

Title: Foundations of Quantum Mechanics - Lecture 10

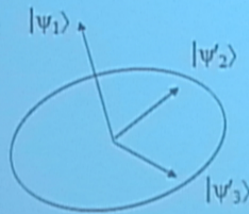
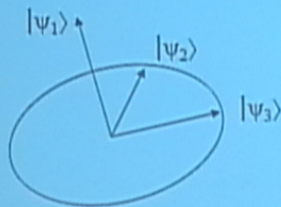
Date: Jan 13, 2012 11:30 AM

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Abstract:

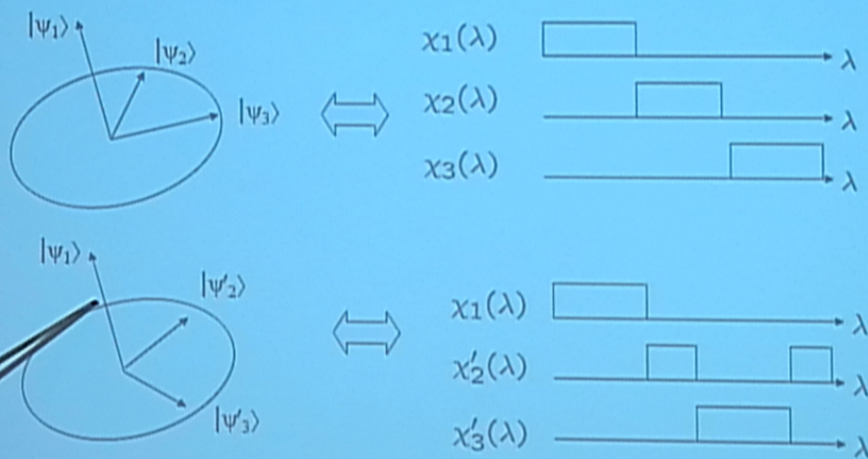
Traditional notion of noncontextuality

A given vector may appear in many different measurements



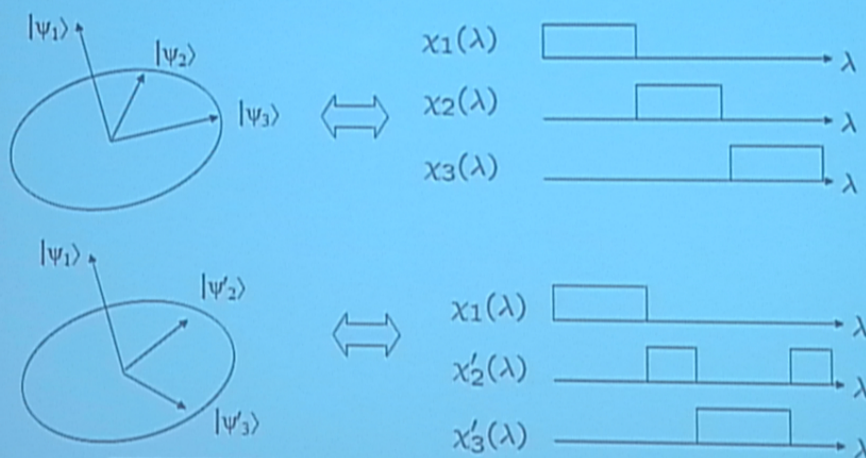
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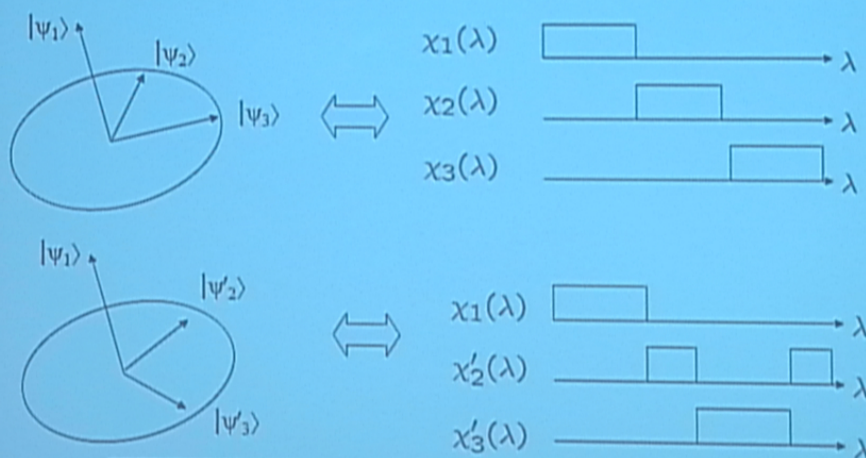
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The traditional notion of noncontextuality:
Every vector is associated with the same $\chi(\lambda)$
regardless of how it is measured (i.e. the context)

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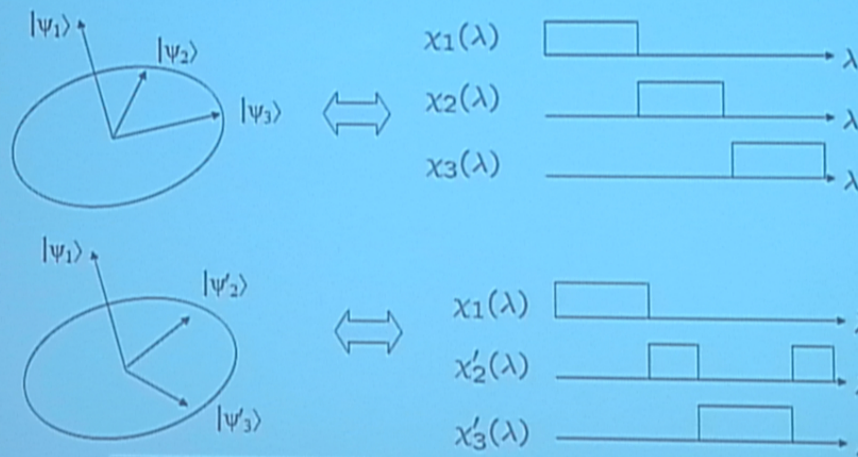
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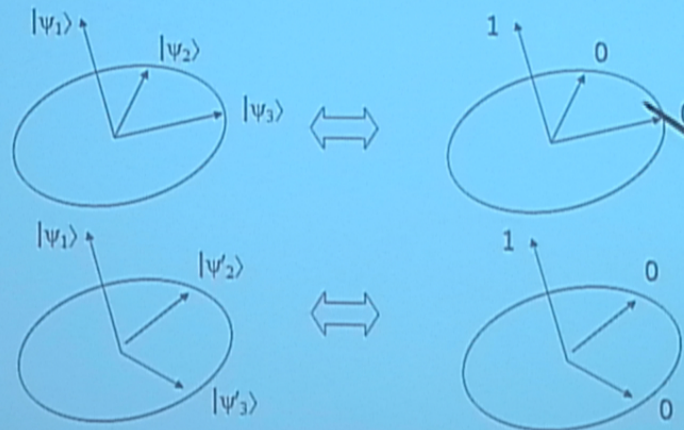
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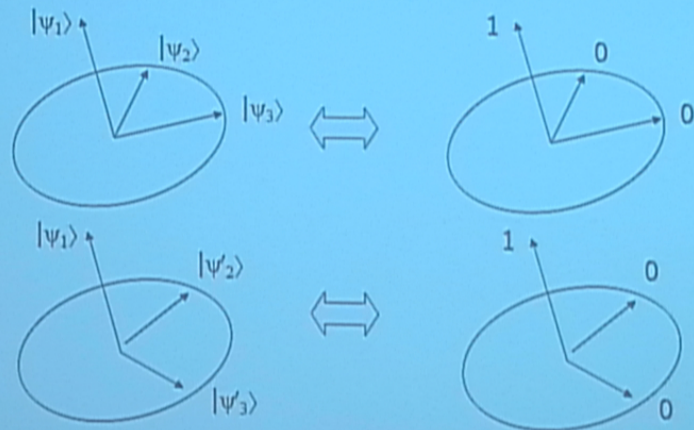
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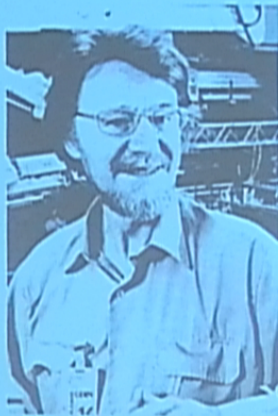
For every λ , every basis of vectors receives a 0-1 valuation, wherein exactly one element is assigned the value 1 (corresponding to the outcome that would occur for λ), and every vector is assigned the same value regardless of the basis it is considered a part (i.e. the context).



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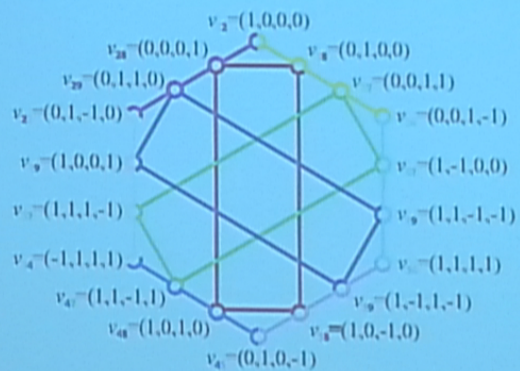
John S. Bell



Ernst Specker (with son) and
Simon Kochen

Bell-Kochen-Specker theorem: A noncontextual hidden variable model of quantum theory for Hilbert spaces of dimension 3 or greater is impossible.

Example: The CEGA algebraic 18 ray proof in 4d:
Cabello, Estebaranz, Garcia-Alcaine, Phys. Lett. A 212, 183 (1996)



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If we list all 9 orthogonal quadruples, each ray appears twice in the list

0,0,0,1	0,0,0,1	1,-1,1,-1	1,-1,1,-1	0,0,1,0	1,-1,-1,1	1,1,-1,1	1,1,-1,1	1,1,1,-1
0,0,1,0	0,1,0,0	1,-1,-1,1	1,1,1,1	0,1,0,0	1,1,1,1	1,1,1,-1	-1,1,1,1	-1,1,1,1
1,1,0,0	1,0,1,0	1,1,0,0	1,0,-1,0	1,0,0,1	1,0,0,-1	1,-1,0,0	1,0,1,0	1,0,0,1
1,-1,0,0	1,0,-1,0	0,0,1,1	0,1,0,-1	1,0,0,-1	0,1,-1,0	0,0,1,1	0,1,0,-1	0,1,-1,0

In each of the 9 quadruples, one ray is assigned 1, the other three 0.
Therefore, 9 rays must be assigned 1

But each ray appears twice and so there must be an even number
of rays assigned 1

CONTRADICTION!

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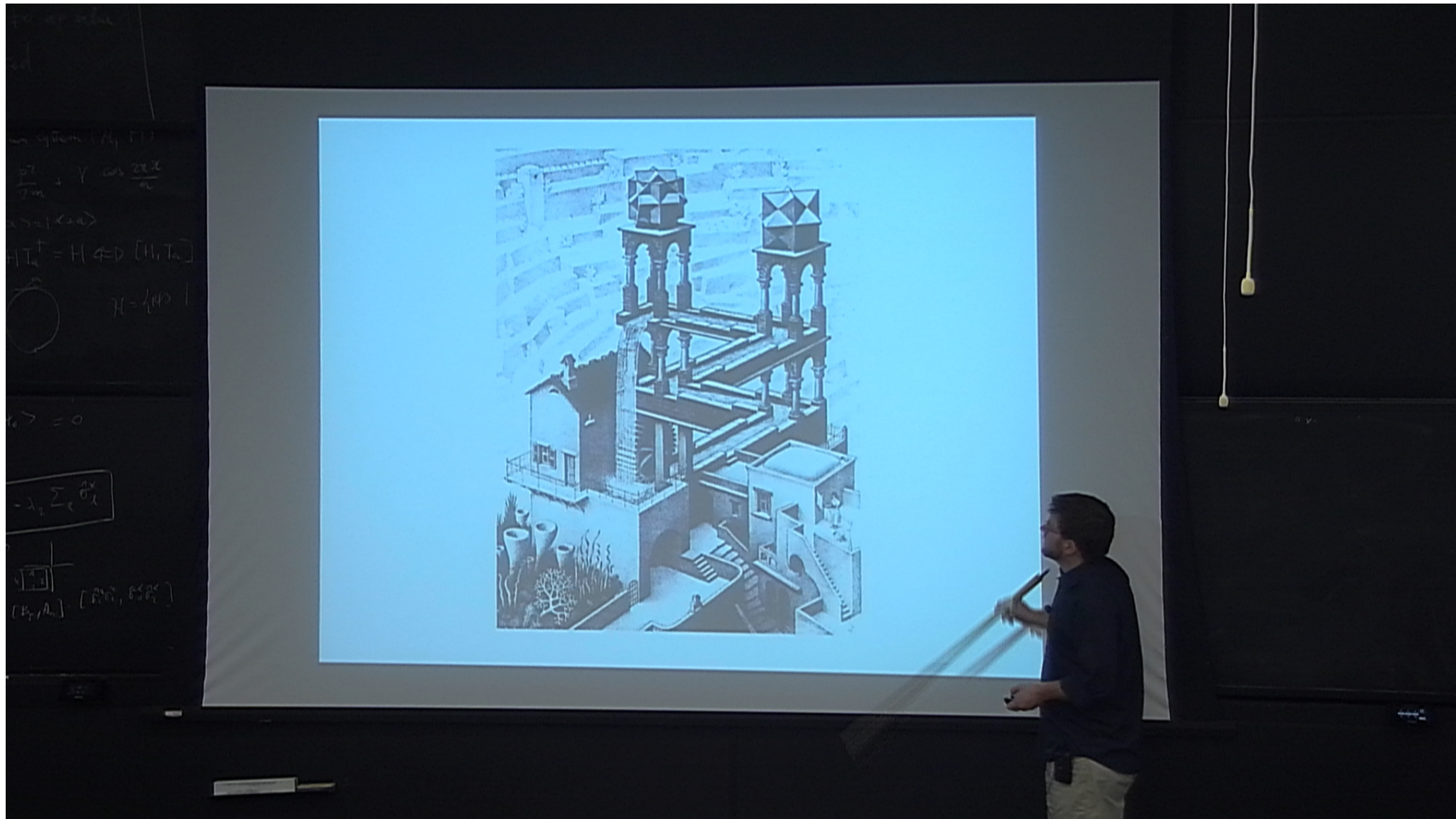
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0,0,1,0	0,1,0,0	1,-1,-1,1	1,1,1,1	0,1,0,0	1,1,1,1	1,1,1,-1	-1,1,1,1	-1,1,1,1
1,1,0,0	1,0,1,0	1,1,0,0	1,0,-1,0	1,0,0,1	1,0,0,-1	1,-1,0,0	1,0,1,0	1,0,0,1
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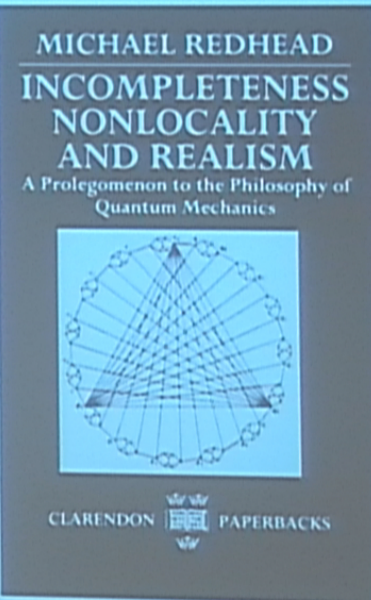
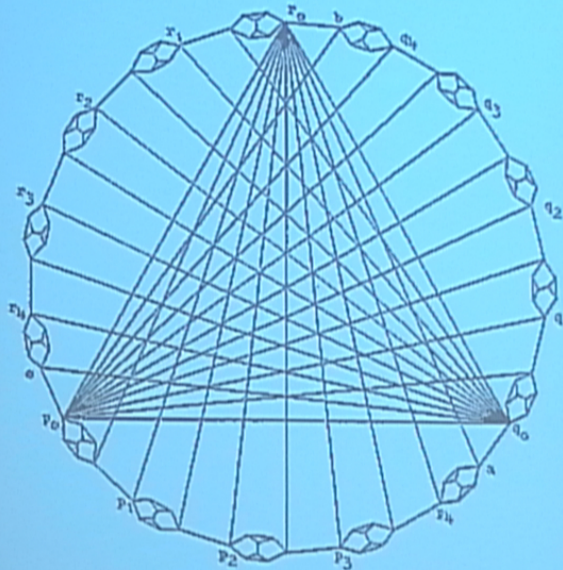
0,0,0,1	0,0,0,1	1,-1,1,-1	1,-1,1,-1	0,0,1,0	1,-1,-1,1	1,1,-1,1	1,1,-1,1	1,1,1,-1
0,0,1,0	0,1,0,0	1,-1,-1,1	1,1,1,1	0,1,0,0	1,1,1,1	1,1,1,-1	-1,1,1,1	-1,1,1,1
1,1,0,0	1,0,1,0	1,1,0,0	1,0,-1,0	1,0,0,1	1,0,0,-1	1,-1,0,0	1,0,1,0	1,0,0,1
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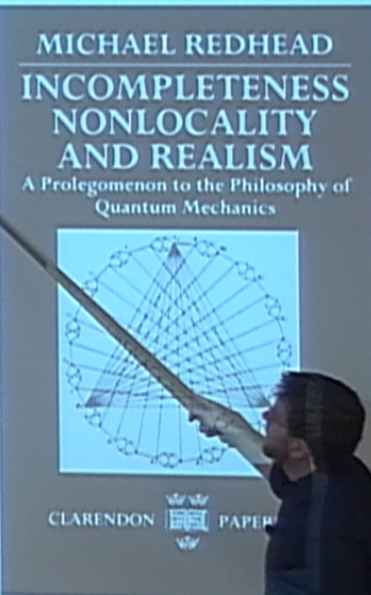
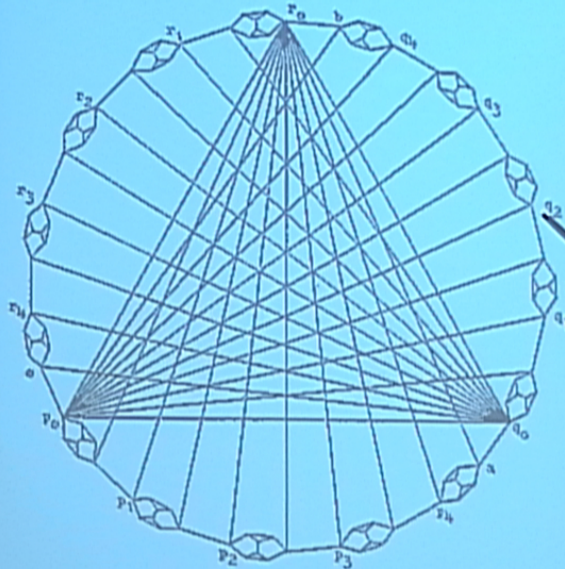
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Example: Kochen and Specker's original algebraic 117 ray proof in 3d



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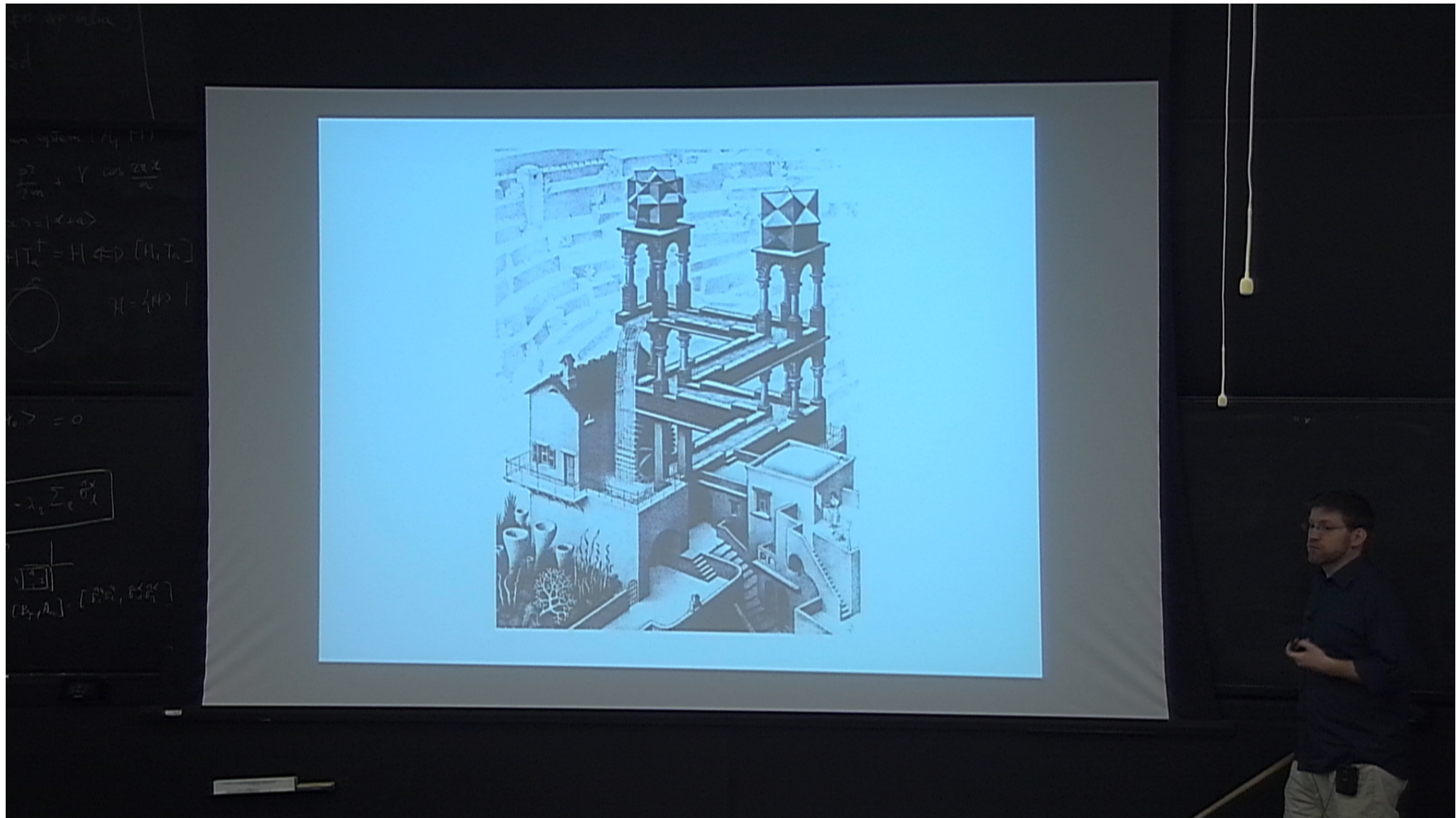
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0,0,1,0	0,1,0,0	1,-1,-1,1	1,1,1,1	0,1,0,0	1,1,1,1	1,1,1,-1	-1,1,1,1	-1,1,1,1
1,1,0,0	1,0,1,0	1,1,0,0	1,0,-1,0	1,0,0,1	1,0,0,-1	1,-1,0,0	1,0,1,0	1,0,0,1
1,-1,0,0	1,0,-1,0	0,0,1,1	0,1,0,-1	1,0,0,-1	0,1,-1,0	0,0,1,1	0,1,0,-1	0,1,-1,0

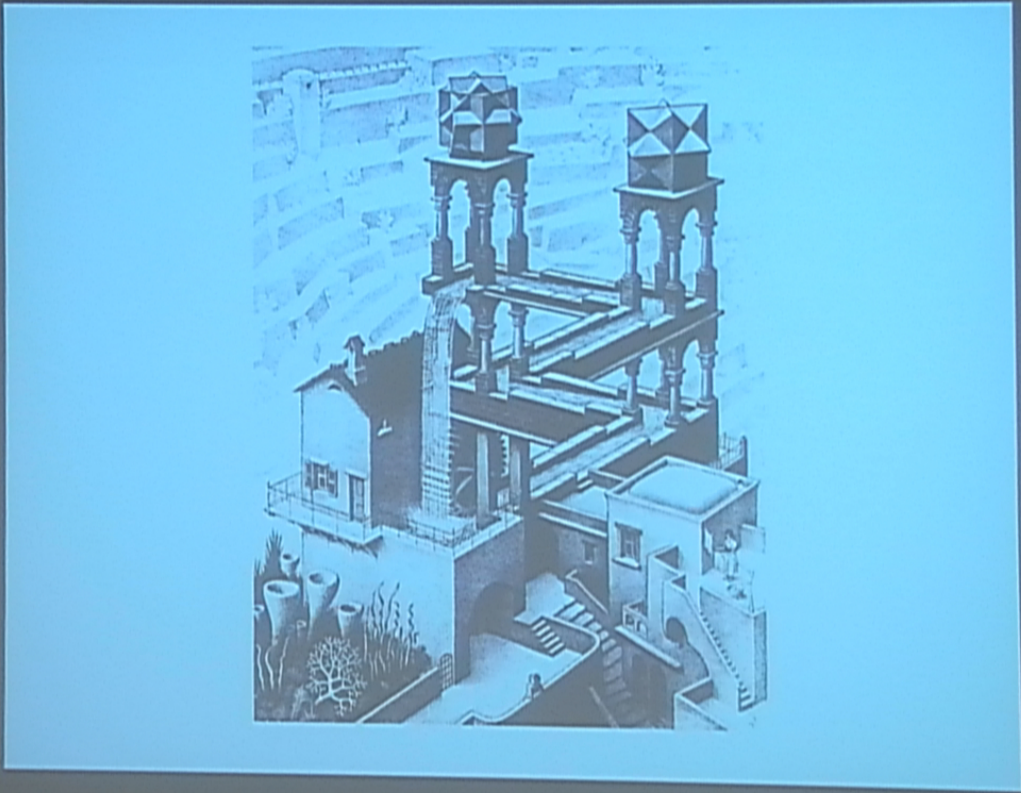
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$\frac{\partial^2}{\partial x^2} + V \cos \frac{2\pi x}{a}$
 $\langle \psi | \langle \psi |$
 $[T_a]^\dagger = H \Leftrightarrow [H, T_a]$
 $\mathcal{H} = \{H\}$
 $\lambda_2 \sum \theta_i^x$
 $[B_1, A_2] \quad [B_2, A_1]$



The traditional notion of noncontextuality:

For every λ , every projector Π is assigned a value 0 or 1 regardless of how it is measured (i.e. the context)

$$v(\Pi) = 0 \text{ or } 1 \quad \text{for all } \Pi$$

Coarse-graining of a measurement implies a coarse-graining of the value (because it is just post-processing)

$$v(\sum_k \Pi_k) = \sum_k v(\Pi_k)$$

Every measurement has *some* outcome

$$v(I) = 1$$

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$$[A, B] = 0 \quad [A, C] = 0 \quad [B, C] \neq 0$$

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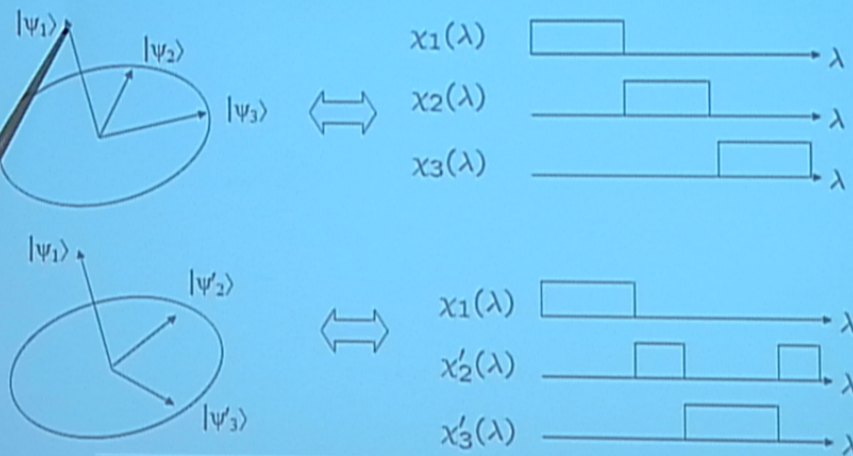
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$v(\Pi_a)$ is independent of context $\rightarrow v(A)$ is independent of context

Functional relationships among commuting Hermitian operators
must be respected by their values

$$\begin{aligned} \text{If } f(L, M, N, \dots) &= 0 \\ \text{then } f(v(L), v(M), v(N), \dots) &= 0 \end{aligned}$$

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Proof: the possible sets of eigenvalues one can simultaneously assign to L, M, N, \dots are specified by their joint eigenstates. By acting the first equation on each of the joint eigenstates, we get the second.

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Example: Mermin's magic square proof in 4d

X_1	X_2	$X_1 X_2$
Y_2	Y_1	$Y_1 Y_2$
$X_1 Y_2$	$Y_1 X_2$	$Z_1 Z_2$

$I \quad I \quad -I$

$$\begin{aligned}
 I \quad & X_1 X_2 X_1 X_2 = I \\
 & Y_1 Y_2 Y_1 Y_2 = I \\
 I \quad & X_1 Y_2 Y_1 X_2 Z_1 Z_2 = I \\
 & X_1 Y_2 X_1 Y_2 = I \\
 & Y_1 X_2 Y_1 X_2 = I \\
 I \quad & X_1 X_2 Y_1 Y_2 Z_1 Z_2 = -I
 \end{aligned}$$

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X_1	X_2	X_1X_2
Y_2	Y_1	Y_1Y_2
X_1Y_2	Y_1X_2	Z_1Z_2

$I \quad I \quad -I$

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 I \quad & X_1 X_2 X_1 X_2 = I \\
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Y_2	Y_1	$Y_1 Y_2$
$X_1 Y_2$	$Y_1 X_2$	$Z_1 Z_2$

I

I

I

I

I

$-I$

$$X_1 X_2 X_1 X_2 = I$$

$$Y_1 Y_2 Y_1 Y_2 = I$$

$$X_1 Y_2 Y_1 X_2 Z_1 Z_2 = I$$

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Example: Mermin's magic square proof in 4d

X_1	X_2	X_1X_2
Y_2	Y_1	Y_1Y_2
X_1Y_2	Y_1X_2	Z_1Z_2

I

I

I

I

I

$-I$

$$X_1 X_2 X_1X_2 = I$$

$$Y_1 Y_2 Y_1Y_2 = I$$

$$X_1Y_2 Y_1X_2 Z_1Z_2 = I$$

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X_1	X_2	$X_1 X_2$
Y_2	Y_1	$Y_1 Y_2$
$X_1 Y_2$	$Y_1 X_2$	$Z_1 Z_2$

I

I

I

I

I

$-I$

$$X_1 X_2 X_1 X_2 = I$$

$$Y_1 Y_2 Y_1 Y_2 = I$$

$$X_1 Y_2 Y_1 X_2 Z_1 Z_2 = I$$

$$X_1 Y_2 X_1 Y_2 = I$$

$$Y_1 X_2 Y_1 X_2 = I$$

$$X_1 X_2 Y_1 Y_2 Z_1 Z_2 = -I$$

Functional relationships among commuting Hermitian operators must be respected by their values

$$\begin{aligned} \text{If } f(L, M, N, \dots) &= 0 \\ \text{then } f(v(L), v(M), v(N), \dots) &= 0 \end{aligned}$$

Proof: the possible sets of eigenvalues one can simultaneously assign to L, M, N, \dots are specified by their joint eigenstates. By acting the first equation on each of the joint eigenstates, we get the second.

Example: Mermin's magic square proof in 4d

X_1	X_2	X_1X_2
Y_2	Y_1	Y_1Y_2
X_1Y_2	Y_1X_2	Z_1Z_2

I

I

I

$I \quad I \quad -I$

$$X_1 X_2 X_1X_2 = I$$

$$Y_1 Y_2 Y_1Y_2 = I$$

$$X_1Y_2 Y_1X_2 Z_1Z_2 = I$$

$$X_1 Y_2 X_1Y_2 = I$$

$$Y_1 X_2 Y_1X_2 = I$$

$$X_1X_2 Y_1Y_2 Z_1Z_2 = -I$$

$$v(X_1) v(X_2) v(X_1X_2) = 1$$

$$v(Y_1) v(Y_2) v(Y_1Y_2) = 1$$

$$v(X_1Y_2) v(Y_1X_2) v(Z_1Z_2) = 1$$

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I

I

I

$I \quad I \quad -I$

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Aside: Local determinism is an instance of traditional noncontextuality where the context is remote

$S_a^A \otimes I^B$ is either measured with $I^A \otimes S_b^B$
or with $I^A \otimes S_{b'}^B$

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Recall traditional noncontextuality:

For Hermitian operators A, B, C satisfying

$$[A, B] = 0 \quad [A, C] = 0 \quad [B, C] \neq 0$$

the value assigned to A should be independent of whether it is measured together with B or together with C (i.e. the context)

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Therefore $v(S_a^A)$ is the same for the two contexts

This is local determinism



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This is local determinism

Every proof of the impossibility of a locally deterministic model is a proof of the impossibility of a traditional noncontextual model

Problems with the traditional definition of noncontextuality:

- applies only to sharp measurements
- applies only to deterministic hidden variable models
- applies only to models of quantum theory

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An operational notion of noncontextuality would determine

- whether any given operational theory admits of a noncontextual model

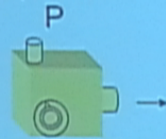
- whether any given experimental data can be explained by a noncontextual model

An Operational Notion of Contextuality

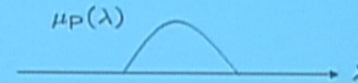
An ontological model of an operational theory

Specifies an ontic state space Λ

Preparation



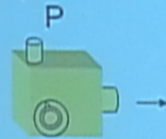
$$\int \mu_P(\lambda) d\lambda = 1$$



An ontological model of an operational theory

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Preparation



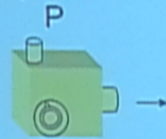
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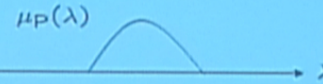
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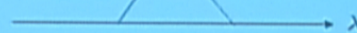
Preparation

P



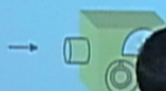
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$\mu_P(\lambda)$



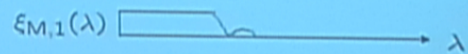
Measurement

M



$$0 \leq \xi_{M,k} \leq 1$$

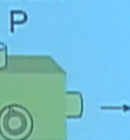
$$\sum_k \xi_{M,k}(\lambda) = 1 \text{ for all } \lambda$$



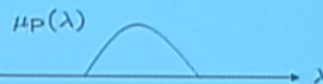
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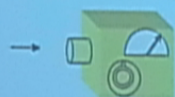
Preparation



$$\int \mu_P(\lambda) d\lambda = 1$$

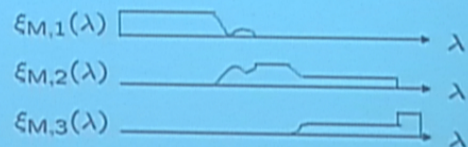


Measurement
M



$$0 \leq \xi_{M,k} \leq 1$$

$$\sum_k \xi_{M,k}(\lambda) = 1 \text{ for all } \lambda$$



$$p(k|P, M) = \int d\lambda \xi_{M,k}(\lambda) \mu_P(\lambda)$$

Generalized definition of noncontextuality:

An ontological model of an operational theory is noncontextual if

Operational equivalence
of two experimental
procedures



Equivalent
representations
in the ontological model

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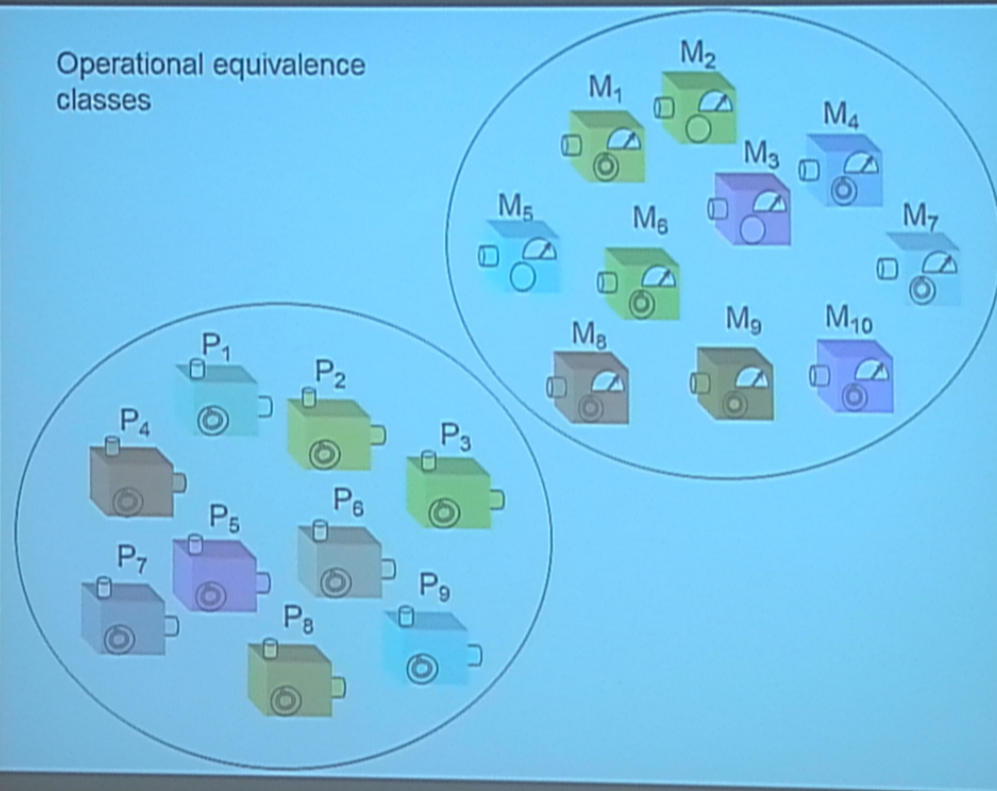
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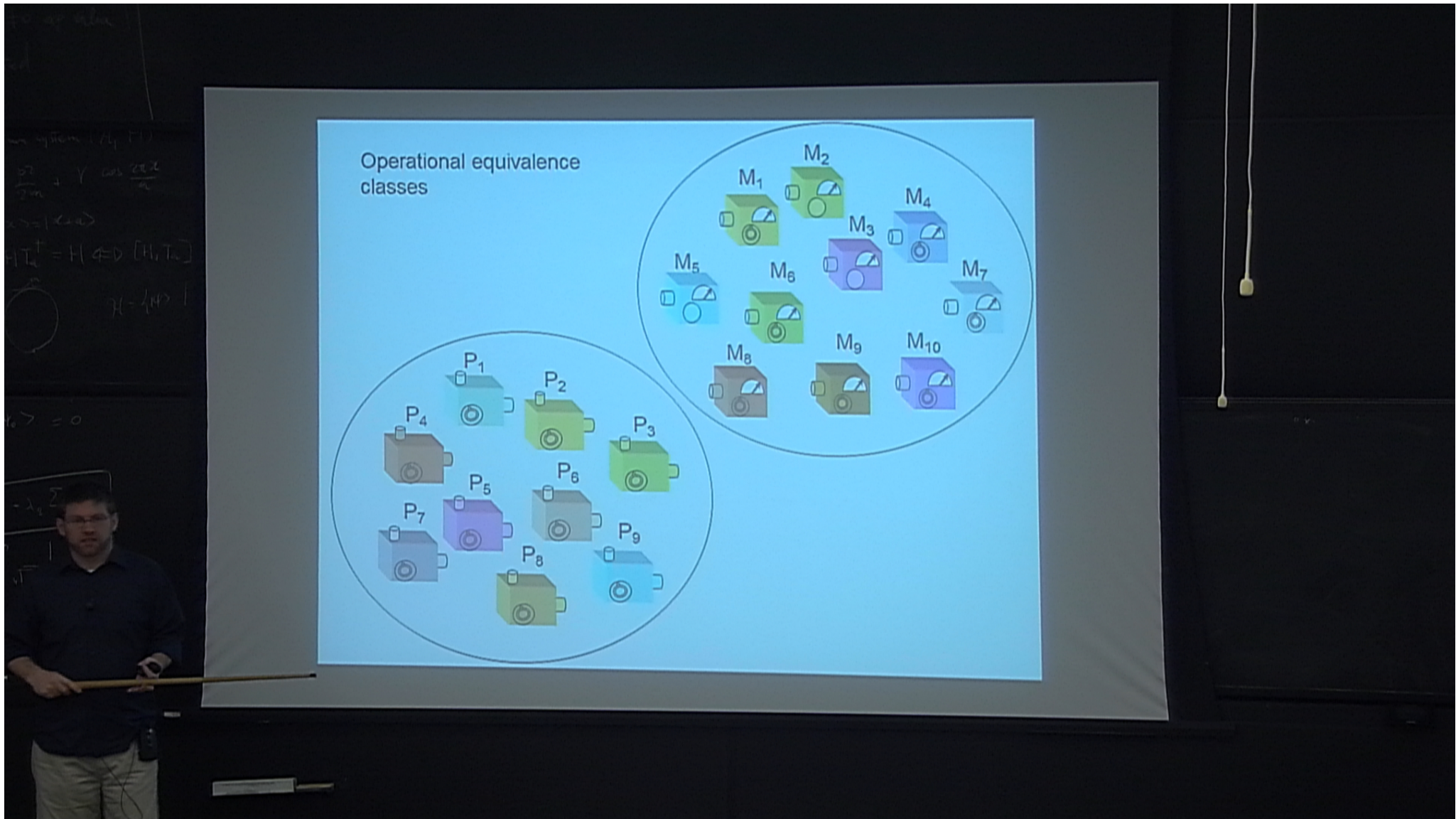
Operational equivalence
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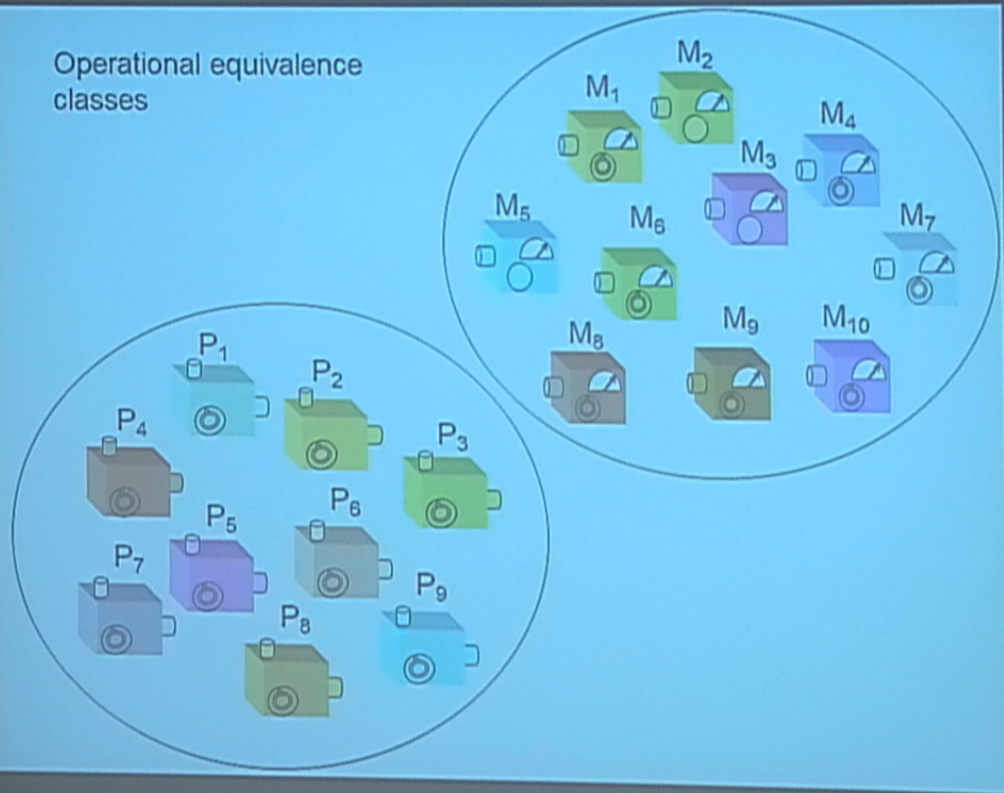
Equivalent
representations
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Operational equivalence classes

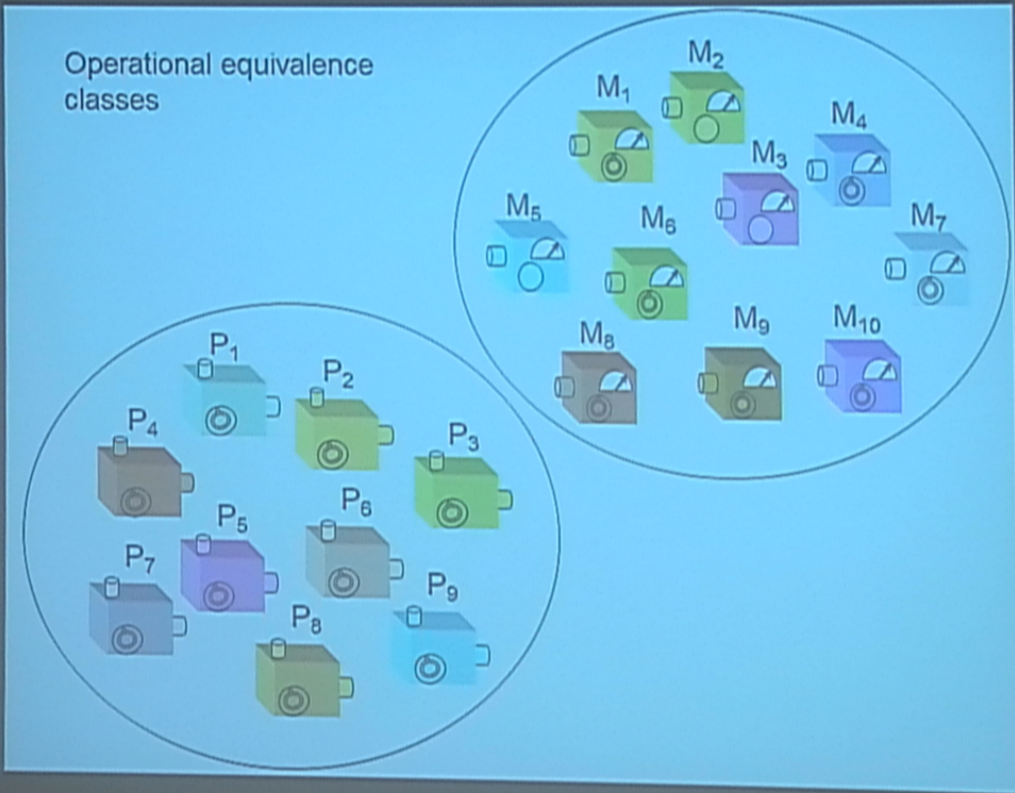


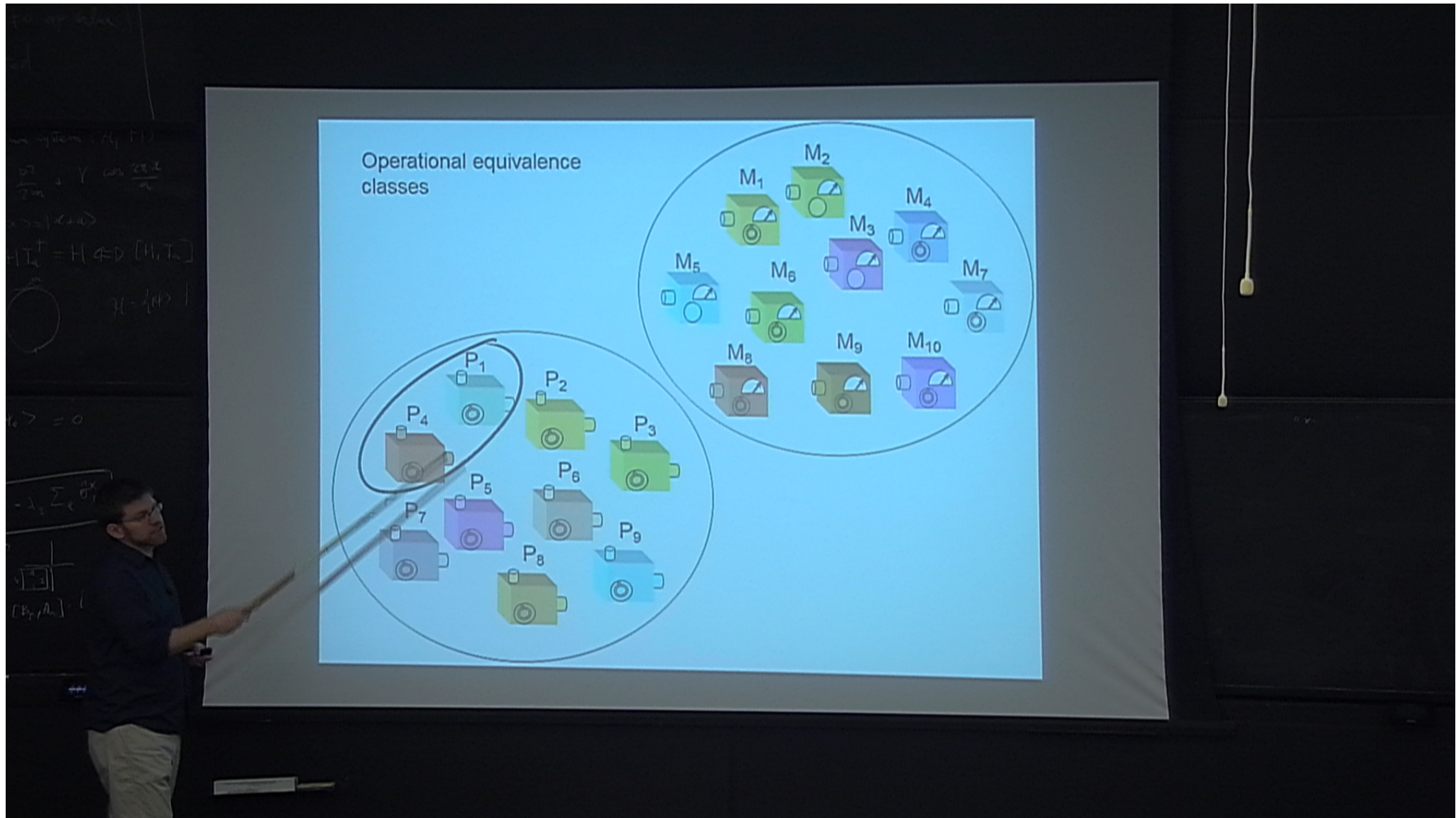


Operational equivalence classes

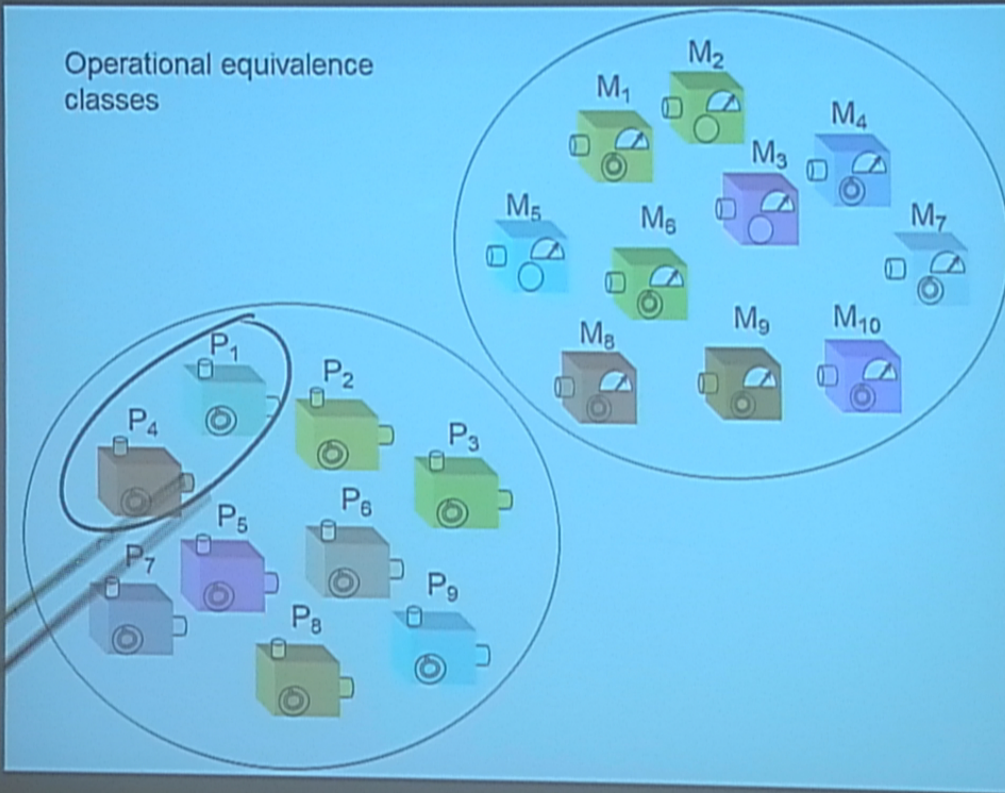


$\frac{d^2}{dt^2} + \gamma \cos \frac{2\pi x}{a}$
 $\langle \psi | \hat{H} | \psi \rangle$
 $[T_0] = H \Leftrightarrow D [H, T_0]$
 $H = \{H\}$
 $\lambda_1 = 0$
 $-\lambda_2$
 $\sqrt{\lambda_1}$
 $[1, 1]$

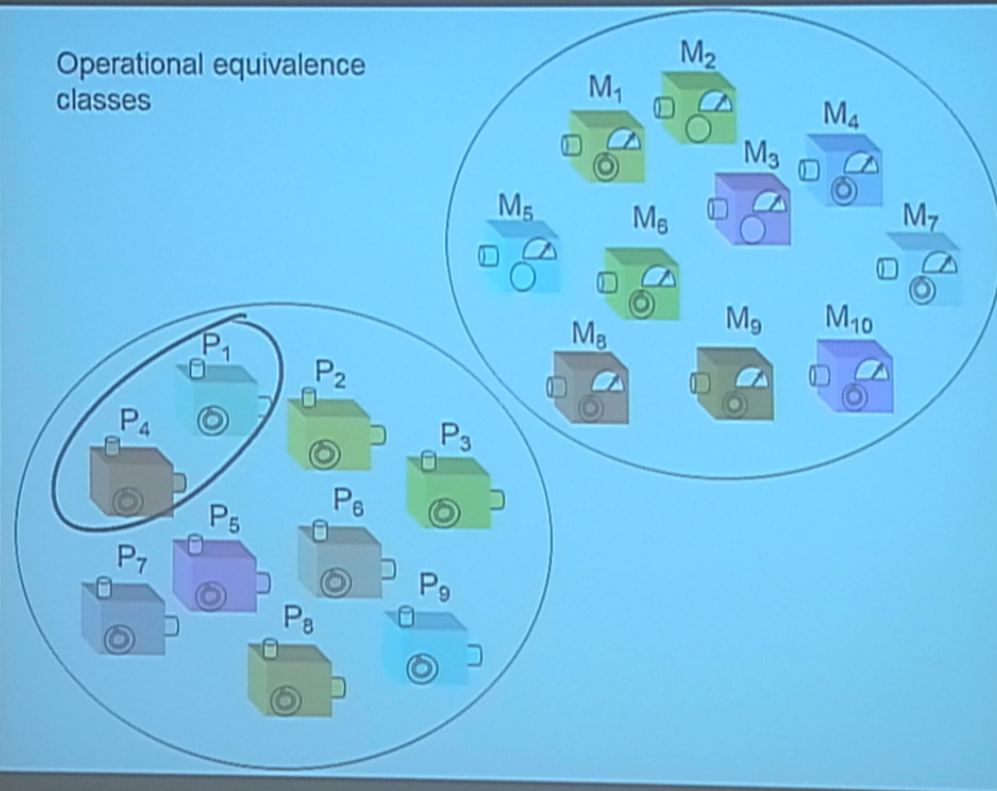




Operational equivalence classes

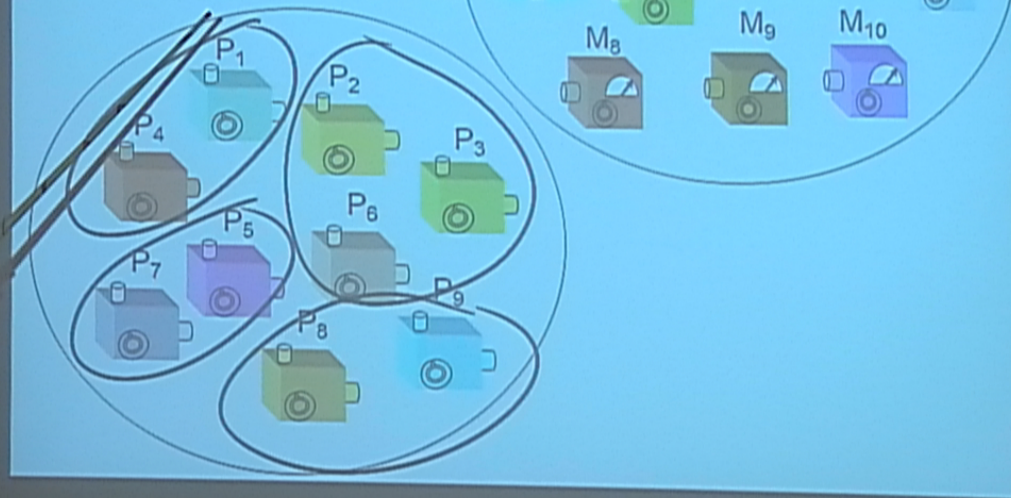


Operational equivalence classes

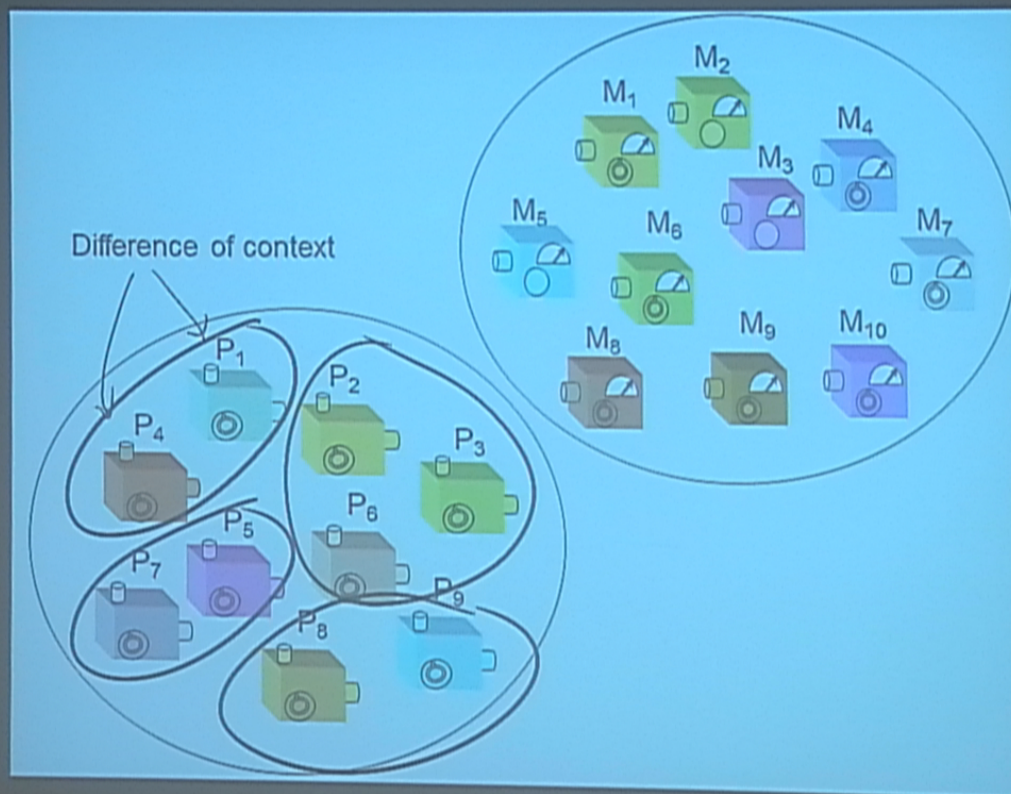


Operational equivalence classes

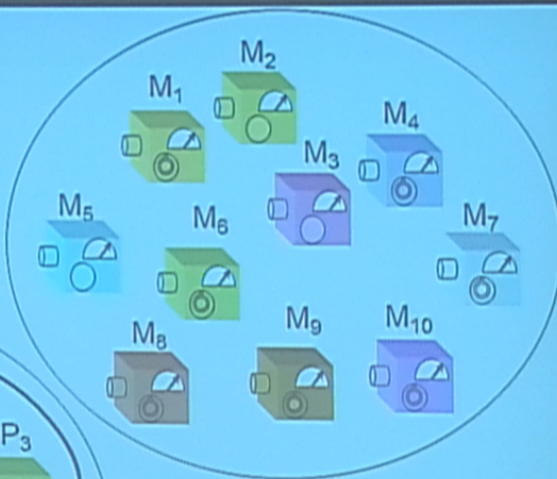
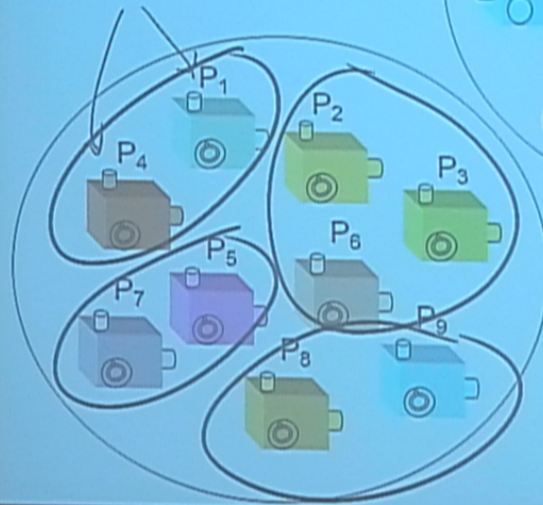
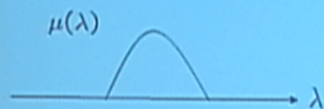
P is equivalent to P' if
 $\forall M \forall k :$
 $p(k|P, M) = p(k|P', M)$

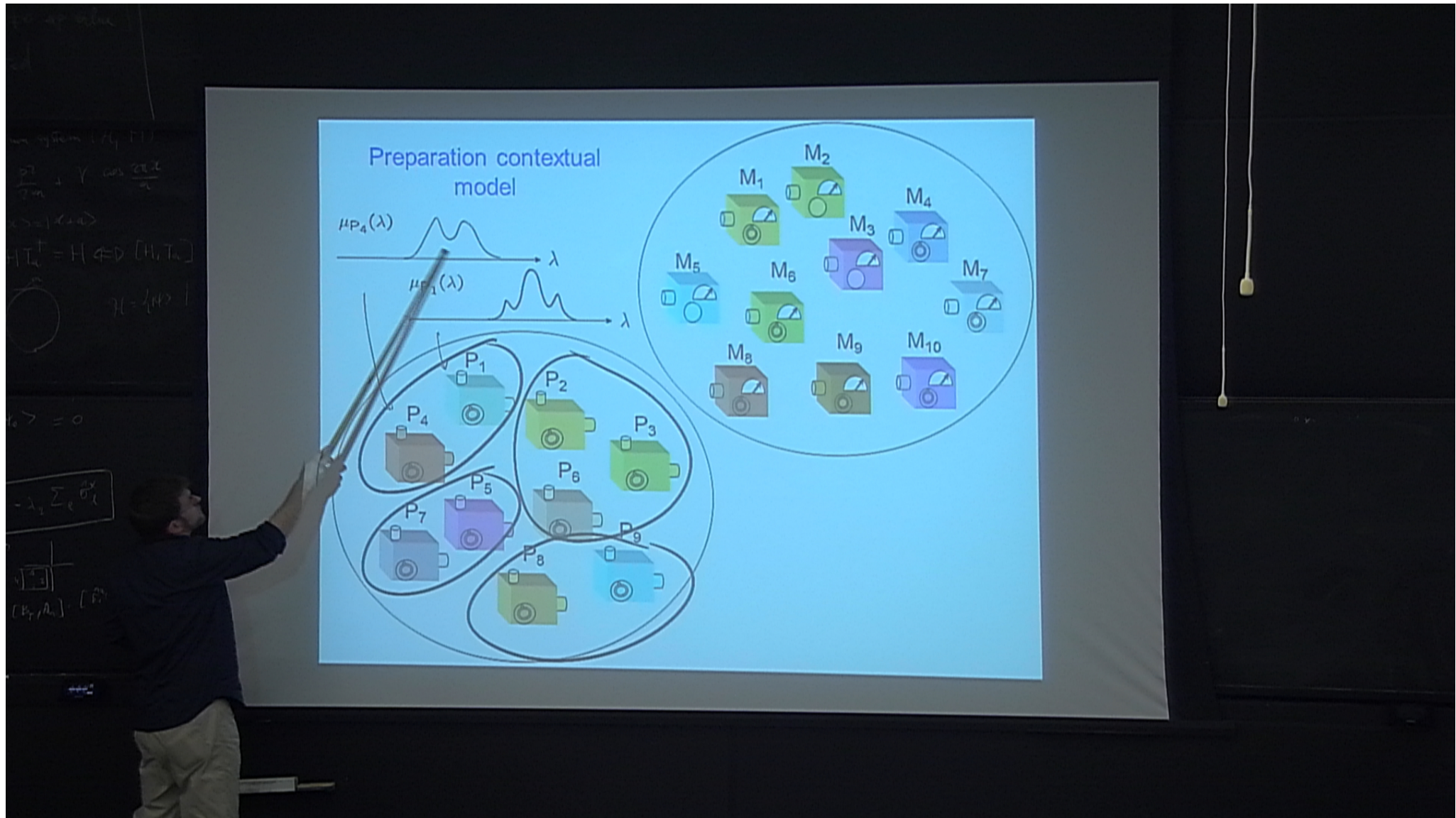


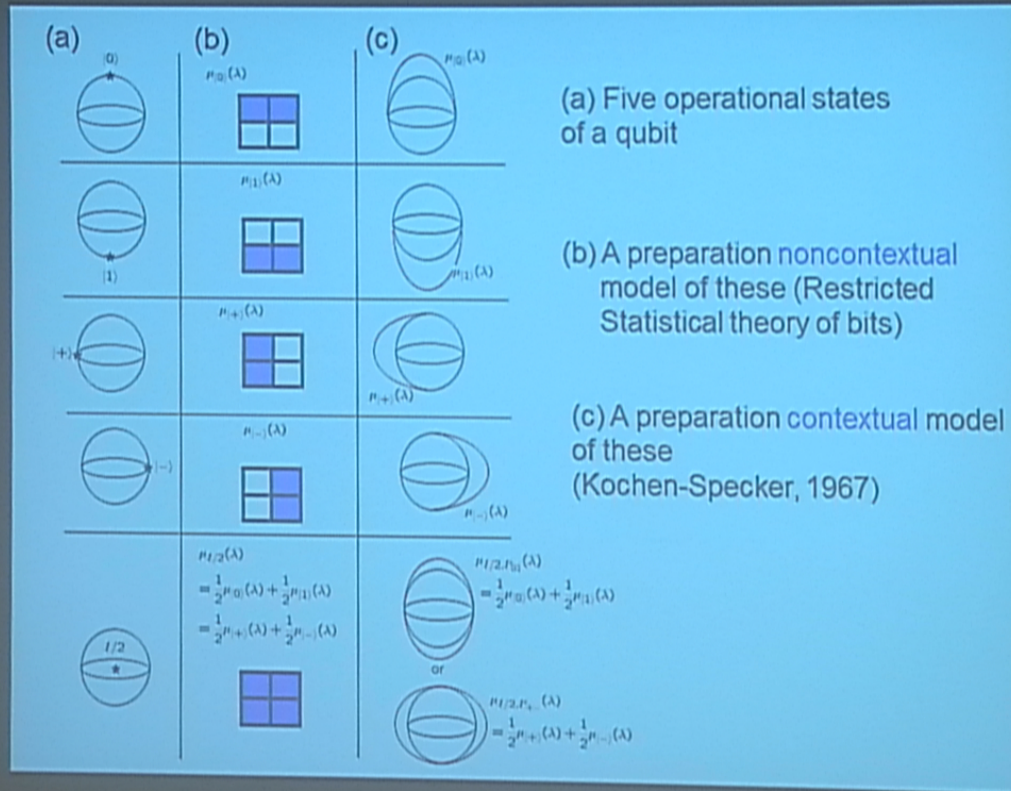
$\frac{\partial^2}{\partial x^2} + \gamma \cos \frac{2\pi x}{a}$
 $\langle \psi | \hat{H} | \psi \rangle$
 $H = \langle H \rangle$
 $\lambda_2 \sum \theta_i^x$
 $[B_T, A_0] = [\theta_1^x, \theta_2^x]$



Preparation noncontextual model



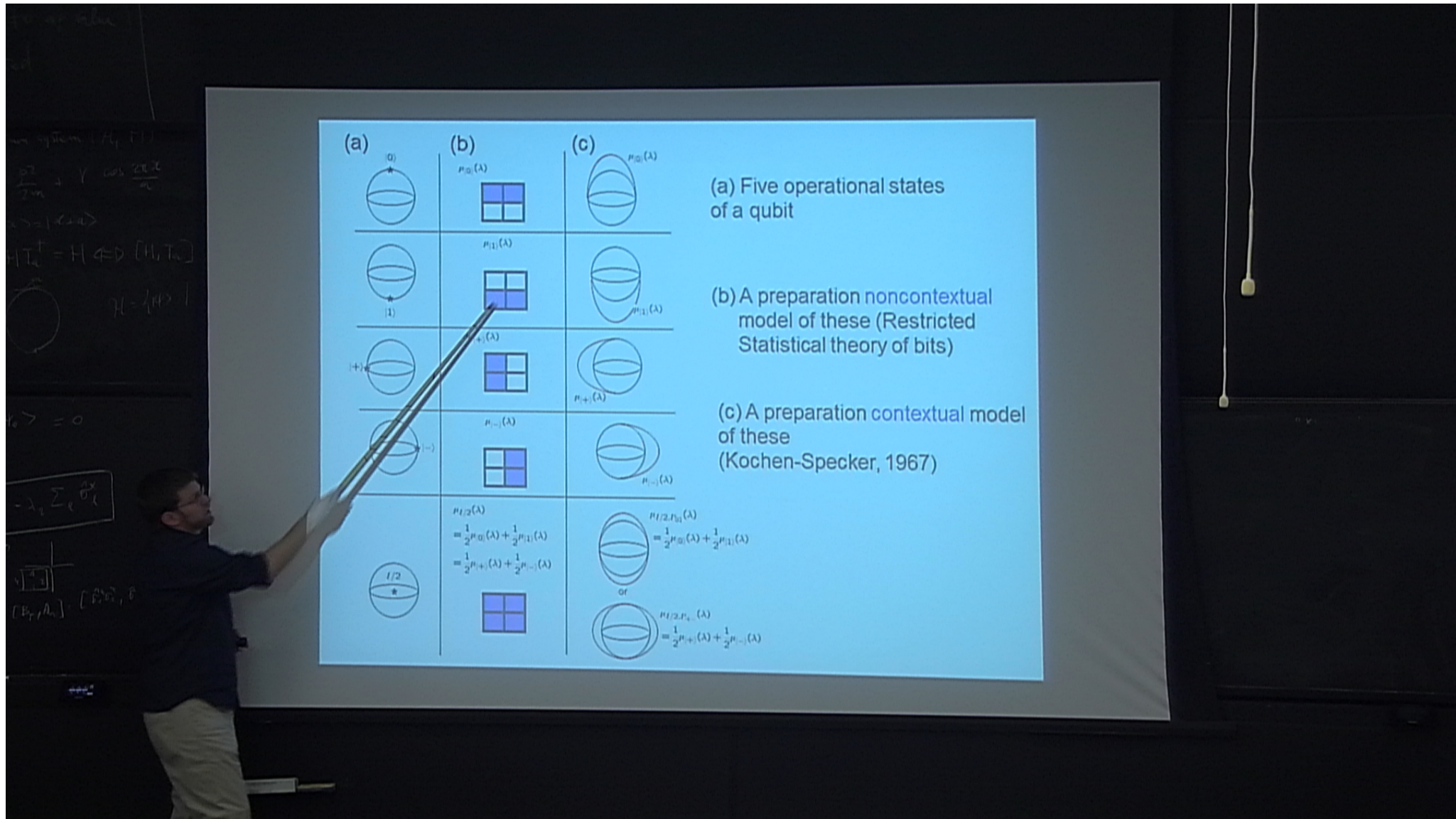


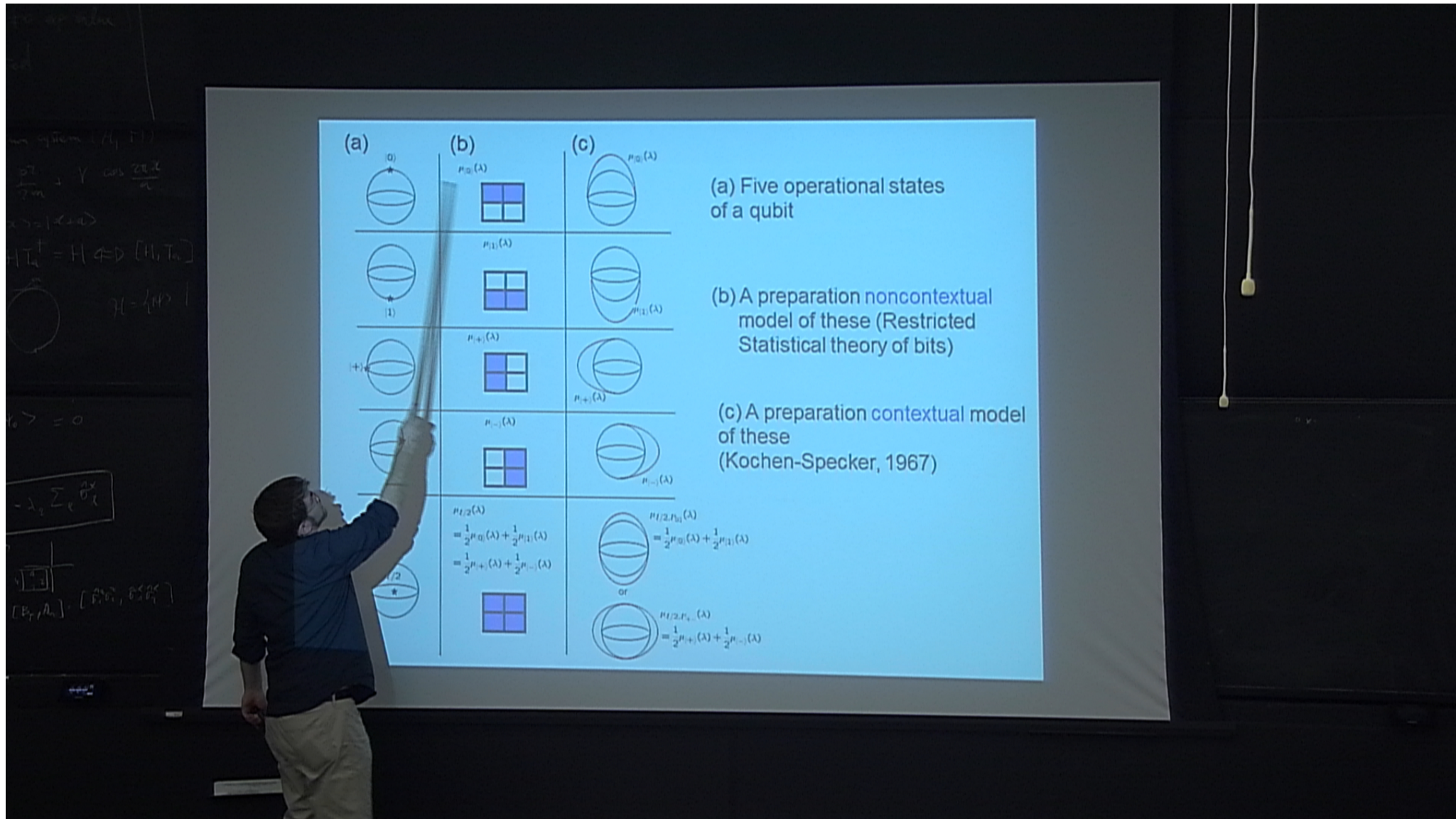


(a) Five operational states of a qubit

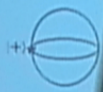
(b) A preparation noncontextual model of these (Restricted Statistical theory of bits)

(c) A preparation contextual model of these (Kochen-Specker, 1967)

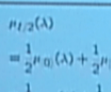
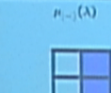
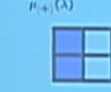
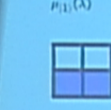
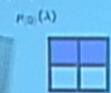




(a)

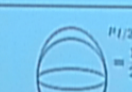
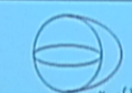
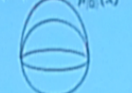


(b)



$$p_{1/2}(A) = \frac{1}{2}p_{|0>}(A) + \frac{1}{2}p_{|1>}(A) = \frac{1}{2}p_{|+>}(A) + \frac{1}{2}p_{|->}(A)$$

(c)

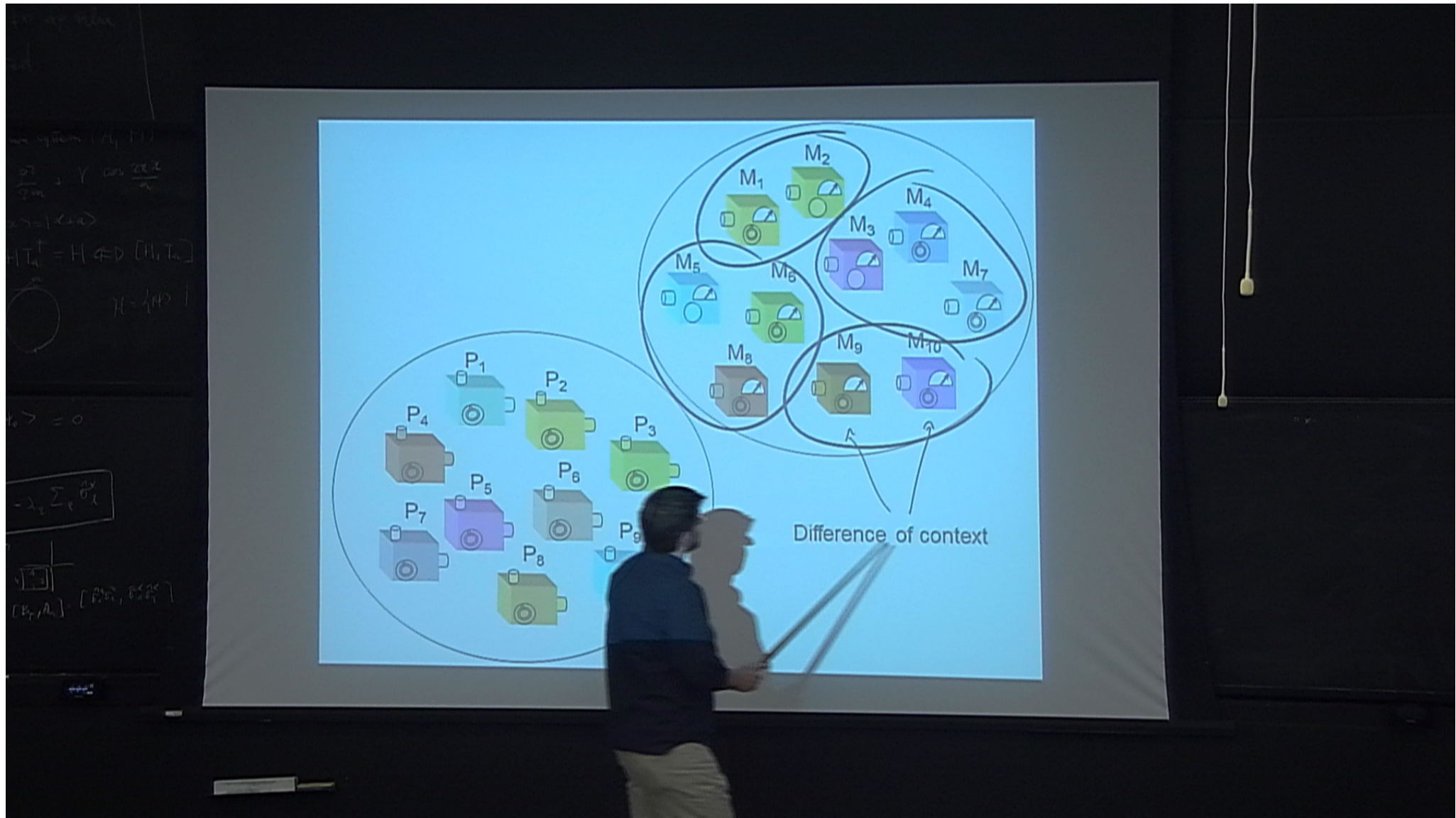


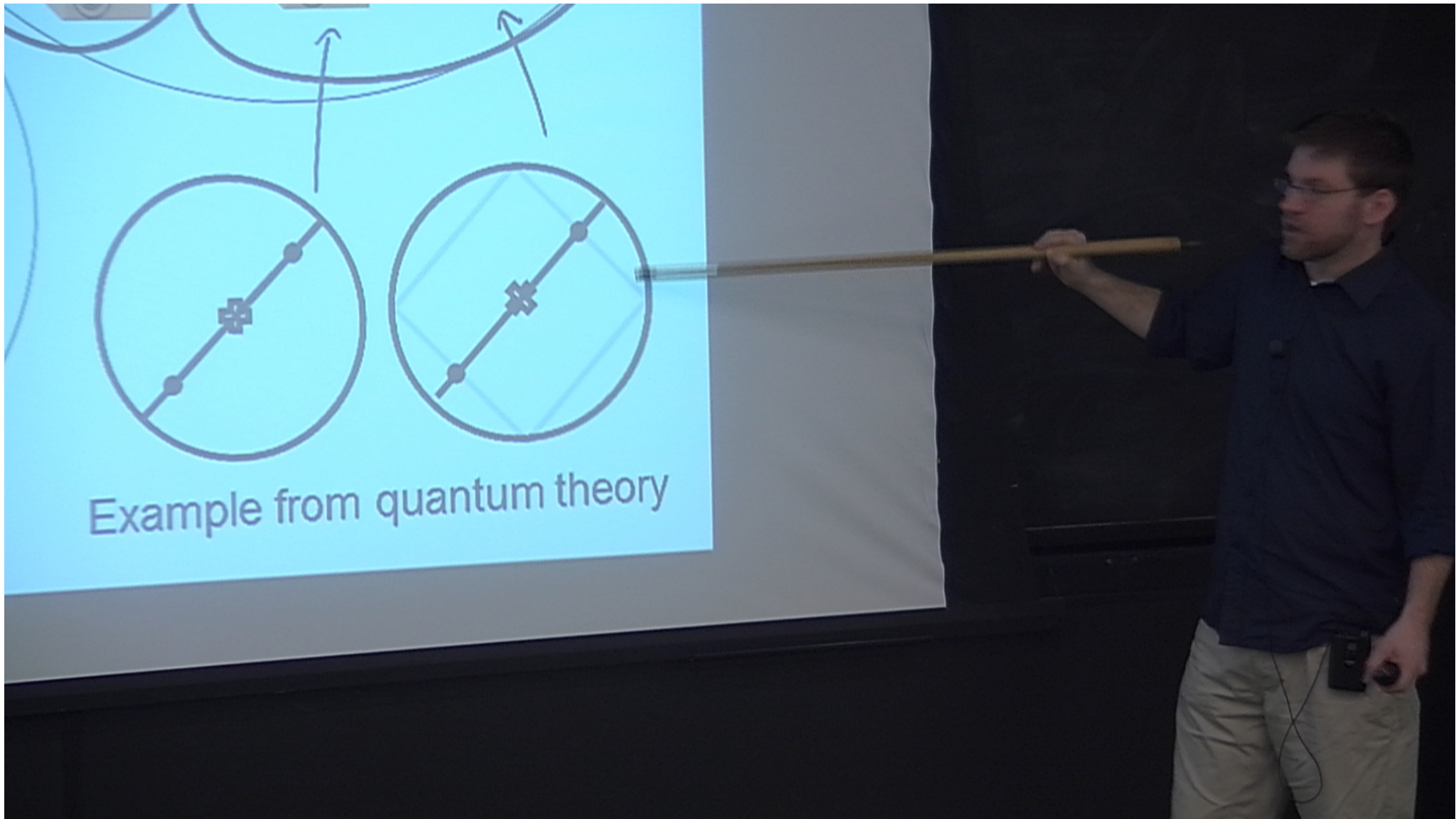
$$p_{1/2, P_{|0>}}(A) = \frac{1}{2}p_{|0>}(A) + \frac{1}{2}p_{|1>}(A) \text{ or } p_{1/2, P_{|+>}}(A) = \frac{1}{2}p_{|+>}(A) + \frac{1}{2}p_{|->}(A)$$

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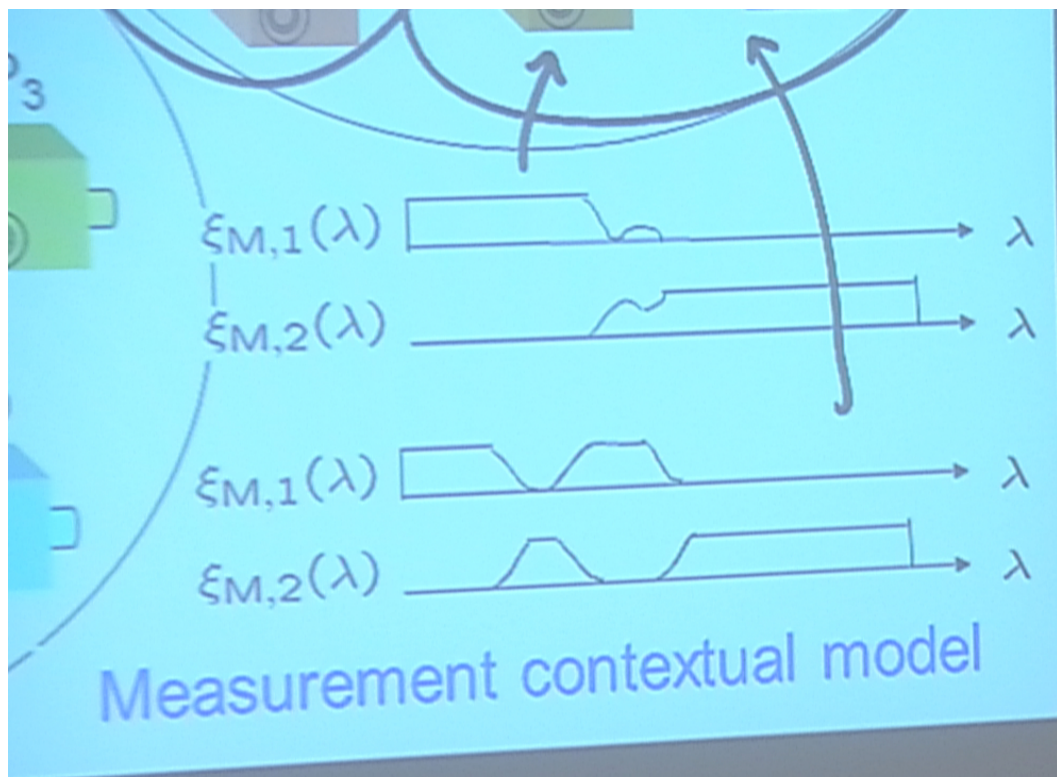


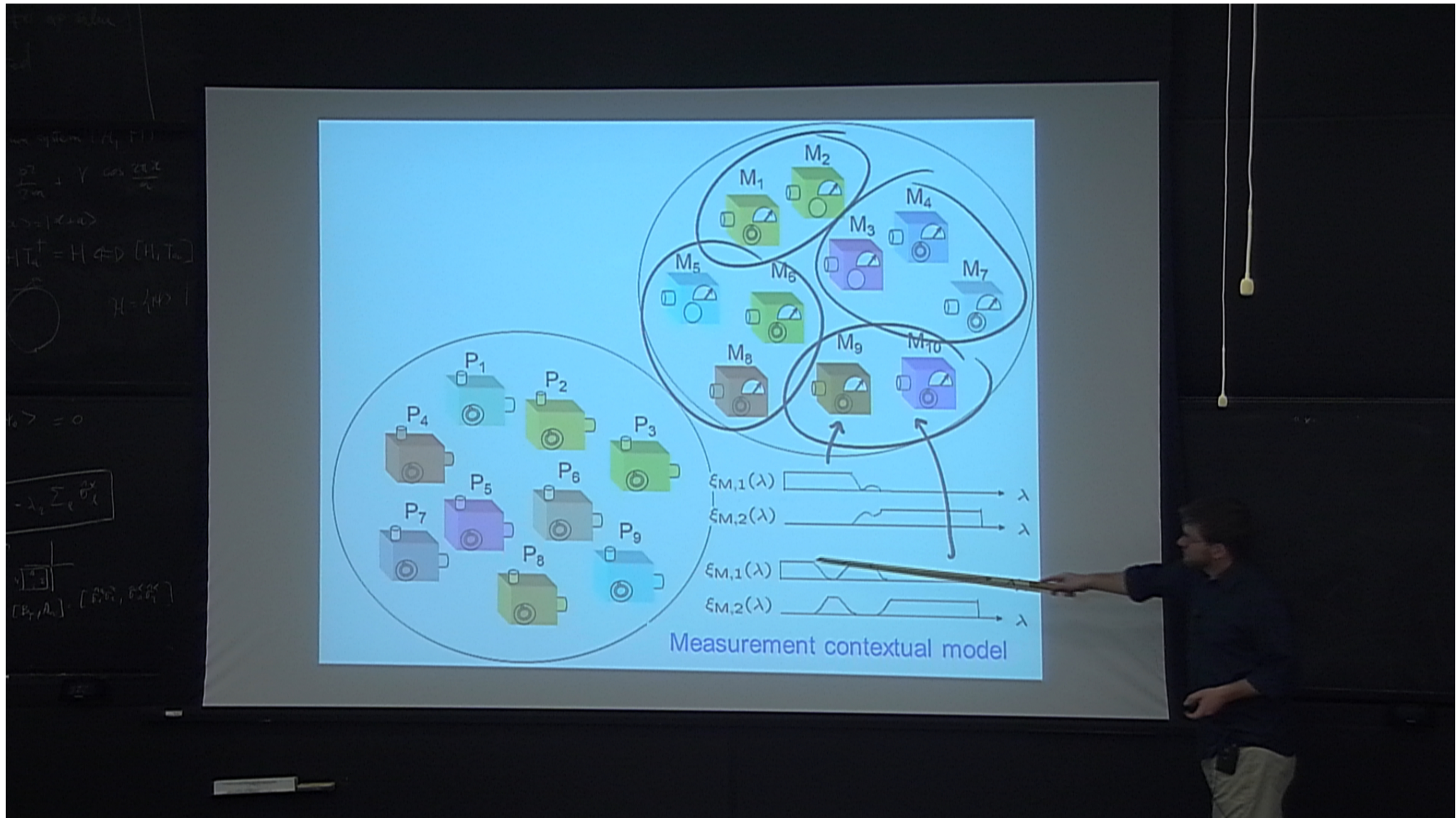


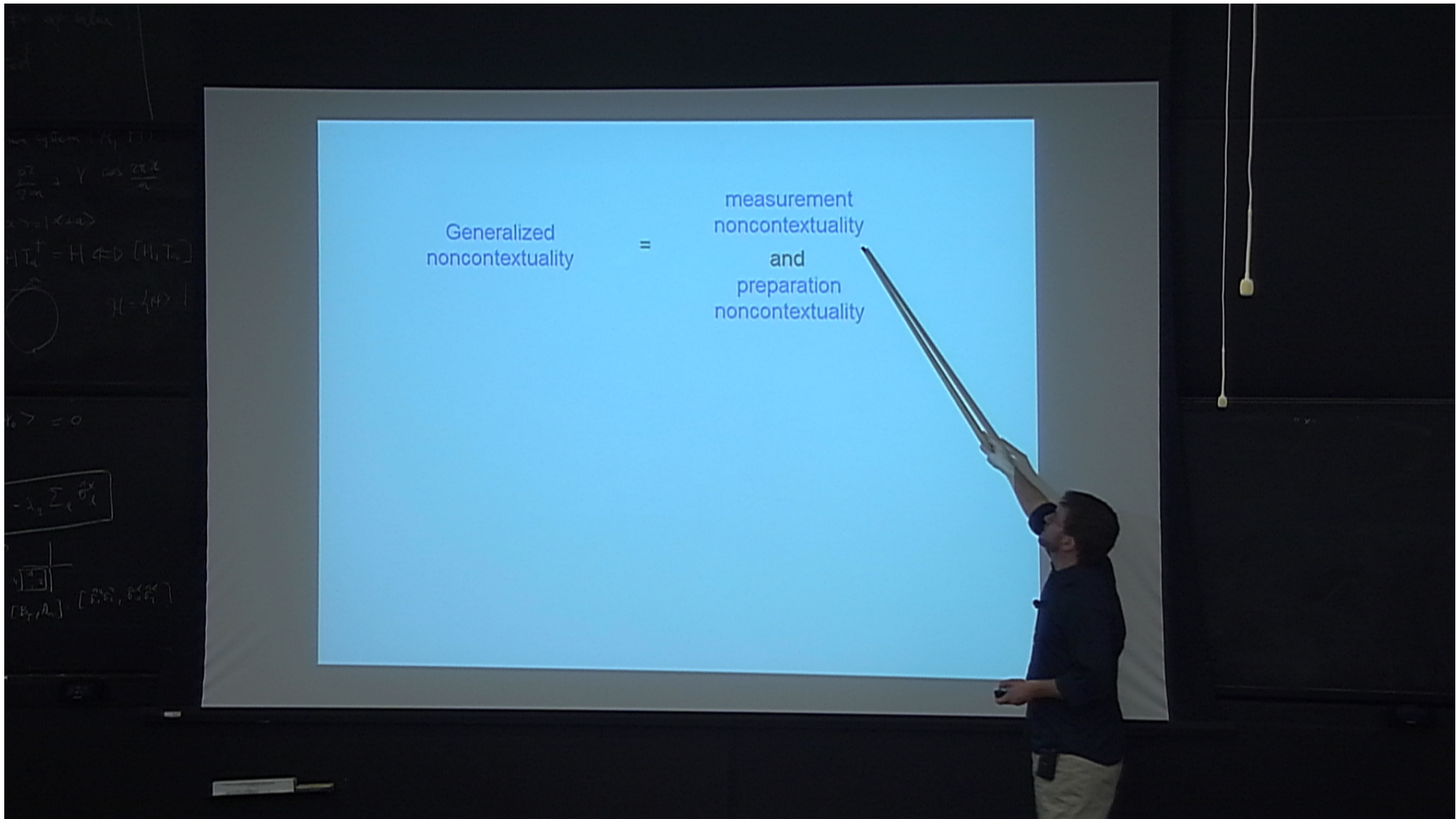
$\frac{\partial^2}{\partial x^2} + V \cos \frac{2\pi x}{a}$
 $\langle x \rangle = \langle x \rangle$
 $[H, T] = 0$
 $H = \{H\}$
 $\lambda_1 = 0$
 $-\lambda_1 \sum \sigma_i^x$
 $[B_1, A_1] = [B_2, A_2]$

$\{E, I - E\}$
 $E = q|\frac{n}{4}\rangle\langle\frac{n}{4}| + (1 - q)\frac{1}{2}I$
 $\{E, I - E\}$
 $E = \frac{1}{2}|0\rangle\langle 0| + \frac{1}{2}|1\rangle\langle 1|$
 Example from quantum theory

$\frac{\partial^2}{\partial x^2} + V \cos \frac{2\pi x}{a}$
 $\langle x \rangle = \langle x \rangle$
 $[H, T] = 0$
 $H = \{H\}$
 $\lambda_1 = 0$
 $-\lambda_1 \sum \sigma_i^x$
 $[B_1, A_1] = [B_2, A_2]$







Generalized
noncontextuality

=

measurement
noncontextuality
and
preparation
noncontextuality

for system (A, T)
 $\frac{\partial^2}{\partial x^2} + V \cos \frac{2\pi x}{a}$
 $\langle \psi | \langle \psi |$
 $[T_A]^\dagger = H \nabla \cdot [H, T_A]$
 $\mathcal{H} = \{H\}$
 $\langle \psi | = 0$
 $-\lambda_2 \sum_i \sigma_i^x$
 $\begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$
 $[B_T, A_n] \quad [B_T, B_n]$

Generalized
noncontextuality

=

measurement
noncontextuality
and
preparation
noncontextuality

Claim: Preparation noncontextuality is as natural (or unnatural) as measurement noncontextuality

Q: Why is noncontextuality plausible at all?

A: The methodological equivalence principle: if a difference in set-up is not distinguished in the observable phenomenon, then it should not be distinguished in the ontological picture either

Generalized
noncontextuality

=

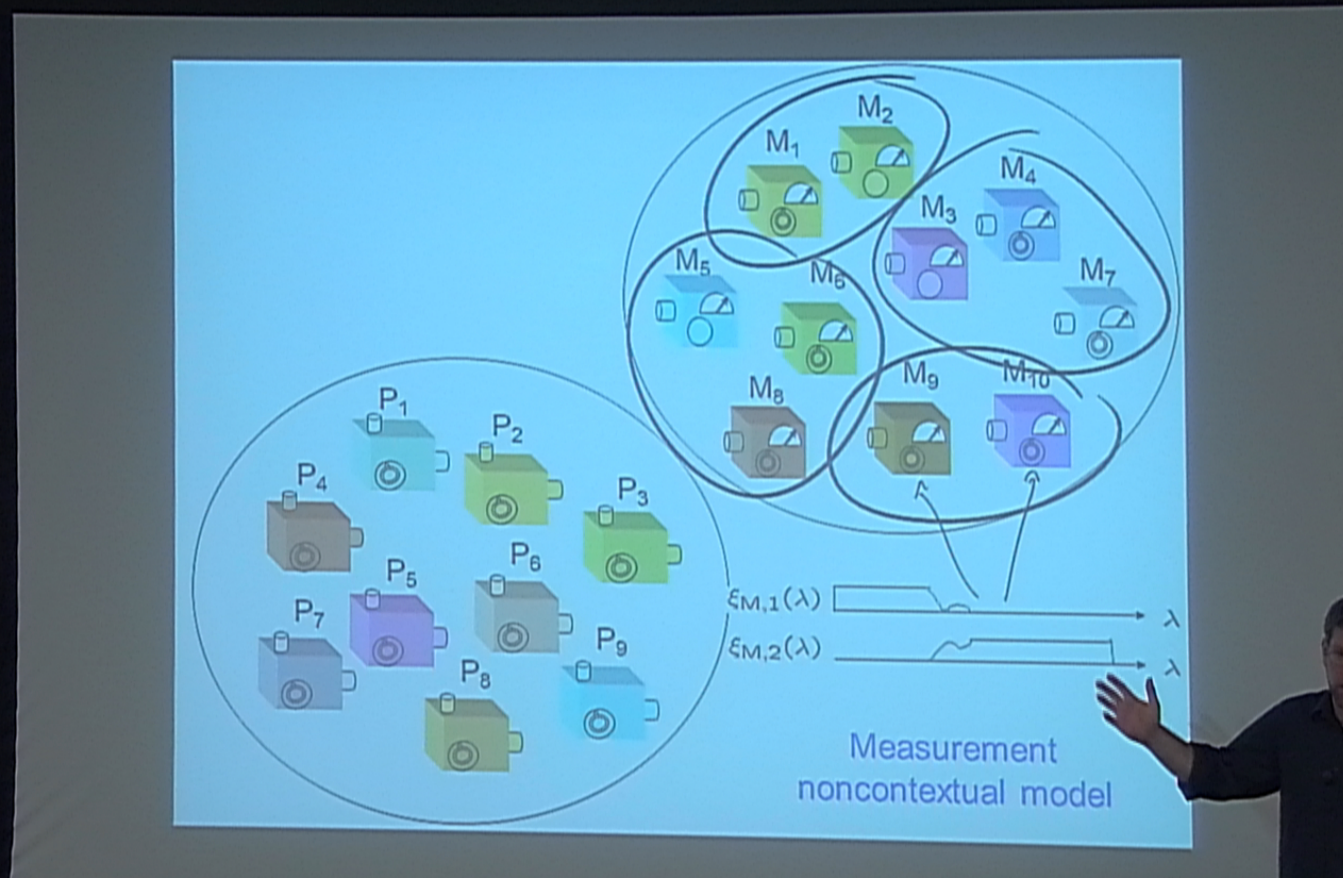
measurement
noncontextuality
and
preparation
noncontextuality

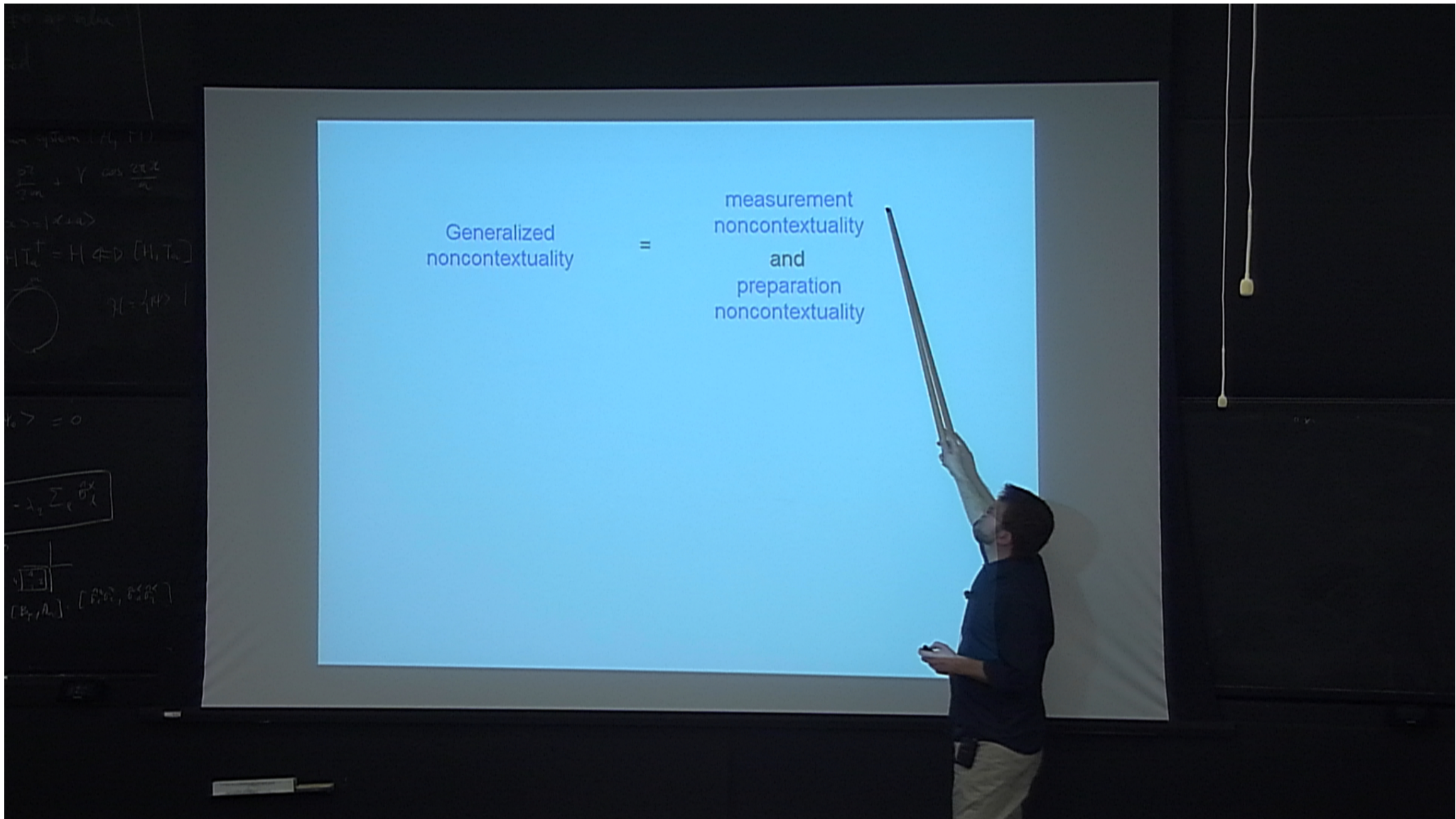
Claim: Preparation noncontextuality is as natural (or unnatural) as measurement noncontextuality

Q: Why is noncontextuality plausible at all?

A: The methodological equivalence principle: if a difference in set-up is not distinguished in the observable phenomena then it should not be distinguished in the ontological picture either

Preparation noncontextuality





Generalized
noncontextuality

=

measurement
noncontextuality
and
preparation
noncontextuality

Generalized
noncontextuality

=

measurement
noncontextuality
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