

Title: Quantum Gravity (Review) - Lecture 10

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URL: <http://pirsa.org/12020002>

Abstract:

PRD 41 (1990) 3785

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problematic when $\mathcal{L}_{\text{phys}}$, $\mathcal{I}_{\text{phys}}$ are nonlinear

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↳ : problematic when $\mathcal{L}_{\text{phys}}$, $\mathcal{I}_{\text{phys}}$ are nonlinear

e.g. "Dirac quantization"

ex. :

ex.: $\mathcal{F} = \mathbb{R}^4$ $(q_1, q_2, q_3, p_1, p_2, p_3)$, $\phi = p_3 =$

ex.: $\mathcal{P} = \mathbb{R}^6$ $(q_1, q_2, q_3, p_1, p_2, p_3)$, $\phi = p_3 = 0$

$\mathcal{H} = L^2(\mathbb{R}^3, dq_i) \ni \psi(q_1, q_2, q_3)$

$\hat{p}_3 \psi = -i\hbar \frac{\partial}{\partial q_3} \psi = 0 \Rightarrow \psi_{\text{phys}} = \psi_{\text{phys}}(q_1, q_2) \in \mathcal{L}_{\text{phys}} = L^2(\mathbb{R}^2, dq_i)$

ex.: $\mathcal{P} = \mathbb{R}^6$ $(q_1, q_2, q_3, p_1, p_2, p_3)$, $\phi = p_3 = 0$

$$\mathcal{H}_{\text{aux}} = L^2(\mathbb{R}^3, dq_i) \ni \psi(q_1, q_2, q_3)$$

$$\hat{\phi}\psi \equiv \hat{p}_3\psi = -i\hbar \frac{\partial}{\partial q_3} \psi = 0 \Rightarrow \psi_{\text{phys}} = \psi_{\text{phys}}(q_1, q_2) \in \mathcal{R}_{\text{phys}} = L^2(\mathbb{R}^2, dq_i)$$

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"Dirac consistency" must have $[\hat{\phi}_a, \hat{\phi}_b] \psi_{\text{phys}} = 0$

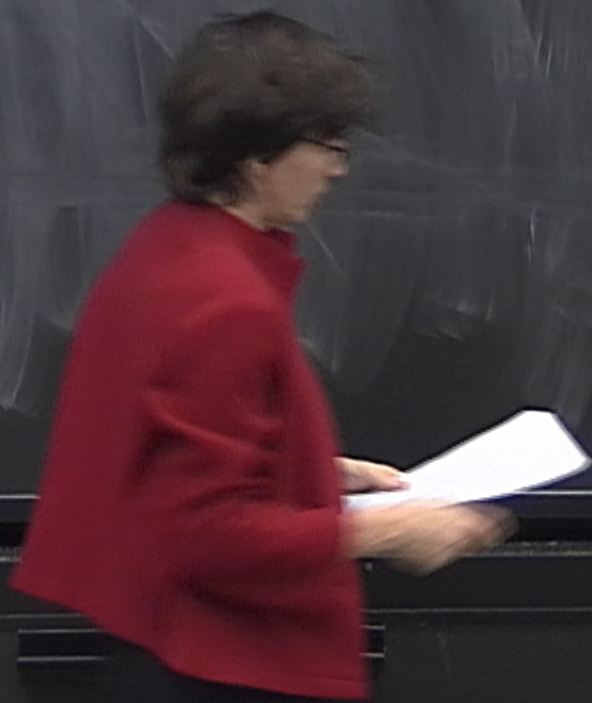
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"Dirac consistency" : must have $[\hat{\phi}_a, \hat{\phi}_b]\psi_{\text{phys}} = 0$

$$\Rightarrow [\hat{\Phi}_a, \hat{\Phi}_b] = i\hbar \hat{f}_{ab}^c \hat{\Phi}_c, \quad \text{in general } [\hat{\Phi}_a, \hat{\Phi}_b] = i\hbar \hat{f}_{ab}^c \hat{\Phi}_c + \hbar^2$$

$$\Rightarrow [\hat{\phi}_a, \hat{\phi}_b] = i\hbar \hat{f}_{ab}^c \hat{\phi}_c, \quad \text{in general } [\hat{\phi}_a, \hat{\phi}_b] = i\hbar \hat{f}_{ab}^c \hat{\phi}_c + \hbar^2 \hat{\mathcal{D}}_{ab}$$

\rightarrow extra conditions $\hat{\mathcal{D}}_{ab} \psi = 0$



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$$\rightarrow \text{extra conditions } \hat{D}_{ab} \psi = 0$$

\Rightarrow gauge invariance is broken at the quantum level,
"gauge anomaly"

Dirac quantization for gravity

- formal "Hilbert space" of $\Psi[h_{ij}]$
- promote PBs to canonical commutators

$$[\hat{h}_{ij}(x), \hat{\pi}^{kl}(y)] = \frac{i\hbar}{2} (\delta_i^k \delta_j^l + \delta_i^l \delta_j^k) \delta^{(3)}(\vec{x}, \vec{y})$$

- quantize $\mathcal{H}_\mu = (\mathcal{H}_\perp, \mathcal{H}_i)$ and demand $\hat{\mathcal{H}}$

quantum level,

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- find a scalar on $\{\psi^{\text{phys}}\}$ to obtain $\mathcal{H}^{\text{phys}}$

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$$(\#_1) \hat{\mathcal{L}}_i \Psi = -2 D_i h_{ik} (-i\hbar) \frac{\delta}{\delta h_{kj}(x)} \Psi = 0$$

$$(\#_2) \hat{\mathcal{L}}_L \Psi = \left(-\kappa^2 \hbar^2 g_{ijkl} \frac{\delta^2}{\delta h_{ij} \delta h_{kl}} - \frac{\sqrt{h}}{\kappa^2} {}^{(3)}R \right) \Psi = 0 \quad \text{"W" D}$$

- find a scalar on ξ^{μ} phys

$$(-i\hbar) \frac{\delta}{\delta h_{\mu\nu}(x)} \Psi = 0$$

$$\delta_{ijkl} \frac{\delta^2}{\delta h_{ij} \delta h_{kl}} - \frac{\sqrt{h}}{\kappa^2} {}^{(3)}R) \Psi = 0 \quad \text{"wheeler-DeWitt equation"} \\ \text{(WdW-equ.)}$$

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(WdW-eqn.)

$$(\#_1) \hat{\mathcal{L}}_i \Psi = -2 D_j h_{ik} (-i\hbar) \frac{\delta}{\delta h_{kj}(x)} \Psi = 0$$

$$(\#_2) \hat{\mathcal{L}}_\perp \Psi = \left(-\kappa^2 \hbar^2 g_{ijkl} \frac{\delta^2}{\delta h_{ij} \delta h_{kl}} - \frac{\sqrt{h}}{\kappa^2} {}^{(3)}R \right) \Psi = 0$$

"wheeler itt eq
(Wdl i.)

$(\#_1)$: $\hat{\mathcal{L}}_i$ generate spatial diffeomorphisms

$\hat{\mathcal{L}}_i \Psi = 0$ is $\Psi[h_{ij}]$ is a diffeomorphism-invariant function

find a solution on $\{q(t)\}$ to obtain $\alpha(t)$

(#2) α must be regularized; this cannot be separated from the factor-ordering issue

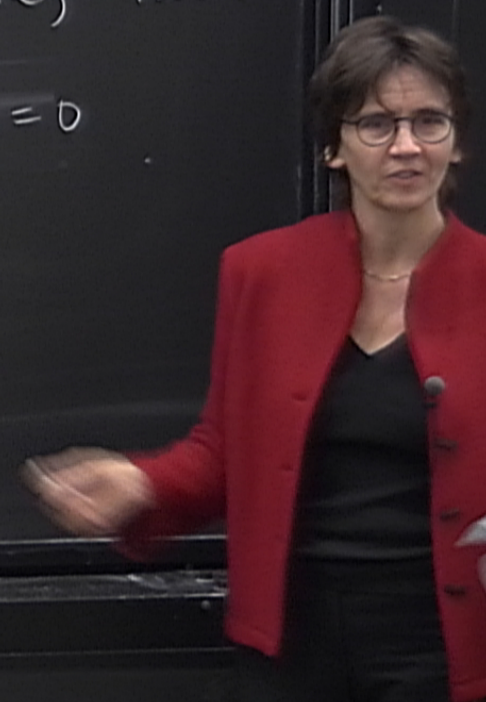
" Wheeler-DeWitt equation"
(WdW-equ.)

• don't know how to solve $\hat{H}_\perp \psi = 0$

→ interpretation?

$$i\hbar \frac{d}{dt} \psi(t) = \hat{H}[N, \dot{N}](t) \psi(t) = 0$$

anal of h_{ij}



find a state on $\{q, p\}$ to obtain α

(#2) \therefore must be regularized; this cannot be separated from the factor-ordering issue

" Wheeler-DeWitt equation"

(WdW-equ.)

\therefore don't know how to solve $\hat{H}_\perp \psi = 0$

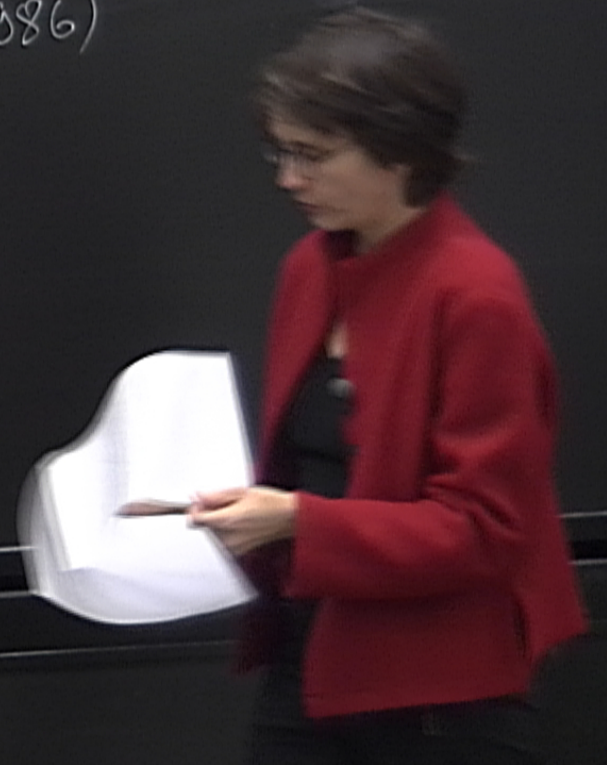
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anal of h_{ij}

Isham, $g_r/g_c - 9210011$

\rightarrow new connection variables (Ashtekar 1986)



Isham, gr/qc - 9210011

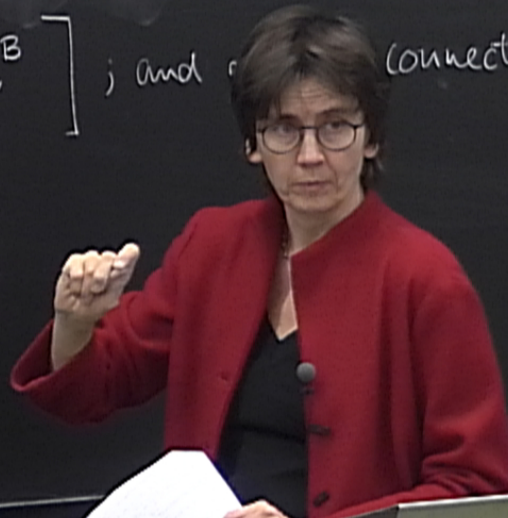
→ new connection variables (Ashtekar 1986)

→ loop quantum gravity



Classically equivalent first-order formulation of GR in terms of an orthonormal frame (tetrad) field $e_\mu^A(x)$, $A=0, \dots, 3$, $\mu=0, \dots, 3$, $A \sim$ internal $so(3,1)$ -index ; $\left[g_{\mu\nu} = \underset{\substack{\uparrow \\ \text{int. Mink.} \\ \text{metric}}}{\eta_{AB}} e_\mu^A e_\nu^B \right]$; and a spin connection $\omega_\mu^A{}_B(x)$

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\uparrow
int. Mink.
metric

\Rightarrow perform a 3+1 decomposition, and partial gauge-fixing $SO(3,1) \rightarrow SO(3)$

Yang-Mills, conjugate variable pairs

$$\{A_i^a(\vec{x}), E_b^i(\vec{y})\} = 8\pi\beta \delta^a_b \delta_i^j \delta^{(3)}(\vec{x}, \vec{y}), \quad \begin{array}{l} i, j = 1, 2, 3 \text{ spatial index} \\ a, b = 1, 2, 3 \text{ so(3)-index} \end{array}$$

↑
Barbero-Immirzi
parameter

$A \ni A_i^a \sim \text{so(3)-connections}$

Yang-Mills, conjugate variable pairs

$$\{A_i^a(\vec{x}), E_b^j(\vec{y})\} = 8\pi\beta \delta^a_b \delta_i^j \delta^{(3)}(\vec{x}, \vec{y})$$

\uparrow
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parameter

$i, j = 1, 2, 3$ spatial index
 $a, b = 1, 2, 3$ $so(3)$ -index

$A \ni A_i^a \sim so(3)$ -connections, $E_b^j = \sqrt{h} e_b^j$ "densitized triad"

\uparrow associated curvature $F_{ij}^a = 2G_N \partial_{[i} A_{j]}^a + G_N^2 \epsilon_{abc} A_i^b A_j^c$