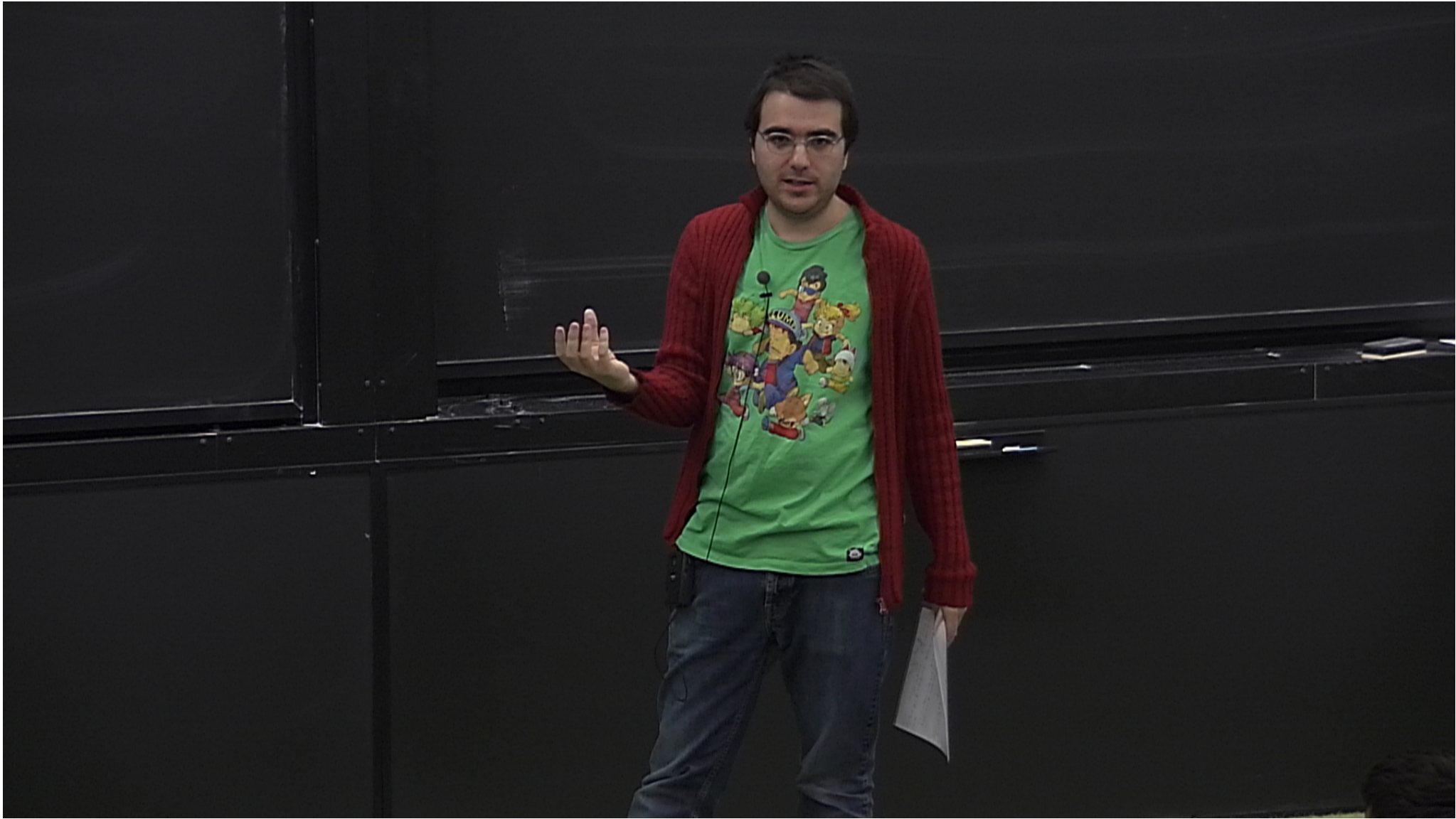


Title: Explorations in String Theory - Lecture 3

Date: Mar 14, 2012 11:30 AM

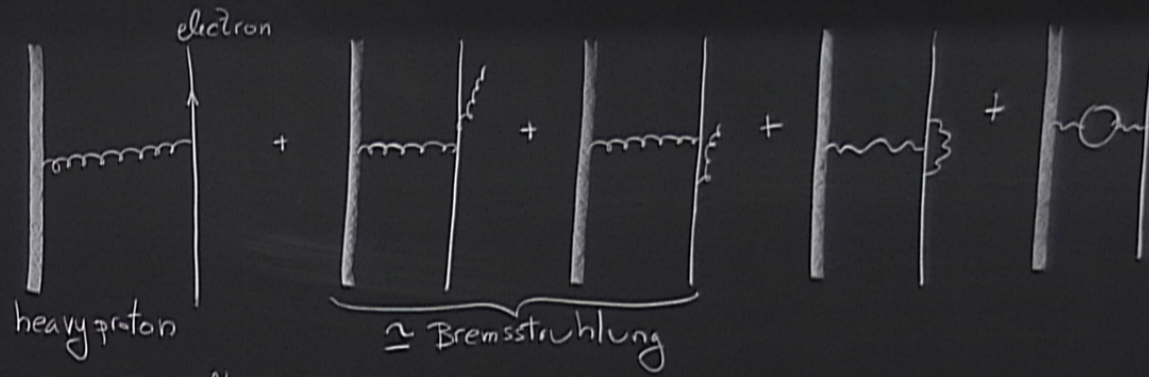
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Abstract:



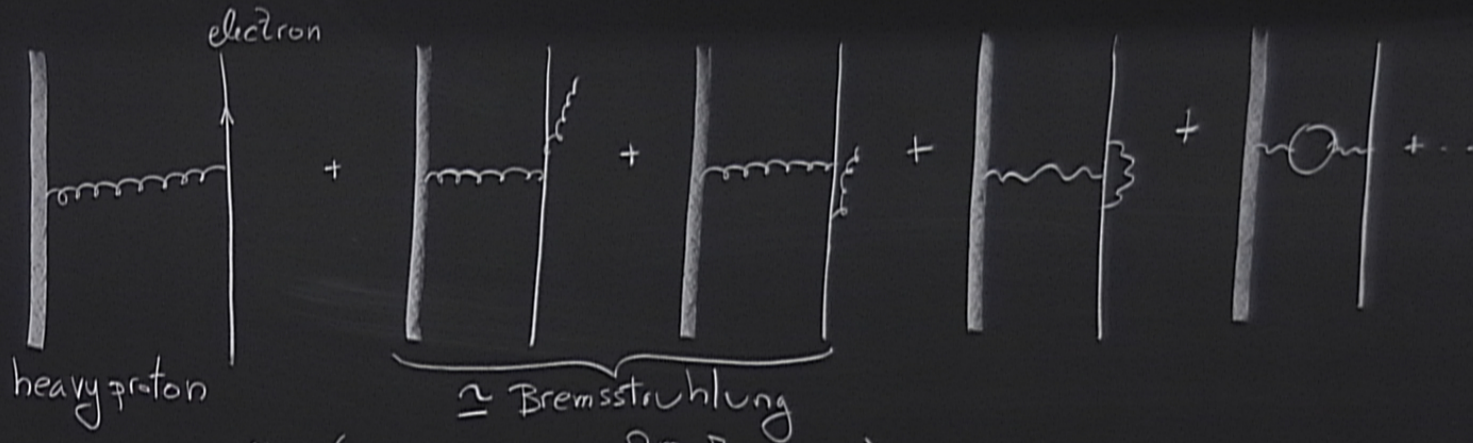
... we are therefore lead to conjecture that

4D, $\mathcal{N}=4$ SYM = type IIB strings in $AdS_5 \times S^5$



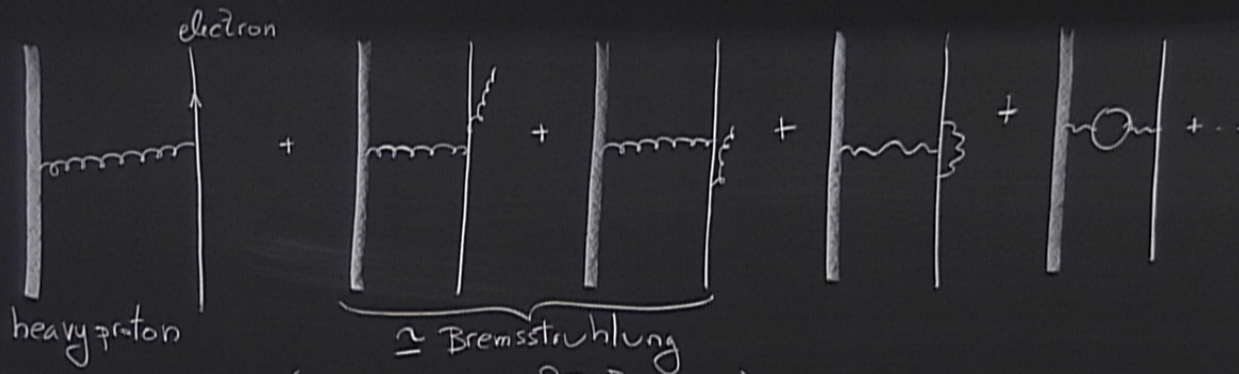
$$V(x) = -\frac{\alpha}{x}$$

Coulomb



$$V(r) = -\frac{\alpha}{r} \left(1 + \frac{\alpha}{4\sqrt{\pi}} \frac{e^{-2m_e r}}{(m_e r)^{3/2}} + \dots \right)$$

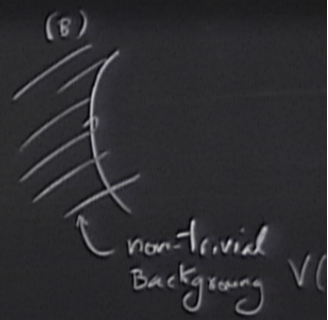
Coulomb



≈ Bremsstrahlung

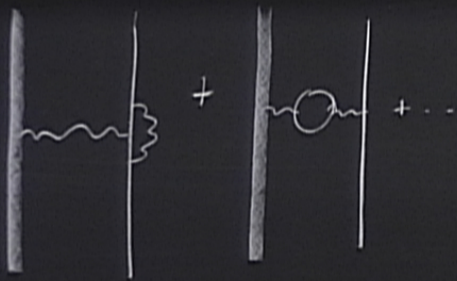
(A)
Feynman
diag
approach

OR



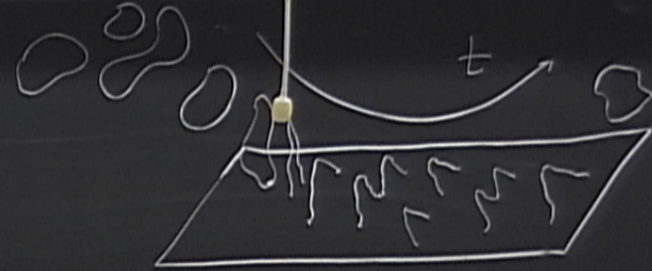
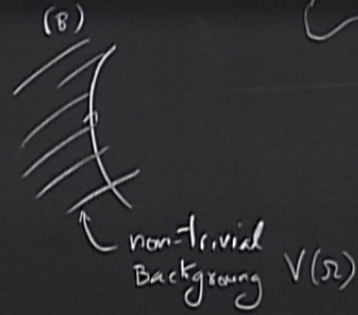
$$V(r) = -\frac{\alpha}{r} \left(1 + \frac{\alpha}{4\sqrt{\pi}} \frac{e^{-2m_e r}}{(m_e r)^{3/2}} + \dots \right)$$

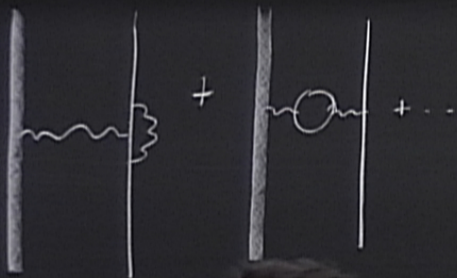
Coulomb



(A)
Feynman
diag
approach

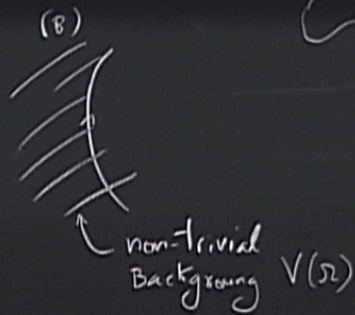
OR





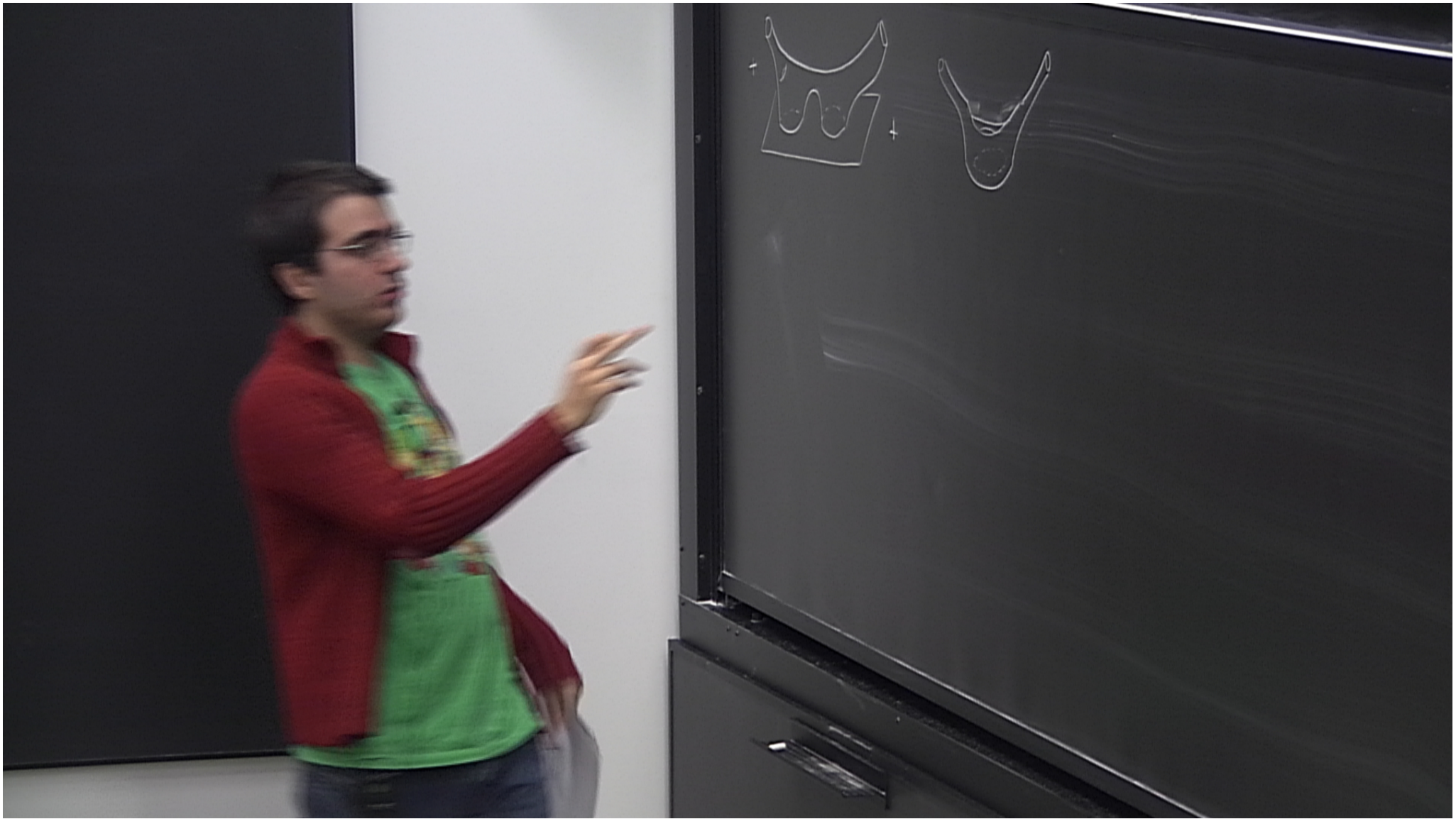
(A)
Feynman
diag
approach

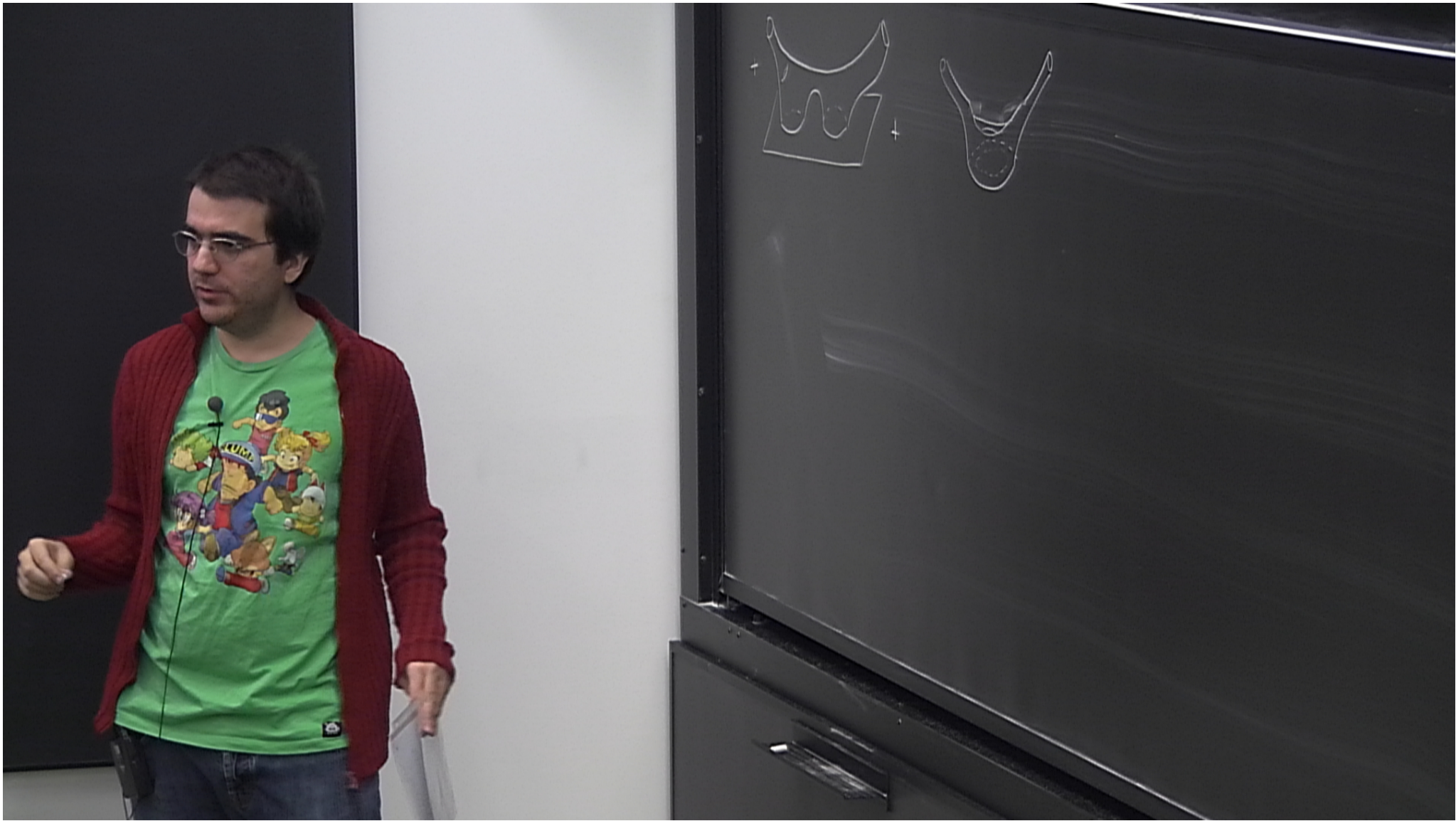
OR

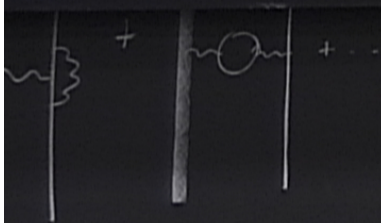


or



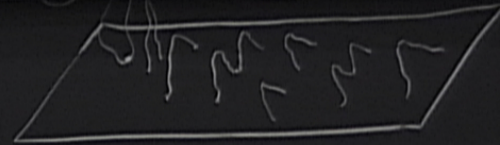




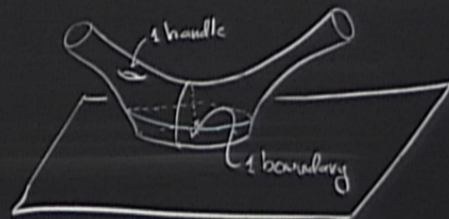


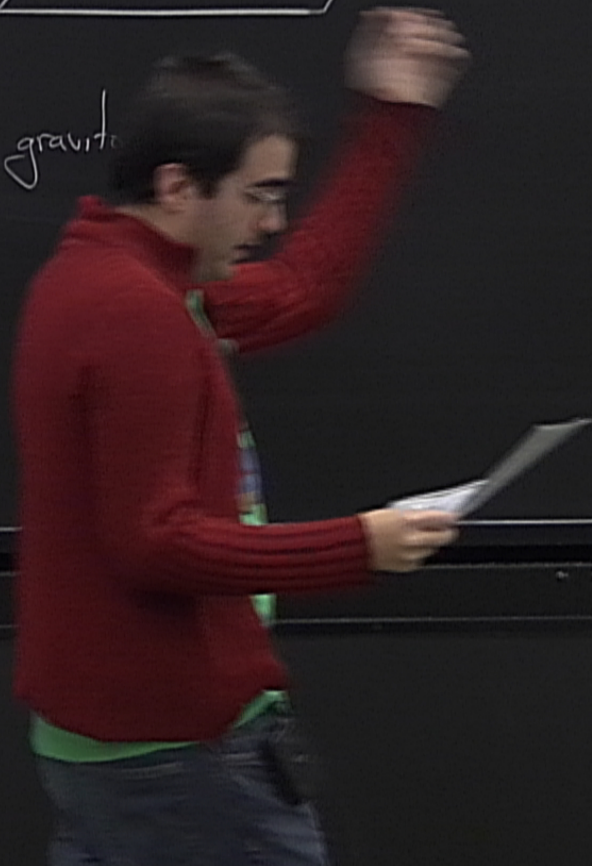
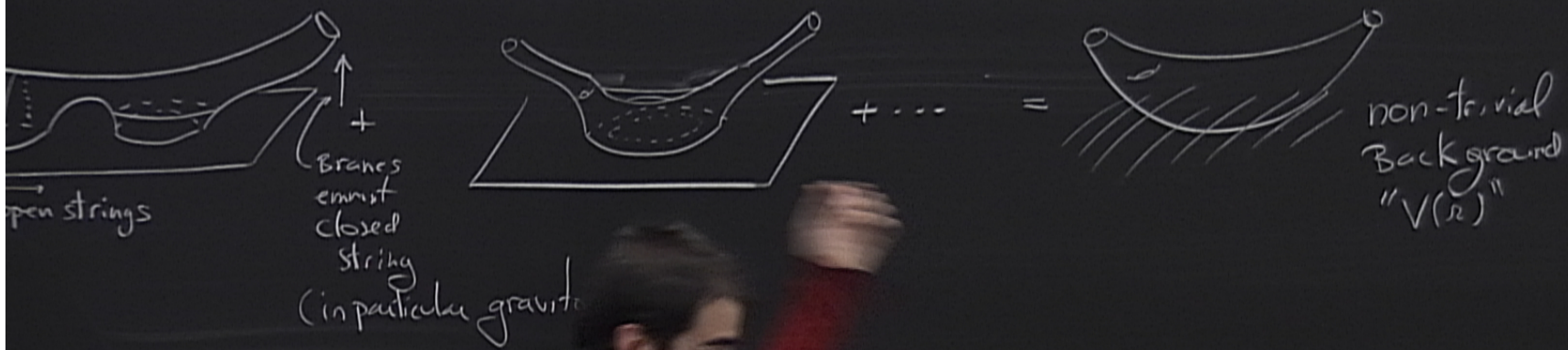
Feynman
diag
approach OR

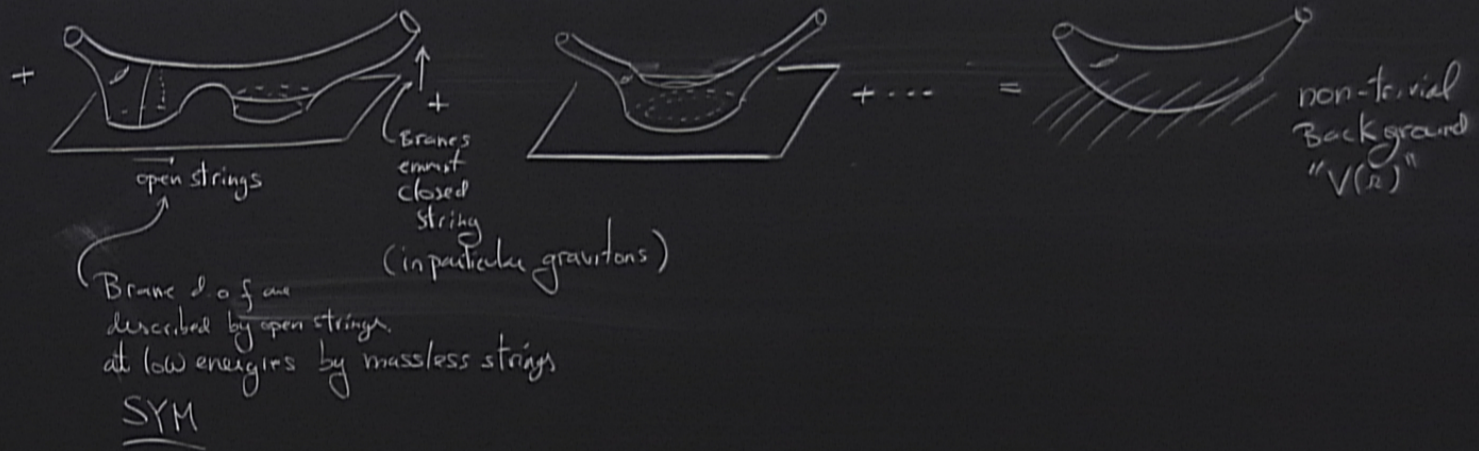
non-trivial
Background $V(\Omega)$

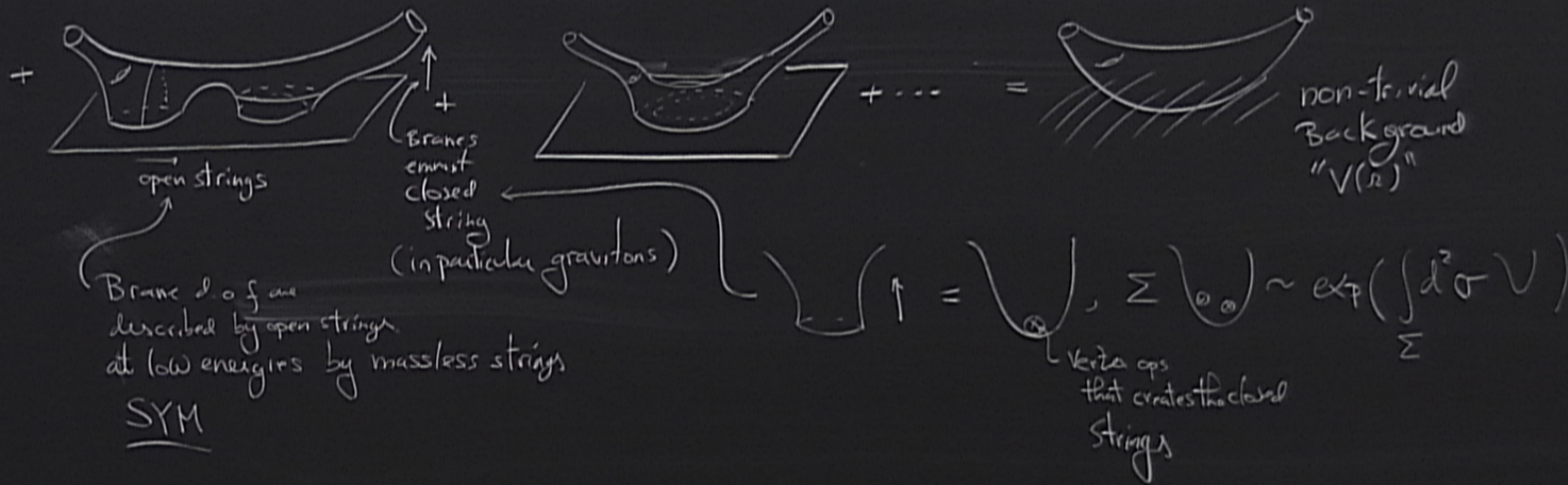


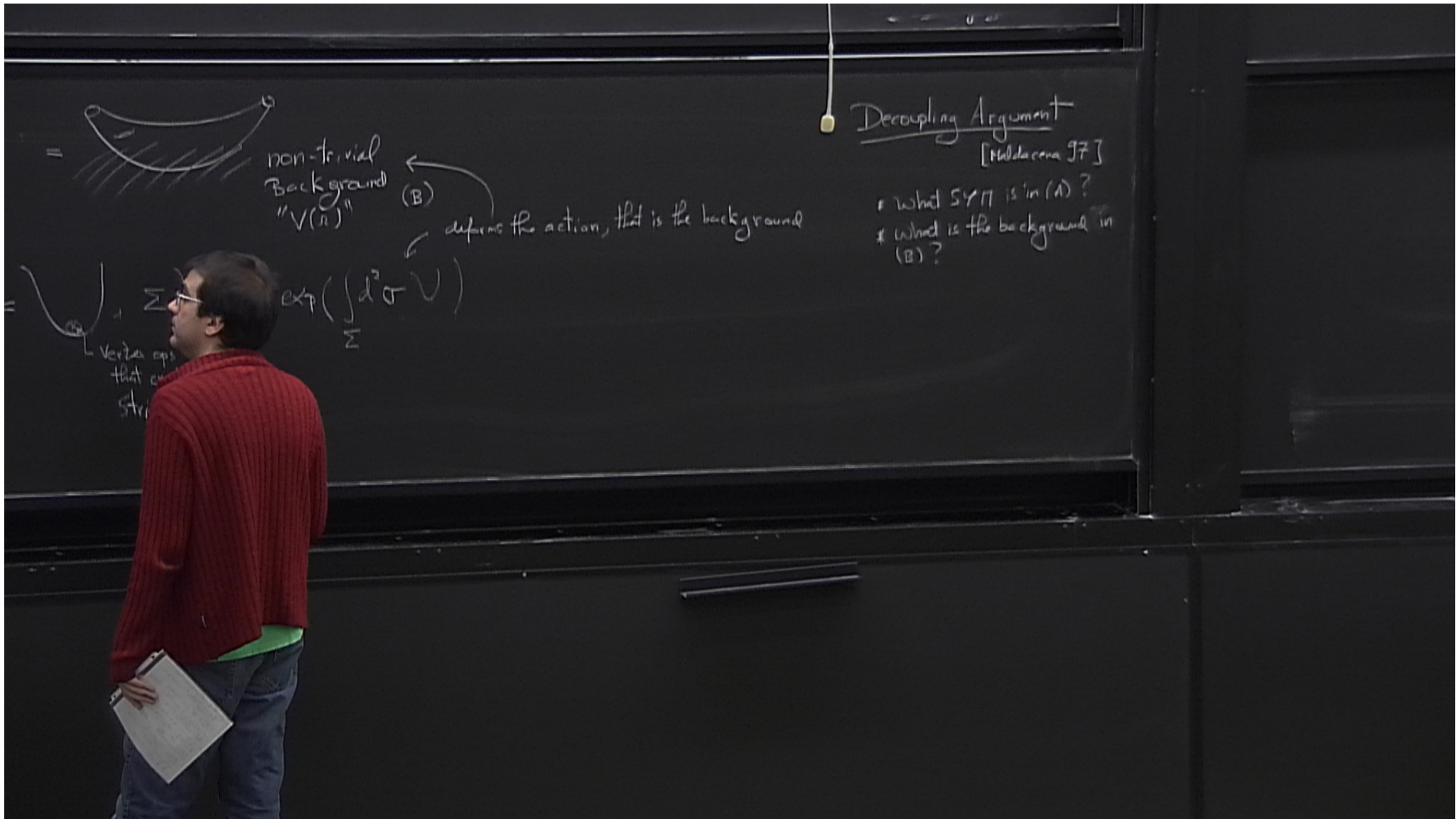
or











action, that is the background

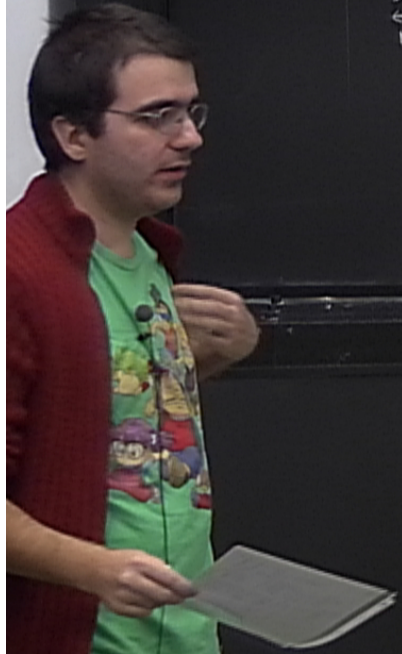
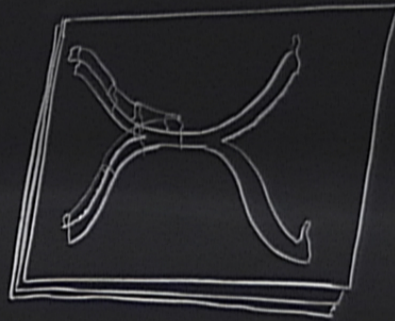
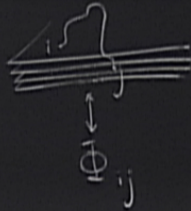
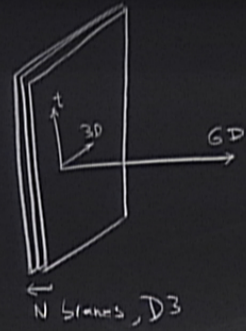
Decoupling Argument

[Maldacena 97]

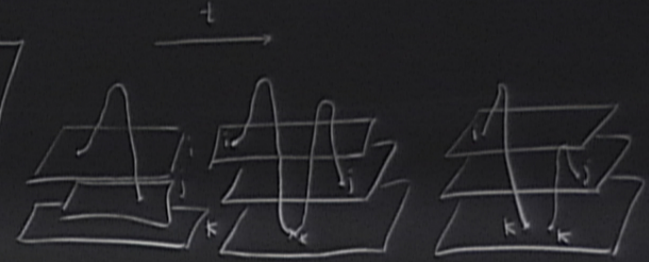
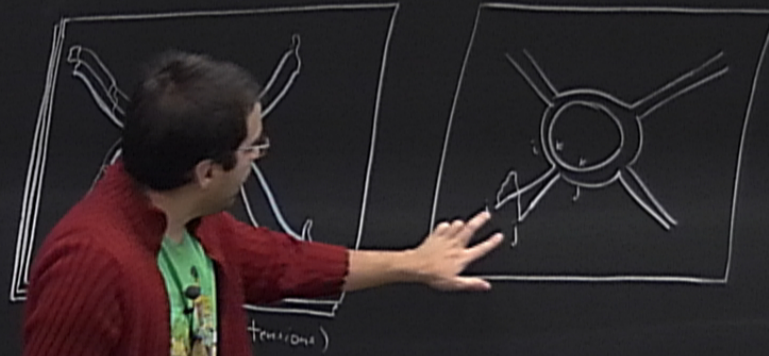
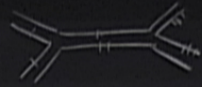
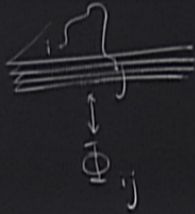
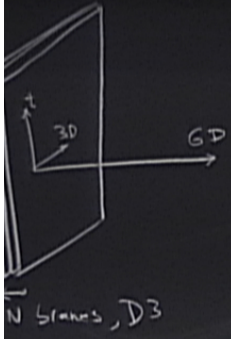
* What SYM is in (A)?

* What is the background in (B)?

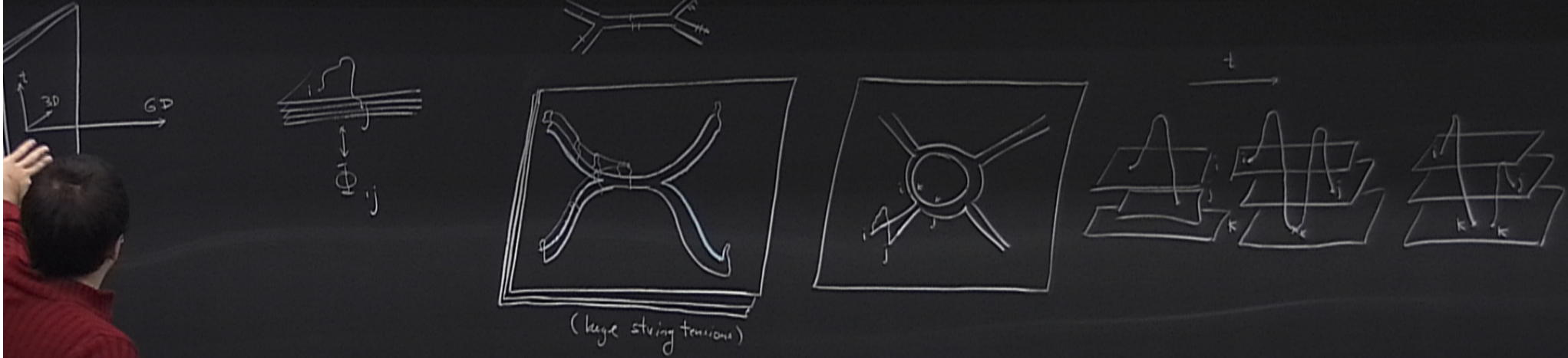
Strings



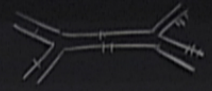
Strings



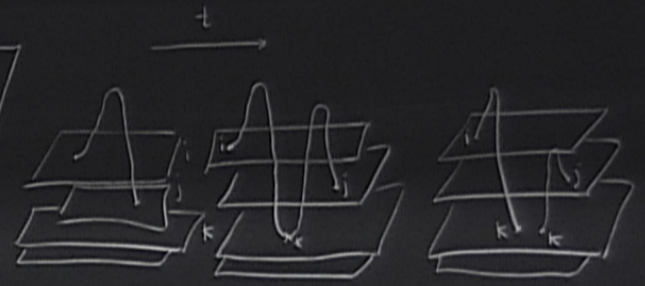
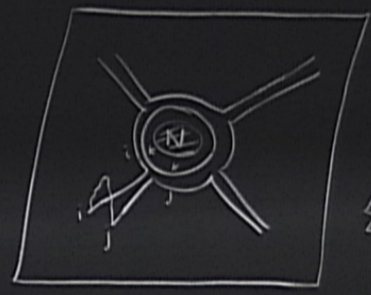
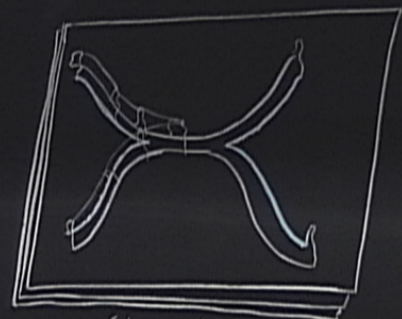
Strings



Strings



CHAN PATTON MECHANISM

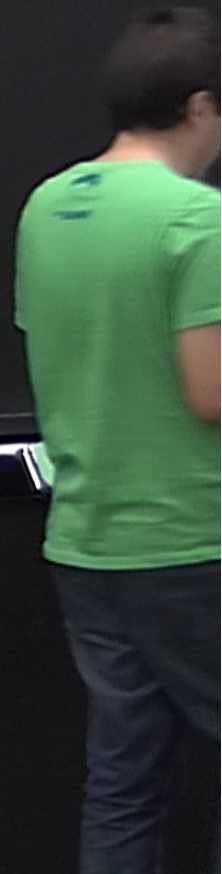


is
part of from
series (Poles)
in total,

$$\lambda = g N$$

$$S = S_{\text{brane}} + S_{\text{int}} + S_{\text{bulk}}$$

↑
SYM + higher derivatives



$$S = S_{\text{brane}} + S_{\text{int}} + S_{\text{bulk}}$$

↑ SYM + higher derivatives
 ↑ IOD SUGRA + higher derivatives
 ↑ $\frac{1}{G_N} \int \sqrt{g} R$
 ↑ $d = 2 + \sqrt{G_N}$

from massive modes

$$S = S_{\text{brane}} + S_{\text{int}} + S_{\text{bulk}}$$

\uparrow SYM + higher derivatives

\uparrow 10D SUGRA + higher derivatives

\uparrow from massive modes

$$\frac{1}{G_N} \int \sqrt{g} R \sim \int (\partial h)^2 + \sqrt{G_N} h (\partial h)^2$$

$d = 2 + \sqrt{G_N} h$

$$S = S_{\text{brane}} + S_{\text{int}} + S_{\text{bulk}}$$

\uparrow SYM + higher derivatives

\uparrow 10D SUGRA + higher derivatives

\uparrow from massive modes

$$\frac{1}{G_N} \int \sqrt{g} R \sim \int (\partial h)^2 + \sqrt{G_N} h (\partial h)^2 + \dots$$

$$g = \eta + \sqrt{G_N} h \quad \sqrt{G_N} \sim g_s \alpha'^2$$

$$S = S_{\text{brane}} + S_{\text{int}} + S_{\text{bulk}}$$

↑
 SYM + higher derivatives
 (coupling λ)
 ↓
 dimensionless

↑
 10D SUGRA + higher derivatives
 ↓
 dimensional (with α')
 couplings

from massive modes

$$\frac{1}{G_N} \int \sqrt{g} R \sim \int (\partial h)^2 + \sqrt{G_N} h (\partial h)^2 + \dots$$

$g = \eta + \sqrt{G_N} h$

$\sqrt{G_N} \sim g_s \alpha'^2$

$$S = S_{\text{brane}} + S_{\text{int}} + S_{\text{bulk}}$$

in fld space

from massive modes

SYM + higher derivatives
(coupling λ)

dimensionless

10D SUGRA + higher derivatives

$$\frac{1}{G_N} \int \sqrt{g} R \sim \int (\partial h)^2 + \sqrt{G_N} h (\partial h)^2 + \dots$$

$$g = \eta + \sqrt{G_N} h$$

$$\sqrt{G_N} \sim g_s \alpha'^2$$

dimensionful couplings (with α')

lets consider a low energy limit

only survives

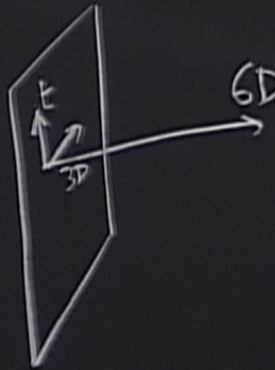
$$\alpha' k^2 \ll 1, \quad |\vec{x} - \vec{x}'| / \alpha' \gg 1$$

in practise this amounts to taking $\alpha' \rightarrow 0$

(A) = 4D SYM ($\mathcal{N}=4$) \oplus 10D SUGRA

lets move to point of view (B)

"Gen. Einstein eqs" \leftarrow Eff. action



$$g_{\mu\nu} = ?$$

$$ds^2 = \sum_{\mu\nu} dx^\mu dx^\nu$$

$$dx^m dx_m$$

lets move to point of view (B)

"Gen. Einstein eqs" \leftarrow Eff. action



$g_s \propto r^2$

$g_{\mu\nu} = ?$

$$ds^2 = \sum_{\mu, \nu} \frac{dx^\mu dx^\nu}{\tilde{f}(r)}$$

$0 \dots 3$

$$+ f(r) dx^m dx_m$$

$4 \dots 9$

$$r^2 = X^m X_m$$

one gets $\tilde{f}(r) = f(r) = \sqrt{1 + \frac{L^4}{r^4}}$

lets move to point of view (B)

"Gen. Einstein eqs" \leftarrow Eff. action



$$g_{\mu\nu} = ?$$

$$ds^2 = \sum_{\mu\nu} \frac{dx^\mu dx^\nu}{\tilde{f}(\mathcal{R})} + f(\mathcal{R}) dx^m dx_m$$

$\begin{matrix} \text{0...3} & & \text{4...9} \\ \curvearrowright & & \curvearrowright \end{matrix}$

one gets $\tilde{f}(\mathcal{R}) = f(\mathcal{R}) = \sqrt{1 + \frac{L^4}{\mathcal{R}^4}}$

$$\mathcal{R}^2 = X^m X_m$$

$$L^4 = g \int_S N 4\pi \alpha'$$

\uparrow N units of flux sourced by the N branes

good, useful, simple, for $g_s N = \lambda \ll 1$

$$(A) = 4D \text{ SYM } (\omega^2 = 4) \oplus 10D \text{ SUGRA}$$

in practise this amounts to taking $\alpha' \rightarrow 0$

$\pi \gg L$
(far away from the Branes)

$\pi \rightarrow 0$
(close to the branes)

$$ds^2 \rightarrow dx_\mu dx_\nu \eta^{\mu\nu} + dx_m dx^m \in M^{10}$$

in practise this amounts to taking $\alpha' \rightarrow 0$

$\pi \gg L$
(far away from the Branes)

$$ds^2 \rightarrow dx_\mu dx_\nu \eta^{\mu\nu} + dx_m dx^m \in M^{10}$$

$\pi \rightarrow 0$
(close to the branes)

$$ds^2 \rightarrow \frac{\pi^2}{L^2} \eta_{\mu\nu} dx^\mu dx^\nu + \frac{L^2}{\pi^2} dt^2 + L^2 d\Omega_5^2$$

4D, \mathcal{N}

$r \gg L$
(far away from the Branes)

$$ds^2 \rightarrow dx_\mu dx_\nu \eta^{\mu\nu} + dx_m dx^m \in M^{10}$$

$r \rightarrow 0$
(close to the branes)

$$ds^2 \rightarrow \underbrace{\frac{r^2}{L^2} \eta_{\mu\nu} dx^\mu dx^\nu}_{\text{metric of AdS}_5} + \underbrace{\frac{L^2}{r^2} dr^2 + L^2 d\Omega_{S^5}^2}_{\text{5 sphere, } S^5}$$

from $dx_\mu dx_\nu \eta^{\mu\nu}$
from the $dx_m dx^m$ part

$$dx_{\mu} dx^{\mu} = -dt^2 + d\vec{x}^2$$

time measured by obs.
at infinity.

good, useful, simple, for $g_s N = \lambda \ll 1$

(A) = 4D SYM ($d=4$) \oplus 10D SUGRA

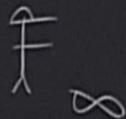
$$dx_\mu dx^\mu = -dt^2 + d\vec{x}^2$$

↑
time measured by obs.
at infinity.



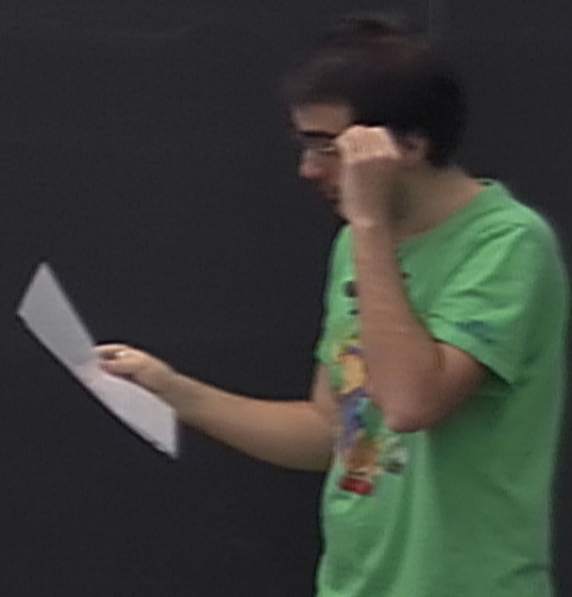
$$\Delta t_{\text{PROPER}} = \frac{\sqrt{2}}{L} \Delta t_\infty$$

$$\Delta E_{\text{PROPER}} = \frac{L}{\sqrt{2}} \Delta E_\infty$$



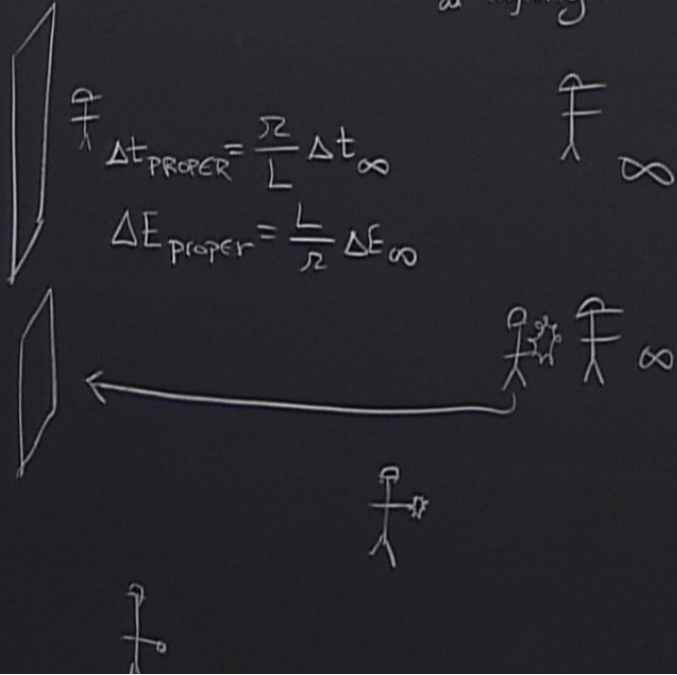
good, useful, simple, for $g_s N = \lambda \ll 1$

(A) = 4DSYM ($\mathcal{N}=4$) \oplus 10D SUGRA



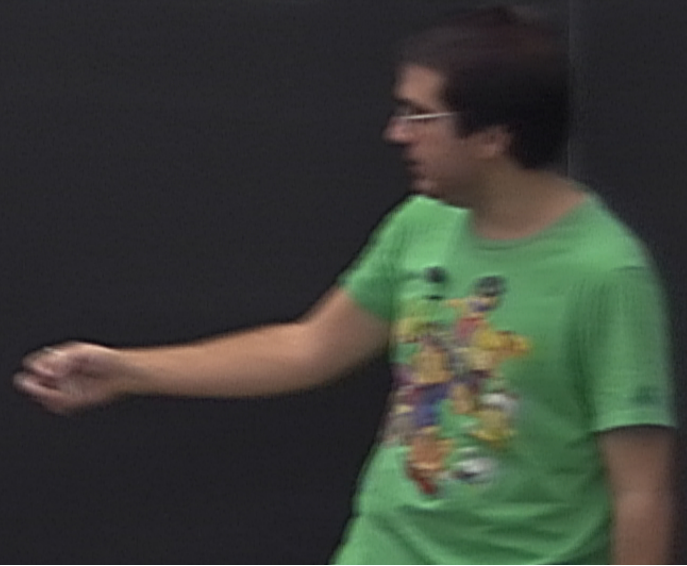
$$dx_\mu dx^\mu = -dt^2 + d\vec{x}^2$$

time measured by obs.
at infinity.



good, useful, simple, for $g_{5N} = \lambda \ll 1$

(A) = 4DSYM ($\mathcal{N}=4$) \oplus 10D SUGRA



$$dx_{\mu} dx^{\mu} = -dt^2 + d\vec{x}^2$$

↑
time measured by obs.
at infinity.

good, useful, simple, for $g_s N = \lambda \ll 1$

(A) = 4DSYM ($d=4$) \oplus ~~10D SUGRA~~

