invariant under usual ADM-BRST transformations **and**

a hidden BRST-invariance of S_{gf} under

$$s_{ESD}g_{ab} = \frac{P}{\sqrt{g}}g_{ab}$$
 $s_{ESD}\pi^{ab} = "long"$
 $s_{ESD}\eta = -\frac{1}{\sqrt{g}}(\pi + \frac{1}{2}P_c\eta^c)$ $s_{ESD}P = 0$
 $s_{ESD}\eta^a = -F^a + \frac{P}{2\sqrt{g}}\eta^a$ $s_{ESD}P_a = \frac{P}{2\sqrt{g}}P_a$

due to extended Shape Dynamics.

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Construction of Doubly General Relativity (III)

Interpretation

The Hamiltonian of Doubly General Relativity is

$$H_{DGR} = S(\frac{\pi}{\sqrt{|g|}} + \lambda) + H(F^{a}) + \mathcal{O}(\eta)$$

The ghost-free part is neither the "frozen Hamiltonain" H=0 nor the CMC-Hamiltonian $H=S(N_{CMC}[g,\pi])$, but a generator of dynamics for $\lambda+\frac{\pi}{\sqrt{|g|}}>0$.



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Consequences 1: Classical Theory and Observations

Refined definition of a gravity theory (eff. field th. reasoning):

Gravity = local action for g_{ab} , π^{ab} , η , P, η^{a} , P_{a} at gh. number 0 that is invariant under ADM- and ESD- BRST-transformations s_{ADM} , s_{ESD} also: dimensional analysis in IR

construction ppl. for classical Doubly General Relativity

Possible Observable Consequences

- Effective field theory for GR: all higher derivative curvature invariants are allowed (just suppressed at low energies)
- these are generally not compatible with Extended Shape Dynamics
 DGR can be experimentally distinguished from usual GR (but only beyond Einstein-Hilbert)

This theory space has **not** been explored!

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What about Quantum Gravity?

Slavnov-Taylor Identities (assuming no annom.):

assuming an invariant path integral measure

BRST-variations yield:

$$\langle s_{ADM}\phi_A\rangle\frac{\delta_L\Gamma}{\delta\phi_A}=0 \text{ and } \langle s_{ESD}\phi_A\rangle\frac{\delta_L\Gamma}{\delta\phi_A}=0$$

BRST-variations are nonlinear -> difficult Legendre transform

Nonlinearity obstructs use of two Zinn-Justin equations

$$(\Gamma, \Gamma)_1|_{\hat{\phi}_2=0} = 0 = (\Gamma, \Gamma)_2|_{\hat{\phi}_1=0}$$

seems unfeasible in metric formulation.

Current Directions:

 Find a formulation of DGR where enough transformations are linearly realized:

This makes prediction about counter terms very feasible

Find a gauge fixing with improved power counting.
 (Problem: Observables need Dictionary)

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Conclusions

- Two Machian principles as foundations of gravity theory
- Symmetry trading is generic and gives equivalent gauge theories
- Symmetry trading implies symmetry doubling in BRST formalism
- Equivalence of Shape Dynamics and GR ⇒ Doubly General Relativity
- OGR implies a new theory space for gravity. To explore:
 - are there semiclassical predictions (beyond E-H-action)?
 - universality classes on this revised theory space (FRGE methods)?
 - new view on dualities?

"Doubly General Relativity" in one line:

There is a hidden BRST-invariance in gravity due to Shape Dynamics.

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