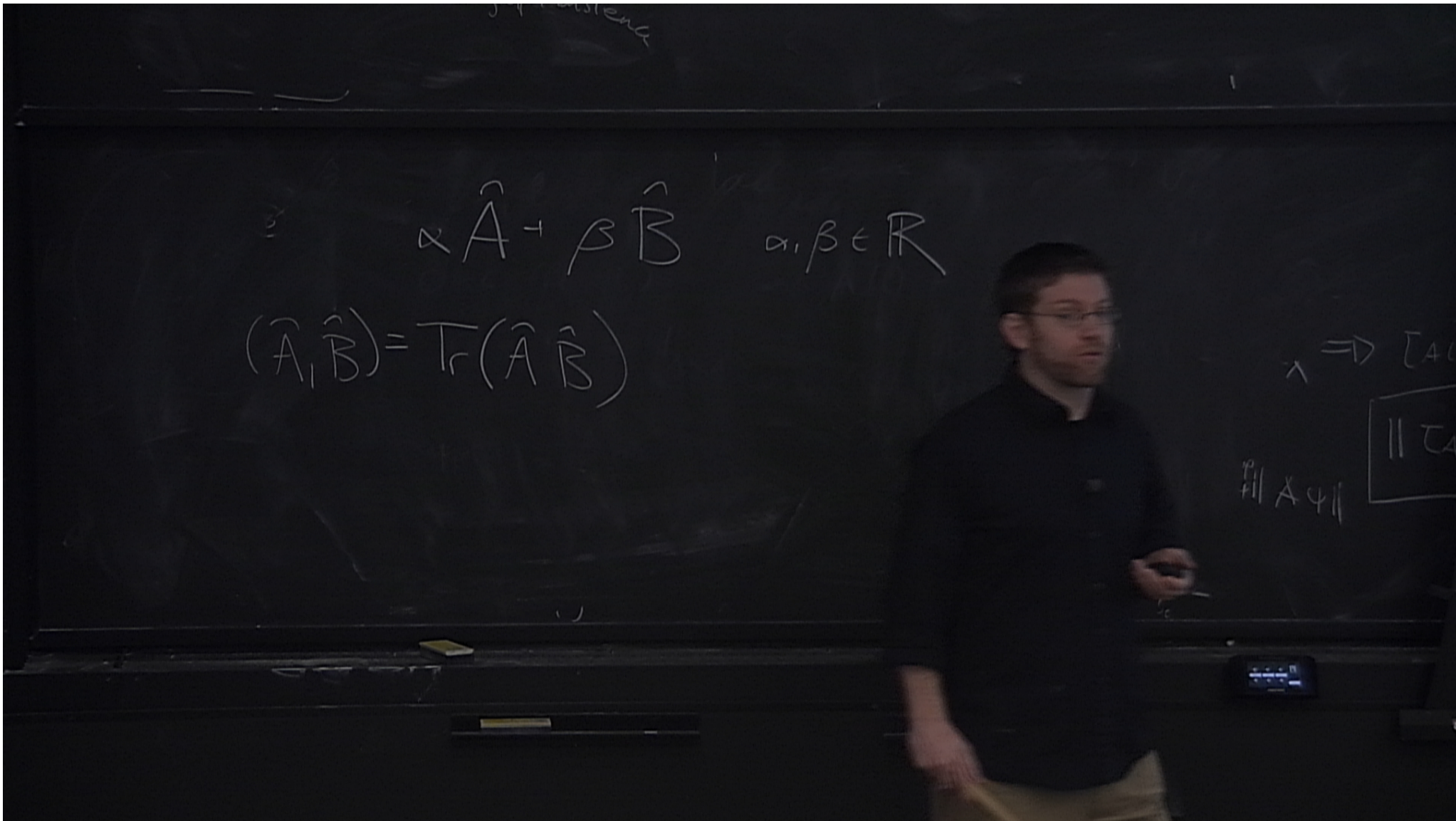


Title: Foundations of Quantum Mechanics - Lecture 5

Date: Jan 06, 2012 11:30 AM

URL: <http://pirsa.org/12010042>

Abstract:

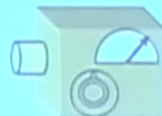


A framework for convex operational theories

See: L. Hardy, quant-ph/0101012



Preparation
P



Measurement
M

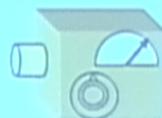
$$s_P = \begin{pmatrix} \Pr(1|M_1, P) \\ \Pr(2|M_1, P) \\ \Pr(1|M_2, P) \\ \Pr(2|M_2, P) \\ \Pr(3|M_2, P) \\ \vdots \end{pmatrix}$$

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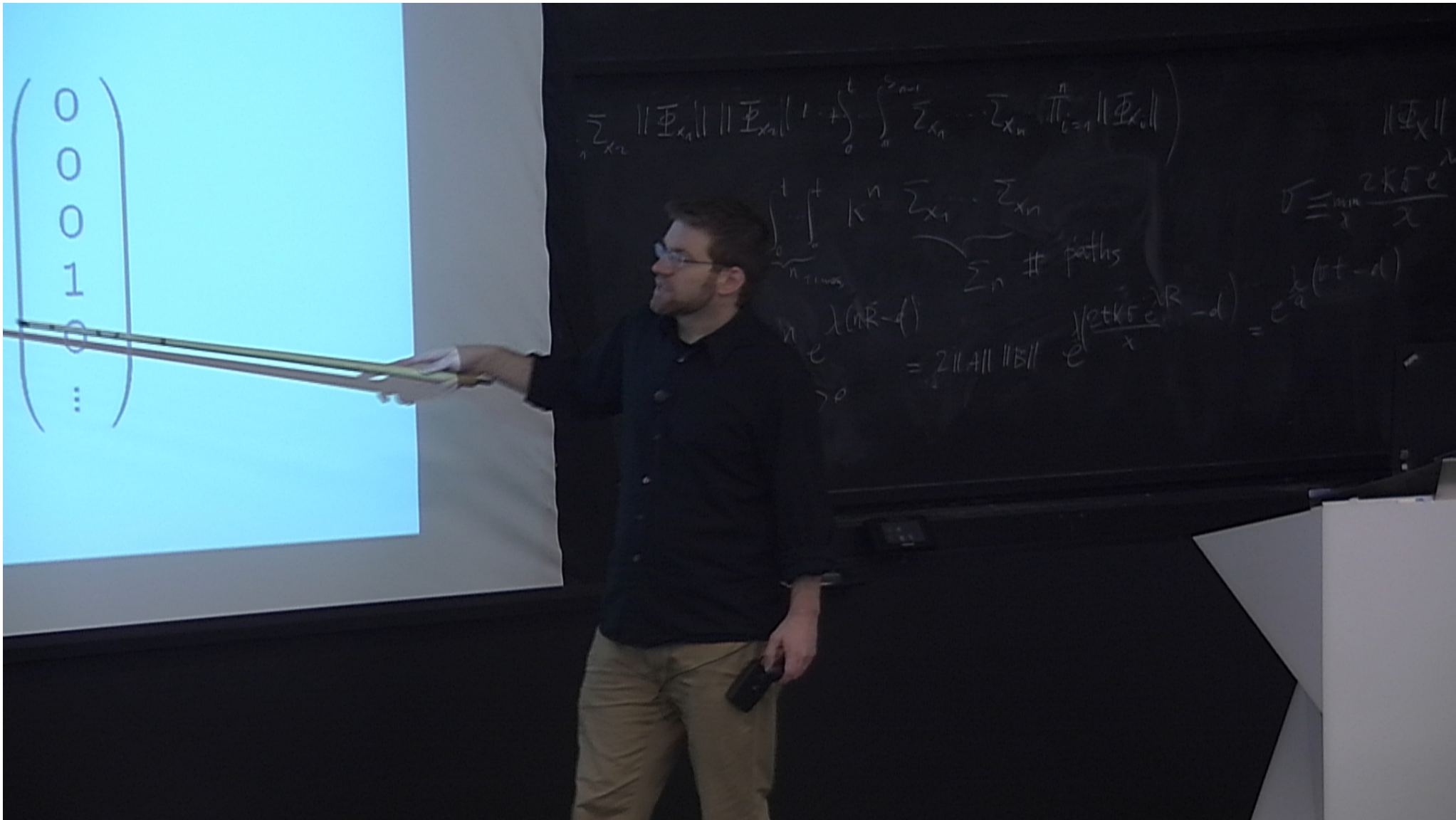
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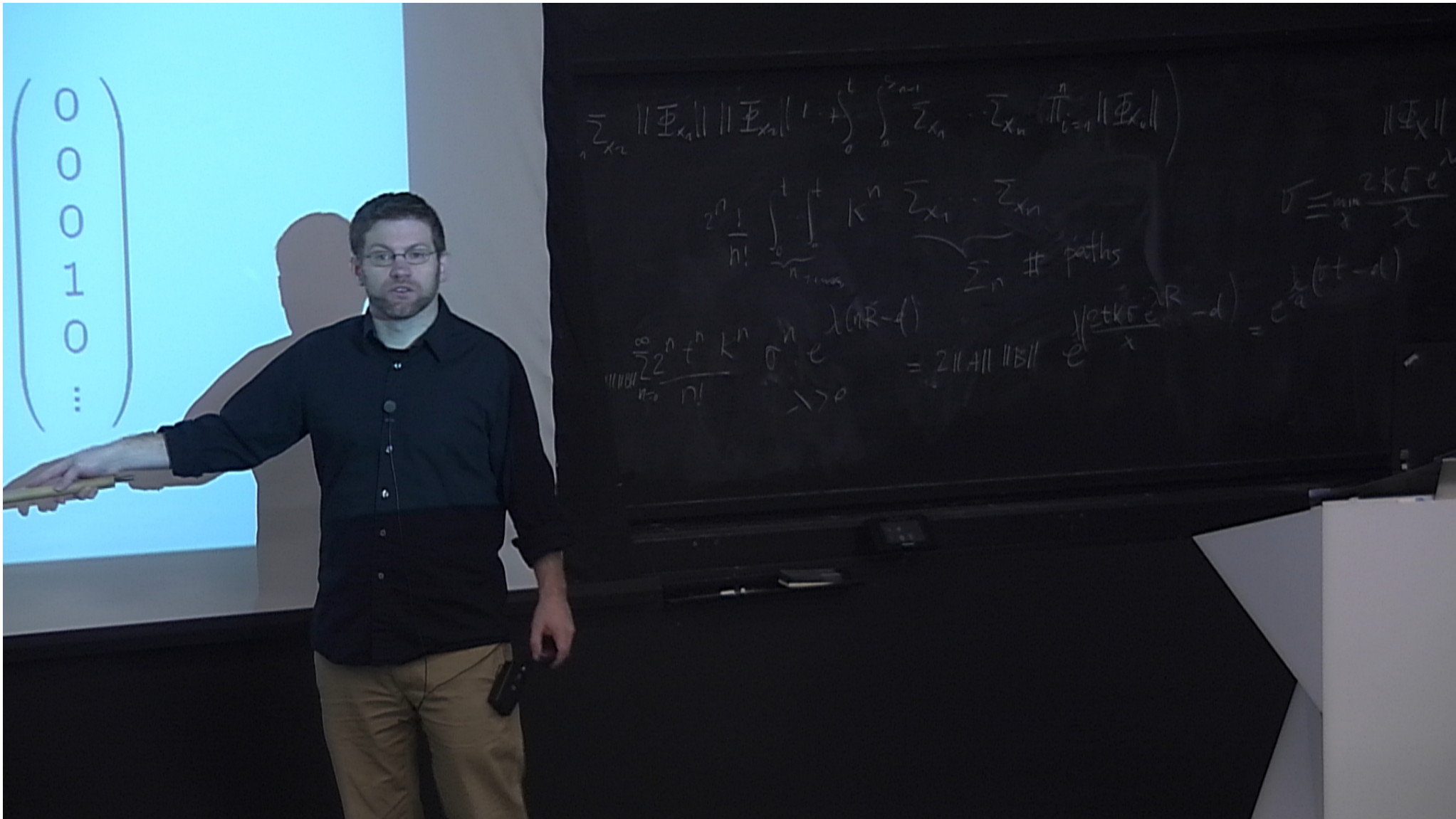


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$$\mathbf{r}_{M,k} = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 1 \\ 0 \\ \vdots \end{pmatrix}$$





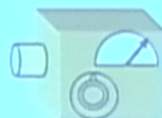
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See: L. Hardy, [quant-ph/0101012](https://arxiv.org/abs/quant-ph/0101012)



Preparation

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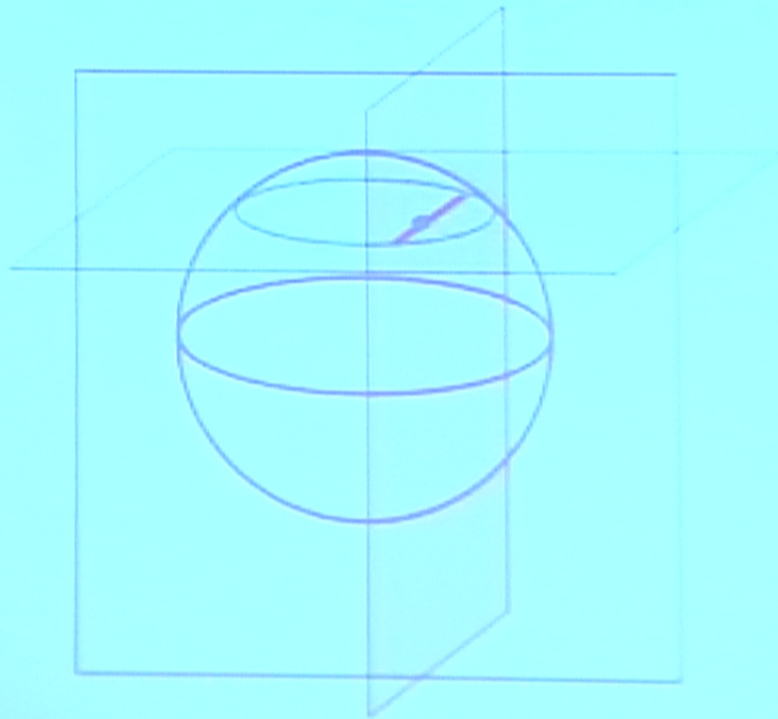


Measurement

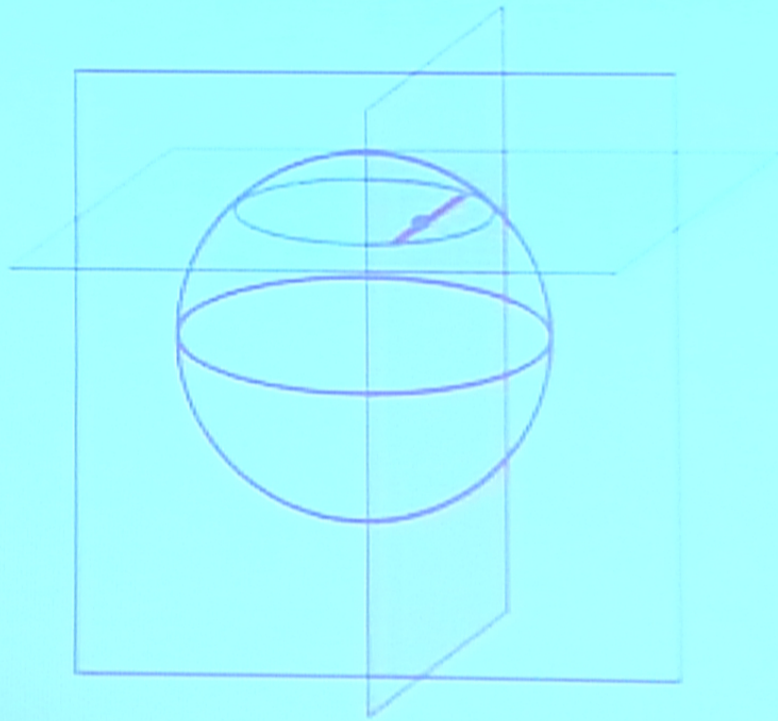
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Suppose there are K fiducial measurements (pass-fail mmts from which one can infer the statistics for all mmts)

State tomography for a single qubit



State tomography for a single qubit



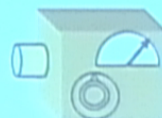
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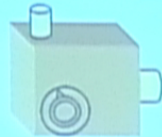
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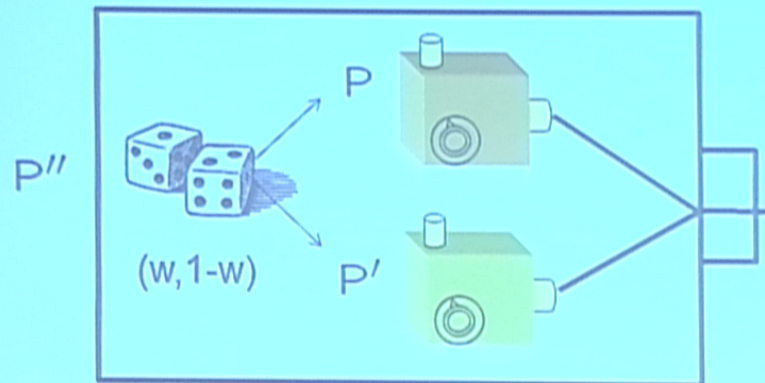
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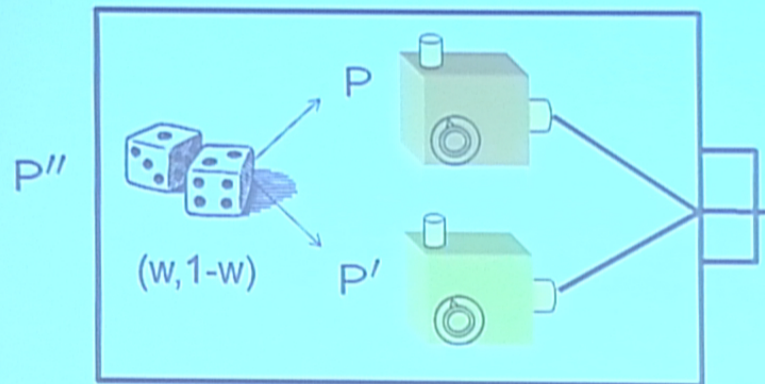
$$s_P = \begin{pmatrix} \Pr(\text{pass}|M_1, P) \\ \Pr(\text{pass}|M_2, P) \\ \vdots \\ \Pr(\text{pass}|M_K, P) \end{pmatrix} \quad \text{"operational state"}$$

$$\Pr(k|P, M) = f_{M,k}(s_P)$$

Operational states form a convex set

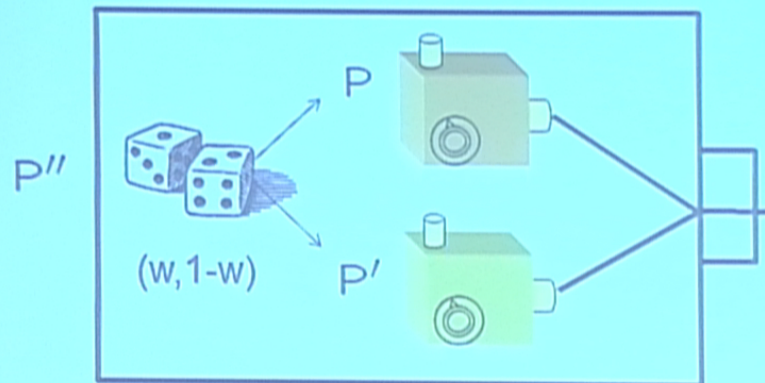


Operational states form a convex set



$$\forall M, k : p(k|M, P'') = w p(k|M, P) + (1-w) p(k|M, P')$$

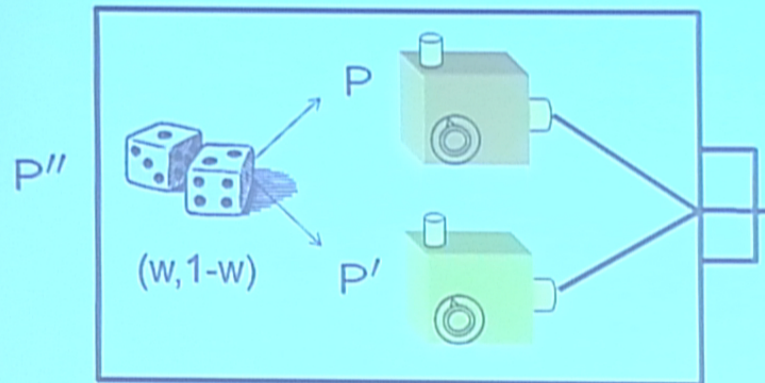
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$$f(s_{P''}) = w f(s_P) + (1-w) f(s_{P'})$$

Operational states form a convex set



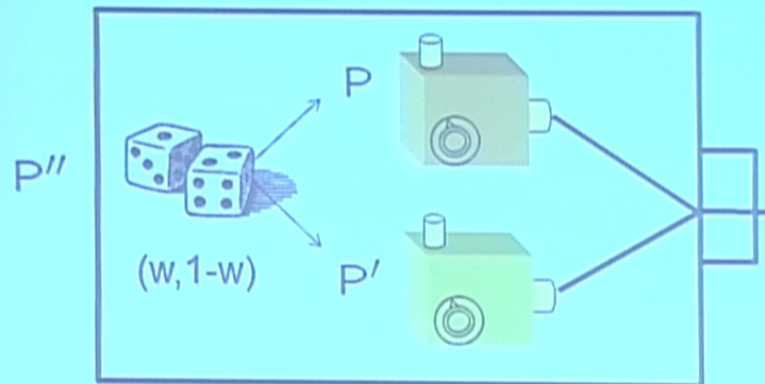
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Also true for fiducial mmts, so $s_{P''} = w s_P + (1-w) s_{P'}$



Operational states form a convex set



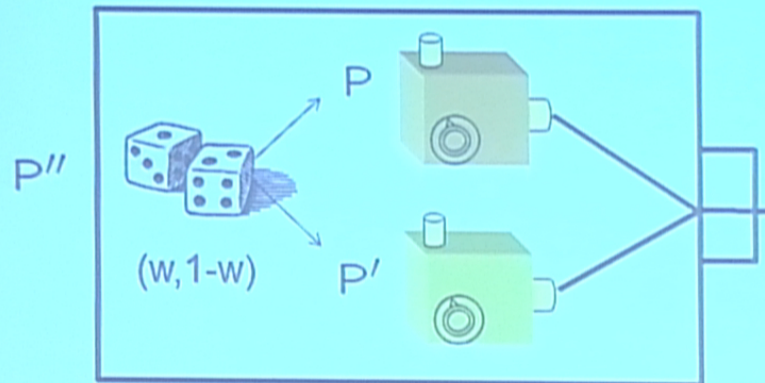
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Closed under convex combination \rightarrow a convex set

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$$f(w s_P + (1-w) s_{P'}) = w f(s_P) + (1-w) f(s_{P'}) \quad \text{Convex linear}$$

Convex linearity implies linearity

If f is convex linear on opt'l states

$$s = \sum_i w_i s_i \Rightarrow f(s) = \sum_i w_i f(s_i) \quad 0 \leq w_i \leq 1 \text{ and } \sum_i w_i = 1$$

Then f is linear on opt'l states

$$s = \sum_i \alpha_i s_i \Rightarrow f(s) = \sum_i \alpha_i f(s_i) \quad \alpha_i \in \mathbb{R}$$

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Proof: $s = \sum_i \alpha_i s_i$

$$s + \sum_{j \in I_-} |\alpha_j| s_j = \sum_{i \in I_+} |\alpha_i| s_i$$

Considering the trivial mmt, $1 = \sum_i \alpha_i$

$f(s)=1$ for all s , we have

$$1 + \sum_{j \in I_-} |\alpha_j| = \sum_{i \in I_+} |\alpha_i| \equiv \mathcal{N}$$

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$$1 + \sum_{j \in I_-} |\alpha_j| = \sum_{i \in I_+} |\alpha_i| \equiv N$$

$$\text{Thus: } \frac{1}{N} s + \sum_{j \in I_-} \frac{|\alpha_j|}{N} s_j = \sum_{i \in I_+} \frac{|\alpha_i|}{N} s_i$$

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$$f(s)=1 \text{ for all } s, \text{ we have } 1 + \sum_{j \in I_-} |\alpha_j| = \sum_{i \in I_+} |\alpha_i| \equiv \mathcal{N}$$

$$\text{Thus: } \frac{1}{\mathcal{N}} s + \sum_{j \in I_-} \frac{|\alpha_j|}{\mathcal{N}} s_j = \sum_{i \in I_+} \frac{|\alpha_i|}{\mathcal{N}} s_i$$

$$\frac{1}{\mathcal{N}} f(s) + \sum_{j \in I_-} \frac{|\alpha_j|}{\mathcal{N}} f(s_j) = \sum_{i \in I_+} \frac{|\alpha_i|}{\mathcal{N}} f(s_i)$$

$$f(s) = \sum_i \alpha_i f(s_i)$$

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Therefore $\exists r : f(s) = r \cdot s$

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A convex operational theory



Preparation
P

$$s_P \in S$$

"operational states"

S = Convex set



Measurement
M

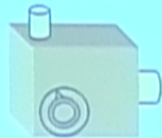
$$r_{M,k} \in R$$

"operational effects"

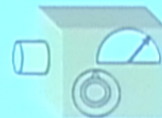
R = Interval of
positive cone

$$Pr(k|P, M) = r_{M,k} \cdot s_P$$

A convex operational theory



Preparation
P



Measurement
M

$$s_P \in S$$

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S = Convex set

$$r_{M,k} \in R$$

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S and R characterize the operational theory!

$$Pr(k|P, M) = r_{M,k} \cdot s_P$$

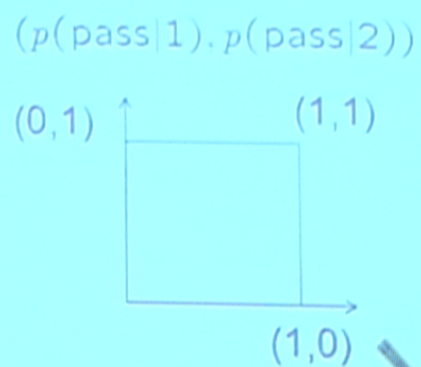
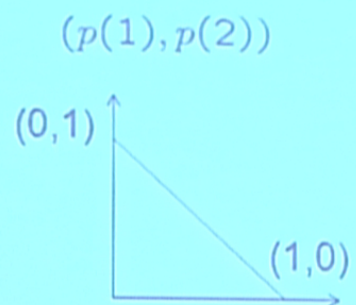
Operational classical theory

\mathbf{s} can be any probability distribution

S = a simplex

\mathbf{r} can be any vector of conditional probabilities

R = the unit hypercube



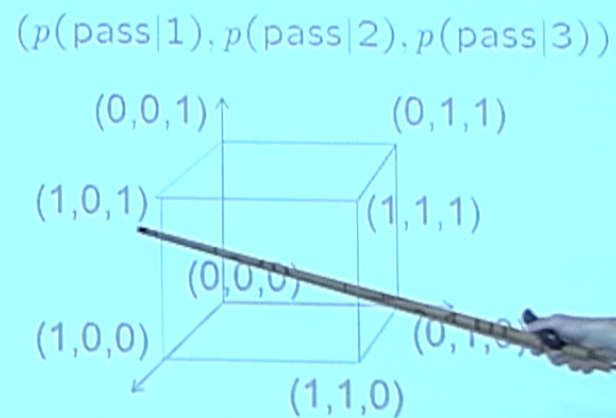
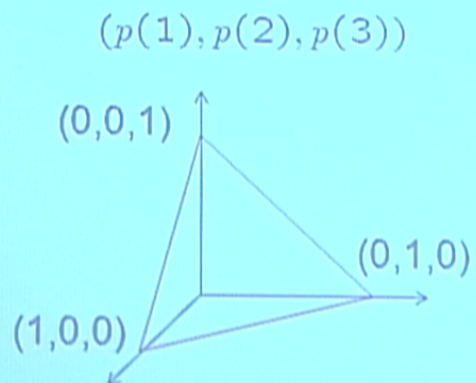
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Operational quantum theory

Recall: The Hermitian operators on a Hilbert space of dimension d form a real Euclidean vector space of dimension d^2

s can be any trace one positive operator

S = the convex set of such operators

existence

$$\alpha \hat{A} + \beta \hat{B} \quad \alpha, \beta \in \mathbb{R}$$

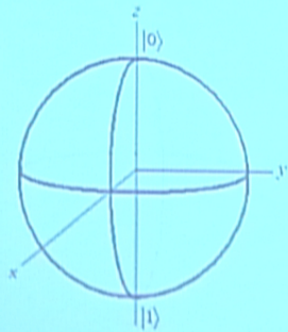
$$(\hat{A}, \hat{B}) = \text{Tr}(\hat{A} \hat{B})$$

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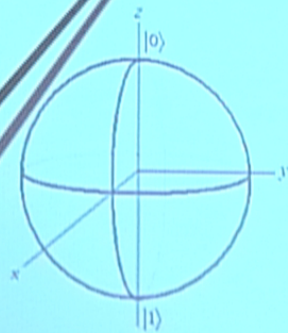
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\mathcal{R} can be any positive operator less than identity

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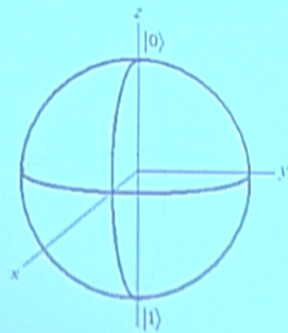
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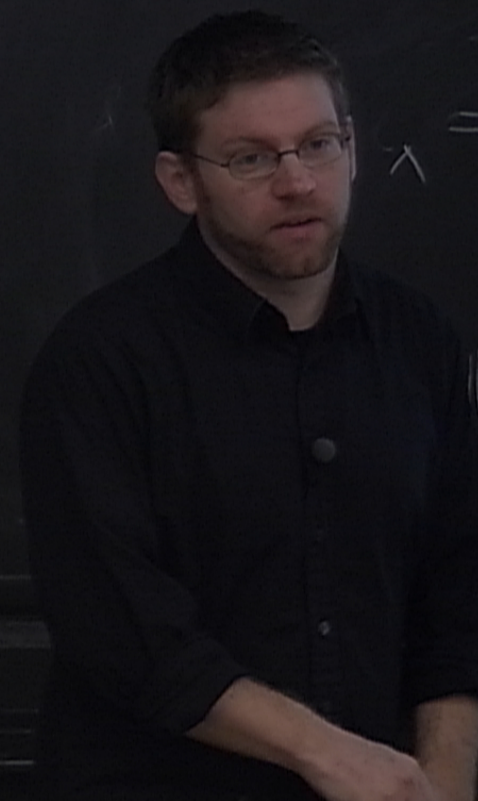
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A little bit of axiomatics

Suppose one takes as given that

S = the convex set of positive trace-one operators

Suppose one assumes that every logically possible measurement is physically possible

Allow all $\{\mathbf{r}_k\}$ such that $\mathbf{r}_k \cdot \mathbf{s} \geq 0 \quad \forall \mathbf{s} \in S$

$$\sum_k \mathbf{r}_k \cdot \mathbf{s} = 1 \quad \forall \mathbf{s} \in S$$

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The real vector space is the space of Hermitian operators
(these are closed under linear combination and scalar multiplication)

The inner product is $(A, B) = \text{Tr}(AB)$

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Allow all $\{E_k\}$ such that $\text{Tr}(E_k \rho) \geq 0 \quad \forall \rho \in S(\mathcal{H})$

$$\sum_k \text{Tr}(E_k \rho) = 1 \quad \forall \rho \in S(\mathcal{H})$$

$$\text{Tr}(\rho E_k) \geq 0 \quad \forall \rho \in \mathcal{S}(\mathcal{H})$$

$$\rightarrow \langle \psi | E_k | \psi \rangle \geq 0 \quad \forall |\psi\rangle \in \mathcal{H}$$

$\rightarrow E_k$ is a positive operator

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$\rightarrow E_k$ is a positive operator

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$$\rightarrow \langle \psi | (\sum_k E_k) | \psi \rangle = 1 \quad \forall |\psi\rangle \in \mathcal{H}$$

$$\rightarrow \sum_k E_k = I$$

The logically possible measurements correspond to the POVMs!

Operational formulation of quantum theory

Every preparation P is associated with a density operator ρ

Every logically possible measurement is physically possible

↳ Every measurement M is associated with a positive operator-valued measure $\{E_k\}$. The probability of M yielding outcome k given a preparation P is $Pr(k|P, M) = \text{Tr}(\rho E_k)$

Every transformation is associated with a trace-preserving completely-positive linear map $\rho \rightarrow \rho' = \mathcal{T}(\rho)$

Every measurement outcome k is associated with a trace-nonincreasing completely-positive linear map \mathcal{T}_k such that

$$\rho \rightarrow \rho_k = \frac{\mathcal{T}_k(\rho)}{\text{Tr}[\mathcal{T}_k(\rho)]} \quad \text{where} \quad \mathcal{T}_k^\dagger(I) = E_k$$

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The inner product is $(A, B) = \text{Tr}(AB)$

Each \mathbf{s} is a density operator ρ

Each set $\{\mathbf{r}_k\}$ is a set of Hermitian operators $\{E_k\}$

$\mathbf{r}_k \cdot \mathbf{s} = (E_k, \rho) = \text{Tr}(E_k \rho) \leftarrow \text{the form of the Born rule}$

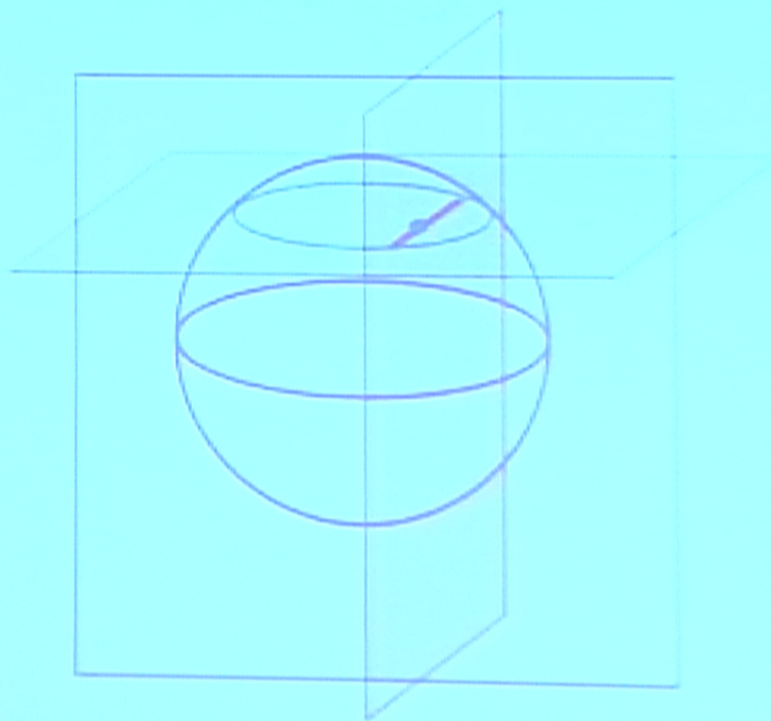
Allow all $\{E_k\}$ such that $\text{Tr}(E_k \rho) \geq 0 \quad \forall \rho \in S(\mathcal{H})$

$$\sum_k \text{Tr}(E_k \rho) = 1 \quad \forall \rho \in S(\mathcal{H})$$

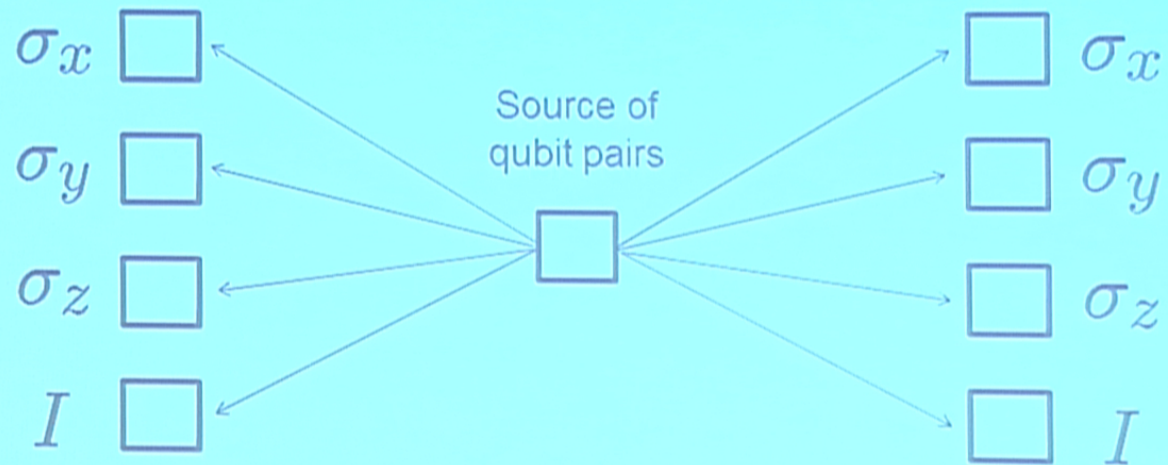
Real versus complex field

	real case	complex case
Pure preparations	rays in \mathbb{R}^N	rays in \mathbb{C}^N
Complete repeatable measurements	Bases for \mathbb{R}^N	Bases for \mathbb{C}^N
Reversible transformations	Orthogonal ($\det = 1$)	Unitary
Mixed preparations	Positive unit-trace real matrix	Positive unit-trace complex matrix
Composition rule	Tensor product	Tensor product

State tomography for a single qubit



State tomography for two qubits

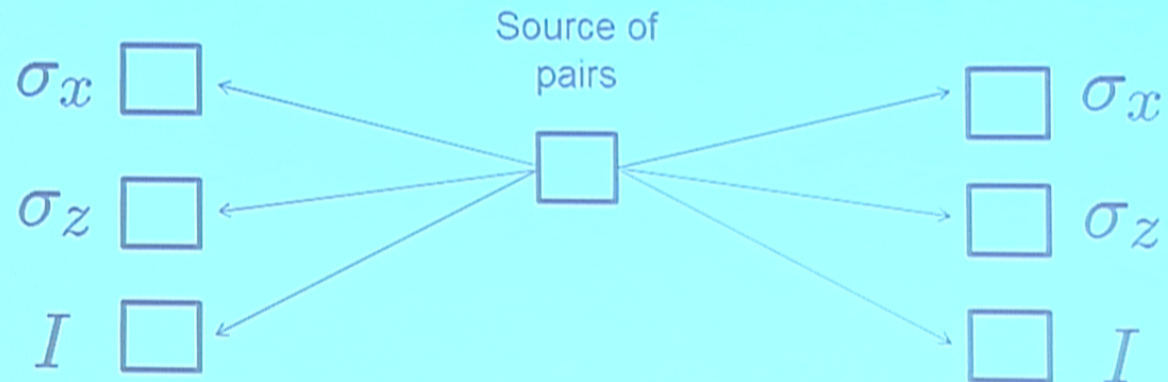


We need $4^2 - 1 = 15$ parameters

We obtain $4^2 - 1 = 15$ parameters

The mixed state of two qubits can be determined from local measurements

State tomography for two real-amplitude qubits



We need $4(4+1)/2 - 1 = 9$ parameters

We obtain $3^2 - 1 = 8$ parameters

$\sigma_y \otimes \sigma_y$ must be accessed globally

The mixed state of two real-amplitude qubits
cannot be determined from local measurements -- a kind of holism