

Title: Explorations in Numerical Relativity - Lecture 6

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URL: <http://pirsa.org/12040046>

Abstract:

$$\frac{1}{2} \left(f(\tilde{P}_{i+1/2}^R) + f(\tilde{P}_{i+1/2}^L) - \sum_{\alpha} |\lambda_{\alpha}| w_{\alpha} \Gamma_{\alpha} \right)$$



$$\bar{P}_i + \sigma_i (X_{i+1/2} - X_i)$$

$$= \bar{P}_{i+1} + \sigma_{i+1} (X_{i+1/2} - X_{i+1})$$

$$\left. \begin{aligned} \sigma_i &= \min \text{mod} (S_{i-1/2}, S_{i+1/2}) \\ \min \text{mod} (a, b) &\begin{cases} 0 & \text{if } ab < 0 \\ a & \text{if } |a| < |b| \text{ \& } a, b > 0 \\ b & \text{if } |b| < |a| \text{ \& } a, b > 0 \end{cases} \end{aligned} \right\} S_{i+1/2} = \frac{\bar{P}_{i+1} - \bar{P}_i}{X_{i+1} - X_i}$$

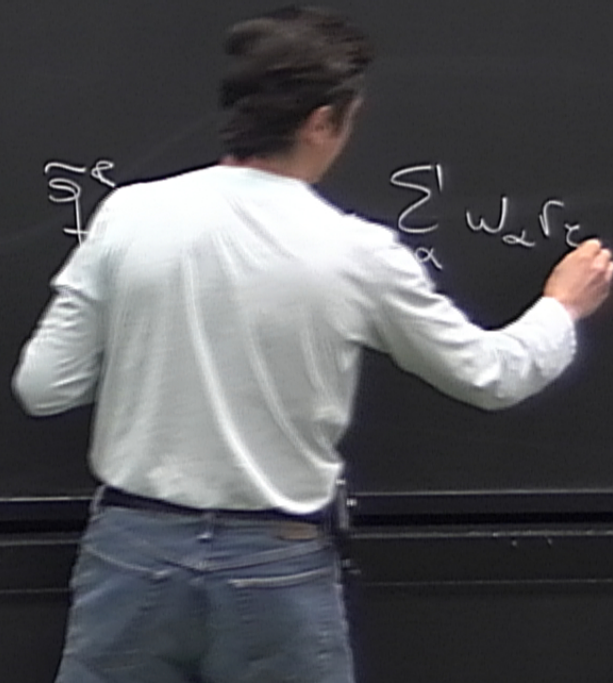
$(+1 \quad +1 \quad +1 \quad +1)$ $(-1 \quad -1 \quad -1 \quad -1)$ $a = \frac{1}{\sqrt{2}}(|a\rangle + |b\rangle) \in \mathbb{C}^2$
 $b = \frac{1}{\sqrt{2}}(|a\rangle - |b\rangle) \in \mathbb{C}^2$

$$\frac{1}{\sqrt{2}} \begin{pmatrix} \tilde{\sigma}_x^L + \tilde{\sigma}_x^R \\ \tilde{\sigma}_y^L + \tilde{\sigma}_y^R \end{pmatrix}$$

λ_α eigenvectors r_α ; jumps w_α defined as.

$$A r_\alpha = \lambda_\alpha r_\alpha$$

$$\tilde{\sigma}_y^R = \sum_\alpha w_\alpha r_\alpha$$



(t_1, t_2, t_3, \dots)

$a = \frac{1}{2}(|a\rangle\langle a| + |a\rangle\langle b| + |b\rangle\langle a| + |b\rangle\langle b|)$
 $b = \frac{1}{2}(|b\rangle\langle a| + |a\rangle\langle b| + |a\rangle\langle a| + |b\rangle\langle b|)$

2-Step method

$$\bar{q}^{n+1/2} = \bar{q}^n + \frac{\Delta t}{2} L(\bar{q}^n)$$

$$\bar{q}^{n+1} = \bar{q}^n + \Delta t L(\bar{q}^{n+1/2})$$

$$\bar{q}^{n+1/2}$$

eigenvectors r_α ; jumps w_α defined as

$$\bar{q}^{n+1} - \bar{q}^{n+1/2} = \sum_{\alpha} w_{\alpha} r_{\alpha}$$

$$A r_{\alpha} = \lambda_{\alpha} r_{\alpha}$$

(t_1, t_2, \dots)

$a = |a\rangle\langle a| + \dots$
 $b = |b\rangle\langle b| + \dots$

2-Step method

$$\bar{q}^{n+1/2} = \bar{q}^n + \frac{\Delta t}{2} L(\bar{q}^n)$$

$$\bar{q}^{n+1} = \bar{q}^n + \Delta t L(\bar{q}^{n+1/2})$$

\bar{q}^k
 $\bar{q}^{i+1/2}$

eigenvectors r_α ; jumps w_α

$$L(\bar{q}^n) = \frac{F_{i+1/2}^{Roe}(\bar{q}^n) - F_{i-1/2}^{Roe}(\bar{q}^n)}{\Delta x} + \hat{\Psi}_i(\bar{q}^n)$$

$$A r_\alpha = \lambda_\alpha r_\alpha$$

$$\bar{q}_{i+1/2}^{PR} - \bar{q}_{i+1/2}^{LC} = \sum_{\alpha} w_\alpha r_\alpha$$

$$\rightarrow \frac{\partial p}{\partial t} + \frac{\partial f}{\partial x} = 0$$

$$p_{i+1/2} = p_i + v_i (x_{i+1/2} - x_i)$$

$$\tilde{p}_{i+1/2}^R = \bar{p}_{i+1/2} + \sigma_{i+1/2} (x_{i+1/2} - x_{i+1})$$

$$v_i = \text{minmod} \left((s_{i-1/2})^{-i+1/2}, \dots, \frac{x_{i+1} - x_i}{\dots} \right)$$

$$\text{minmod}(a, b) = \begin{cases} 0 & \text{if } a < 0 \\ a & \text{if } |a| < |b| \text{ \& } a, b > 0 \\ b & \text{if } |b| < |a| \text{ \& } a, b > 0 \end{cases}$$

$$A|_{i+1/2} \equiv \frac{\partial f}{\partial q} \Big|_{q = \frac{1}{2}(\tilde{q}_{i+1/2}^L + \tilde{q}_{i+1/2}^R)}$$

→ Compute eigenvalues

λ_α , eigenvectors r_α ; jumps w_α

$$A r_\alpha = \lambda_\alpha r_\alpha$$

$$\tilde{p}_{i+1/2}^R - \tilde{p}_{i+1/2}^L = \sum_{\alpha} w_\alpha r_\alpha$$

2-step method

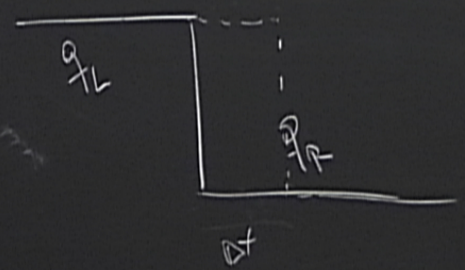
$$\bar{q}^{n+1/2} = \bar{q}^n + \frac{\Delta t}{2} L(\bar{q}^n)$$

$$\bar{q}^{n+1} = \bar{q}^n + \Delta t L(\bar{q}^{n+1/2})$$

$$L(\bar{q}^n) = \frac{F_{i+1/2}^{Roe}(\bar{q}^n) - F_{i-1/2}^{Roe}(\bar{q}^n)}{\Delta x} + \hat{\Psi}_i(\bar{q}^n)$$

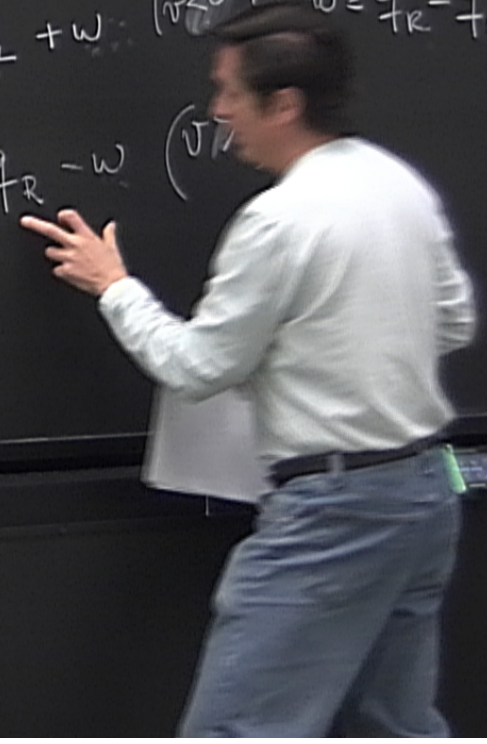
$$q_{T_L} + v q_{T_R} = 0$$

Riemann problem



At later time

$$q(t) = \begin{cases} q_{T_L} + w & (v < 0) \\ q_{T_R} - w & (v > 0) \end{cases} \quad w = q_{T_R} - q_{T_L}$$



$$\frac{q^R}{\alpha}$$

$$q^R - (q^R - q^L) = q^L$$

Burger's eqn

$$q_{,t} + q q_{,x} = 0 \rightarrow q_{,t} + \left(\frac{q^2}{2}\right)_{,x} = 0$$

$$F_{i+1/2}^{n+1} =$$

$$q_{,t} + 2x q = 0$$

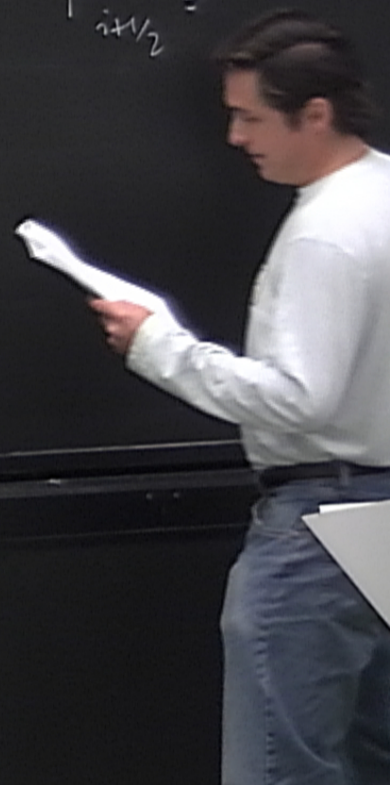
$$q_{,x} = \frac{1}{2} q^2, \quad \psi = 0$$

$$\lambda = \frac{\partial q}{\partial q} = q$$

$$r_\alpha = 1$$

$$\lambda = q$$

$$w = q^R - q^L$$



$$q_R - (q_R - q_L) = q_L$$

$$f_{,t} + f f_{,x} = 0 \rightarrow f_{,t} + \left(\frac{q^2}{2}\right)_{,x} = 0$$

$$f = \frac{1}{2} q^2, \quad \psi = 0$$

$$r_2 = 1, \quad \lambda = f, \quad w = q^R - q^L$$

$$F_{i+1/2}^{j+1/2} = \frac{1}{2} [f(q^L) + f(q^R) - |x| r w]_{i+1/2}$$

$$= \frac{1}{2} \left[q^L + q^R - |q| \left(\frac{q^R - q^L}{2} \right) \right]$$

↓
 $\frac{1}{2}(q^R + q^L)$

GR + Hydro.

1) $G_{ab} = 8\pi T_{ab} \rightarrow \{g_{ab}$

2) $\nabla_a T^{ab} = 0 \rightarrow$ perfect

fluid

(3)

(no visc
or
heat exchange)

$$T_{ab} = (\rho + P) u_a u_b + P g_{ab}$$

→ ξ_{gas}

perfect fluid visc stress exchange

$$T_{ab} = (\rho + P) u_a u_b + P g_{ab}$$

$$\rho = \rho_0 (1 + \epsilon)$$

- ρ_0 (1)
- P (1)
- u^a (4)
- ϵ (1) specific internal energy

Constraint $u^a u_a = -1$

Eqn of state $P = P(\rho, \epsilon)$

GR + Hydro.

1) $G_{ab} = 8\pi T_{ab} \rightarrow \{g_{ab}\}$

2) $\nabla_a T^{ab} = 0 \rightarrow$ perfect fluid

(3) $\nabla_a (p u^a) = 0$
(no visc
or
heat exchange)

$$T_{ab} = (\rho + p) u_a u_b + p g_{ab}$$

GR + Hydro.

1) $G_{ab} = 8\pi T_{ab} \rightarrow \{g_{ab}\}$

2) $\nabla_a T^{ab} = 0$ perfect fluid

(3) $\nabla_a (\rho u^a) = 0$

$T_{ab} = (\rho + P) u_a u_b + P g_{ab}$

p. part of $G_{ab} \rightarrow g_{ab}$

p. part of $\nabla T = 0$

$\nabla(\rho u) = 0$

$$\xi \quad g_{ab} = \eta_{ab}$$

$$\partial_a (\xi_a u^a) = 0 \rightarrow \partial_t (\xi_a u^0) + \partial_i (\xi u^i) = 0$$

$$\partial_b (T^{ab}) = 0$$

New variables

$$\xi \quad g_{ab} = \eta_{ab}$$

to solving

Nas variables - Combination

$$\partial_a (\xi_a u^a) = 0 \rightarrow \partial_t (\xi_a u^0) + \partial_i (\xi u^i) = 0$$

$$\partial_b (T^{ab}) = 0$$

to solve ρ

Nas variables - combinations of "primitive variables" (ρ, ε, u^i)
↳ "conserved / conservation variables"

$$\rightarrow \partial_t (\rho_0 u^0) + \partial_i (\rho u^i) = 0 \quad \rho_0 = 1$$

* Simplify → ultra-relativistic limit: $\rho_0 \varepsilon \gg \rho_0$

GR + Hydro.

1) $G_{ab} = 8\pi T_{ab} \rightarrow \{g_{ab}\}$

2) $\nabla_a T^{ab} = 0 \rightarrow$ perfect fluid

(3) $\nabla_a (\beta u^a) = 0$
(no visc or heat exchange)

$$T_{ab} = (\rho + P) u_a u_b + P g_{ab}$$

p. part of $\underline{G_{ab}} \rightarrow \partial_a g_{ab} \rightarrow (\Gamma)$

p. part of $\nabla T = 0$
 $\nabla(\rho u) = 0$

$$\partial(\text{fluid}) = \dots$$

$$\partial_a (P_0 u^a) = 0 \rightarrow \partial_t (P_0 u^0) + \partial_i (P_0 u^i) = 0 \rightarrow \text{"conserved / conservation variables"}$$

$$\partial_b (T^{ab}) = 0$$

* Simplify \rightarrow ultra-relativistic limit: $P_0 \epsilon \gg P_0$

* Symmetry: "slab" $P = P_0 (1 + \epsilon) = P_0 + P_0 \epsilon \Rightarrow P_0 \epsilon \approx P$

Symmetry $f \neq f(y, z)$

Notice = $u^a = (\gamma t, \gamma v_i)$ $\gamma = \frac{1}{\sqrt{1 - v^2}}$ $v = \frac{u^x}{u^t}$

x Symmetry: "sleb"
 Symmetry, $f \neq f(y)$

Notice - $u^a = (\gamma^1, \gamma^1 v_i)$ $\gamma^1 = \gamma = \frac{1}{\sqrt{1-v^2}}$ $v = \frac{u^x}{u^t}$

New vars · $\tau = (\gamma + p) \gamma^2 - p$
 $S = v \gamma^2 (\tau + p)$

$T^{tt} = \tau$
 $T^{tx} = T^{xt} = S$
 $T^{xx} = S v + P =$

* Symmetry: "slab" $J = J_0(t, z) = J_0(t) \delta(z)$
 Symmetry $f \neq f(y, z)$

$(\gamma^i, \gamma^i v_i)$

$$\gamma^i = \gamma^i v_i = \frac{1}{\sqrt{1-v^2}} \quad v = \frac{u^x}{u^t}$$

$$\begin{aligned} & (\tau + p) \gamma^2 - p \\ & v \gamma^2 (\tau + p) \end{aligned}$$

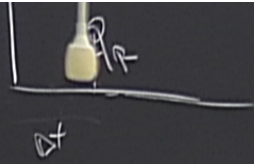
$$T^{tt} = \tau$$

$$T^{tx} = T^{xt} = \tau v$$

$$T^{xx} = \tau v^2 + p = \tau v^2 + p$$

$$1) \tau_{,t} + \partial_x S = 0$$

$$2) S_{,t} + \partial_x (S v + P) = 0$$



$$q_R - (q_R - q_L) = q_L$$

$$q = \begin{bmatrix} \tau \\ s \end{bmatrix}; f(q) = \begin{bmatrix} S \\ 5\sigma + P \end{bmatrix}; \psi = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Assume $P = (R-1)S$

$$A = \begin{bmatrix} 0 & 1 \\ -\sigma^2 + (1-\sigma^2) \frac{\partial P}{\partial \tau} & 2\sigma + (1-\sigma^2) \frac{\partial P}{\partial S} \end{bmatrix}$$

