

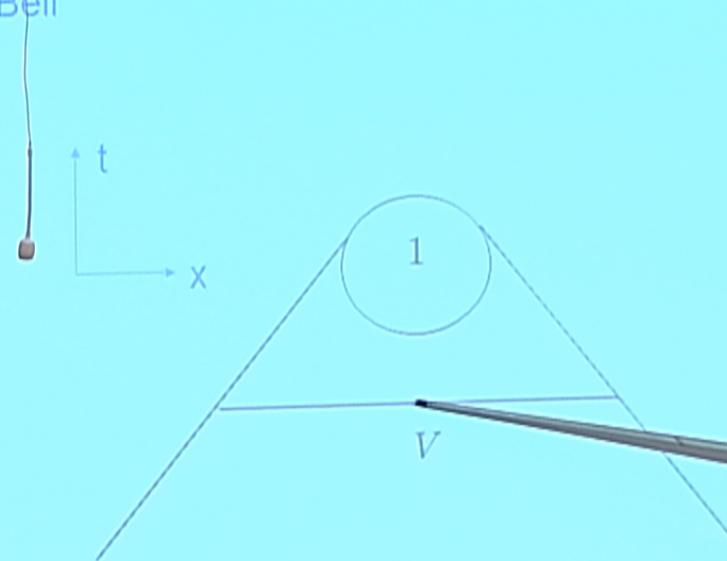
Title: Foundations of Quantum Mechanics - Lecture 9

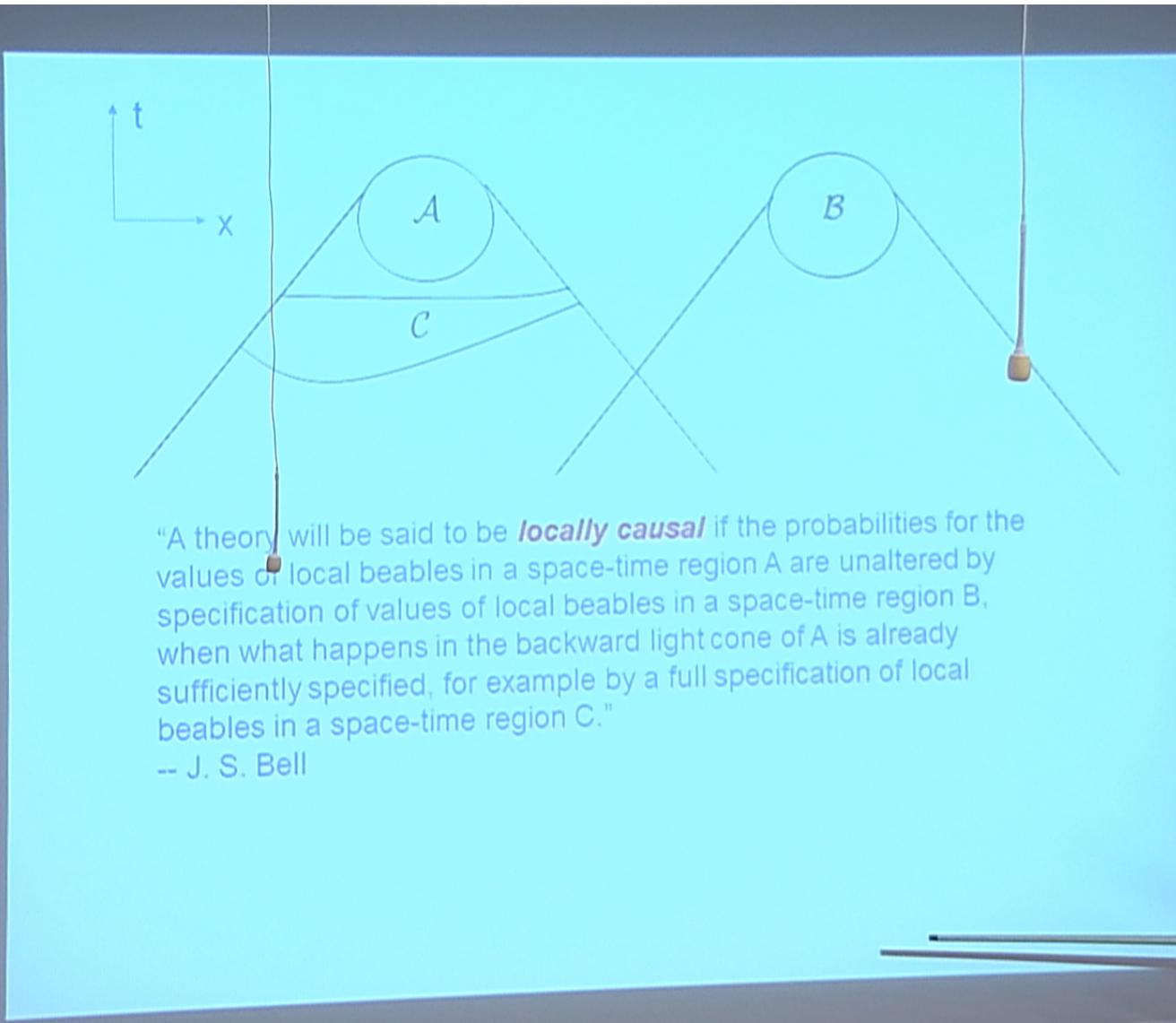
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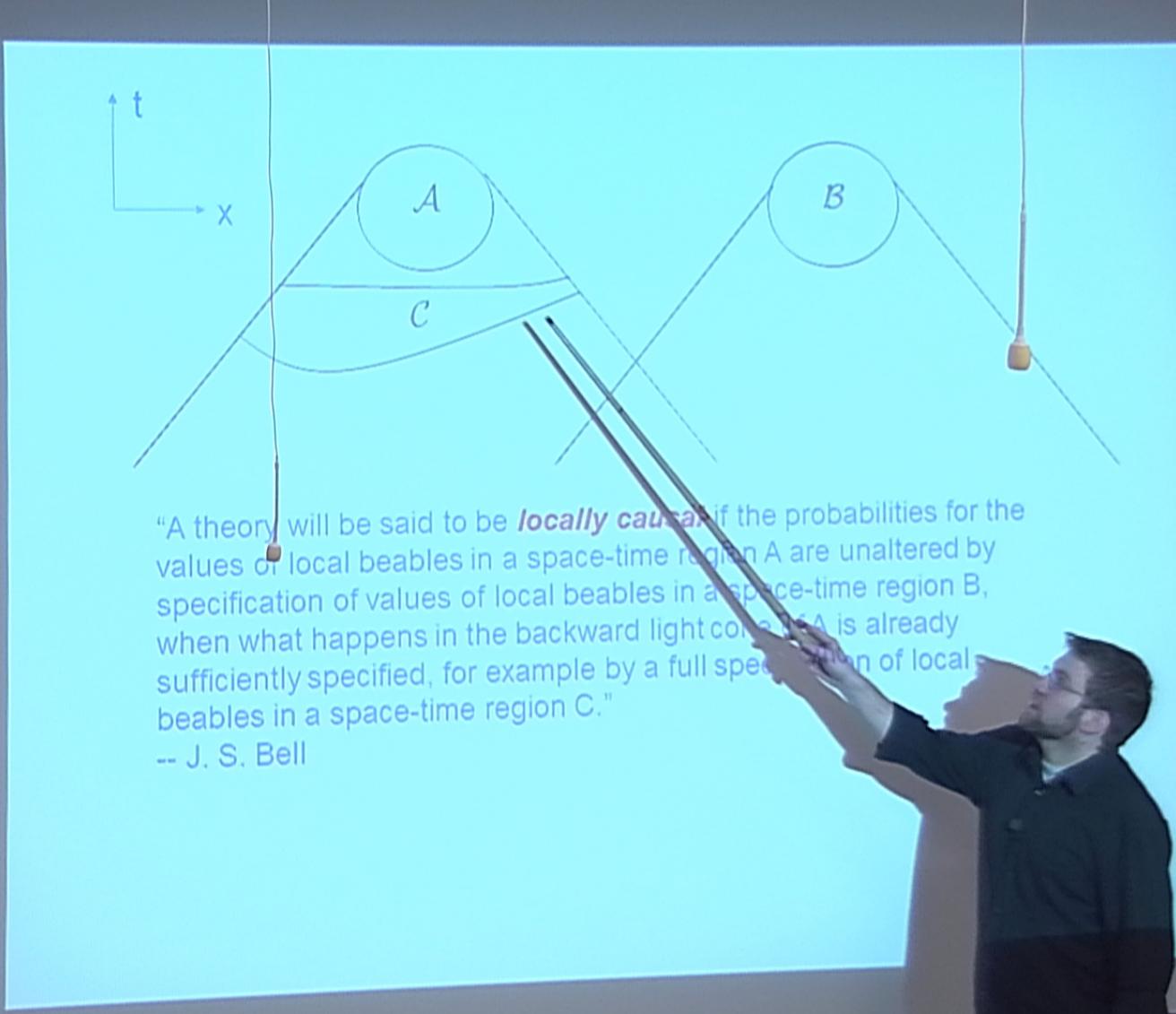
URL: <http://pirsa.org/12010048>

Abstract:

"The [beables/ontic parameters] in any space-time region 1 are determined by those in any space region V , at some time t , which fully closes the backward light cone of 1. Because the region V is limited, localized, we will say the theory exhibits *local determinism*.
-- J.S. Bell

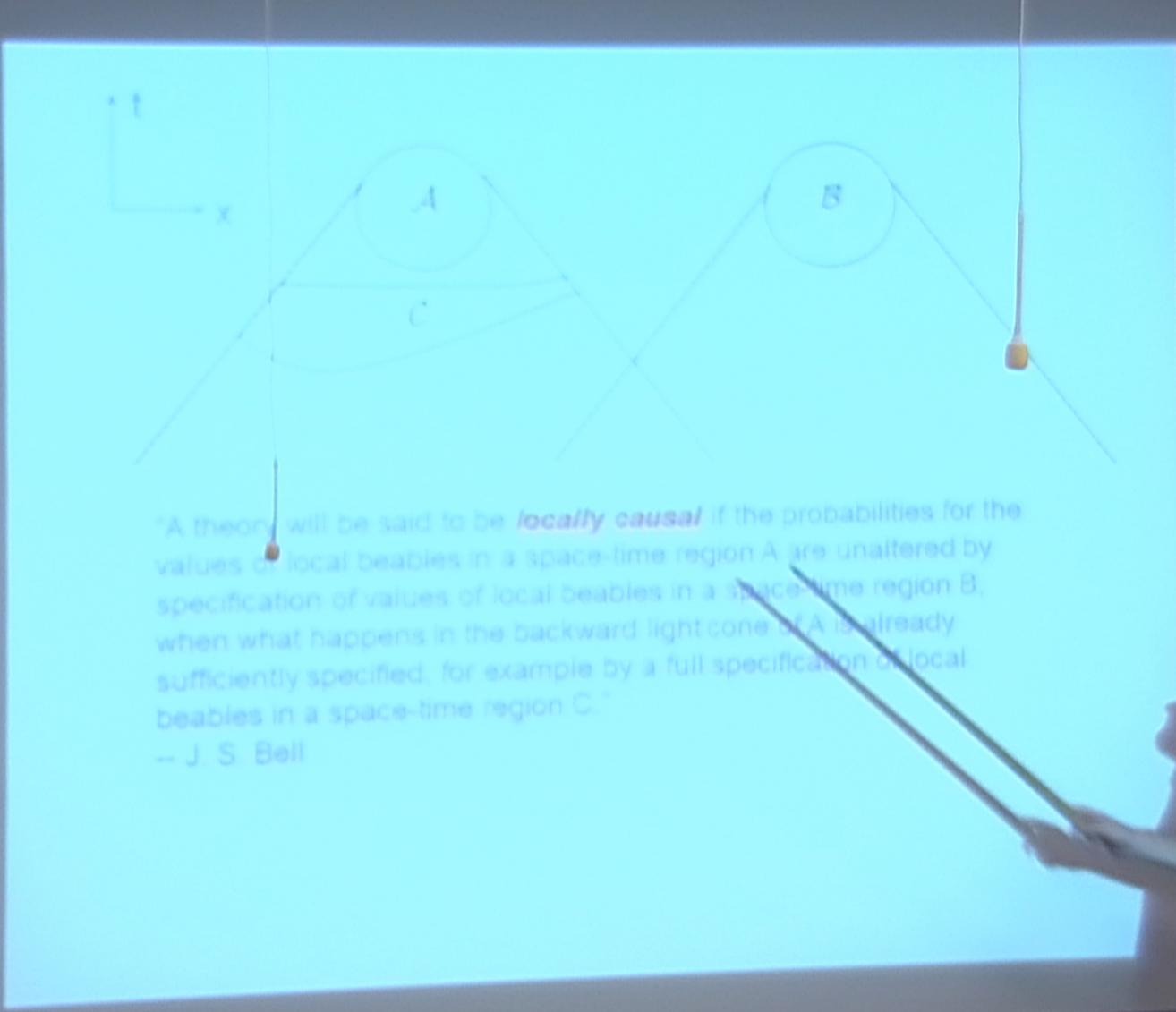






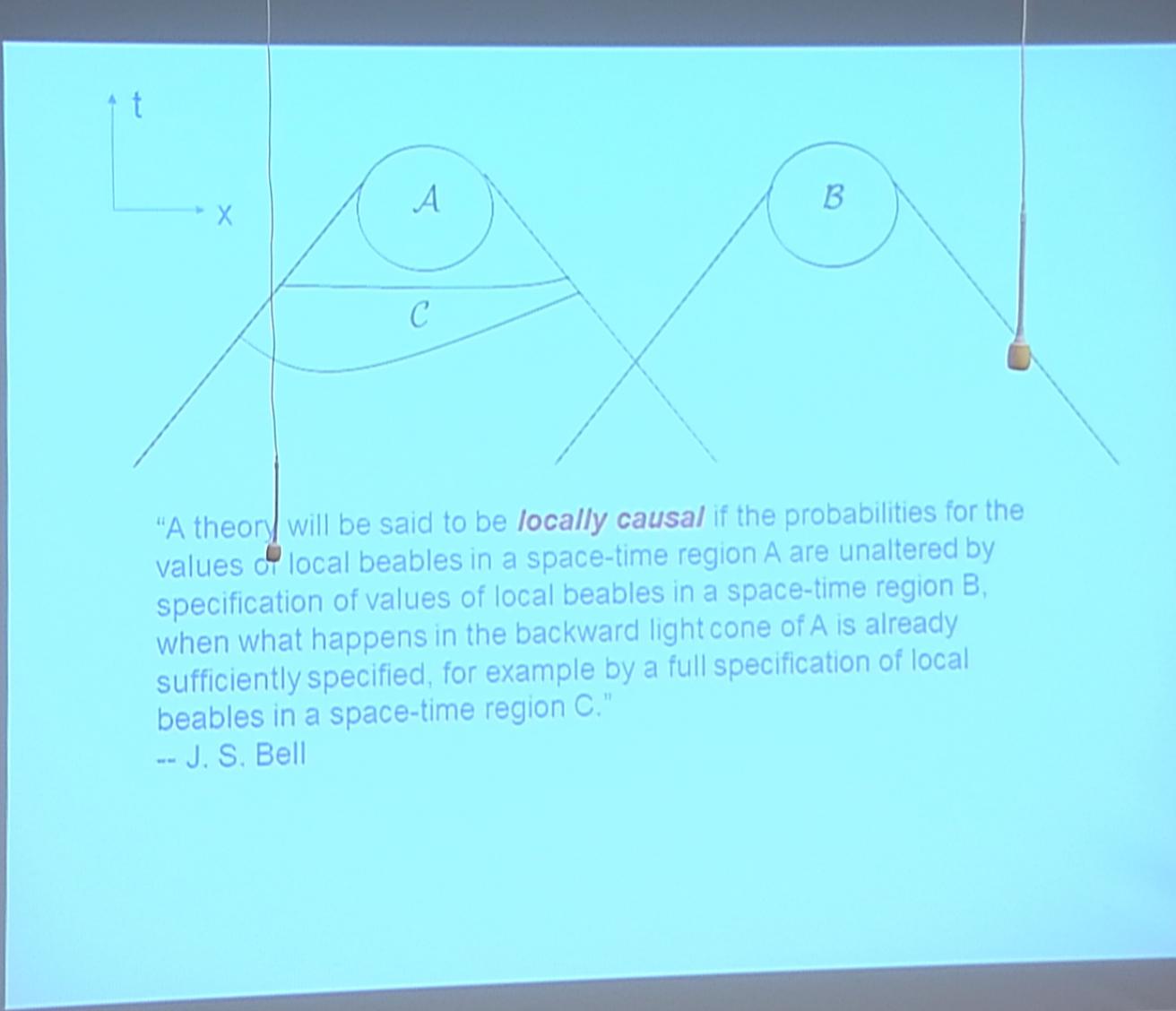
"A theory will be said to be *locally causal* if the probabilities for the values of local beables in a space-time region A are unaltered by specification of values of local beables in a space-time region B, when what happens in the backward light cone of A is already sufficiently specified, for example by a full specification of local beables in a space-time region C."

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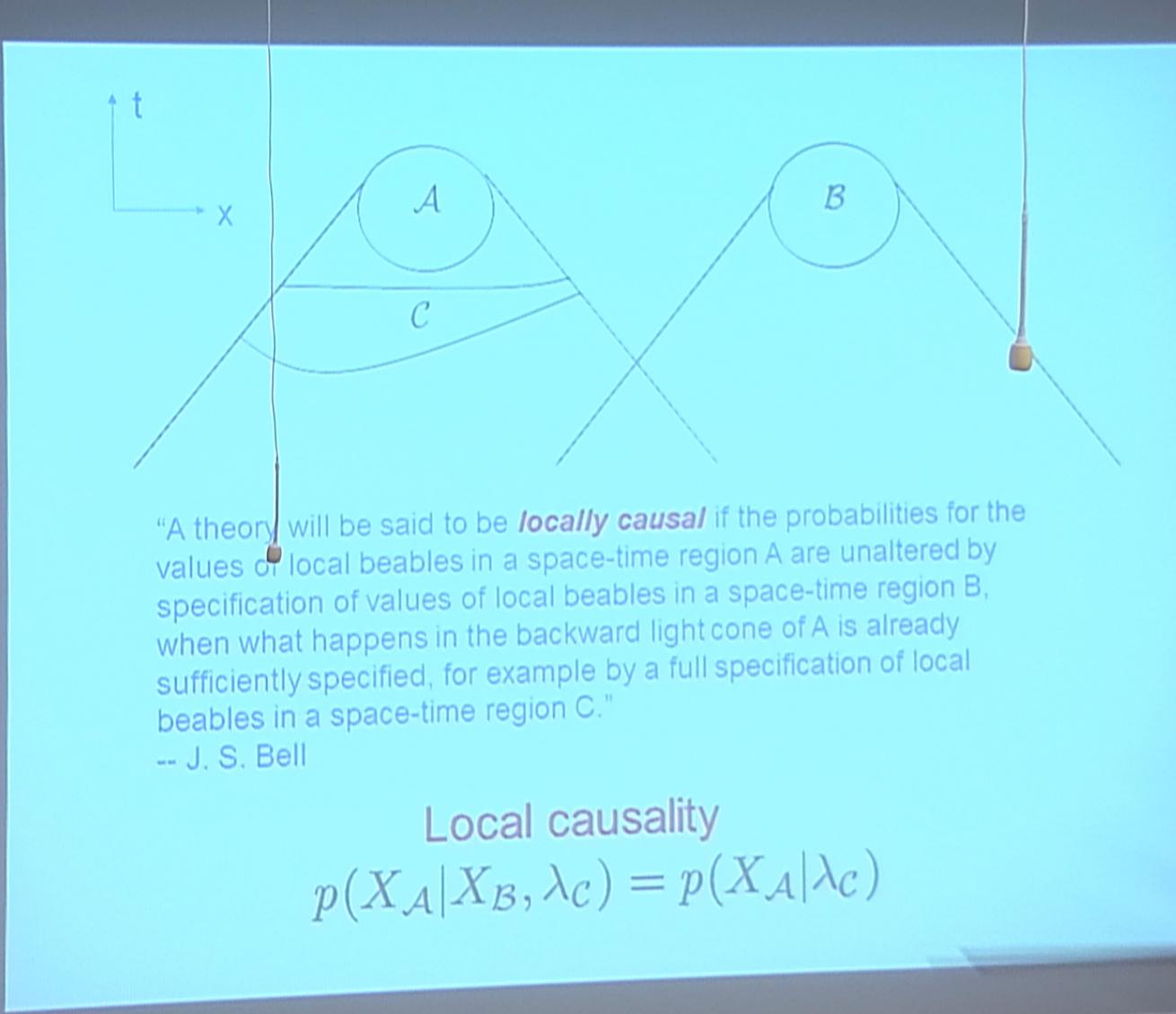
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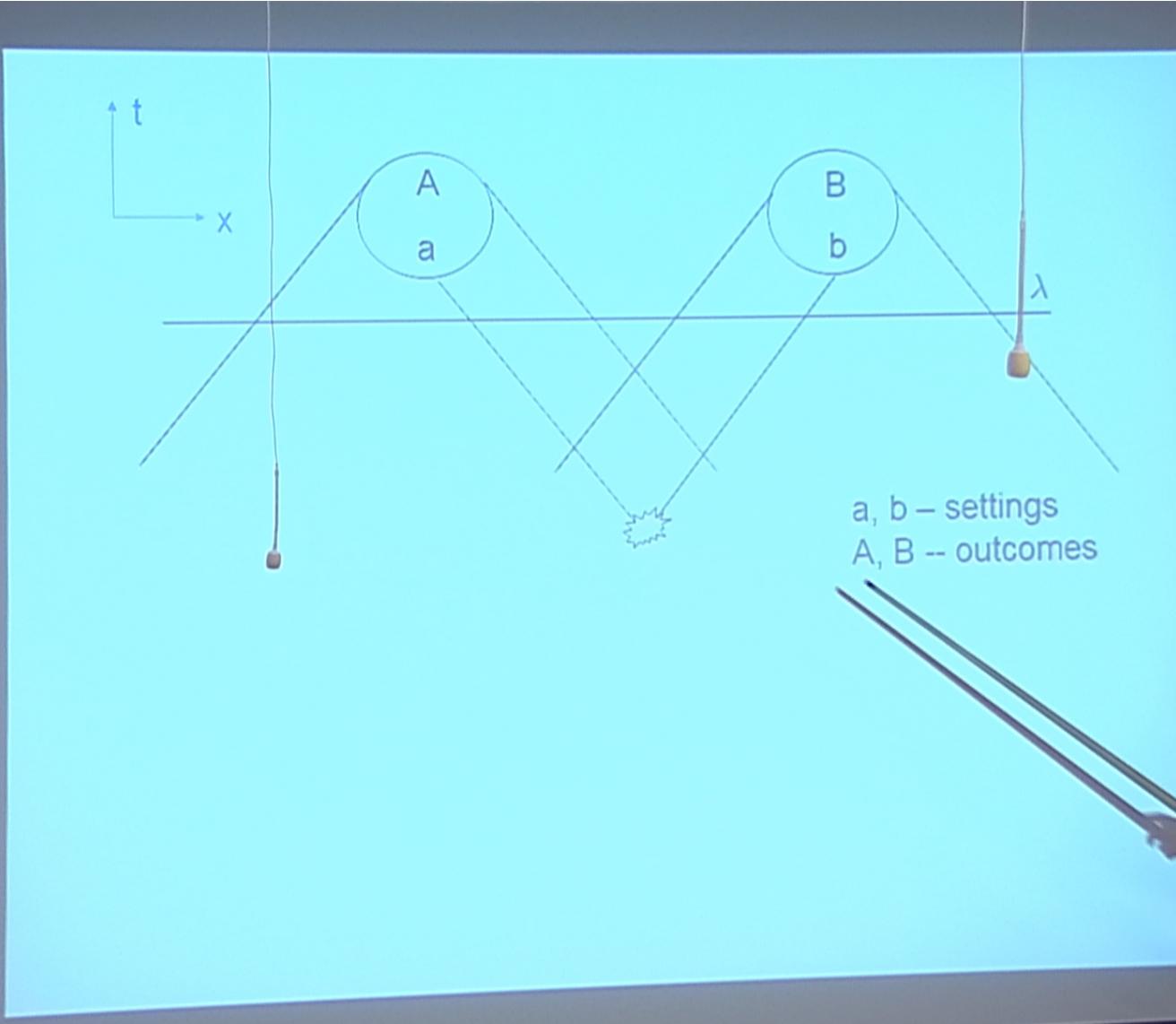


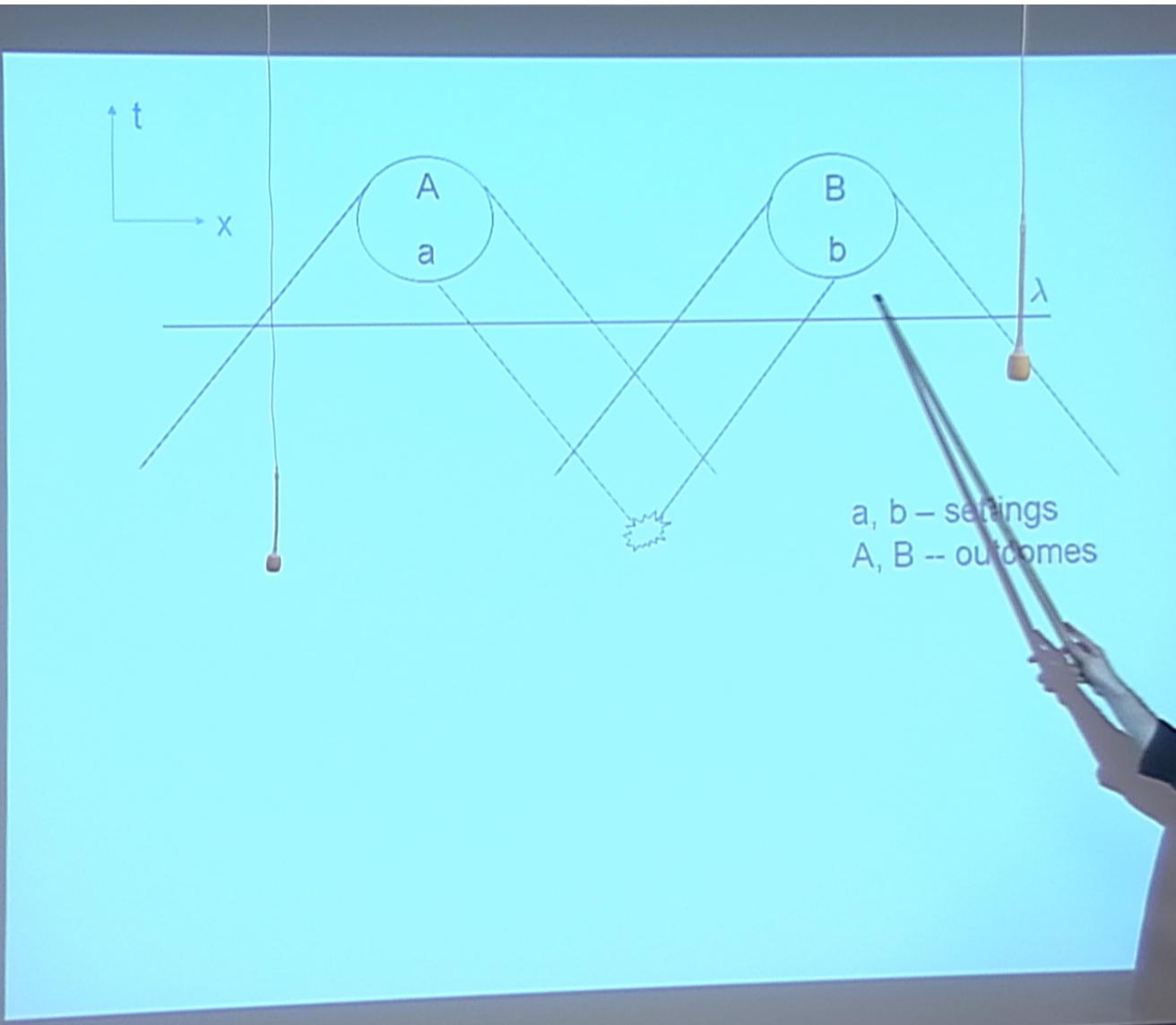
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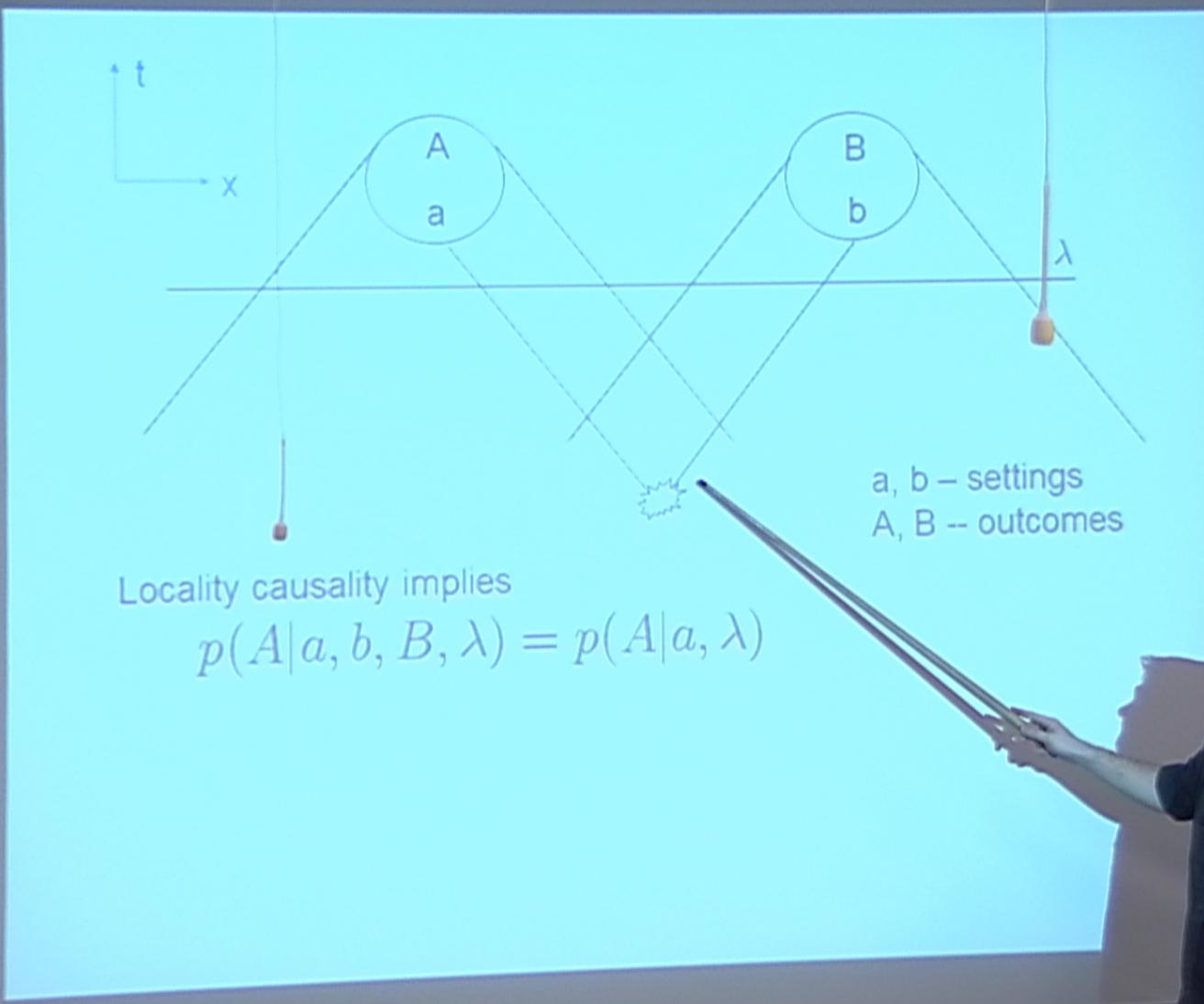
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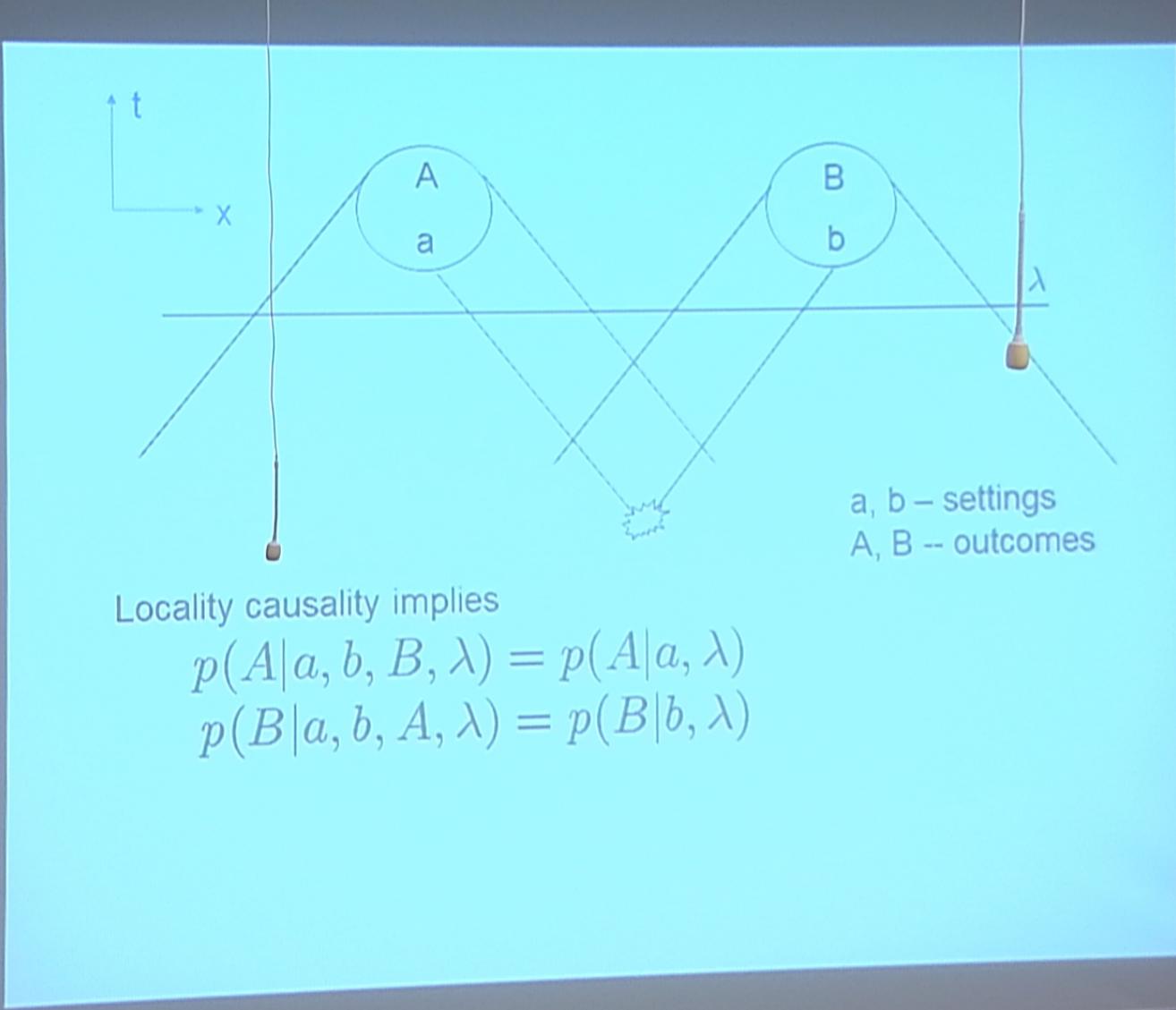
Local causality

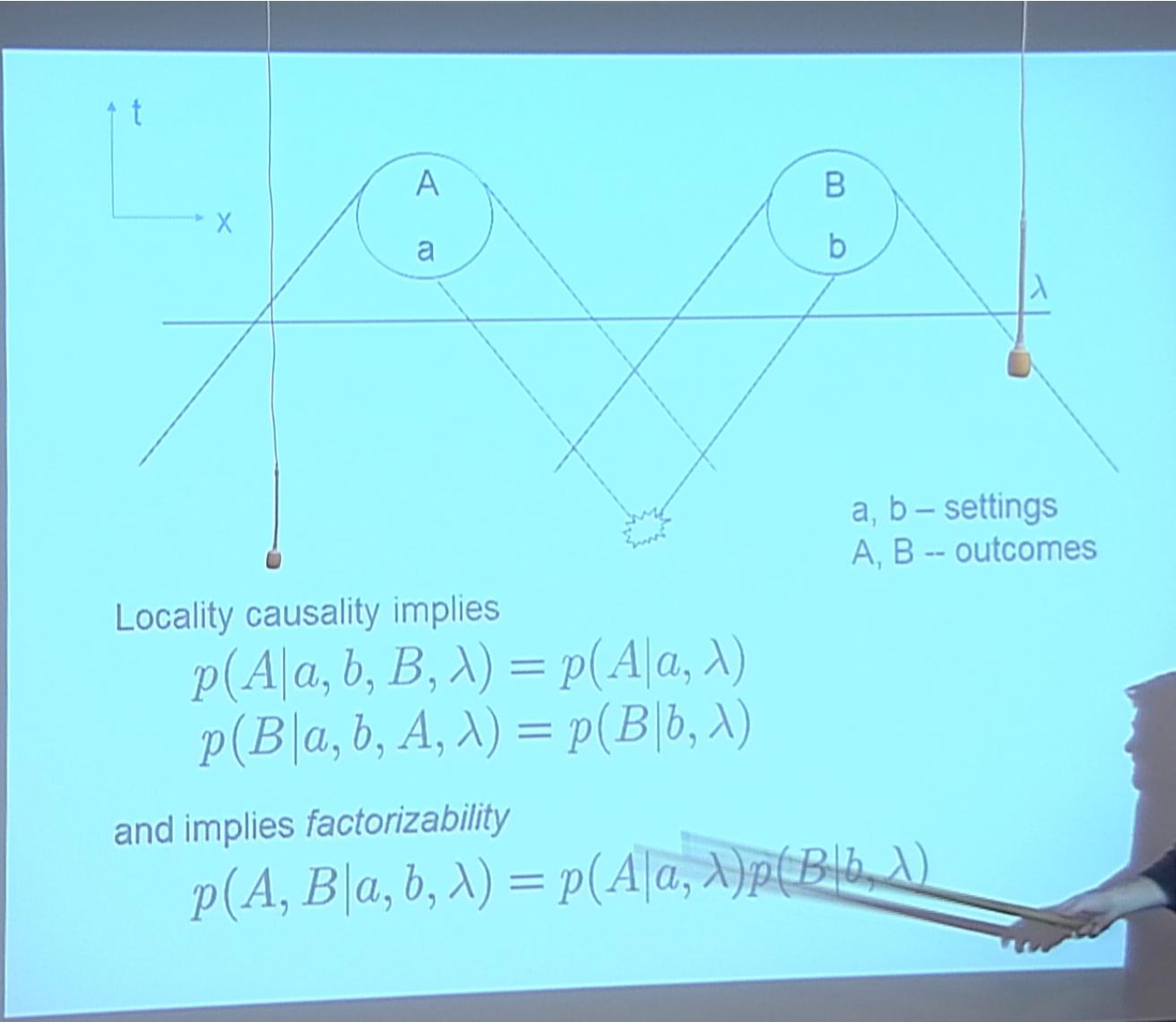
$$p(X_A | X_B, \lambda_C) = p(X_A | \lambda_C)$$

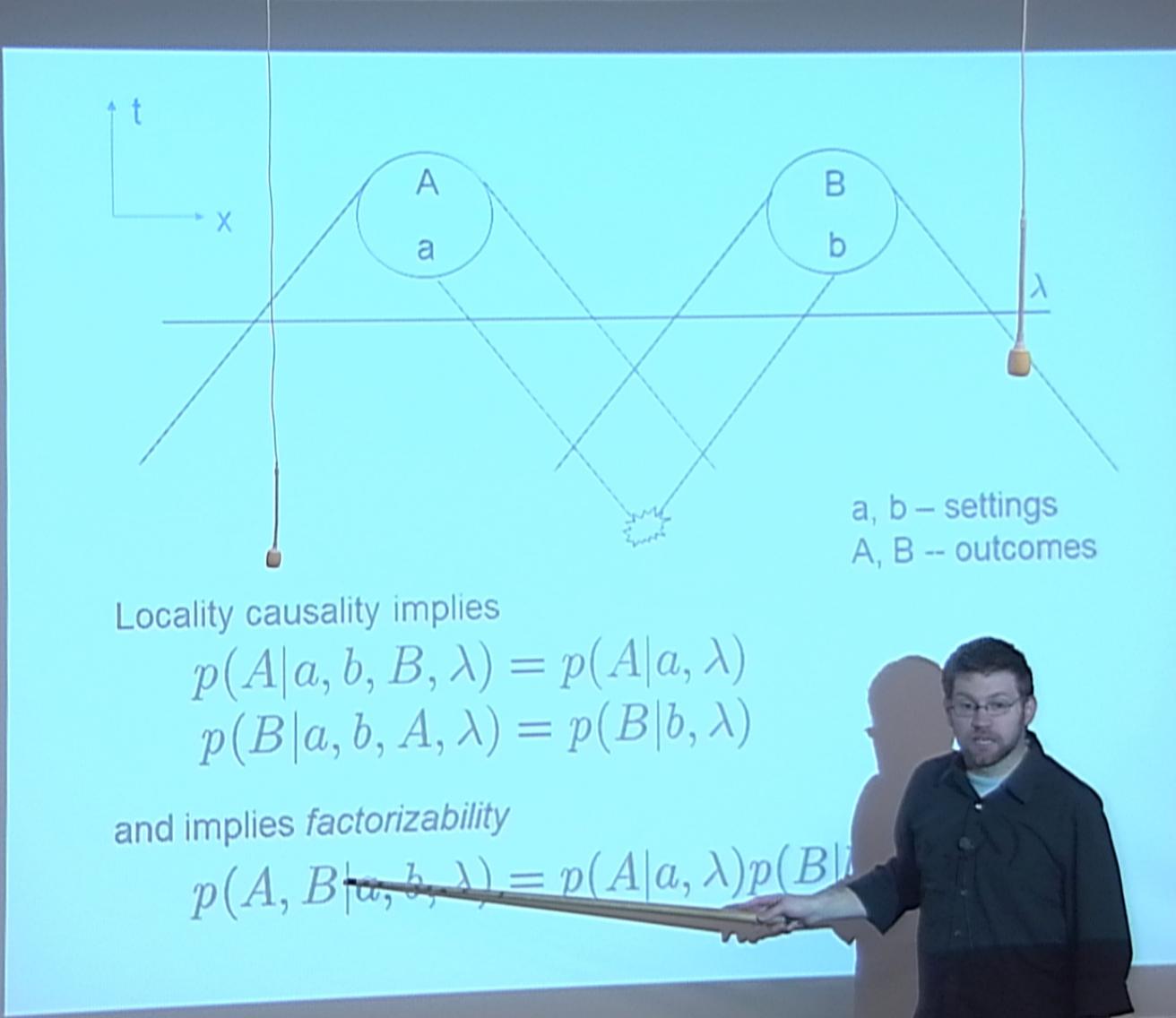


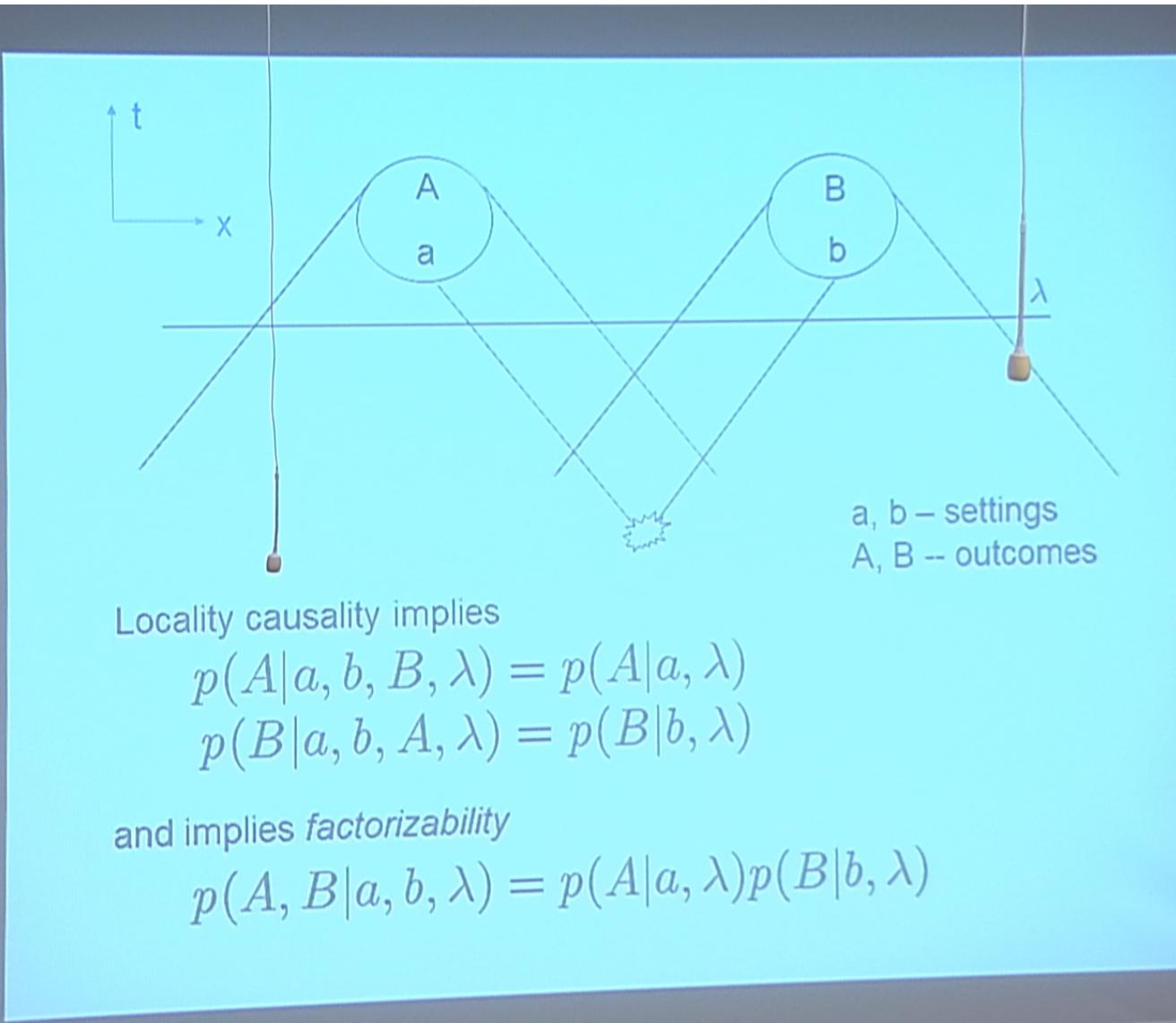












Factorizability from local causality

Recall the relation between joint and conditional probability

$$p(A, B) = p(A|B)p(B)$$

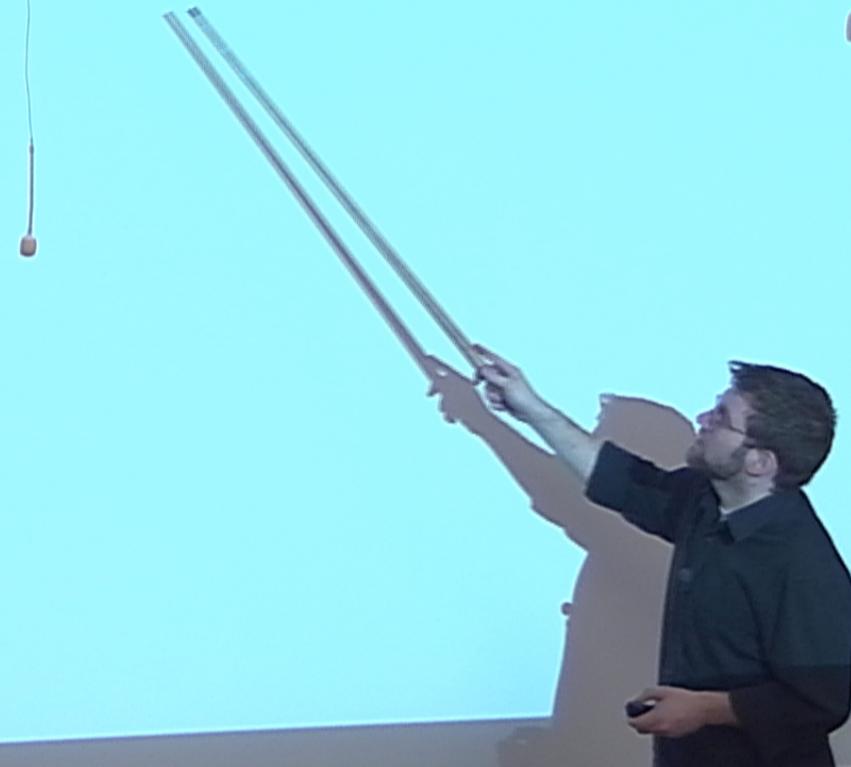
$$p(A, B|C) = p(A|B, C)p(B|C)$$

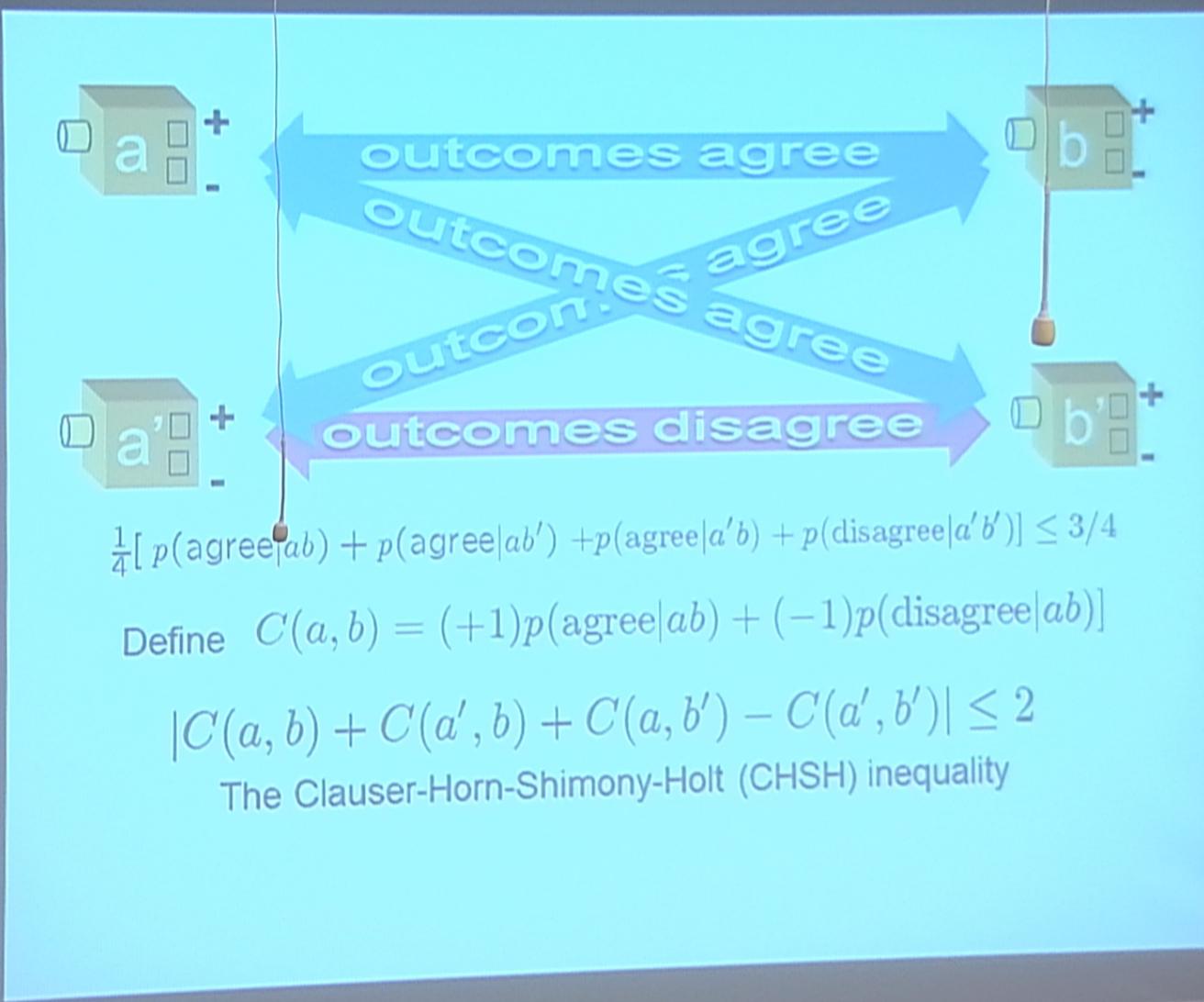
therefore

$$p(A, B|a, b, \lambda) = p(A|B, a, b, \lambda)p(B|a, b, \lambda)$$

Freedom of settings (no superdeterminism)

$$p(a, b, \lambda) = p(a)p(b)p(\lambda)$$

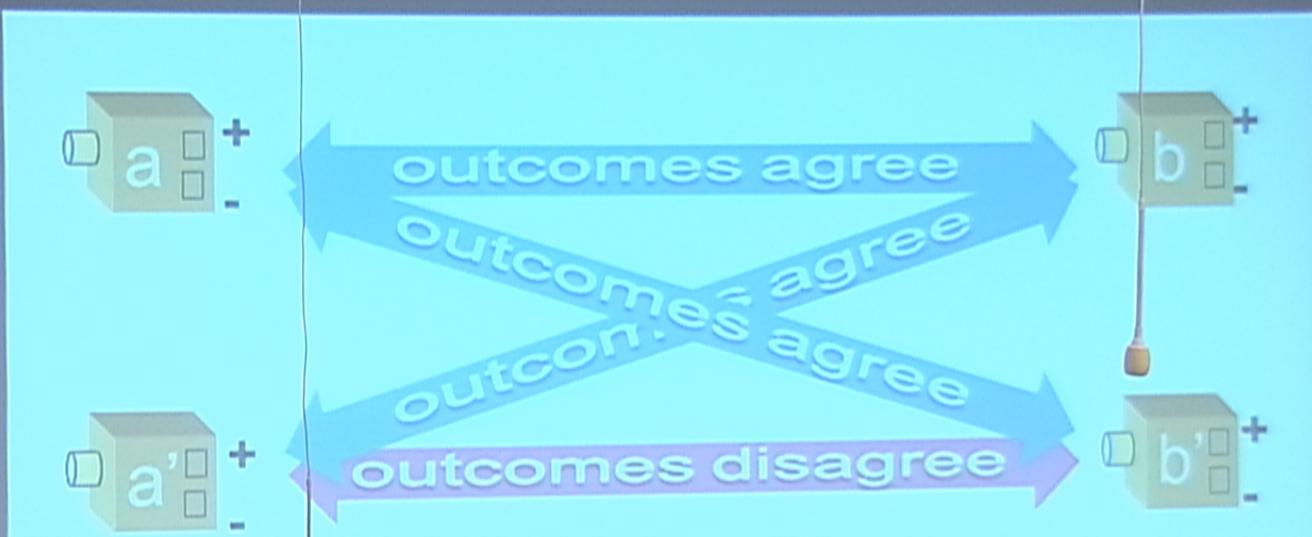




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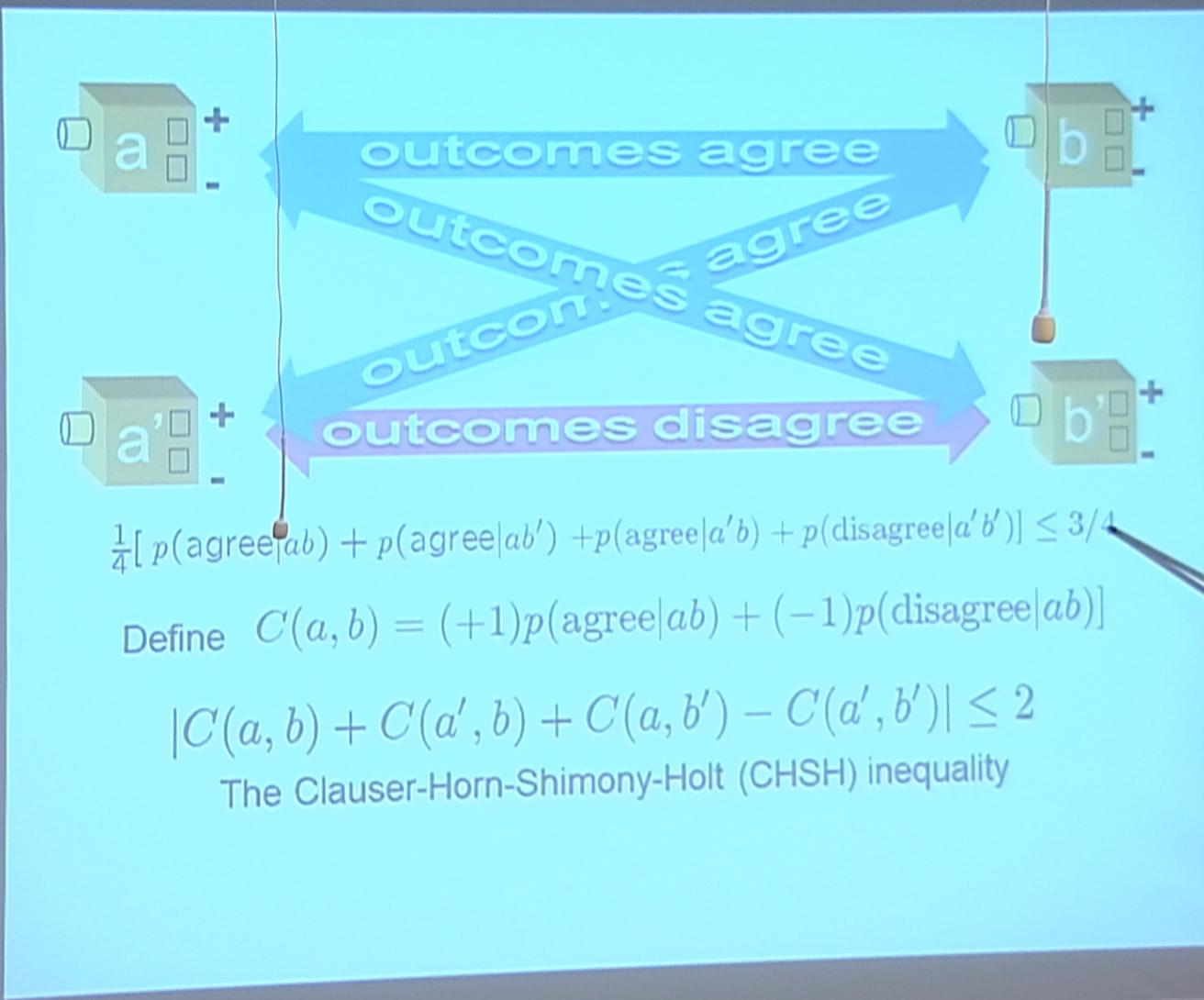


$$\frac{1}{4} [p(\text{agree}|ab) + p(\text{agree}|ab') + p(\text{agree}|a'b) + p(\text{disagree}|a'b')] \leq 3/4$$

Define $C(a, b) = (+1)p(\text{agree}|ab) + (-1)p(\text{disagree}|ab)$

$$|C(a, b) + C(a', b) + C(a, b') - C(a', b')| \leq 2$$

The Clauser-Horn-Shimony-Holt (CHSH) inequality



Locality causality: $p(A|a, b, B, \lambda) = p(A|a, \lambda)$
 $p(B|a, b, A, \lambda) = p(B|b, \lambda)$

No superluminal signalling: $p(A|a, b) = p(A|a)$
 $p(B|a, b) = p(B|b)$

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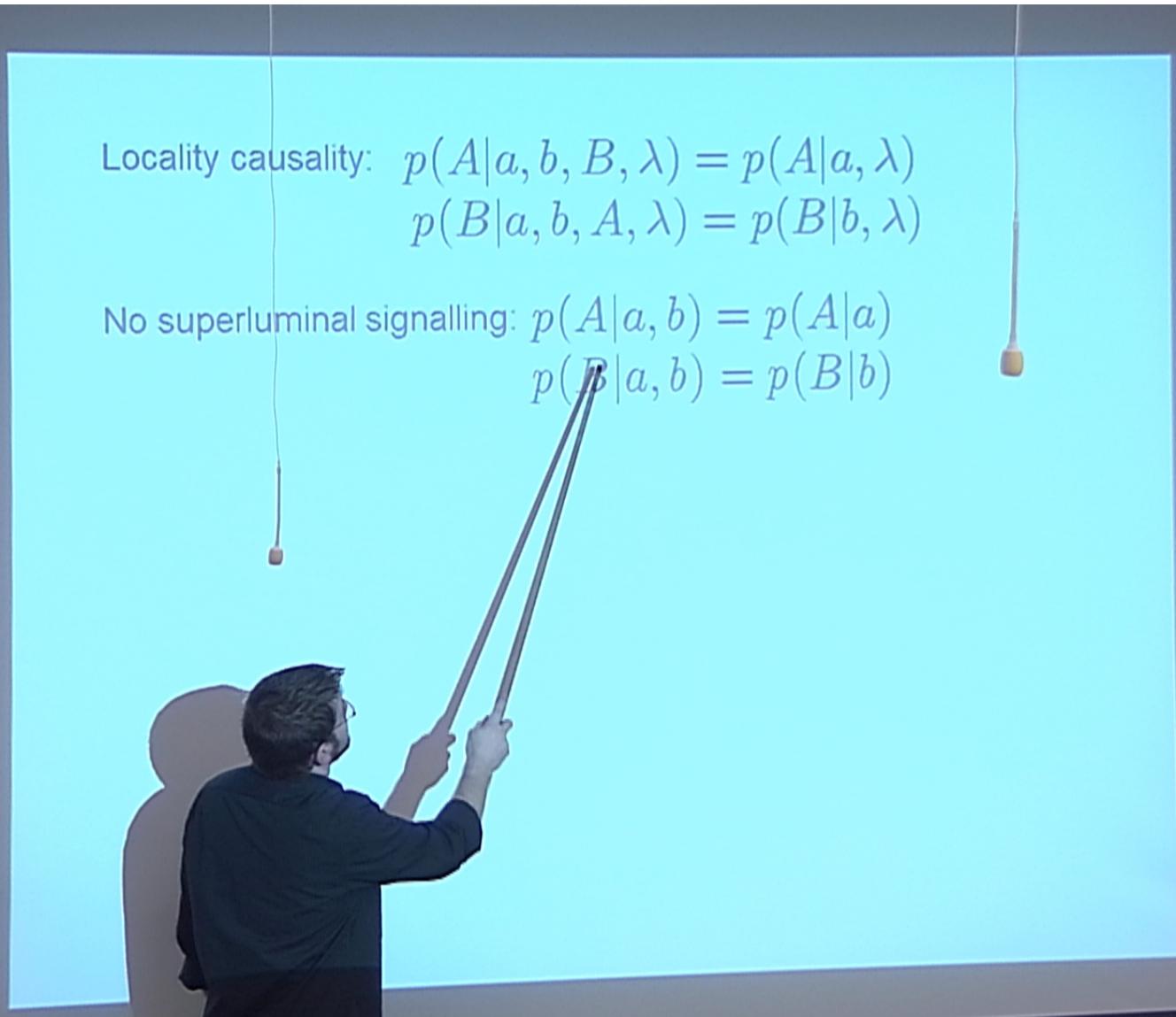
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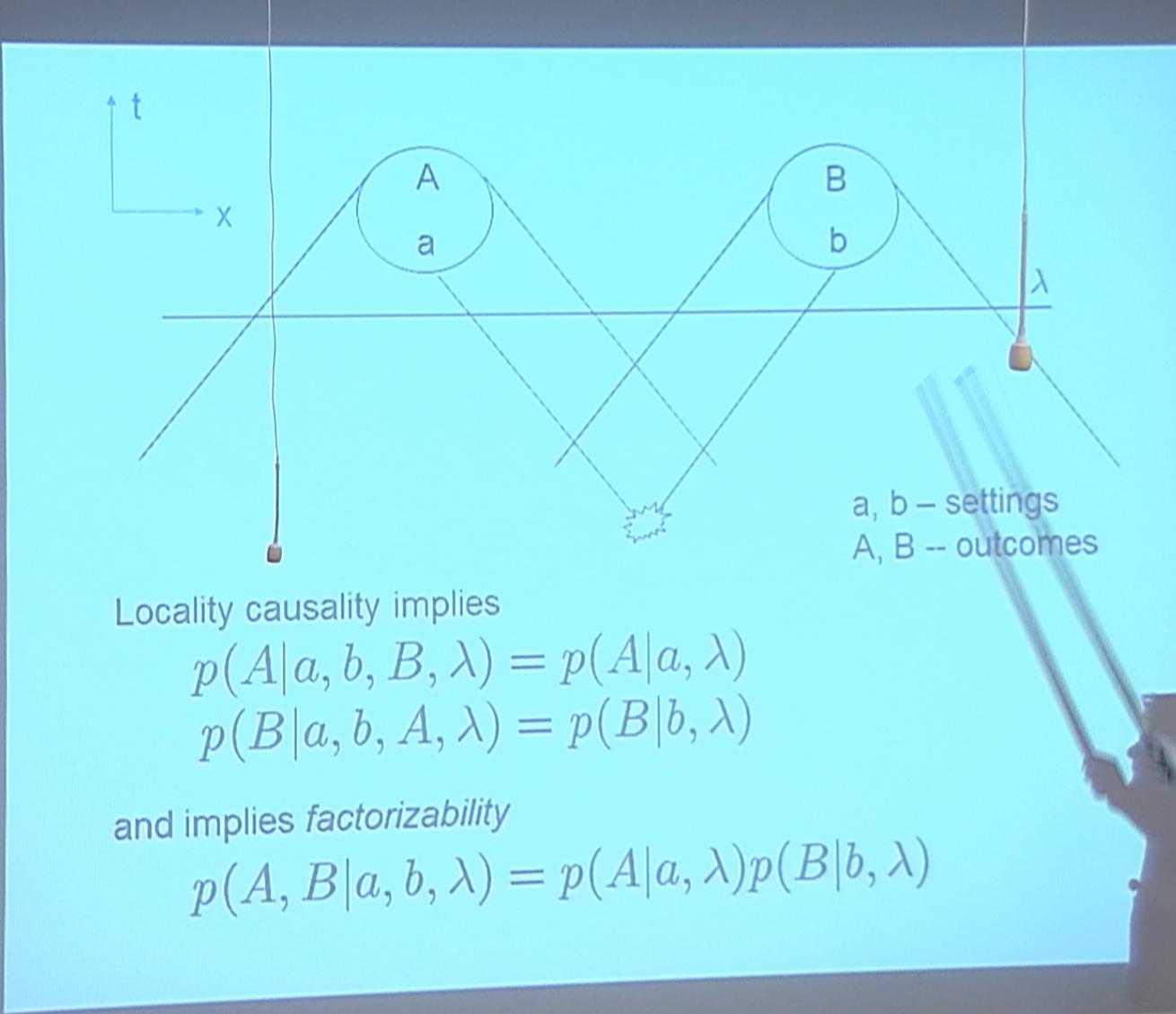
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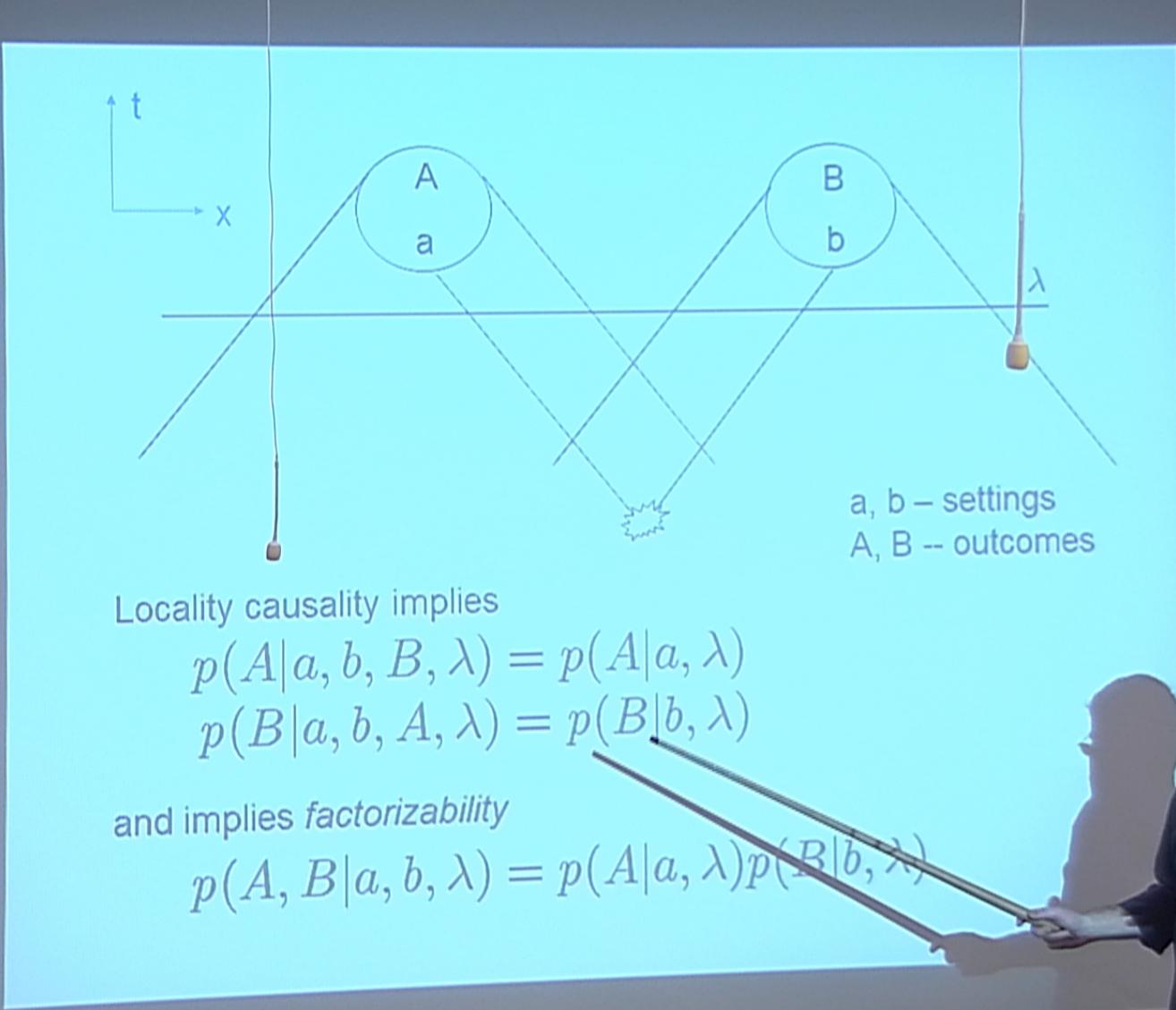
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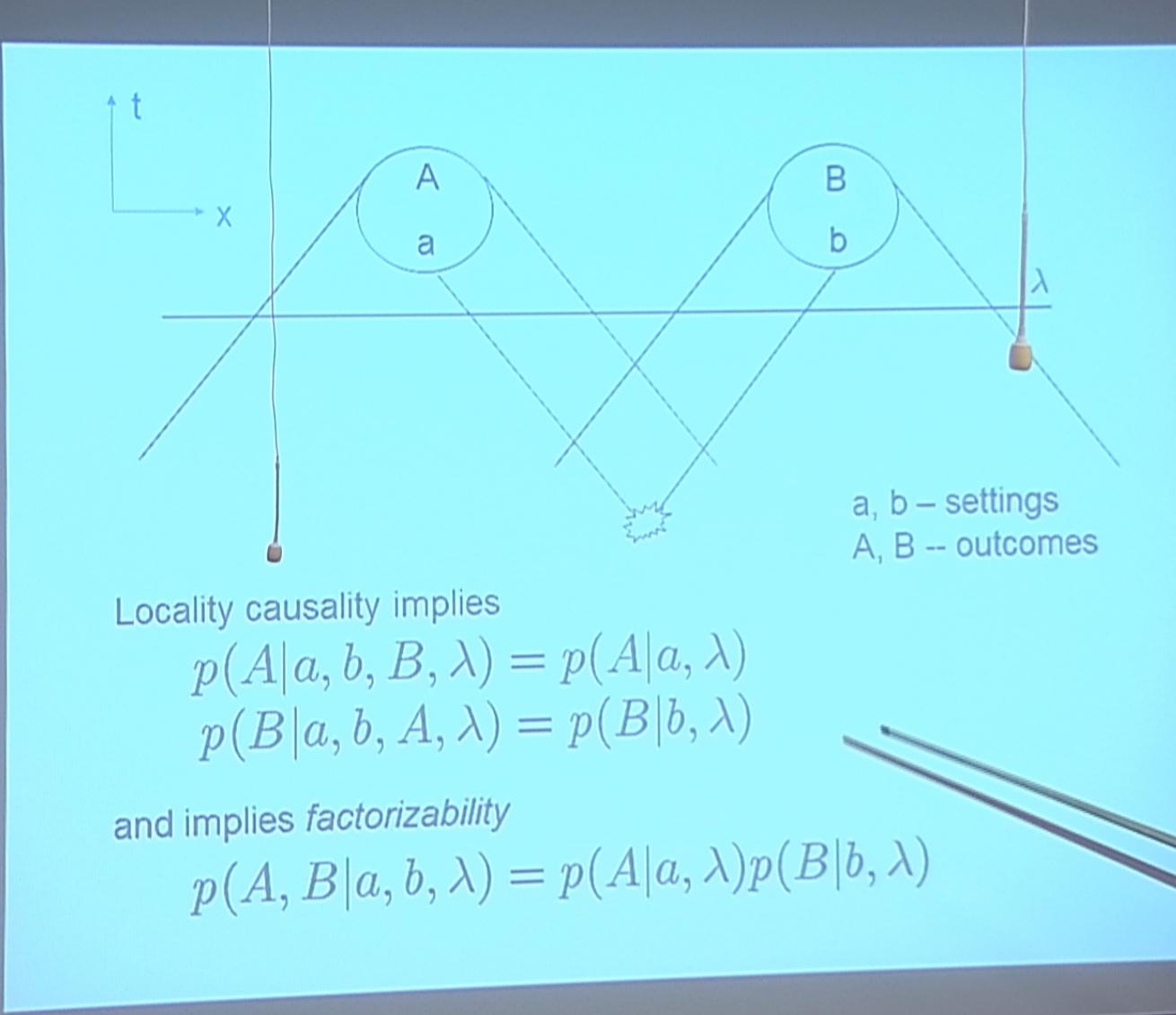
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Does the notion of "local causality" capture the content of relativity?







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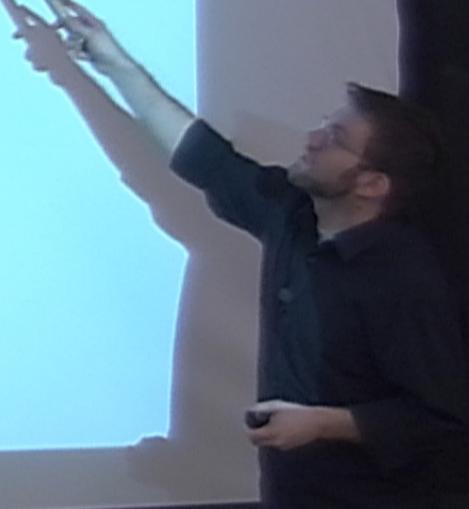
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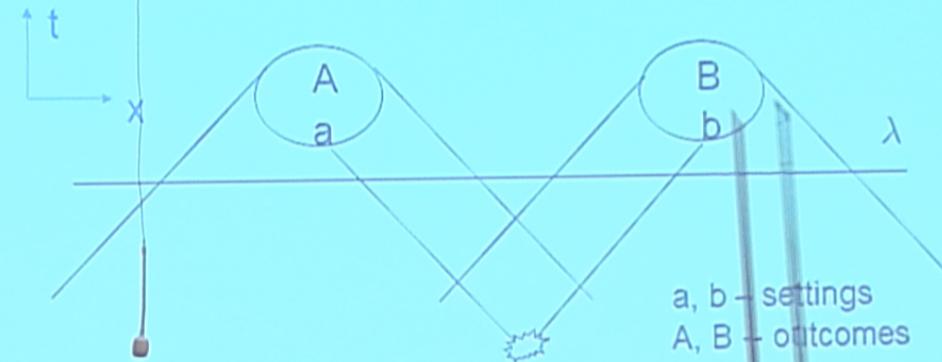
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Does the notion of "local causality" capture the content of relativity?

Failure of predictability from Bell-inequality violations and no signalling

Acin, Masanes and Gisin, Cavalcanti, Seevinck, Aharonov



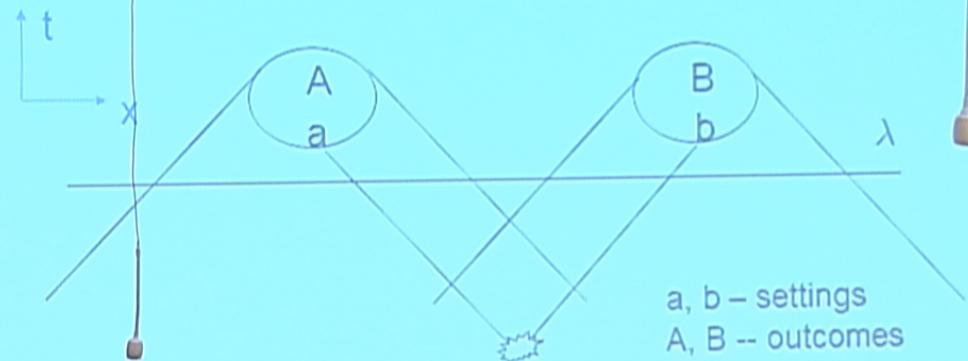
Predictability: $p(A, B|a, b) \in \{0, 1\}$

No signalling: $p(A|a, b) = p(A|a)$ and $p(B|a, b) = p(B|b)$

Thm: No signalling + Bell-inequality violation \rightarrow unpredictability

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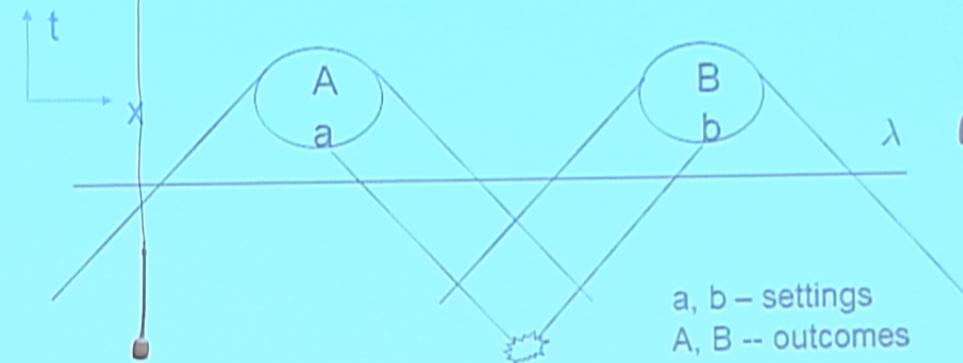
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a, b – settings
A, B -- outcomes

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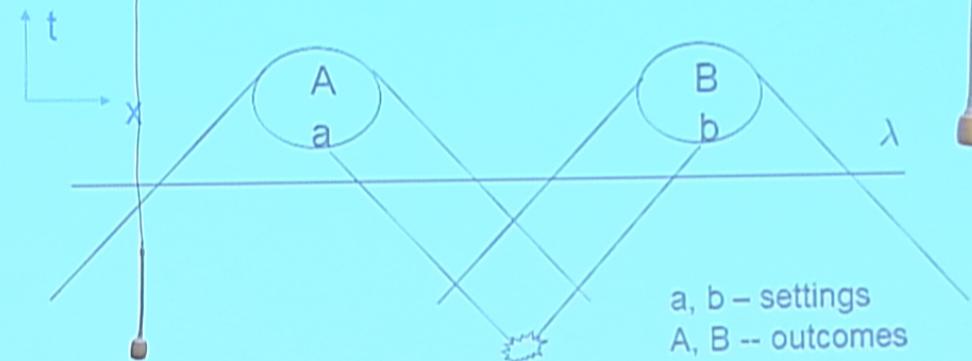
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Proof: $p(A, B|a, b, \lambda) = p(A, B|a, b) = p(A|a, b) p(B|a, b)$ (by predictability)



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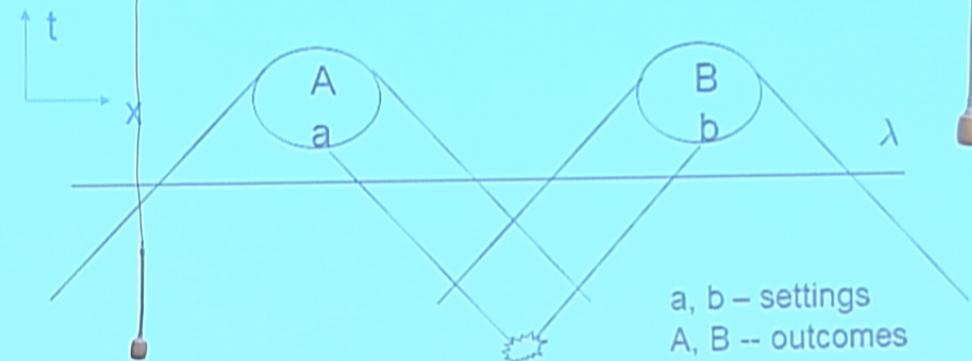
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Proof: $p(A, B|a, b, \lambda) = p(A, B|a, b) = p(A|a, b) p(B|a, b)$ (by predictability)
 $= p(A|a) p(B|b)$ (by no signalling)

But this is factorizability, from which the Bell inequalities follow.

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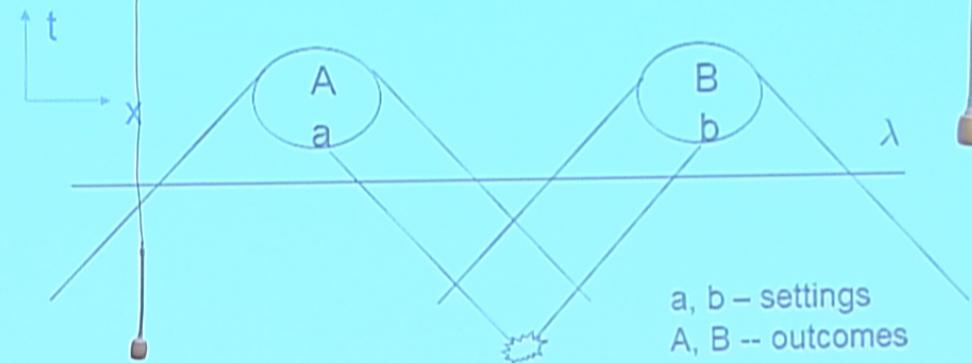
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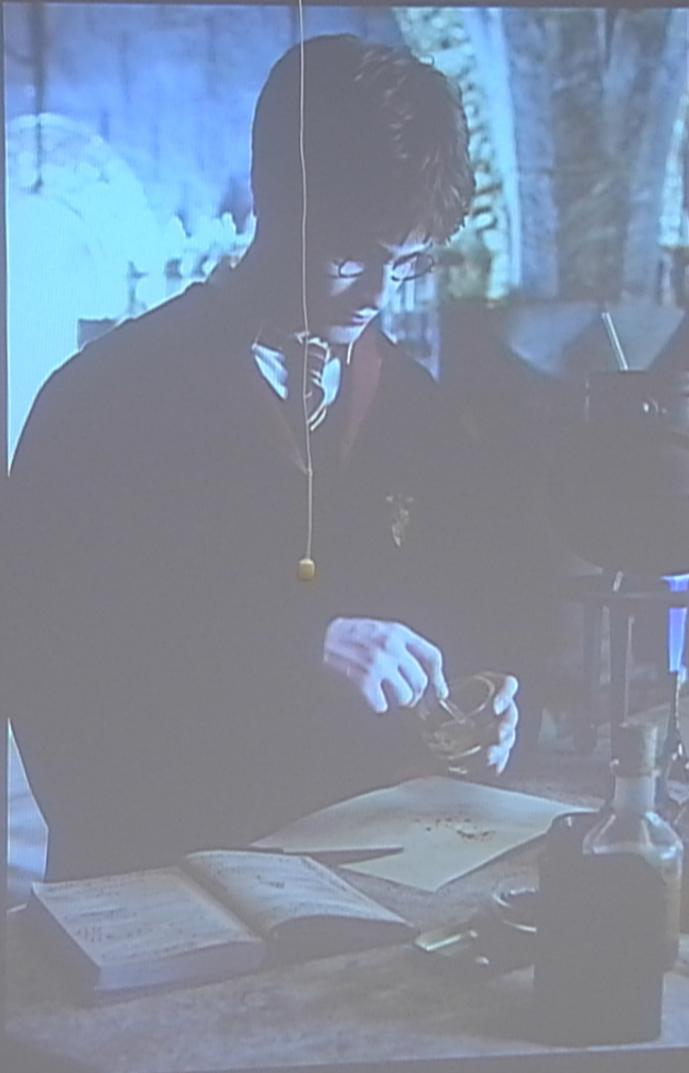
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Applications of nonlocality



Magic is a natural force that can be used to override the usual laws of nature.

-- Harry Potter entry in wikipedia

Bell-inequality violations are natural phenomena that can be used to override the usual (classical-like) laws of nature

Quantum Spellcraft

Based on Bell-inequality violation

Reduction in communication complexity

Buhrman, Cleve, van Dam, SIAM J.Comput. 30 1829 (2001)

Brassard, Found. Phys. 33, 1593 (2003)

Device-independent secure key distribution

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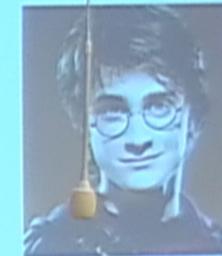
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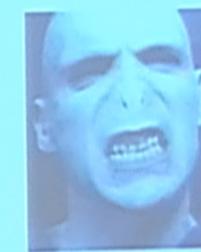
Monogamy of Bell-inequality violating correlations



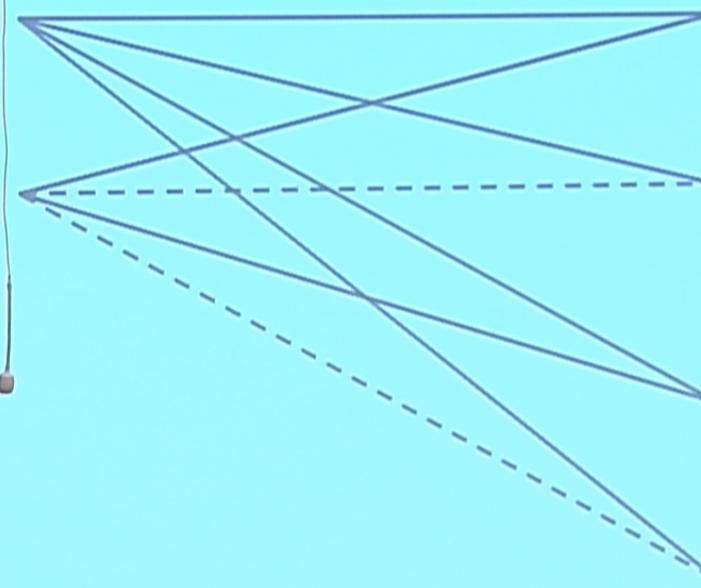
Alice



Bob



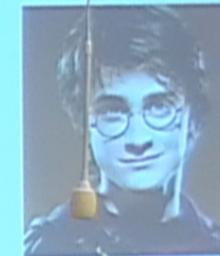
Adversary



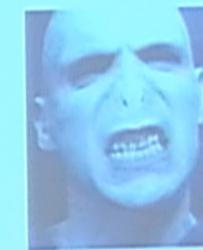
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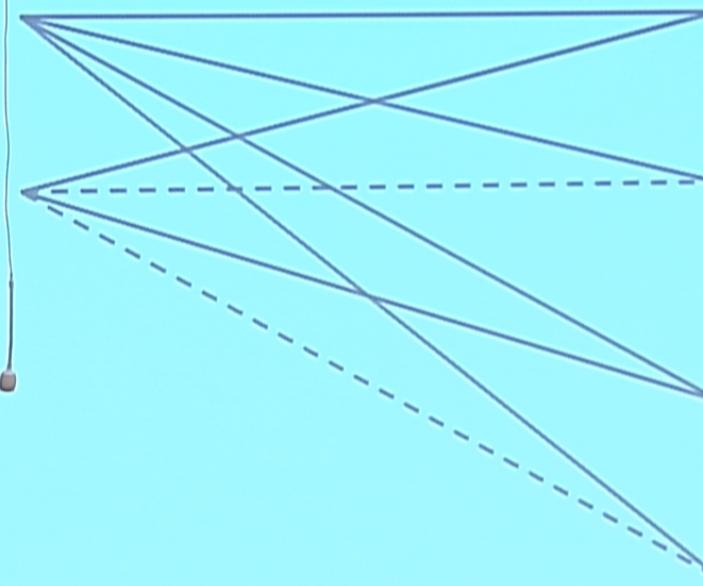
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Bob



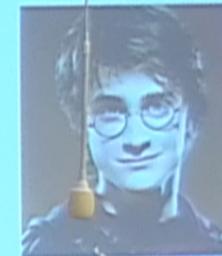
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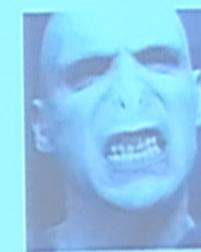
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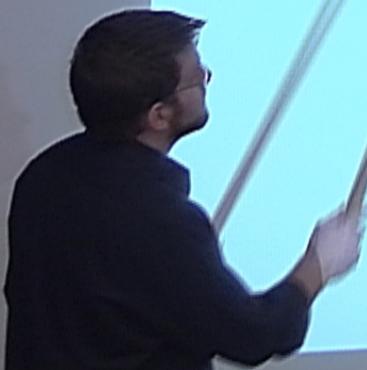
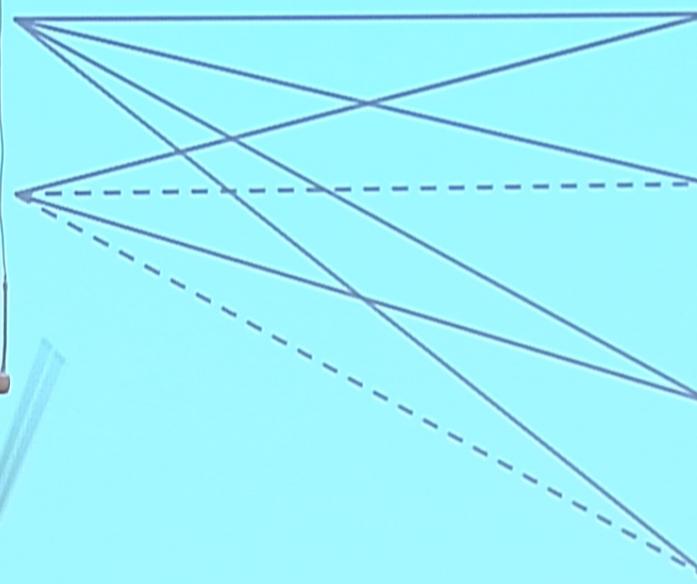
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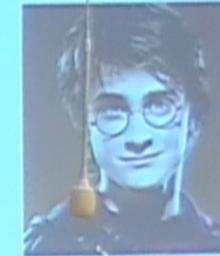
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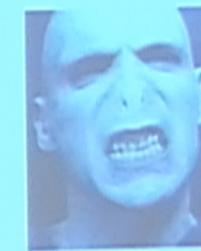
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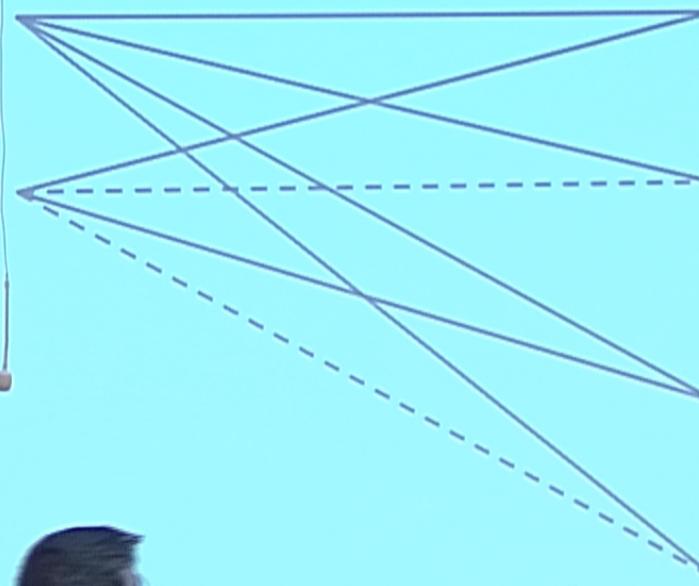
Alice



Bob



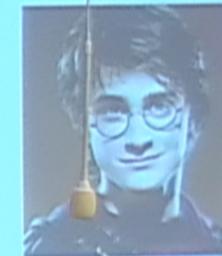
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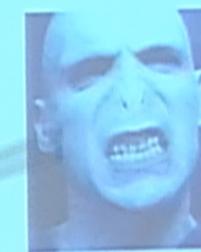
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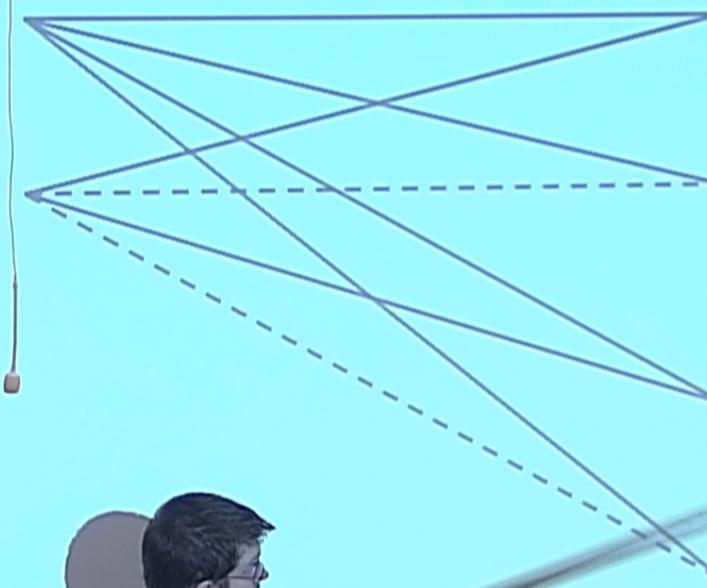
Alice



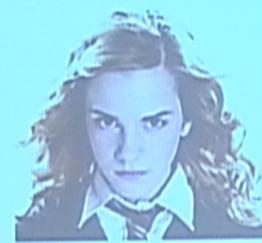
Bob



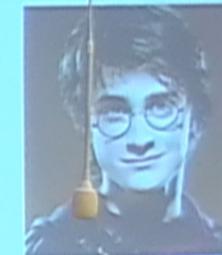
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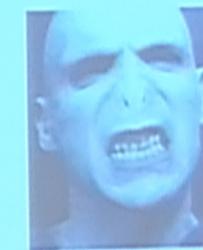
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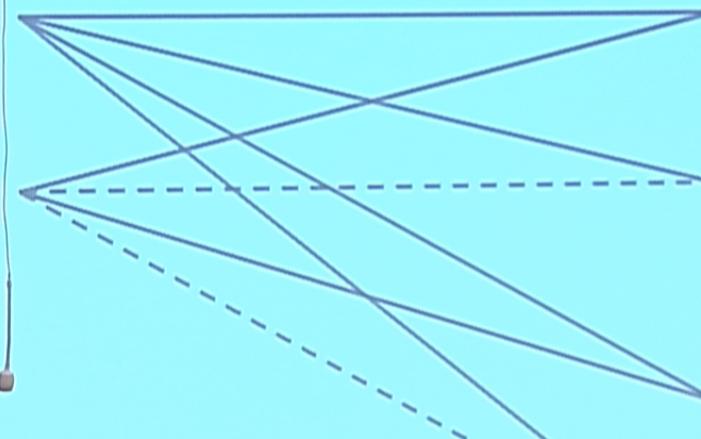
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Adversary



Recent trend in axiomatization:
Why isn't the world *more* nonlocal?

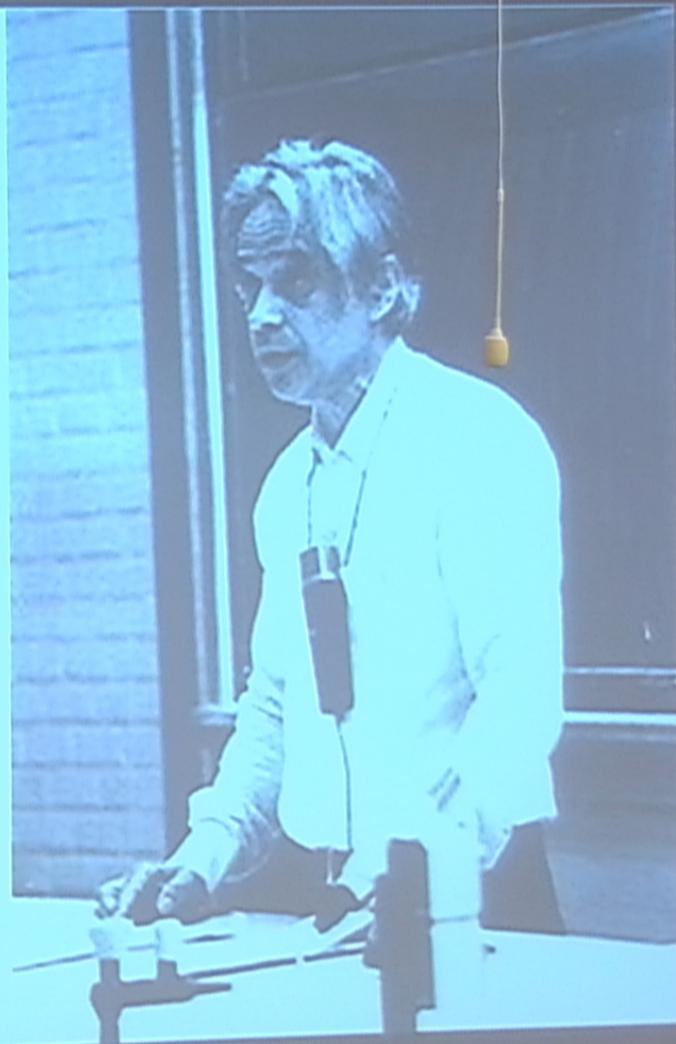
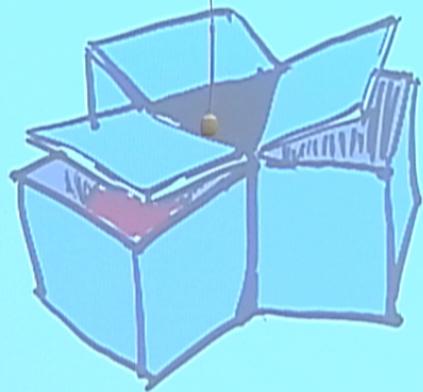
Recent trend in axiomatization:
Why isn't the world *more* nonlocal?



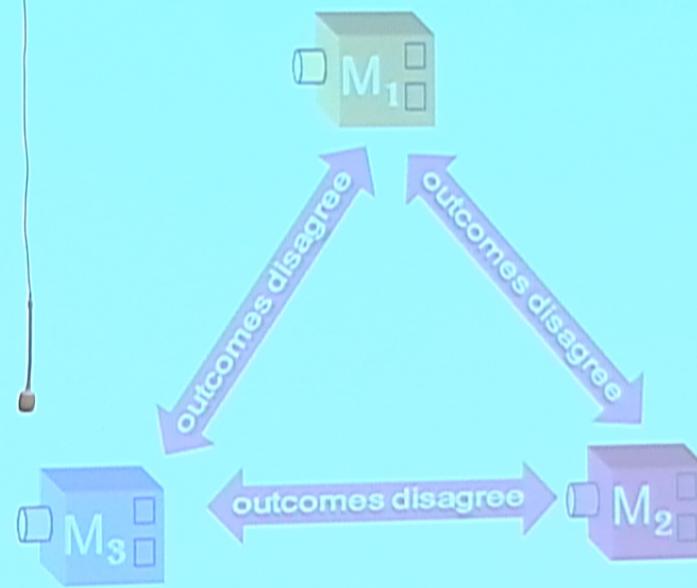
Contextuality

Contextuality

Ernst Specker, "The logic of propositions which are not simultaneously decidable", *Dialectica* 14, 239 (1960).

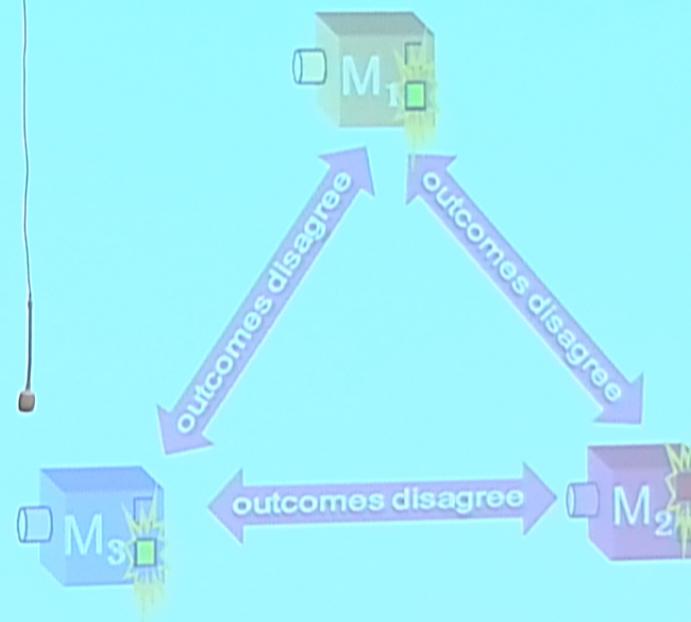


Specker's example



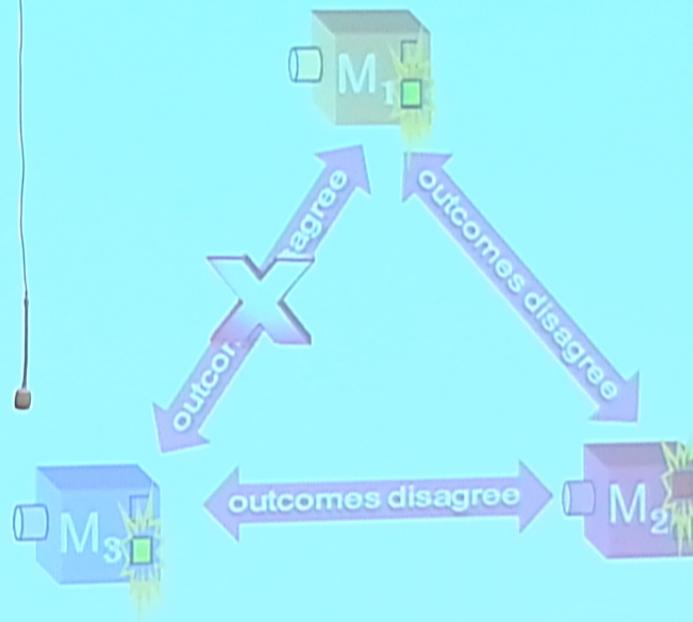
$$p(\text{success}) \leq \frac{2}{3}$$

Specker's example



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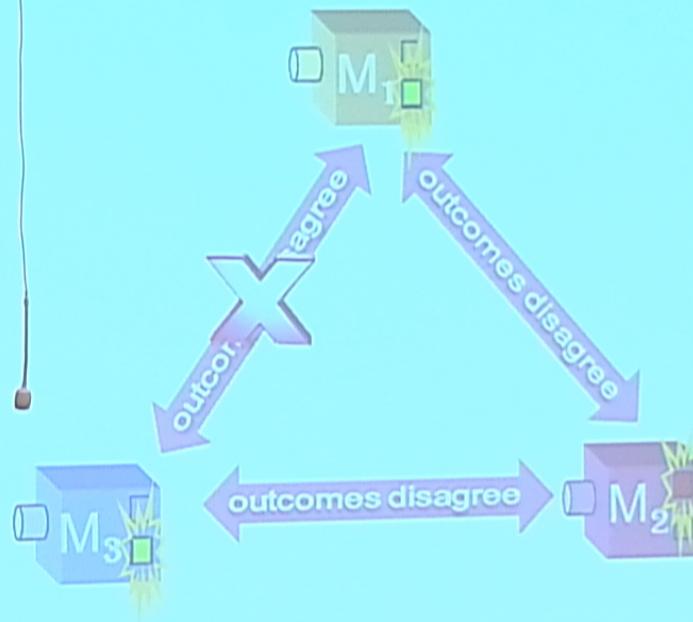
Specker's example



If the outcomes are fixed deterministically by the ontic state and are independent of the context in which the measurement is performed, then

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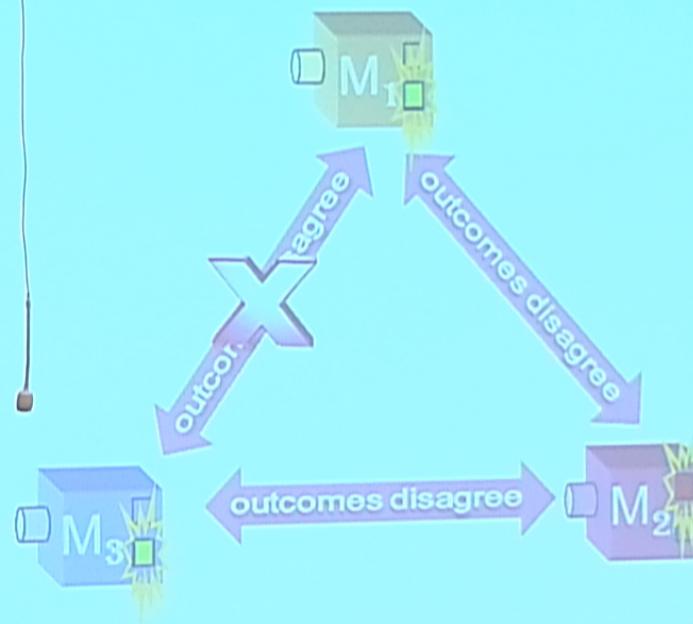
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Frustrated Networks

Nodes are binary variables

Edges imply joint measurability

- Outcomes agree -- Perfect correlation
- - - Outcomes disagree --Perfect anti-correlation

Frustration = no valuation satisfying all correlations



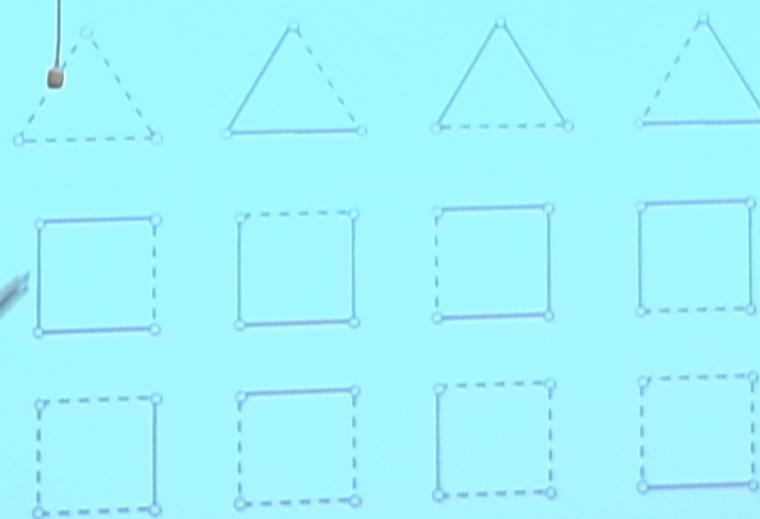
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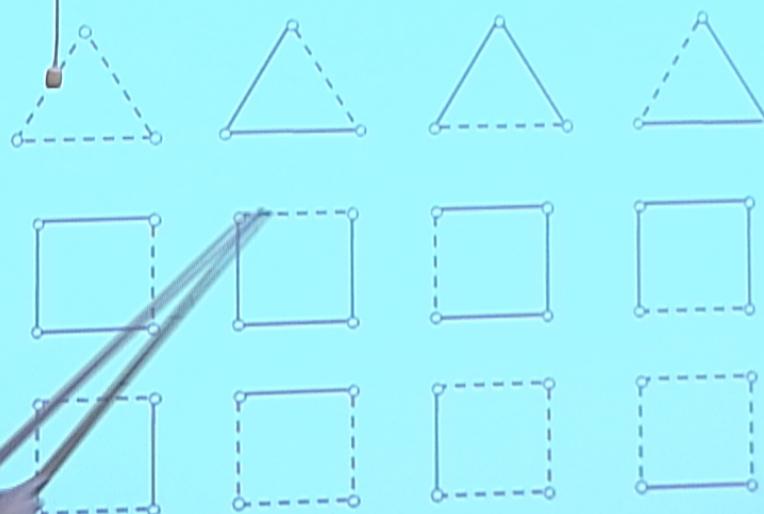
Frustrated Networks

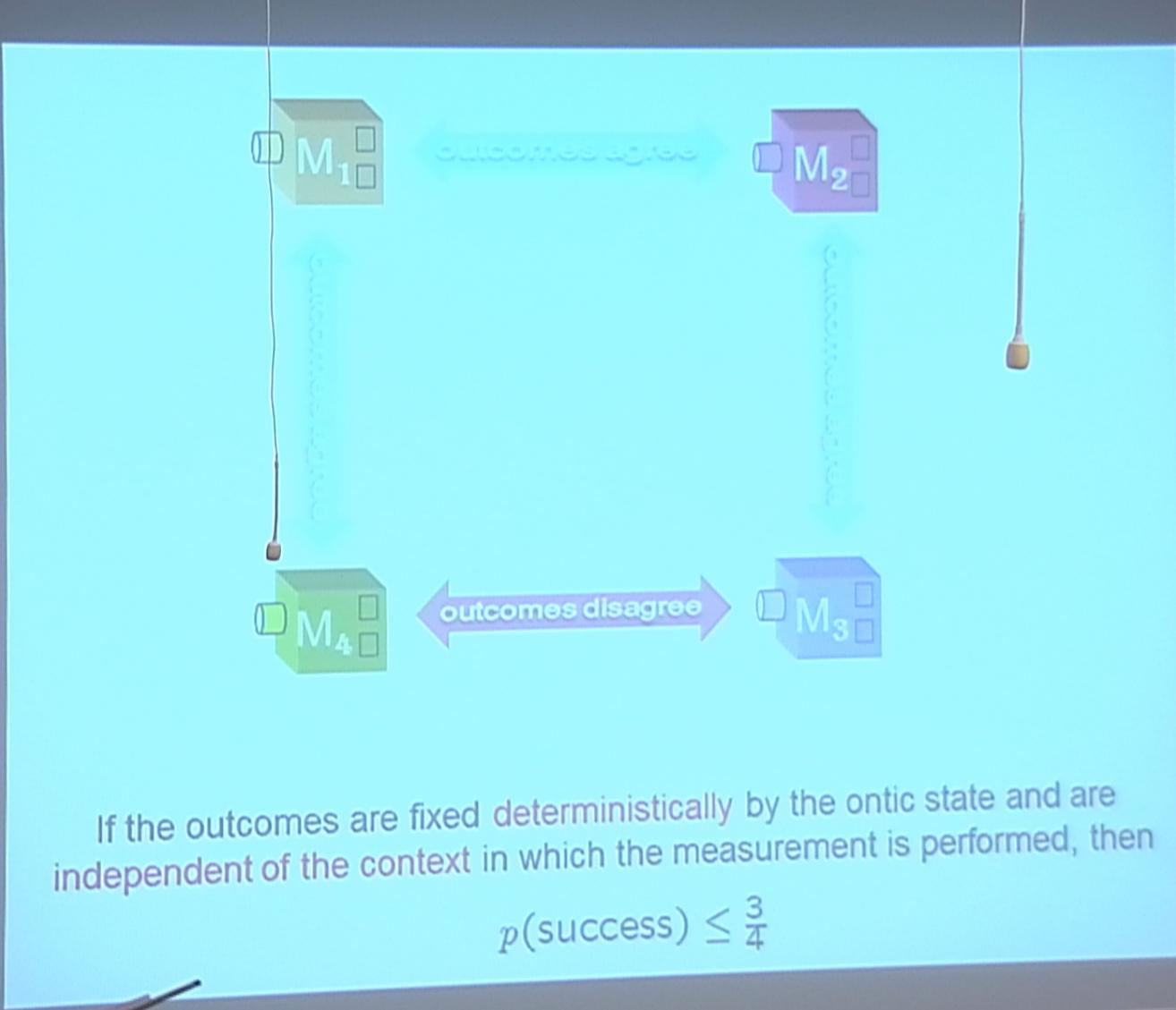
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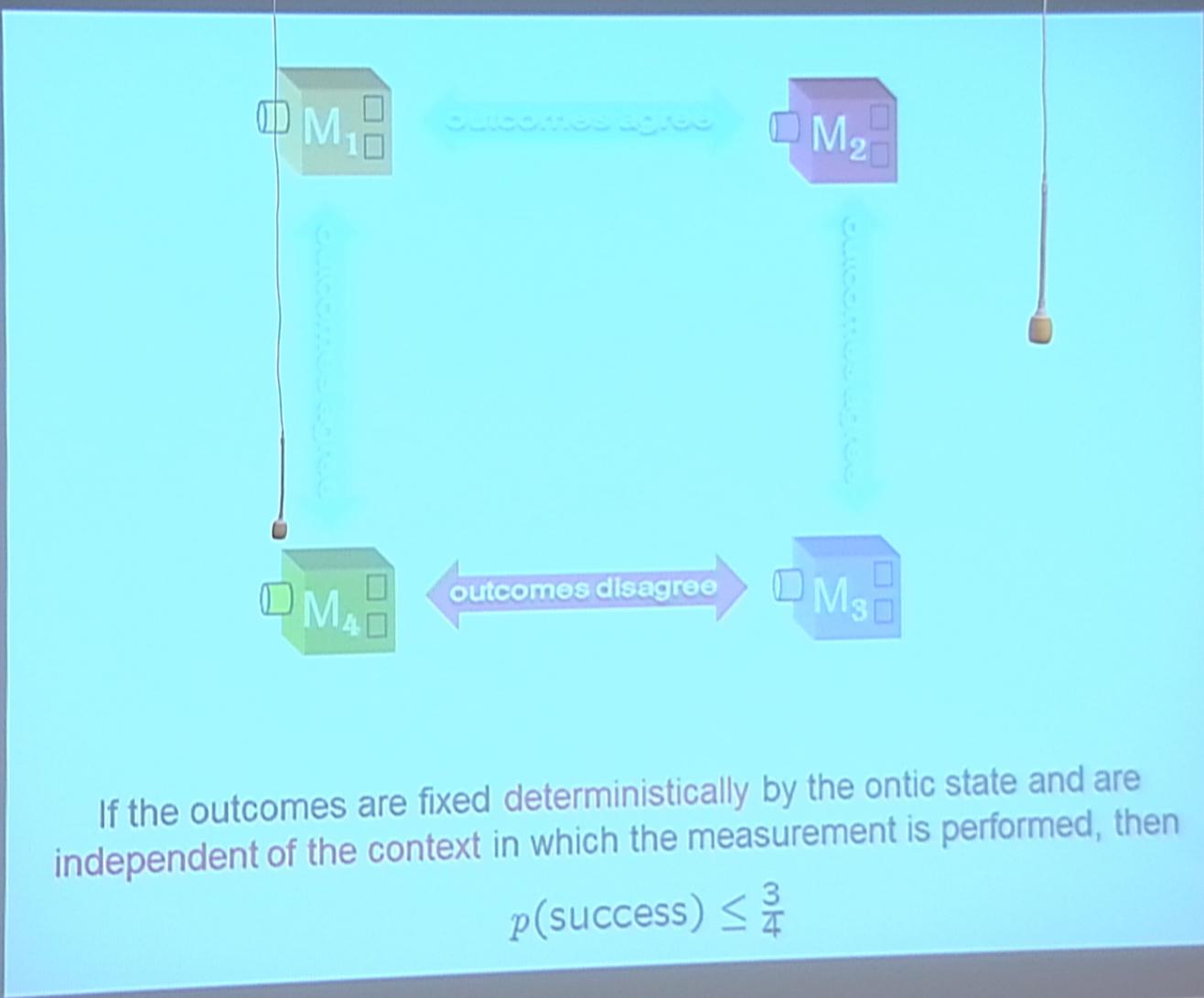
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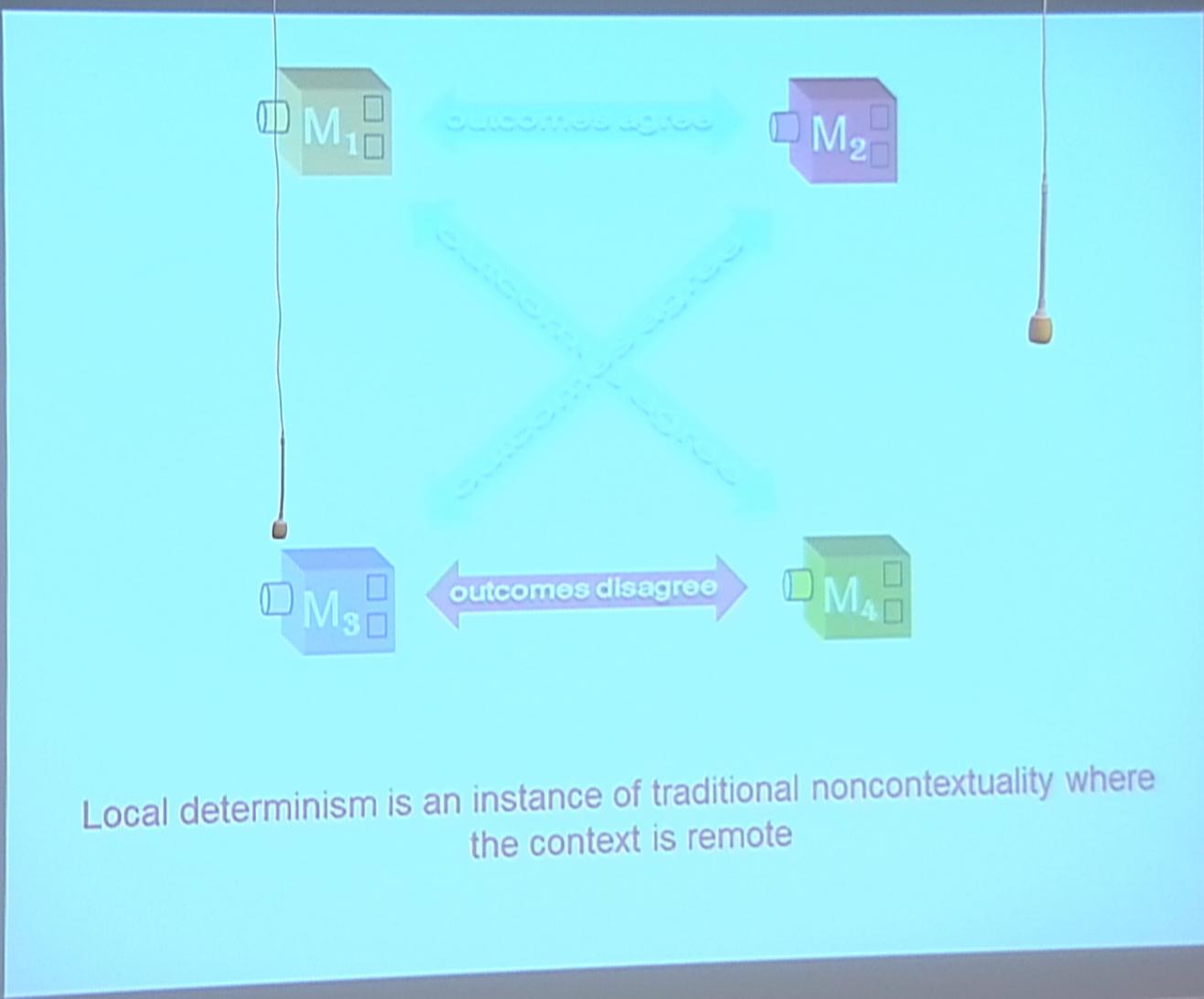
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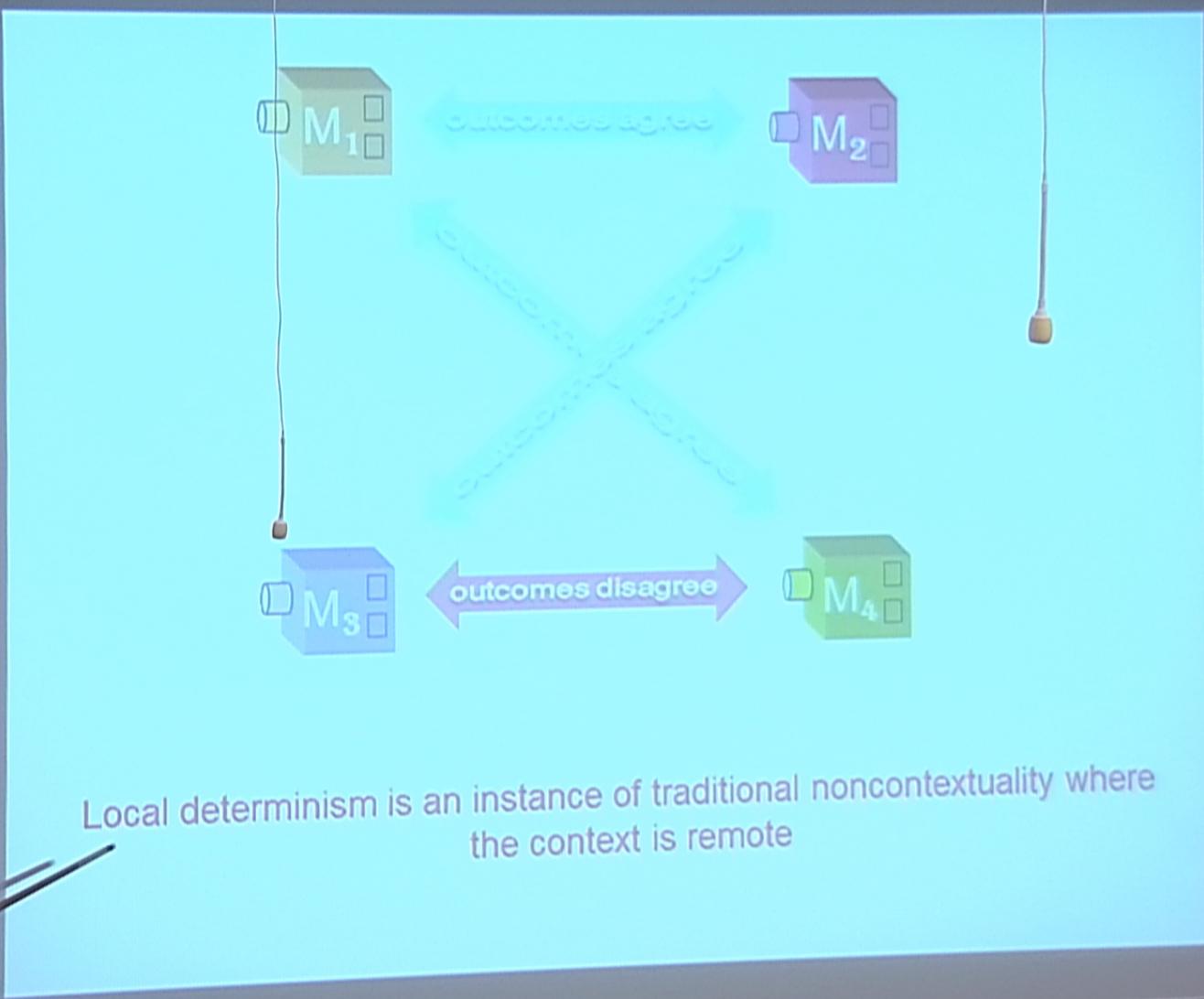
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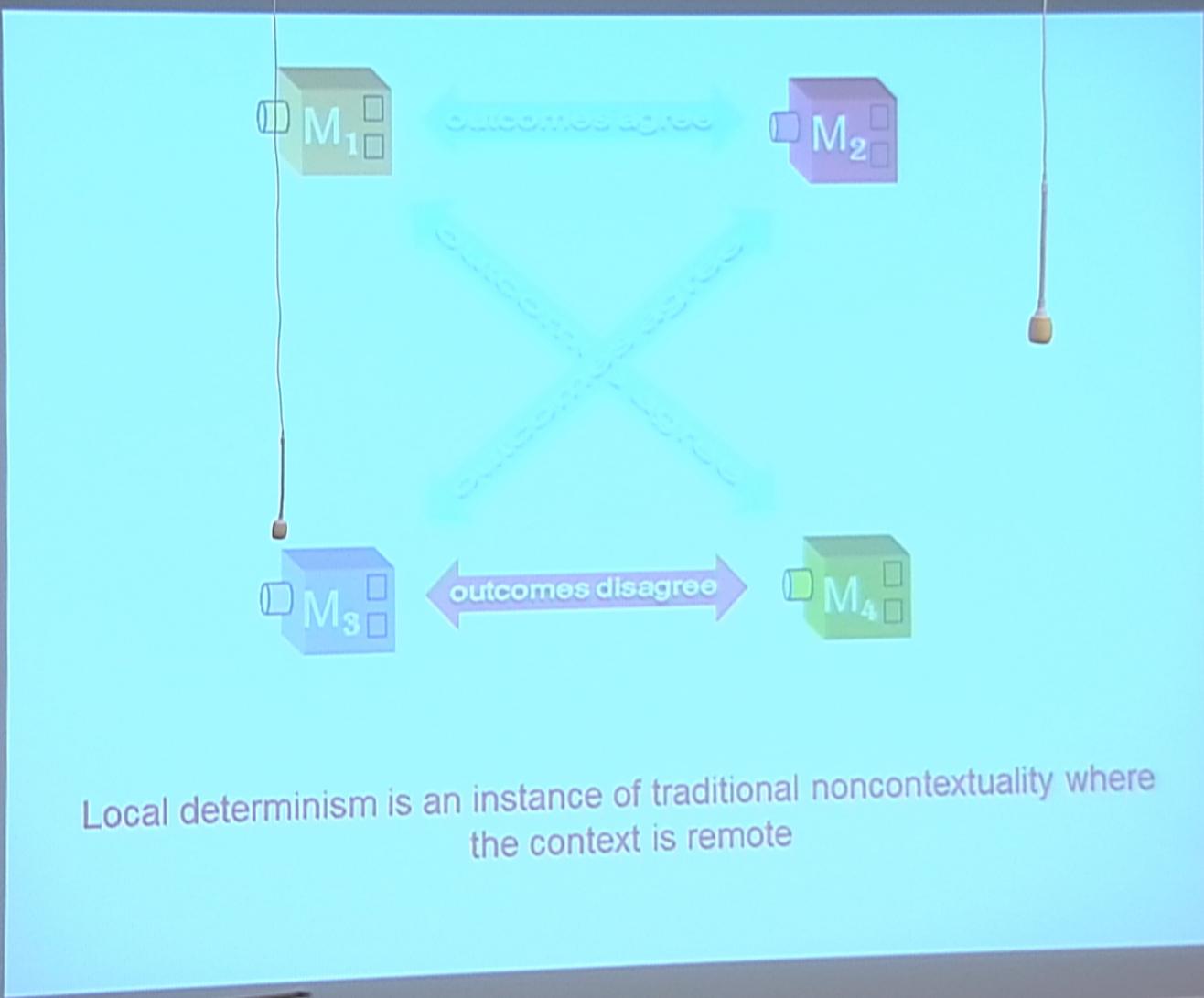








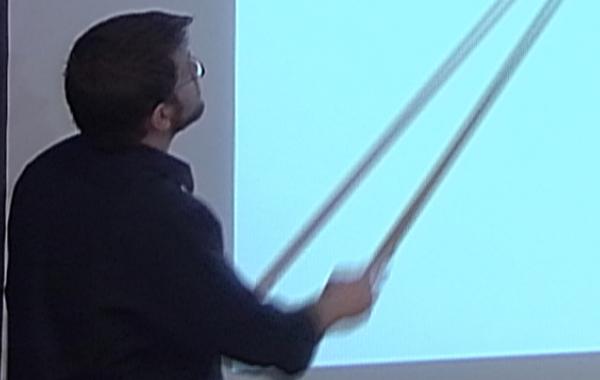




Klyachko's example



$$p(\text{success}) \leq \frac{4}{5}$$



Klyachko's example



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5 projective mmts:

$$\{|l_1\rangle\langle l_1|, I - |l_1\rangle\langle l_1|\}$$

$$\{|l_2\rangle\langle l_2|, I - |l_2\rangle\langle l_2|\}$$

$$\{|l_3\rangle\langle l_3|, I - |l_3\rangle\langle l_3|\}$$

$$\{|l_4\rangle\langle l_4|, I - |l_4\rangle\langle l_4|\}$$

$$\{|l_5\rangle\langle l_5|, I - |l_5\rangle\langle l_5|\}$$

where $\langle l_i | l_{i \oplus 1} \rangle = 0 \quad i \in \{1, \dots, 5\}$

Klyachko's example

$$\text{dashed hexagon} = \text{dashed star} \quad p(\text{success}) \leq \frac{4}{5}$$

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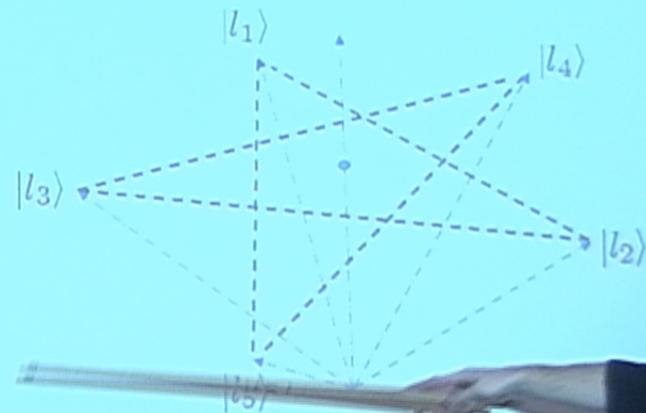
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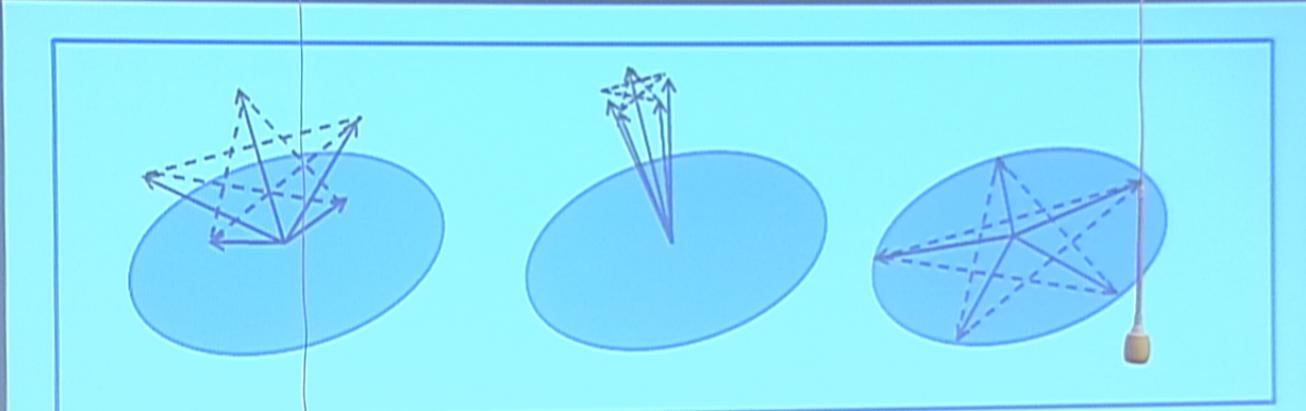
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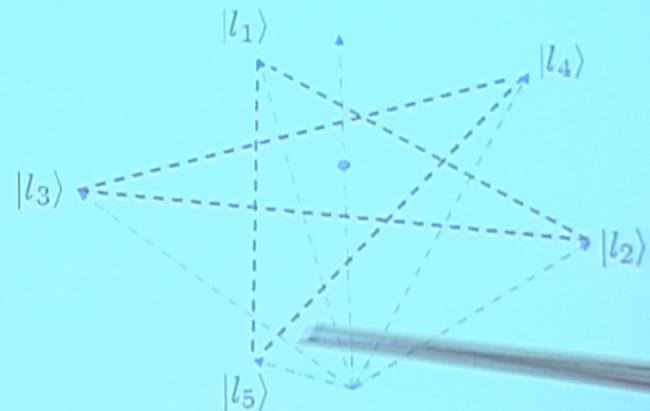
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Klyachko's example

$$\cos^2 \theta = \frac{1}{\sqrt{5}}$$

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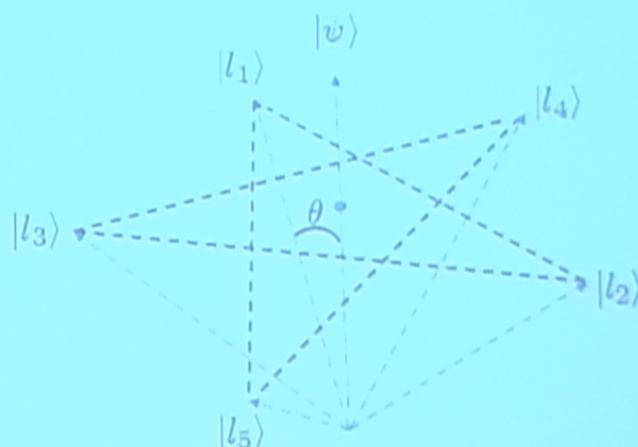
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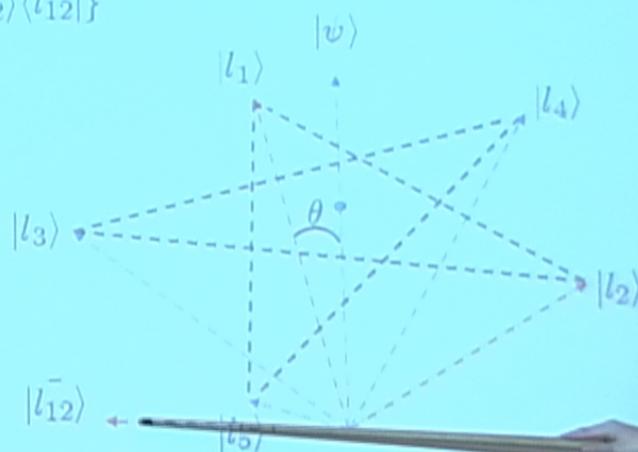
Preparation: the ψ that lies on the symmetry axis



Klyachko's example

$$\cos^2 \theta = \frac{1}{\sqrt{5}}$$

Consider: $\{|l_1\rangle\langle l_1|, |l_2\rangle\langle l_2|, |l_{12}\rangle\langle l_{12}|\}$



Klyachko's example

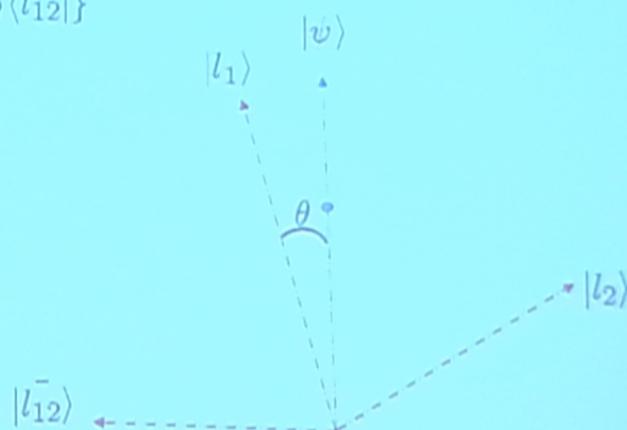
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$$\left. \begin{array}{l} v(|l_1\rangle\langle l_1|) = 1 \\ v(|l_2\rangle\langle l_2|) = 0 \\ v(|\bar{l}_{12}\rangle\langle \bar{l}_{12}|) = 0 \end{array} \right\} \text{prob. } |\langle \psi | l_1 \rangle|^2 = \frac{1}{\sqrt{5}}$$

$$\left. \begin{array}{l} v(|l_1\rangle\langle l_1|) = 0 \\ v(|l_2\rangle\langle l_2|) = 1 \\ v(|\bar{l}_{12}\rangle\langle \bar{l}_{12}|) = 0 \end{array} \right\} \text{prob. } |\langle \psi | l_2 \rangle|^2 = \frac{1}{\sqrt{5}}$$

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Klyachko's example

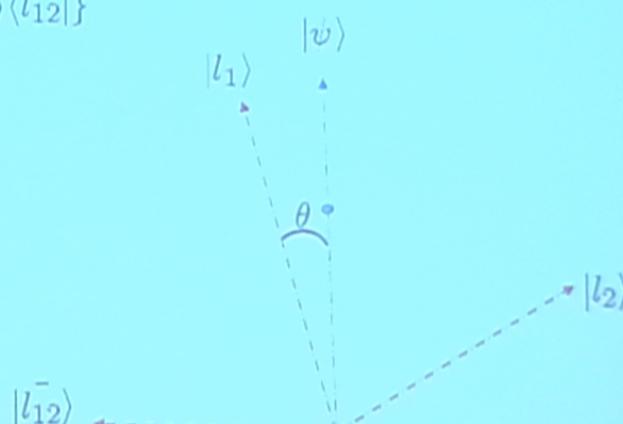
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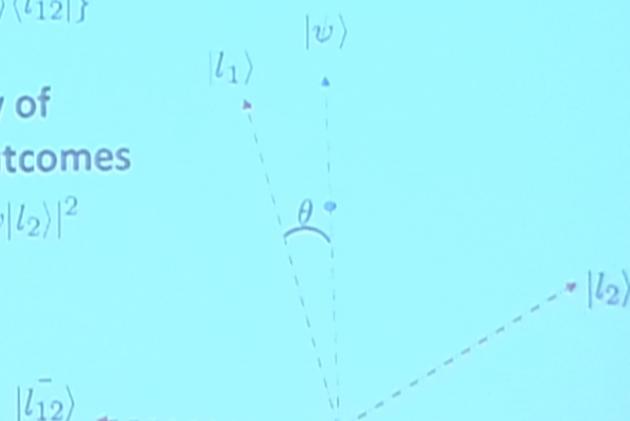
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Probability of
anticorrelated outcomes

$$|\langle \psi | l_1 \rangle|^2 + |\langle \psi | l_2 \rangle|^2$$

$$= \frac{2}{\sqrt{5}}$$



$$|\langle \psi | \bar{l}_{12} \rangle|^2$$

$$\text{prob.} |\langle \psi | \bar{l}_{12} \rangle|^2$$

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Klyachko's example

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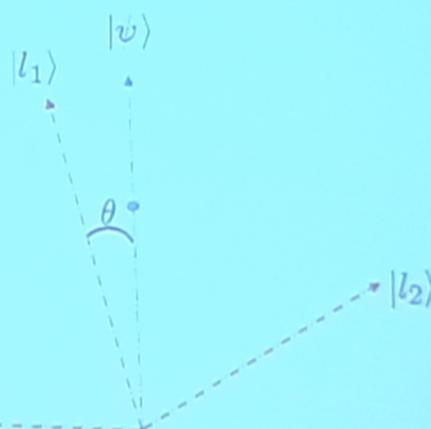
$$v(|l_2\rangle\langle l_2|) = 0$$

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Probability of
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$$|\langle \psi | l_1 \rangle|^2 + |\langle \psi | l_2 \rangle|^2 = \frac{2}{\sqrt{5}}$$

$$|\bar{l}_{12}\rangle$$



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$$\cos^2 \theta = \frac{1}{\sqrt{5}}$$

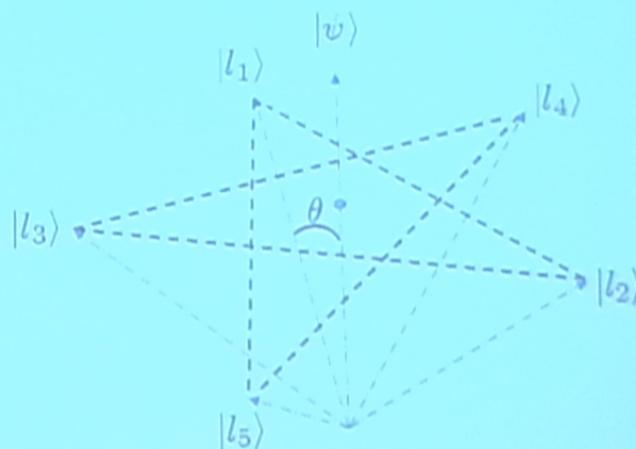
Similarly for any pair of measurements...

Probability of anticorrelated outcomes

$$= \frac{2}{\sqrt{5}}$$

Quantum probability of success

$$p(\text{success}) = \frac{2}{\sqrt{5}} \sim 0.89 > \frac{4}{5}$$



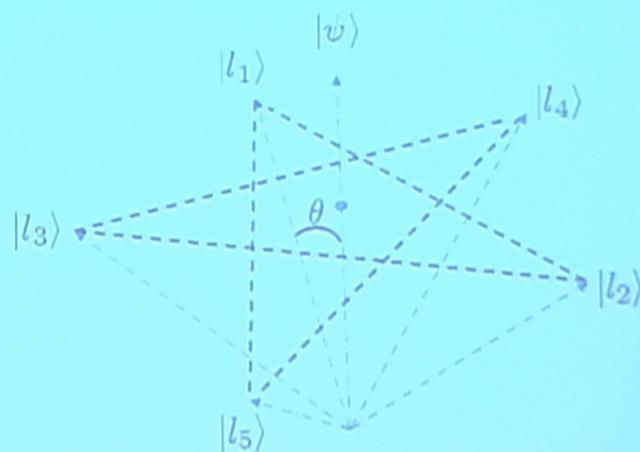
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