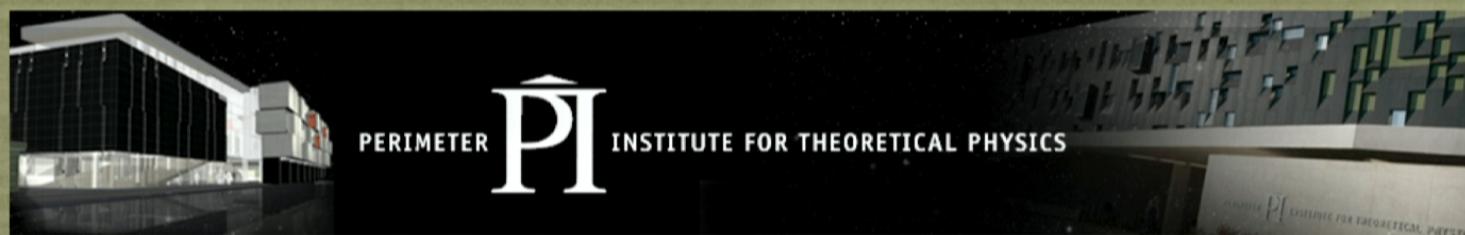


Title: G-Bounce

Date: Mar 19, 2012 02:00 PM

URL: <http://pirsa.org/12030103>

Abstract: I will discuss a wide class of models which realise a bounce in a spatially flat Friedmann universe in standard General Relativity. The key ingredient is a noncanonical, minimally coupled scalar field<br>belonging to the class of theories with Kinetic Gravity Braiding/Galileon-like self-couplings. In these models, the universe smoothly evolves from contraction to expansion, suffering neither from ghosts<br>nor gradient instabilities around the turning point. The end-point of the evolution can be a standard radiation-domination era or an inflationary phase.<br>The talk is based on arXiv:1109.1047v2 [hep-th], JCAP11(2011)021.



Alexander Vikman

(CERN)

**19.03.2012**

# THIS TALK IS BASED ON

- *Imperfect Dark Energy from Kinetic Gravity Braiding*  
**arXiv:1008.0048**, JCAP 1010:026, 2010
- *The Imperfect Fluid behind Kinetic Gravity Braiding*  
**arXiv:1103.5360**, JHEP 1111 (2011) 156
- *G-Bounce*, <sup>3</sup>**arXiv:1109.1047**, JCAP 1111:021, 2011

IN COLLABORATION WITH

Cédric Deffayet,<sup>1</sup> Damien Easson,<sup>3</sup>  
Oriol Pujolàs,<sup>1,2</sup> Ignacy Sawicki<sup>1,2,3</sup>

# PLAN

- Motivation
- General class of models - perfect imperfect fluids
- G - Bounce
- Open problems

# ORIGIN OF THE UNIVERSE?

- Was there a beginning of time i.e. of the quasiclassical universe? Was there a strong quantum gravity époque in our past?
- If there was a beginning, was the universe collapsing or expanding immediately afterwards?
- If the universe experienced an early period of inflation, which *all* observations currently perfectly confirm, what happened *before inflation*? Indeed, *inflation* is not past-complete - Börde, Guth, Vilenkin (2001)

CAN ONE CONSTRUCT  
A **STABLE**  
(WITH REAL SOUND SPEED AND WITHOUT GHOSTS)  
**CLASSICAL MODEL**  
WHERE A  
**SPATIALLY FLAT**  
FRIEDMANN UNIVERSE  
**BOUNCES:**  
GOES FROM COLLAPSE TO EXPANSION ???

**OBSTACLE:**  
TO BOUNCE ONE HAS TO VIOLATE  
NULL ENERGY CONDITION (NEC),  
BUT NORMAL FIELDS AND PERFECT  
FLUIDS CANNOT DO IT!



*go to imperfect fluids and to even less canonical fields!*



**GO G!**

# **G STANDS FOR GALILEON**

*Nicolis, Rattazzi , Trincherini, (2008)*

In current literature: ***Galileons*** = scalar-tensor theories with

***higher derivatives in the action***

***but with equations of motion which are *all* of the second order - NO Ostrogradsky's ghosts***



***the most general theory of this type was derived by Horndeski (1974), and rederived by Deffayet, Gao, Steer, Zahariade (2011)***

*International Journal of Theoretical Physics*, Vol. 10, No. 6 (1974), pp. 363-384

## **Second-Order Scalar-Tensor Field Equations in a Four-Dimensional Space**

GREGORY WALTER HORNDESKI

*Department of Applied Mathematics, University of Waterloo, Waterloo, Ontario,  
Canada*

*Received: 10 July 1973*

## Second-Order Scalar-Tensor Field Equations in a Four-Dimensional Space

GREGORY WALTER HORNDESKI

*Department of Applied Mathematics, University of Waterloo, Waterloo, Ontario,  
Canada*

*Received: 10 July 1973*

$$\begin{aligned}\mathcal{L} = & K + \mathcal{G}\phi_{;\mu}^{;\mu} + \\ & + \mathcal{G}_2 R + \mathcal{G}'_2 \left[ (\phi_{;\mu}^{;\mu})^2 - \phi_{;\mu}^{;\nu} \phi_{;\nu}^{;\mu} \right] + \\ & + \mathcal{G}_3 G_\nu^\mu \phi_{;\mu}^{;\nu} - \frac{1}{6} \mathcal{G}'_3 \left[ (\phi_{;\mu}^{;\mu})^3 - 3\phi_{;\lambda}^{;\lambda} \phi_{;\mu}^{;\nu} \phi_{;\nu}^{;\mu} + 2\phi_{;\alpha}^{;\mu} \phi_{;\beta}^{;\alpha} \phi_{;\mu}^{;\beta} \right]\end{aligned}$$

*Kobayashi, Yamaguchi, Yokoyama (2011)*

where we have 4 free functions!

$$K(\phi, X) \text{ and } \mathcal{G}_i(\phi, X)$$

$$\mathcal{G}'_i = \frac{\partial \mathcal{G}_i}{\partial X}$$

$$X \equiv \frac{1}{2} g^{\mu\nu} \nabla_\mu \phi \nabla_\nu \phi$$

## WHY IS THIS **G** INTERESTING?

One can **stably violate** the most basic of the classical energy conditions-Null Energy Condition (NEC):  $T_{\mu\nu}n^\mu n^\nu \geq 0$  which in cosmology reduces to:  $p + \varepsilon \geq 0$  or  $\dot{H} \leq 0$

*Pandora's  
box*



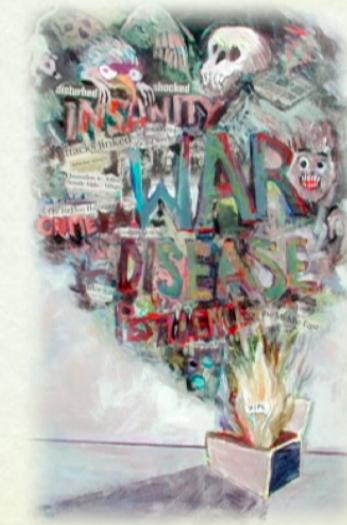
# WHY IS THIS **G** INTERESTING?

One can **stably violate** the most basic of the classical energy conditions-Null Energy Condition (NEC):  $T_{\mu\nu}n^{\mu}n^{\nu} \geq 0$  which in cosmology reduces to:  $p + \varepsilon \geq 0$  or  $\dot{H} \leq 0$



*Pandora's  
box*

- *Healthy and testable*  
Phantom Dark Energy,  
Deffayet, Pujolas, Sawicki, AV, 2010
- *Bouncing Cosmology*  
for a *spatially flat* Friedmann universe  
Ciminelli, Nicolis, Trincherini 2010;  
Easson, Sawicki, AV;  
Taotao Qiu, Eslin, Cai, Li, Zhang 2011
- *Superinflation with blue spectra of*  
gravity waves  
Kobayashi, Yamaguchi, Yokoyama 2010



# WHY IS THIS **G** STILL INTERESTING?

The scalar field can model ***imperfect superfluid*** and should not be necessarily considered more fundamental as a velocity potential or a “phase”

Pujolas, Sawicki, AV 2011



- New linear and nonlinear effects in cosmological perturbations - theoretical laboratory to test **DE /  $\Lambda$**

Similarly to gravity, these theories are highly nonlinear and a UV completion they can admit should be by the ***classicalization*** or something more exotic - nonlocal, etc..

Dvali et al. 2010

## SIMPLEST INTERESTING SUBSECTOR OF GALILEONS /HORNDESKI'S THEORIES -

### ***Kinetic Gravity Braiding***

$$S_\phi = \int d^4x \sqrt{-g} [K(\phi, X) + G(\phi, X) \square \phi]$$

kinetic mixing / braiding  
 $\partial\phi \quad \partial g$

Armendariz-Picon, Damour, Mukhanov, Steinhardt 1999/2000

Pujolàs, Deffayet, Sawicki, AV, 2010

where  $X \equiv \frac{1}{2}g^{\mu\nu}\nabla_\mu\phi\nabla_\nu\phi$

**Minimal** coupling to gravity  $S_{\text{tot}} = S_\phi + S_{\text{EH}}$

However, derivatives of the metric are coupled to the derivatives of the scalar, provided  $G_X \neq 0$

*No Galilean symmetry!*  $\cancel{\partial_\mu\phi \rightarrow \partial_\mu\phi + c_\mu}$

# EQUATION OF MOTION I

$$L^{\mu\nu} \nabla_\mu \nabla_\nu \phi + (\nabla_\alpha \nabla_\beta \phi) Q^{\alpha\beta\mu\nu} (\nabla_\mu \nabla_\nu \phi) + \\ + Z - G_X R^{\mu\nu} \nabla_\mu \phi \nabla_\nu \phi = 0$$

**Braiding**

**EOM is of the second order:**  $L_{\mu\nu}$ ,  $Q^{\alpha\beta\mu\nu}$ ,  $Z$

**constructed from field and it's first derivatives**

**$Q^{\alpha\beta\mu\nu}$  is such that EOM is a 4D Lorentzian  
generalization of the Monge-Ampère Equation,  
always linear in  $\ddot{\phi}$**

SHOULD  $\phi$  BE FUNDAMENTAL?  
NO NOT AT ALL  
 $\phi$  CAN MODEL SOME  
HYDRODYNAMICS !

$K(X)$  for equation of state,  $G(X)$ : transport coefficient

# EQUATION OF MOTION II

- Shift-Charge Noether Current: - interpret as “particle” current  $J_\mu$
- New Equivalent Lagrangian:  $\mathcal{P}$  pressure!
- Equation of motion is a “conservation law”:

$$\nabla_\mu J^\mu = \mathcal{P}_\phi$$

# IMPERFECT FLUID FOR TIMELIKE GRADIENTS

- **Four velocity:**  $u_\mu \equiv \frac{\nabla_\mu \phi}{\sqrt{2X}}$   $\phi$  is an **internal clock**
- **projector:**  $\perp_{\mu\nu} = g_{\mu\nu} - u_\mu u_\nu$
- **Time derivative:**  $(\dot{\phantom{x}}) \equiv \frac{d}{d\tau} \equiv u^\lambda \nabla_\lambda$
- **Expansion:**  $\theta \equiv \perp_\mu^\lambda \nabla_\lambda u^\mu = \dot{V}/V$   
  
**comoving volume**

*Shift-symmetry*  
 $\phi \rightarrow \phi + c$   
*violates*  
 $\phi \rightarrow -\phi$   
*and introduces*  
**arrow of time**

## SIMPLEST INTERESTING SUBSECTOR OF GALILEONS /HORNDESKI'S THEORIES -

### ***Kinetic Gravity Braiding***

$$S_\phi = \int d^4x \sqrt{-g} [K(\phi, X) + G(\phi, X) \square \phi]$$

Armendariz-Picon, Damour, Mukhanov, Steinhardt 1999/2000

kinetic mixing / braiding  
 $\partial\phi$        $\partial g$

Pujolàs, Deffayet, Sawicki, AV, 2010

where  $X \equiv \frac{1}{2}g^{\mu\nu}\nabla_\mu\phi\nabla_\nu\phi$

**Minimal** coupling to gravity  $S_{\text{tot}} = S_\phi + S_{\text{EH}}$

However, derivatives of the metric are coupled to the derivatives of the scalar, provided  $G_X \neq 0$

*No Galilean symmetry!*  $\cancel{\partial_\mu\phi \rightarrow \partial_\mu\phi + c_\mu}$

# EQUATION OF MOTION II

- Shift-Charge Noether Current: - interpret as “particle” current  $J_\mu$

$$J_\mu = (\mathcal{L}_X - 2G_\phi) \nabla_\mu \phi - G_X \nabla_\mu X$$

- New Equivalent Lagrangian:  $\mathcal{P}$  pressure!

- Equation of motion is a “conservation law”:

$$\{\nabla_\mu J^\mu = \mathcal{P}_\phi\}$$

# EQUATION OF MOTION II

- Shift-Charge Noether Current: - interpret as “particle” current  $J_\mu$

$$J_\mu = (\mathcal{L}_X - 2G_\phi) \nabla_\mu \phi - G_X \nabla_\mu X$$

- New Equivalent Lagrangian:  $\mathcal{P}$  pressure!

$$\mathcal{P} = K - 2XG_\phi - G_X \nabla^\lambda \phi \nabla_\lambda X$$

- Equation of motion is a “conservation law”:

$$\{\nabla_\mu J^\mu = \mathcal{P}_\phi\}$$

# IMPERFECT FLUID FOR TIMELIKE GRADIENTS

- **Four velocity:**  $u_\mu \equiv \frac{\nabla_\mu \phi}{\sqrt{2X}}$   $\phi$  is an **internal clock**
- **projector:**  $\perp_{\mu\nu} = g_{\mu\nu} - u_\mu u_\nu$
- **Time derivative:**  $(\dot{\phantom{x}}) \equiv \frac{d}{d\tau} \equiv u^\lambda \nabla_\lambda$
- **Expansion:**  $\theta \equiv \perp_\mu^\lambda \nabla_\lambda u^\mu = \dot{V}/V$   
 **comoving volume**

*Shift-symmetry*  
 $\phi \rightarrow \phi + c$   
*violates*  
 $\phi \rightarrow -\phi$   
*and introduces*  
**arrow of time**

# IMPERFECT FLUID FOR TIMELIKE GRADIENTS

- **Four velocity:**  $u_\mu \equiv \frac{\nabla_\mu \phi}{\sqrt{2X}}$   $\phi$  is an **internal clock**
- **projector:**  $\perp_{\mu\nu} = g_{\mu\nu} - u_\mu u_\nu$
- **Time derivative:**  $(\dot{\phantom{x}}) \equiv \frac{d}{d\tau} \equiv u^\lambda \nabla_\lambda$
- **Expansion:**  $\theta \equiv \perp_\mu^\lambda \nabla_\lambda u^\mu = \dot{V}/V$   
 **comoving volume**

*Shift-symmetry*  
 $\phi \rightarrow \phi + c$   
*violates*  
 $\phi \rightarrow -\phi$   
*and introduces*  
**arrow of time**

# IMPERFECT FLUID FOR TIMELIKE GRADIENTS

- **Four velocity:**  $u_\mu \equiv \frac{\nabla_\mu \phi}{\sqrt{2X}}$   $\phi$  is an **internal clock**
- **projector:**  $\perp_{\mu\nu} = g_{\mu\nu} - u_\mu u_\nu$
- **Time derivative:**  $(\dot{\phantom{x}}) \equiv \frac{d}{d\tau} \equiv u^\lambda \nabla_\lambda$
- **Expansion:**  $\theta \equiv \perp_\mu^\lambda \nabla_\lambda u^\mu = \dot{V}/V$   
 **comoving volume**

*Shift-symmetry*  
 $\phi \rightarrow \phi + c$   
*violates*  
 $\phi \rightarrow -\phi$   
*and introduces*  
**arrow of time**

# EFFECTIVE MASS & CHEMICAL POTENTIAL

$$\kappa \equiv 2XG_X$$

- charge density:  $n \equiv J^\mu u_\mu = n_0 + \kappa \theta$   
“Braiding”
- energy density:  $\mathcal{E} \equiv T^{\mu\nu} u_\mu u_\nu = \mathcal{E}_0 + \theta \dot{\phi} \kappa$
- effective mass per shift-charge / chemical potential:

$$m \equiv \left( \frac{\partial \mathcal{E}}{\partial n} \right)_{V,\phi} = \sqrt{2X} = \dot{\phi}$$

## SHIFT-CURRENT AND “DIFFUSION”

$$J_\mu = n u_\mu - \frac{\kappa}{m} \perp_\mu^\lambda \nabla_\lambda m$$

“Diffusion”

§ 59, L&L, vol. 6

$$\kappa \equiv 2XG_X$$

Is a “diffusivity”/  
transport coefficient

Particle / charge current is not parallel to energy flow!

# IMPERFECT FLUID ENERGY-MOMENTUM TENSOR

- Pressure

$$\mathcal{P} = P_0 - \kappa \dot{m}$$

- Energy Flow

*No Heat Flux!*

$$q_\mu = -\kappa \perp_\mu^\lambda \nabla_\lambda m = m \perp_\mu^\lambda J_\lambda$$

- Energy Momentum Tensor

$$T_{\mu\nu} = \mathcal{E} u_\mu u_\nu - \perp_{\mu\nu} \mathcal{P} + 2u_{(\mu} q_{\nu)}$$

Solving for  $\dot{m}$  for small gradients or small  $\kappa$  one obtains “bulk viscosity”  
but change of frame results in a perfect fluid with vorticity up to  $\mathcal{O}(\kappa^2)$

# DIFFUSION OF CHARGE?

For incompressible motion  $\theta \equiv 0$   
equation of motion is:

$$\dot{n} = -\bar{\nabla}_\mu (\mathfrak{D} \bar{\nabla}^\mu n) + \mathfrak{D} a^\mu \bar{\nabla}_\mu n$$

where the diffusion constant:

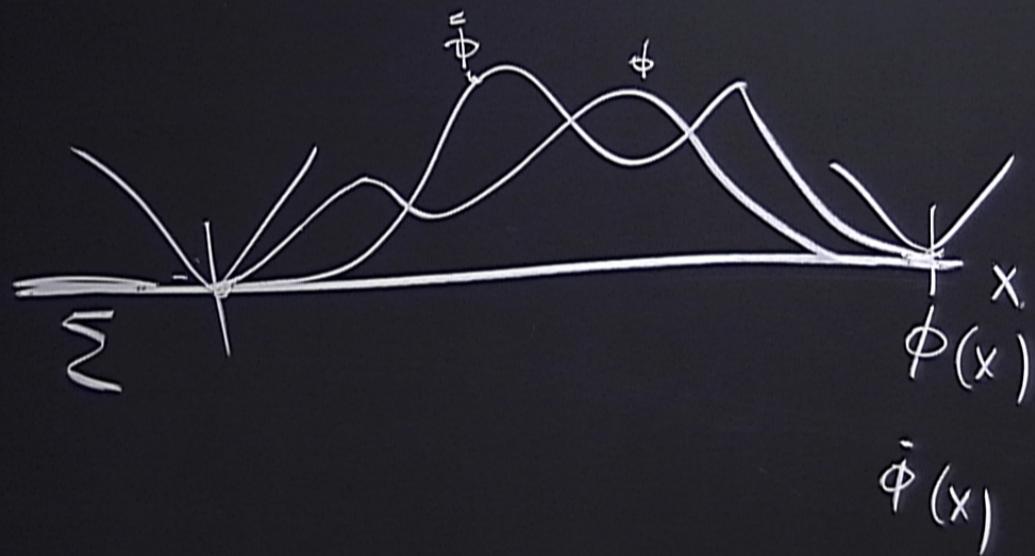
c.f. § 59, *L&L*, vol. 6, p 232

$$\mathfrak{D} \equiv -\frac{\kappa}{n_m m}$$

4-acceleration:  $a^\mu \equiv \dot{u}^\mu$

spatial gradient:  $\bar{\nabla}_\mu \equiv \perp_\mu^\nu \nabla_\nu$

$$K(x,t) + G(x,t) \Box \phi.$$



## ENERGY CONSERVATION IN COMOVING VOLUME

Energy conservation:  $u_\nu \nabla_\mu T^{\mu\nu} = 0$



$$dE = -\mathcal{P}dV + m d\mathcal{N}_{\text{dif}}$$

Euler relation:  $\mathcal{E} = mn - P_0$



Momentum conservation:

$$\perp_{\mu\nu} \nabla_\lambda T^{\lambda\nu} = 0$$

# COSMOLOGY

$$q_\mu = 0 \text{ and } \theta = 3H$$

**Friedmann Equation:**

$$H^2 = \kappa m H + \frac{1}{3} (\mathcal{E}_0 + \rho_{\text{ext}})$$

$$r_c^{-1} = \kappa m \quad \text{"crossover" scale in DGP}$$

# EQUATION OF MOTION IN COSMOLOGY

## (CHARGE CONSERVATION)

$$\dot{n} + 3Hn = \mathcal{P}_\phi$$

If there is shift-symmetry then

$$\mathcal{P}_\phi = 0$$



$$n \propto a^{-3}$$

# EQUATION OF MOTION IN COSMOLOGY

## (CHARGE CONSERVATION)

$$\dot{n} + 3Hn = \mathcal{P}_\phi$$

If there is shift-symmetry then

$$\mathcal{P}_\phi = 0$$



$$n \propto a^{-3}$$

# ACTION FOR THE COSMOLOGICAL PERTURBATIONS

$$S_2 = \int d^3x dt A \left[ \dot{\zeta}^2 - \frac{c_s^2}{a^2} (\partial_i \zeta)^2 \right]$$

where  $A = \frac{X a^3}{(H - m\kappa/2)^2} D$

$$D = \frac{\mathcal{E}_m - 3H\kappa}{m} + \frac{3}{2}\kappa^2$$



Controls “ghosts”  $D > 0$  No ghosts!

# SOUND SPEED

$$c_s^2 = \frac{\mathcal{P}_m + 2\dot{\kappa} + \kappa(4H - \kappa m/2)}{\mathcal{E}_m - 3\kappa(H - \kappa m/2)} \neq \frac{\dot{\mathcal{P}}}{\dot{\mathcal{E}}}$$

The relation between the equation of state, the sound speed and the presence of ghosts is very different from the k-*essence* & perfect fluid.



A manifestly stable *Phantom* ( $w_X < -1$ ) is possible even with a *single* degree of freedom and *minimal* coupling to gravity

## ACTION FOR THE COSMOLOGICAL PERTURBATIONS

$$S_2 = \int d^3x dt A \left[ \dot{\zeta}^2 - \frac{c_s^2}{a^2} (\partial_i \zeta)^2 \right]$$

where  $A = \frac{X a^3}{(H - m\kappa/2)^2} D$

$$D = \frac{\mathcal{E}_m - 3H\kappa}{m} + \frac{3}{2}\kappa^2$$



Controls “ghosts”  $D > 0$  No ghosts!

# **G BOUNCE IDEA I**

- Consider matter with **constant equation of state - radiation, dust, spatial curvature etc...**  $p_{\text{ext}} = w\rho_{\text{ext}}$  with  $w = \text{const}$
- Shift-symmetric Lagrangian for the scalar field  $K (\cancel{X})$  and  $G (\cancel{X})$
- Phase space is two dimensional  $(m, \rho_{\text{ext}})$
- Go to new coordinates  $(m, H)$  by solving the Friedmann equation
- Integral of motion 1st integral

$$I(m, H) = \frac{n^{1+w}}{\rho_{\text{ext}}}$$

# **G BOUNCE IDEA I**

- Consider matter with **constant equation of state - radiation, dust, spatial curvature etc...**  $p_{\text{ext}} = w\rho_{\text{ext}}$  with  $w = \text{const}$
- Shift-symmetric Lagrangian for the scalar field  $K (\cancel{X})$  and  $G (\cancel{X})$
- Phase space is two dimensional  $(m, \rho_{\text{ext}})$
- Go to new coordinates  $(m, H)$  by solving the Friedmann equation
- Integral of motion 1st integral

$$I(m, H) = \frac{n^{1+w}}{\rho_{\text{ext}}}$$

# **G** BOUNCE IDEA II

- In new coordinates  $(m, H)$  at  $H = 0$  pose conditions

$$\rho_{\text{ext}} > 0, D > 0, \dot{H} > 0, c_s^2 > 0$$

- These conditions are on the range of chemical potential  $\Delta m$  and on the equation of state  $K(m)$  along with the transport coefficient  $\kappa(m)$  or on  $K(X)$  and  $G(X)$



**“Healthy” Bounce!**

IS IT POSSIBLE  
TO SATISFY  
ALL THESE CONDITIONS ?



# SIMPLE HIERARCHY

**if**  $K(m)$  and  $\kappa(m)$  satisfy the hierarchy:

$$K > mK_m > \frac{1}{2}(m\kappa)^2 > 0 > \frac{1}{2}m^2K_{mm} > -\frac{3}{4}(m\kappa)^2$$

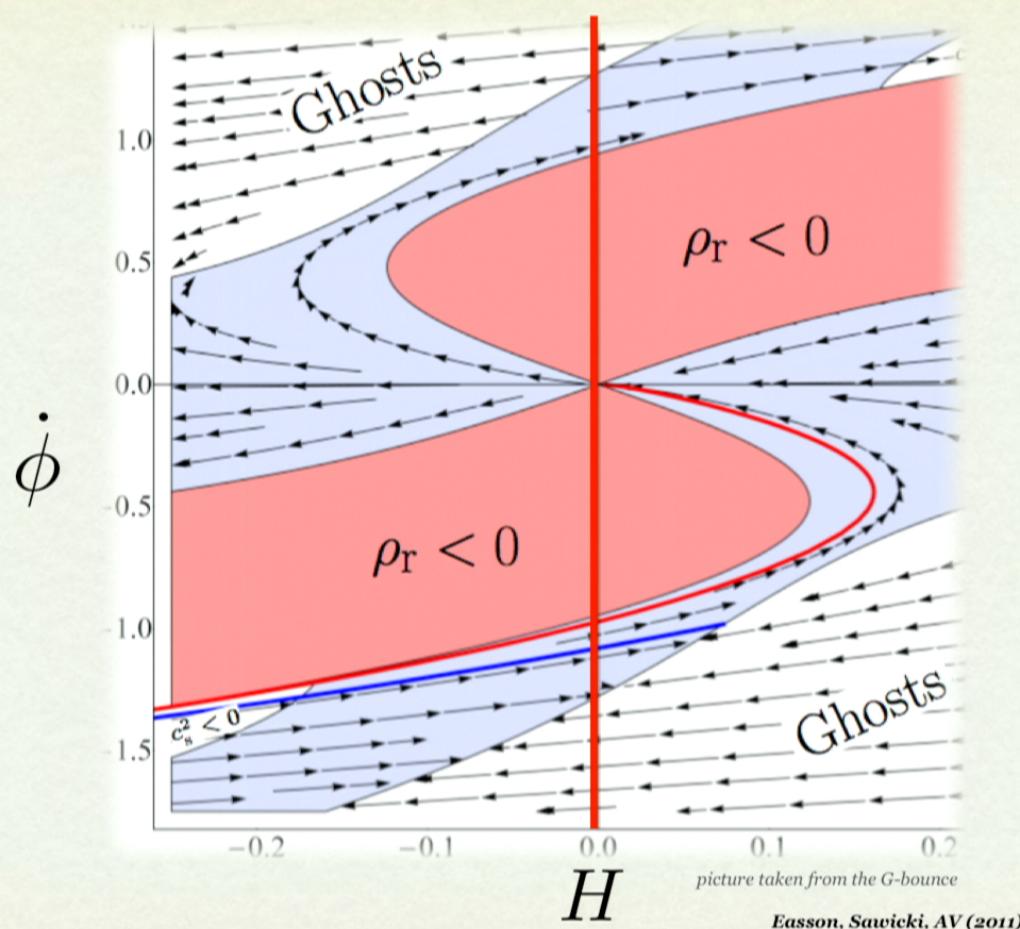
**then** at  $H = 0$  :

$$\rho_{\text{ext}} > 0, D > 0, \dot{H} > 0, c_s^2 > 0$$

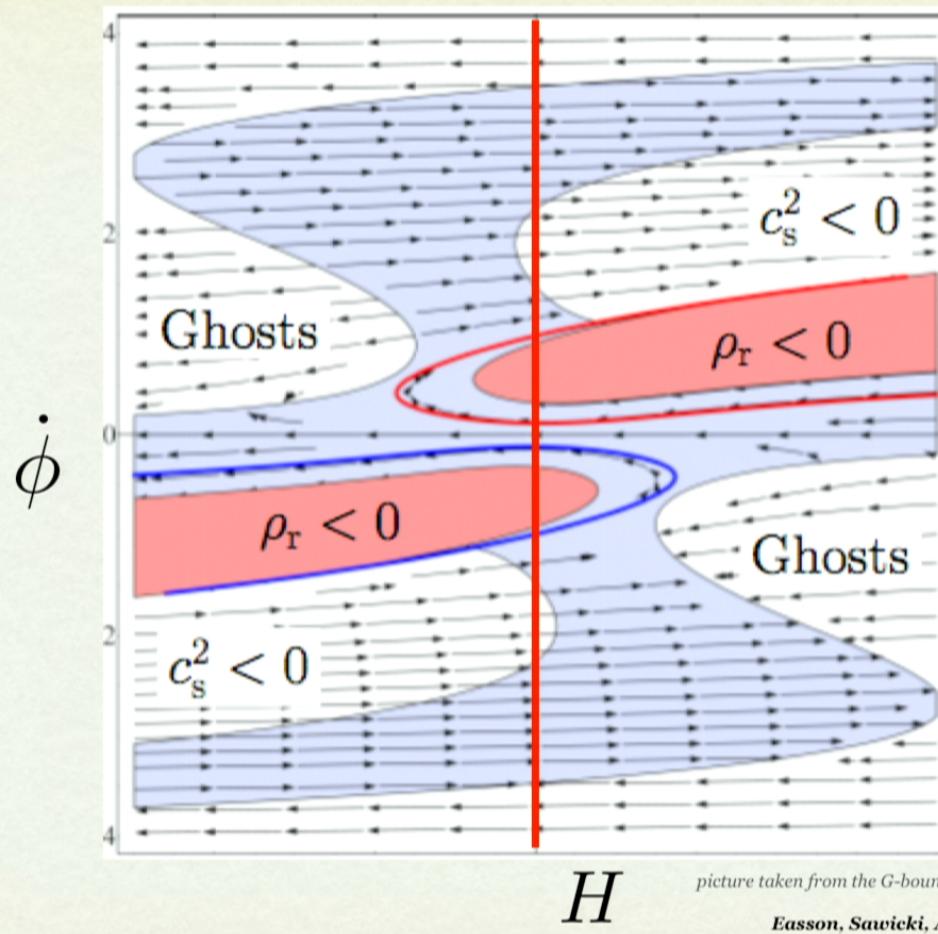
does not  
depend on  
external equation  
of state  $w$ !



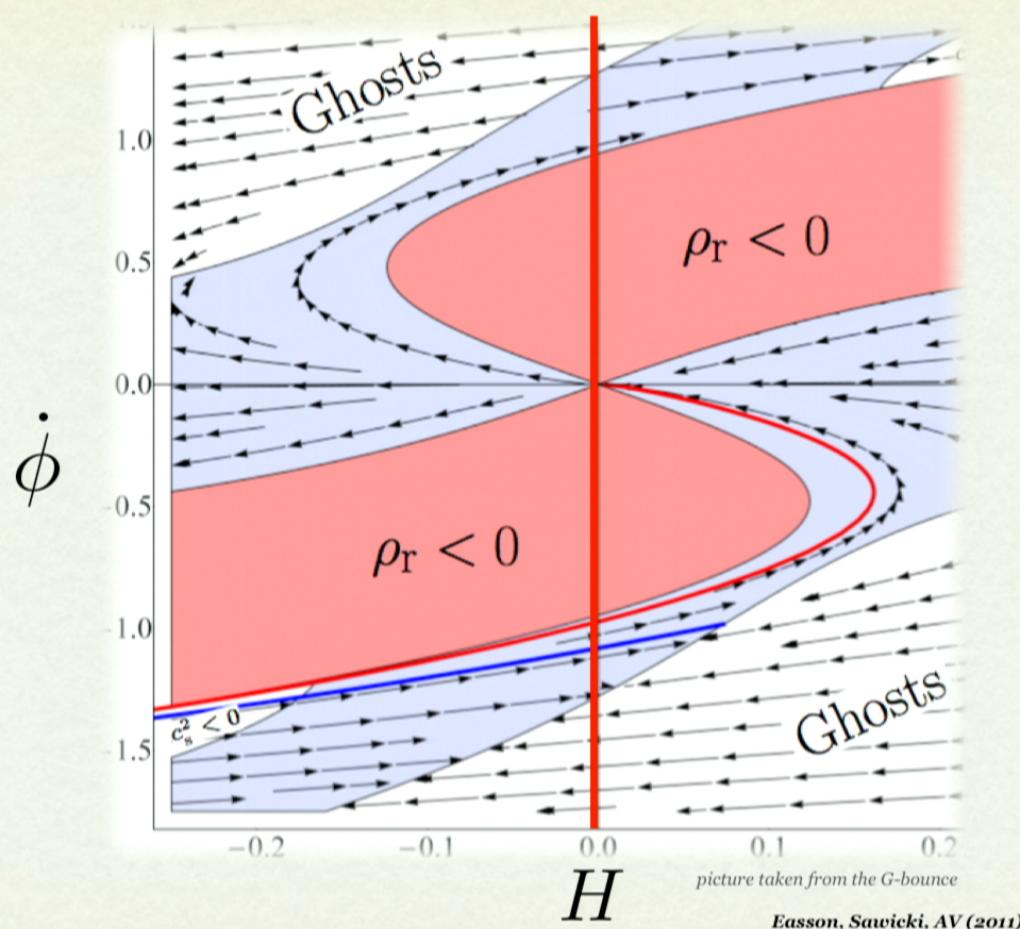
**“Healthy” Bounce!**



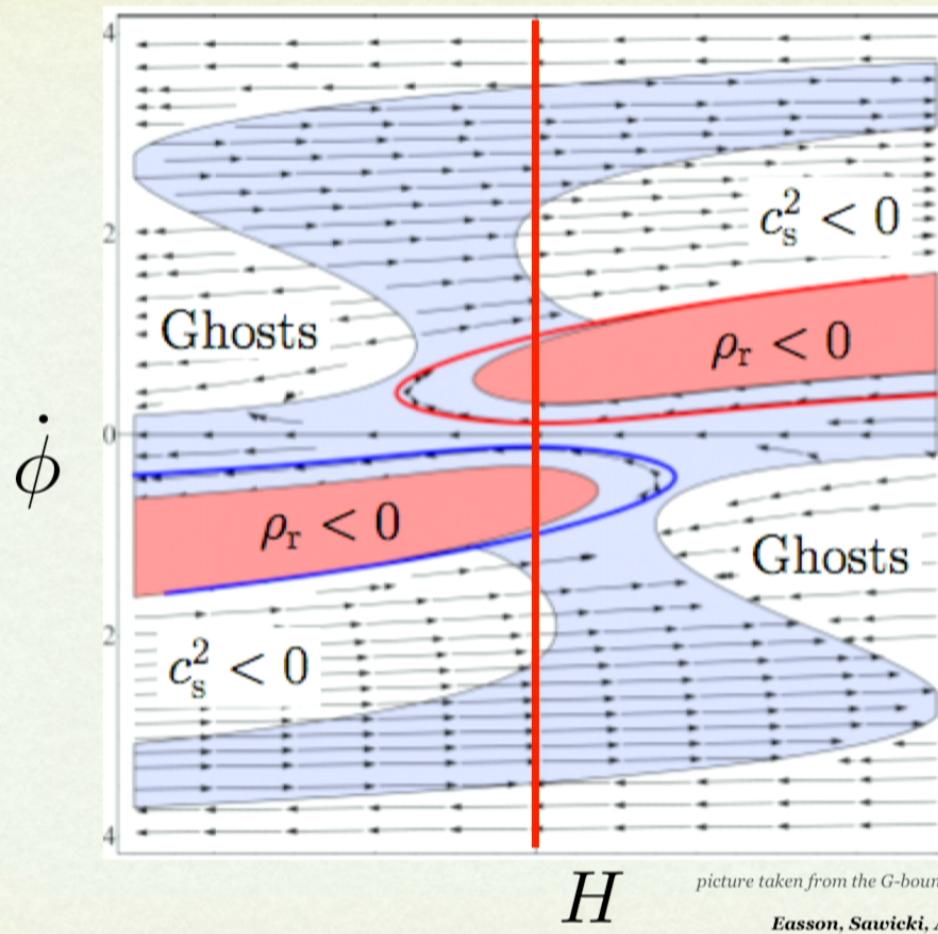
**G-Bounce:**  $\mathcal{L} = X - \alpha X^3 + \varkappa X \square \phi + \text{Radiation } \rho_r$



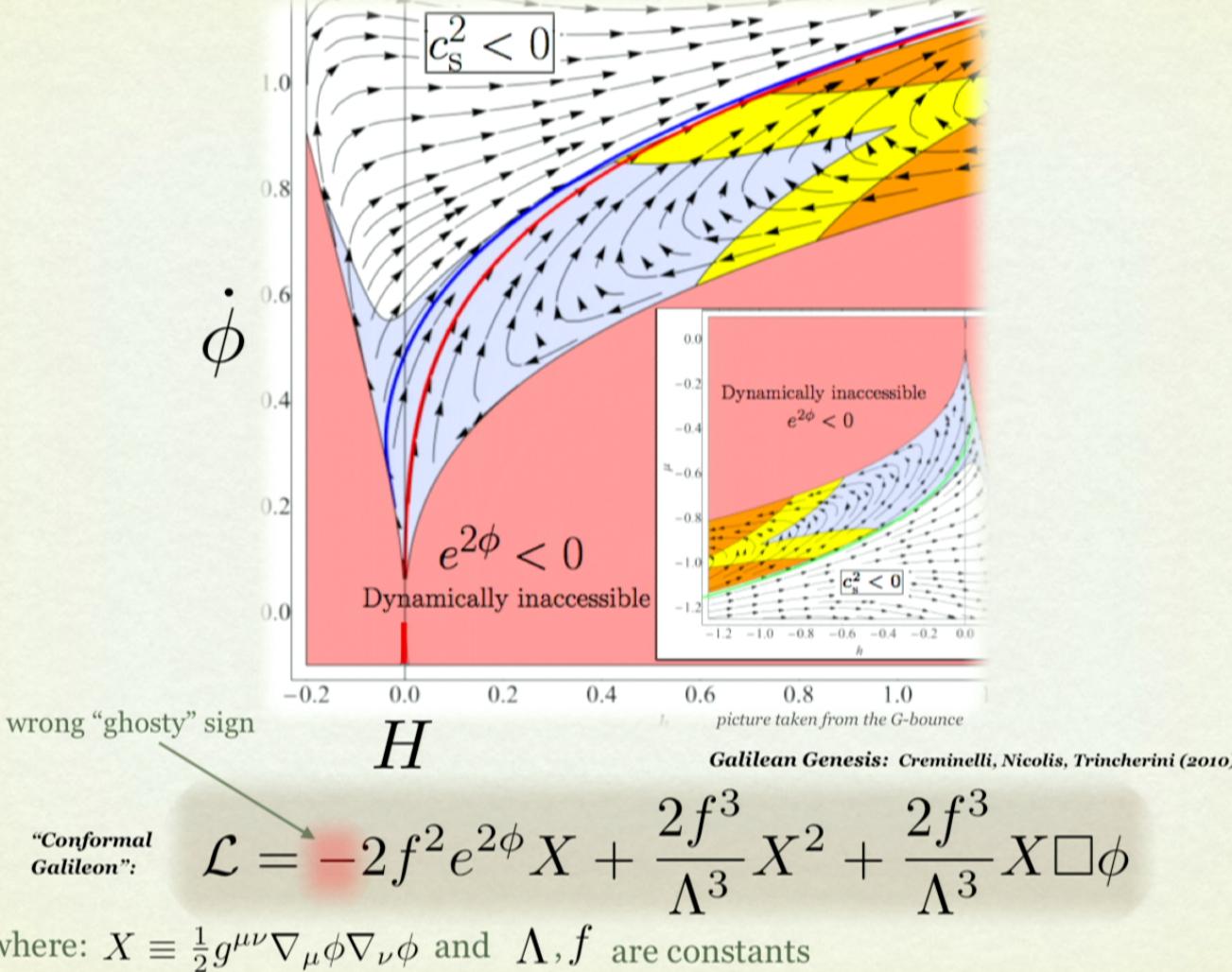
$$\kappa = 2\beta X + \gamma X^2 \quad K = -\Lambda + X - X^3 + \text{Radiation} \quad \rho_r$$



**G-Bounce:**  $\mathcal{L} = X - \alpha X^3 + \varkappa X \square \phi + \text{Radiation } \rho_r$



$$\kappa = 2\beta X + \gamma X^2 \quad K = -\Lambda + X - X^3 + \text{Radiation} \quad \rho_r$$



# OPEN QUESTIONS

- Strong coupling?
- Possible anisotropy? Too strong tachyonic / Jeans instabilities?
- Can one arrange a cyclic i.e. oscillating evolution?
- Can one avoid all singularities and troubles for the past?
- Perturbations?
- Any *realistic* scenarios? Smooth transition to standard cosmology?

# OPEN QUESTIONS

- Strong coupling?
- Possible anisotropy? Too strong tachyonic / Jeans instabilities?
- Can one arrange a cyclic i.e. oscillating evolution?
- Can one avoid all singularities and troubles for the past?
- Perturbations?
- Any *realistic* scenarios? Smooth transition to standard cosmology?

THANKS A LOT FOR  
YOUR ATTENTION!