

Title: Search for New Non-standard Decay Modes of SM-like Higgs at the LHC

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Abstract: We present strategies of searching for supersymmetric non-standard decays of Standard Model (SM)-like Higgs bosons (h_2) at the Large Hadron Collider (LHC), motivated by "Dark Light Higgs" (DLH) scenario. The DLH scenario represents a limit of the nearly-Peccei-Quinn-symmetric Next-to-Minimal Supersymmetric Standard Model, where there naturally co-exist two light singlet-like particles: a scalar (h_1), a pseudoscalar (a_1), and a light singlino-like DM candidate (χ_1^-), all with masses of order 10 GeV or below. In this scenario, the SM-like Higgs boson typically decays dominantly into a pair of neutralinos, allowing themselves to be as light as below 100 GeV. We systematically study the searches of both the SM-like and the light Higgs bosons at the LHC, using the supersymmetric non-standard decay chains: $h_2 \rightarrow \chi_1^- \chi_2^-$, $\chi_2^- \rightarrow \chi_1^-$ (h_1 , a_1) with the h_1 or a_1 further decayed into a pair of fermions (including diMuon, diTau and b bbar, etc).



Search for New Non-standard Decays of the SM-like Higgs at the LHC



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Dec. 6th 2011



Outline

- Motivation
- Dark Light Higgs (DLH) Model
- Model Independent DLH Search
 - Di-Muon channel
 - B-bbar channel
 - Di-Tau channel
- Conclusion



Dark Light Higgs Scenario

Draper, Liu, Wagner, Wang, Zhang, PRL. 106 121805 (2011)

$$\mathbf{W} = \lambda \mathbf{N} \mathbf{H_u} \mathbf{H_d} + \frac{1}{3} \kappa \mathbf{N}^3,$$
$$V_{soft} = m_{H_d}^2 |H_d|^2 + m_{H_u}^2 |H_u|^2 + m_N^2 |N|^2$$
$$- (\lambda A_\lambda H_u H_d N + h.c.) + \left(\frac{\kappa}{3} A_\kappa N^3 + h.c. \right)$$

- κN^3 explicitly breaks Peccei-Quinn symmetry
- Dark light higgs scenario:
nearly PQ limit of NMSSM
($\kappa/\lambda \rightarrow 0$, $A_\kappa \rightarrow 0$, moderate or small λ)
- Three CP-even higgs (h_1, h_2, h_3); two CP-odd higgs (a_1, a_2)



Masses of the Higgses

h₂ is SM-like: $h_2 \sim h_u + h_d \cot \beta - \frac{2\varepsilon v m_Z}{m_Z^2 + \mu^2} h_n$

$$\varepsilon = \frac{\lambda\mu}{m_Z} \left(\frac{A_\lambda}{\mu \tan \beta} - 1 \right)$$

h₁ is the lightest CP-even scalar, Singlet Like:

$$m_{h_1}^2 \approx -4\varepsilon^2 v^2 + \frac{4\lambda^2 v^2}{\tan^2 \beta} + \frac{\kappa A_\kappa \mu}{\lambda} + \frac{4\kappa^2 \mu^2}{\lambda^2} \quad \Delta m_{h_1}^2 \approx \frac{\lambda^2 \mu^2}{2\pi^2} \log \frac{\mu^2 \tan^3 \beta}{m_Z^2} \quad (\text{Loop correction})$$

A light CP-odd Higgs a_1 , **Singlet Like**:

$$m_{a_1}^2 \approx -\frac{3\kappa A_\kappa \mu}{\lambda}$$

A lightest neutralino χ_1 , **Singlino Like**:

$$m_{\chi_1} \approx \frac{\lambda^2 v^2}{\mu} \sin 2\beta + \frac{2\kappa \mu}{\lambda}$$

dark matter particle

Comparison: in the R-symmetry limit, h_1 and χ_1 are typically not so light and h_1 is SM-like

B. A. Dobrescu et al., PRD 63, 075003 (2001);
 R. Dermisek et al., PRL 95, 041801 (2005).



Dark Light Higgs Scenario

Draper, Liu, Wagner, Wang, Zhang, PRL. 106 121805 (2011)

$$\begin{aligned}\mathbf{W} &= \lambda \mathbf{N} \mathbf{H_u} \mathbf{H_d} + \frac{1}{3} \kappa \mathbf{N}^3, \\ V_{soft} &= m_{H_d}^2 |H_d|^2 + m_{H_u}^2 |H_u|^2 + m_N^2 |N|^2 \\ &\quad - (\lambda A_\lambda H_u H_d N + h.c.) + \left(\frac{\kappa}{3} A_\kappa N^3 + h.c. \right)\end{aligned}$$

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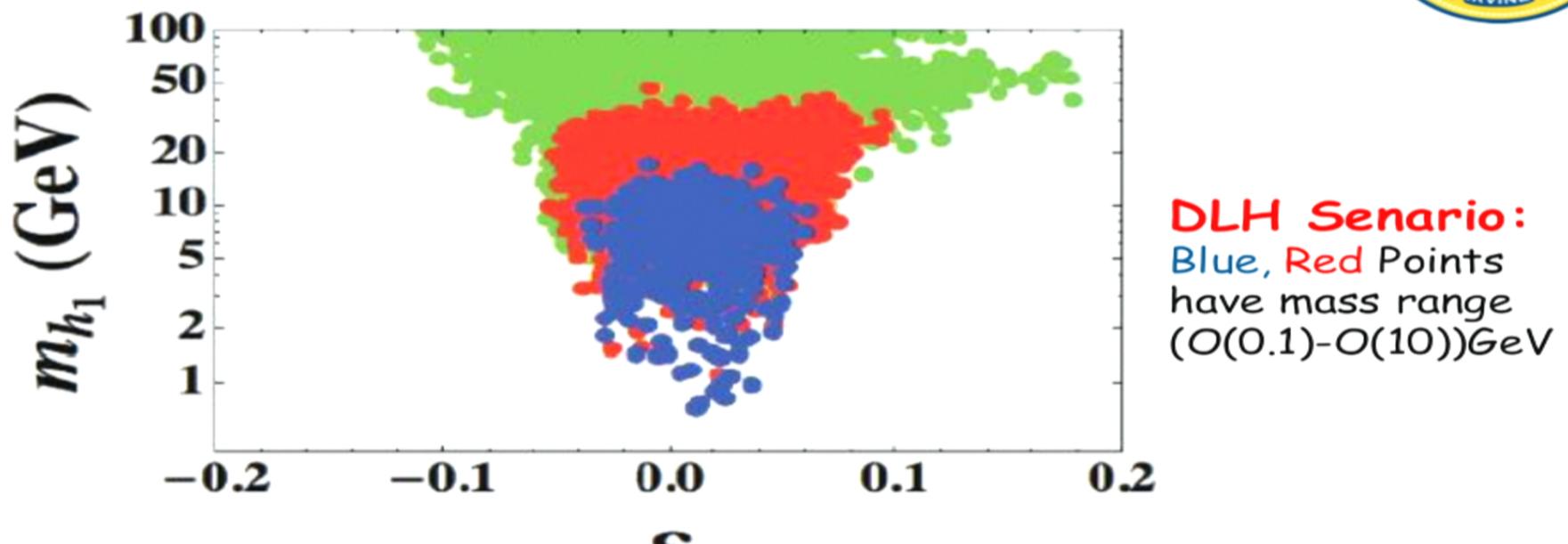
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Parameter Scan



Vacuum stability sets a small upper bound on ε

No points near $\varepsilon \rightarrow 0$ because of the vacuum stability requirement

$$5 \leq \tan \beta \leq 50, \quad 0.05 \leq \lambda \leq 0.5, \quad 0.0005 \leq \kappa \leq 0.05, \quad -0.8 \leq \varepsilon' \leq 0.8, \quad -40 \text{ GeV} \leq A_\kappa \leq 0, \quad 0.1 \text{ TeV} \leq \mu \leq 1 \text{ TeV}$$

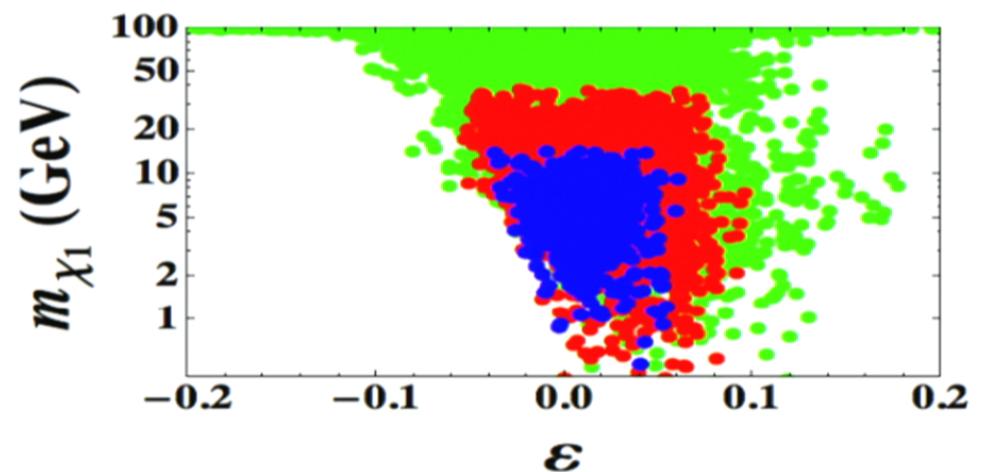
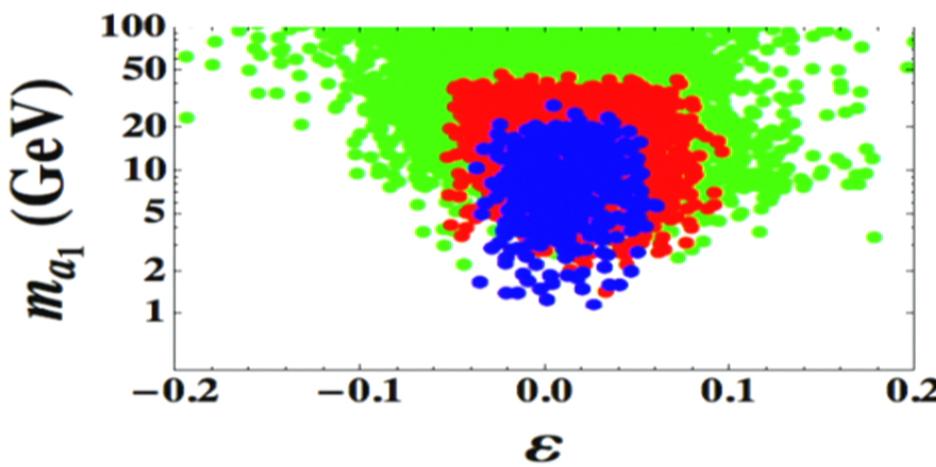
► $\lambda < 0.30, \quad \kappa/\lambda < 0.05, \quad \mu < 400 \text{ GeV}$

$\lambda < 0.15, \quad \kappa/\lambda < 0.03, \quad \mu < 250 \text{ GeV}$



Parameter Scan (cont.)

Light pseudoscalar and neutralino masses





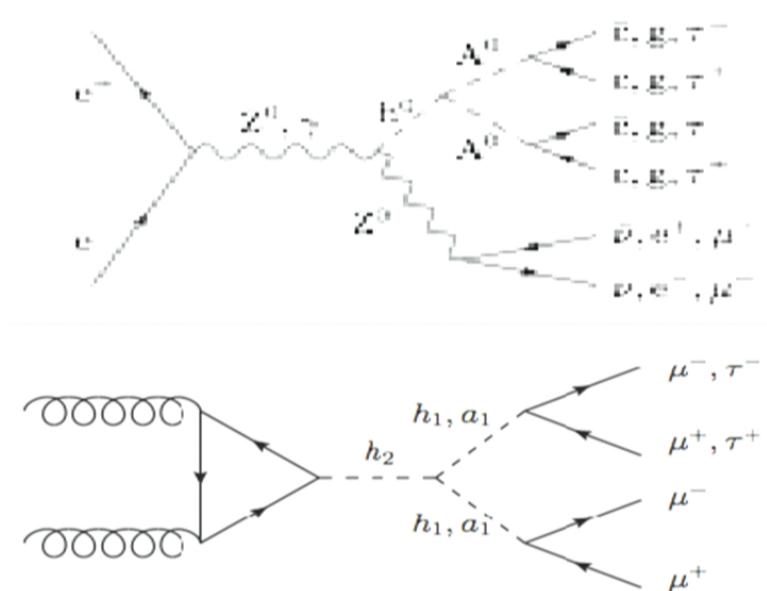
Constraints from $h_2 \rightarrow h_1 h_1, a_1 a_1$

► LEP searches:

- (1) $(h_2 \rightarrow a_1 a_1)$, $a_1 \rightarrow 2b$ (S. Schael et al. [ALEPH, DELPHI, L3, and OPAL Collaborations], Eur. Phys. J. C 47(2006); S. Schael et al. [ALEPH Collaboration], JHEP 1005 (2010));
- (2) Z-associated Higgs production, with Z leptonically decayed (S. Schael et al. [ALEPH Collaboration], JHEP 1005 (2010); G. Abbiendi et al. [The OPAL Collaboration], Eur. Phys. J. C 27, (2003)).

► Tevatron searches:

$h_2 \rightarrow a_1 a_1, h_1 h_1 \rightarrow 4\mu, 2\mu 2\tau$ (V. M. Abazov et al. [D0 Collaboration], Phys. Rev. Lett. 103 (2009))



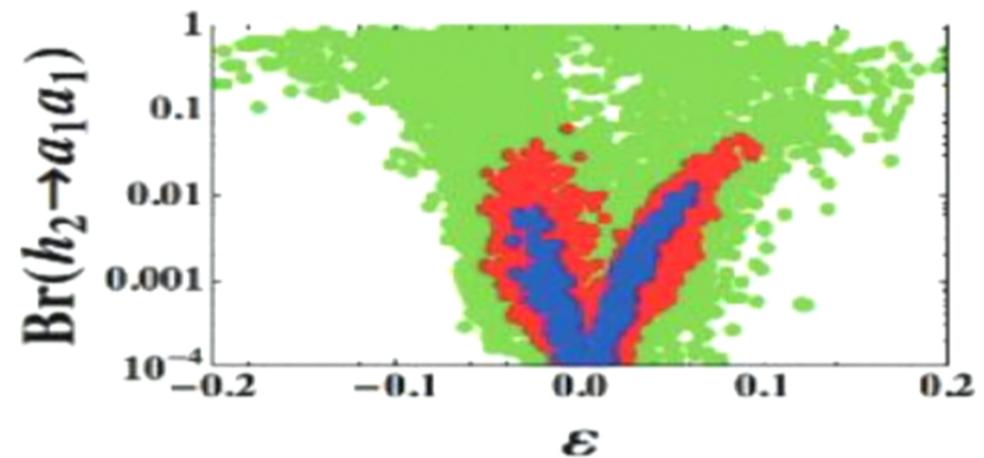
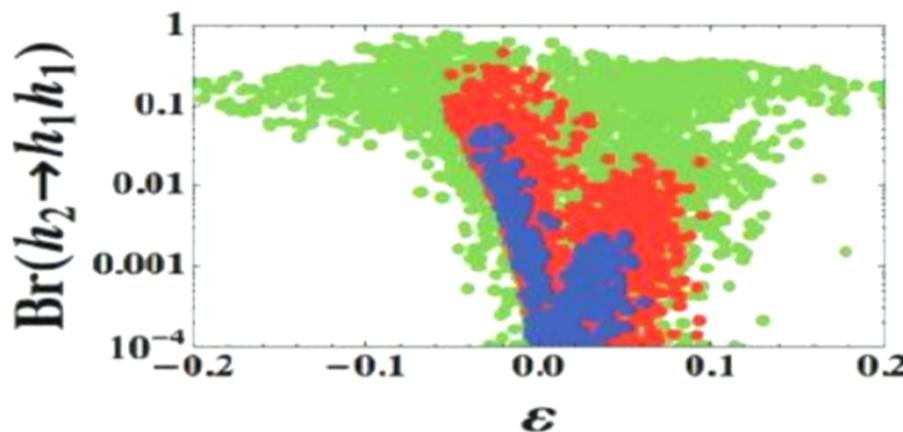


$h_2 \rightarrow h_1 h_1$, $a_1 a_1$ modes

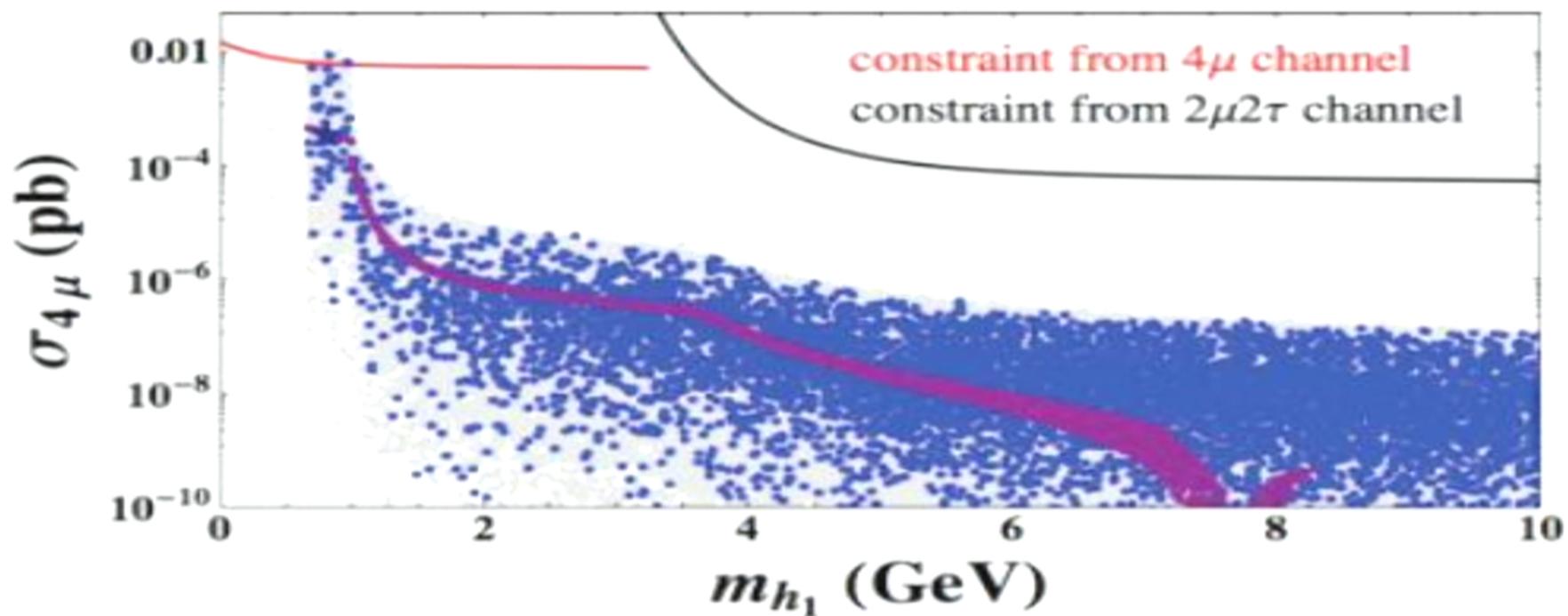
DLH scenario has $h_2 \rightarrow h_1 h_1$ and $h_2 \rightarrow a_1 a_1$ decay channels as well, but highly suppressed

Exp. Constraints can be easily satisfied.

$$|y_{h_2 h_1 h_1}| = |y_{h_2 a_1 a_1}| = \frac{\lambda v m_Z \varepsilon}{\sqrt{2} \mu}$$

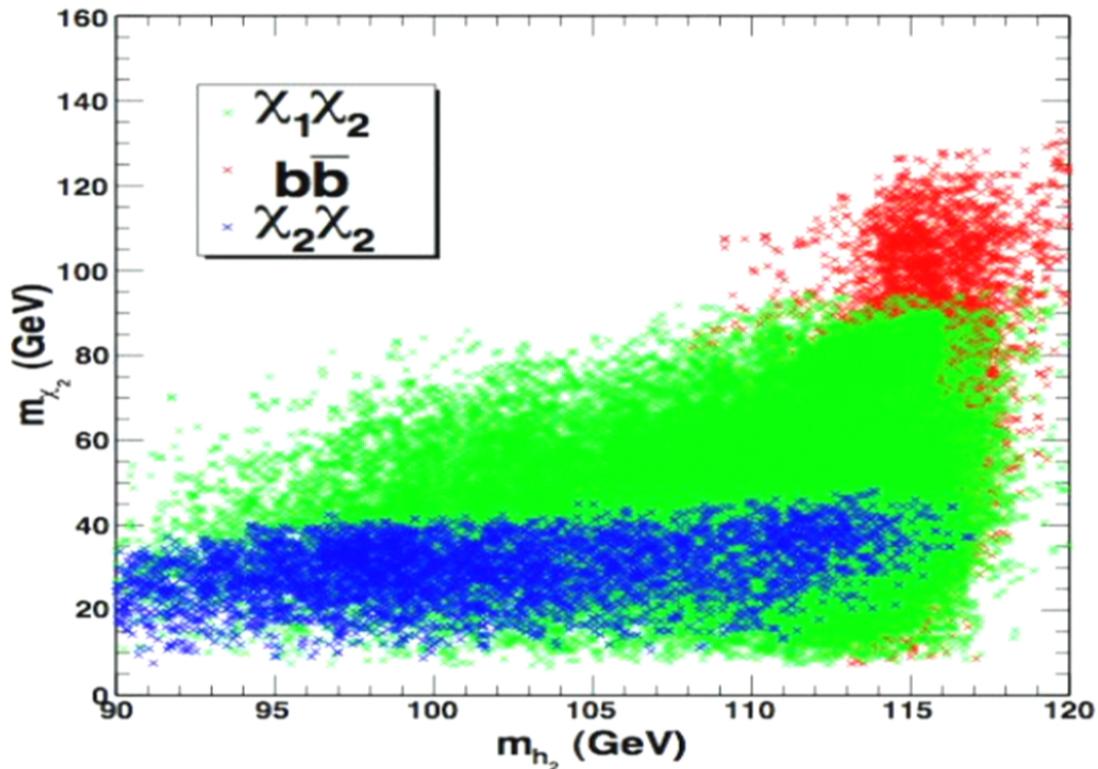


Constraints from $h_2 \rightarrow h_1 h_1, \alpha_1 \alpha_1$





h_2 decay modes

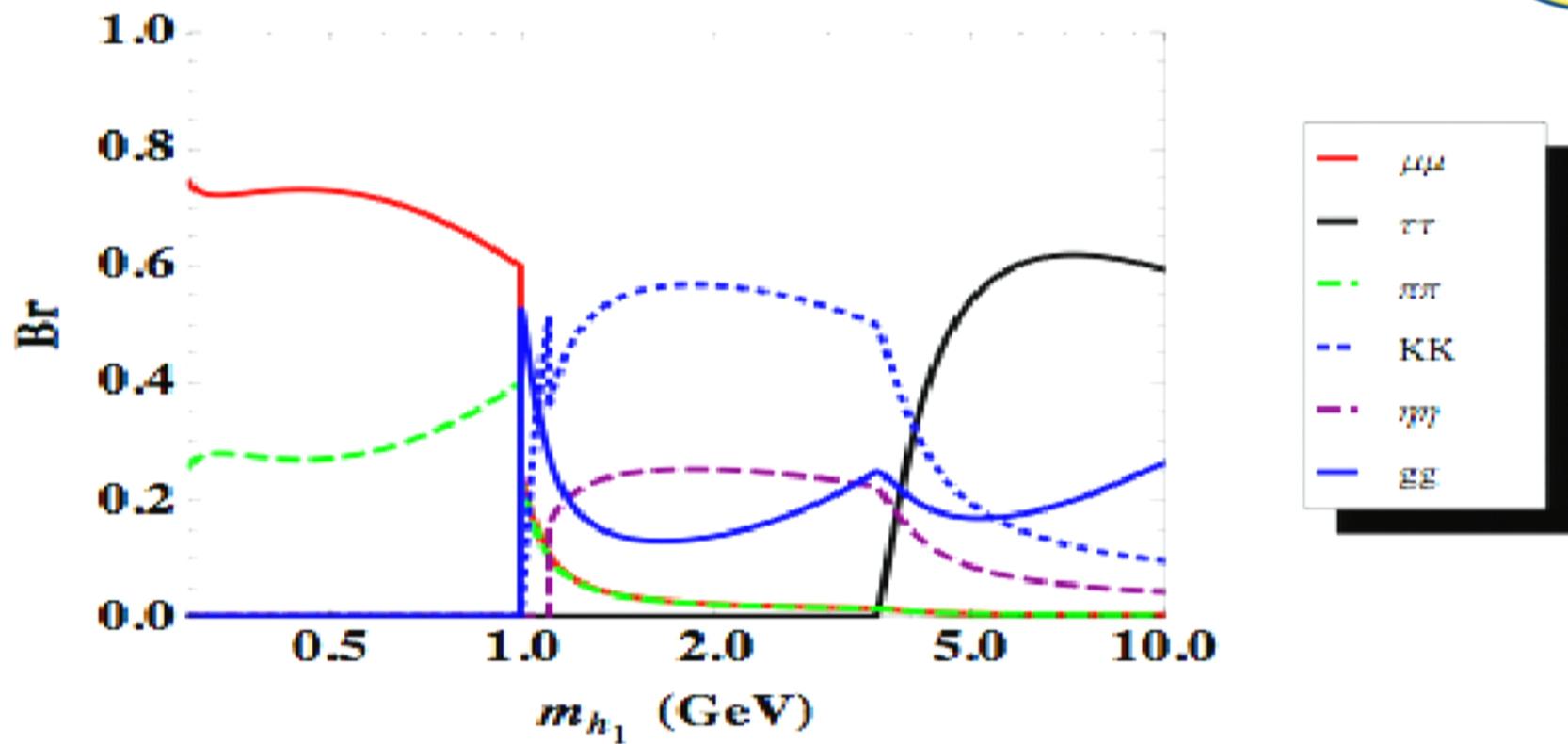


$h_2 \rightarrow \chi_1 \chi_2$ is typically dominant as long as it is kinematically allowed, and it is corresponding to the **GREEN** points.

$h_2 \rightarrow bb$ mode can be dominant sometimes, but NOT generic.

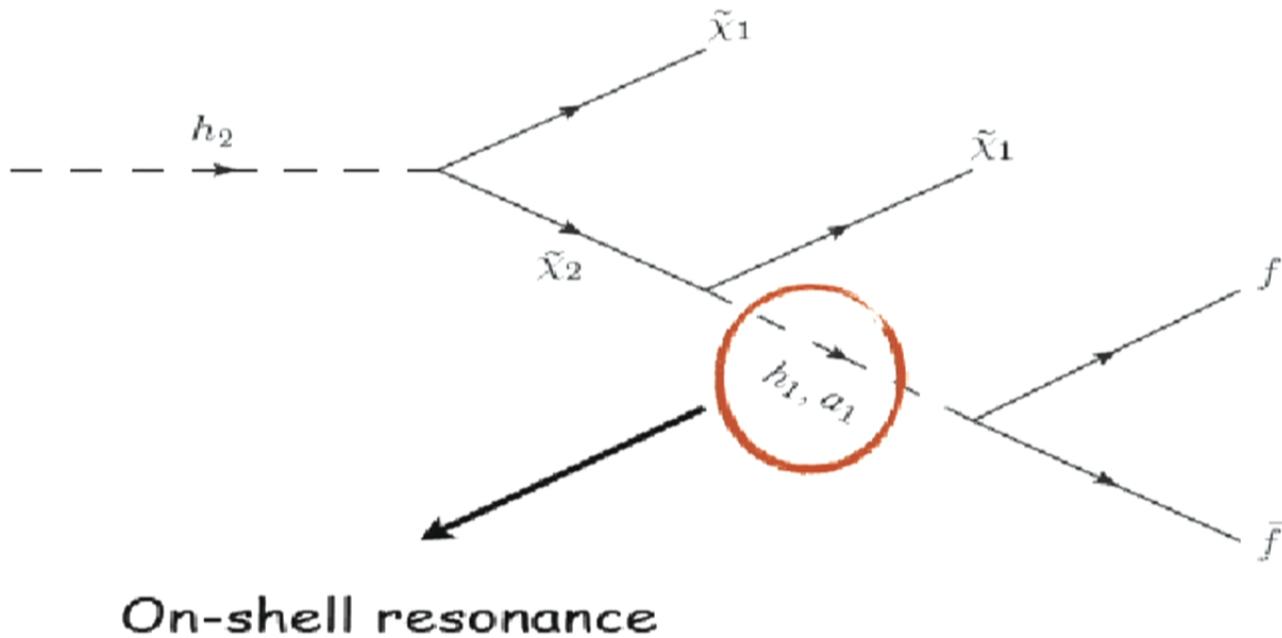


h_1 decay modes





Dark Light higgs search



Signal: Collimated Fermion pairs + MET



Benchmark points

► **Assumption:**

$\text{Br}(h_2 \rightarrow \chi_1 \chi_2) = 100\%$, $\text{Br}(\chi_2 \rightarrow \chi_1 h_1) = 100\%$, $\text{Br}(h_1 \rightarrow ff) = 100\%$

► **Parameters:**

- $m_{h_2} = 115\text{GeV}$ ($95\text{GeV} \sim 135\text{GeV}$)

- $m_{\chi_2} = 80\text{GeV}$

- $m_{\chi_1} = 10\text{GeV}$

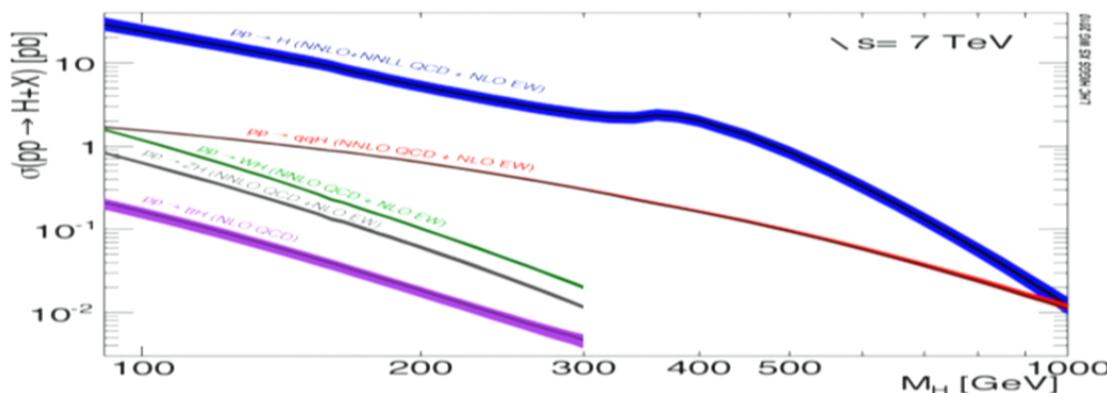
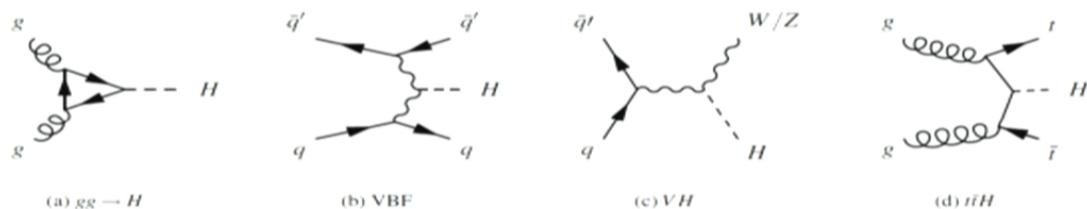
- $m_{h_1} = 1\text{GeV}(\mu\mu)$

$6\text{GeV}(\tau\tau)$

$15\text{GeV}(bb)$



SM Higgs Production



$h_2 (m_{h_2} = 115 \text{ GeV})$	$\sigma (\text{pb}) @ 7 \text{ TeV}$	$\sigma (\text{pb}) @ 14 \text{ TeV}$
Gluon Fusion	18.35	59.37
W/Z Fusion	1.393	4.771
Wh_2	0.7546	1.952
Zh_2	0.4107	1.130
$tt\bar{t} h_2$	0.1106	0.7699

Handbook of LHC Higgs, CERN-2011-002



Di-Muon Channel @ 7TeV

- ▶ Fairly Straight-forward:
Two close muons + MET + narrow invariant dimuon mass peak around m_{h_1}
- ▶ $Z h_2$ with $Z \rightarrow l\bar{l}$, $h_1 \rightarrow \mu\mu$
Almost no irreducible background
- ▶ $W h_2$ with $W \rightarrow l\nu$, $h_1 \rightarrow \mu\mu$
Also very easy to be discovered and the dominant background is from $W + (\gamma^* \rightarrow \mu\mu)$
- ▶ Event Generation
MG5/ME4 + pythia + PGS

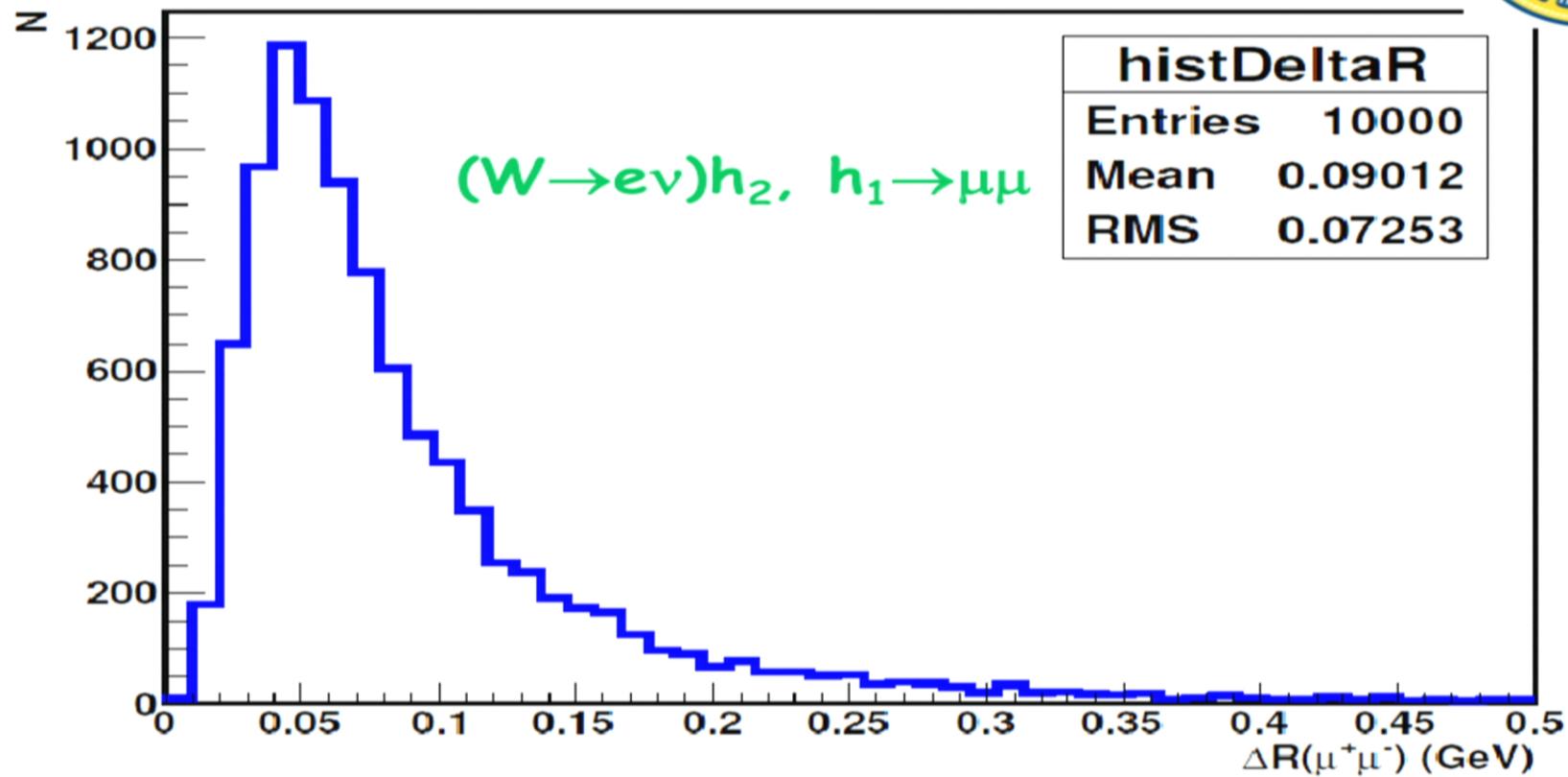


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$\Delta R(\mu\mu)$

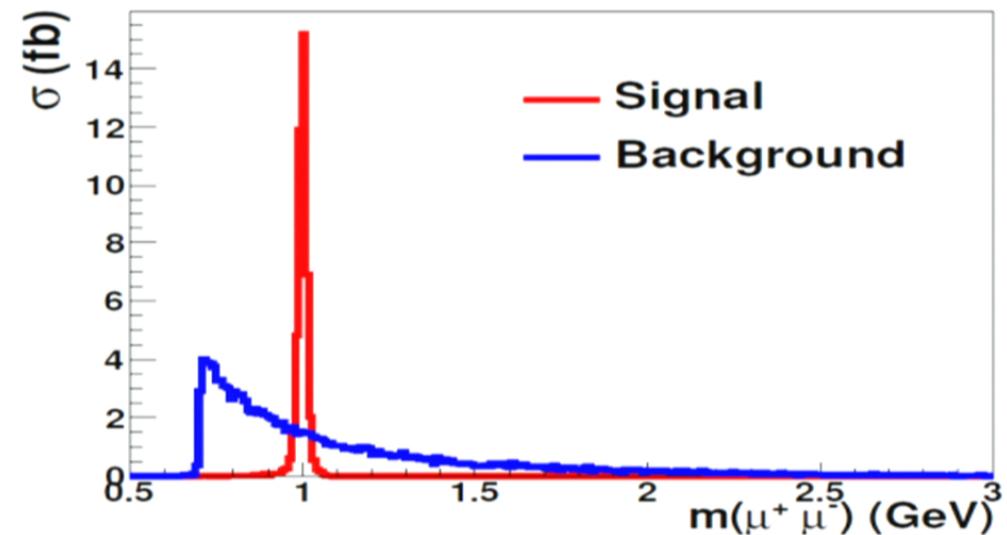
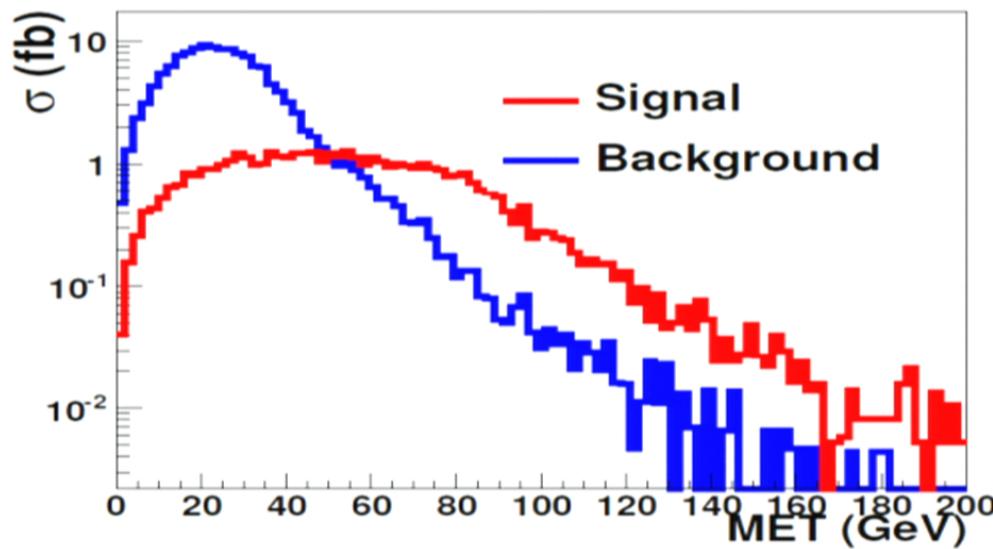




MET + $m(\mu\mu)$

The two most effective variables that can reduce the SM background are MET and $m(\mu\mu)$

$(W \rightarrow l\nu) h_2, h_1 \rightarrow \mu\mu$





Cut-Flow-Table

Basic Cuts	MET Cut	$m(\mu\mu)$
Reco (l); $ \eta_l \leq 2.4$; $PT\mu \geq 10\text{GeV}$; $PT_{l3} \geq 20\text{GeV}$; $\Delta R(\mu\mu) < 0.2$; $PT_{iso}(\mu) < 5\text{GeV}$	$\text{MET} \geq 40\text{GeV}$	$0.9\text{GeV} \leq m(\mu\mu) \leq 1.1\text{GeV}$

$(W \rightarrow l\nu) h_2, \quad h_1 \rightarrow \mu\mu$

Cut	Signal (0.162 pb)	Background (11.3 pb)
Basic Cuts	27.926%	1.2328%
MET	18.3396%	0.1764%
$m(\mu\mu)$	18.1546%	0.0365%

Note: Some preselection cuts have been applied



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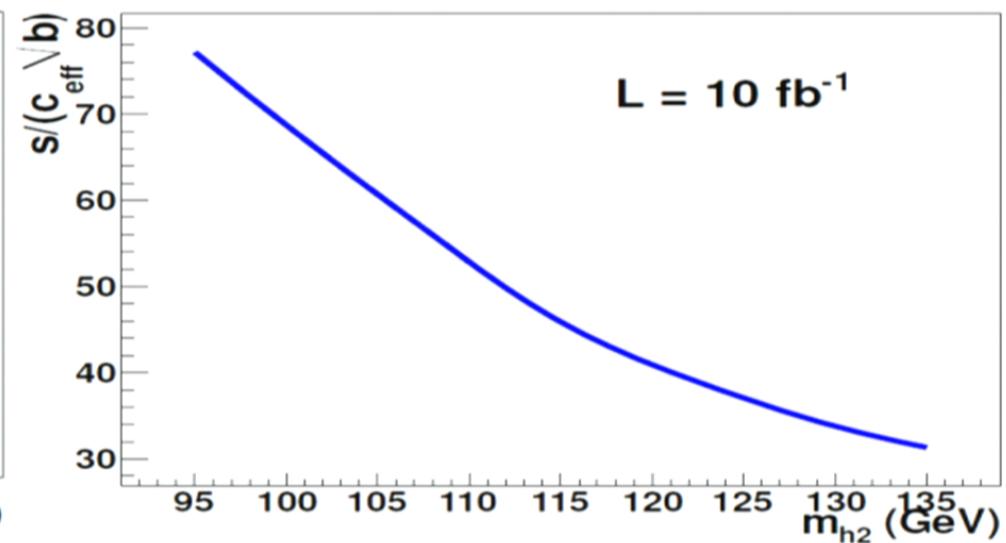
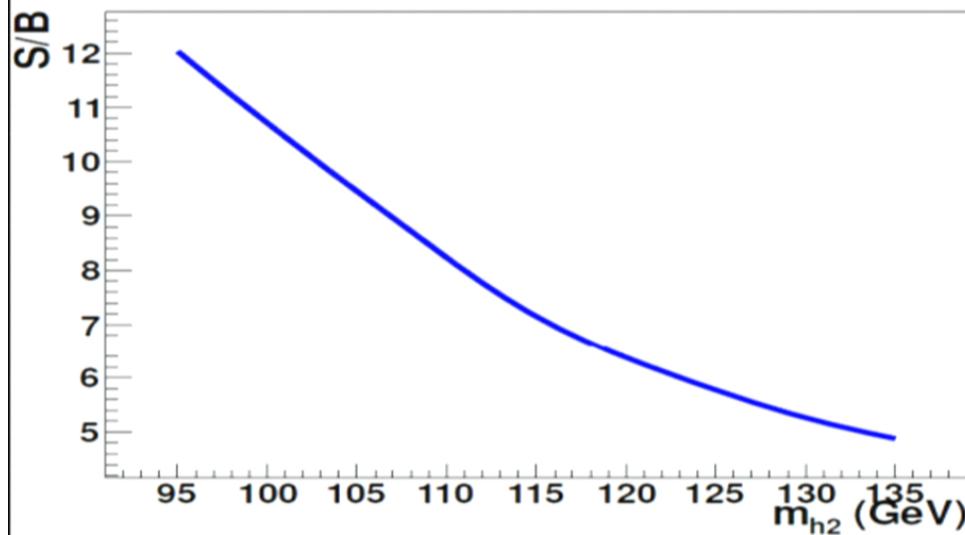
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Discovery Potential



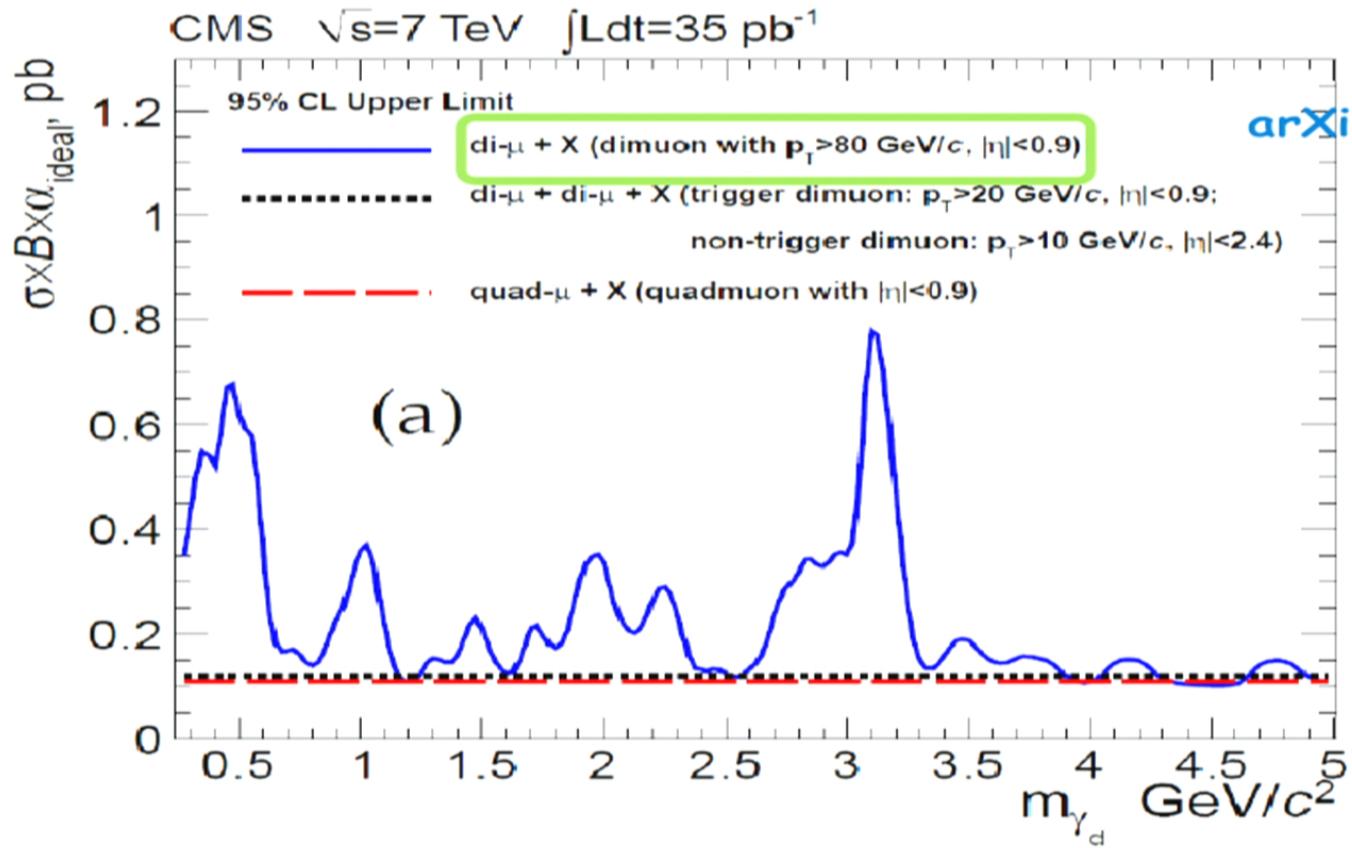
$(W \rightarrow l\nu) h_2, h_1 \rightarrow \mu\mu$



C_{eff} includes Branching ratio, K factors, etc



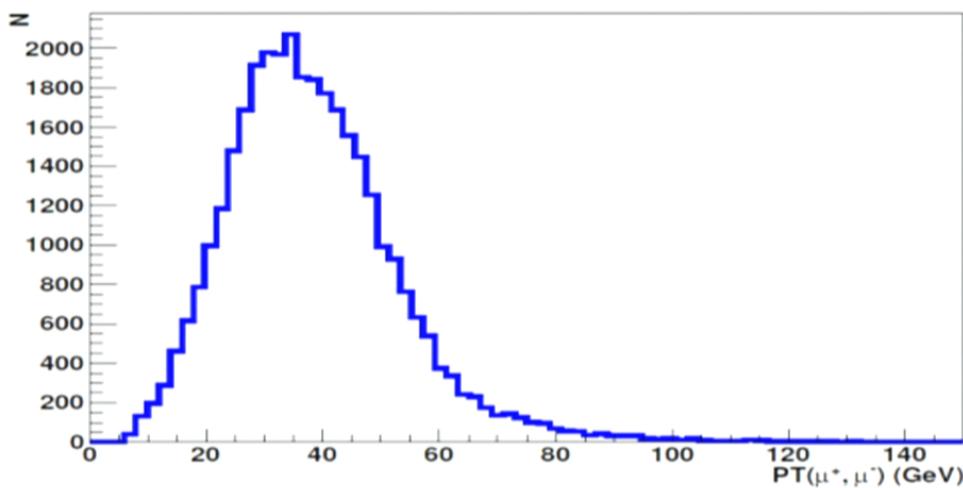
Dark Photon Search



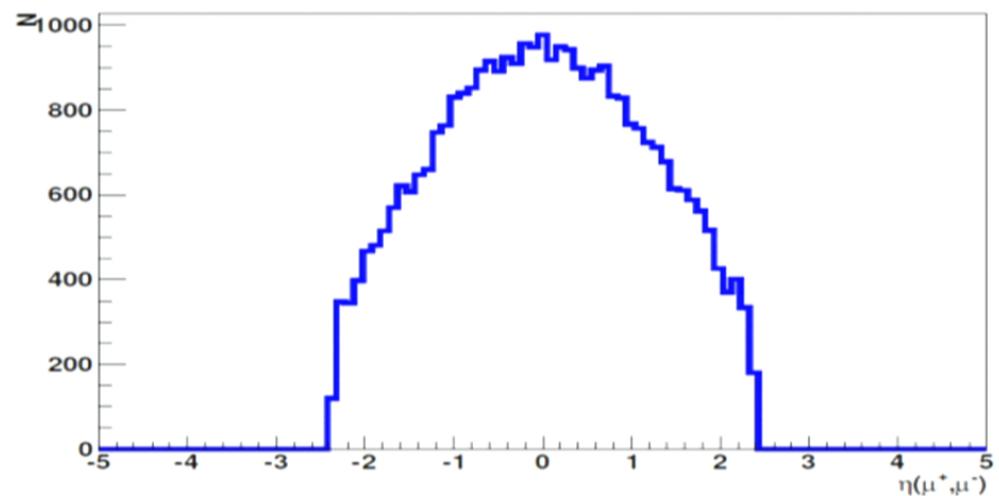
arXiv:1106.2375



Check $gg \rightarrow h_2, h_1 \rightarrow \mu\mu$



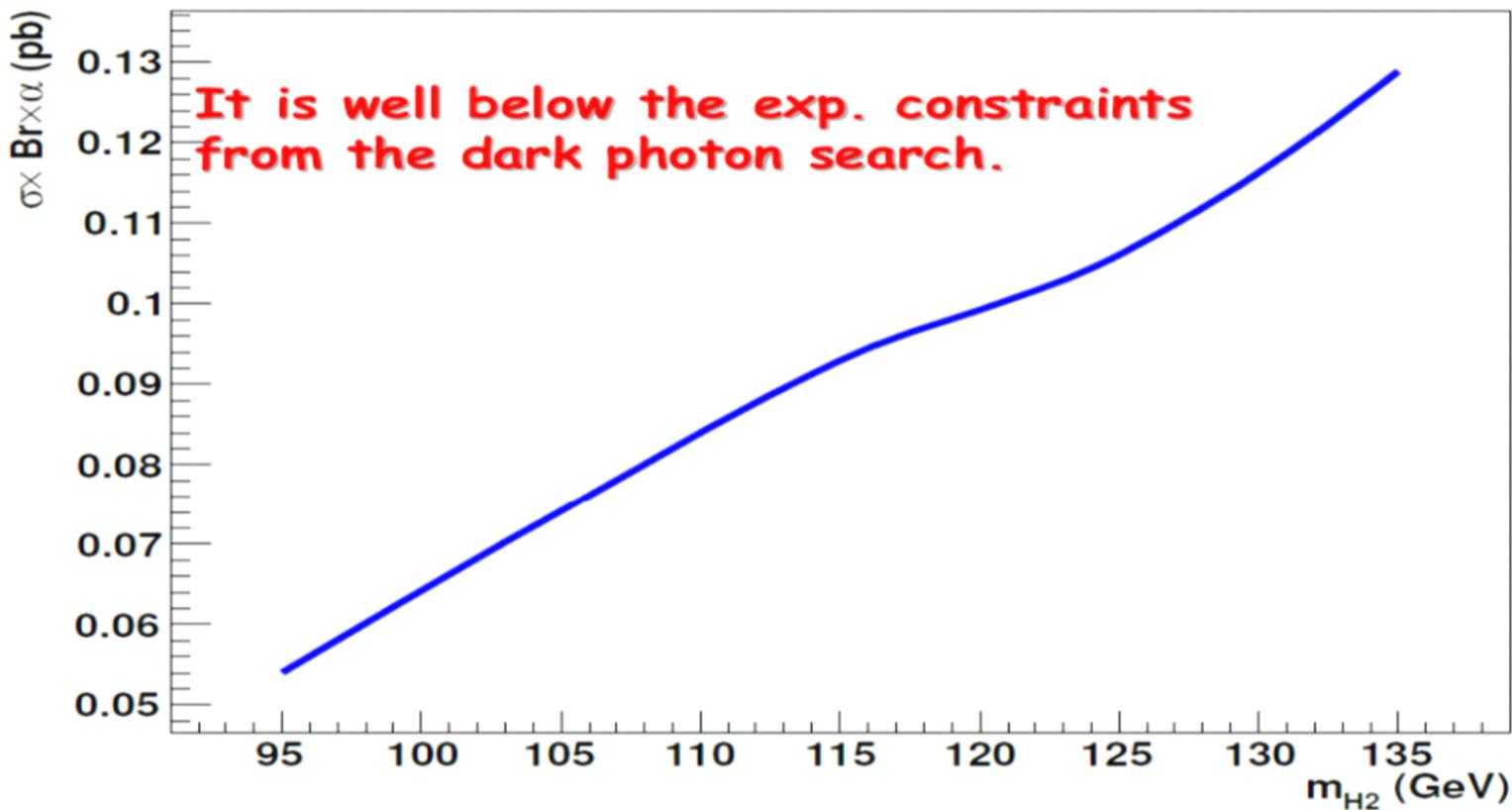
$\text{PT}(\mu\mu) > 80\text{GeV}$



$|\eta(\mu\mu)| < 0.9$



Check $gg \rightarrow h_2, h_1 \rightarrow \mu\mu$ (Cont.)





B-bbar channel @ 14TeV

- ▶ Bbbar channel is much more difficult
- ▶ ggfusion: bb + MET signal
 - overwhelmed by ttbar background
- ▶ VBF: bb + jets + MET
 - overwhelmed by ttbar background
- ▶ Wh₂: bb + ℓ+MET
 - overwhelmed by ttbar semileptonic background
- ▶ Zh₂: bb + ℓ+ℓ- + MET
 - can use Z mass window cut to control ttbar fully leptonic background
 - remaining Zg(g → bb) background is reduced by MET requirement
- ▶ tth₂: bb(+bb) + ℓ+MET or bb(+bb) + ℓ+ℓ- + MET
 - can isolate inclusive ttbar sample, and use MET and additional b-tag requirements to isolate signal



Z $h_2(t\bar{t}b\bar{b})$, ($h_1 \rightarrow b\bar{b}$)

- ▶ Z h_2 is the more promising channel
- ▶ Event Generation
 - Generate events using MG5/ME4, shower and hadronize with Pythia, cluster with FastJet (anti- k_T with $R = 1$)
 - Minimal detector simulation



Background

- Backgrounds
 - Z+jets
 - ttbar+jets
- Generate background using MG5/ME4
 - Z+jets for 0, 1, 2, and 3+jets
 - ttbar+jets for 0, 1, and 2+jets





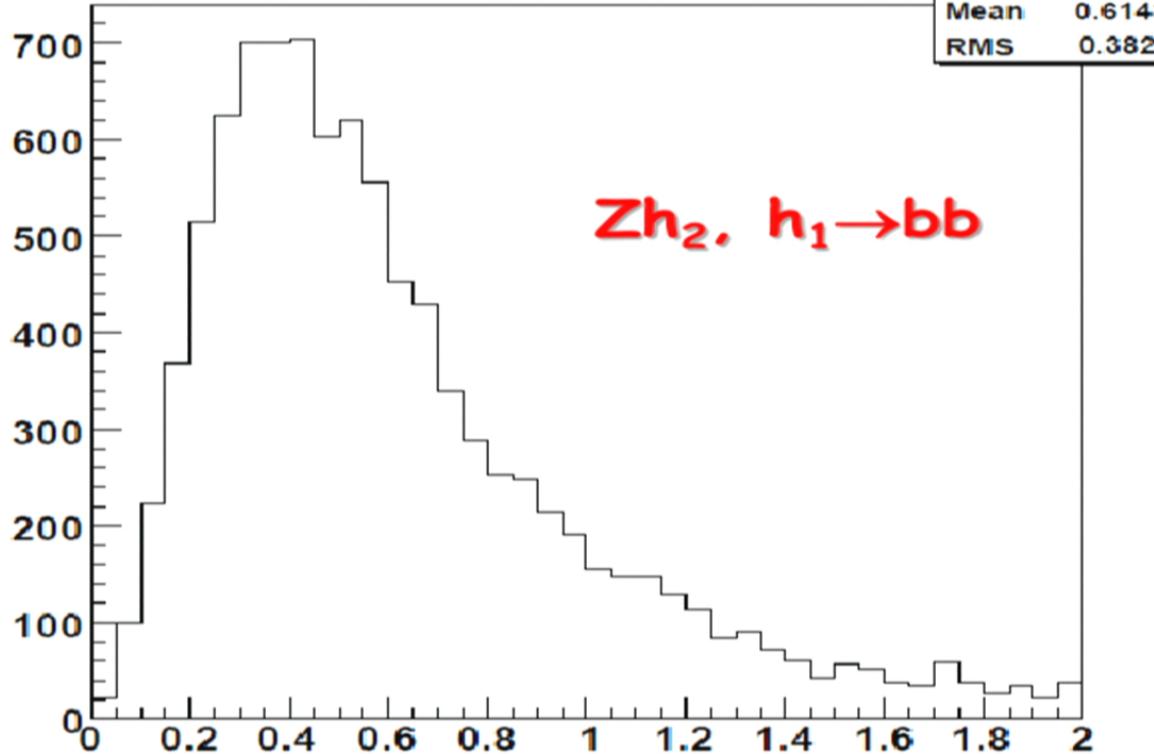
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 - ttbar+jets for 0, 1, and 2+jets



$\Delta R(bb)$

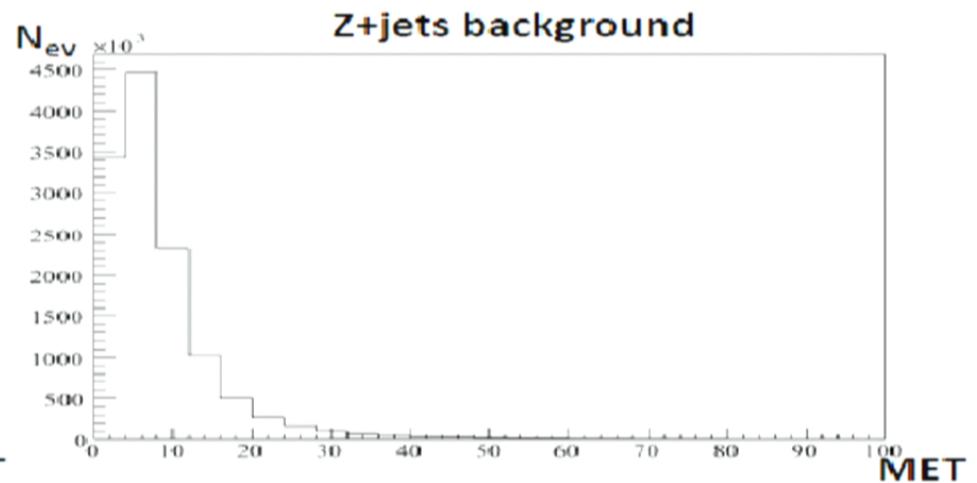
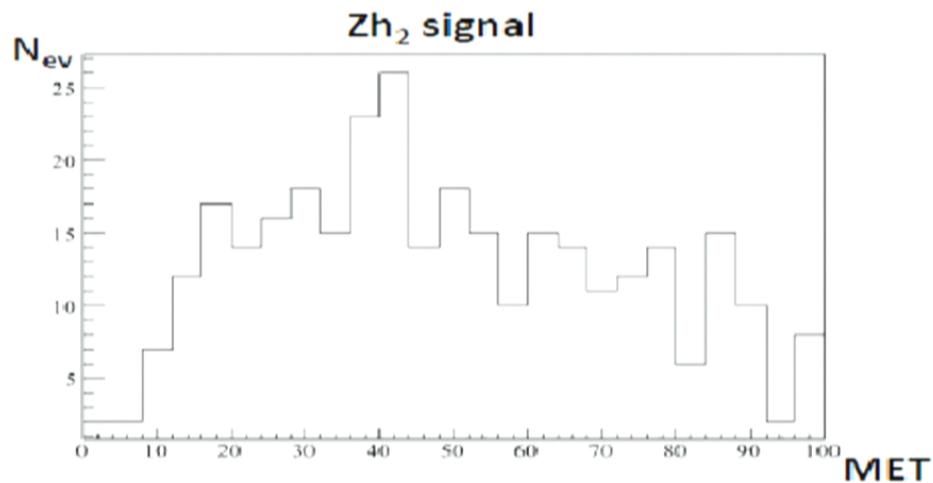
dR



MET distribution Comparison



10fb^{-1} @ 14TeV





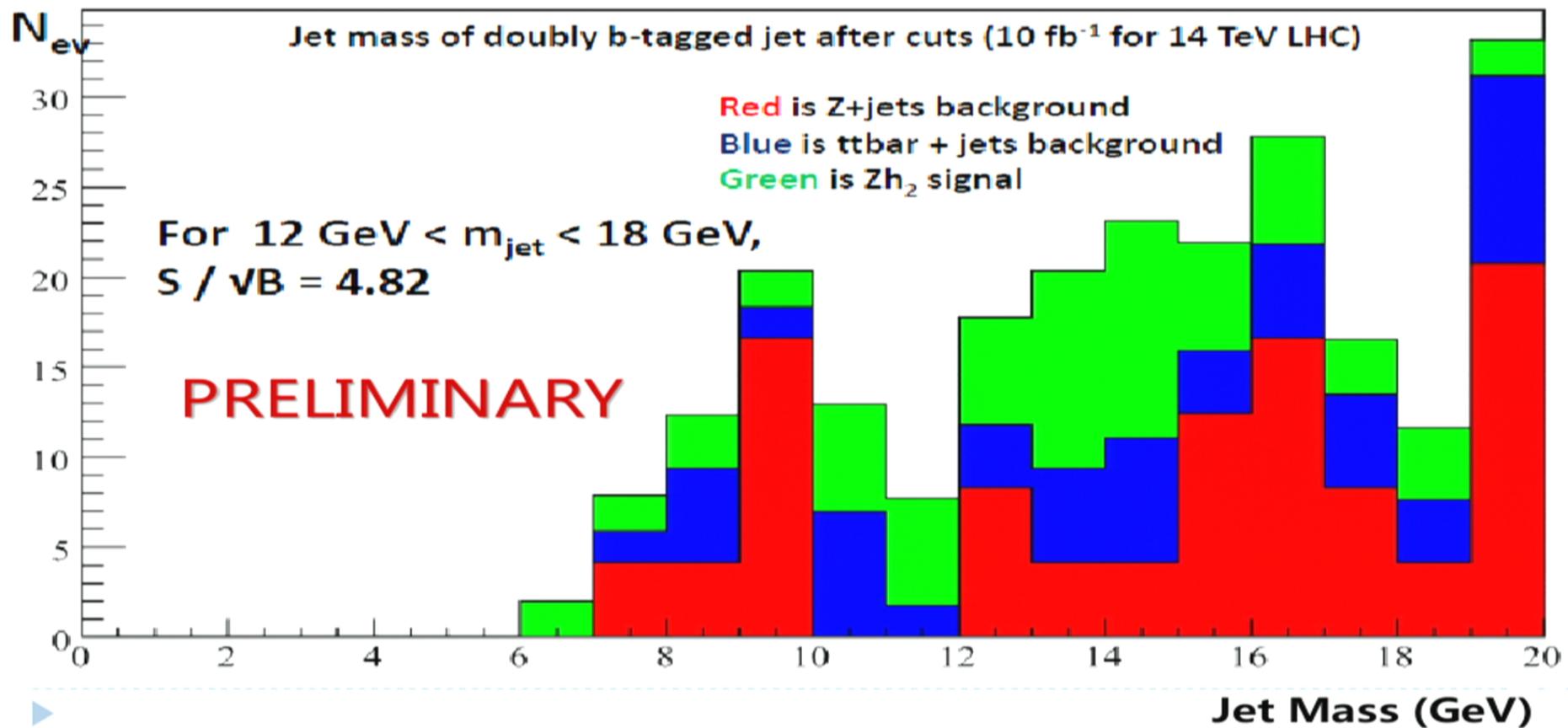
Cut-Flow-Table

Cut	Zh ₂ , Z→ll (0.114pb)	Z+jets, Z→ll (3002pb)	tt +jets (88.4pb)
=2l, Pt>20GeV	48.291%	51.775%	31.097%
Same flavor	48.028%	51.578%	15.544%
Opp. sign	47.940%	51.537%	14.364%
m _{ll} -m _Z <5GeV	37.248%	40.527%	1.056%
MET>40GeV	24.803%	1.015%	0.829%
Doubly b-tagged jet, P _T >20GeV	10.692	1.7433E-5	0.104%
P _T (h ₂ [*]) - P _T (ll) < 15GeV	7.274%	7.6095E-6	0.025%
N _{evt} (10fb ⁻¹ @ 14TeV)	83	55	127

Note: (80% b-tagging efficiency -cf. [ATLAS-CONF-2011-100](#))



Jet mass





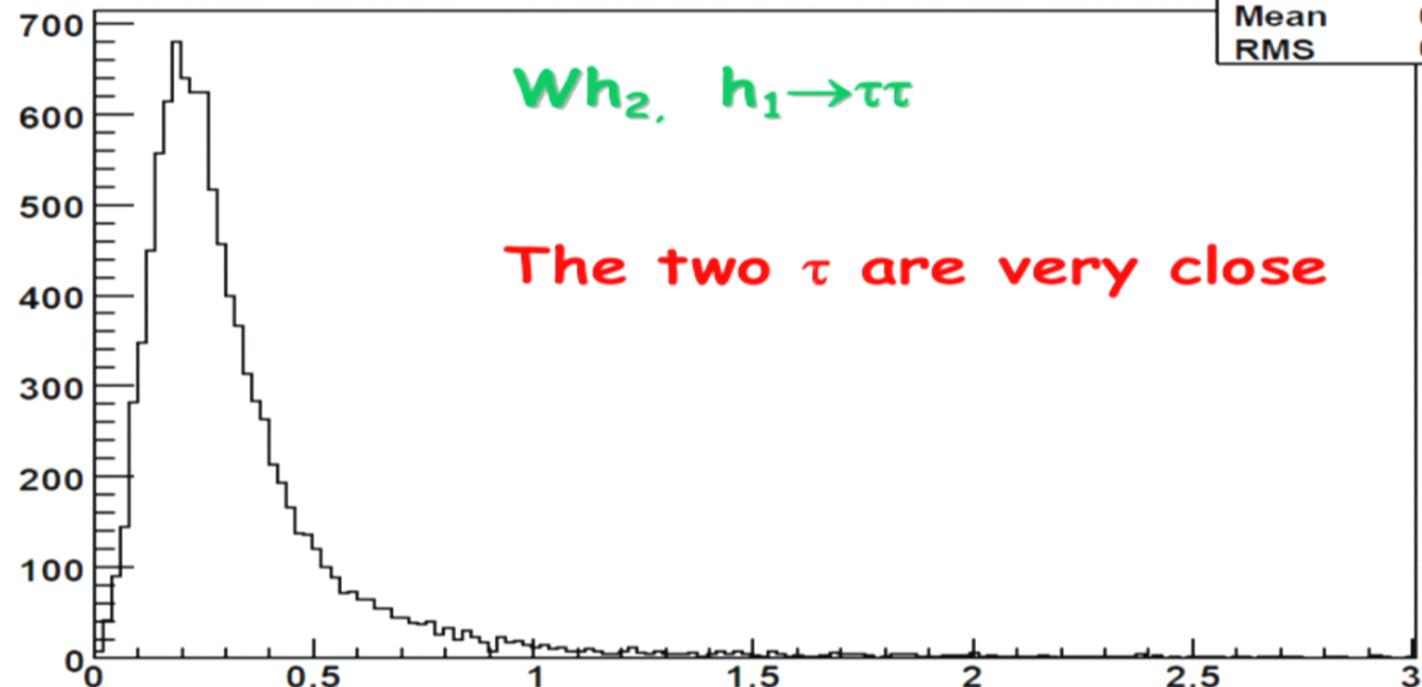
Di-tau Channel

- ▶ $\tau\tau$ channel is also much more difficult
- ▶ The two taus are very close to each other and the tau decay products are fairly soft, and we can NOT use the standard approach to identify the taus
- ▶ We treat di-tau as one jet and look for the jet substructure

$\Delta R(\tau\tau)$



dR

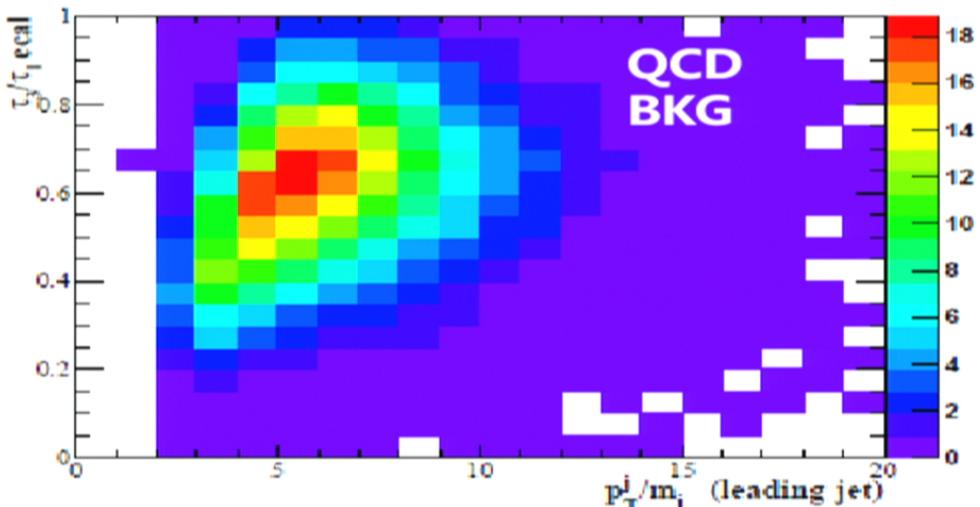
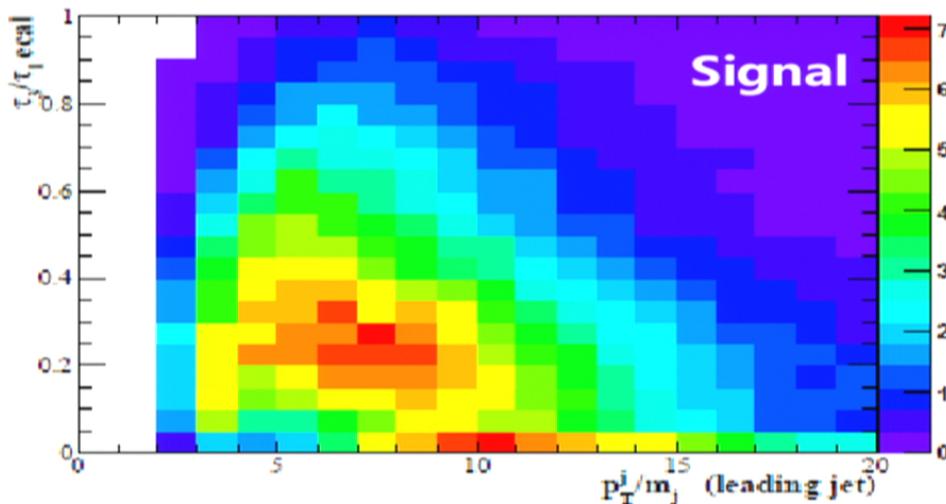




Di-Tau jet identification

N-Subjettiness: C. Englert, T. Roy, M. Spannowsky, arXiv:1106.4545;
J. Thaler, K. Van Tilburg, JHEP 1103 (2011) 015.

$$\tau_N = \frac{\sum_k P_{T,k} \min(\Delta R(1,k), \dots, \Delta R(N,k))}{\sum_k P_{T,k} R}$$



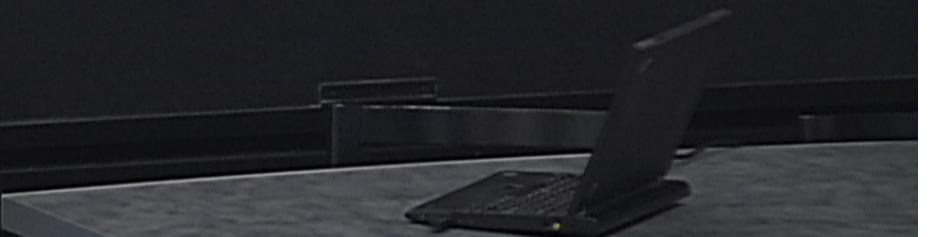
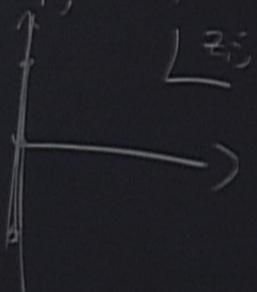
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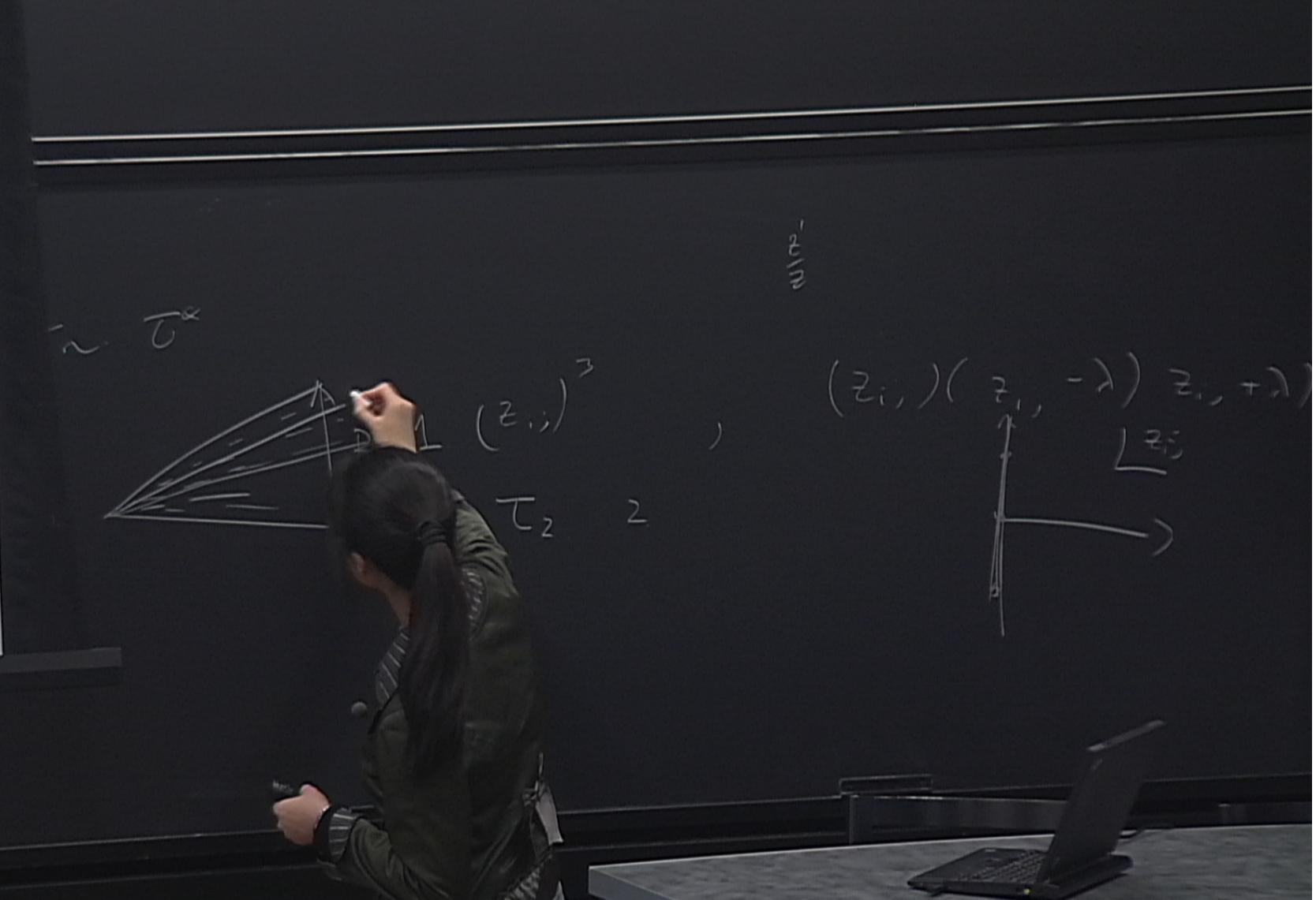
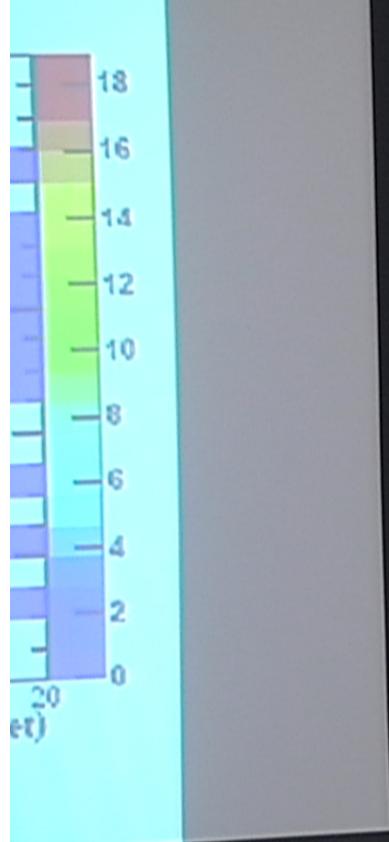


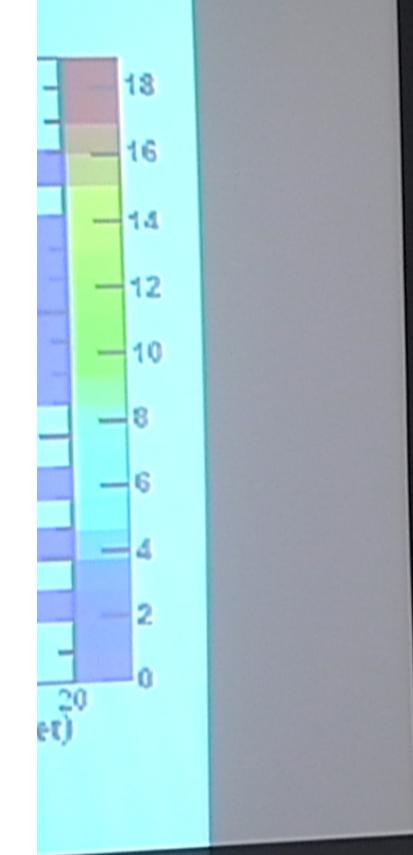
$$(z_{ij})^3$$

$\frac{z}{z_1}$

$$(z_{ij})(z_{ij} - \lambda) = z_{ij} + \lambda$$





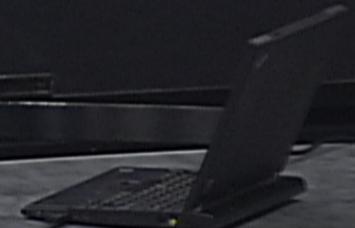
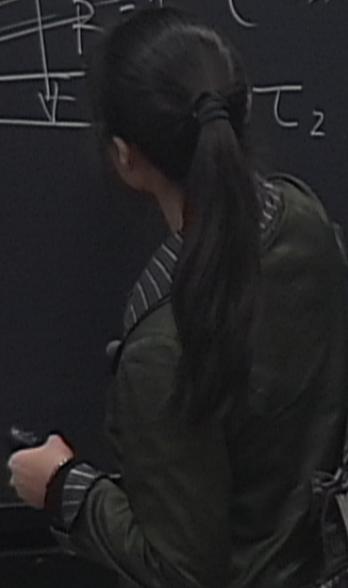
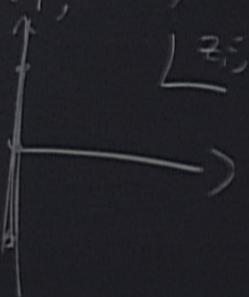


$\sim \tau^{\infty}$

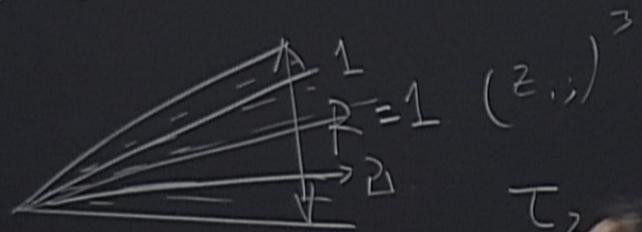
$$\begin{matrix} 1 \\ R-1 \\ 4 \end{matrix} \quad (z_{:,j})^3 \quad \tau_2 \quad 2$$

$\frac{1}{\tau}$

$$(z_{:,j})(z_{:,j} - \lambda) z_{:,j} + \lambda$$



$\sim \bar{\tau}^\infty$

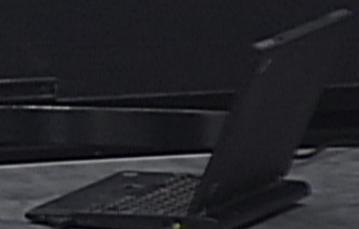
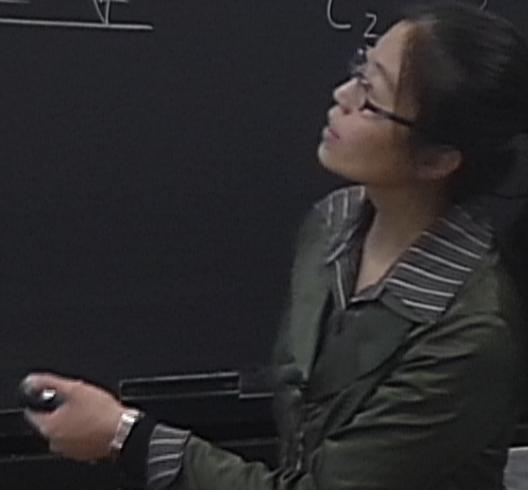
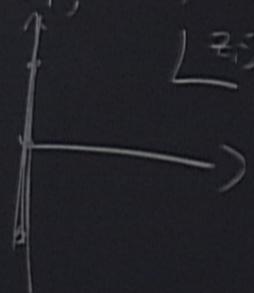


τ_2

$$(z_{\pm})^3$$

$$(z_+)(z_-, -\lambda) z_+ + \lambda)$$

$\underline{z_-}$





Conclusions

- ▶ DLH scenario provides a theoretical framework for studying non-standard Higgs phenomenology
- ▶ Many interesting channels to consider
 - $\mu\mu$ and bb preliminary results presented
 - $\tau\tau$ is underway
 - aim to provide a comprehensive LHC search strategy for SM-like higgs and light scalar resonances.

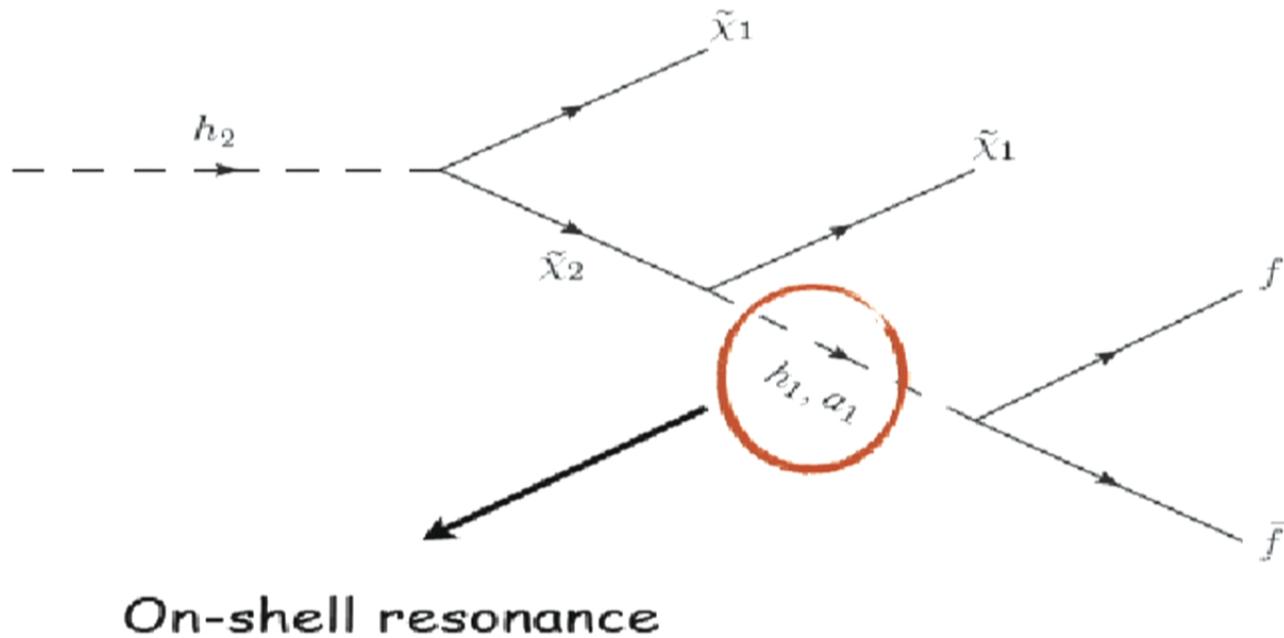


Conclusions

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Dark Light higgs search

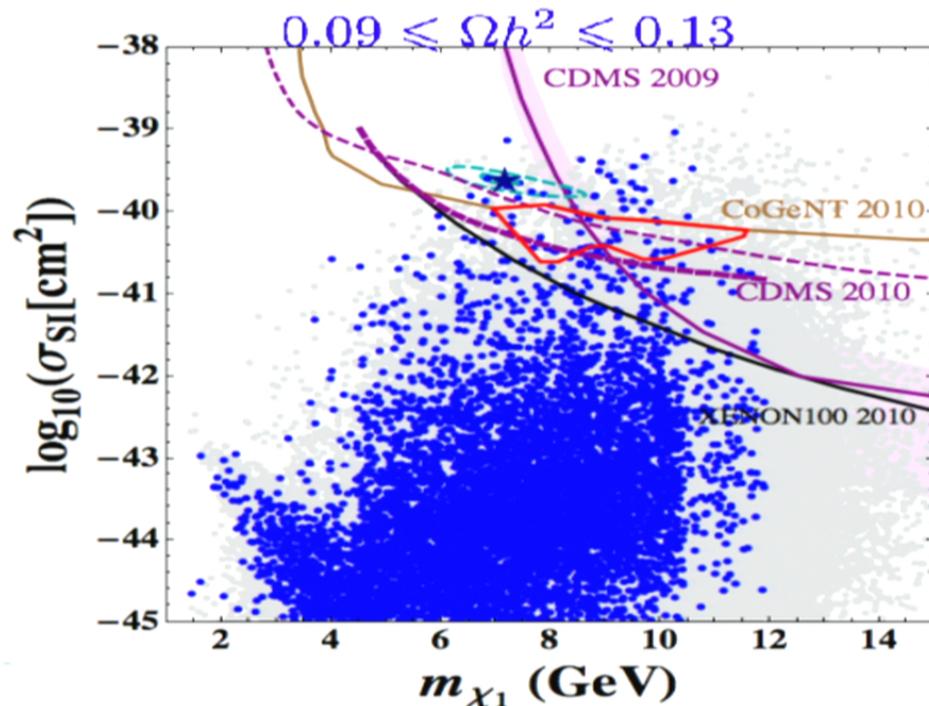


Signal: Collimated Fermion pairs + MET



Numerical Results

λ	$\kappa(10^{-3})$	$A_\lambda(10^3)$	A_κ	μ	$\tan \beta$	m_{h_1}
0.1205	2.720	2.661	-24.03	168.0	13.77	0.811
m_{α_1}	m_{χ_1}	m_{h_2}	Br hh	Braa	Ωh^2	$\sigma_{SI}(10^{-40})$
16.7	7.20	116	0.158%	0.310%	0.112	2.34



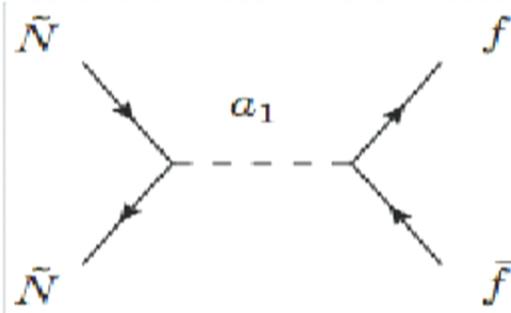
$0.05 \leq \lambda \leq 0.15, \quad 0.001 \leq \kappa \leq 0.005,$
 $|\varepsilon'| \leq 0.25, \quad -30\text{GeV} \leq A_\kappa \leq -15\text{GeV},$
 $5 \leq \tan \beta \leq 50, \quad 100\text{GeV} \leq \mu \leq 250\text{GeV}$

All points have passed the current exp. bounds of flavor physics, meson decays, and collider exp.

The blue points fall in a 3σ range of the observed relic density.

Their σ_{SI} can be as large as above 10^{-40} cm^{-2}

Breit-Wigner Effect



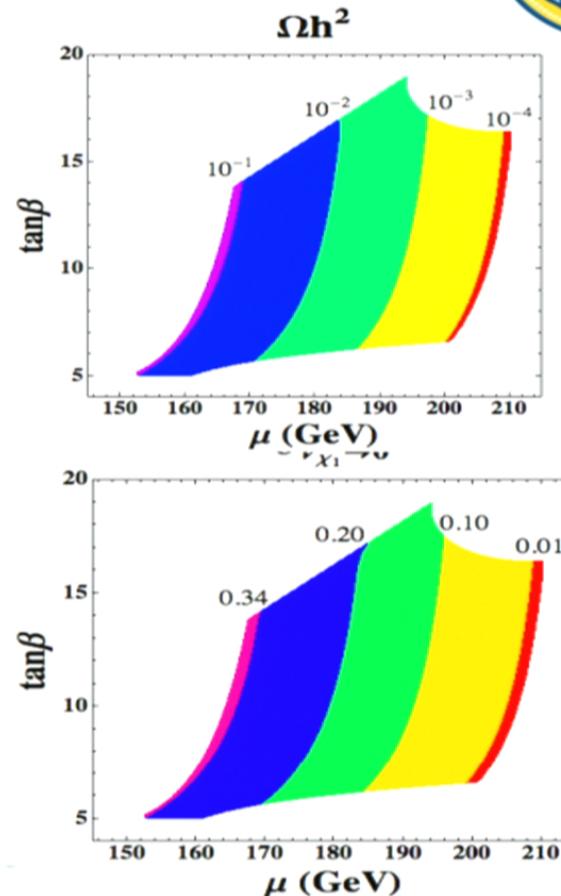
Thermal average of the LSP annihilation section

$$\sigma_{ff} v_{\chi_1} \approx \frac{3 |g_{\chi_1 \chi_1 \chi_1} g_{\chi_1 f f}|^2 (1 - m_f^2/m_{\chi_1}^2)^{1/2}}{32 \pi m_{\chi_1}^2 \left(\delta^2 + \left| \frac{\Gamma_{\chi_1} m_{\chi_1}}{4 m_{\chi_1}^2} \right|^2 \right)}$$

$$\delta = |(1 - v_{\chi_1}^2/4)^{-1} - m_{a_1}^2/(4 m_{\chi_1}^2)|$$

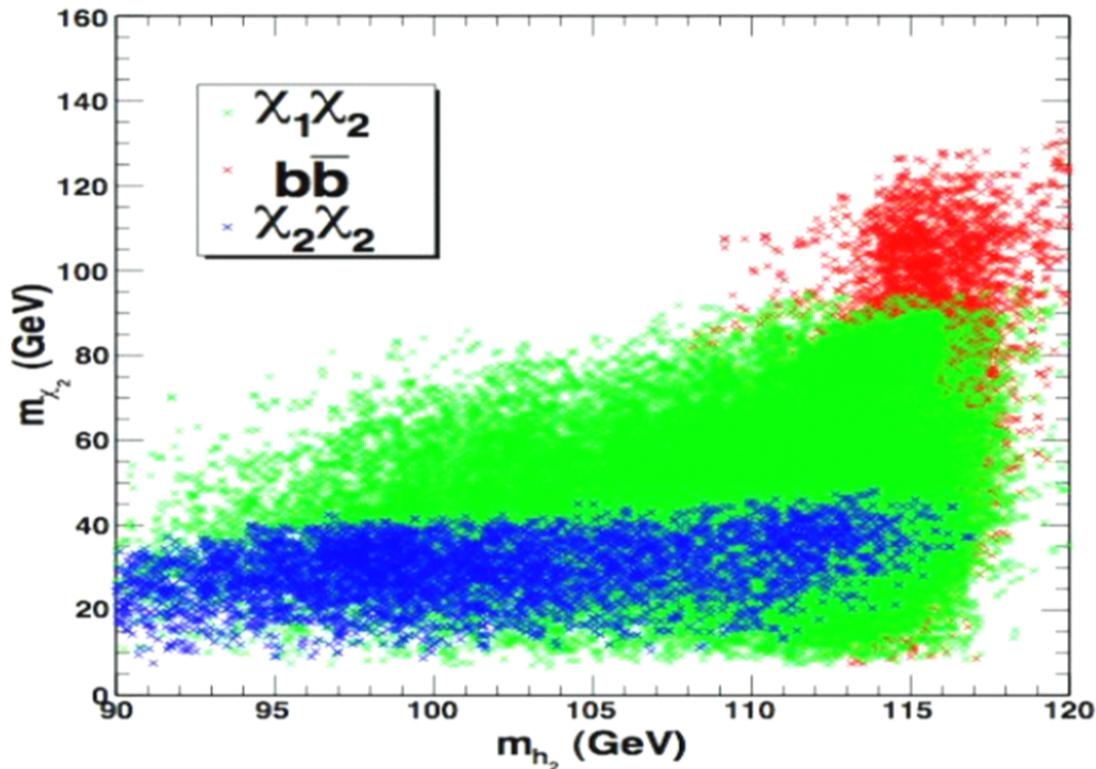
Relic density

$$\Omega h^2 \approx \frac{0.1 \left(\frac{m_{\chi_1}}{15 \text{ GeV}} \right) \left(\frac{\Gamma_{a_1}}{10^{-5} \text{ GeV}} \right) \left(\frac{\mu}{v} \right)^2 \left(\frac{0.003}{\kappa} \right)^2 \left(\frac{0.1}{\lambda} \right)^2}{\text{erfc} \left(\frac{2 m_{\chi_1}}{m_{a_1}} \sqrt{x_f \delta_{v_{\chi_1} \rightarrow 0}} \right) / \text{erfc}(2.2)}$$





h_2 decay modes



$h_2 \rightarrow \chi_1 \chi_2$ is typically dominant as long as it is kinematically allowed, and it is corresponding to the **GREEN** points.

$h_2 \rightarrow b\bar{b}$ mode can be dominant sometimes, but NOT generic.