

Title: Quantum Gravity (Review) - Lecture 7

Date: Jan 31, 2012 10:15 AM

URL: <http://pirsa.org/12010065>

Abstract:

dynamics on phase space (a "symplectic manifold" (\mathcal{P}, ω))

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standard QM on \mathbb{R}^n : unique representation of Heisenberg algebra

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dynamics on phase space (a "symplectic manifold" (\mathcal{S}, ω))

standard QM on \mathbb{R}^n : unique representation of Heisenberg algebra of $q_i, p_i, 1$

$$\{q_i, p_j\} = \delta_{ij}, \text{ and add a Hamiltonian } H \xrightarrow{?} \hat{H}$$

(a "symplectic manifold" (\mathcal{P}, ω))

the representation of Heisenberg algebra of $q, p; \hbar$

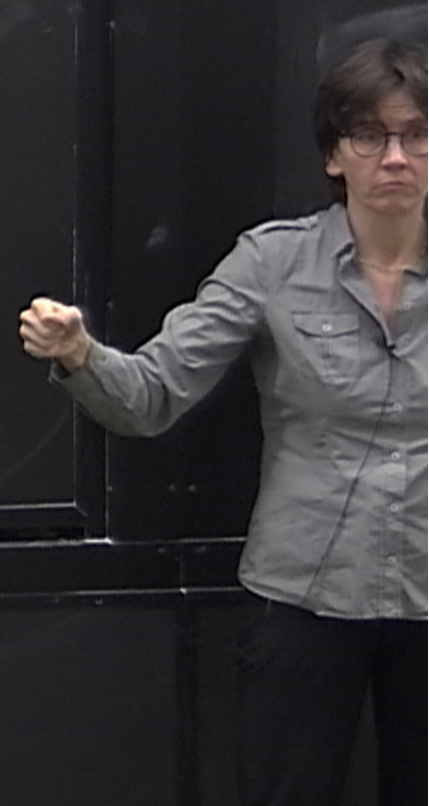
and a Hamiltonian $H \mapsto \hat{H}$

Constrained Hamiltonian S

o))

g algebra of $q, p; 1$

Constrained Hamiltonian systems



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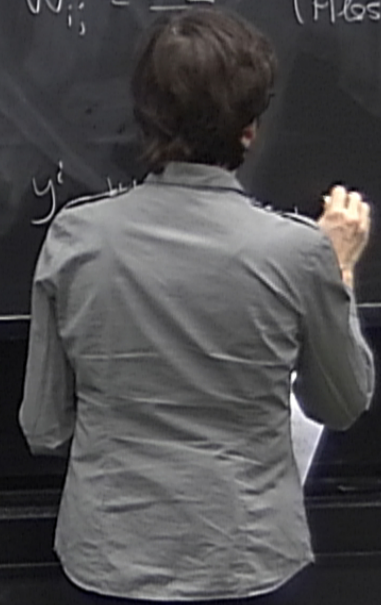
g algebra of $q_i, p_i, 1$

Constrained Hamiltonian systems

A n -dimensional Lagrangian system is called singular if $0 = \det W_{ij}$, $W_{ij} = \frac{\partial^2 L}{\partial \dot{q}^i \partial \dot{q}^j}$ (Hessian)

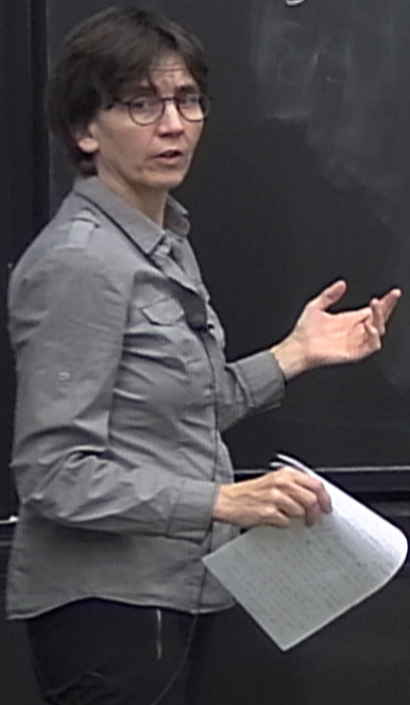
$\text{rank}(W) = n - m \Rightarrow$

\exists m zero-eigenvectors



\Rightarrow Legendre transform $(q_i, \dot{q}_i) \mapsto (q_i, p_i)$ is

Singular $|W_{ij}| = \left| \frac{\partial^2 L}{\partial \dot{q}_i \partial \dot{q}_j} \right| = \left| \frac{\partial p_j}{\partial \dot{q}_i} \right| = 0$



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can't solve for all $\dot{q}_i = \dot{q}_i(q_i, p_i) \Rightarrow$

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can't solve for all $\dot{q}_i = \dot{q}_i(q_i, p_i) \Rightarrow$

\Rightarrow have m relations $\phi_a(q_i, p_i) = 0$, $a = 1, \dots, m$

\Rightarrow so-called "constraints"

form $(q_i, \dot{q}_i) \mapsto (q_i, p_i)$ is

$$= \left| \frac{\partial^2 L}{\partial \dot{q}_i \partial \dot{q}_i} \right| = \left| \frac{\partial p_i}{\partial \dot{q}_i} \right| = 0$$

$\Rightarrow (2n-m)$ -dim. subspace, "constraint surface"

ex: $L = \frac{1}{2}$

for all $\dot{q}_i = \dot{q}_i(q, p) \Rightarrow$

as $\phi_a(q, p) = 0, a = 1, \dots, m$
 "constraints" (assumed indept.)

is

$\Rightarrow (2n-m)$ -dim. subspace, "constraint surface" $f_c \subset f$

ex: $L = \frac{1}{2} \dot{q}_1^2 + c \dot{q}_2 q_1 - \frac{m}{2} (q_1^2 + q_2^2 + q_3^2)$ ($n=3$)

$$\frac{\partial L}{\partial q_1} = \dot{q}_1, \quad p_2 = \frac{\partial L}{\partial \dot{q}_2} = c q_1, \quad p_3 = \frac{\partial L}{\partial \dot{q}_3} = 0$$

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$$W_{ij} = \begin{pmatrix} & & \\ & & \\ & & \end{pmatrix}$$

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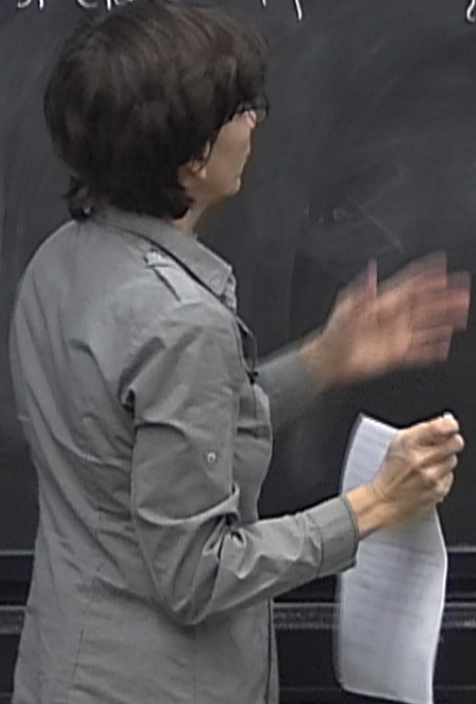
$$W_{ij} = \begin{pmatrix} 1 & & \\ & 0 & \\ & & 0 \end{pmatrix}, \quad \text{rank } W_{ij} = 1$$

constraints are $\phi_1 := p_2 - c q_1 = 0$, $\phi_2 := p_3 = 0 \Rightarrow \mathcal{F}_c$ is 4-dim.

constraints (assumed indept.)

constraints

A Hamiltonian system with m constraints $\phi_a(q, p) = 0$ is called "first class" if
$$\{\phi_a, \phi_b\} = f_{ab}^c \phi_c$$



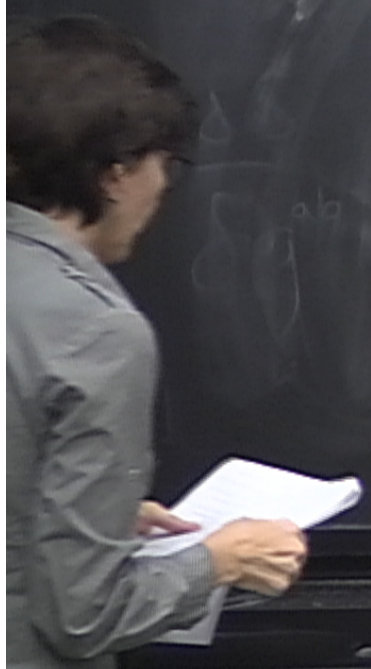
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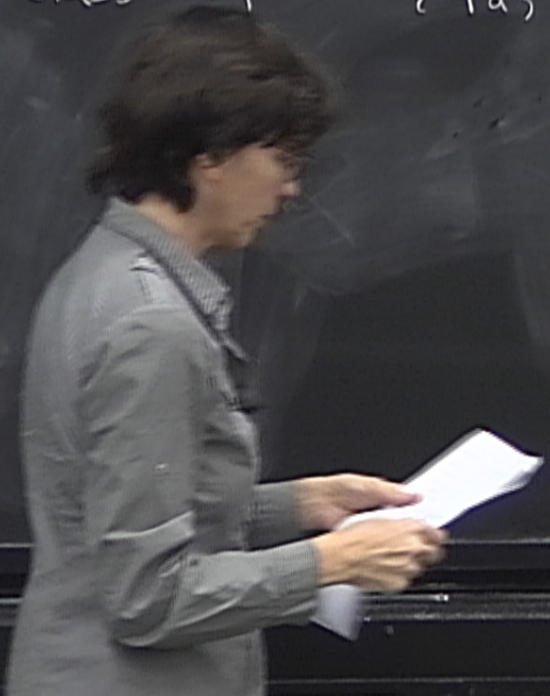
constraints (assumed indept.)

constraints a

A Hamiltonian system with m constraints $\phi_a(q,p) = 0$ is called "first class" if $\{\phi_a, \phi_b\} = f_{ab}^c \phi_c$ and $\{\phi_a, H\} = d_a^b \phi_b$

↑ needn't be constants

↑ Hamiltonian



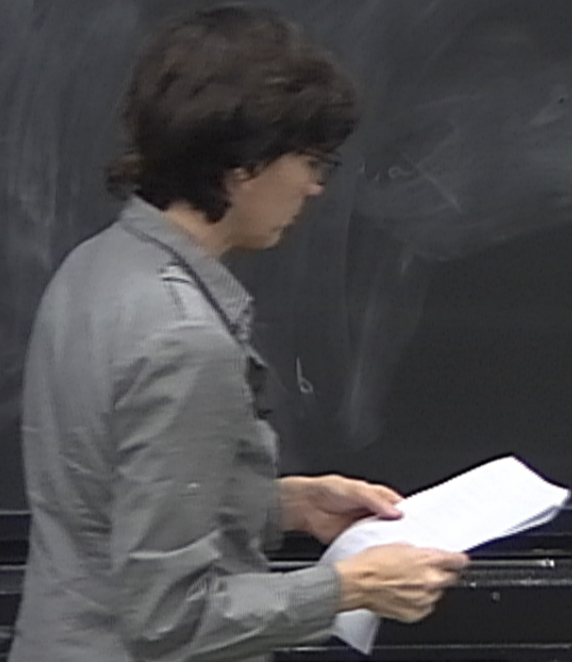
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constraints a

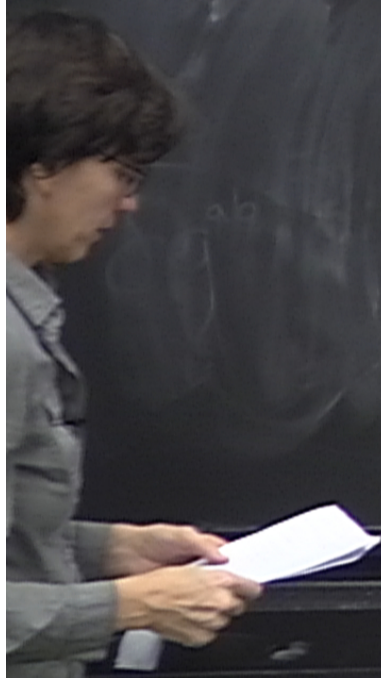
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$$\{\phi_a, \phi_b\} = f_{ab}^c \phi_c \quad (*)$$

and $\{\phi_a, H\} = d_a^b \phi_b \quad (**)$

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(*) needn't be constants

(**) Hamiltonian

[notation: " \approx " means "equality on \mathcal{J} "]

$$\{\phi_a, \phi_b\} \approx 0$$

constraints (assumed indept.)

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↑
needn't be constants

↑
Hamiltonian

[notation: " \approx " means "equality on \mathcal{P}_c "]

$$\{\phi_a, \phi_b\} \approx 0$$

constraints are $\Phi_1 := p_2 - cq_1 = 0$, $\Phi_2 := p_3 = 0 \Rightarrow J_c$ is 4-dim.

Hamiltonian is not unique $H^* = H + c_m(q,p) \Phi_m(q,p)$

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(**)
 $\{\Phi_a, H\} = d_a^b \Phi_b$
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$\dot{q} = \{q, H^*\}$

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\dot{q}_i

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$$\dot{q} = \{q, H^*\} = \{q, H\} + c_m \{q, \Phi_m\} + \underbrace{\Phi_m \{q, c_m\}}_{\approx 0}$$

$$\dot{q}_i \approx \frac{\partial H}{\partial p_i} + c_m \frac{\partial \Phi_m}{\partial p_i}, \quad \dot{p}_i \approx -\frac{\partial H}{\partial q_i} - c_m \frac{\partial \Phi_m}{\partial q_i}$$

\Rightarrow e.o.m. contain m arbitrary phase space functions c_m .

Consistency: constraints must be preserved under time evolution.

$$\dot{\Phi}_m \approx 0 \Leftrightarrow \{\Phi_m, H\} +$$

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$$\dot{\phi}_m \approx 0 \Leftrightarrow \left\{ \phi_m, H \right\} + c_n \left\{ \phi_m, \phi_n \right\} \approx 0$$

\uparrow
 $(*) \approx 0$ $(*) \approx 0$

Consistency: constraints must be preserved under time evolution.

$$\dot{\Phi}_m \approx 0 \Leftrightarrow \underbrace{\{\Phi_m, H\}}_{(*) \approx 0} + c_n \underbrace{\{\Phi_m, \Phi_n\}}_{(*) \approx 0} \approx 0$$

$$\{\Phi_m, H\} \approx 0$$

(i)

Consistency: constraints must be preserved under time evolution.

$$\dot{\phi}_m \approx 0 \Leftrightarrow \begin{cases} \{\phi_m, H\} + c_n \{\phi_m, \phi_n\} \approx 0 \\ \uparrow \\ (*, *) \approx 0 \quad (*) \approx 0 \end{cases}$$

$$\{\phi_m, H\} \approx 0$$

(i) flow of Hamiltonian stays on \mathcal{P}_c

$$\{\phi_m, H\} = \sum_i \left(-\frac{\partial H}{\partial q_i} \frac{\partial \phi_m}{\partial p_i} + \frac{\partial H}{\partial p_i} \frac{\partial \phi_m}{\partial q_i} \right)$$

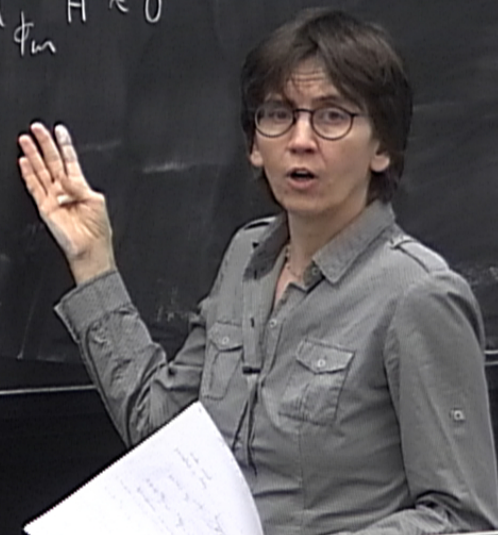
$$=: X_H \phi_m \approx 0$$

(ii) H is unchanged when m

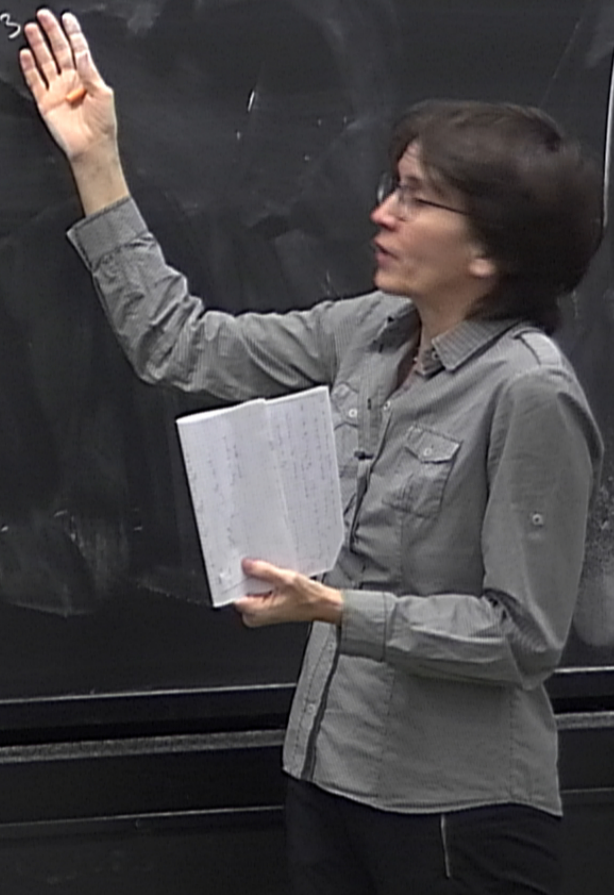
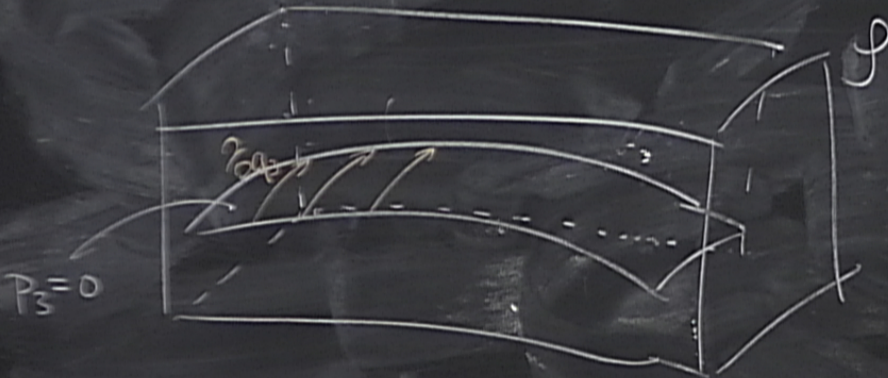
under time evolution.

(ii) H is unchanged when moving along the vector field X_{ϕ_m} on \mathcal{J}_c

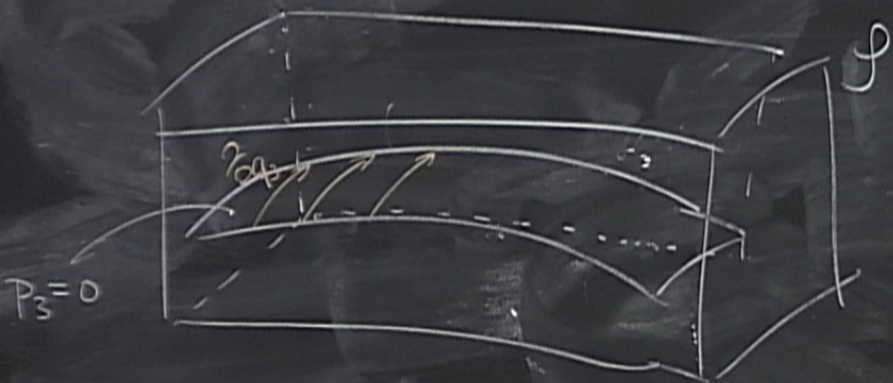
$$\{\phi_m, H\} = \sum_i \left(\frac{\partial \phi_m}{\partial q_i} \frac{\partial}{\partial p_i} - \frac{\partial \phi_m}{\partial p_i} \frac{\partial}{\partial q_i} \right) H = X_{\phi_m} H \approx 0$$



Ex. $\phi = p_3 = 0$, $X_{p_3} = -\frac{\partial}{\partial q_3} \Rightarrow \frac{\partial}{\partial q_3} H \approx 0$



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