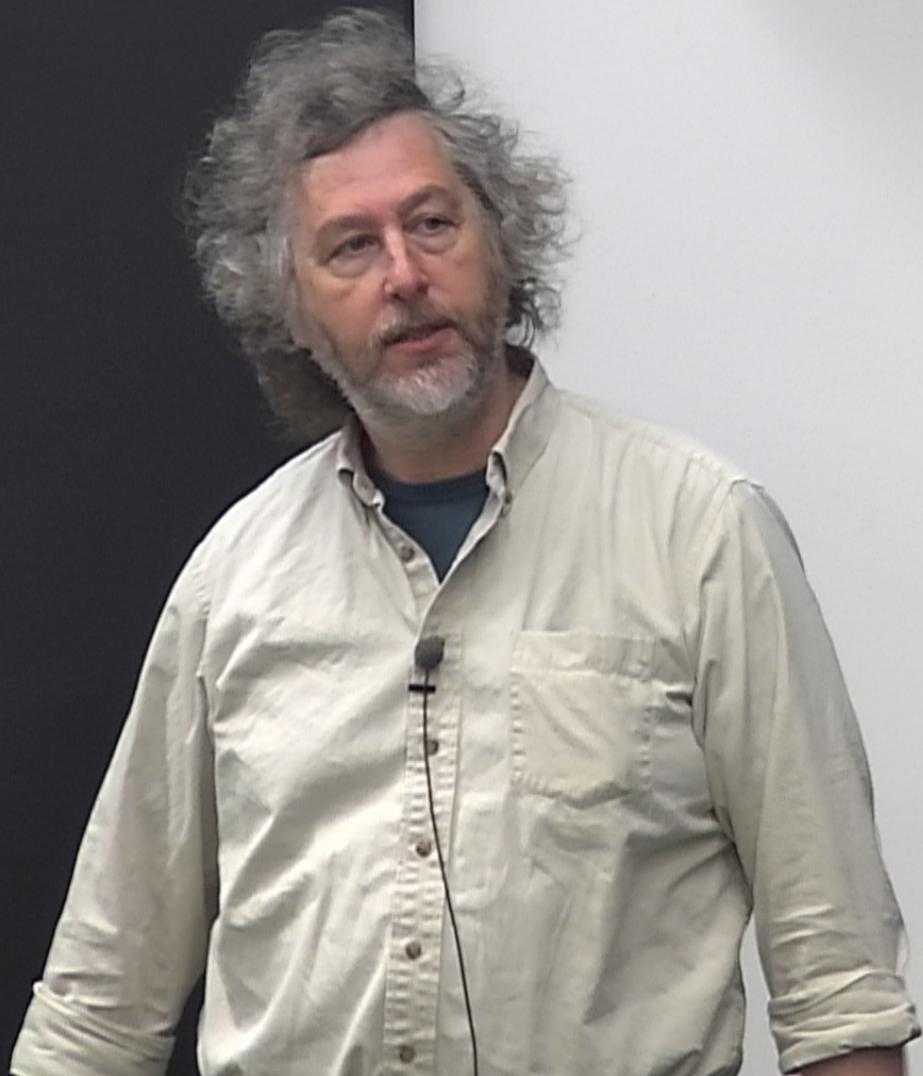


Title: Explorations in Quantum Information - Lecture 13

Date: Mar 29, 2012 09:00 AM

URL: <http://pirsa.org/12030020>

Abstract:



$$\begin{aligned} m \ddot{\vec{r}}_1 &= \nabla V : \vec{Y}(\vec{r}_1, t) \\ \frac{d}{dt} \left\{ \vec{Y}(t) \right\} &= \nabla V \left\{ \vec{Y}(t) \right\} \\ \dot{\vec{Y}} &= \sum_{i=1}^3 \frac{d^2}{dx_i^2} + U(x) \quad \text{Vector} \\ \frac{\partial^2 \vec{Y}}{\partial x_i^2} &= -\frac{1}{a^2} \left[Y(x_{i+1}) - 2Y(x_i) + Y(x_{i-1}) \right] \end{aligned}$$

Finite difference Method

$$\text{it } \frac{\partial}{\partial t} \psi = \mathcal{H} \psi \quad ; \quad \psi(\vec{r}, t)$$

$$\text{it } \frac{d}{dt} \left\{ \psi(t) \right\} = H_m \left\{ \psi(t) \right\}$$

$\left\{ \psi(x_0), \psi(x_1), \dots \right\}$
 $x_n = na ; n = \text{integer}$

vector

$$f'(x) \simeq \frac{f(x+h) - f(x)}{h}$$

$$f''(x) \simeq \frac{f(x+h) - 2f(x) + f(x-h)}{h^2}$$

Finite difference Method

$$\text{ith } \frac{\partial}{\partial t} \psi = \mathcal{H} \psi : \psi(\vec{r}, t)$$

$$\text{ith } \frac{\partial}{\partial t} \{ \psi(t) \} = H_m \{ \psi(t) \}$$

$\{ \psi(x_0), \psi(x_1), \dots \}$

$$\mathcal{H} = \frac{-\hbar^2}{2m} \frac{d^2}{dx^2} + U(x)$$

Vector

$$\left(\frac{\partial^2 \psi}{\partial x^2} \right) = \frac{1}{a^2} \left[\psi(x_{n+1}) - 2\psi(x_n) + \psi(x_{n-1}) \right]$$

$x = x_n$

$$\boxed{x_n = na ; n = \text{integer}}$$

$$\boxed{t_0 = \frac{\hbar^2}{2ma^2}}$$

$$\text{ith } \frac{d \psi_i}{dt} =$$

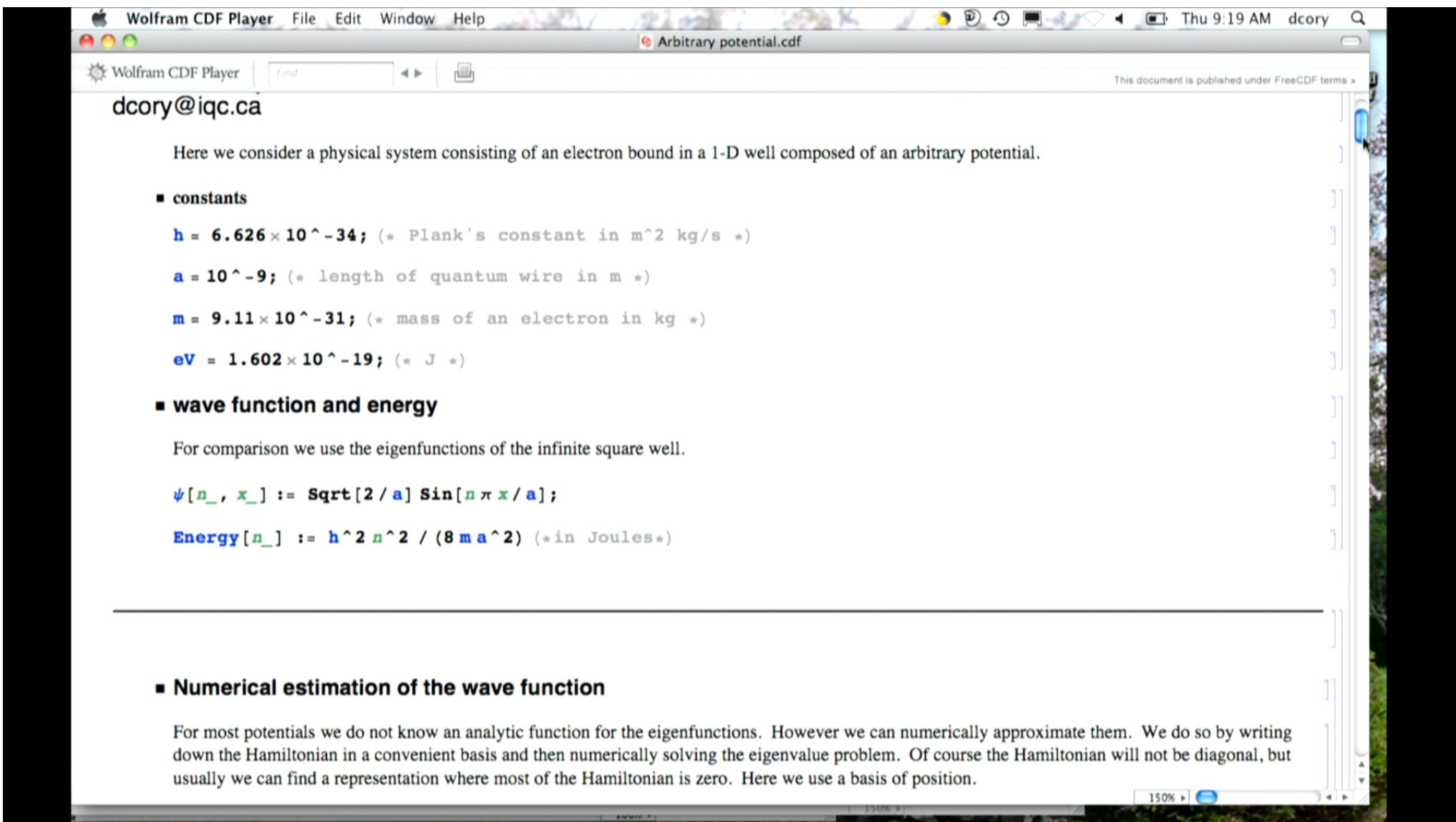
$$\left. \begin{aligned} i\hbar \frac{d\psi_n}{dt} &= (U_n + 2t_0)\psi_n - t_0 \psi_{n-1} - t_0 \psi_{n+1} \\ \psi(x_i) \dots \} &= \sum_m \left[(U_n + 2t_0) \delta_{nm} - t_0 \delta_{n,n-1} - t_0 \delta_{n,n+1} \right] \psi_n \end{aligned} \right|$$

$x_n = na$; $n = \text{integer}$

$$t_0 = \frac{\hbar^2}{2m\alpha^2}$$

$$(x_{n-1})$$

$$\begin{aligned}
 & \left. \frac{dt}{d\tau} \frac{d\psi_n}{dt} = (\mathcal{U}_n + 2t_0) \psi_n - t_0 \psi_{n-1} - t_0 \psi_{n+1} \right\} \\
 & \quad \left. \psi(x_i) \dots \right\} \quad n = \text{integer} \\
 & \quad \left. \begin{array}{l} x_n = na \\ t_0 = \frac{\hbar^2}{2m\omega^2} \end{array} \right\} \\
 & \quad = \sum_m \left[(\mathcal{U}_n + 2t_0) \underbrace{\psi_{nm}}_{\mathcal{H}} - \underbrace{t_0 \psi_{n,n-1}}_{\mathcal{H}} - \underbrace{t_0 \psi_{n,m+1}}_{\mathcal{H}} \right] \psi_n \\
 & f'(x) = \frac{f(x+h) - f(x)}{h} \\
 & f''(x) = \frac{f(x+h) - 2f(x) + f(x-h)}{h^2}
 \end{aligned}$$

Wolfram CDF Player File Edit Window Help **Arbitrary potential.cdf** Thu 9:19 AM dcory 

dcory@iqc.ca

This document is published under FreeCDF terms >

Here we consider a physical system consisting of an electron bound in a 1-D well composed of an arbitrary potential.

■ constants

```
h = 6.626 × 10 ^ -34; (* Plank's constant in m^2 kg/s *)  
a = 10 ^ -9; (* length of quantum wire in m *)  
m = 9.11 × 10 ^ -31; (* mass of an electron in kg *)  
ev = 1.602 × 10 ^ -19; (* J *)
```

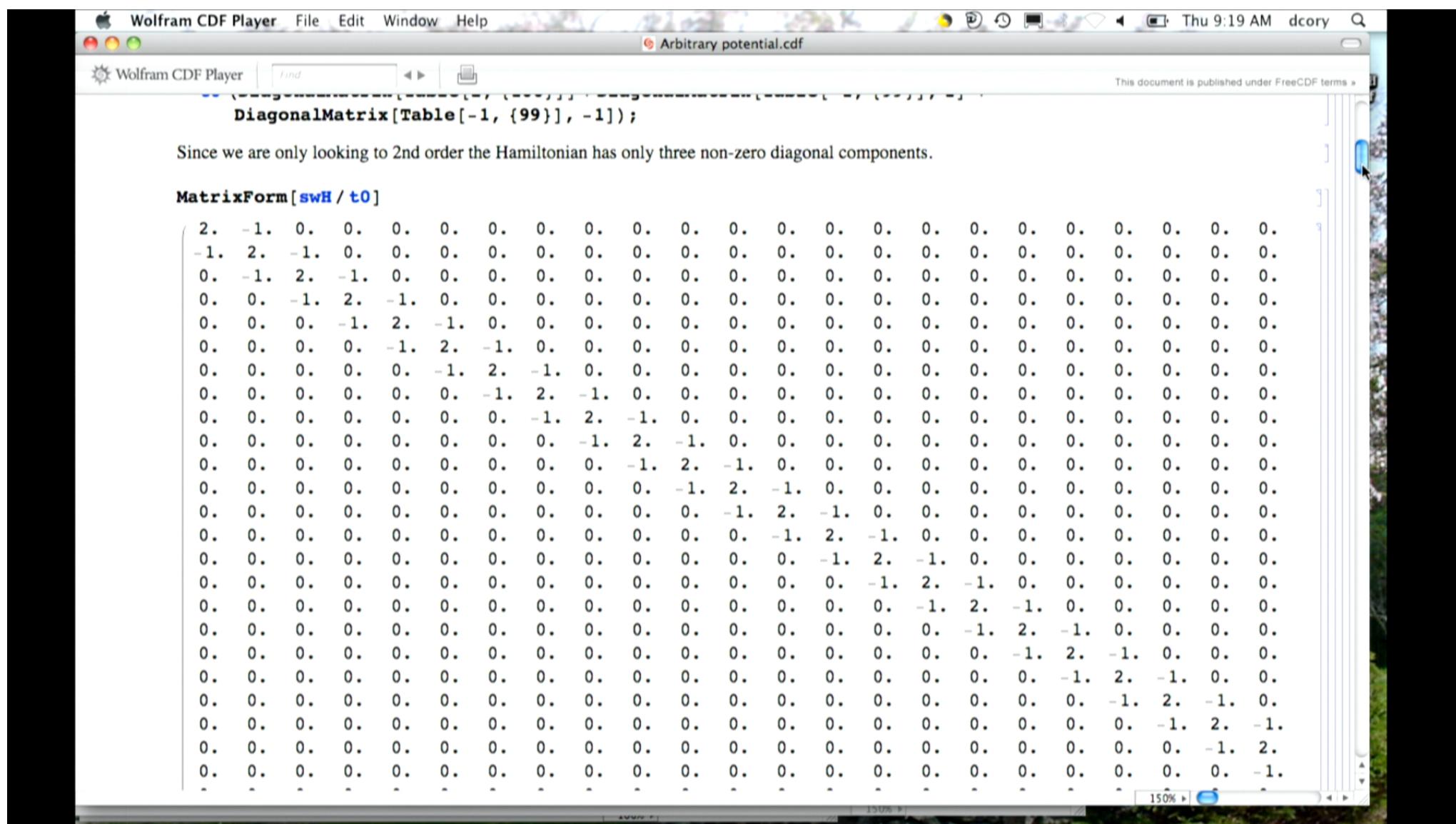
■ wave function and energy

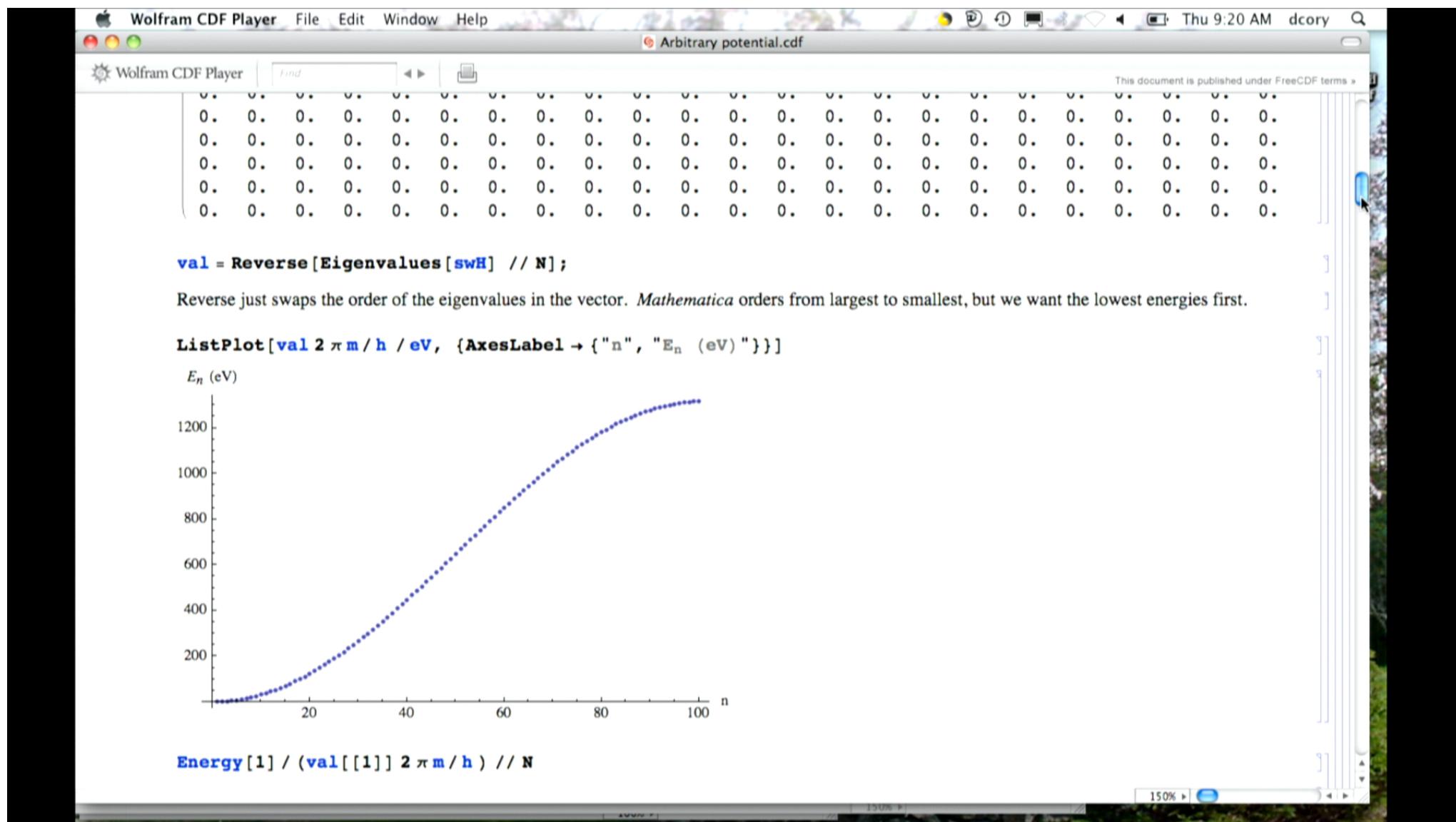
For comparison we use the eigenfunctions of the infinite square well.

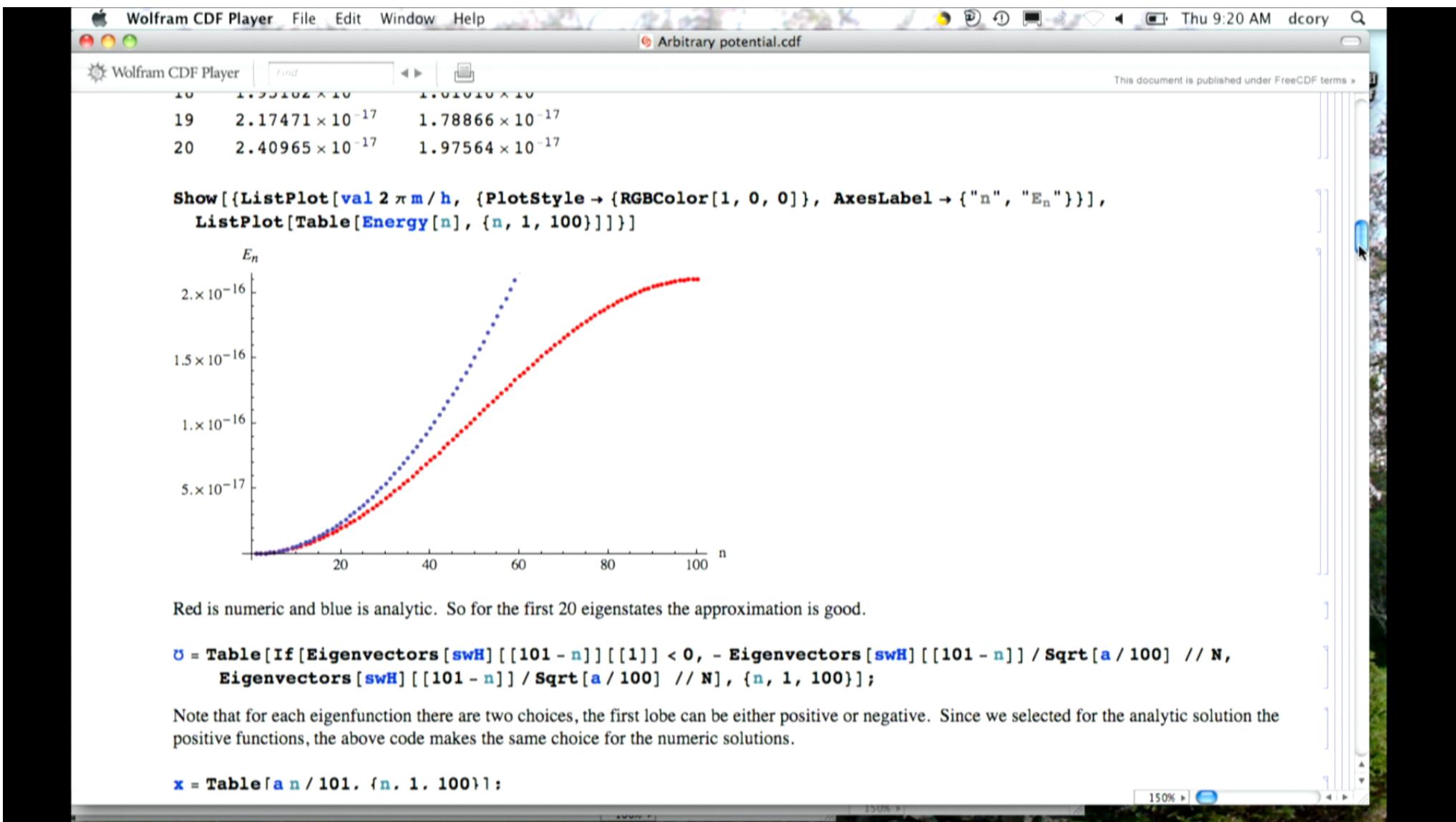
```
ψ[n_, x_] := Sqrt[2 / a] Sin[n π x / a];  
Energy[n_] := h ^ 2 n ^ 2 / (8 m a ^ 2) (*in Joules*)
```

■ Numerical estimation of the wave function

For most potentials we do not know an analytic function for the eigenfunctions. However we can numerically approximate them. We do so by writing down the Hamiltonian in a convenient basis and then numerically solving the eigenvalue problem. Of course the Hamiltonian will not be diagonal, but usually we can find a representation where most of the Hamiltonian is zero. Here we use a basis of position.

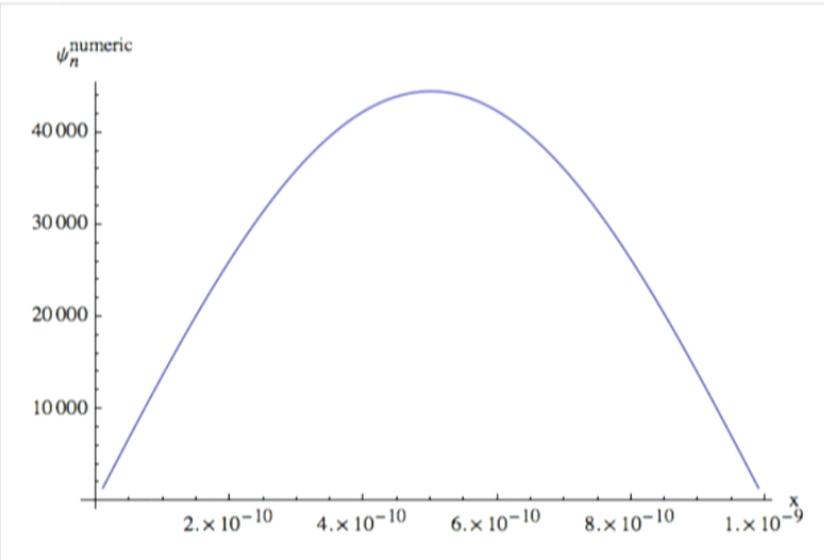






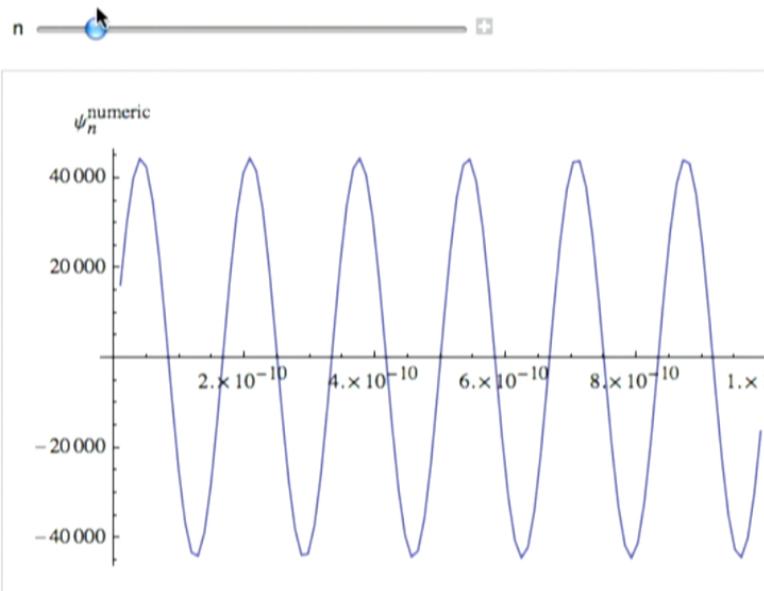
```
Manipulate[ListPlot[Transpose[{x, ψ[[n]]}], {AxesLabel -> {"x", "ψnnumeric"}, PlotJoined -> True}],
{n, 1, 100, 1}, SaveDefinitions -> True]
```

n

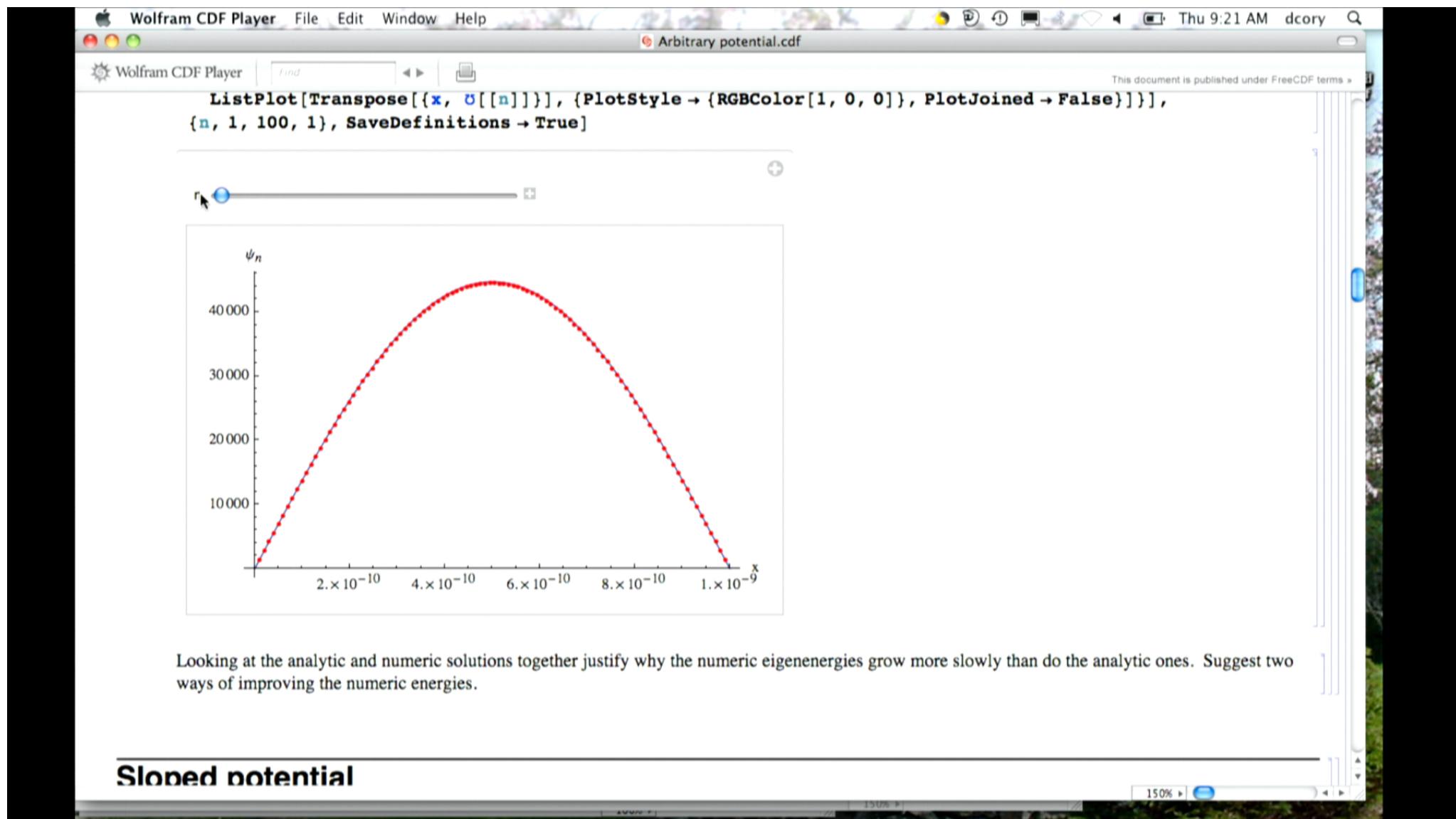


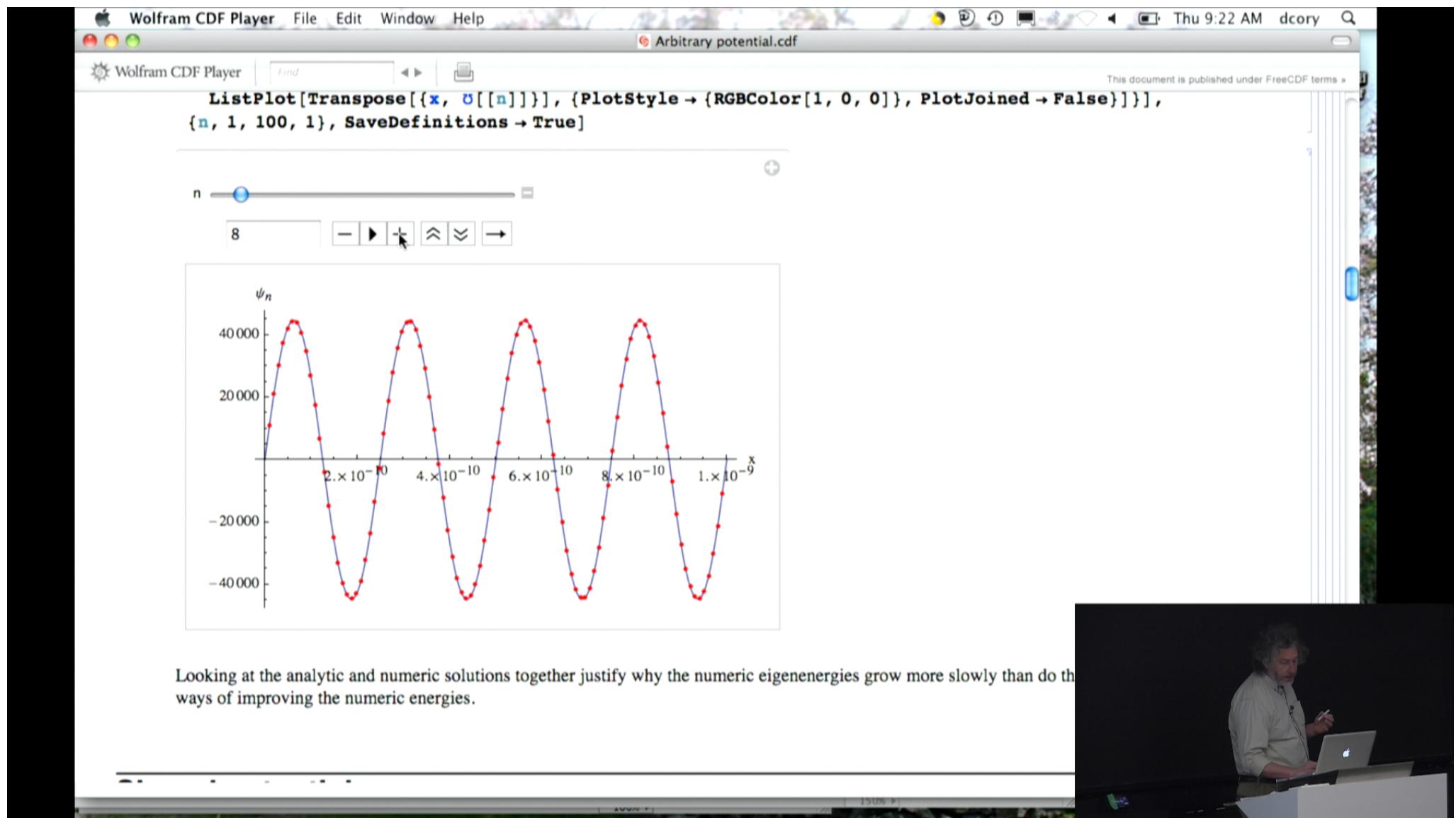
```
Manipulate[
Show[{Plot[ψ[n, x], {x, 0, a}, {AxesLabel -> {"x", "ψn"}, PlotRange -> All}],
ListPlot[Transpose[{x, ψ[[n]]}], {PlotStyle -> {RGBColor[1, 0, 0]}, PlotJoined -> False}]},
{n, 1, 100, 1}, SaveDefinitions -> True]
```

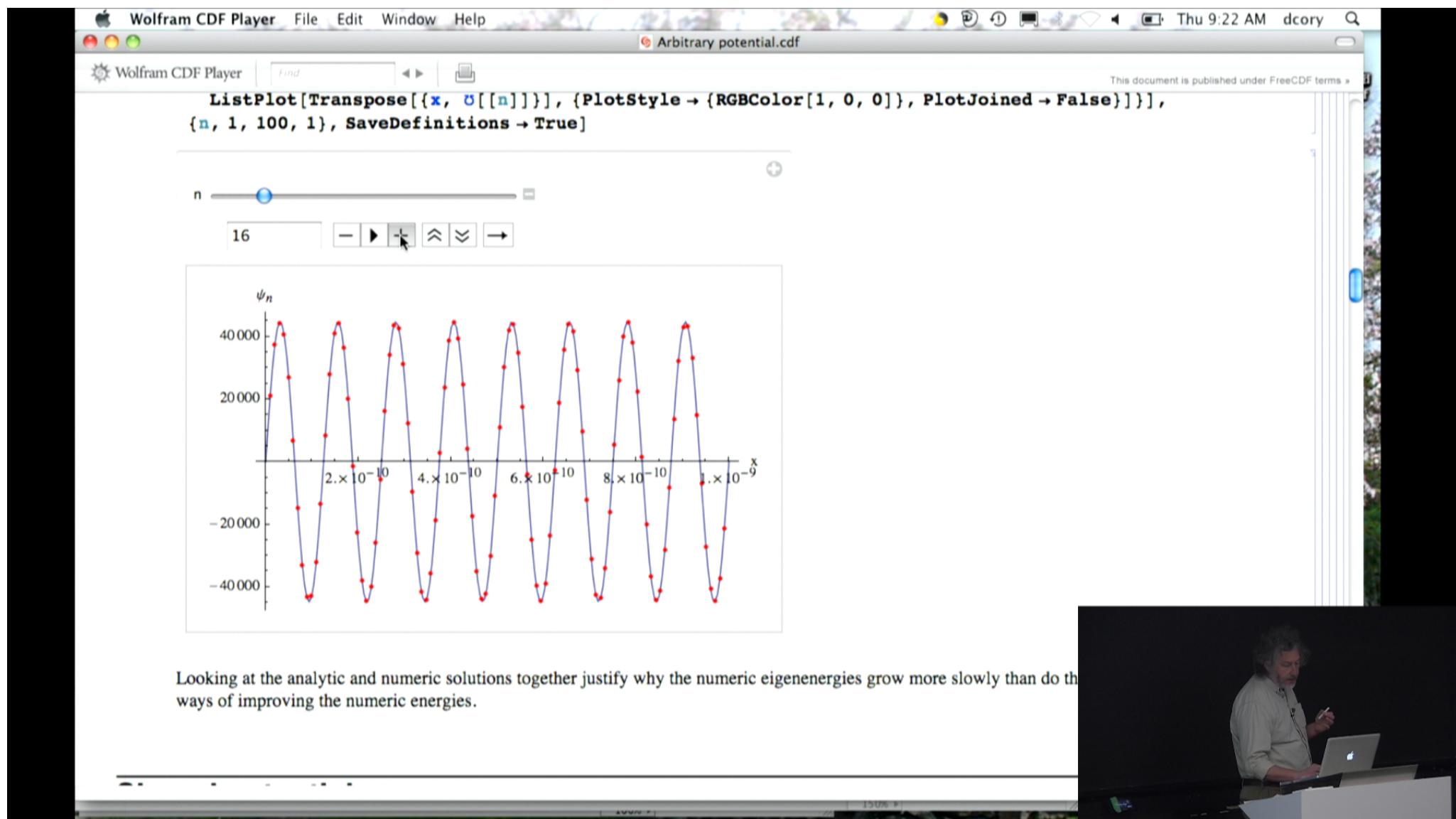
```
Manipulate[ListPlot[Transpose[{x, ψ[[n]]}], {AxesLabel -> {"x", "ψnnumeric"}, PlotJoined -> True}],
{n, 1, 100, 1}, SaveDefinitions -> True]
```

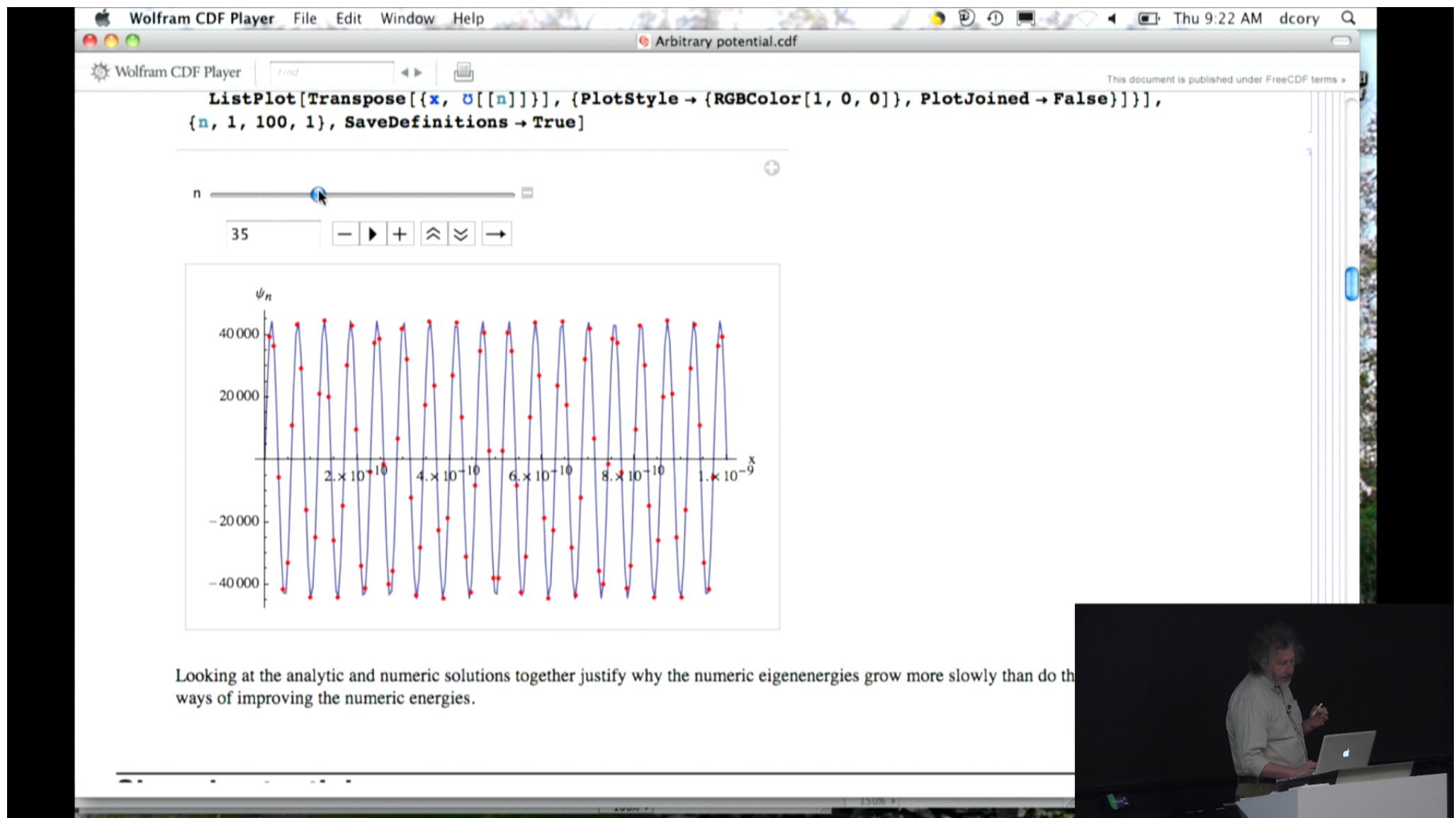


```
Manipulate[
Show[{Plot[ψ[n, x], {x, 0, a}, {AxesLabel -> {"x", "ψn"}, PlotRange -> All}],
ListPlot[Transpose[{x, ψ[[n]]}], {PlotStyle -> {RGBColor[1, 0, 0]}, PlotJoined -> False}]},
{n, 1, 100, 1}, SaveDefinitions -> True]
```



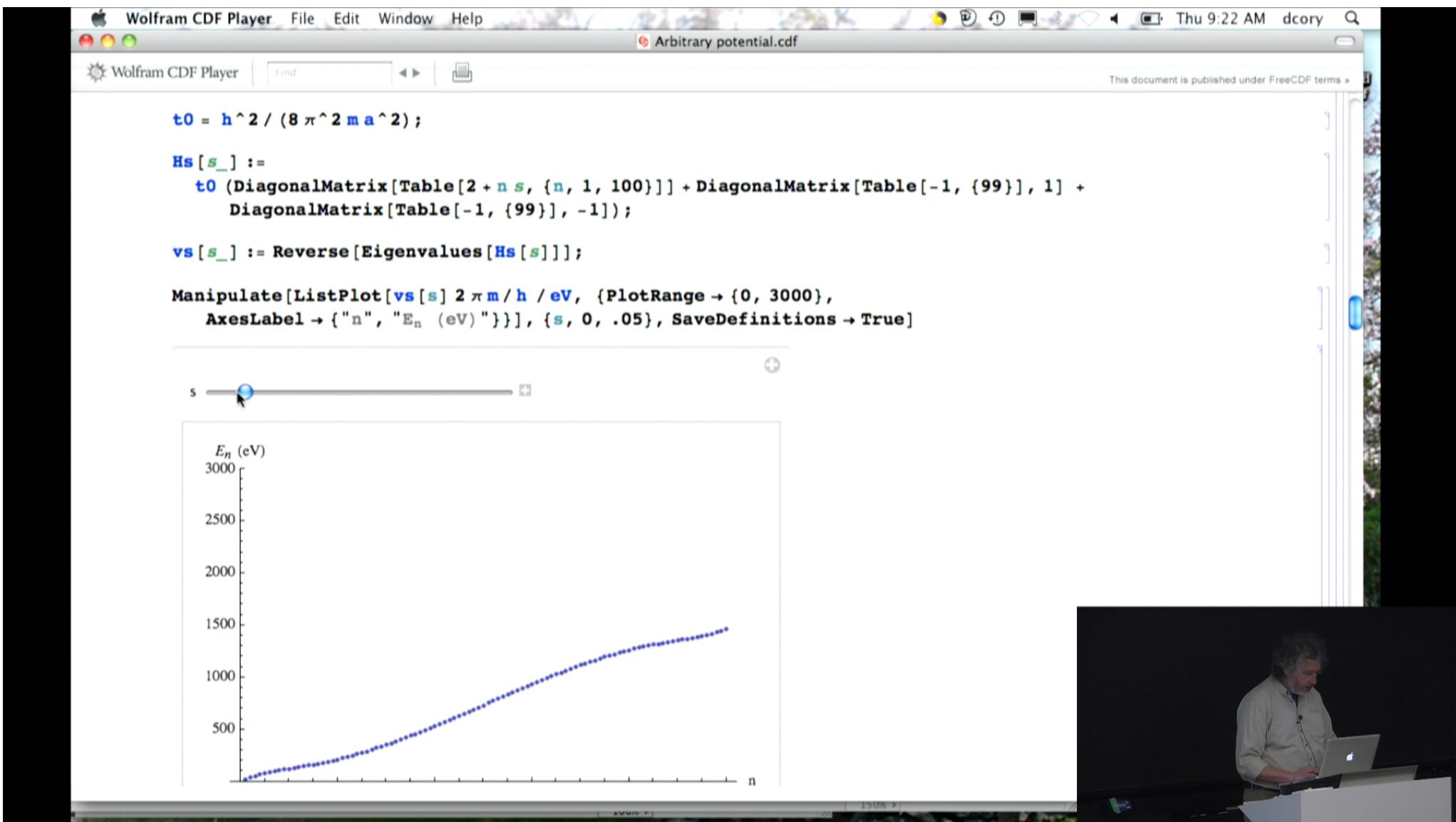


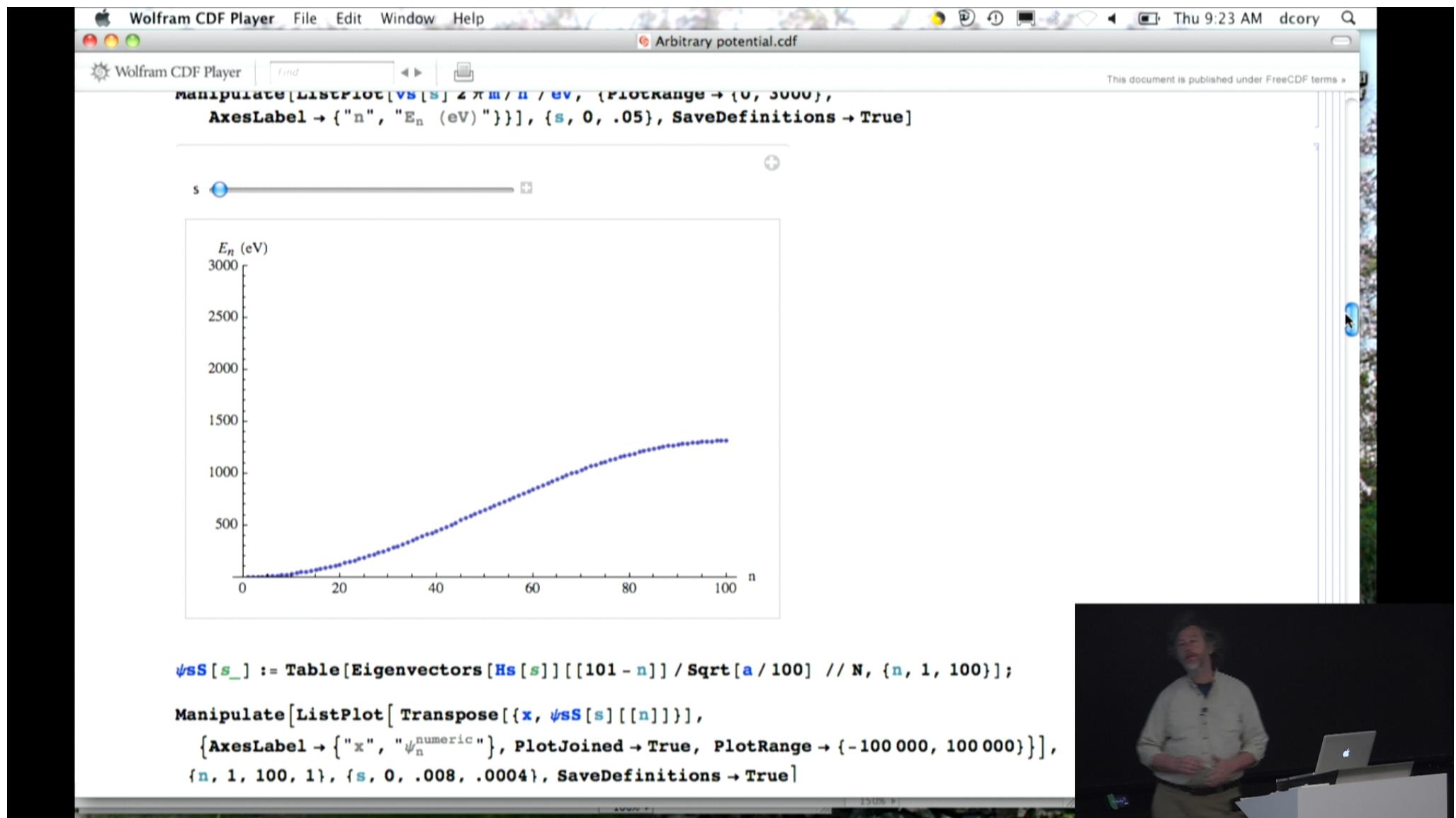


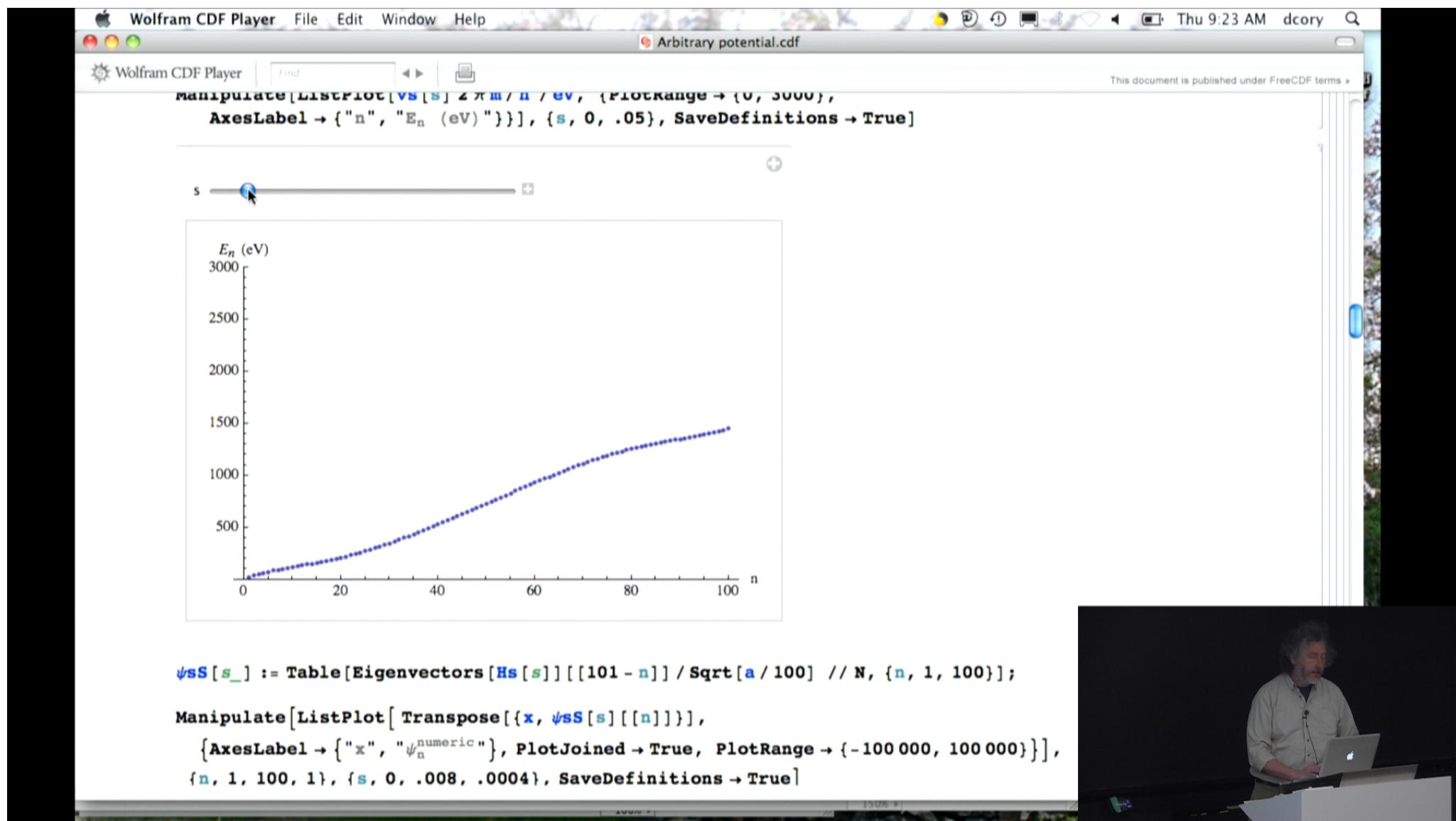


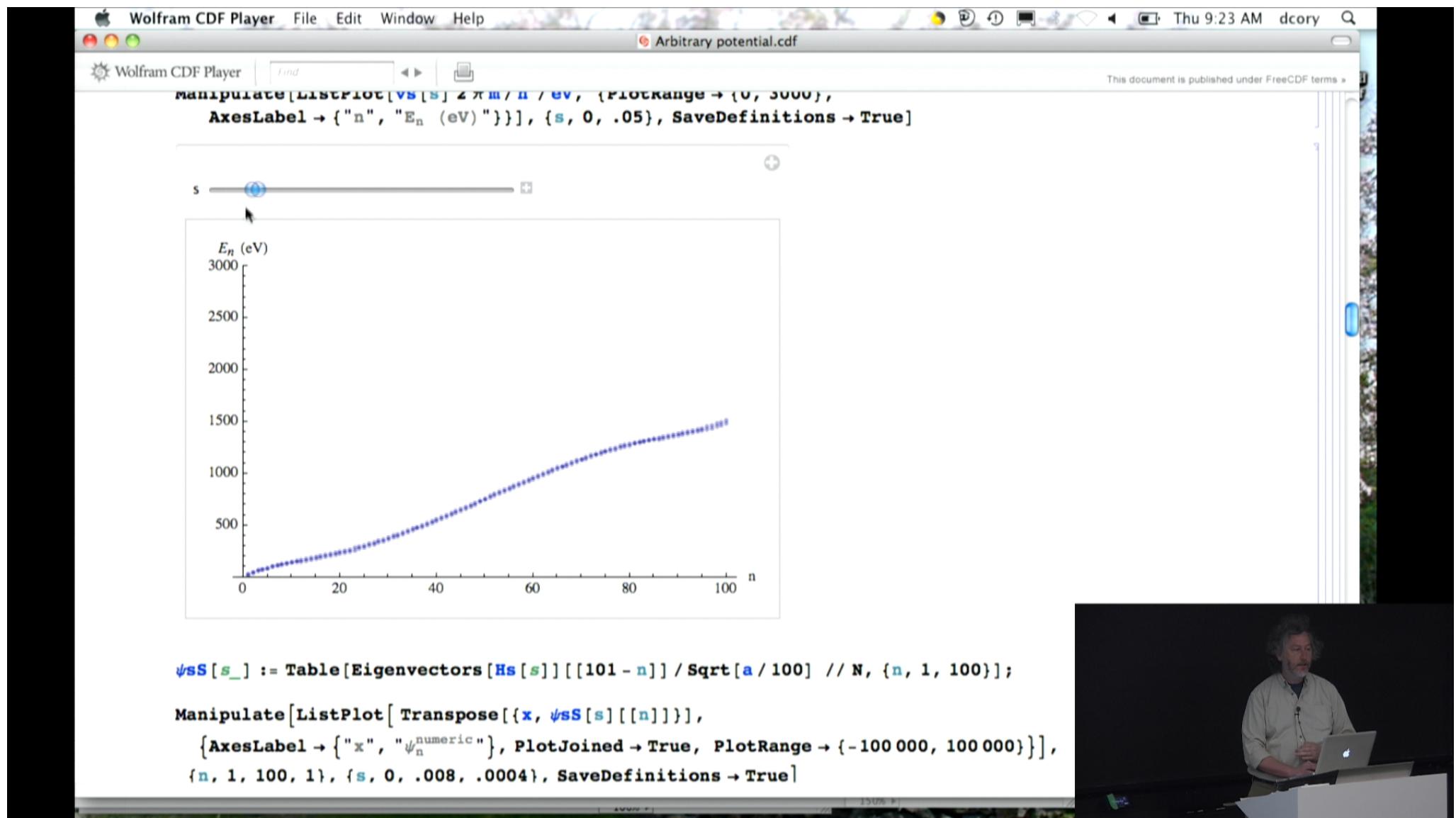
Looking at the analytic and numeric solutions together justify why the numeric eigenenergies grow more slowly than do the ways of improving the numeric energies.

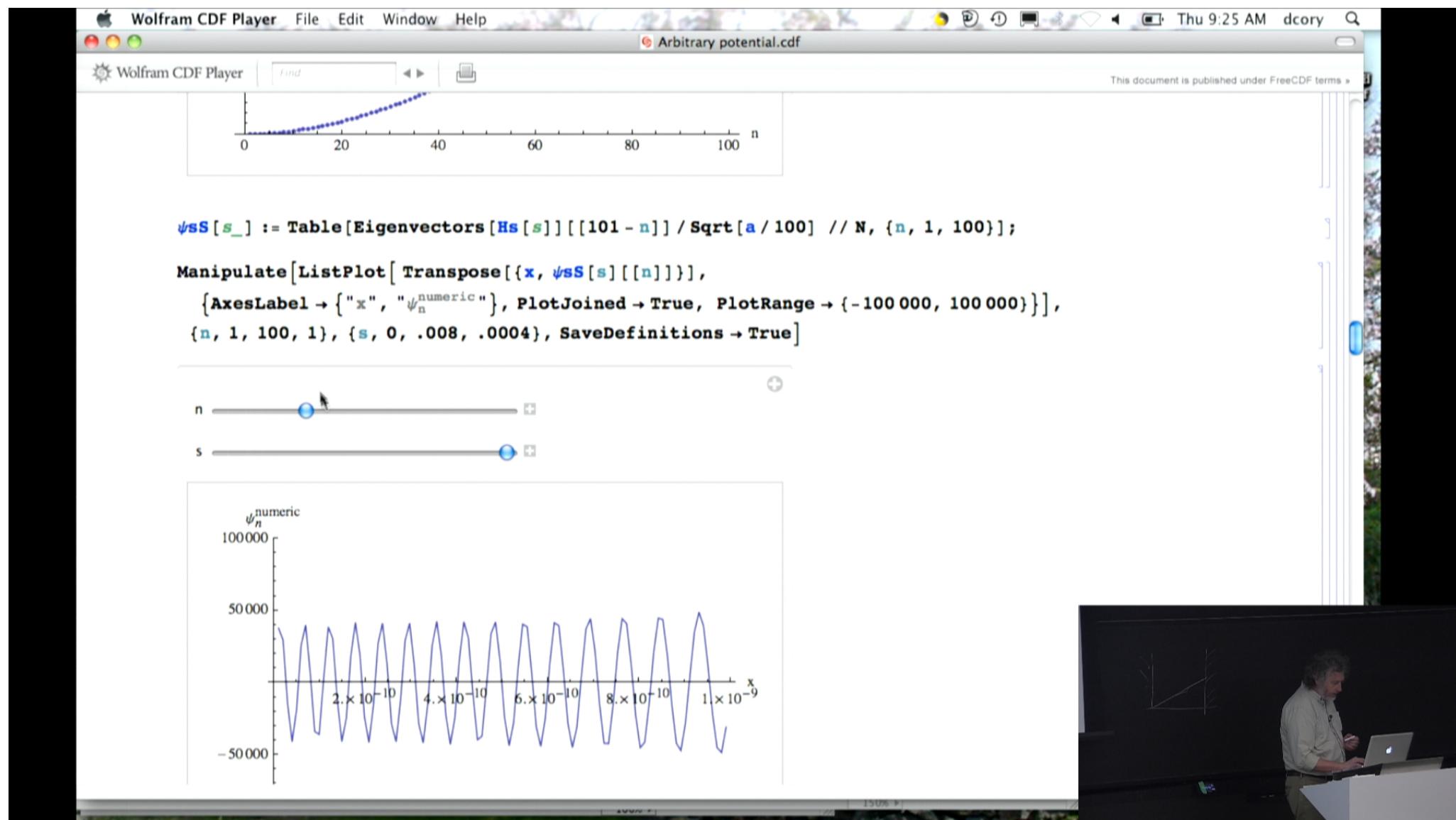


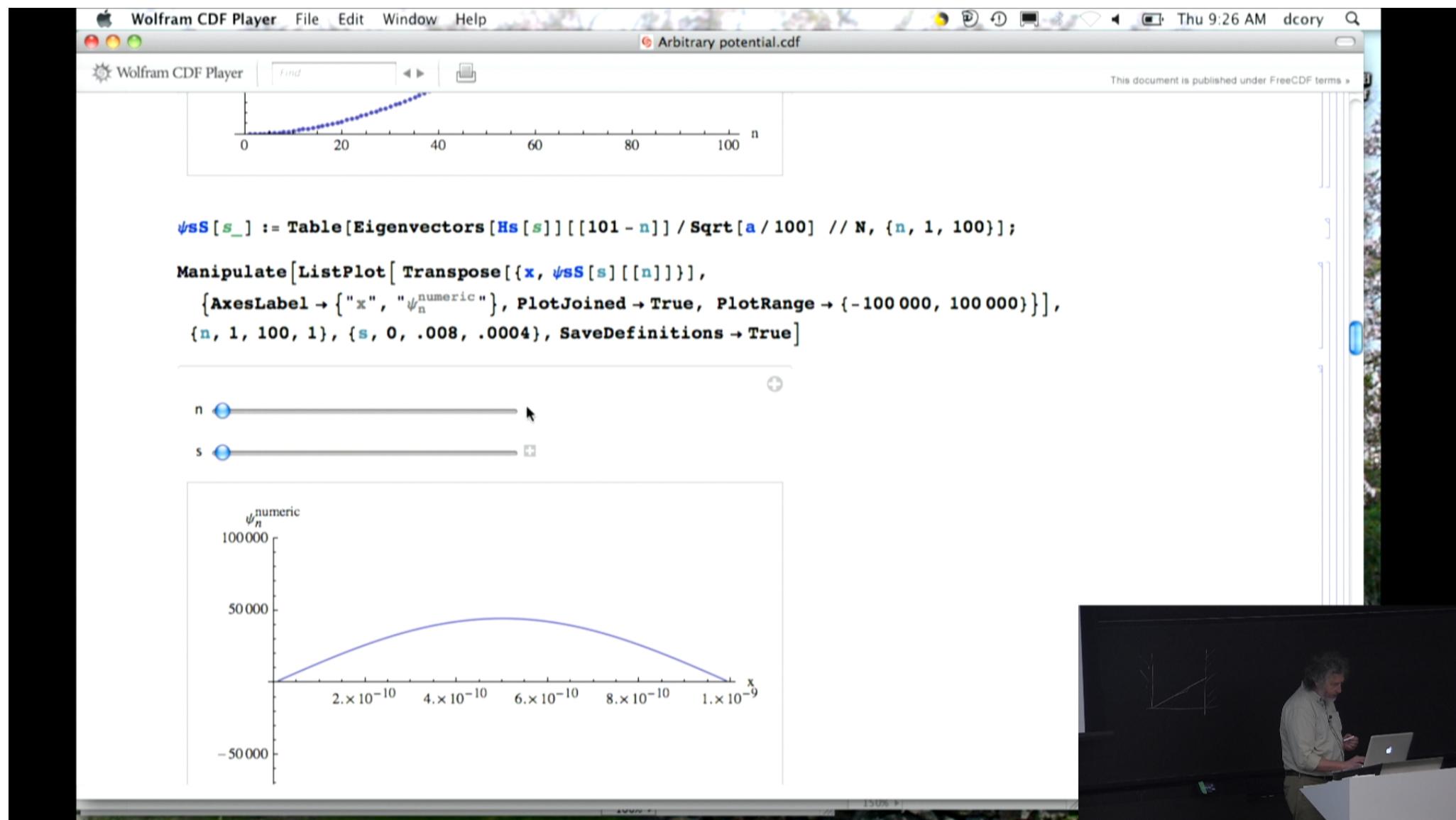


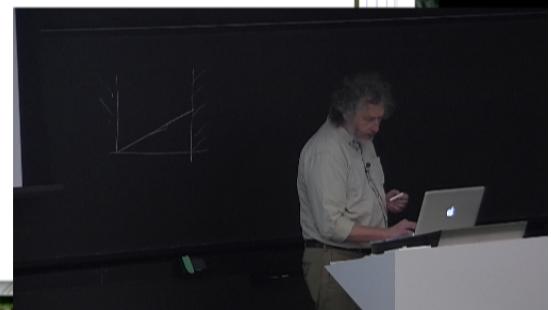
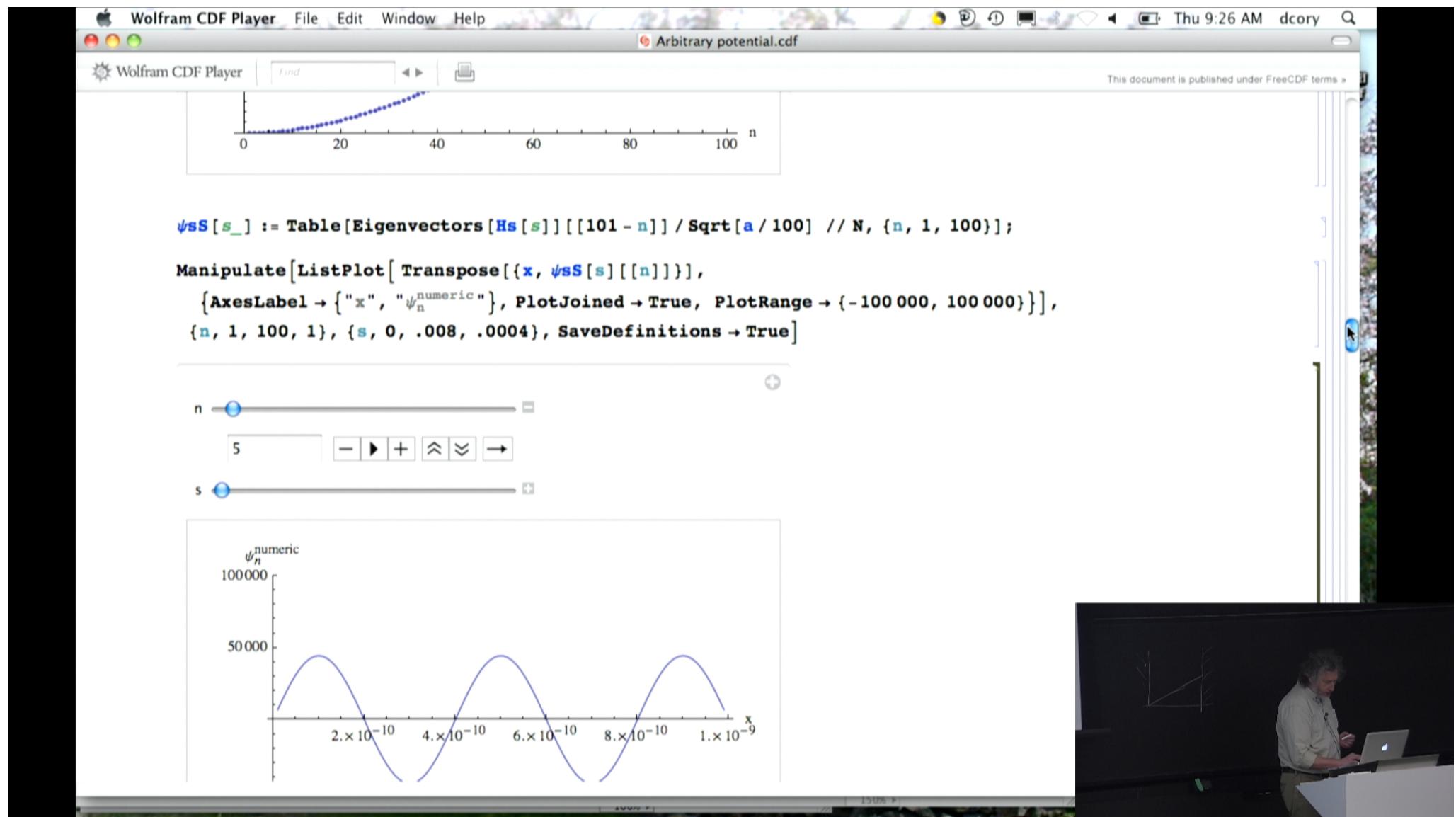












Wolfram CDF Player File Edit Window Help

Arbitrary potential.cdf

This document is published under FreeCDF terms >

```
{n, 1, 100, 1}, {s, 0, .008, .0004}, SaveDefinitions → True]
```

n = 5

s = 0.008

ψ_n^{numeric}

x

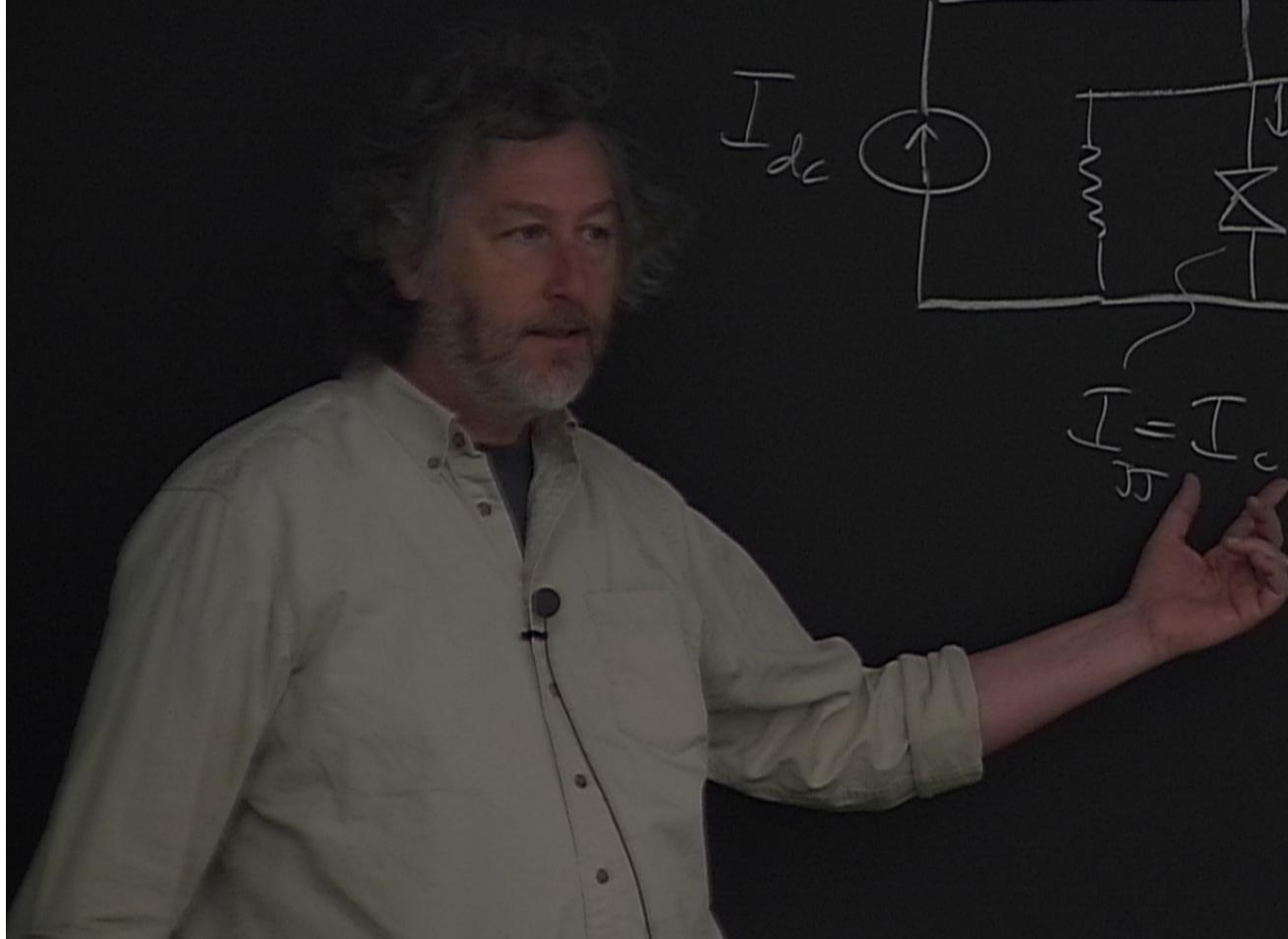
$\psi[1] = \text{Table}[\text{Eigenvectors}[\text{Hs}[0]][[101 - n]] / \text{Sqrt}[a/100] // N, \{n, 1, 100\}]$;

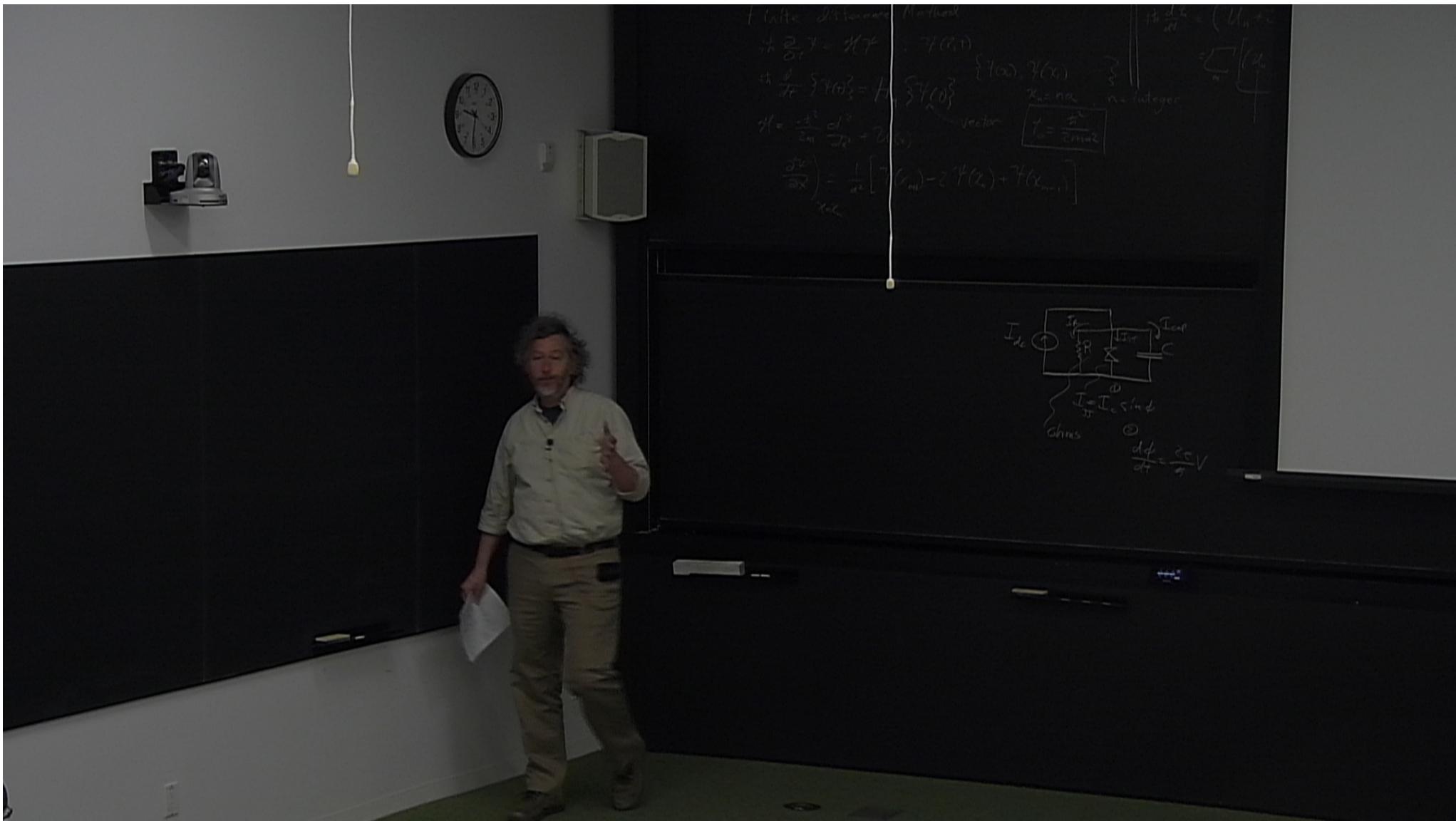
$\psi[2] = \text{Table}[\text{Eigenvectors}[\text{Hs}[0.0001]][[101 - n]] / \text{Sqrt}[a/100] // N, \{n, 1, 100\}]$;





$$I_{JJ} = I_c \sin \phi$$







$$\text{chms} \quad \textcircled{1} \quad \frac{I}{I_c} = \sin \phi$$

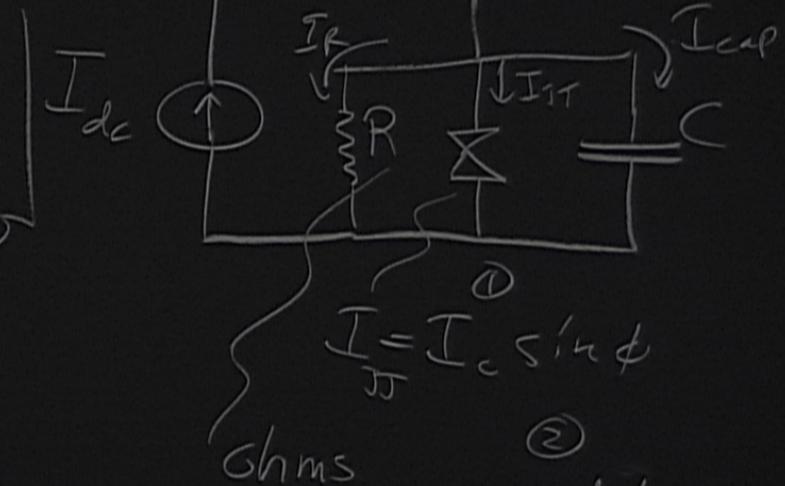
②

$$\frac{d\phi}{dt} = \frac{2\pi}{T} V$$

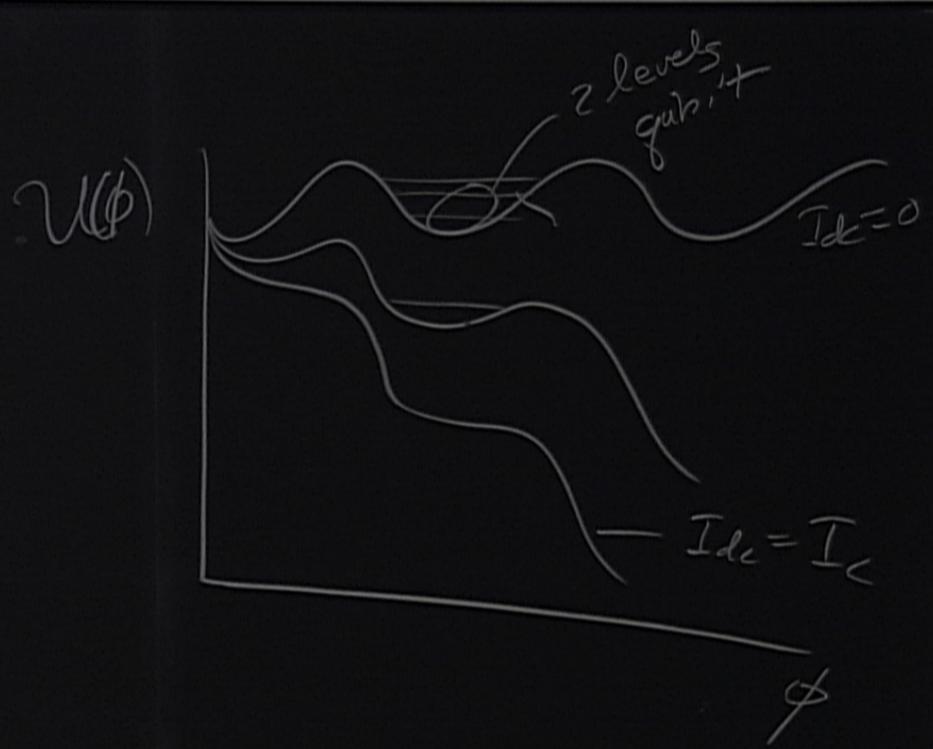
$$\frac{\hbar C}{2e} \frac{d^2\phi}{dt^2} + \frac{\hbar}{2eR} \frac{d\phi}{dt} = -I_c \sin\phi - I_{dc}$$

$$U(\phi) = -\cos\phi - \frac{I_{dc}}{I_c} \phi$$

Force



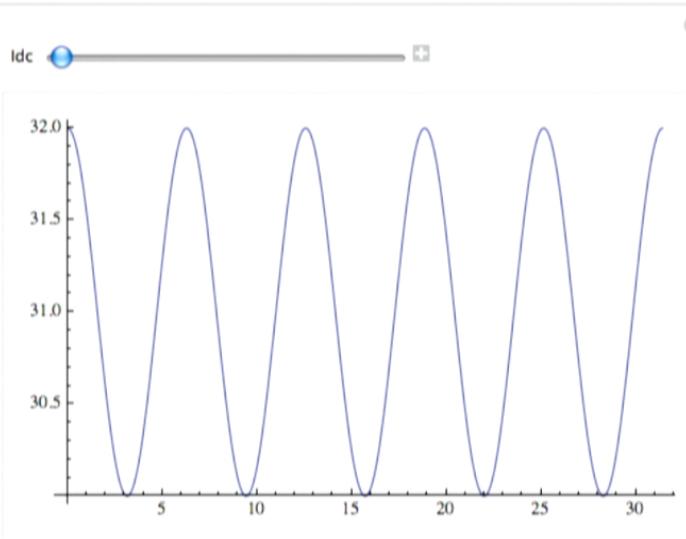
$$I_{dc} = I_R + I_{JJ} + I_{cap} \quad \frac{d\phi}{dt} = \frac{2e}{\hbar}$$



Potential for current biased JJ

```
In[1]:= U[\phi_, Ic_, Idc_] := - Ic \left( \frac{Idc}{Ic} \phi - \cos[\phi] \right) + 31
```

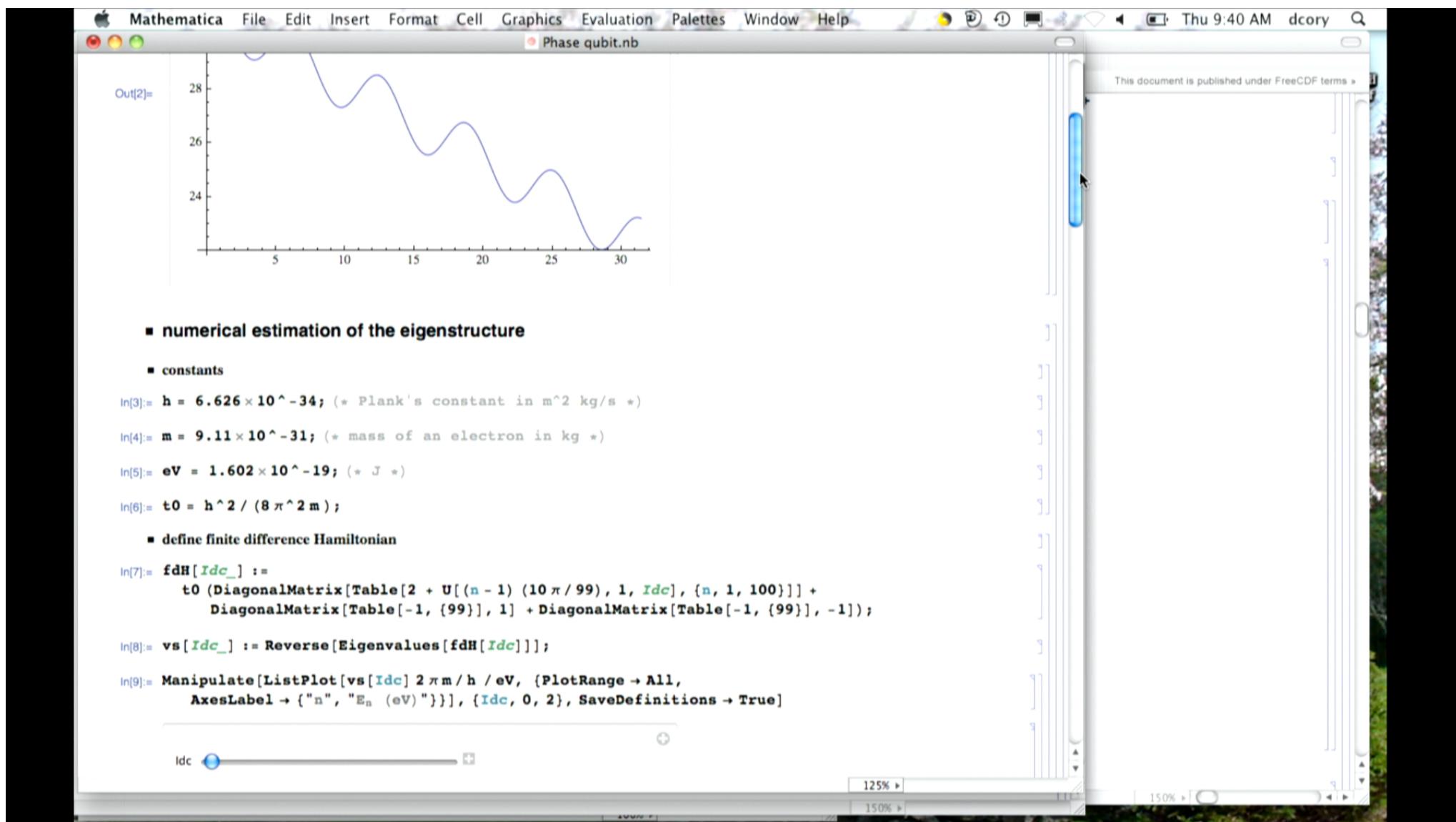
```
In[2]:= Manipulate[Plot[U[\phi, 1, Idc], {\phi, 0, 10 \pi}], {Idc, 0, 3}]
```



- numerical estimation of the eigenstructure

- constants

```
In[3]:= h = 6.626 \times 10^{-34}; (* Plank's constant in m^2 kg/s *)
```



Mathematica File Edit Insert Format Cell Graphics Evaluation Palettes Window Help

Phase qubit.nb

```
In[6]:= t0 = h^2 / (8 π^2 m);

* define finite difference Hamiltonian

In[7]:= fdH[Idc] :=
  t0 (DiagonalMatrix[Table[2 + U[(n - 1) (10 π/99), 1, Idc], {n, 1, 100}]] +
    DiagonalMatrix[Table[-1, {99}], 1] + DiagonalMatrix[Table[-1, {99}], -1]);

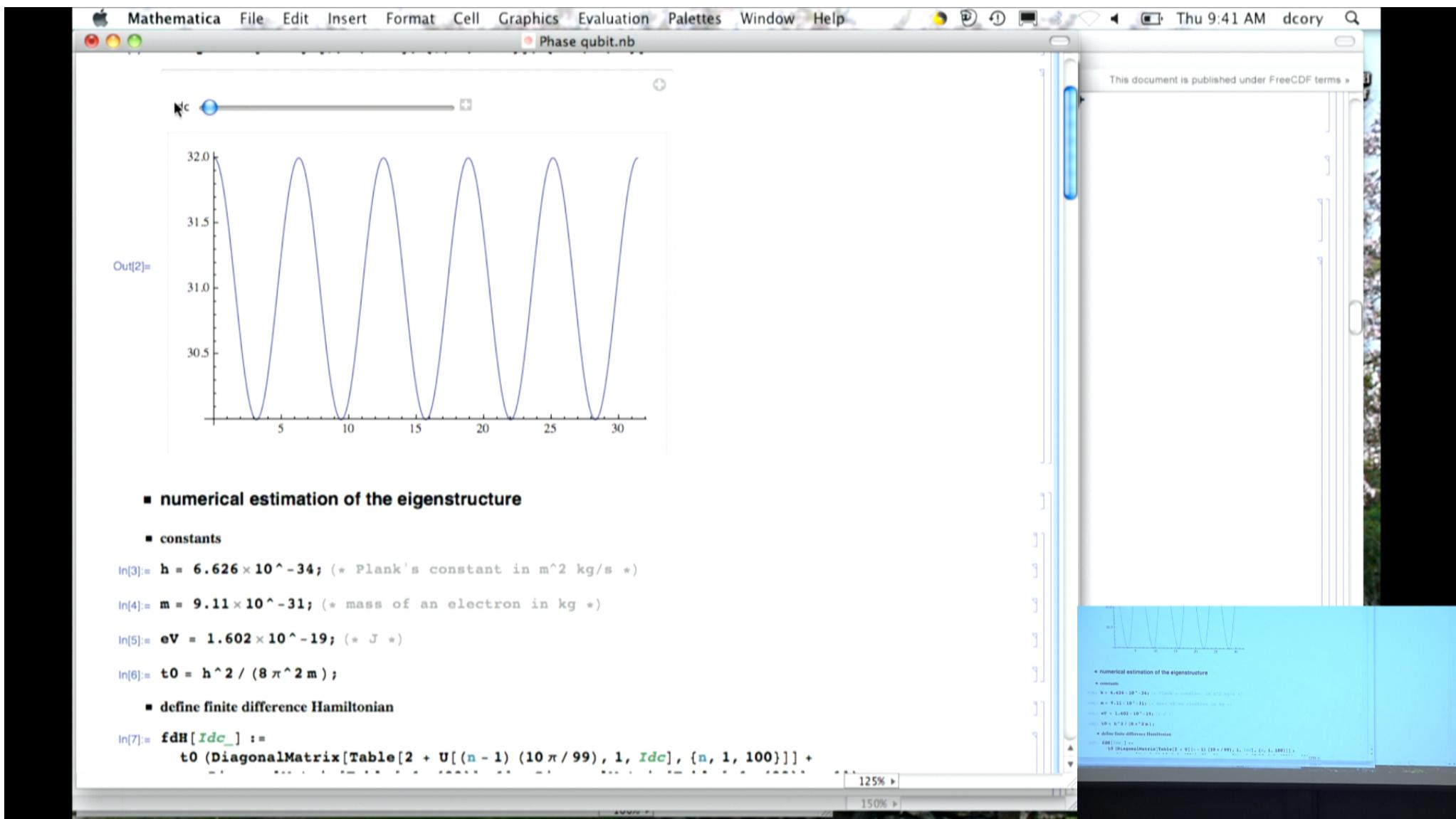
In[8]:= vs[Idc] := Reverse[Eigenvalues[fdH[Idc]]];

In[9]:= Manipulate[ListPlot[vs[Idc] 2 π m/h/eV, {PlotRange -> All,
  AxesLabel -> {"n", "En (eV)"}}], {Idc, 0, 2}, SaveDefinitions -> True]
```

Out[9]=

In[10]:= ψ[*Idc*] := Table[Eigenvectors[fdH[*Idc*]][[101 - n]] / Sqrt[1/100] // N, {n, 1, 100}];

In[11]:= Manipulate[ListPlot[Transpose[{x, ψ[*Idc*][[n]]}],
 {AxesLabel -> {"x", "ψ_n^{numeric}"}, PlotJoined -> True, PlotRange -> All}], {n, 1, 100, 1}, SaveDefinitions -> True]



Mathematica File Edit Insert Format Cell Graphics Evaluation Palettes Window Help

Phase qubit.nb

```
In[6]:= t0 = h^2 / (8 \pi^2 m);
* define finite difference Hamiltonian

In[7]:= fdH[_Idc_] :=
  t0 (DiagonalMatrix[Table[2 + U[(n - 1) (10 \pi / 99), 1, _Idc], {n, 1, 100}]] +
    DiagonalMatrix[Table[-1, {99}], 1] + DiagonalMatrix[Table[-1, {99}], -1]);

In[8]:= vs[_Idc_] := Reverse[Eigenvalues[fdH[_Idc]]];

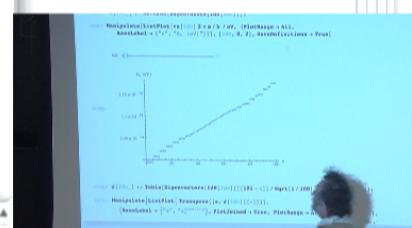
In[9]:= Manipulate[ListPlot[vs[_Idc] 2 \pi m/h/eV, {PlotRange -> All,
  AxesLabel -> {"n", "E_n (eV)"}}], {_Idc, 0, 2}, SaveDefinitions -> True]
```

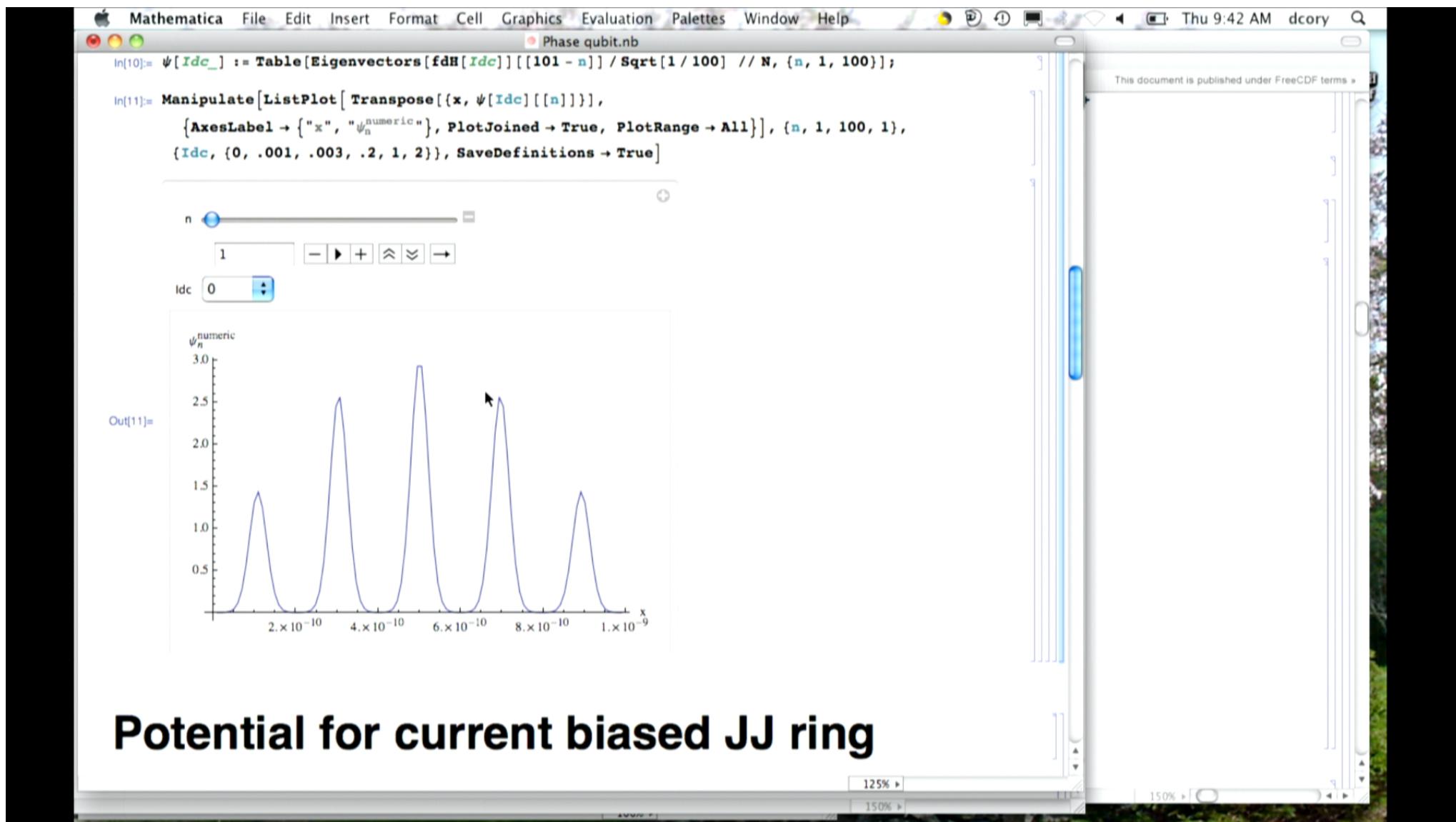
Out[9]=

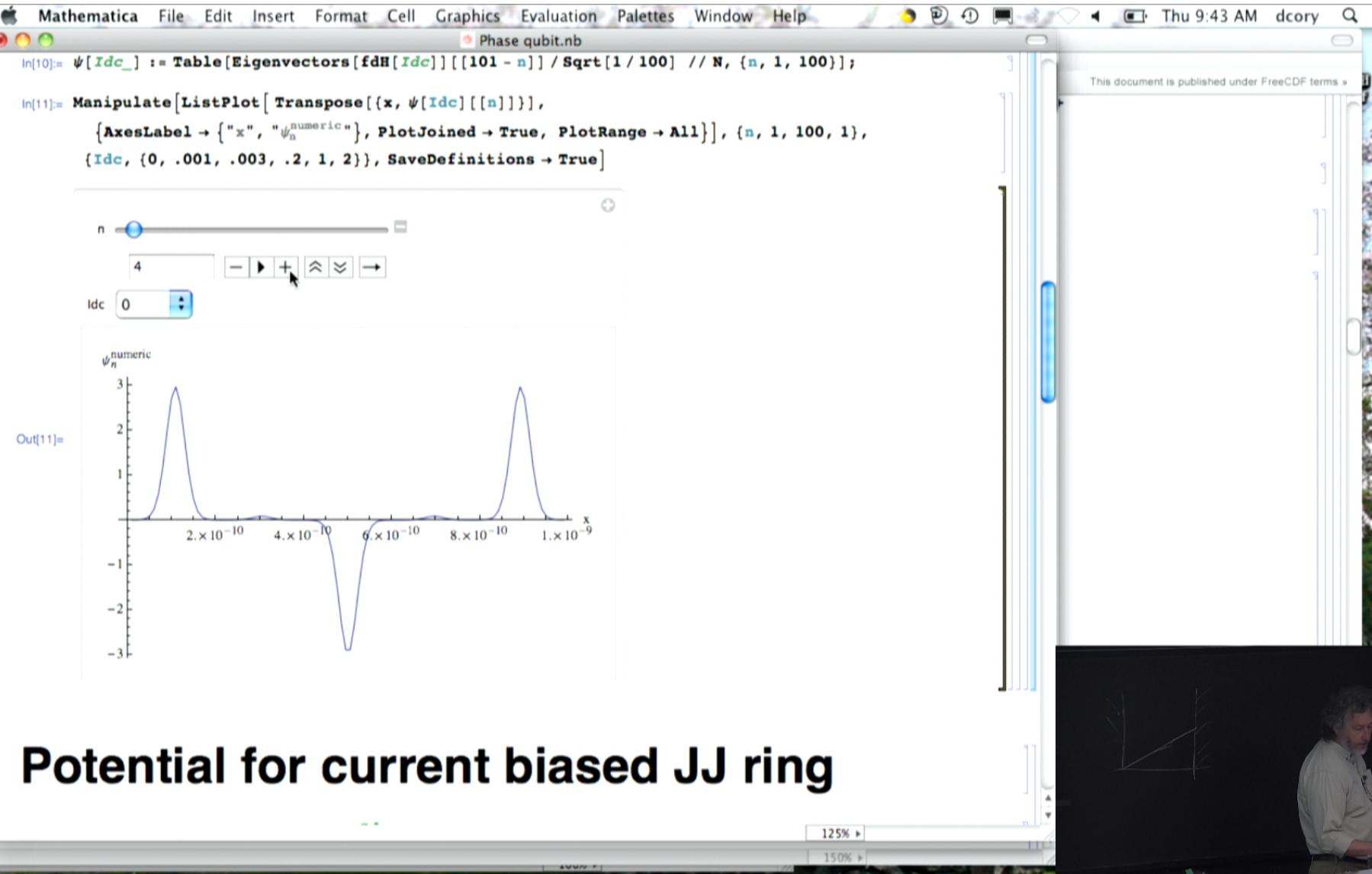
In[10]:= \psi[_Idc_] := Table[Eigenvectors[fdH[_Idc]][[101 - n]] / Sqrt[1/100] // N, {n, 1, 100}];

In[11]:= Manipulate[ListPlot[Transpose[{x, \psi[_Idc][[n]]}],
 {AxesLabel -> {"x", "\psi_n^{numeric}"}, PlotJoined -> True, PlotRange -> All}], {n, 1, 100, 1},

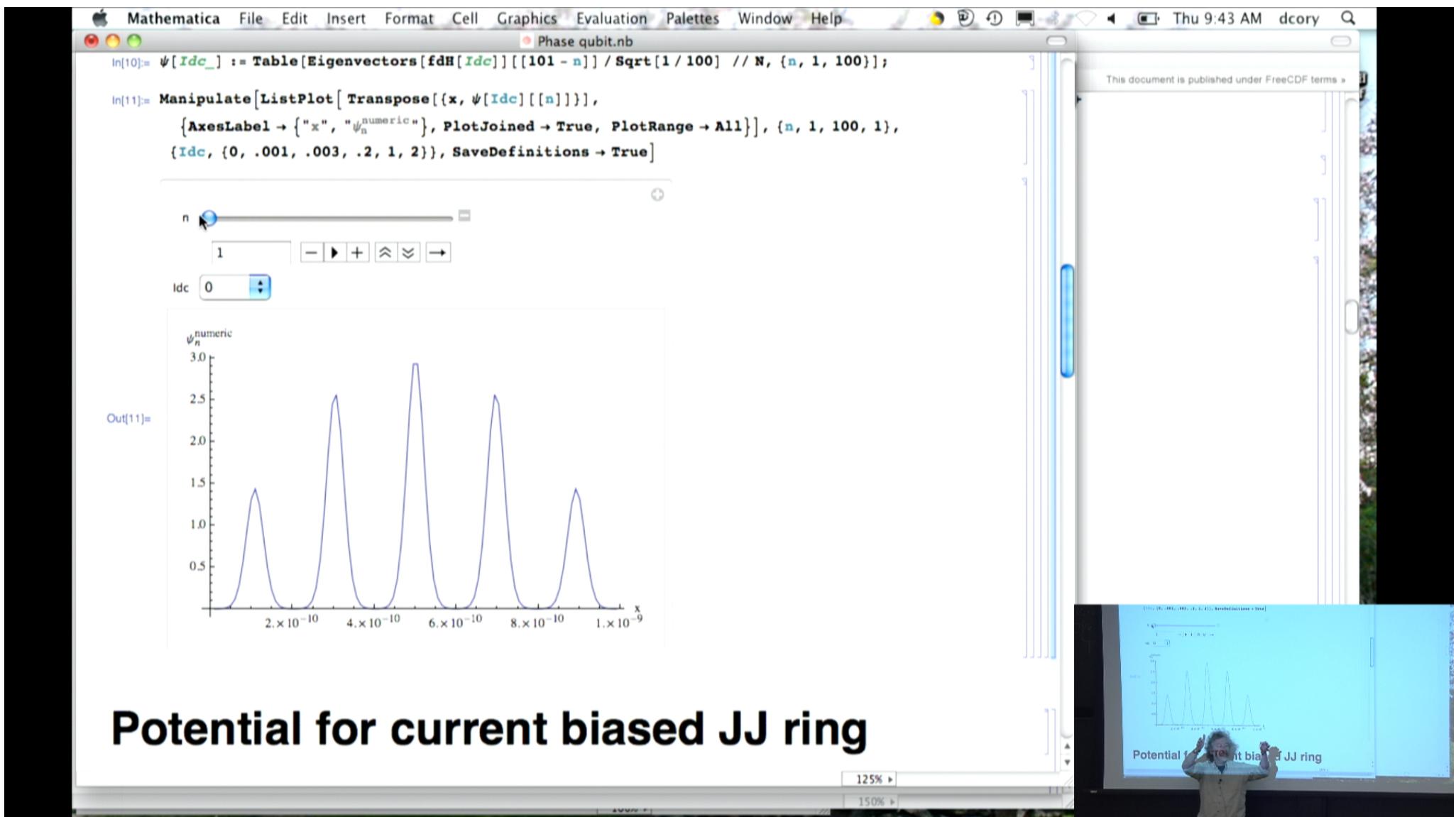
This document is published under FreeCDF terms.

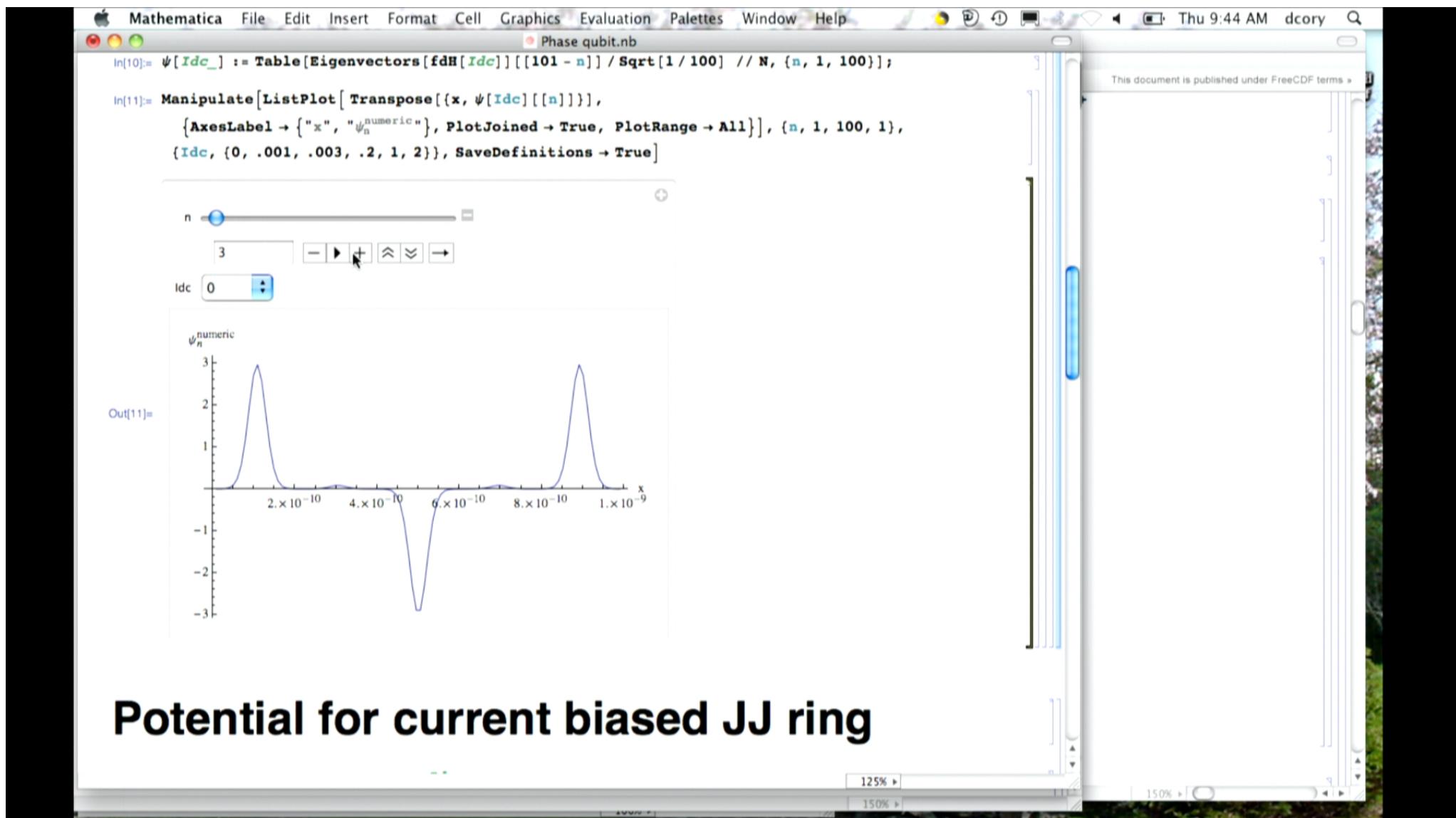




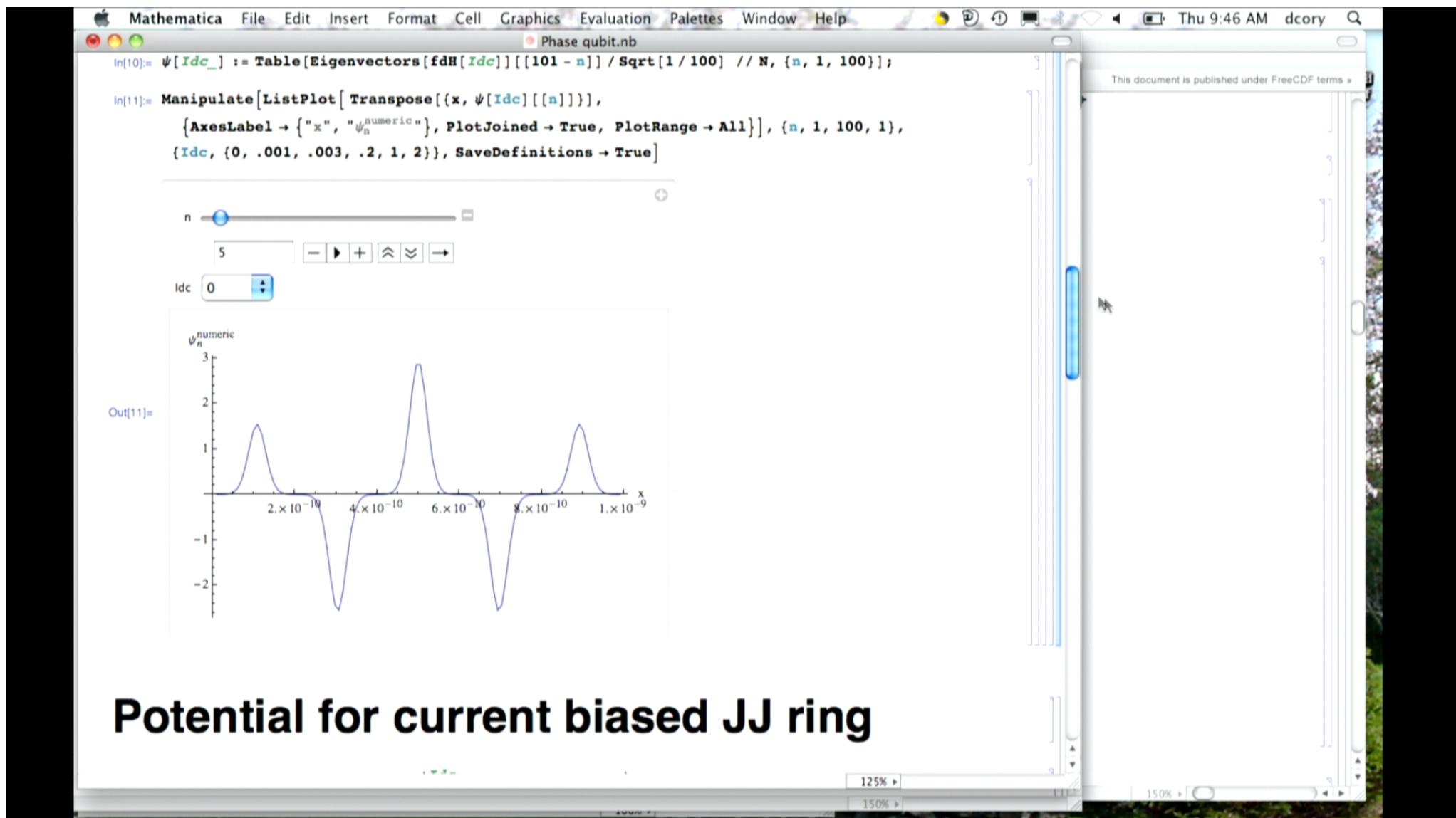


Potential for current biased JJ ring

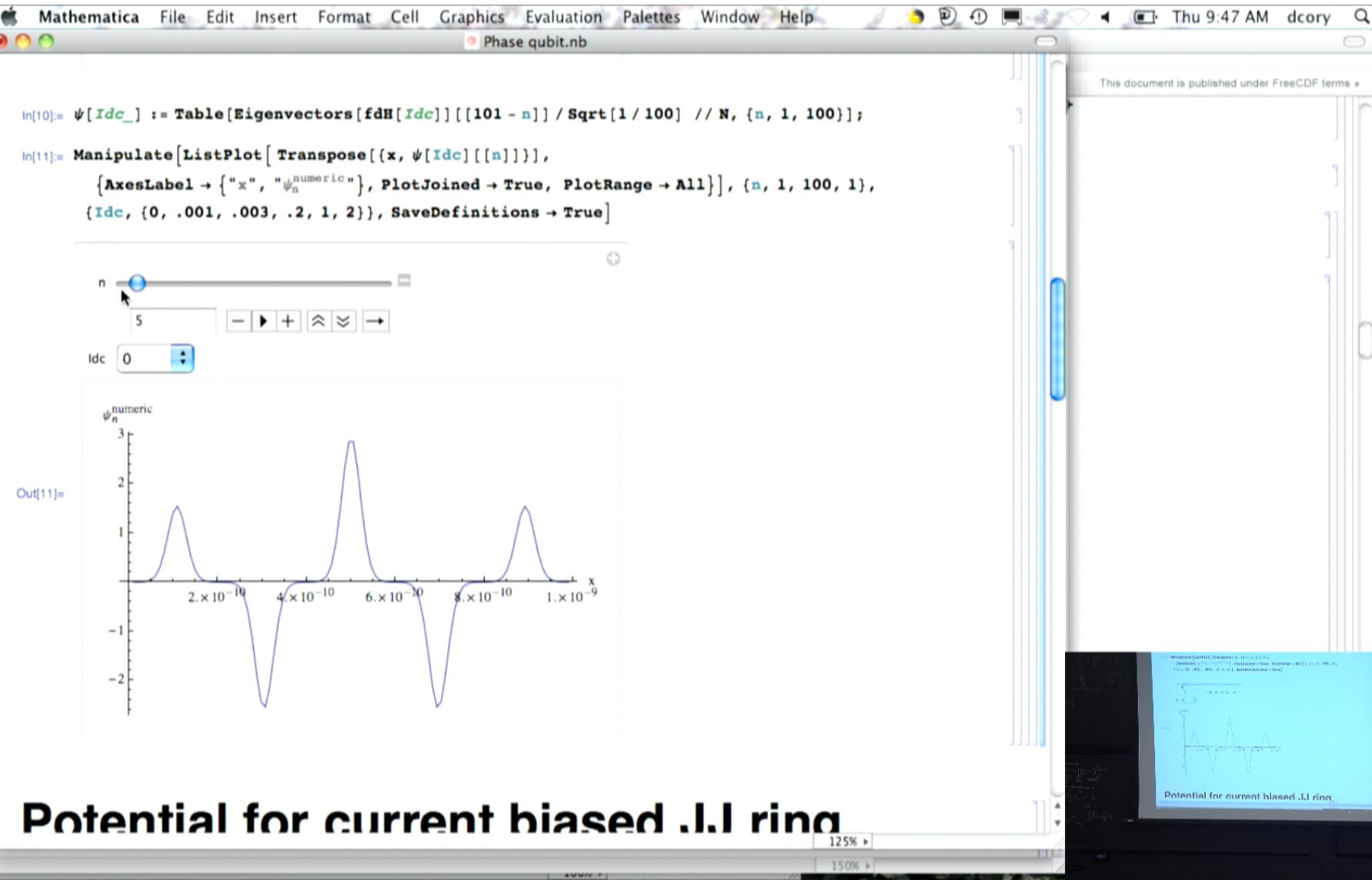




Potential for current biased JJ ring



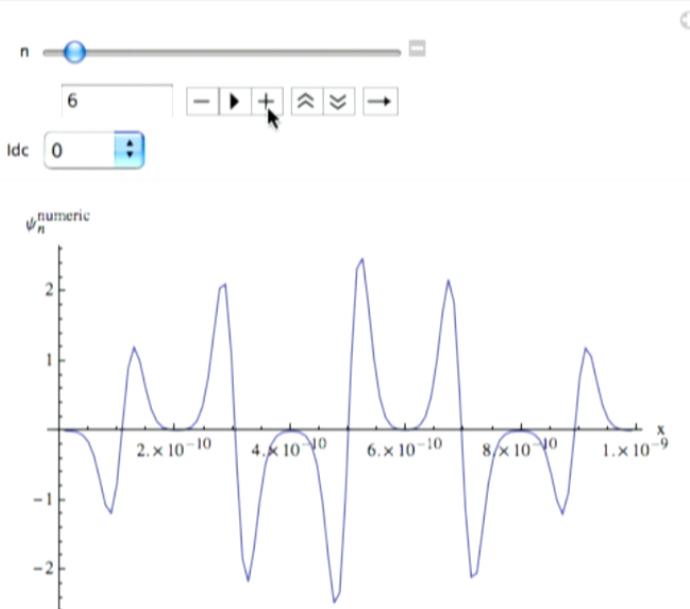
Potential for current biased JJ ring



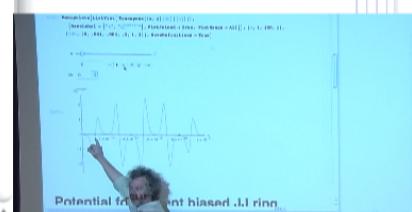
```
In[10]:=  $\psi[\text{Idc}_n] := \text{Table}[\text{Eigenvectors}[\text{fdH}[\text{Idc}]][[101 - n]] / \text{Sqrt}[1/100] // \text{N}, \{n, 1, 100\}]$ ;
In[11]:= Manipulate[ListPlot[Transpose[{x,  $\psi[\text{Idc}][[n]]$ }],  

    {AxesLabel -> {"x", " $\psi_n^{\text{numeric}}$ "}, PlotJoined -> True, PlotRange -> All}], {n, 1, 100, 1},  

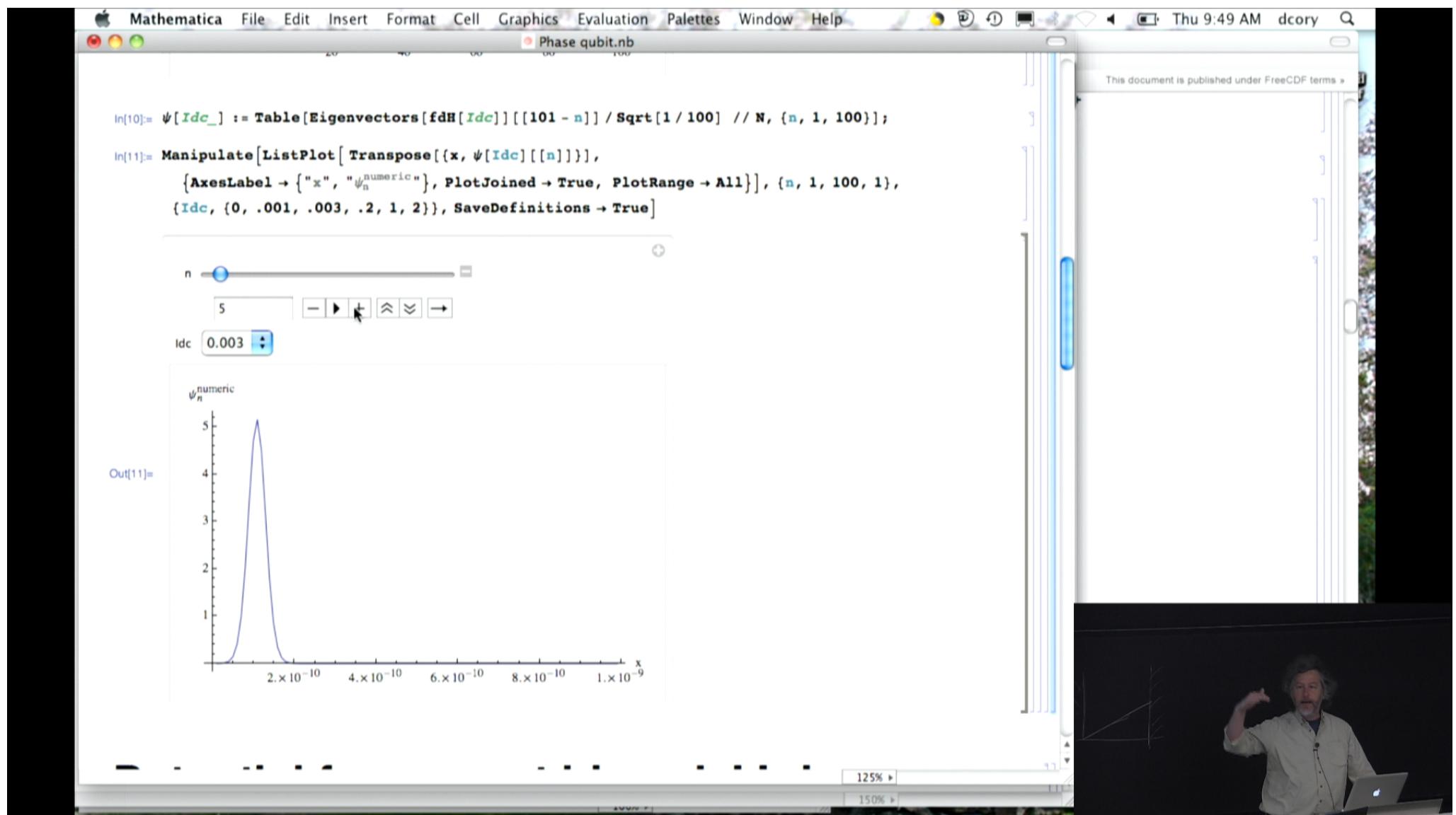
    {Idc, {0, .001, .003, .2, 1, 2}}, SaveDefinitions -> True]
```



Out[11]=



Potential for current biased J.J ring



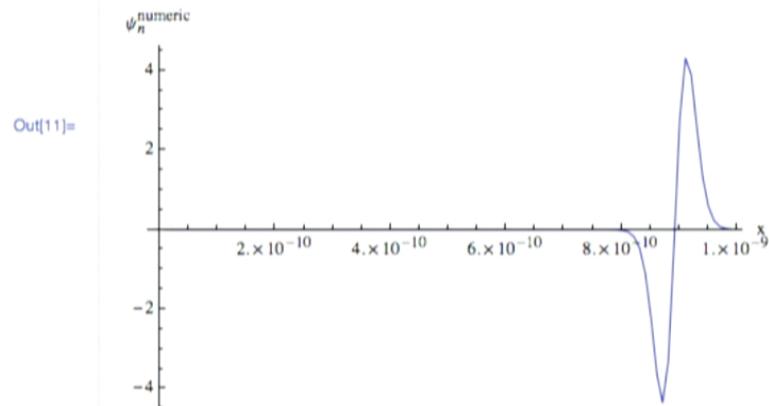
```
In[10]:=  $\psi[\text{Idc}_\perp] := \text{Table}[\text{Eigenvectors}[\text{fdH}[\text{Idc}]][[101 - n]] / \text{Sqrt}[1/100] // \text{N}, \{n, 1, 100\}]$ ;
In[11]:= Manipulate[ListPlot[Transpose[{x,  $\psi[\text{Idc}][[n]]$ }]],
 {AxesLabel -> {"x", " $\psi_n^{\text{numeric}}$ "}, PlotJoined -> True, PlotRange -> All}], {n, 1, 100, 1},
 {Idc, {0, .001, .003, .2, 1, 2}}, SaveDefinitions -> True]
```

n

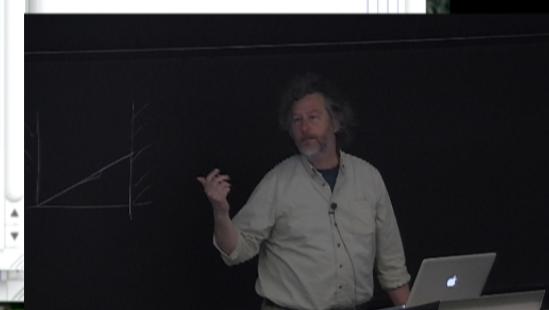
6

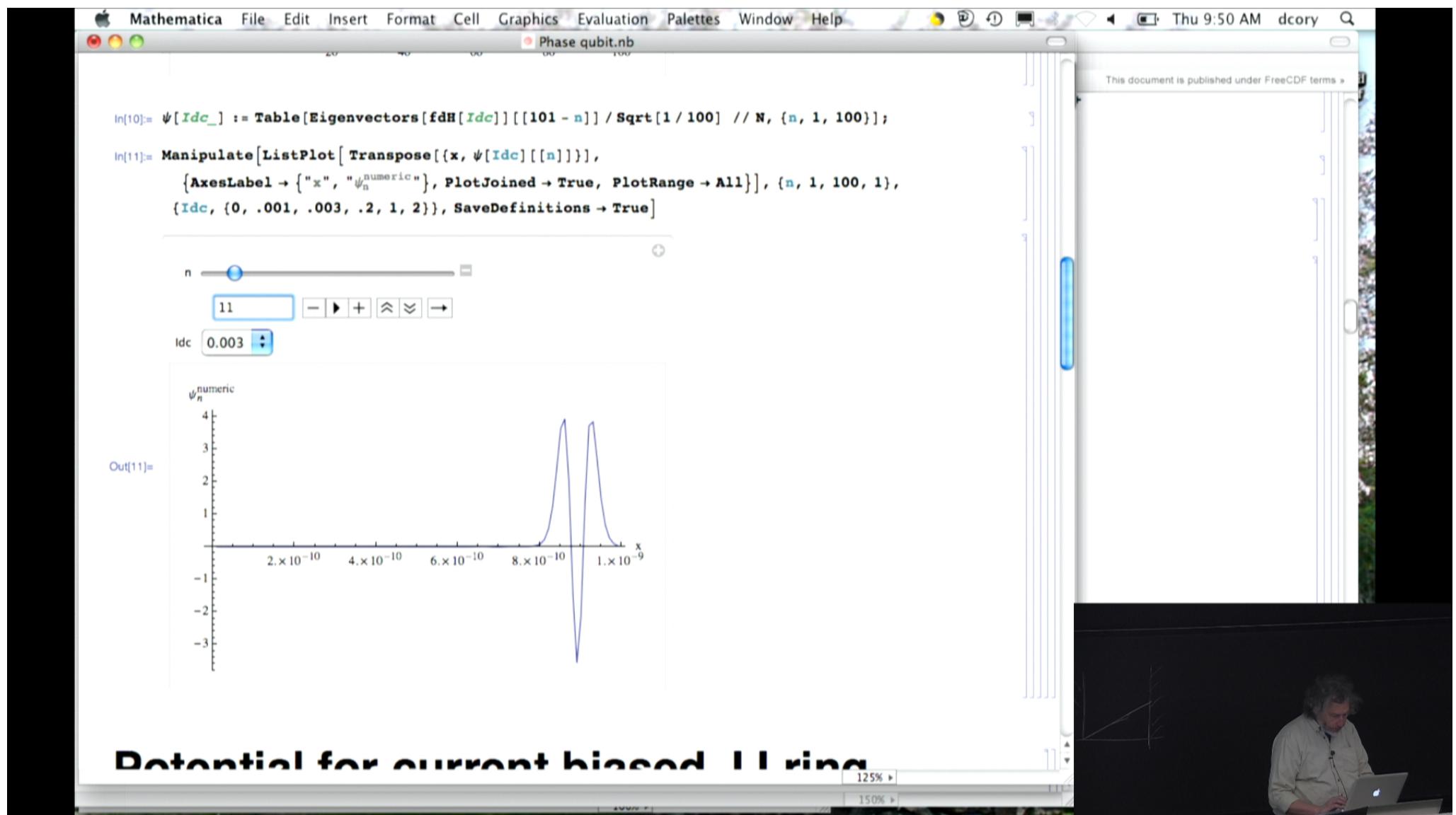
Idc

0.003



Potential for current biased IJ rings





```
In[10]:=  $\psi[\text{Idc}_\_\_] := \text{Table}[\text{Eigenvectors}[\text{fdH}[\text{Idc}]][[101 - n]] / \text{Sqrt}[1/100] // \text{N}, \{n, 1, 100\}]$ ;
In[11]:= Manipulate[ListPlot[Transpose[{x,  $\psi[\text{Idc}][[n]]$ }]],
 {AxesLabel -> {"x", " $\psi_n^{\text{numeric}}$ "}, PlotJoined -> True, PlotRange -> All}], {n, 1, 100, 1},
 {Idc, {0, .001, .003, .2, 1, 2}}, SaveDefinitions -> True]
```

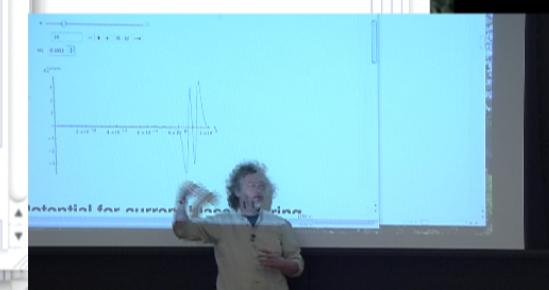
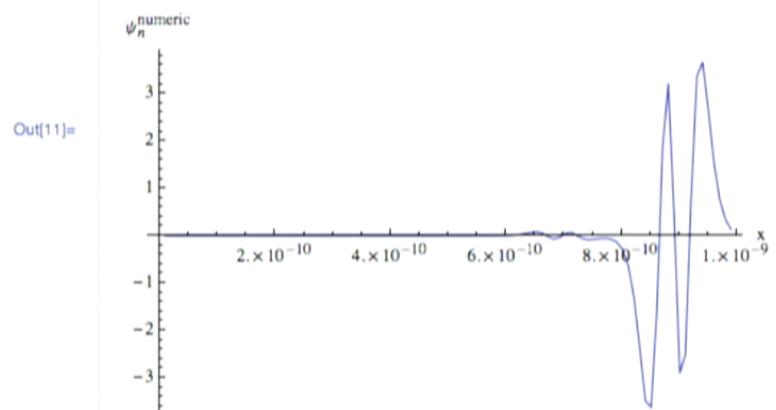
n

16

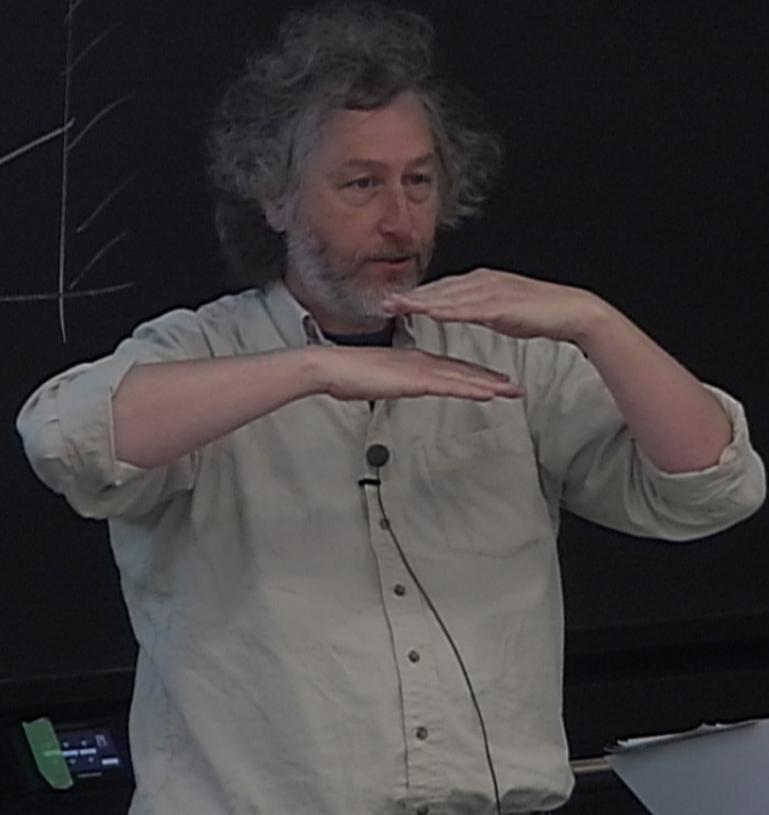
- ▶ + ⌂ ⌃ →

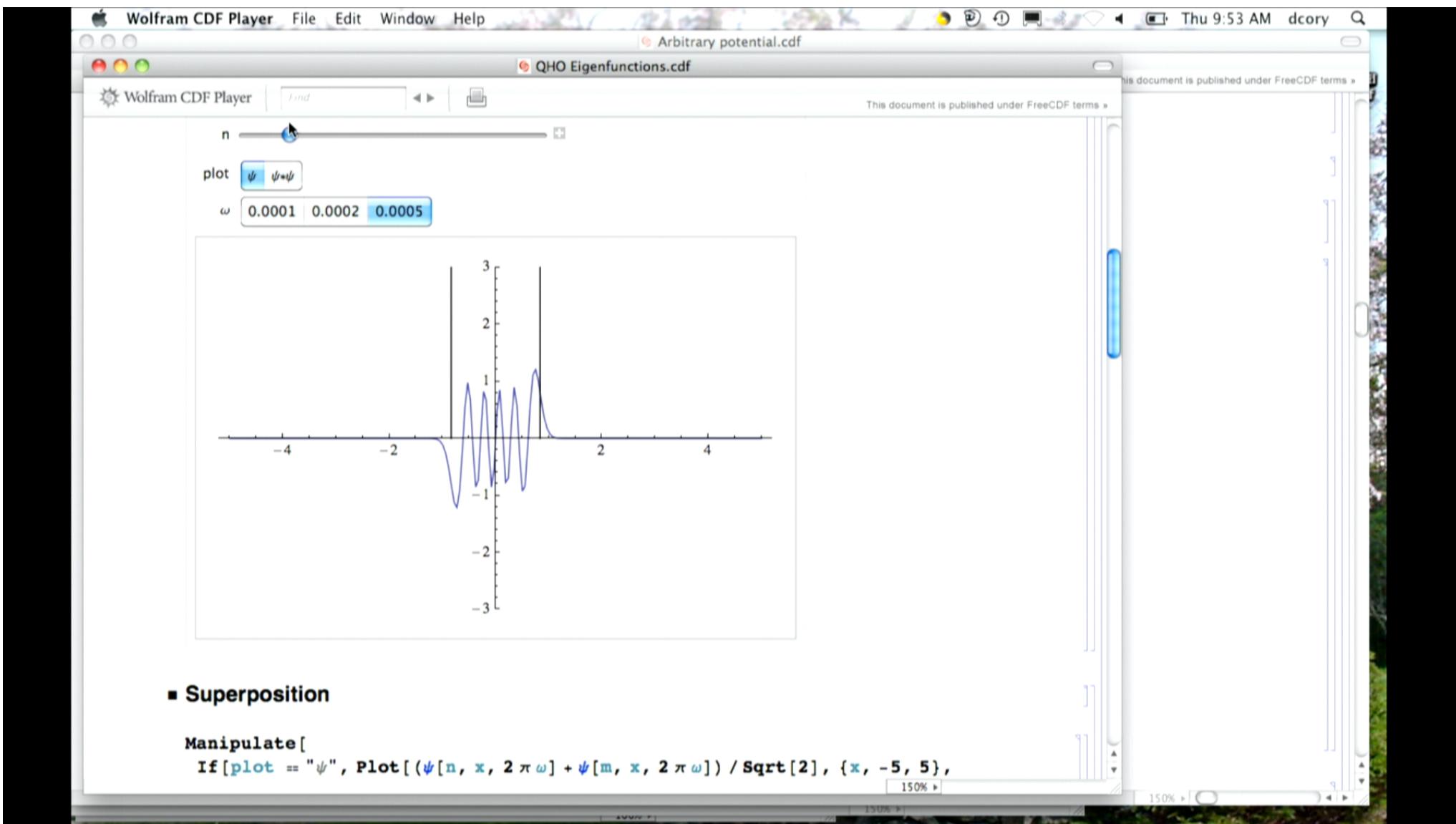
Idc

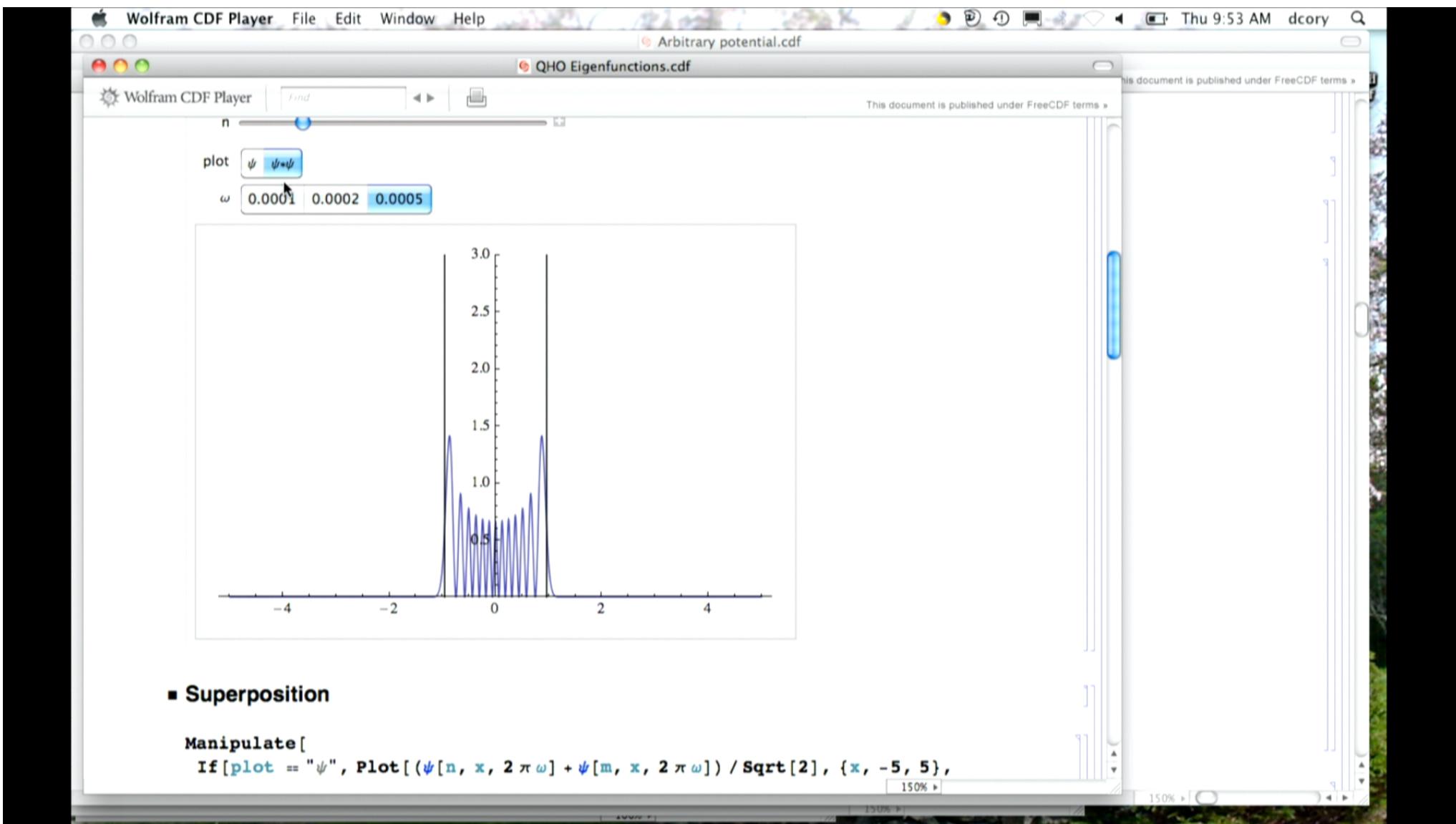
0.003

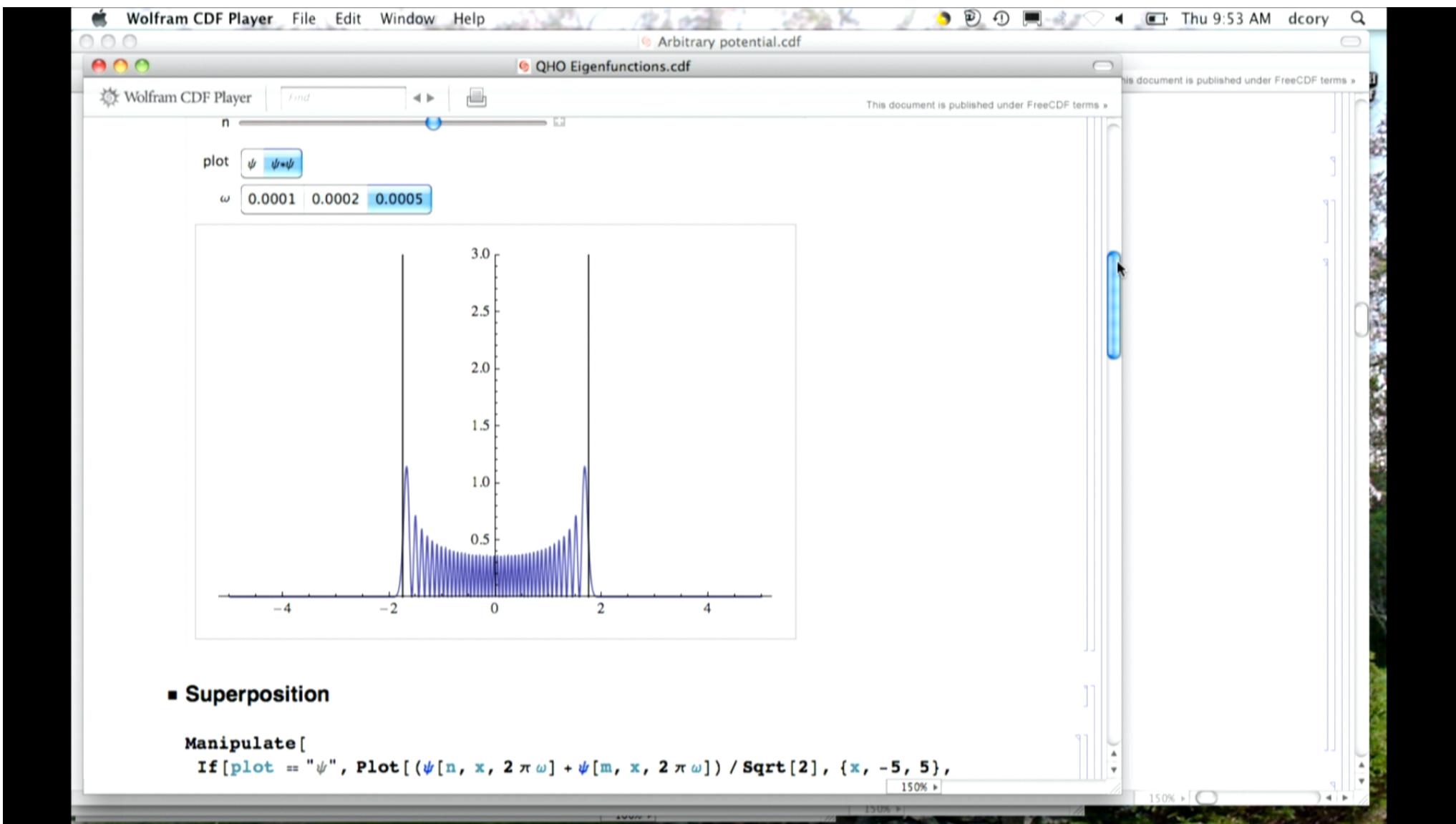


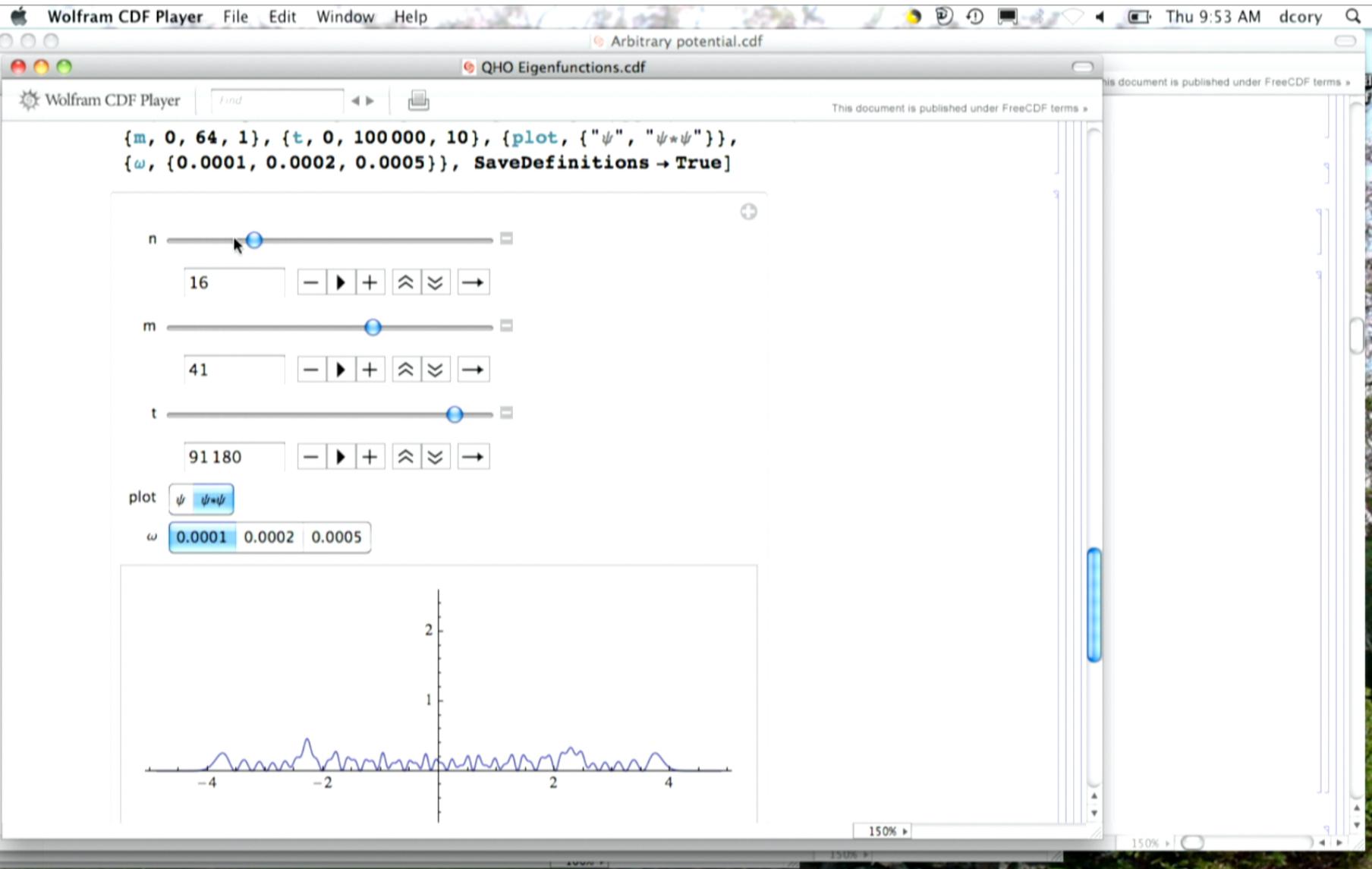
Potential for current biased I I rings

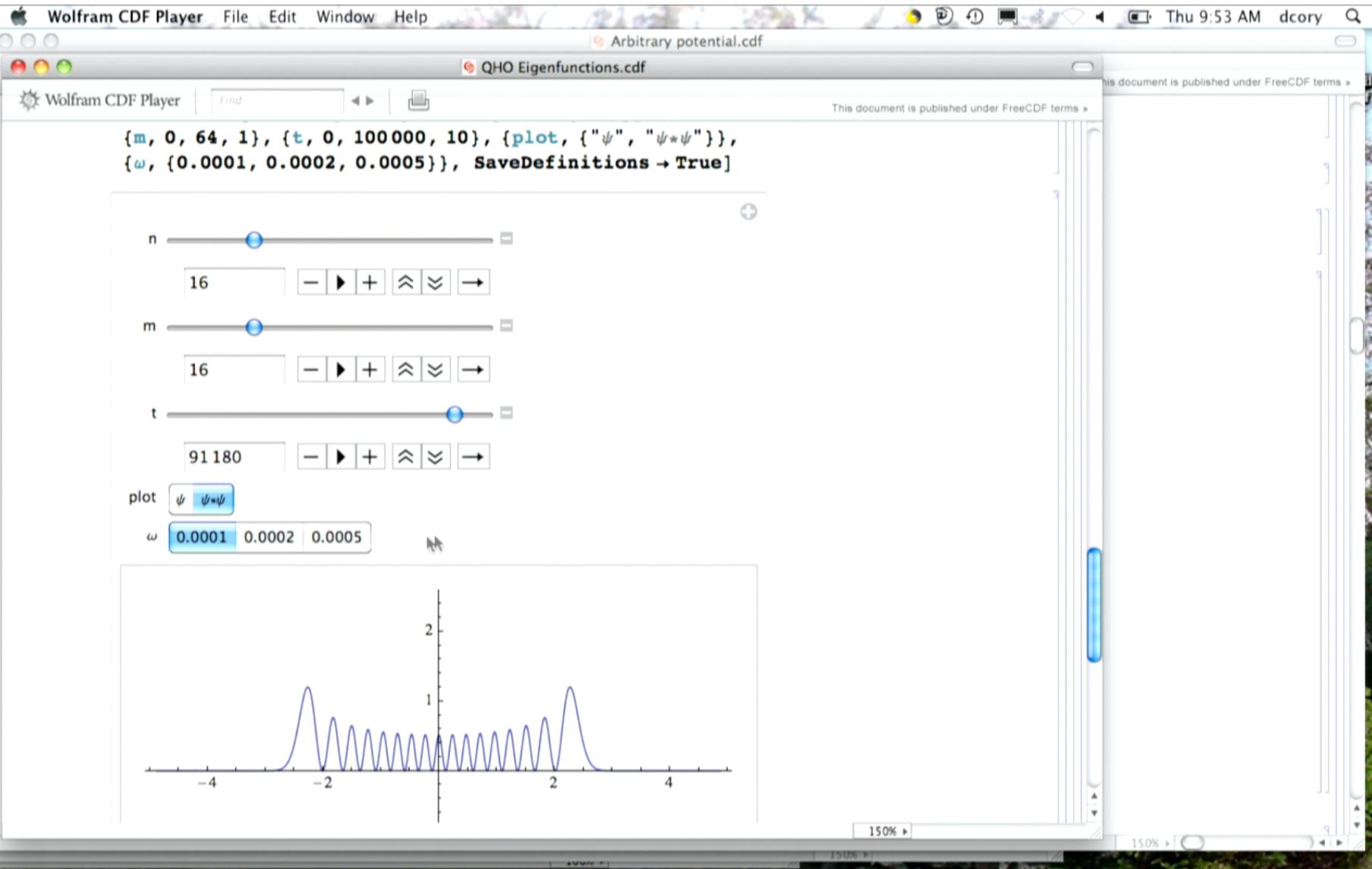


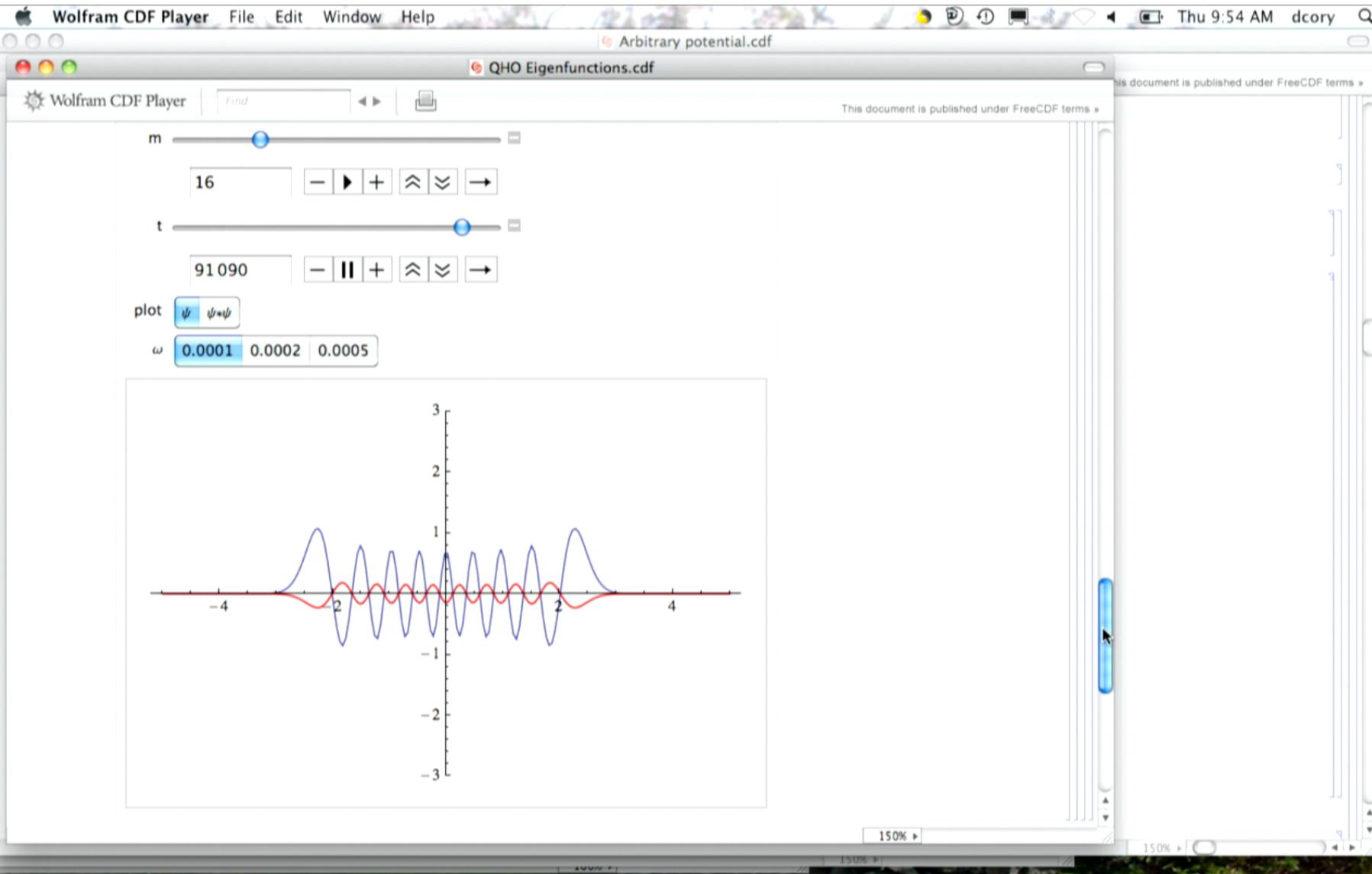












This document is published under FreeCDF terms »

QHO Eigenfunctions.cdf

Wolfram CDF Player

Find

 m

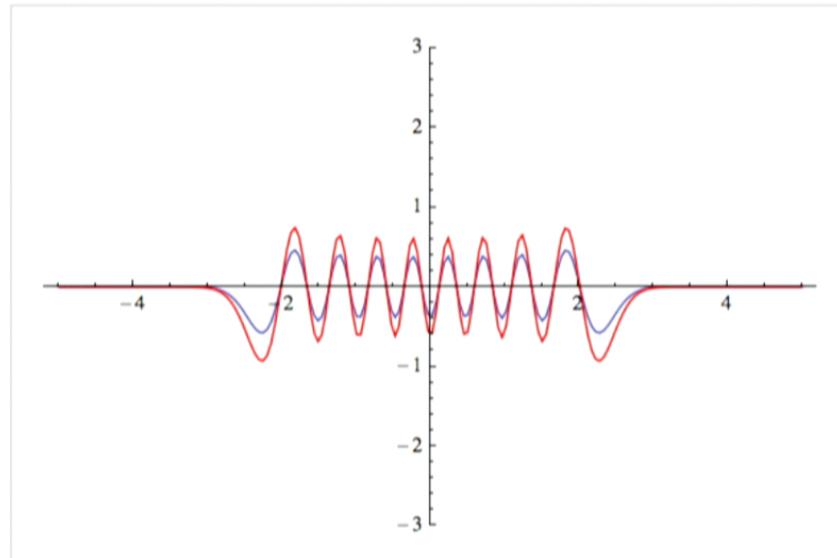
16

- ▶ + ⌈ ⌉ →

 t

10030

- II + ⌈ ⌉ →

plot ψ $\psi\#\psi$ ω 0.0001 0.0002 0.0005

150% ↑

150%

Start $\psi_n(\alpha)$ ^{$t \rightarrow t$} "local"
bound state

$t \rightarrow t$ $\alpha \rightarrow \alpha + \varepsilon$

new eigenstates, $\{\psi_n(\alpha + \varepsilon)\}$

$$\psi_n(\alpha) = \sum c_k \psi_k(\alpha + \varepsilon)$$

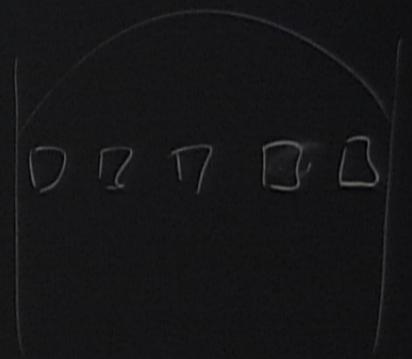
$$\psi(t) = \sum c_k \psi_k(\alpha + \varepsilon) e^{-i E_k \frac{\varepsilon}{\hbar} t}; \text{ determine these}$$

start $\psi_n(\alpha)$ ^{tilt} ^{"local"} = bound state

$$t/t \rightarrow \alpha + \varepsilon$$

New eigenstates: $\{\psi_n(\alpha + \varepsilon)\}$

$$\begin{aligned} \psi_n(\alpha) &= \sum c_k \psi_k(\alpha + \varepsilon) & \text{change in slope} = \frac{\varepsilon}{\hbar} \\ \underline{\psi(t)} &= \sum c_k \psi_k(\alpha + \varepsilon) e^{-E_k \frac{\varepsilon}{\hbar} t} ; \underline{\delta \text{ time}} & \text{determine these} \\ V(t) & \end{aligned}$$



$$\frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{x}_i} \right) - \frac{\partial \mathcal{L}}{\partial x_i} = 0$$

$$\frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{x}_i} \right) - \frac{\partial \mathcal{L}}{\partial x_i} = T_{ii} - \underbrace{T_{ij} \omega_j}_{\text{Coriolis}} + \underbrace{T_{ij} \frac{d\omega_j}{dt}}_{\text{Centrifugal}} + \underbrace{\frac{\partial \mathcal{L}}{\partial x_i}}_{\text{Lagrange}}$$

$$\text{Time dep: } T_{ii}(t) = e^{-i\omega t} T_{ii}(0)$$

