

Title: Condensed Matter (Review) - Lecture 13

Date: Jan 18, 2012 10:15 AM

URL: <http://pirsa.org/12010099>

Abstract:





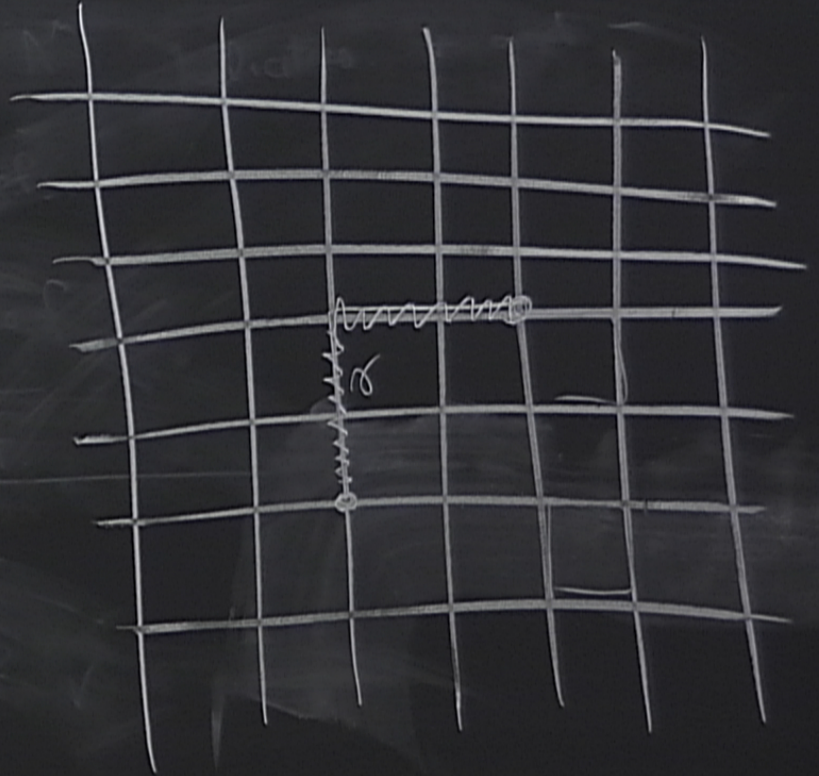
## $\mathbb{Z}_2$ Topological Order

$$H_{TC} = -U \sum_s A_s - J \sum_p B_p$$

$$\begin{cases} A_s = \prod_{j \in s} \sigma_j^x \\ B_p = \prod_{j \in p} \sigma_j^z \end{cases}$$

$$W_\gamma^z = \prod_{j \in \gamma} \sigma_j^z$$

$B_p$  is such a string operator





# $Z_2$ Topological Order

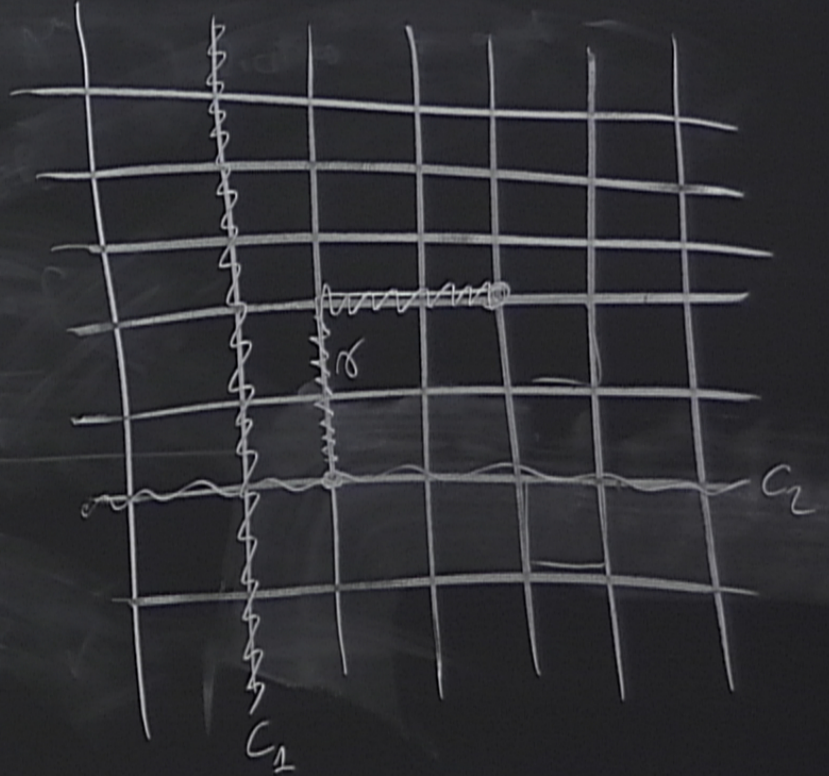
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$$W_1^x = \prod_{j \in C_1} \sigma_j^x$$





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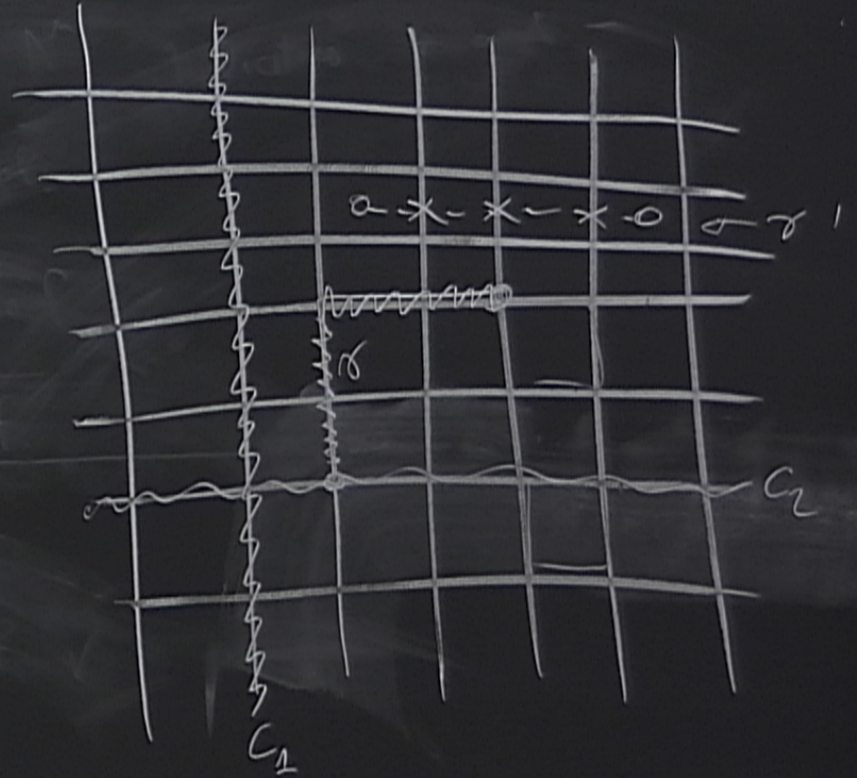
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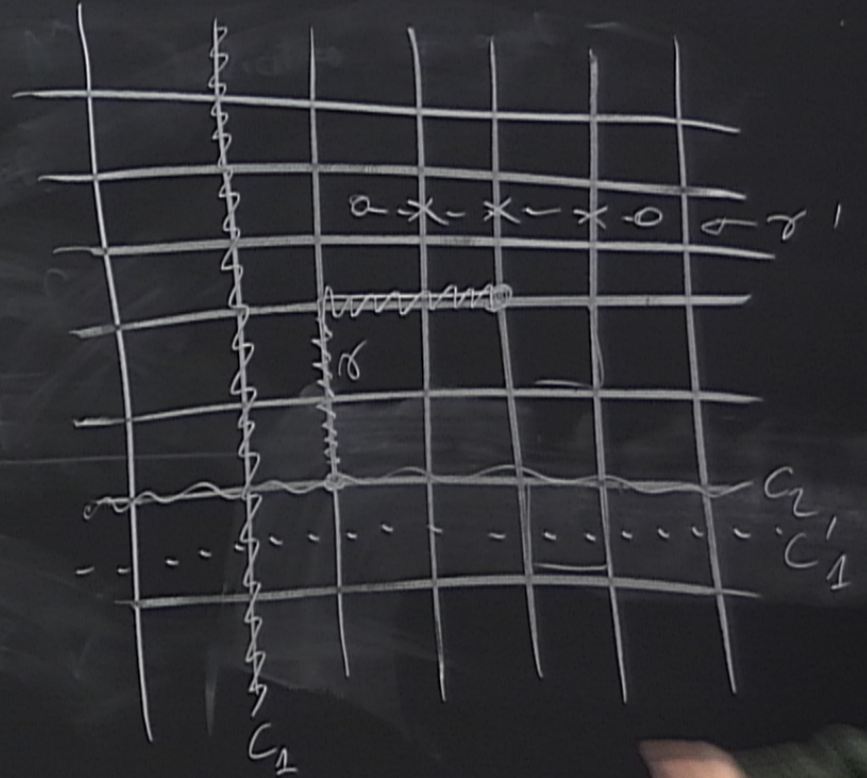
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# $\mathbb{Z}_2$ Topological Order

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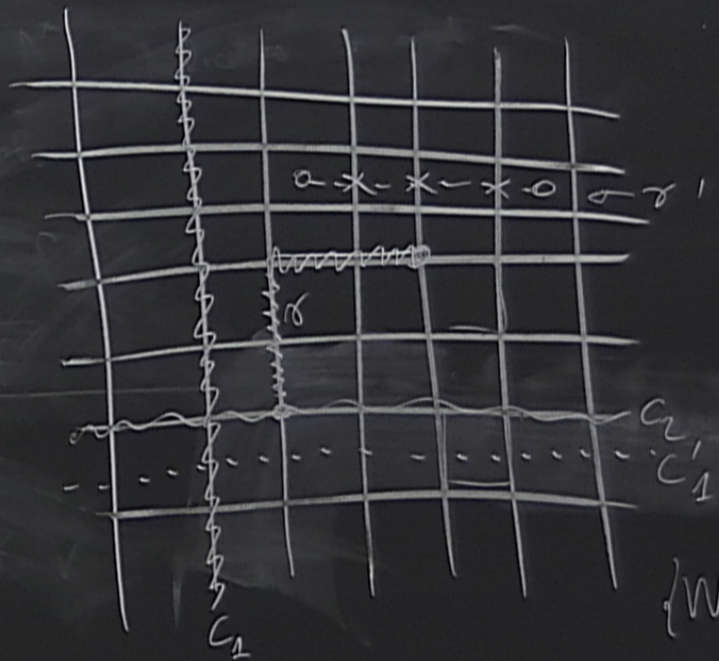
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$$W_{\gamma'}^x = \prod_{j \in \gamma'} \sigma_j^x$$



$$\{W_{C_1}^z, W_{C_1'}^z\}$$



$$G = \langle A_{1,1}, \dots, A_{L^2-1} \rangle$$

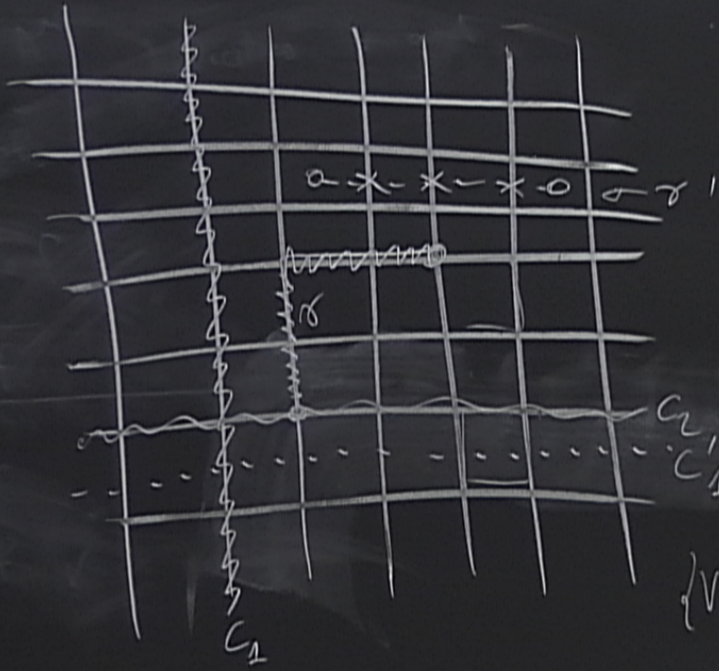
$$|G| = 2^{L^2-1}$$

$$\mathcal{L} = \text{span} \{ \phi_0, W_1 \phi_0, W_2 \phi_0, W_3 \phi_0, W_4 \phi_0 \}$$

$$\phi_0 = \frac{1}{\sqrt{|G|}} \sum_{g \in G} |g\rangle$$

$$|\uparrow\uparrow\uparrow\uparrow \dots \uparrow\rangle$$

$$\{W_1^x, W_1^z\}$$





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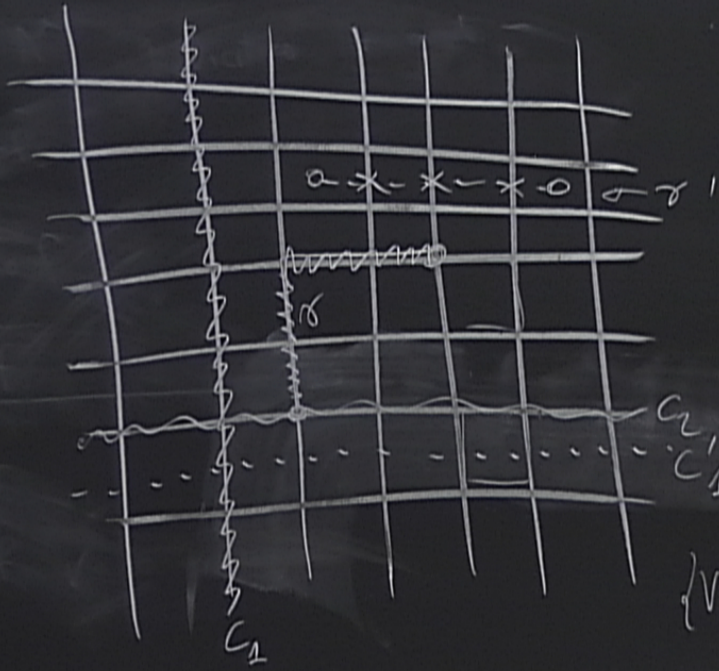
$$\mathcal{L} = \text{span} \{ \phi_0, W_1 \phi_0, W_2 \phi_0, \dots, W_{L^2-1} \phi_0 \}$$

$$\phi_0 = \frac{1}{\sqrt{|G|}} \sum_{g \in G} |g\rangle$$

$$\psi \in \mathcal{L} \Rightarrow g\psi = \psi$$

$| \uparrow \uparrow \uparrow \dots \uparrow \rangle$

$$\{ W_1^x, W_1^z \}$$





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$$\phi_0 = \frac{1}{\sqrt{|G|}} \sum_{g \in G} g |\delta\rangle$$

$$\psi \in \mathcal{L} \Rightarrow \boxed{g\psi = \psi}$$

$(\uparrow \uparrow \uparrow \dots \uparrow)$

$$\Rightarrow \begin{aligned} \phi_0 &\rightarrow \int_A \\ W_1 \phi_0 &\rightarrow \int_A' \end{aligned}$$

$$\int_A = \int_A'$$

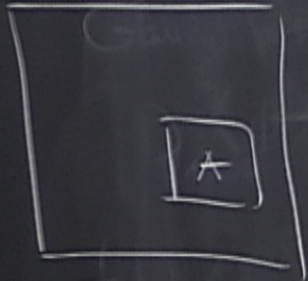
A is a subregion

$$\{W_1^x, W_1^z\}$$



Entanglement in T.O. ground state

$$\phi_0 = |G|^{-1/2} \sum_{j \in G} |j\rangle$$

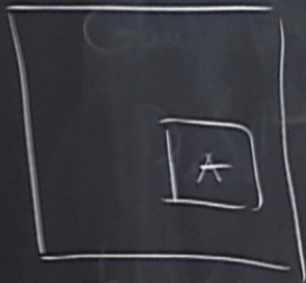


$$\rho_A \rightarrow S(\rho_A) = -\text{Tr} \rho_A \log \rho_A$$
$$= -\sum_i \lambda_i \log \lambda_i$$



Entanglement in T.O. ground state

$$\phi_0 = |G|^{-1/2} \sum_{g \in G} |g\rangle$$



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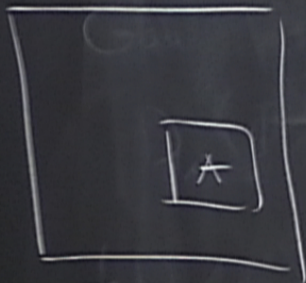
$$= -\sum_i \lambda_i \log \lambda_i \sim \alpha |A| + \gamma$$

non-universal



Entanglement in T.O. ground state

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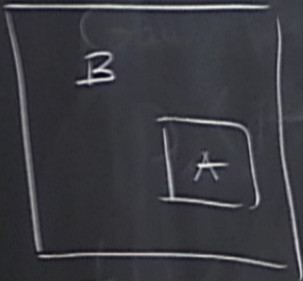


Entanglement in T.O. ground state

$$\phi_0 = |G|^{-1/2} \sum_{g \in G} g |\bar{0}\rangle$$

$$G_A = \{g \in G \mid g = g_A \otimes 1_B\}$$

$$G_B$$



$$S_A \rightarrow S(\rho_A) = -\text{Tr} \rho_A \log \rho_A$$

$$= -\sum_i \lambda_i \log \lambda_i \sim \alpha |\partial A| + \gamma$$

$S$  is invariant under local unitaries

non-universal

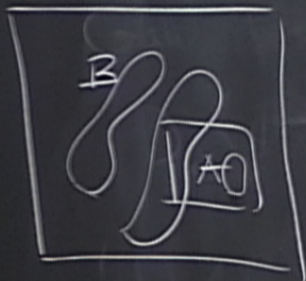


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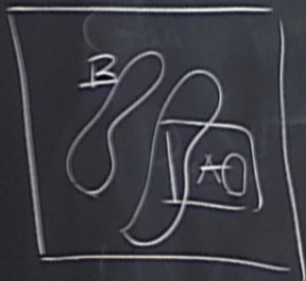
Entanglement in T.O. ground state

$$\phi_0 = |G|^{-1/2} \sum_{g \in G} |g\rangle$$

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$$G \neq G_A \times G_B$$



$$S_A \rightarrow S(\rho_A) = -\text{Tr} \rho_A \log \rho_A$$

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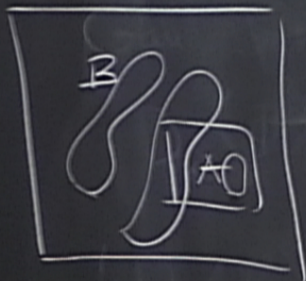
Entanglement in T.O. ground state

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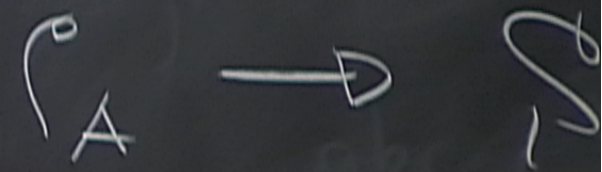
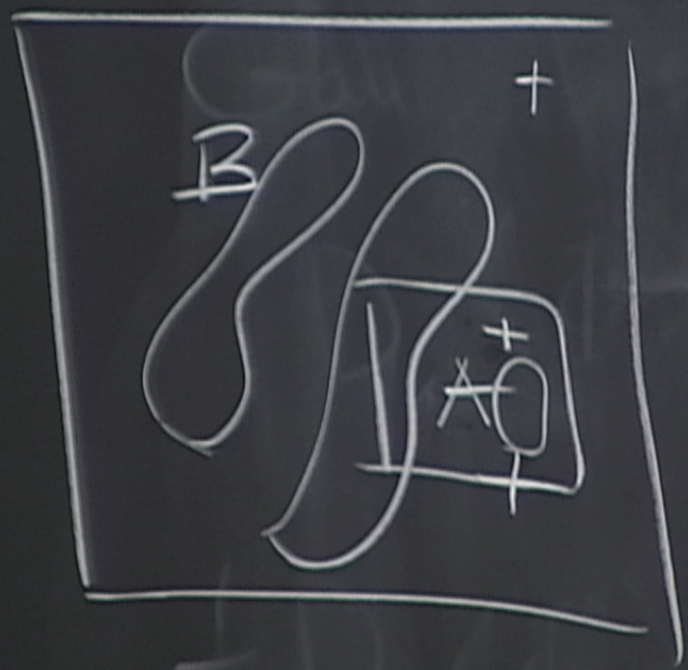
$$= -\sum_i \lambda_i \log \lambda_i \sim \alpha |\partial A| + \gamma$$

$S$  is invariant under local unitaries

non-universal



$$\phi_0 = |G|^2 \int_G$$



$\int_S$  is invariant  
local th



$$S = \log \frac{|G|}{|G_A| |G_B|} =$$

$\log |G| \sim$  number of A operators

$\log |G_A| + \log |G_B| \sim$  number of A operators that just in A or B

$$G_{AB} = \frac{G}{G_A \times G_B}$$

$\log |G_{AB}|$  is the # of A operators near



$$S = \log \frac{|G|}{|G_A| |G_B|} =$$

$\log |G| \sim$  number of A operators

$\log |G_A|, |G_B| \sim$  number of A operators that just in A or B

$$G_{AB} = \frac{G}{G_A \times G_B}$$

$\log |G_{AB}| \rightarrow$  the # of A operators near the boundary



$$S = \log \frac{|G|}{|G_A| |G_B|} = \log_2 |Z_A| - \frac{1}{2} \log_2 |Z| = |Z_2| \text{ min}$$

$d=1$

$\log |G| \sim$  number of A operators

$\log |G_A| |G_B| \sim$  number of A operators that just in A or B

$$G_{AB} = G$$



$\{W_1, W_2\}$

$$\int_A = \int_A'$$

$$S = \log \frac{|G|}{|G_A| |G_B|} = |2A| - \gamma$$

$$\log_2 |Z| = |Z_2|$$

$\log |G| \sim$  number of A operators

D dimension of local gauge group

$\log |G_A| |G_B| \sim$  number of A operators that just in A or B

$$\gamma = -\log D$$

$$S = D^\gamma$$

$$G_{AB} = \frac{G}{G_A \times G_B}$$

$\log |G_{AB}| \sim$  the # of A operators near the boundary



$Z_2$   $|0\rangle, |1\rangle$

$$\sigma^x |0\rangle = |1\rangle, \sigma^x |1\rangle = |0\rangle$$

$$L_g^+ |\tilde{g}\rangle = |g\tilde{g}\rangle$$
$$L_g^- |\tilde{g}\rangle = |\tilde{g}g^{-1}\rangle$$

D dimensional

$$|0\rangle = |e\rangle$$
$$|1\rangle = |g_1\rangle$$
$$\vdots$$
$$|D-1\rangle = |g_{D-1}\rangle$$



$Z_2$   $|0\rangle, |1\rangle$

$$\sigma^x |0\rangle \equiv |1\rangle \equiv |e\rangle$$
$$\sigma^x |1\rangle \equiv |0\rangle$$

D dimensional

$$|0\rangle \equiv |e\rangle$$
$$|1\rangle \equiv |g_1\rangle$$
$$\vdots$$
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$$L_g^+ |\tilde{g}\rangle = |g\tilde{g}\rangle$$
$$L_g^- |\tilde{g}\rangle = |\tilde{g}g^{-1}\rangle$$

$$T_h^+ |g\rangle = \int_{g,h} |g\rangle$$
$$T_h^- |g\rangle = \int_{g,h^{-1}} |g\rangle$$



$$A_g(s) = \prod_{j \in S} L_j^+$$

$$A(s) = \bar{Z}_g A_g(s)$$

$$h_4 \left[ \begin{array}{c} h_3 \\ p \\ h_2 \end{array} \right] h_1$$

$$B_p = \sum_{h_1 h_2 h_3 h_4 = e} T_{h_1} T_{h_2} T_{h_3} T_{h_4}$$

$$H = \bar{Z}_p (1 - B_p) + \bar{Z}_s (1 - A_s)$$



$\log |G_{AB}|$  is the # of A operators near the boundary

$$A(s) = \prod_{j \in s} L_j^+ \quad A(s) = \bar{\sum}_j A_j(s)$$

$$H = \bar{\sum}_p (1 - B_p) + \bar{\sum}_s (1 - A_s) \quad \text{Quantum Double}$$

$$B_p = \sum_{h_1, h_2, h_3, h_4 = e} T_{h_1} T_{h_2} T_{h_3} T_{h_4}$$

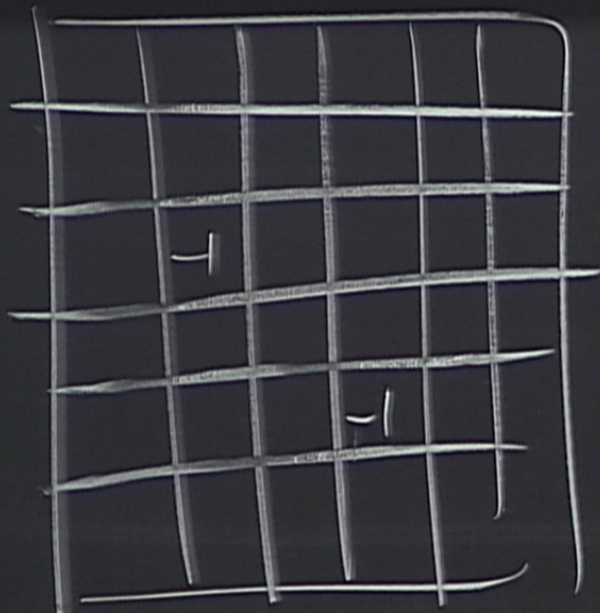
$$\gamma = \log D$$





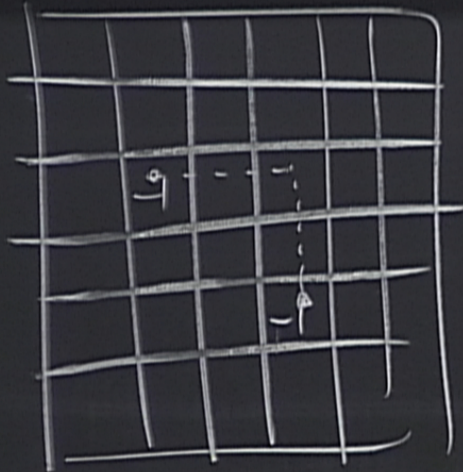


$\phi_0$



$$\frac{l^2 (l^2 - 1)}{2}$$

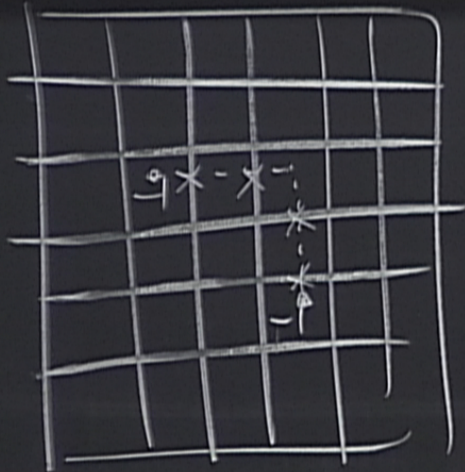




$\frac{L^2(L^2-1)}{2}$  degeneracy of 1st excited state

$$\prod_p B_p = +1$$





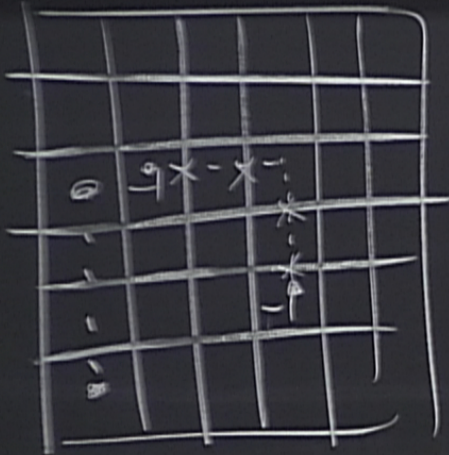
$\frac{L^2(L^2-1)}{2}$  degeneracy of 1st excited state

$$\prod_p B_p = +1$$

$$N_z = \prod_{j \in z} \sigma_j^z$$

$\prod_j |\psi_0\rangle$  is an eigenstate of  $H$  w/ energy





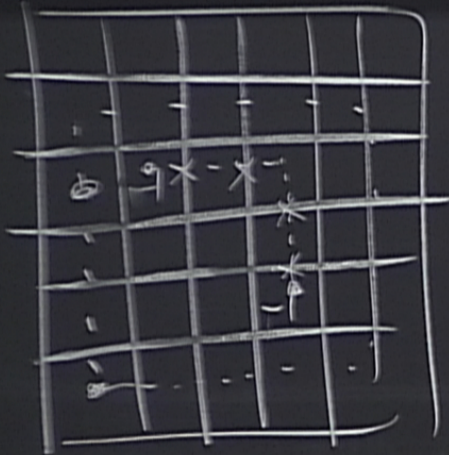
$4 \times \frac{L(L-1)}{2}$  degeneracy of 1st excited state

$$\prod_p B_p = +1$$

$$N_z = \prod_{j \in \mathbb{Z}_2} \sigma_j^z$$

$W_\sigma |\phi_0\rangle$  is an eigenstate of  $H$  w/ energy  $4J$





$4 \times \frac{L^2(L^2-1)}{2}$  degeneracy of 1st excited state

$$g|4\rangle = |4\rangle$$

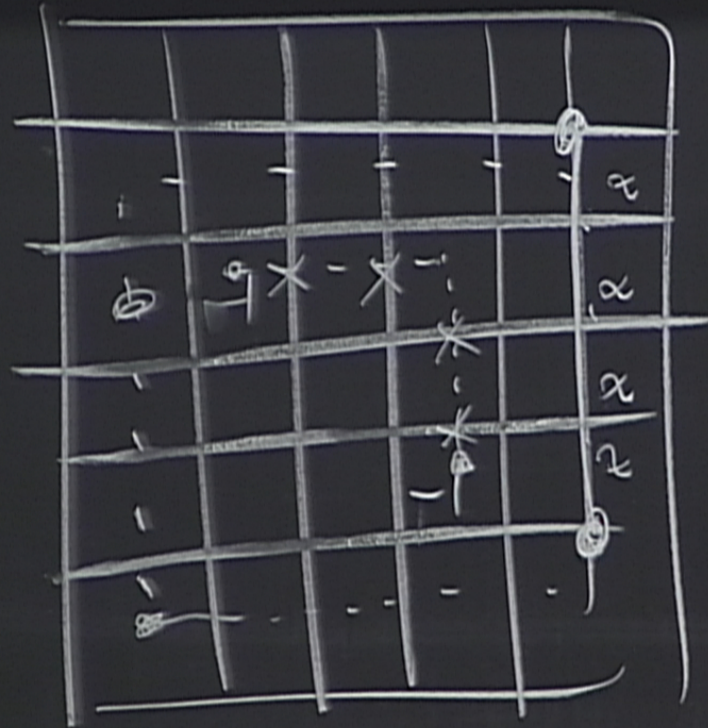
$$\prod_p B_p = +1$$

$$N_z = \sum_{j \in z} \sigma_j^z$$

$g W_\sigma |\phi_0\rangle$  is an eigenstate of  $H$  w/ energy  $4J$



$\phi_0$



$$4 \times \frac{l^2(l^2-1)}{2} \text{ deg}$$

$$\prod_p B_p = +1$$

$W = \pi^2 \dots$   $W | \phi \rangle$  is



$$\begin{array}{|c|c|c|c|c|} \hline 0 & * & - * & * & - * & * & 0 \\ \hline \end{array}$$

degeneracy of 1st excited state

$$g |4\rangle = |4\rangle$$

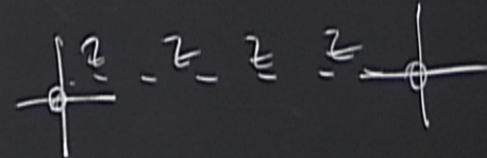
+1

is an eigenstate of  $H$  w/ energy  $4J$





degeneracy of 1st excited state

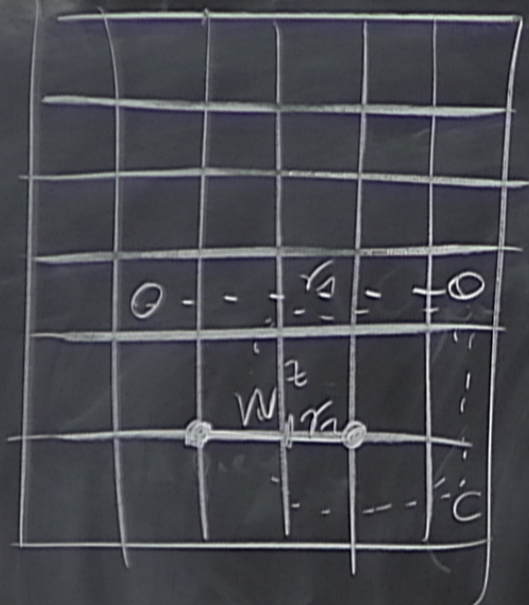


$$g(|4\rangle) = |4\rangle$$

+1

is an eigenstate of  $H$  w/ energy  $4J$

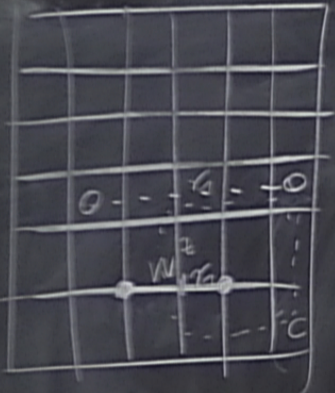




$$\psi_{in} = W_{\gamma_2}^z W_{\gamma_1}^x |\phi_0\rangle$$

$$\psi_{fin} = W_c^x \psi_{in} = W_c^x W_{\gamma_2}^z W_{\gamma_1}^x |\phi_0\rangle$$

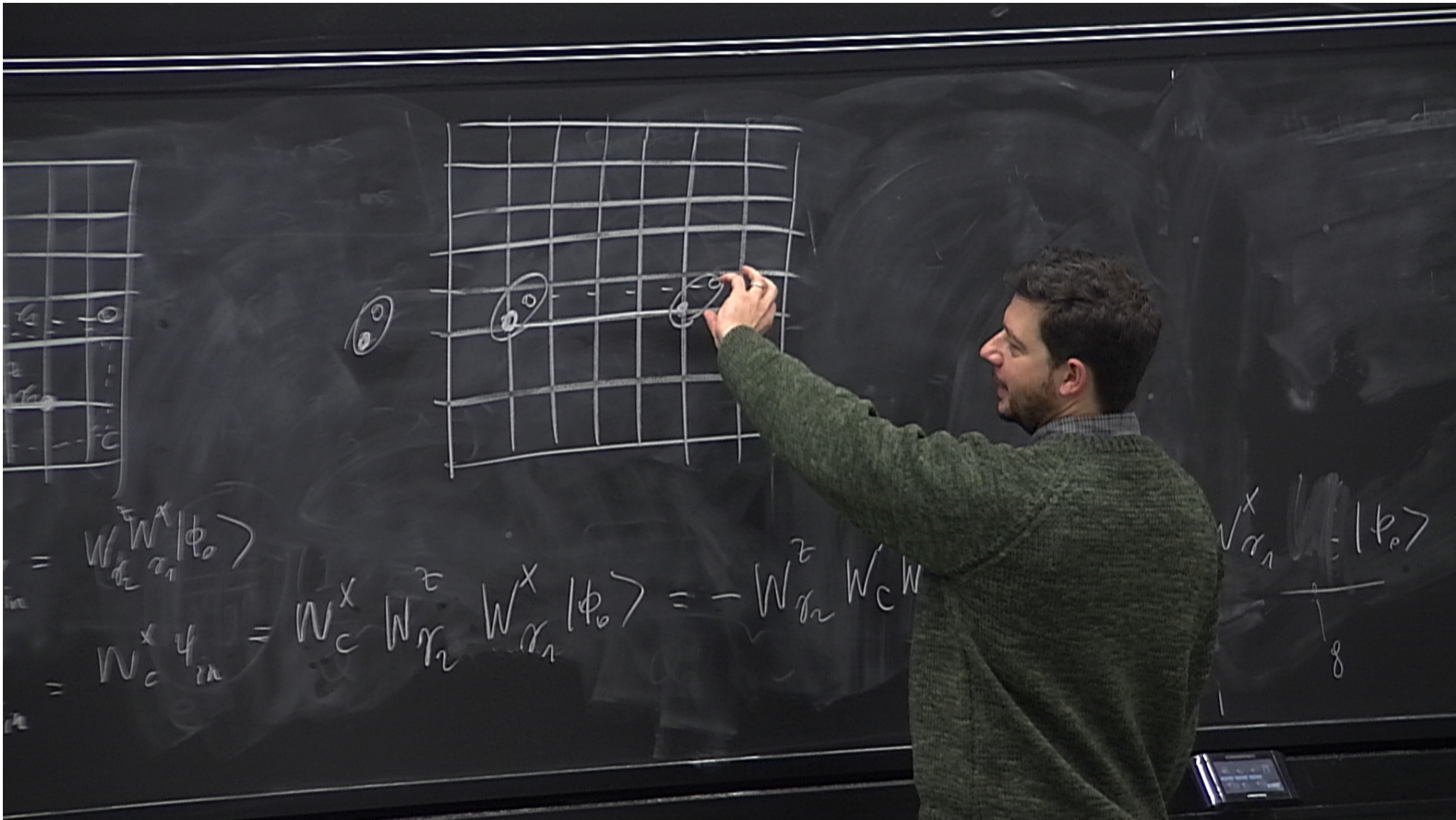




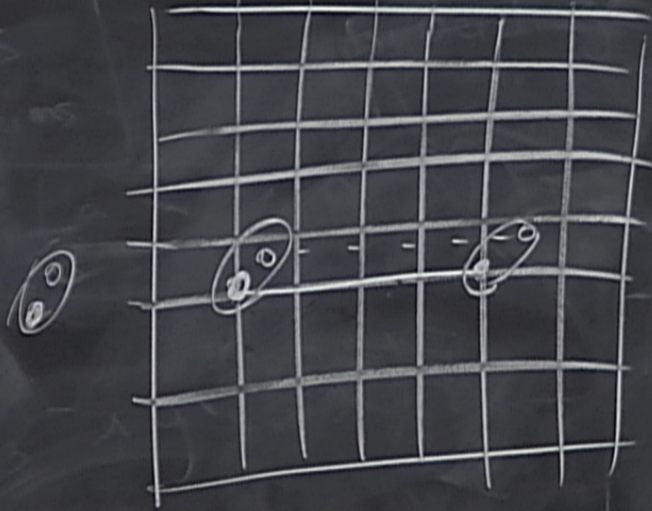
$$\psi_{in} = W_{\alpha_2}^z W_{\alpha_1}^x |\phi_0\rangle$$

$$\psi_{fin} = W_c^x \psi_{in} = W_c^x W_{\alpha_2}^z W_{\alpha_1}^x |\phi_0\rangle = -W_{\alpha_2}^z W_c^x W_{\alpha_1}^x |\phi_0\rangle = -W_{\alpha_2}^z W_{\alpha_1}^x W_c^x |\phi_0\rangle = -\psi_{in}$$









Emergence of fermions

$$\begin{aligned}
 &= W_{\alpha_2}^z W_{\alpha_1}^x |\phi_0\rangle \\
 W_c^x \psi_m &= W_c^x W_{\alpha_2}^z W_{\alpha_1}^x |\phi_0\rangle = -W_{\alpha_2}^z W_c^x W_{\alpha_1}^x |\phi_0\rangle = -W_{\alpha_2}^z W_{\alpha_1}^x W_c^x |\phi_0\rangle
 \end{aligned}$$