

Title: Explorations in Particle Theory - Lecture 6

Date: Apr 10, 2012 09:00 AM

URL: <http://pirsa.org/12040008>

Abstract:

Non-Thermal DM

SUSY \rightarrow extension of Poincaré

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UGRA \rightarrow extension of GR.

$g = \text{graviton} \leftrightarrow \Psi_{\mu} = \text{gravitino}$

$$S=2$$

$$S=\frac{3}{2}$$

Non-Thermal DM

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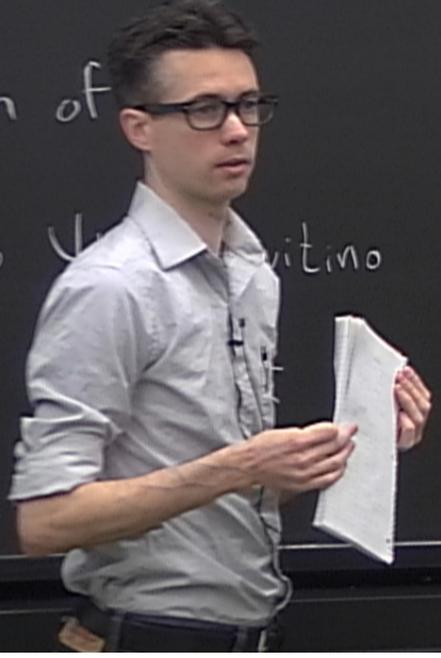
SUSY : $[F]_M = 2$

$$m_{3/2} = \frac{F}{M_{\text{Pl}}}$$

of Poincaré
on of
→ ψ (gritino)

SUSY : $[F]_M = 2$ ($F=0 \Rightarrow$ SUSY is unbroken)

$$m_{3/2} = \frac{F}{M_{Pl}} \approx 2.4 \times 10^{18} \text{ GeV}$$



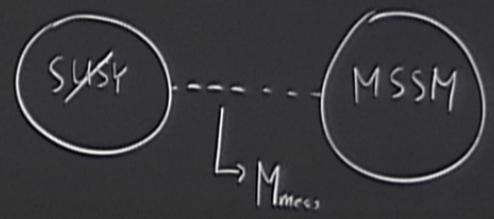
of Poincaré
on of GR.

→ Ψ_{μ} = gravitino
 $S = \frac{3}{2}$

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$$m_{\text{soft}} = \frac{F}{M_{\text{mess}}} = \text{MSSM soft term.}$$



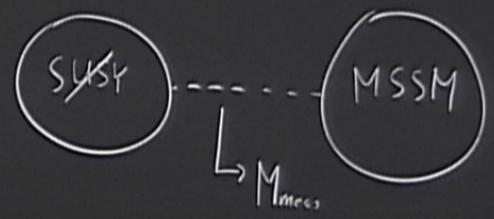
of Poincaré
on of G.R.

$\Psi_\mu = \text{gravitino}$
 $S =$

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$M_{\text{mess}} \sim \begin{cases} M_{\text{Pl}} \rightarrow \text{gravity mediated SUSY} \\ M_* \ll M_{\text{Pl}} \rightarrow \text{gauge mediation} \end{cases}$

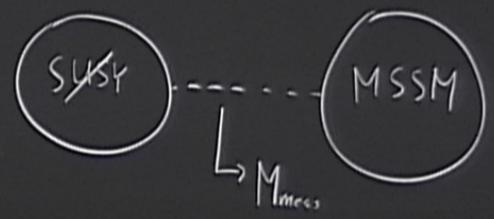
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$M_{\text{mess}} \sim \left\{ \begin{array}{l} M_{\text{Pl}} \rightarrow \text{gravity mediated SUSY} \Rightarrow m_{3/2} \approx m_{\text{soft}} \\ M_* \ll M_{\text{Pl}} \rightarrow \text{gauge mediation} \Rightarrow m_{3/2} \ll m_{\text{soft}} \end{array} \right.$

$g = \text{graviton} \leftrightarrow \Psi_{\frac{3}{2}} = \text{gravitino}$
 $S=2$ $S=\frac{3}{2}$

$M_{\text{mass}} \sim \begin{cases} M_{\text{Pl}} \rightarrow \text{gravity mediated SUSY} \Rightarrow m_{3/2} = M_{\text{soft}} \Rightarrow \text{can} \\ M_{\text{g}} \ll M_{\text{Pl}} \rightarrow \text{gauge mediation} \Rightarrow m_{3/2} \ll M_{\text{soft}} \Rightarrow \text{LSP} \end{cases}$

eg 1 SuperWIMP DM

- MSSM LSP freezes out χ
- $\chi \rightarrow \Psi + (\text{SM})$
- Ψ is now the DM

$$\Omega_{\Psi} h^2 =$$

SP freezes out: χ

$$\Omega_{\Psi_m} h^2 = \left(\frac{m_{3/2}}{m_{LSP}} \right) \underbrace{\Omega_{LSP} h^2}$$

relic density if χ were stable

DM

LSP freezes out: χ

+ (SM)

low the DM

$$\Omega_{\Psi_{\mu}} h^2 = \left(\frac{m_{3/2}}{m_{LSP}} \right) \underbrace{\Omega_{LSP} h^2}$$

relic density if χ were stable

→ decay happens late \Rightarrow can mess up Nucleosynthesis.

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- MSSM LSP
 - Ψ_μ
 - SM
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$$\Omega_\Psi h^2 = \left(\frac{m_{3/2}}{m_{\text{LSP}}} \right) \underbrace{\Omega_{\text{LSP}} h^2}_{\text{relic density if } \chi \text{ were stable}}$$

decay happens late \Rightarrow can mess up Nucleosynthesis

eg 2 Direct Ψ_n DM

$$-\mathcal{L} = \frac{1}{M_{\text{Pl}}} \bar{\Psi}_n \gamma^\mu \gamma^{\nu} (\partial_\nu \tilde{F}^*) f$$

Rate is too slow for Ψ_n to equilibrate unless $T \sim M_{\text{Pl}}$

eg 2 Direct Ψ_n DM

$$-L = \frac{1}{M_{\text{Pl}}} \bar{\Psi}_n \gamma^\mu \gamma^\nu (\partial_\nu \tilde{F}^\mu) f$$

Massive vector $-\mathcal{L}_{\text{int}} + \frac{p_\mu p_\nu}{M_W^2}$
↑
enhancement

Rate is too slow for Ψ_n to equilibrate unless $T \sim M_{\text{Pl}}$.

↳ Ψ_n never equilibrates

Direct Ψ_μ DM

$$\hookrightarrow \frac{1}{M_{\text{Pl}}} \bar{\Psi}_\mu \gamma^\mu \gamma^\nu (\partial_\nu \tilde{f}) f$$

is too slow for Ψ_μ to equilibrate unless $T \sim M_{\text{Pl}}$.

Ψ_μ never equilibrates.

$$\text{distino component} \Rightarrow \left(1 + \frac{m_{\text{soft}}^2}{m_{3/2}^2} \right) \text{factor} \times |M|^2$$

Massive vector $-\mathcal{L}_{\text{uv}} + \frac{p_\mu p_\nu}{M_W^2}$
↑
enhancement

Direct Ψ_μ DM

$$\rho \approx \frac{1}{M_{\text{Pl}}} \bar{\Psi}_\mu \gamma^\mu \gamma^0 (\partial_\nu \tilde{f}) f$$

is too slow for Ψ_μ to equilibrate unless $T \sim M_{\text{Pl}}$.

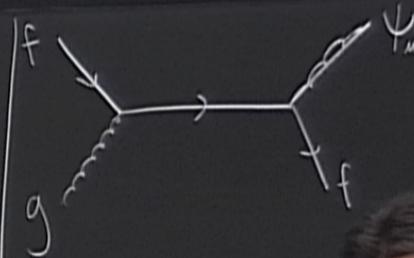
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$$\text{distino component} \Rightarrow \left(1 + \frac{m_{\text{soft}}^2}{m_{3/2}^2} \right) \text{factor} \times |M|^2$$

Massive vector

$$- \mathcal{N}_{\mu\nu} + \frac{p_\mu p_\nu}{M_W^2}$$

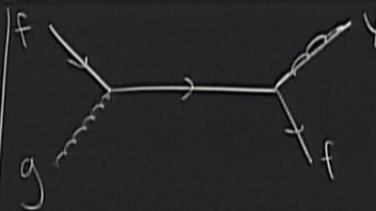
↑ enhancement



Massive vector $-N_{UV} + \frac{p_U p_V}{M_W^2}$
 ↑
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equilibrate unless $T \sim M_{Pl}$

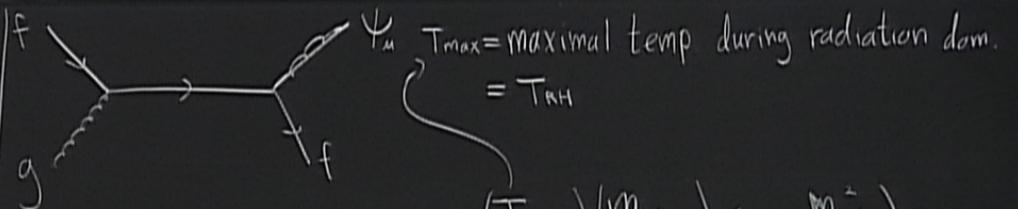
$\left(1 + \frac{m_{soft}^2}{m_{3/2}^2}\right)$ factor $\times |M|^2$



$T_{max} = \text{maximal temp during radiation dom.}$
 $= T_{RH}$

$$\Omega_\psi h^2 \approx (0.5) \alpha_3 \left(\frac{T_{max}}{10^{10} \text{ GeV}}\right) \left(\frac{m_{3/2}}{100 \text{ GeV}}\right) \left(1 + \frac{m_{soft}^2}{m_{3/2}^2}\right)$$

Massive vector $-N_{UV} + \frac{p_U p_V}{M_W^2}$
 ↑ enhancement



equilibrate unless T

$$\Omega_{\psi} h^2 \approx (0.5) \alpha_3 \left(\frac{T_{max}}{10^{16} \text{ GeV}} \right) \left(\frac{m_{3/2}}{100 \text{ GeV}} \right) \left(1 + \frac{m_{soft}^2}{m_{3/2}^2} \right)$$

↳ assumption is that $\psi_k = \text{LSP}$.

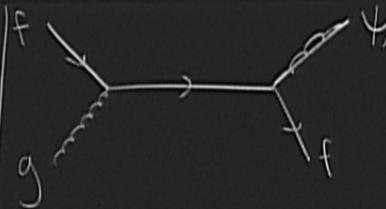
$+ \frac{m_{soft}^2}{m_{3/2}^2}$ factor $\times |M|$

Massive vector $-N_{UV} + \frac{P_U P_V}{M_W^2}$
 ↑ enhancement

equilibrate unless $T \sim M_{Pl}$

$\alpha_3 = \text{QCD coupling}$
 $= g_s^2 / 4\pi$
 ↳ evaluated at T_{max}

$+ \frac{m_{soft}^2}{m_{3/2}^2}$ factor $\times (M_{Pl}^2)$



$T_{max} = \text{maximal temp during radiation dom.}$
 $= T_{RH}$

$$\Delta \Omega_\psi h^2 \approx (0.5) \alpha_3 \left(\frac{T_{max}}{10^{10} \text{ GeV}} \right) \left(\frac{m_{3/2}}{100 \text{ GeV}} \right) \left(1 + \frac{m_{soft}^2}{m_{3/2}^2} \right)$$

↳ assumption is that $\Psi_k = \text{LSP.}$
 ↳ QCD

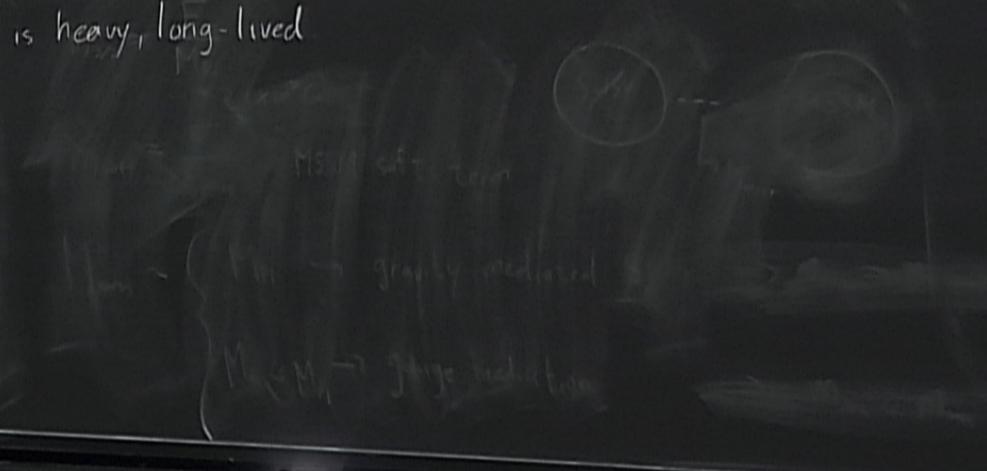
e'g 3 DM from Heavy Particle Decay

Suppose we have a heavy particle P that is heavy, long-lived

$U(1) \rightarrow U(1) \times U(1)$

$g = g_1 + g_2$
 \rightarrow gaugino \rightarrow gravitino

$F=0 \Rightarrow$ SUSY is unbroken



e'g 3 DM from Heavy Particle Decay

Suppose we have a heavy particle P that is heavy, long-lived

P becomes the dominant contribution^{to} ρ .

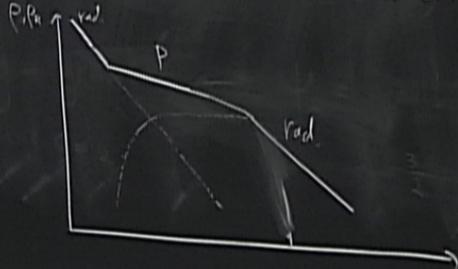
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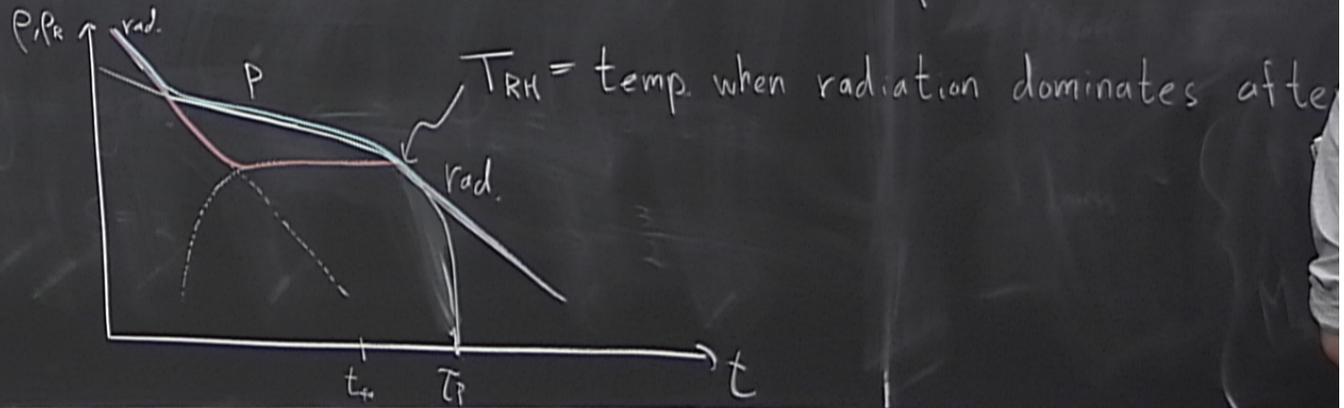
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e'g 3 DM from Heavy Particle Decay

Suppose we have a heavy particle P that is heavy, long-lived

P becomes the dominant contribution ρ .



easy, long-lived.

dominates after P decays

$\epsilon =$ fraction of P decay products that are the DM χ .
 $\ll 1$

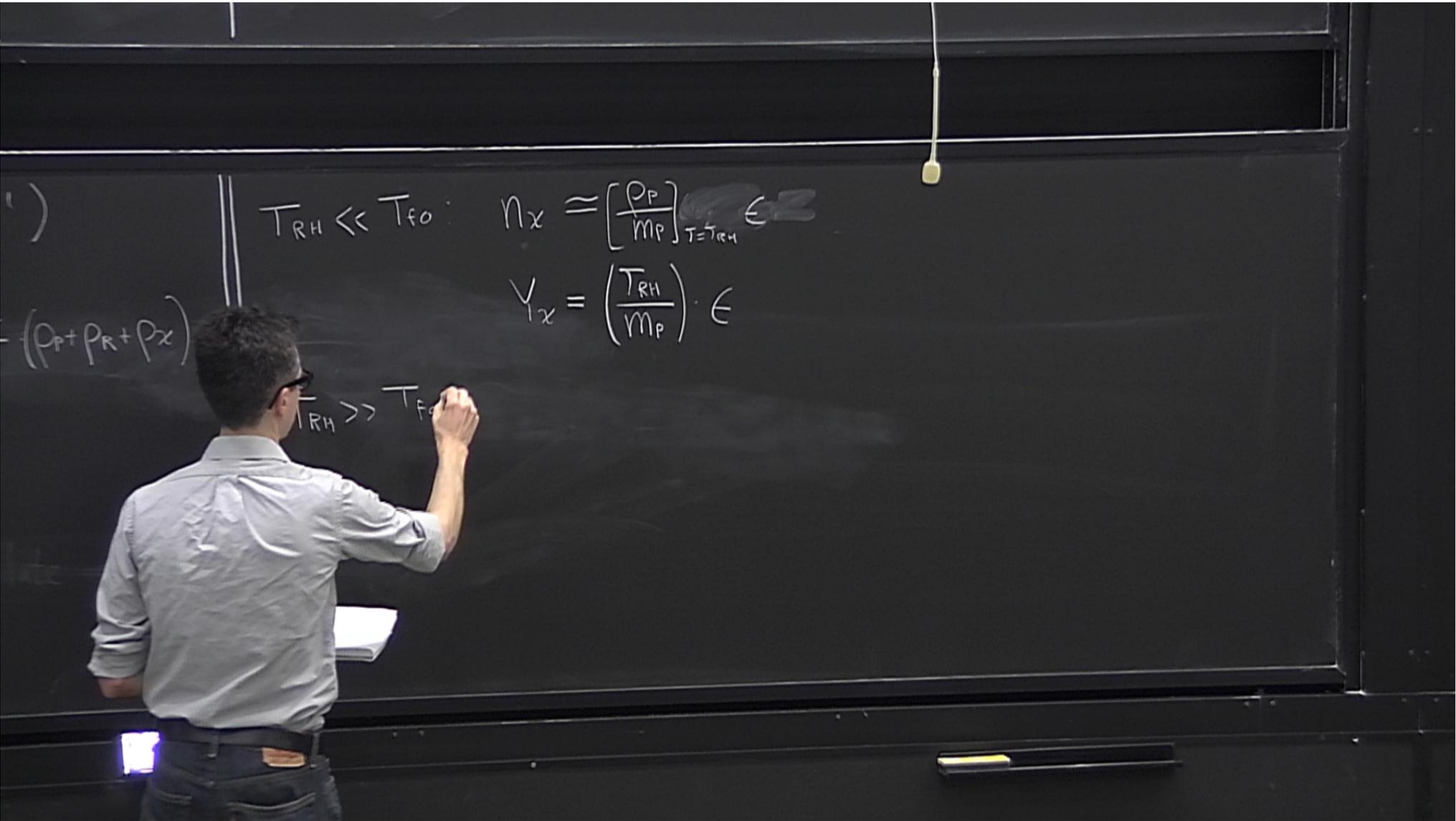
$(1-\epsilon) =$ fraction of decays that end up as radiation.

Non-thermal $m_P \gg m_\chi$

$T_{RH} \ll T_{Fo} \leq m_\chi$

t_* τ_P t

$$\left\{ \begin{aligned} \frac{d\rho_P}{dt} + 3H\rho_P &= -\Gamma_P \rho_P & (\rho_P = m_P n_P, \Gamma_P = \tau_P^{-1}) \\ \frac{d\rho_R}{dt} + 4H\rho_R &= (1-\epsilon)\Gamma_P \rho_P, & H^2 = \frac{8\pi G}{3} (\rho_P + \rho_R + \rho_X) \\ \frac{dn_X}{dt} + 3Hn_X &= \epsilon\Gamma_P \rho_P - \langle \sigma v \rangle (n_X^2 - n_{X_0}^2) \end{aligned} \right.$$



)
($\rho_P + \rho_R + \rho_X$)

$$T_{RH} \ll T_{fo} \quad n_X \approx \left[\frac{\rho_P}{m_P} \right]_{T=T_{RH}} \in \quad \longrightarrow \text{non-thermal}$$

$$Y_X = \left(\frac{T_{RH}}{m_P} \right) \in$$

$$T_{RH} \gg T_{fo} \quad Y_X = \text{usual freeze out value} \longrightarrow \text{thermal}$$

\hookrightarrow DM scattering equilibrates χ density.

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($\rho_P + \rho_R + \rho_X$)

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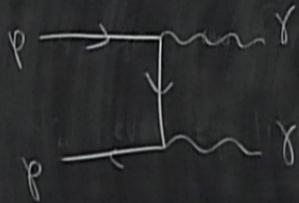
$$Y_X = \left(\frac{T_{RH}}{m_P} \right) \cdot \epsilon$$

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\hookrightarrow DM scattering equilibrates χ density.

eg 4 Asymmetric DM

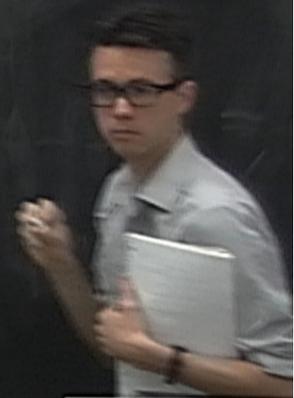
Baryons p, \bar{p}, n, \bar{n}

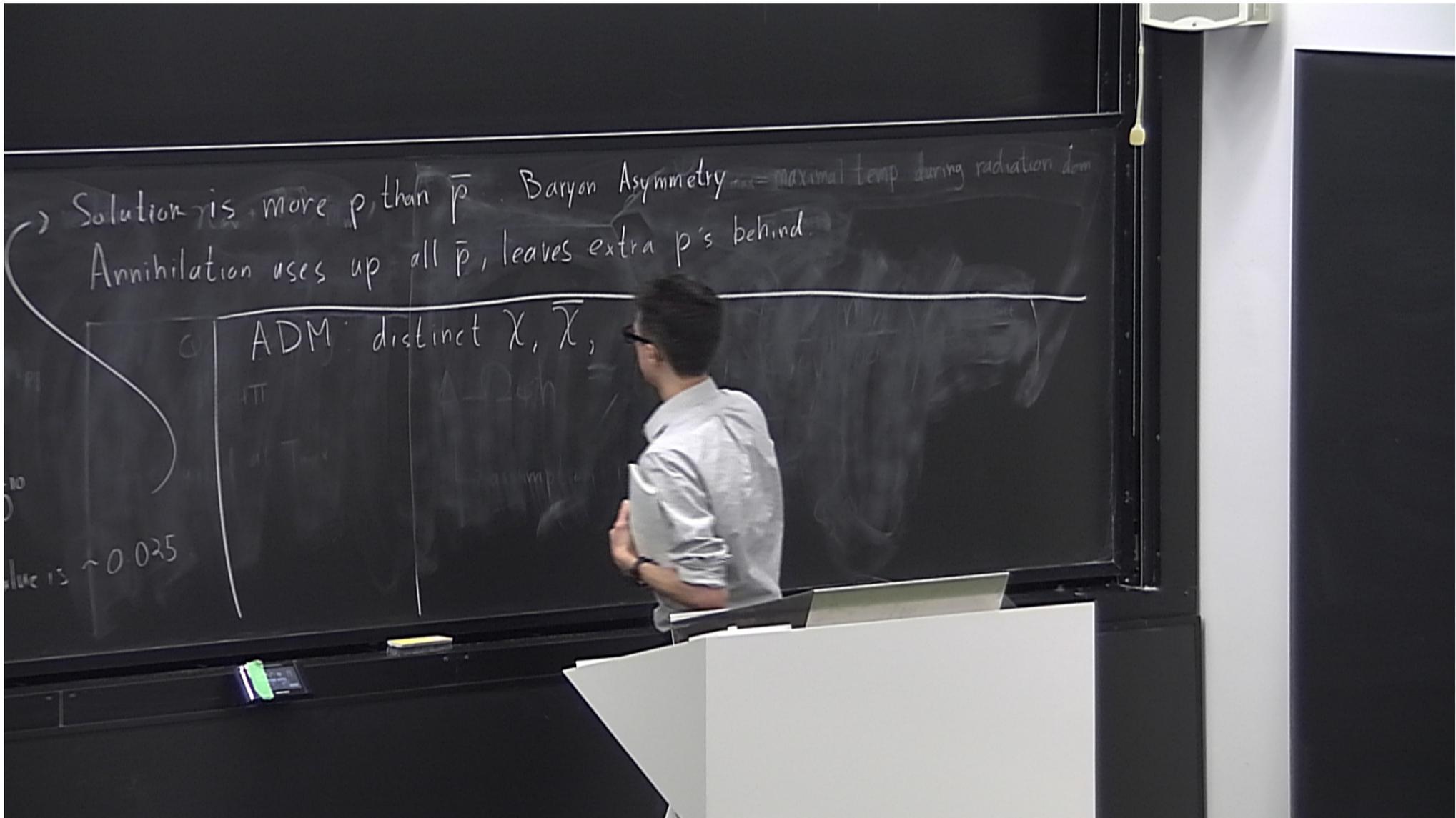


\Rightarrow very efficient
($T \sim 50$ MeV)

Equal # of p, \bar{p} : $\Omega_{bh}^2 = 10^{-10}$
 \hookrightarrow observed value is ~ 0.025

M \rightarrow Solution is more p than \bar{p}





Solution is more p than \bar{p} Baryon Asymmetry - maximal temp during radiation dom
Annihilation uses up all \bar{p} , leaves extra p 's behind.

ADM: distinct χ , $\bar{\chi}$,

value is ~ 0.025

Solution is more p than \bar{p} . Baryon Asymmetry \dots = maximal temp during radiation dom
Annihilation uses up all \bar{p} , leaves extra p 's behind.

ADM: distinct $\chi, \bar{\chi}$; more χ than $\bar{\chi}$

χ charge is related to B charge.

$$n_\chi \approx n_b \Rightarrow \Omega_{DM} h^2 \approx 0.1 \approx \left(\frac{m_\chi}{m_p}\right) \Omega_b h^2 \Rightarrow m_\chi \sim 5 \text{ GeV}$$

value is ~ 0.025