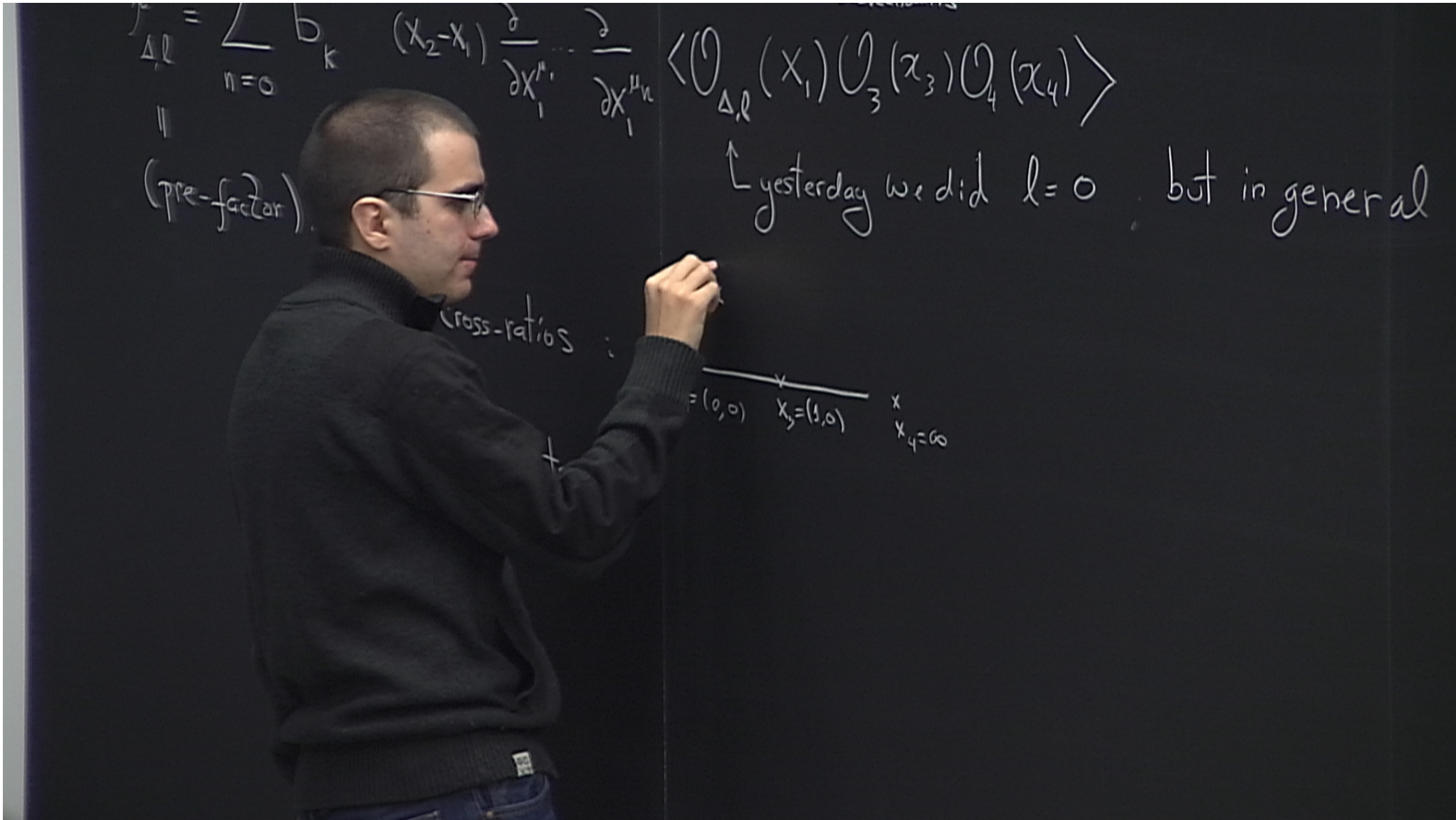


Title: Explorations in String Theory - Lecture 10

Date: Mar 23, 2012 11:30 AM

URL: <http://pirsa.org/12030053>

Abstract:



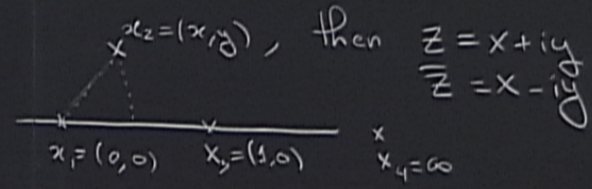
$$\langle \mathcal{O}_1(x_1) \dots \mathcal{O}_4(x_4) \rangle = \sum_K C_{12K} C_{K34} \mathcal{F}_K^{(12)(34)}$$

\uparrow primaries \uparrow conformal block containing all descendants

$$\mathcal{F}_{\Delta, l} = \sum_{n=0}^{\infty} b_K^{\mu_1 \dots \mu_n} (x_2 - x_1)^{\mu_1} \frac{\partial}{\partial x_1^{\mu_1}} \dots \frac{\partial}{\partial x_1^{\mu_n}} \langle \mathcal{O}_{\Delta, l}(x_1) \mathcal{O}_3(x_3) \mathcal{O}_4(x_4) \rangle$$

\uparrow yesterday we did $l=0$ but in general

$(\text{pre-factor}) \times G_{\Delta, l}(z, \bar{z})$
 \uparrow cross-ratios
 Some (product of) Hypergeometrics



so 1 complex CR
 or
 2 scalar CR's

\uparrow
 slight confusion yesterday?

general

$$0 \neq \langle \mathcal{O}_1(x_1) \mathcal{O}_2(x_2) \mathcal{O}(x_3) \rangle$$

CR

CR's

sign? day

fixed by
Conf
Symmetry

only they can appear for
 \mathcal{O}_1 and \mathcal{O}_2 Scalars

$\mu_1 \dots \mu_l$

Sym traceless

spin l
(scalar is $l=0$)

Casimir of the Conformal group

$$\vec{J}^2 = -D^2 - \frac{1}{2} (K_\mu P_\mu + P_\mu K_\mu) + \frac{1}{2} M_{\mu\nu} M^{\mu\nu}$$

$$\vec{J}^2 |O_{\Delta, l}\rangle = (\Delta(\Delta - d) + l(l + d - 2)) |O_{\Delta, l}\rangle$$

$$J = (P, K, P, M)$$

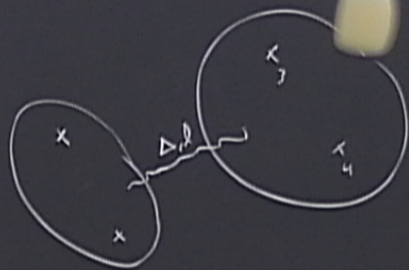
$$M_{\mu\nu} M^{\mu\nu}$$

$$|U_{\Delta, e}\rangle$$

$$\left[J_A^{(1)} + J_A^{(2)} + J_A^{(3)} + J_A^{(4)} \right] \langle 0, \dots, 0_4 \rangle = 0 \quad \leftarrow \text{Conf invariance}$$

$$\mathcal{F}_K^{(12)(34)} = 0$$

acts on X_4



$$\left[J^{(3)} + J^{(4)} \right]^2 \mathcal{F}_K^{(12)(34)} = \left(J^{(3)} + J^{(4)} \right)^2 \text{ acting \#}$$

$$\langle \mathcal{O}_1(x_1) \dots \mathcal{O}_4(x_4) \rangle = \sum_K C_{12K} C_{K34} \mathcal{F}_K^{(1234)}$$

↑ primaries Δ, l

↑ conformal block containing all descendants

fixed by conf. sym

$$\mathcal{F}_{\Delta, l} = \sum_{n=0}^{\infty} b_K^{m_1 \dots m_n} (x_2 - x_1)^{\frac{\partial}{\partial x_1^{m_1}} \dots \frac{\partial}{\partial x_1^{m_n}}} \langle \mathcal{O}_{\Delta, l}(x_1) \mathcal{O}_3(x_3) \mathcal{O}_4(x_4) \rangle$$

#

$$\langle \mathcal{O}_{\Delta, l}(x_1) \mathcal{O}_3(x_3) \mathcal{O}_4(x_4) \rangle$$

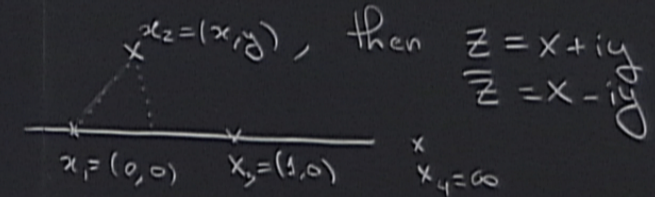
↑ yesterday we did $l=0$ but in

||

(pre-factor) $\times G_{\Delta, l}(z, \bar{z})$

cross-ratios

Some (product of) Hypergeometrics



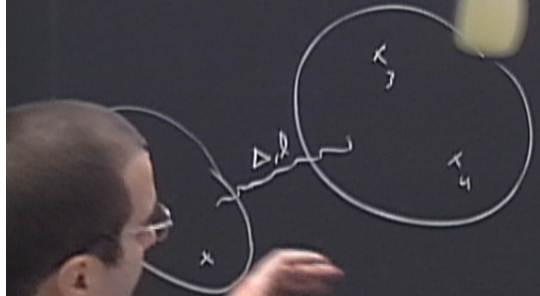
so 1 complex or 2 real

↑ slight con yest

$$\left[J_A^{(1)} + J_A^{(2)} + J_A^{(3)} + J_A^{(4)} \right] \langle 0, \dots, 0_4 \rangle = 0 \quad \leftarrow \text{Conf invariance}$$

$$\mathcal{F}_K^{(12)(34)} = 0$$

acts on X_4



$$\left[J^{(3)} + J^{(4)} \right]^2 \mathcal{F}_K^{(12)(34)} = \left(J^{(3)} + J^{(4)} \right)^2 \text{ acting \#} = \left(J^{(1)} \right)^2 \text{ acting on \#}$$

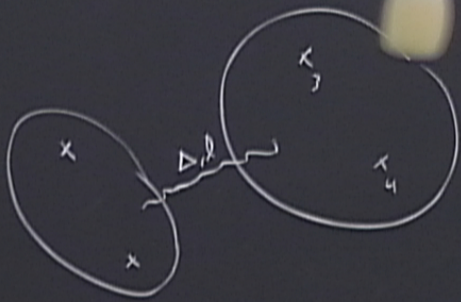
$$\text{but } \left(J_A^{(3)} + J_A^{(4)} + J_A^{(1)} \right) \langle 0_{\Delta, l}(x_1) 0_3(x_3) 0_4(x_4) \rangle = 0$$

$$\text{but } \left(J^{(1)} \right)^2 \langle 0_{\Delta, l} \dots \rangle = \left(\Delta(\Delta-d) + l(l+d-z) \right) \langle 0_{\Delta, l} \dots \rangle$$

$$\left[\mathbb{J}_A^{(1)} + \mathbb{J}_A^{(2)} + \mathbb{J}_A^{(3)} + \mathbb{J}_A^{(4)} \right] \langle 0, \dots, 0_4 \rangle = 0 \quad \leftarrow \text{Conf invariance}$$

$$\mathbb{F}_K^{(12)(34)} = 0$$

acts on x_4



Somewhat eq

$$\left[\mathbb{J}^{(3)} + \mathbb{J}^{(4)} \right]^2 \mathbb{F}_K^{(12)(34)} = \left(\mathbb{J}^{(3)} + \mathbb{J}^{(4)} \right)^2 \text{ acting on } \# = \left(\mathbb{J}^{(1)} \right)^2 \text{ acting on } \#$$

$$\text{but } \left(\mathbb{J}_A^{(3)} + \mathbb{J}_A^{(4)} + \mathbb{J}_A^{(1)} \right) \langle 0_{\Delta, l}(x_1) 0_3(x_3) 0_4(x_4) \rangle = 0$$

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$$[\text{Dif op}] \mathbb{F}_{\Delta, l} = C_{\Delta, l} \mathbb{F}_{\Delta, l}$$

$$\left[J_A^{(1)} + J_A^{(2)} + J_A^{(3)} + J_A^{(4)} \right] \langle 0, \dots, 0_4 \rangle = 0$$

acts on X_4

$$\left. \begin{aligned} & (J_3^A + J_4^A) (J_3^A + J_4^A + J_1^A - J_1^A) \\ & = -J_1^A (J_3^A - J_4^A) = +J_1^A J_1^A \end{aligned} \right\}$$

← Conf invariance

$$\mathcal{F}_K^{(12)(34)} = 0$$

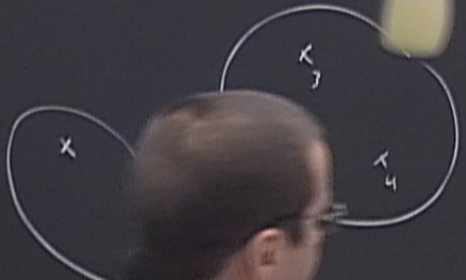
Somewhat eq

$$\left[J^{(3)} + J^{(4)} \right]^2 \mathcal{F}_K^{(12)(34)} = (J^{(3)} + J^{(4)})^2 \text{ acting } \# = (J^{(1)})^2 \text{ acting on } \#$$

$$\text{but } (J_A^{(3)} + J_A^{(4)} + J_A^{(1)}) \langle 0_{\Delta, l}(x_1) 0_3(x_3) 0_4(x_4) \rangle = 0$$

$$\text{but } (J^{(1)})^2 \langle 0_{\Delta, l} \dots \rangle = (\Delta(\Delta-d) + l(l+d-2)) \langle 0_{\Delta, l} \dots \rangle$$

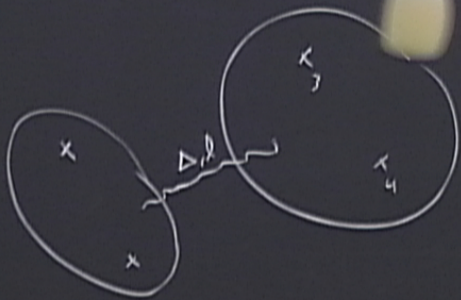
$$\Delta, l = C_{\Delta, l} \mathcal{F}_{\Delta, l}$$



$$\left[\mathbb{J}_A^{(1)} + \mathbb{J}_A^{(2)} + \mathbb{J}_A^{(3)} + \mathbb{J}_A^{(4)} \right] \langle 0, \dots, 0_4 \rangle = 0$$

acts on X_4

$$\mathbb{F}_K^{(12)(34)} = 0$$



Somewhat eq

$$\left[\mathbb{J}^{(3)} + \mathbb{J}^{(4)} \right]^2 \mathbb{F}_K^{(12)(34)} = \left(\mathbb{J}^{(3)} + \mathbb{J}^{(4)} \right)^2 \text{ acting on } \# = \left(\mathbb{J}^{(1)} \right)^2 \text{ acting on } \#$$

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$$[\text{Dif op}] \mathbb{F}_{\Delta, l} = C_{\Delta, l} \mathbb{F}_{\Delta, l}$$

$$\left. \begin{aligned} & \left(\mathbb{J}_3^A + \mathbb{J}_4^A \right) \left(\mathbb{J}_3^A + \mathbb{J}_4^A + \mathbb{J}_1^A - \mathbb{J}_1^A \right) \\ & = -\mathbb{J}_1^A \left(\mathbb{J}_3^A - \mathbb{J}_4^A \right) = +\mathbb{J}_1^A \mathbb{J}_1^A \end{aligned} \right\}$$

← Conf invariance

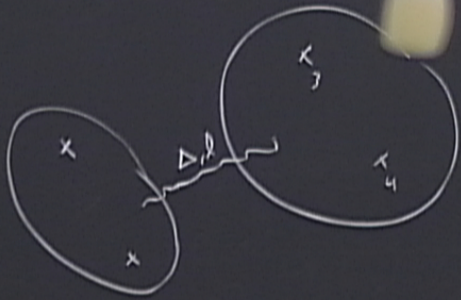
$$\left[J_A^{(1)} + J_A^{(2)} + J_A^{(3)} + J_A^{(4)} \right] \langle 0, \dots, 0_4 \rangle = 0$$

acts on X_4

$$\left. \begin{aligned} & (J_3^A + J_4^A) (J_3^A + J_4^A + J_1^A - J_1^A) \\ & = -J_1^A (J_3^A - J_4^A) = +J_1^A J_1^A \end{aligned} \right\}$$

← Conf invariance

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Somewhat eq

$$\left[J^{(3)} + J^{(4)} \right]^2 \mathcal{F}_K^{(12)(34)} = (J^{(3)} + J^{(4)})^2 \text{ acting on } \# = (J^{(1)})^2 \text{ acting on } \#$$

$$\text{but } (J_A^{(3)} + J_A^{(4)} + J_A^{(1)}) \langle 0_{\Delta, l}(x_1) 0_3(x_3) 0_4(x_4) \rangle = 0$$

$$\text{but } (J^{(1)})^2 \langle 0_{\Delta, l} \dots \rangle = (\Delta(\Delta-d) + l(l+d-z)) \langle 0_{\Delta, l} \dots \rangle$$

$$[\text{Dif op}] \mathcal{F}_{\Delta, l} = C_{\Delta, l} \mathcal{F}_{\Delta, l}$$

$$\mathcal{D} = z^2 (1 - \bar{z}^2) \frac{z}{z^2} + \frac{(\Delta_1 - \Delta_2)(\Delta_3 - \Delta_4)}{4} (z + \bar{z}) + \dots$$

$$\mathcal{D} G_{\Delta, l} = C_{\Delta, l} G_{\Delta, l}$$

Casimir of the Conformal group

$$J = (P, K, P, M)$$

spin l
(l=0)

$$\vec{J}^2 = -D^2 - \frac{1}{2} (K_\mu P_\mu + P_\mu K_\mu) + \frac{1}{2} M_{\mu\nu} M^{\mu\nu}$$

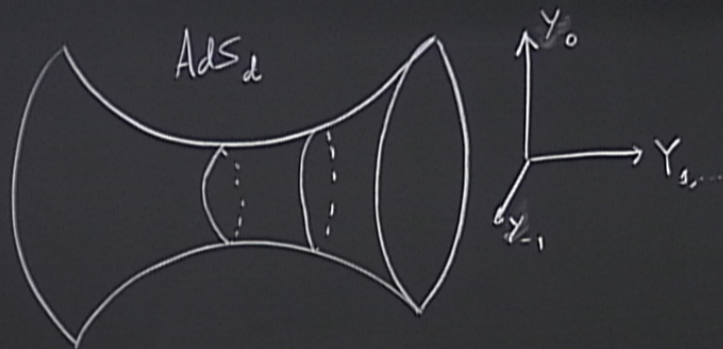
$$\vec{J}^2 |U_{\Delta, l}\rangle = (\Delta(\Delta - d) + l(l + d - 2)) |U_{\Delta, l}\rangle$$

$\leftarrow C_{\Delta, l}$

* Conf Block for odd d are not known.



[Diso
DG



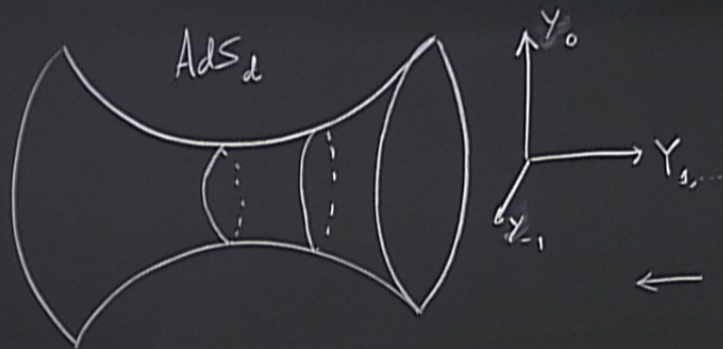
$$Y_{-1}^2 + Y_0^2 - Y_1^2 - \dots - Y_{d-1}^2 = 1$$

in $\mathbb{R}^{2, d-1}$

Poincaré Coordinates

$$ds^2 = \frac{dz^2 + dx_\mu dx^\mu}{z^2}$$





$$Y_{-1}^2 + Y_0^2 - Y_1^2 - \dots - Y_{d-1}^2 = 1$$

in $\mathbb{R}^{2,d-1}$

$$ds^2 = d\vec{Y} \cdot d\vec{Y} \Big|_{\vec{Y}^2 = 1}$$

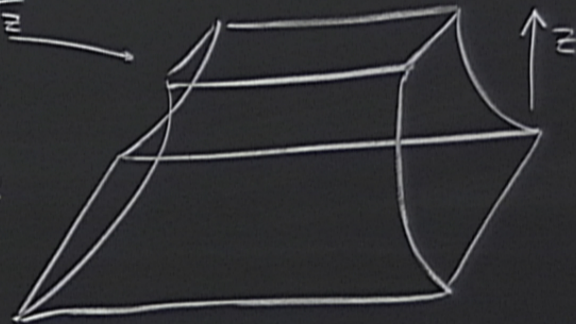
$$Y_{-1} + Y_{d-1} = \frac{1}{z}$$

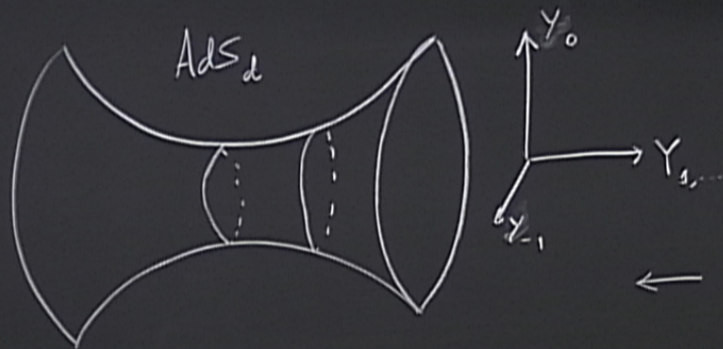
$$Y_{\mu} = \frac{x^{\mu}}{z}$$

$$\mu = 0 \dots d-2$$

Poincaré Coordinates

$$ds^2 = \frac{dz^2 + dx_{\mu} dx^{\mu}}{z^2}$$





$$Y_{-1}^2 + Y_0^2 - Y_1^2 - \dots - Y_{d-1}^2 = 1$$

in $\mathbb{R}^{2, d-1}$

$$ds^2 = d\vec{Y} \cdot d\vec{Y} \Big|_{\vec{Y}^2 = 1}$$

$$Y_{-1} + Y_{d-1} = \frac{1}{z} > 0$$

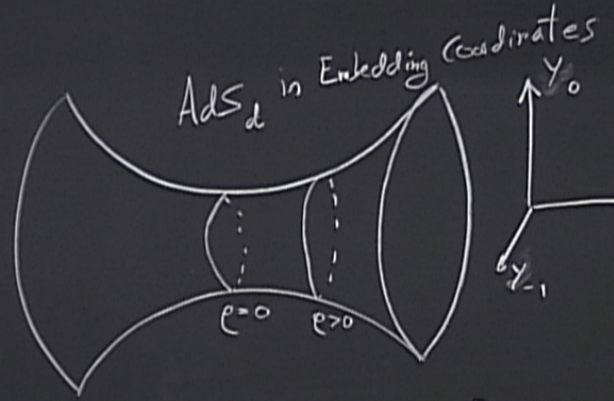
$$Y_{\mu} = \frac{x^{\mu}}{z}$$

$$\mu = 0 \dots d-2$$

Poincaré Coordinates

$$ds^2 = \frac{dz^2 + dx_{\mu} dx^{\mu}}{z^2}$$





$$Y_{-1}^2 + Y_0^2 - Y_1^2 - \dots - Y_{d-1}^2 = 1$$

in $\mathbb{R}^{2, d-1}$

$$ds^2 = d\vec{Y} \cdot d\vec{Y} \Big|_{Y^2=1}$$

$$Y_{-1} + Y_{d-1} = \frac{1}{z} > 0$$

$$Y_{\mu} = \frac{x^{\mu}}{z}$$

$$\mu = 0 \dots d-2$$

AdS in Poincaré Coordinates

$$ds^2 = \frac{dz^2 + dx_{\mu} dx^{\mu}}{z^2}$$



Y_{-1}

Y_{d-1}

Global AdS

$$Y_{-1} + iY_0 = \cosh \rho e^{it}$$

$$Y_i = \sinh \rho \Omega_i$$

\uparrow unit vector, S^{d-2}

\uparrow $1 \dots d-1$

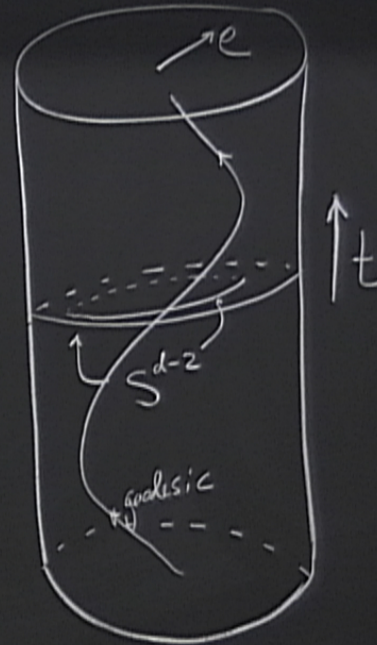
Global AdS

$$Y_{-1} + iY_0 = \cosh \rho e^{it}$$

$$Y_i = \sinh \rho \Omega_i$$

\uparrow unit vector, S^{d-2}

$$ds^2 = -\cosh^2 \rho dt^2 + d\rho^2 + \sinh^2 \rho d\Omega_{S^{d-2}}^2$$



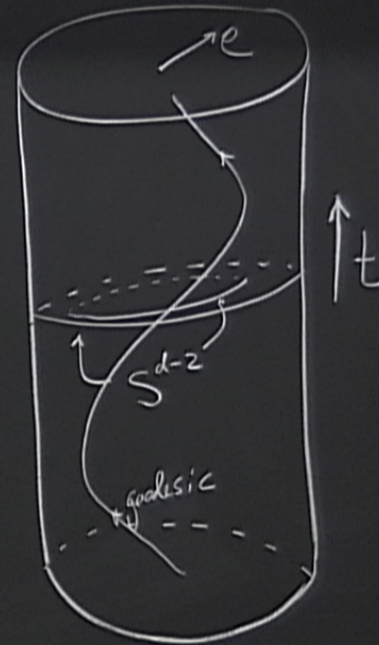
Global AdS

$$Y_{-1} + iY_0 = \cosh \rho e^{it} \quad \text{namely } t \sim t + 2\pi$$

$$Y_i = \sinh \rho \Omega_i \quad \text{unit vector, } S^{d-2}$$

\uparrow

$$ds^2 = -\cosh^2 \rho dt^2 + d\rho^2 + \sinh^2 \rho d\Omega_{S^{d-2}}^2$$



AdS_d in Embedding Coordinates

$Y_0^2 + Y_1^2 - Y_2^2 - \dots - Y_{d-1}^2 = 1$

$ds^2 = d\varphi^2 - d\vec{x}^2$ (at $Y_0 = 1$)

AdS in Poincaré Coordinates

$ds^2 = \frac{dz^2 + dx_\mu dx^\mu}{z^2}$

$Y_{-1} + Y_{d-1} = \frac{1}{z} > 0$

$Y_\mu = \frac{x^\mu}{z}$

$\mu = 0, \dots, d-2$

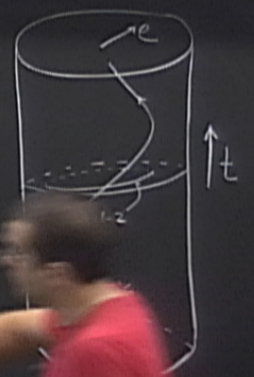
Global AdS

$Y_{-1} + iY_0 = \cosh \rho e^{it}$ (radius ρ , $t \in [t, t+2\pi]$)

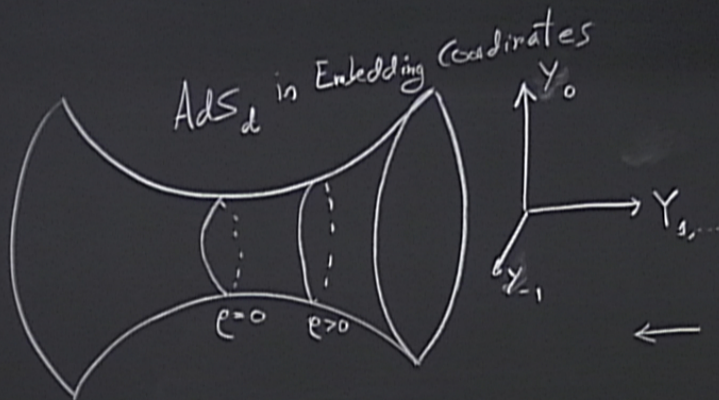
$Y_\mu = \sinh \rho \frac{\Omega_\mu}{r}$ (unit vector, S^{d-2})

$\mu = 1, \dots, d-1$

$ds^2 = -\cosh^2 \rho dt^2 + d\rho^2 + \sinh^2 \rho d\Omega_{d-2}^2$



$\partial(\text{Global AdS}) : e \rightarrow \infty, ds^2 \rightarrow e^{2\varrho} (-dt^2 + d\Omega_{S^{d-2}}^2)$



$$Y_{-1}^2 + Y_0^2 - Y_1^2 - \dots - Y_{d-1}^2 = 1$$

in $\mathbb{R}^{2,d-1}$

$$ds^2 = d\vec{Y} \cdot d\vec{Y} \Big|_{Y^2=1}$$

∂ Poincaré
 $\mathbb{R}^{1,3}$

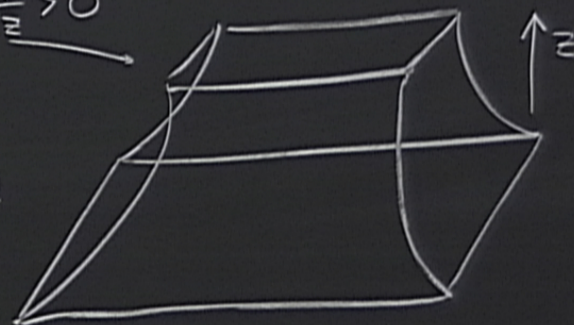
AdS in
Poincaré Coordinates

$$ds^2 = \frac{dz^2 + dx_\mu dx^\mu}{z^2}$$

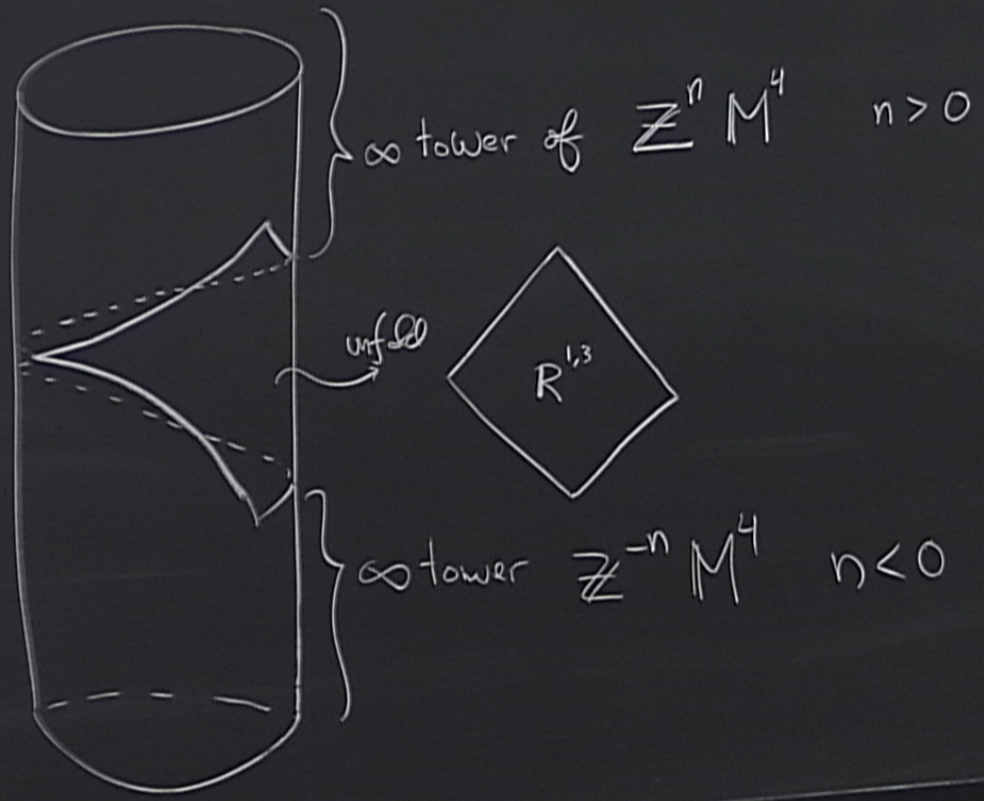
$$Y_{-1} + Y_{d-1} = \frac{1}{z} > 0$$

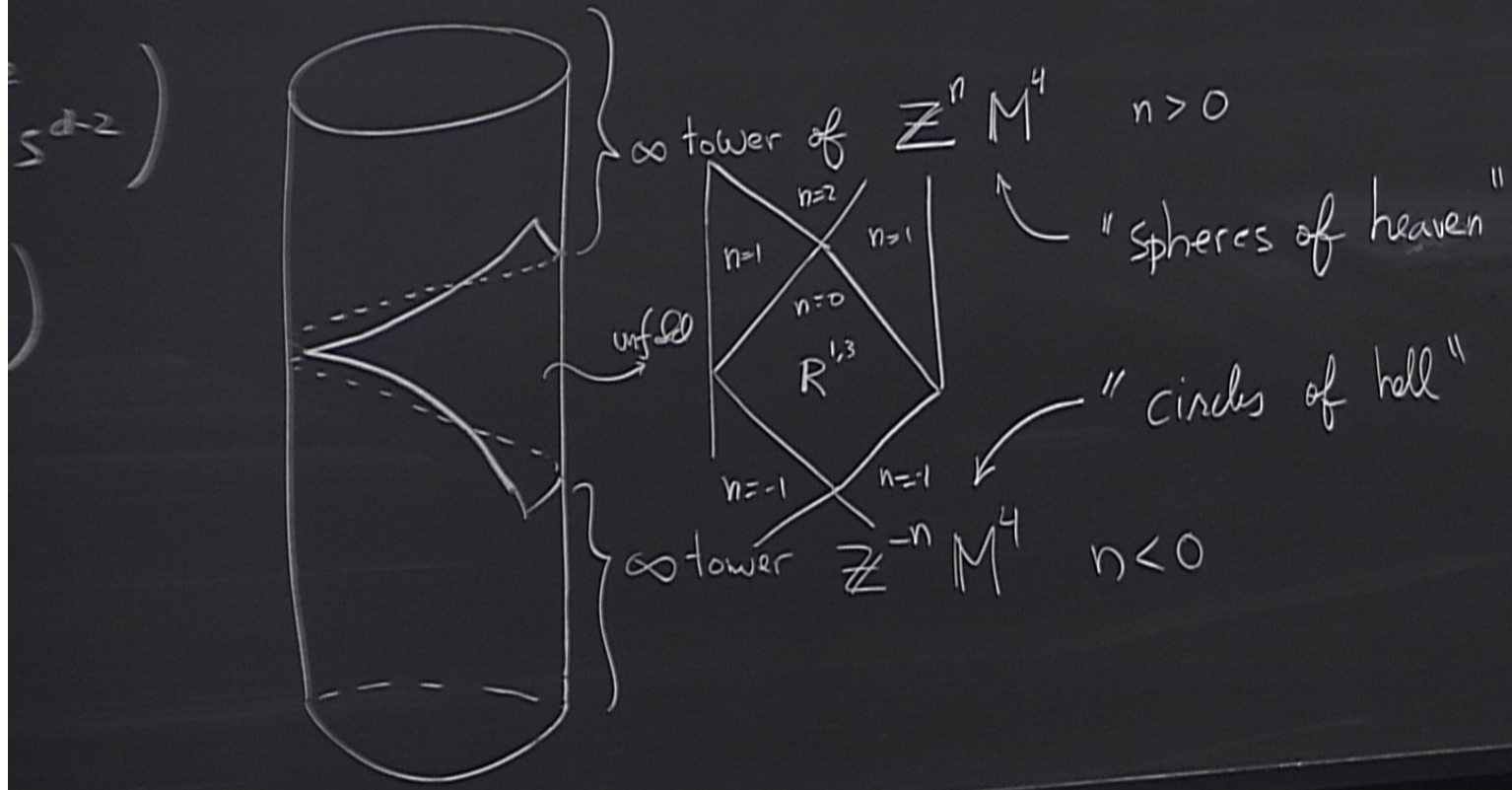
$$Y_\mu = \frac{x^\mu}{z}$$

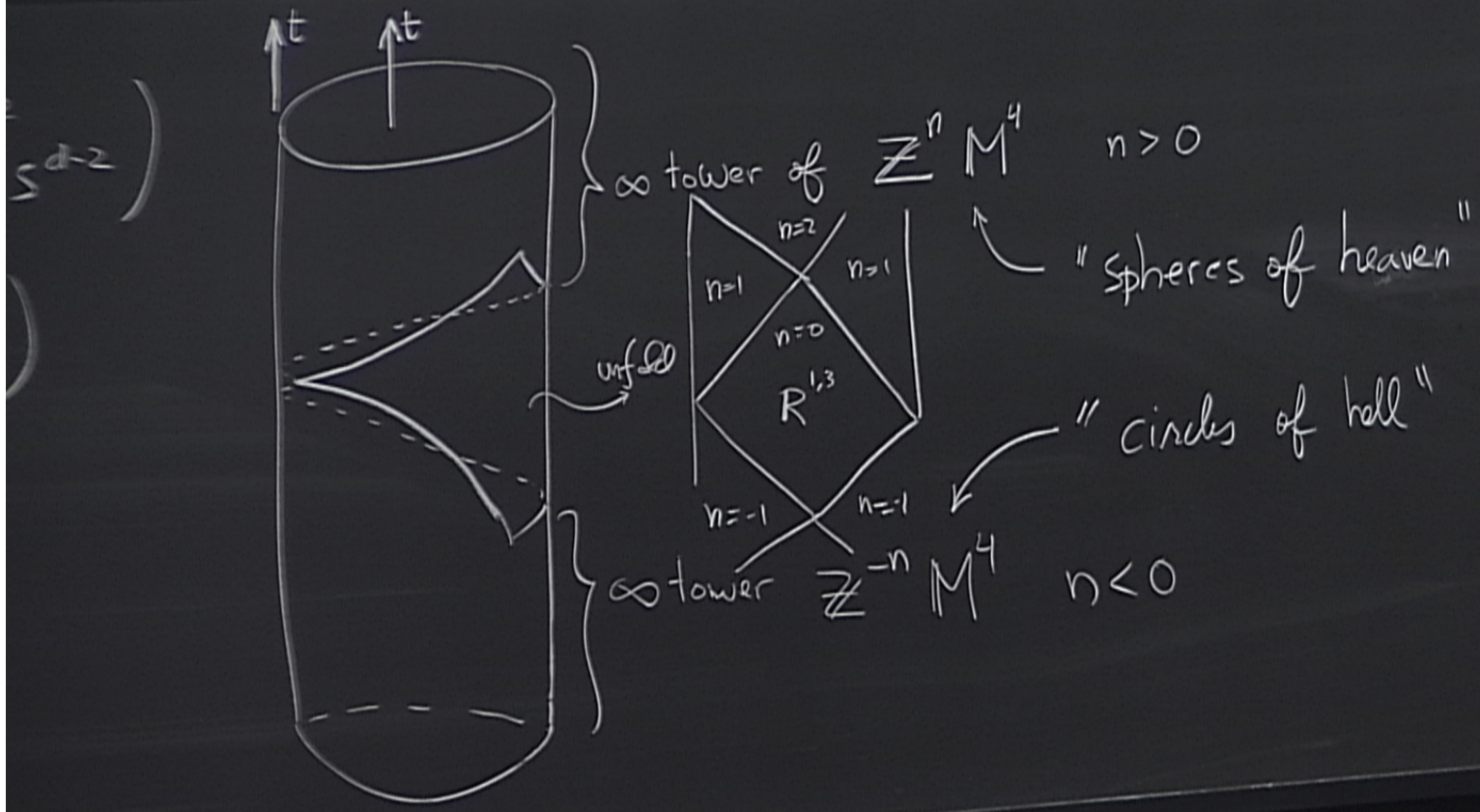
$$\mu = 0 \dots d-2$$



G
 Y_{-1}
 Y_{d-1}
 z







strings in Pohrean AdS $\equiv \mathcal{N}=4$ in \mathbb{R}^4

t in the Bulk = t in the boundary $\mathbb{R} \times S^3$

$\left\{ \begin{array}{l} \text{E string inside} \\ \text{the cylinder} \end{array} \right\} = \left\{ \begin{array}{l} \text{E} \\ \mathcal{N}=4 \text{ in } \mathbb{R} \times S^3 \end{array} \right\} \overset{\substack{\uparrow \\ \text{yesterday}}}{=} \left\{ \Delta \right\} \mathcal{N}=4 \text{ on } \mathbb{R}^4$

strings in Poincare AdS $\equiv \mathcal{N}=4$ in \mathbb{R}^4

t in the Bulk = t in the boundary $\mathbb{R} \times S^3$

$\left\{ \begin{array}{l} \text{E string inside} \\ \text{the cylinder} \end{array} \right\} = \left\{ E_{\mathcal{N}=4 \text{ in } \mathbb{R} \times S^3} \right\} \overset{\substack{\uparrow \\ \text{yesterday}}}{=} \left\{ \Delta \right\} \mathcal{N}=4 \text{ on } \mathbb{R}^4$

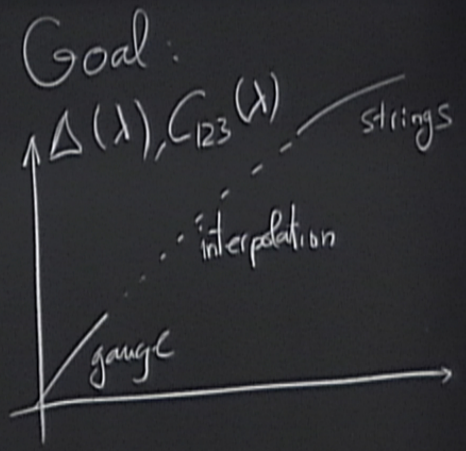
Conclusion

$$\left. \{ \Delta \} \right|_{d=4} = \left. \{ \text{Energies} \} \right|_{\text{Global AdS}}$$

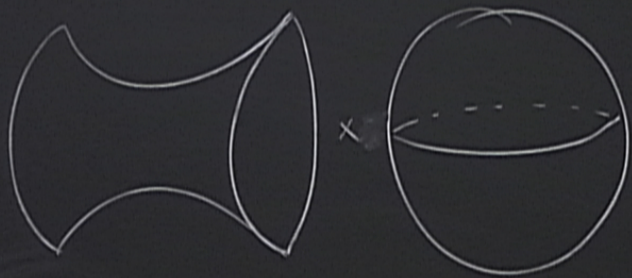
4 on \mathbb{R}^4

Conclusion

$$\left\{ \Delta(\lambda) \right\}_{d\mathcal{P}=4} = \left\{ \text{Energies}(\lambda) \right\}_{\text{Global AdS}}$$

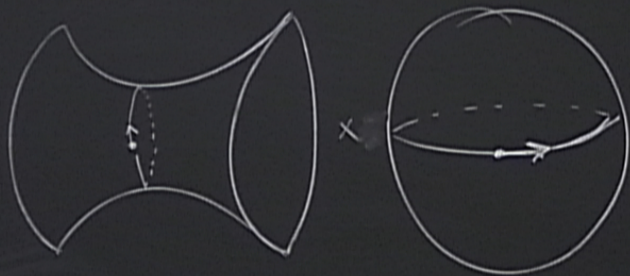


Simplest solution.



in flat space $X^\mu(\sigma, z) = p^\mu z + \text{"modes"}$
no modes

Simplest solution.

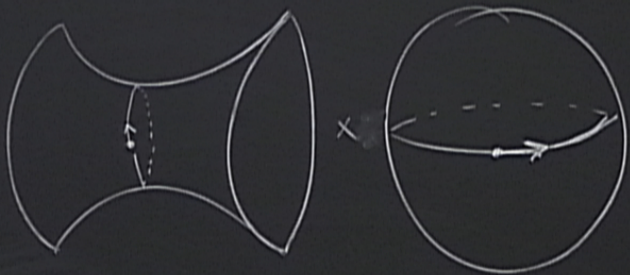


in flat space $X^\mu(\sigma, z) = p^\mu z + \text{"modes"}$
no modes

$$X_1^2 + \dots + X_5^2 = 1$$

$$X_1 + iX_2 = e^{iJz}$$

Simplest solution.



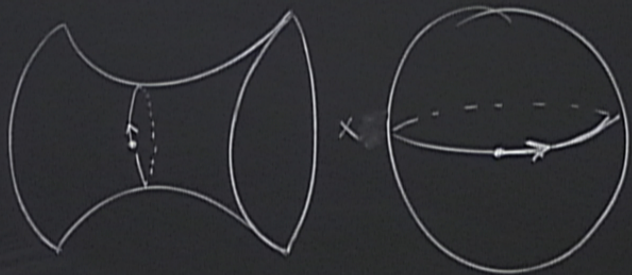
$$Y_{-1} + iY_0 = e^{iKz}$$

in flat space $X^\mu(\sigma, z) = p^\mu z + \text{"modes"}$
no modes

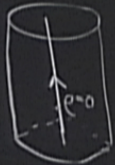
$$X_1^2 + \dots + X_6^2 = 1$$

$$X_1 + iX_2 = e^{i\sigma z}$$

Simplest solution.



$$Y_{-1} + iY_0 = e^{iKz}$$



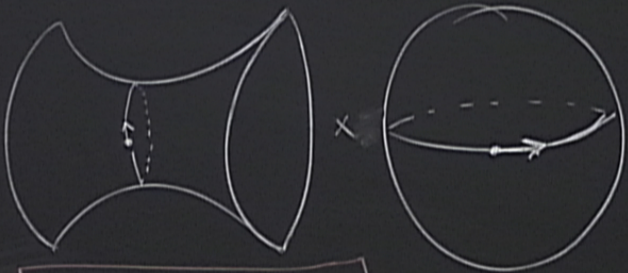
in flat space $X^\mu(\sigma, z) = p^\mu z + \text{"modes"}$
no modes

$$X_1^2 + \dots + X_6^2 = 1$$

$$X_1 + iX_2 = e^{i\sigma z}$$

Simplest solution.

in flat space $X^\mu(\sigma, z) = p^\mu z + \text{"modes"}$
no modes



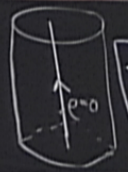
$$X_1^2 + \dots + X_5^2 = 1$$

$$X_1 + iX_2 = e^{i\sigma z}$$

$$\equiv e^{i\phi}$$

$$\phi(\sigma, z) = \sigma z$$

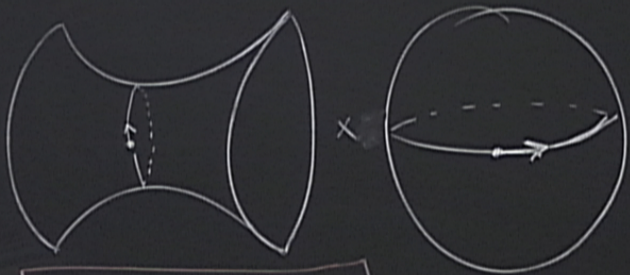
$$Y_{-1} + iY_0 = e^{iKz}$$



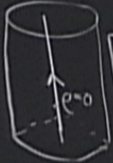
$$t(\sigma, z) = Kz$$

Simplest solution.

in flat space $X^\mu(\sigma, z) = p^\mu z + \text{"modes"}$
no modes



$$Y_{-1} + iY_0 = e^{iKz}$$



$$t(\sigma, z) = Kz$$

$$X_1^2 + \dots + X_5^2 = 1$$

$$X_1 + iX_2 = e^{i\sqrt{2}\sigma}$$

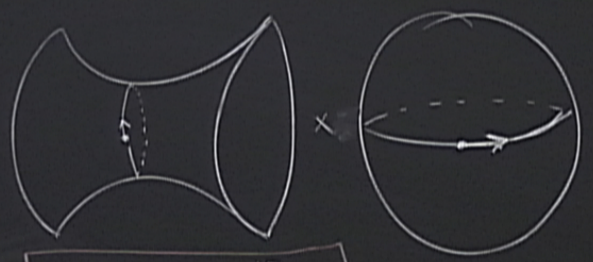
$$\equiv e^{i\phi}$$

$$\phi(\sigma, z) = \sqrt{2}\sigma$$

$$|y| = 1$$

Simplest solution.

in flat space $X^M(\sigma, z) = p^M z + \text{"modes"}$
no modes



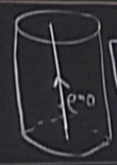
$$X_1^2 + \dots + X_6^2 = 1$$

$$X_1 + iX_2 = e^{i\sigma} z$$

\parallel
 $e^{i\phi}$

$$\phi(\sigma, z) = \sigma$$

$$Y_{-1} + iY_0 = e^{iKz}$$



$$t(\sigma, z) = Kz$$

$$S_{\text{Polyakov in conf gauge}} = \int d\sigma dz \left[\partial_a \vec{X} \cdot \partial_a \vec{X} \right]$$

$$X^\mu(\sigma, z) = p^\mu z + \text{"~~modes~~"}$$

no modes

$$\int_{\text{Polyakov in conf. gauge}} = \int d\sigma dz \left[\partial_a \vec{X} \cdot \partial_a \vec{X} + \lambda_{\text{F}} (\vec{X} \cdot \vec{X} - 1) \right]$$

$$z = \sigma$$

$$X^\mu(\sigma, z) = p^\mu z + \text{"~~modes~~"}$$

no modes

$$\int_{\text{Polyakov in conf. gauge}} d\sigma dz \left[\partial_a \vec{X} \cdot \partial_a \vec{X} + \lambda_{\vec{X}} (\vec{X} \cdot \vec{X} - 1) - \partial_a \vec{Y} \cdot \partial_a \vec{Y} - \lambda_{\vec{Y}} (\vec{Y} \cdot \vec{Y} - 1) \right]$$

$$z = \int \sigma$$

$$X^\mu(\sigma, z) = p^\mu z + \text{"~~modes~~"}$$

no modes

$$\int_{\text{Polyakov in conf. gauge}} d\sigma dz \left[\partial_a \vec{X} \cdot \partial_a \vec{X} + \lambda_{\vec{X}} (\vec{X} \cdot \vec{X} - 1) - \partial_a \vec{Y} \cdot \partial_a \vec{Y} - \lambda_{\vec{Y}} (\vec{Y} \cdot \vec{Y} - 1) \right] + \text{fermions}$$

$$z = \sqrt{\sigma}$$

$$X^\mu(\sigma, z) = p^\mu z + \text{"modes"}$$

~~no modes~~

$$S_{\text{Polyakov in conf. gauge}} = \int d\sigma dz \left[\partial_a \vec{X} \cdot \partial_a \vec{X} + \lambda_{\mathbb{H}} (\vec{X} \cdot \vec{X} - 1) - \partial_a \vec{Y} \cdot \partial_a \vec{Y} - \lambda_{\mathbb{H}} (\vec{Y} \cdot \vec{Y} - 1) \right] + \text{fermions}$$

+++++

2

coordinates
 $dx^0 dx^1 dx^2 dx^3 dx^4 dx^5 dx^6 dx^7 dx^8 dx^9$

Global AdS



$$|\vec{Y}|=1$$

$$\vec{X}(\sigma+2\pi, z) = \vec{X}(\sigma, z), \text{ same for } \vec{Y}$$

$$\text{Virasoro: } (\partial_z \vec{Y} \pm \partial_\sigma \vec{Y}) \cdot (\partial_z \vec{Y} \pm \partial_\sigma \vec{Y}) = \text{something with } \vec{Y} \rightarrow \vec{X}$$

$$\text{E.o.m for } \vec{X}: \partial_a \partial_a \vec{X} - \lambda_X \vec{X} = 0$$

$$\text{E.o.m for } \lambda_X: \vec{X}^2 = 1$$

$$|\vec{Y}|=1$$

$$\vec{X}(\sigma+2\pi, z) = \vec{X}(\sigma, z), \text{ same for } \vec{Y}$$

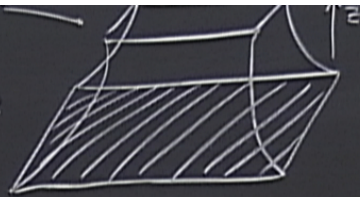
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$$\text{E.o.m for } \vec{X}: \partial_a \partial_a \vec{X} - \lambda_X \vec{X} = 0 \quad \# \quad \vec{X} \cdot \# \Rightarrow \lambda = \vec{X} \cdot \partial^2 \vec{X}$$


$$\text{E.o.m for } \lambda_X: \vec{X}^2 = 1$$

$Y_0^2 + Y_1^2 - Y_2^2 - \dots - Y_{d-1}^2 = 1$
 $\mathbb{R}^{2, d-1}$
 $d^2 = d\vec{Y} \cdot d\vec{Y} \Big|_{\vec{Y}^2=1}$

$Y_\mu = \frac{x^\mu}{z}$
 $\mu=0, \dots, d-2$



$\vec{e}_i = \sinh \rho \vec{e} - \text{unit vector}, S^{d-2}$
 $d\vec{s}^2 = -\cosh^2 \rho dt^2 + d\rho^2 + \sinh^2 \rho d\Omega_{S^{d-2}}$
 $(\sigma, z), t(\sigma, z), \phi(\sigma, z)$



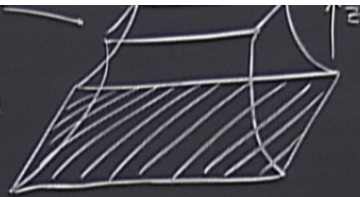
$\vec{X}(\sigma + 2\pi, z) = \vec{X}(\sigma, z)$, same for \vec{Y}

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E.o.m for λ_μ : $\vec{X}^2 = 1$

$Y_0^2 + Y_1^2 - Y_2^2 - \dots - Y_{d-1}^2 = 1$
 $\mu = 0, d-2$
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$d\Omega^2 = -\cosh^2 \rho dt^2 + d\rho^2 + \sinh^2 \rho d\Omega_{S^{d-2}}^2$
 $(\sigma, z), t(\sigma, z), \phi(\sigma, z)$

unit vector, S^{d-2}
 S^{d-2}
 intrinsic

$\vec{X}(\sigma + 2\pi, z) = \vec{X}(\sigma, z)$, same for \vec{Y}

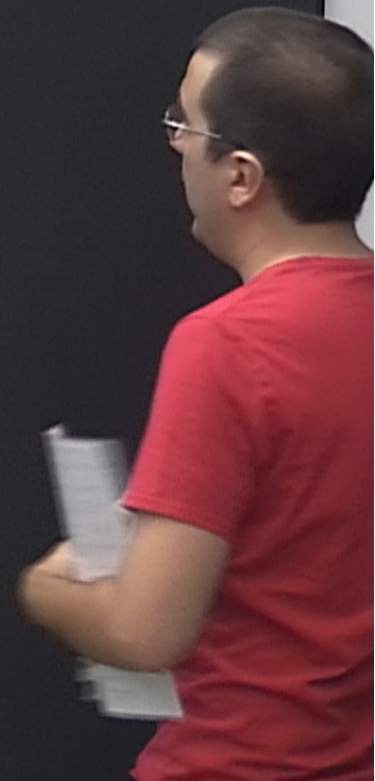
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E.o.m for λ_X : $\vec{X}^2 = 1$

$\partial^2 \vec{X} + (\partial \vec{X} \cdot \partial \vec{X}) \vec{X} = 0$

$$\begin{aligned} \partial \vec{X} \cdot \partial \vec{X} &= J^2 \\ \partial \vec{Y} - \partial \vec{Y} &= K^2 \end{aligned} \quad , \quad \partial^2 \vec{X} + J^2 \vec{X} = 0 \quad ? \quad \checkmark$$



$$\partial \vec{X} \cdot \partial \vec{X} = J^2$$

$$\partial \vec{Y} \cdot \partial \vec{Y} = K^2$$

$$\text{Virasoro: } K = J$$

$$\partial^2 \vec{X} + J^2 \vec{X} = 0 \quad ? \quad \checkmark$$

$$\partial^2 \vec{Y} + K^2 \vec{Y} = 0 \quad \checkmark$$

What is the physical meaning of K, J ?

Noether Charges:

$$\text{Energy} = \frac{\sqrt{\lambda}}{4\pi} \int d\sigma (Y_0 \dot{L}_1 - Y_{-1} \dot{Y}_0) = \frac{\sqrt{\lambda}}{2\pi} \int_0^{2\pi} d\sigma \dot{t}$$

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$$\Delta = \text{Energy} = \sqrt{\lambda} K$$
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Virasoro: $K = J$

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Virasoro: $K = J$ What is the physical meaning of K, J ?

$$\Delta = J_1 \longleftrightarrow \text{dual operator } \mathcal{O}$$

$$\Delta \mathcal{O} \xrightarrow{\lambda \gg 1} J_1$$

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$$\text{Energy} = \frac{\sqrt{\lambda}}{4\pi} \int d\sigma (Y_0 \dot{Y}_1 - Y_1 \dot{Y}_0) = \frac{\sqrt{\lambda}}{2\pi} \int_0^{2\pi} d\sigma \dot{t}$$

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Virasoro: $K = J$

What is the physical meaning of K, J ?

$$\Delta = J_1$$

dual operator $\mathcal{O} = \text{tr}(\phi_1 + i\phi_2)$
 $\Delta(\mathcal{O}) \xrightarrow{\lambda \gg 1} J_1$

Noether Charges:

$$\text{Energy} = \frac{\sqrt{\lambda}}{4\pi} \int d\sigma (Y_0 \dot{Y}_1 - Y_1 \dot{Y}_0) = \frac{\sqrt{\lambda}}{2\pi} \int_0^{2\pi} d\sigma \dot{t}$$

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Virasoro: $K = J$ What is the physical meaning of K, J ?

$\Delta = J_1$ \longleftrightarrow dual operator $O = \text{tr}(\phi_1 + i\phi_2)$

$\Delta(O) \xrightarrow{\lambda \gg 1} J_1$

$\xrightarrow{\lambda \rightarrow 0} J_1 = \# \text{ scalars!}$