

Title: Many-body Entanglement: From Topological Order to Quantum Computation

Date: Nov 29, 2011 09:30 AM

URL: <http://pirsa.org/11110140>

Abstract: Many-body entanglement, the special quantum correlation that exists among a large number of quantum particles, underlies interesting topics in both condensed matter and quantum information theory. On the one hand, many-body entanglement is essential for the existence of topological order in condensed matter systems and understanding many-body entanglement provides a promising approach to understand in general what topological orders exist. On the other hand, many-body entanglement is responsible for the power of quantum computation and finding it in experimentally stable systems is the key to building large scale quantum computers. In this talk, I am going to discuss how our understanding of possible many-body entanglement patterns in real physical systems contributes to the development on both topics. In particular, I am going to show that based on simple many-body entanglement patterns, we are able to (1) completely classify topological orders in one-dimensional gapped systems, (2) systematically construct new topological phases in two and higher dimensional systems, and also (3) find an experimentally more stable scheme for measurement-based quantum computation. The perspective from many-body entanglement not only leads to new results in both condensed matter and quantum information theory, but also establishes tight connection between the two fields and gives rise to exciting new ideas.

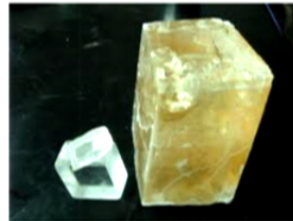


# Outline

- Background and question
  - Topological Order
  - What topological orders exist
- Results
  - Classification of 1D gapped topological order
  - Systematic construction of 2D topological order
- Approach
  - Many-body entanglement
  - Relation to quantum computation
- Outlook

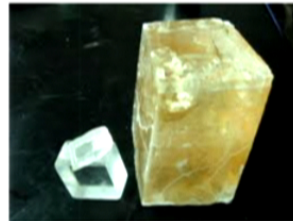
# Phase and Order: classical

Liquid  $\leftrightarrow$  Solid

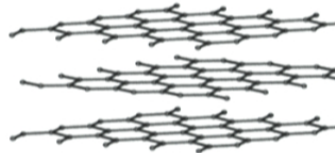
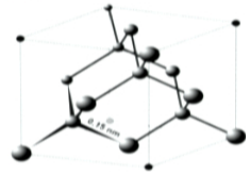


# Phase and Order: classical

Liquid ↔ Solid



Diamond ↔ Graphite



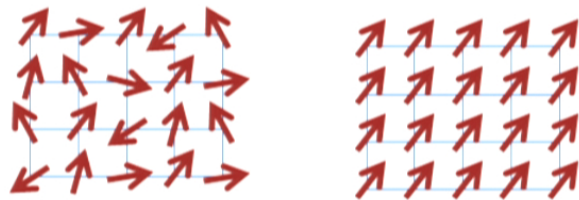
Thermal

Symmetry  
Breaking  
Order

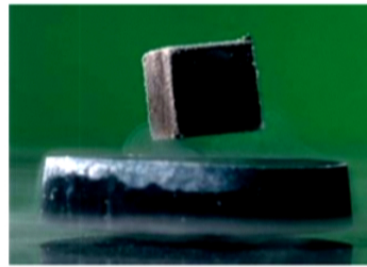


# Phase and Order: quantum

Paramagnet  $\leftrightarrow$  Ferromagnet

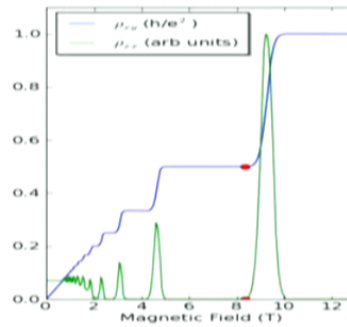
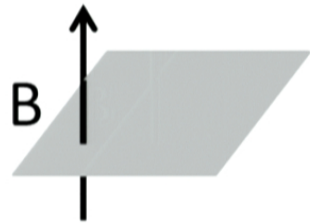


Insulator  $\leftrightarrow$  Superconductor



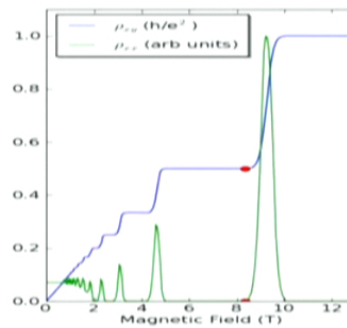
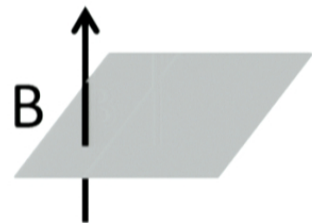
# Phase and Order: topological

Quantum Hall System  $\nu=1 \leftrightarrow \nu=2$



# Phase and Order: topological

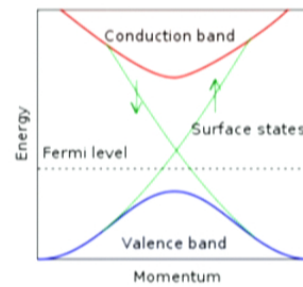
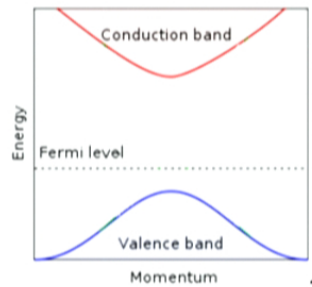
Quantum Hall System  $\nu=1 \leftrightarrow \nu=2$



Zero T

No Symmetry breaking

Normal Insulator  $\leftrightarrow$  Topological Insulator

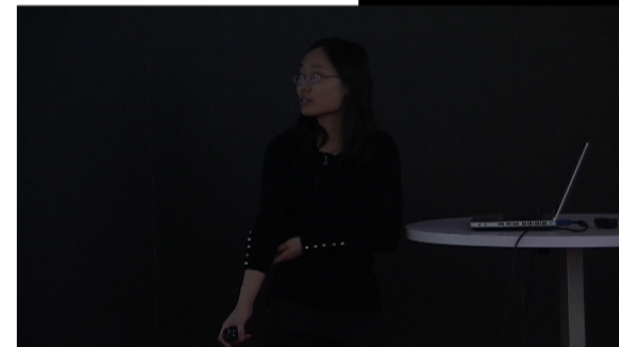




# Topological Order

## Generic Features

- No symmetry breaking



# Topological Order

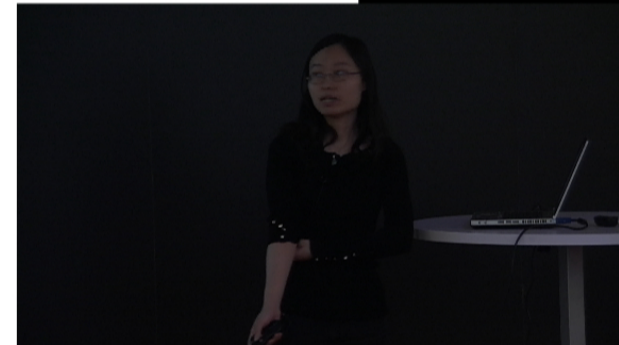
- 1D Haldane chain
  - Theoretical proposal: Haldane, 1983; Affleck, Kennedy, Lieb, Tasaki, 1987
  - Experimental realization: Hirakawa et al, 1986...
- 2D Quantum Hall
  - Experimental discovery: von Klitzing, 1980; Tsui, Stromer, 1982
  - Theoretical understanding: Laughlin, Halperin, Jain, Reed...
- 2D topological insulator
  - Theoretical proposal: Mele, Kane, 2005; Bernevig, Hughes, Zhang, 2006
  - Experimental realization: Konig et al, 2007



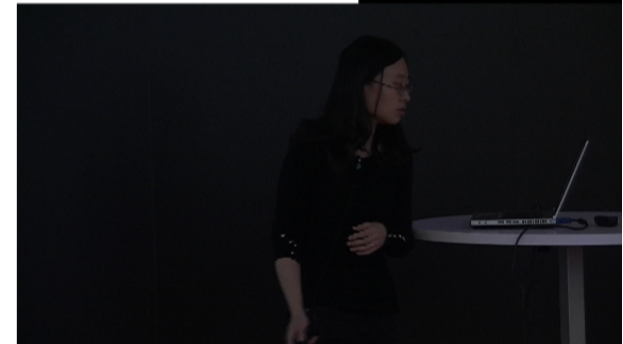
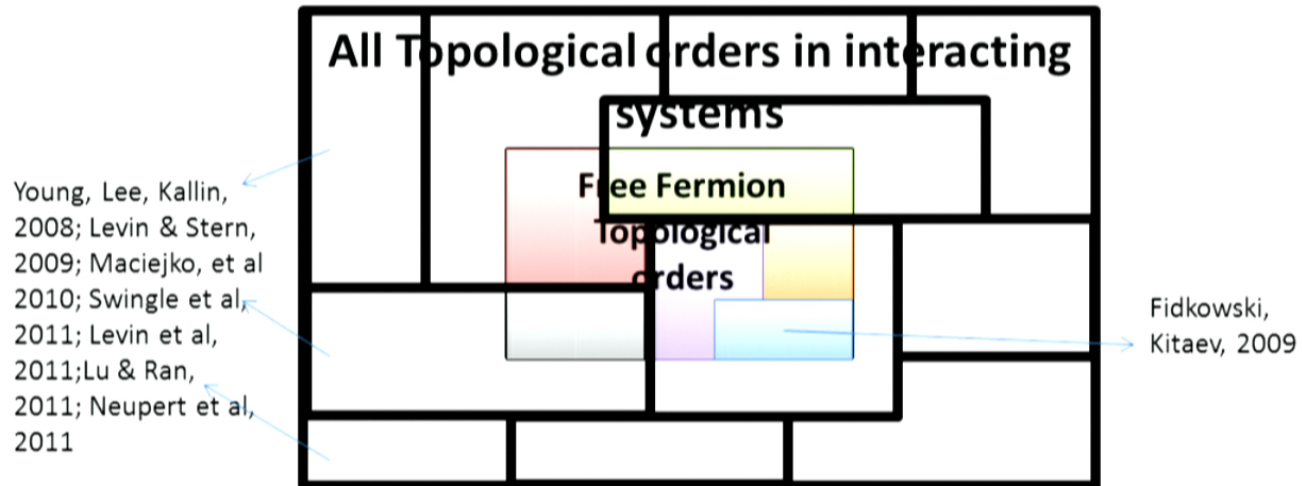
# Topological Order

**All Topological orders in interacting systems**

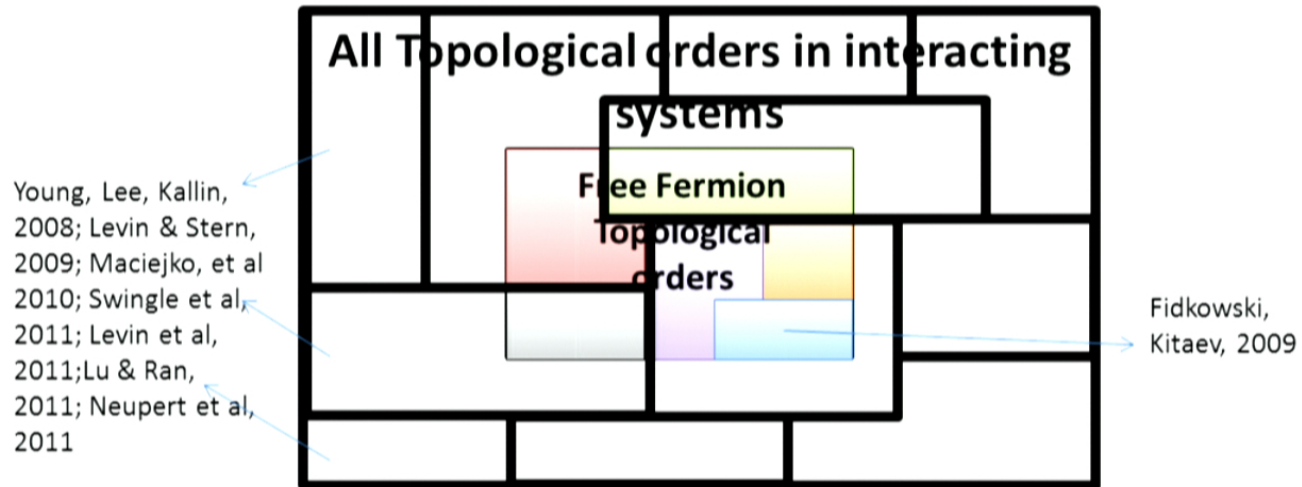
**Free Fermion  
Topological  
orders**



# Topological Order



# Topological Order



What topological orders exist in general **interacting** systems?

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# Classification of 1D topological order

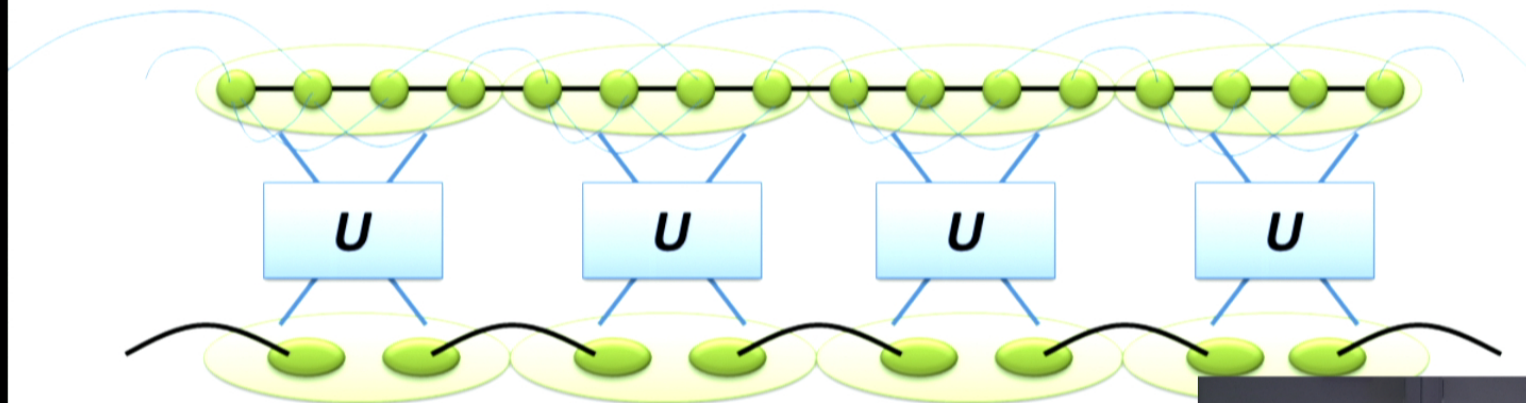


Fannes, Nachergaele, Werner, 1992; Verstraete, Cirac, Latorre, Rico, Wolf, 2004; Hastings, Verstraete, Cirac, 2006



# Classification of 1D topological order

Gapped system, finite correlation length



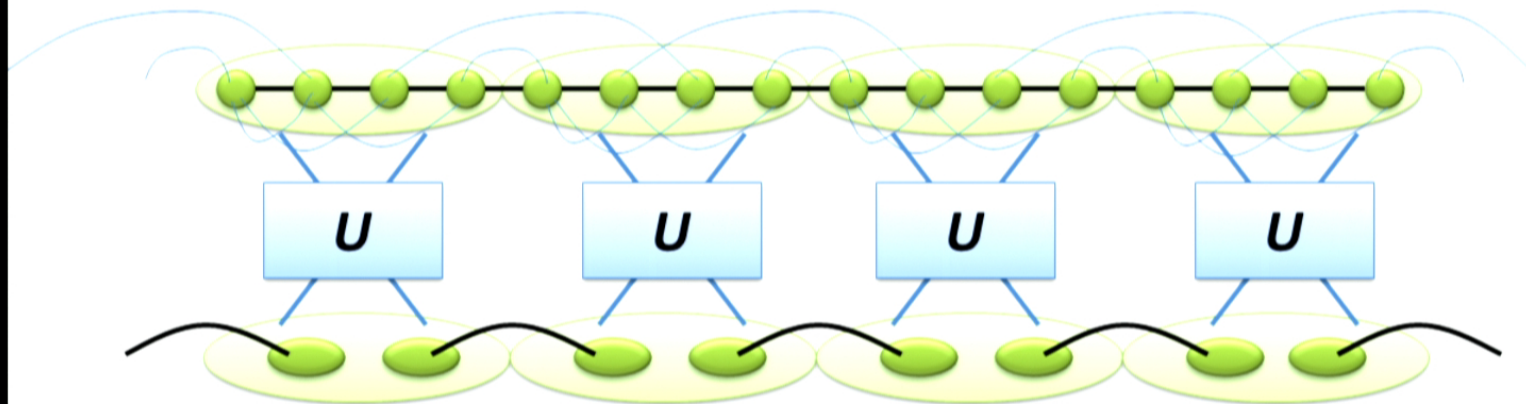
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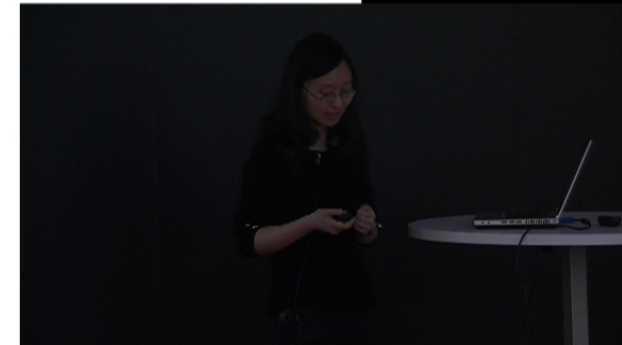


Matrix Product State Representation

Fannes, Nachergaele, Werner, 1992; Verstraete, Cirac, Latorre, Rico, Wolf, 2004; Hastings, 2007;  
Verstraete, Cirac, 2006

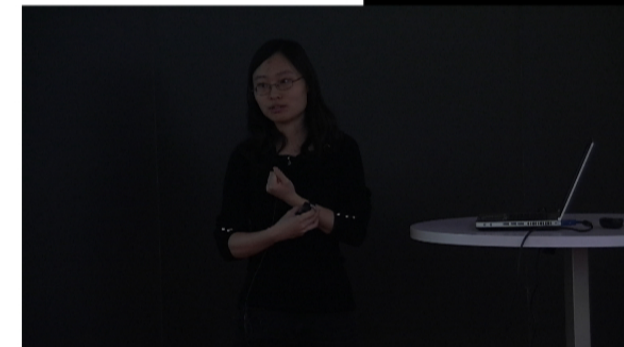
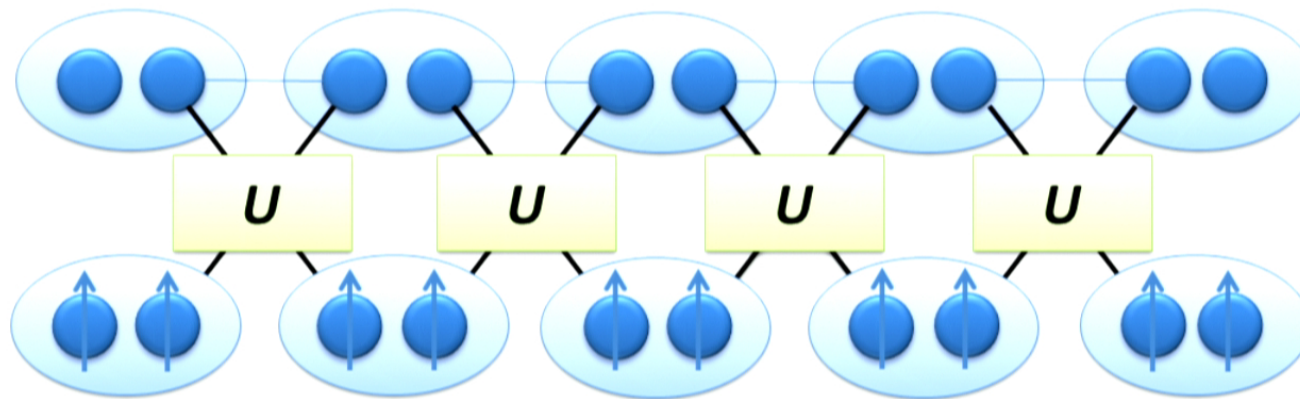
# Classification of 1D topological order

Gapped system, at large length scale



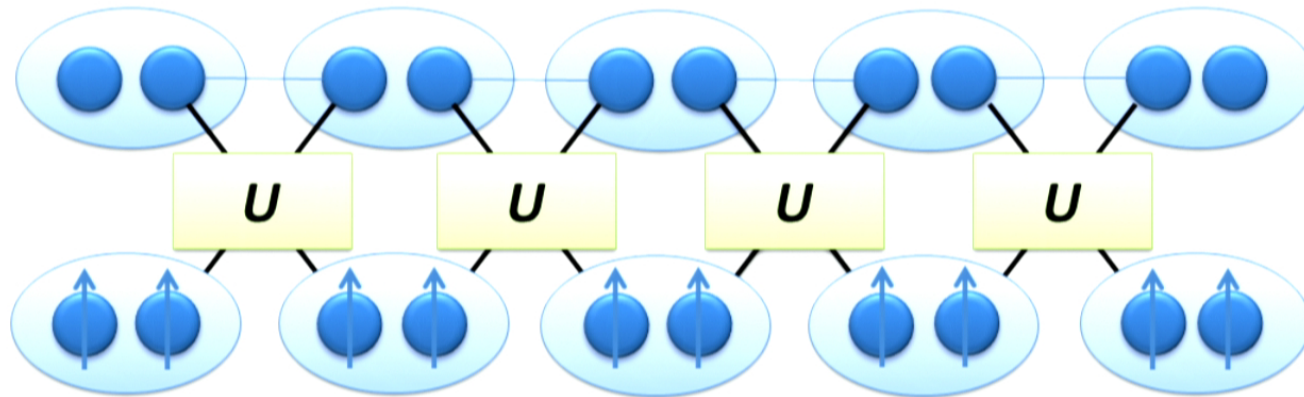
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# Classification of 1D topological order

Gapped system, at large length scale



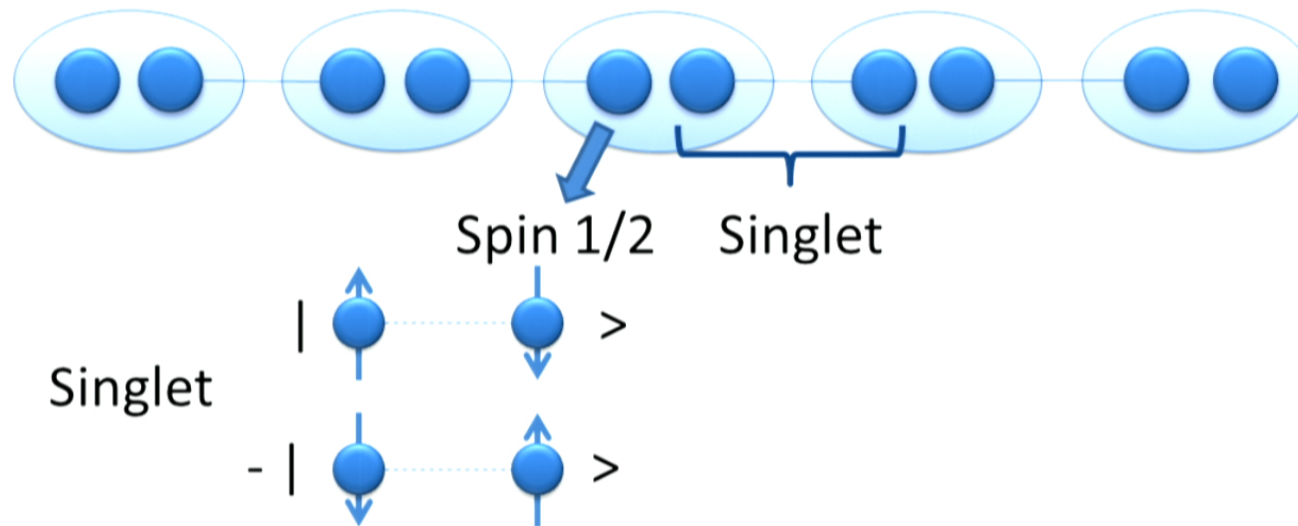
- Can be mapped to a classical state
- No intrinsic topological order in 1D

# Classification of 1D topological order

If symmetry exist

Example: spin rotation symmetry  $SO(3)$

Haldane chain,  
AKLT model



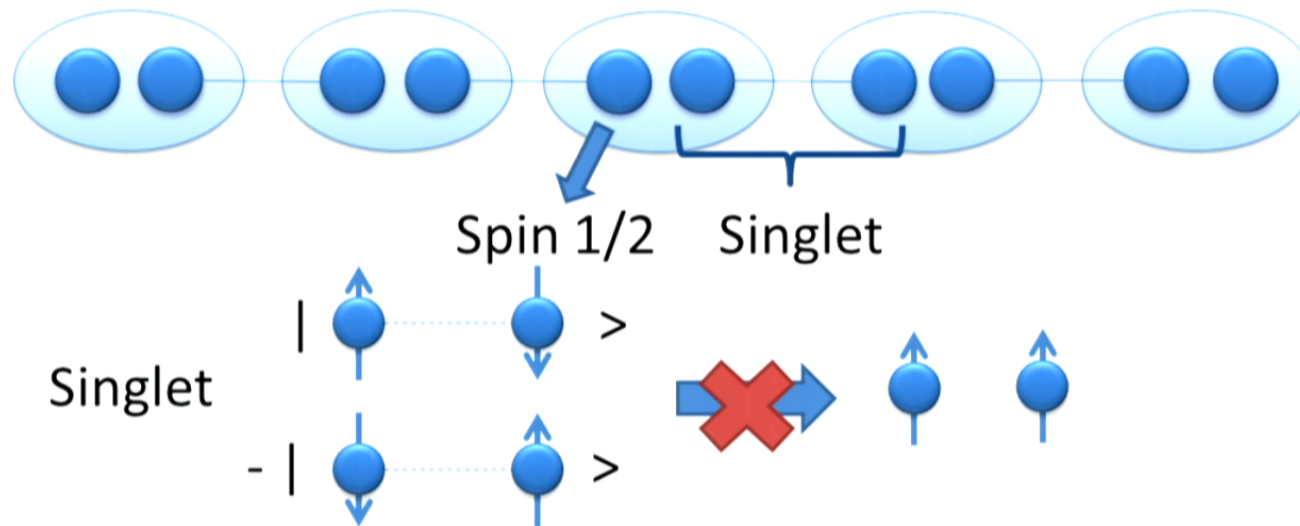
Haldane, 1983; Affleck, Kennedy, Lieb, Tasaki, 1987

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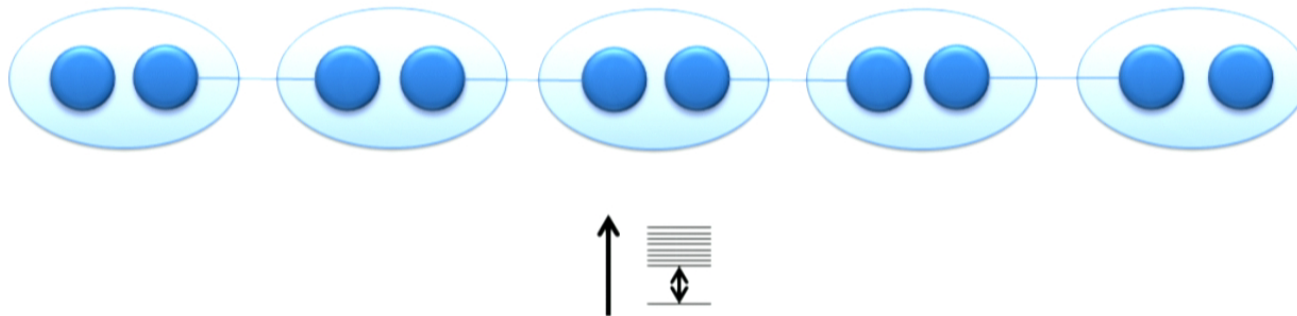
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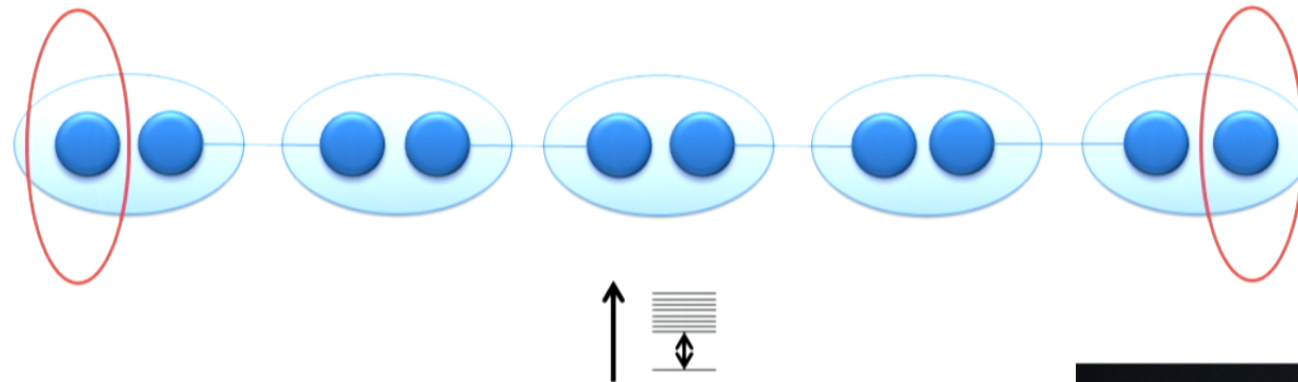
Turner, Pollmann, Berg, 2010; Pollmann, Berg, Turner, Oshikawa 2009

# Classification of 1D topological order

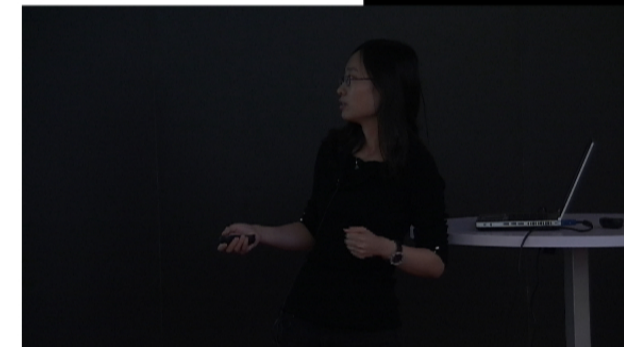
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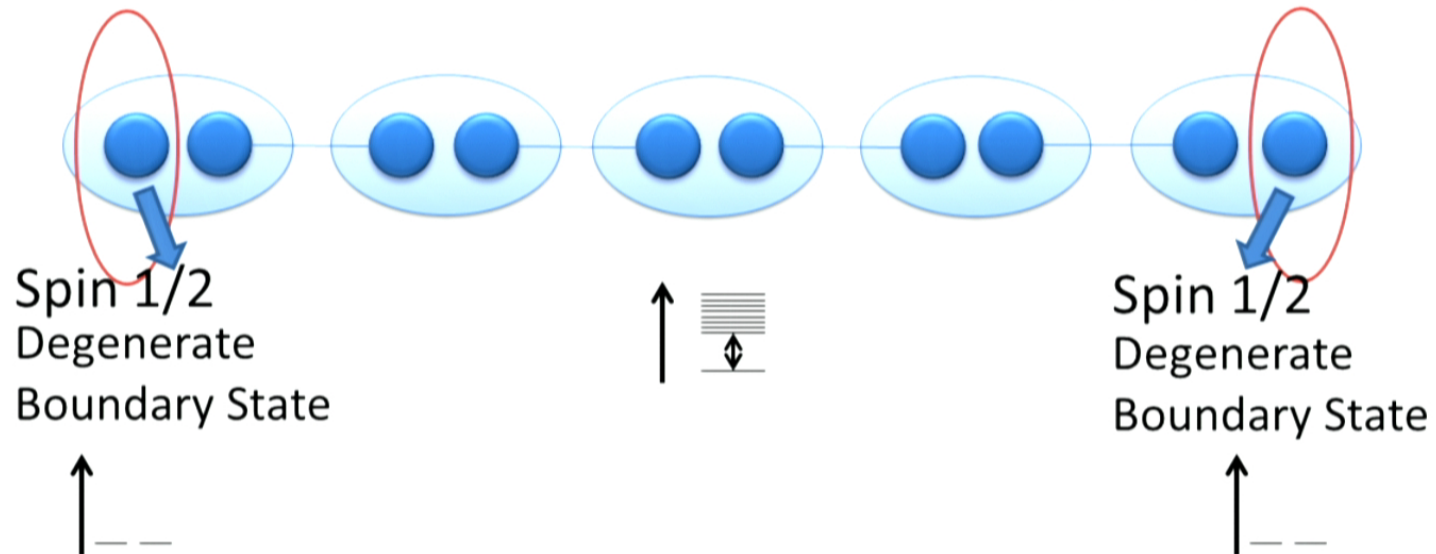


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If symmetry exist

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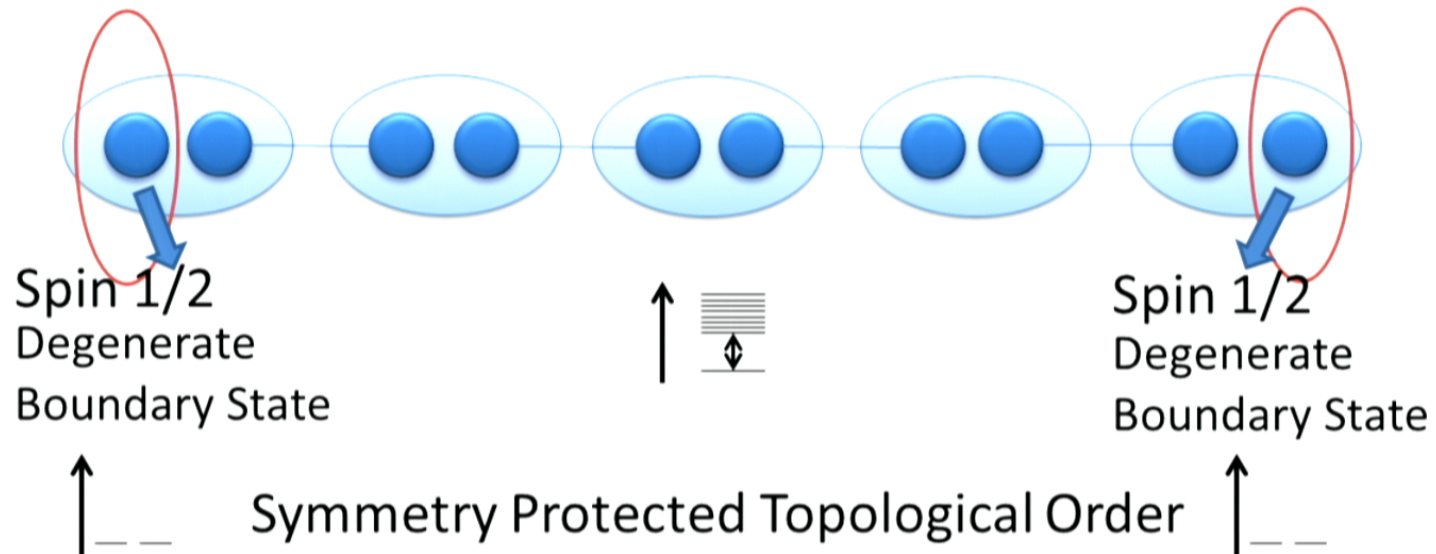
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If symmetry exist

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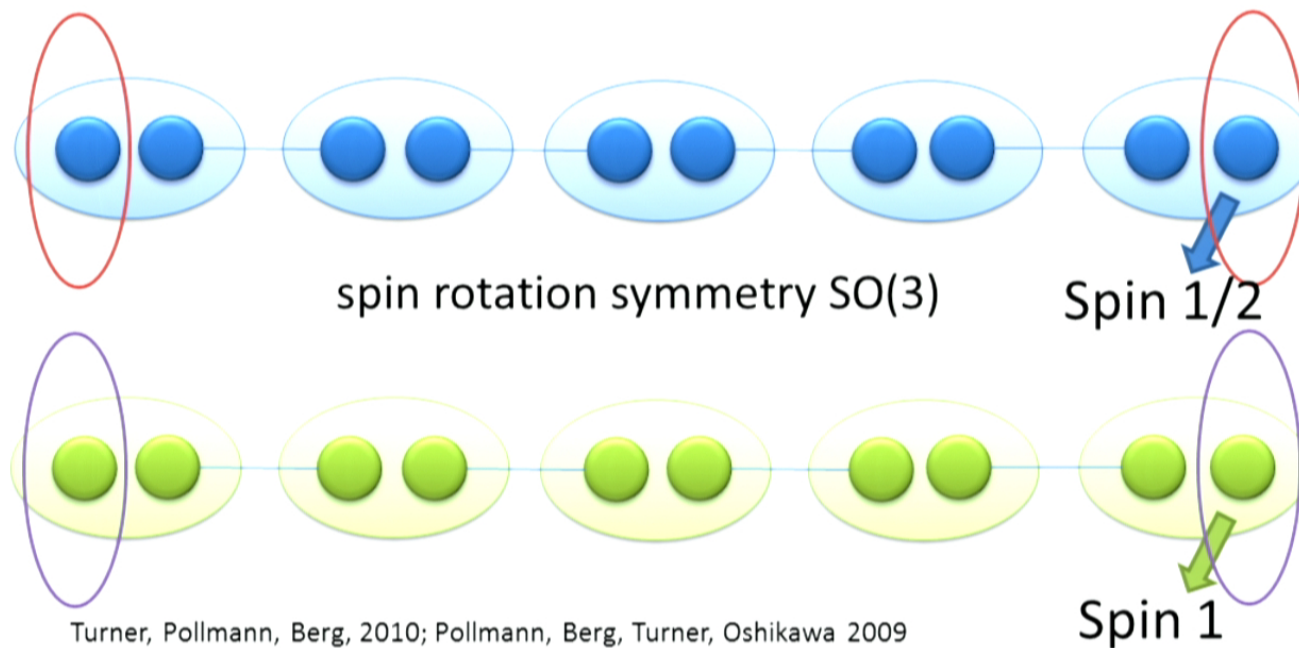
Haldane chain;  
AKLT model



Turner, Pollmann, Berg, 2010; Pollmann, Berg, Turner, Oshikawa 2009

# Classification of 1D topological order

Different boundary degrees of freedom correspond to different topological phases



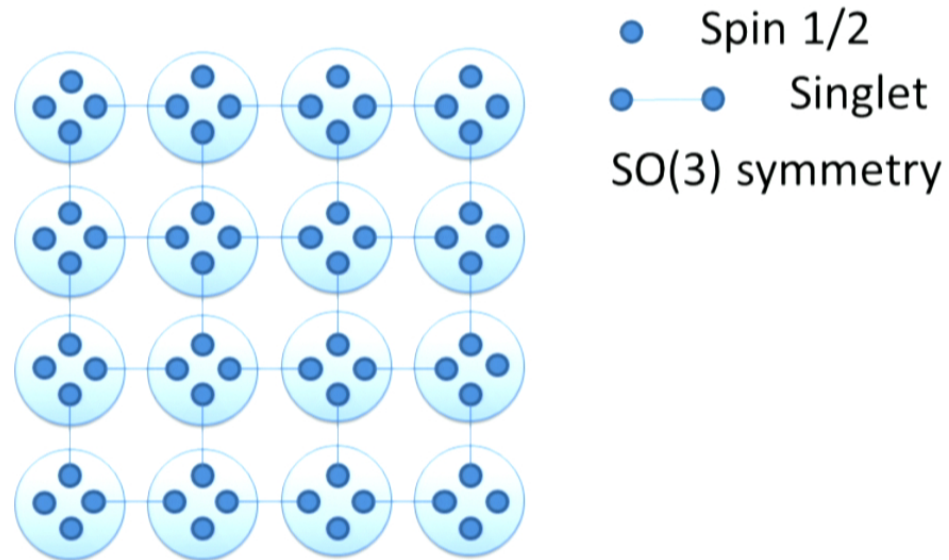
Turner, Pollmann, Berg, 2010; Pollmann, Berg, Turner, Oshikawa 2009

# Classification of 1D topological order

Chen, Gu, Wen, PRB, 83,025107, 2011

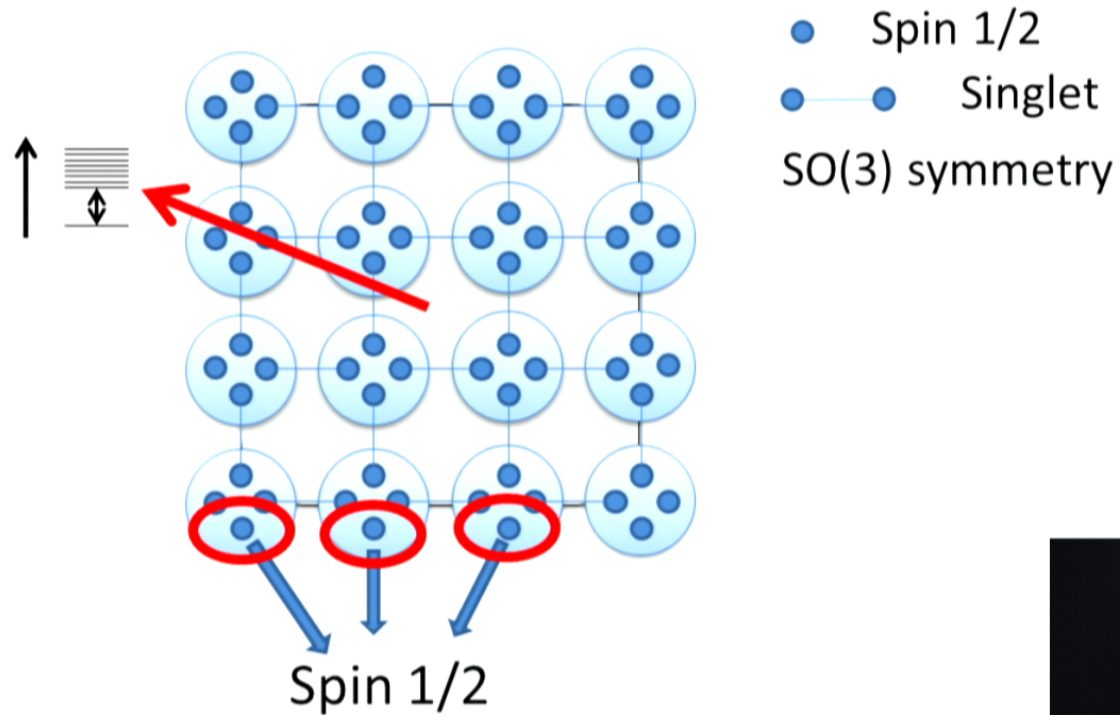
| Symmetry             | No. of Different Topological Orders | Symmetry              | No. of Different Topological Orders |
|----------------------|-------------------------------------|-----------------------|-------------------------------------|
| None                 | 1                                   | SO(3)                 | 2                                   |
| Time reversal        | 2                                   | Parity                | 2                                   |
| $D_2$                | 2                                   | Translation           | 1                                   |
| Translation + Parity | 4                                   | Time reversal + $D_2$ | 16                                  |

# Generalization to 2D: AKLT model



Affleck, Kennedy, Lieb, Tasaki, 1988; Schuch, Perez-Garcia, Cirac, 2010; Chen, Gu, Wen 2011

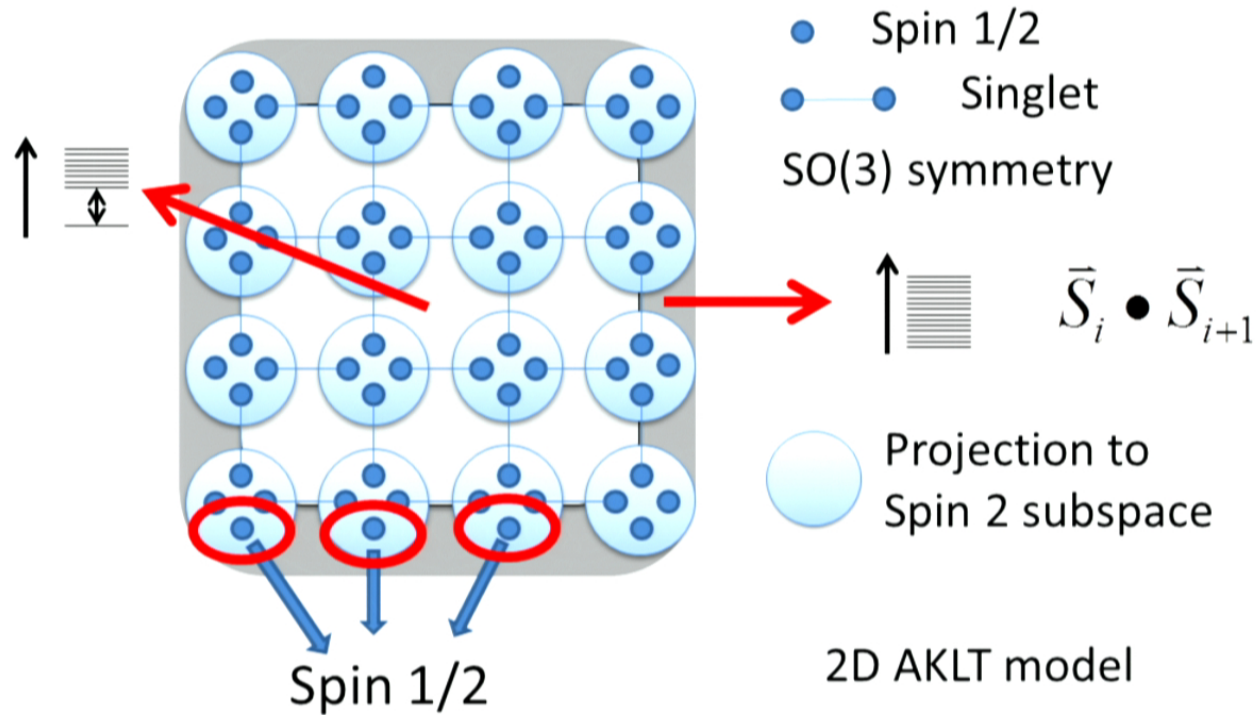
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# Generalization to 2D: AKLT model



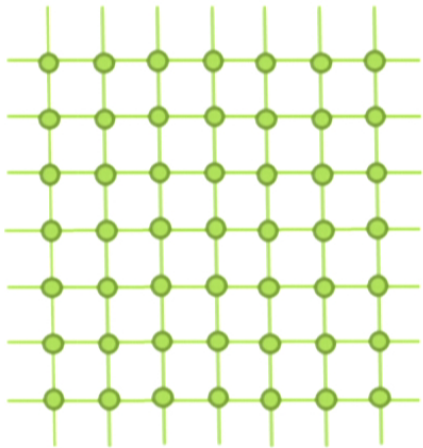
Affleck, Kennedy, Lieb, Tasaki, 1988; Schuch, Perez-Garcia, Cirac, 2010; Chen, Gu, Wen 2011

# Comparison with topological insulator

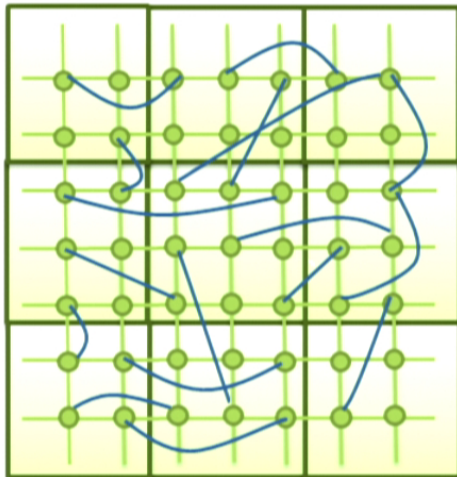
|  | 2D AKLT | Topological Insulator |
|--|---------|-----------------------|
| Gap in the bulk                                      | YES     | YES                   |
| Gapless excitation on boundary                       | YES     | YES                   |
| Boundary modes can be destroyed by breaking symmetry | YES     | YES                   |
| Boundary modes can be destroyed by disorder          | YES     | NO                    |



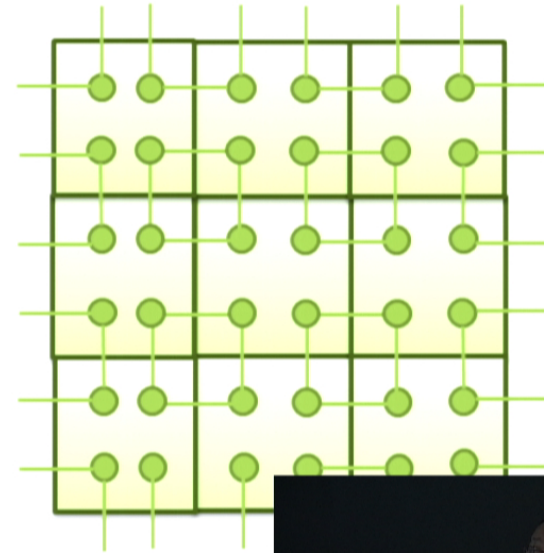
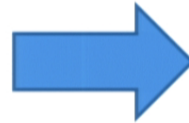
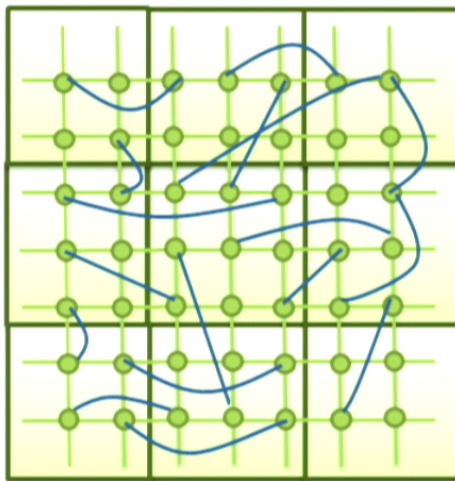
# Generalization to 2D



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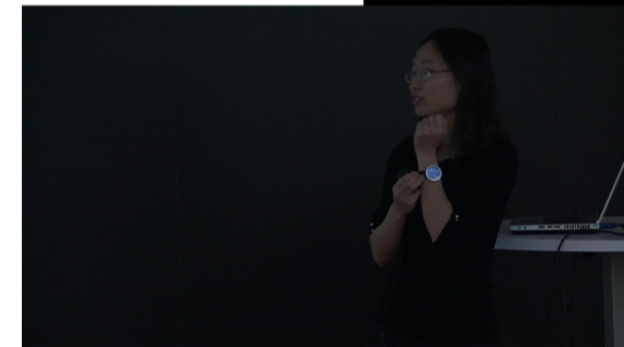
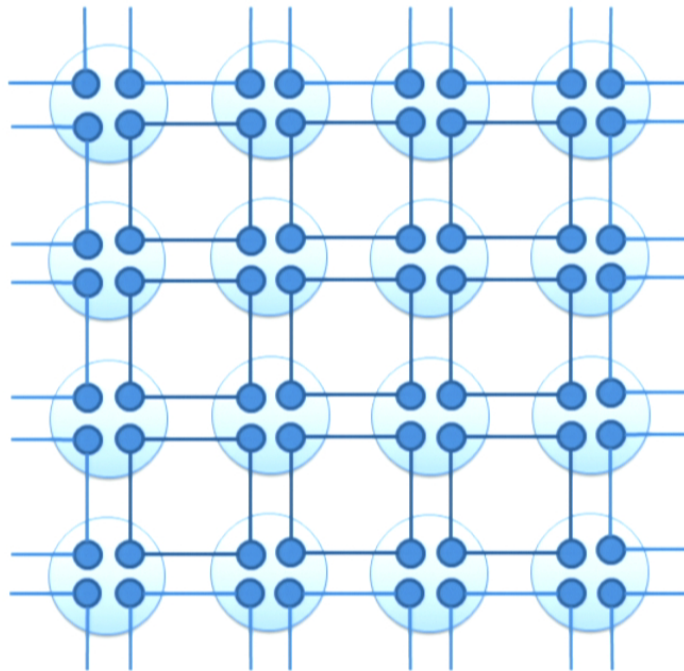


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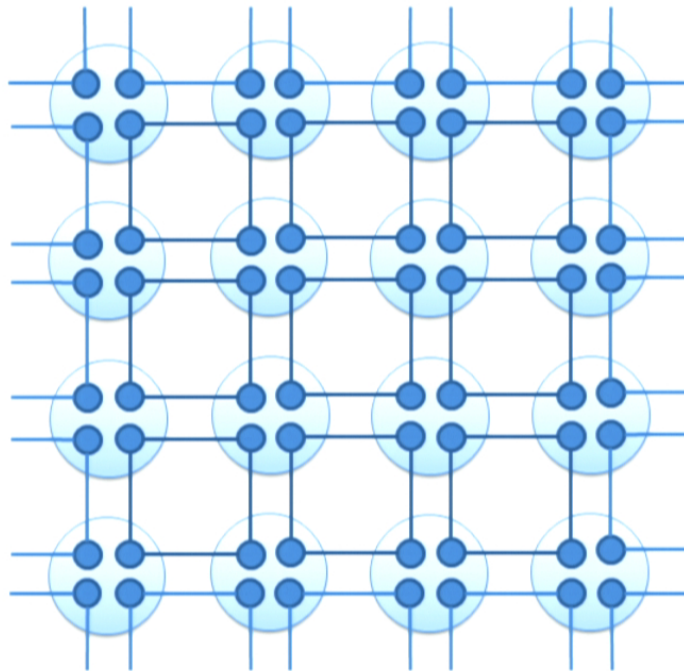
# Generalization to 2D: CZX model

Chen, Liu, Wen, arxiv:1106.4752 ( 2011)

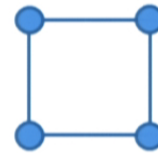


# Generalization to 2D: CZX model

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● 2 level particle  $|0\rangle, |1\rangle$

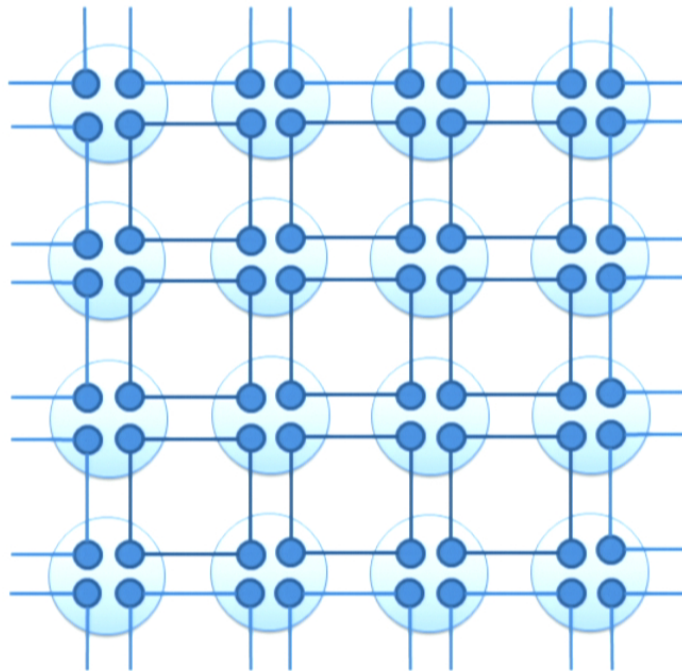


$$\begin{vmatrix} 0 & 0 \\ 0 & 0 \end{vmatrix} + \begin{vmatrix} 1 & 1 \\ 1 & 1 \end{vmatrix}$$



# Generalization to 2D: CZX model

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● 2 level particle  $|0\rangle, |1\rangle$

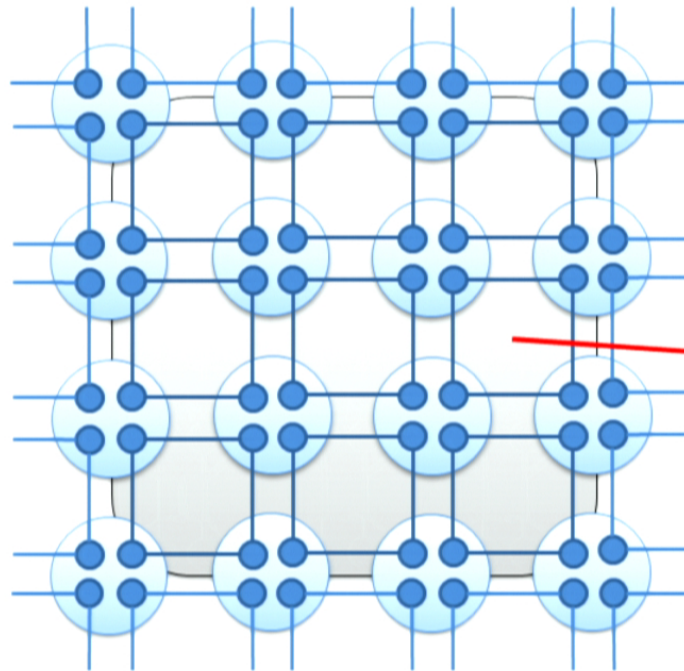
A square diagram with four blue dots at the corners, connected by lines. To its right is the equation:

$$\begin{pmatrix} 0 & 0 \\ 0 & 0 \end{pmatrix} + \begin{pmatrix} 1 & 1 \\ 1 & 1 \end{pmatrix}$$

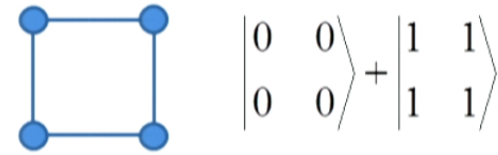
$H = \sum \text{projections onto}$  

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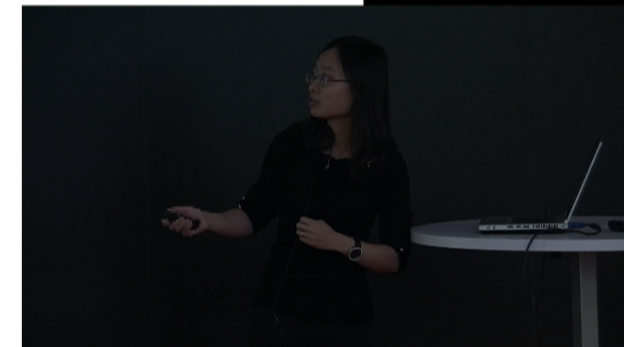
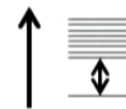
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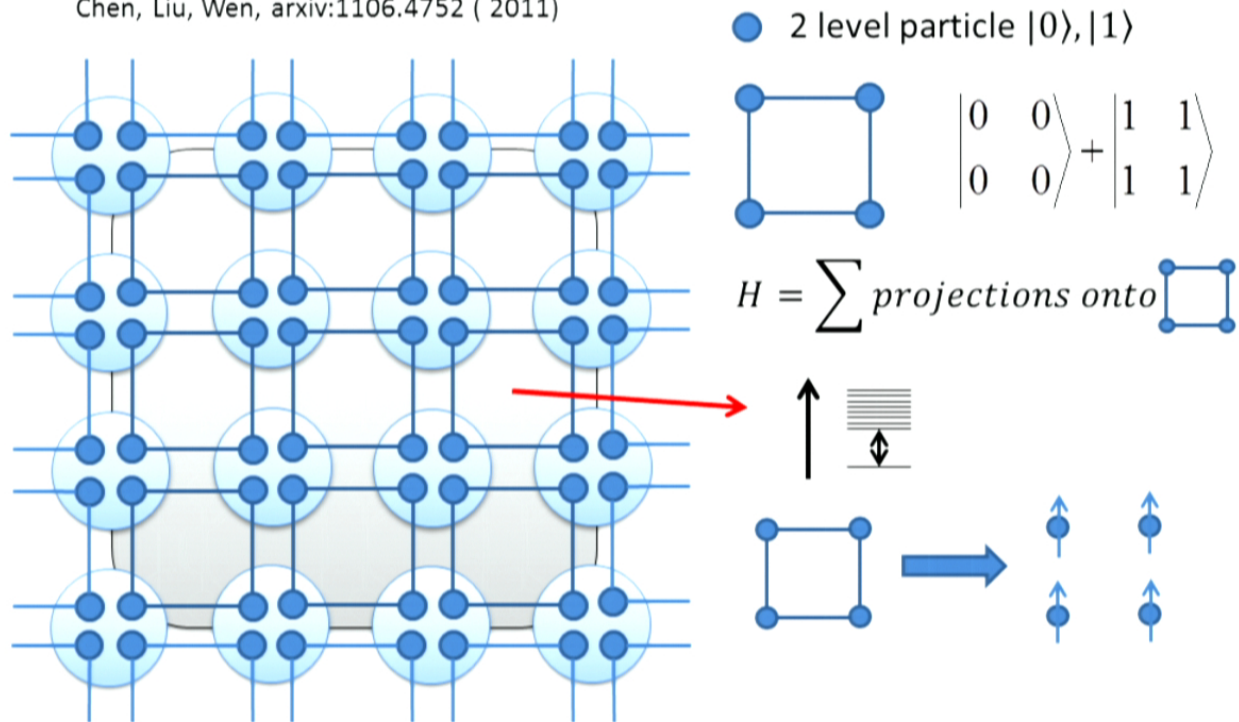


$$H = \sum \text{projections onto } \square$$



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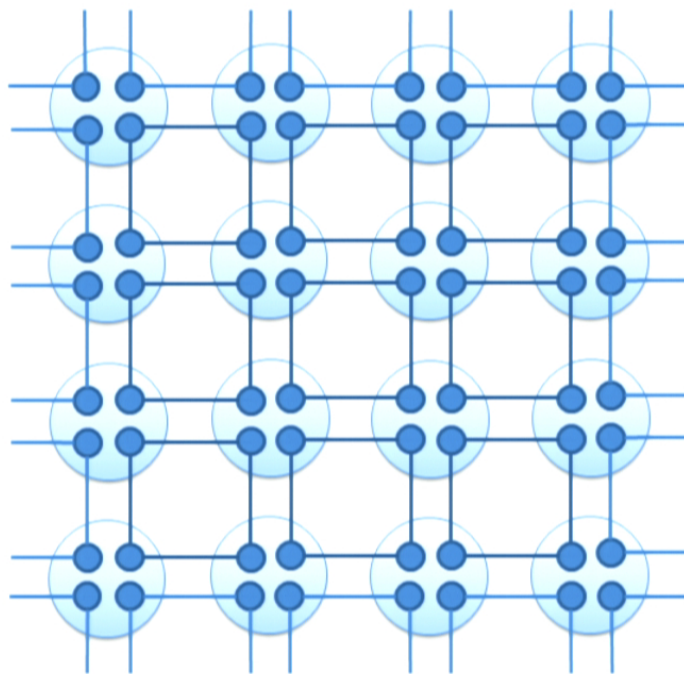
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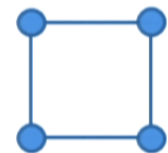


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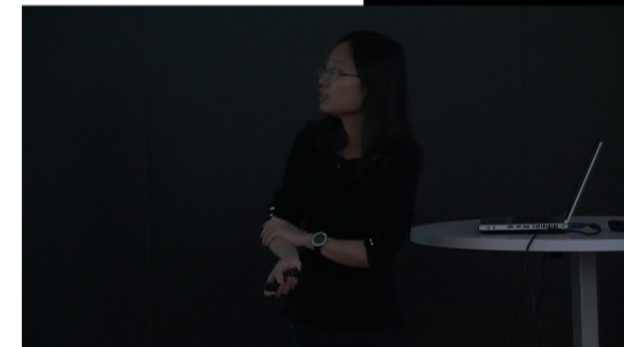


● 2 level particle  $|0\rangle, |1\rangle$


$$\begin{vmatrix} 0 & 0 \\ 0 & 0 \end{vmatrix} + \begin{vmatrix} 1 & 1 \\ 1 & 1 \end{vmatrix}$$

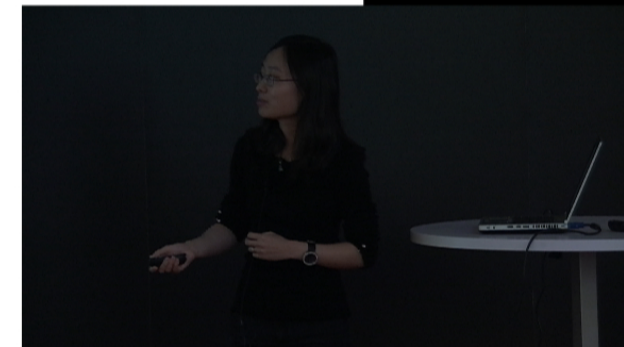
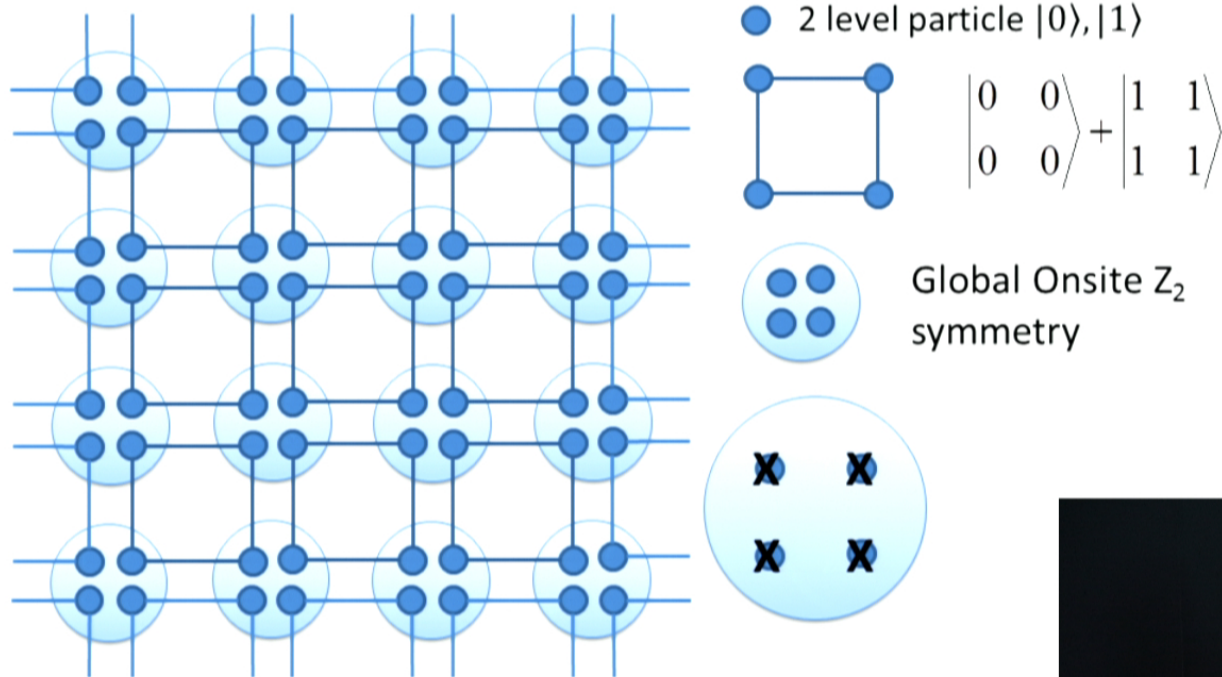


Global Onsite  $Z_2$   
symmetry



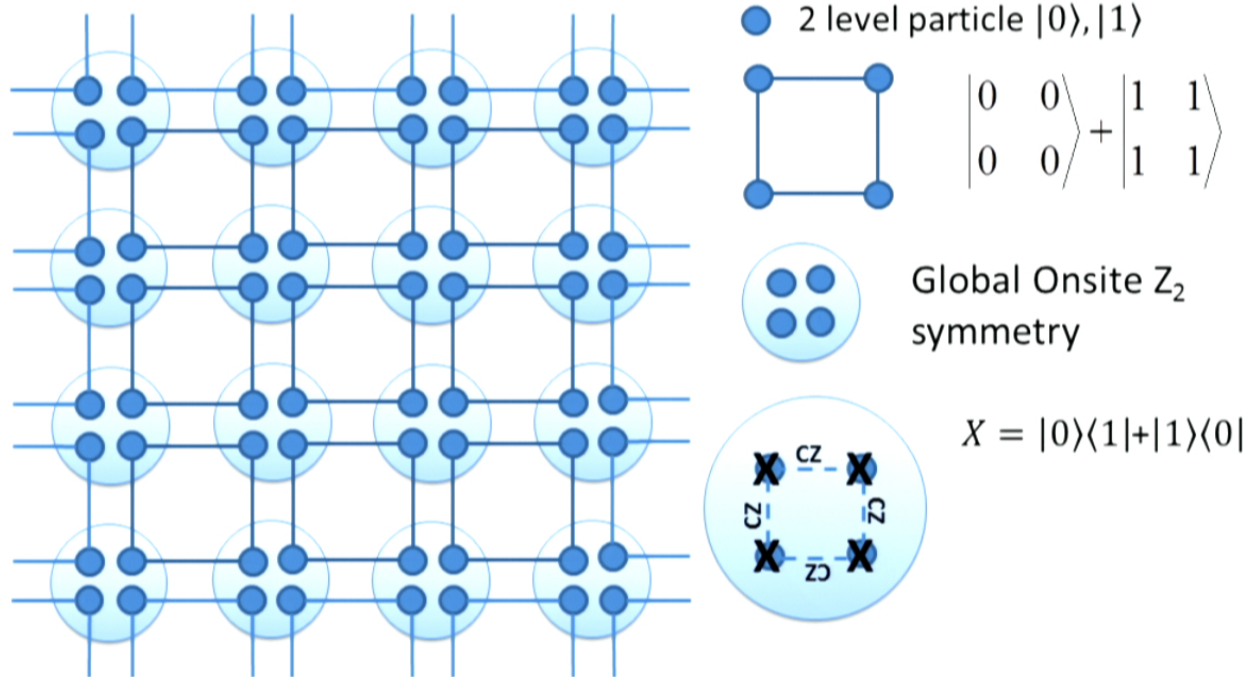
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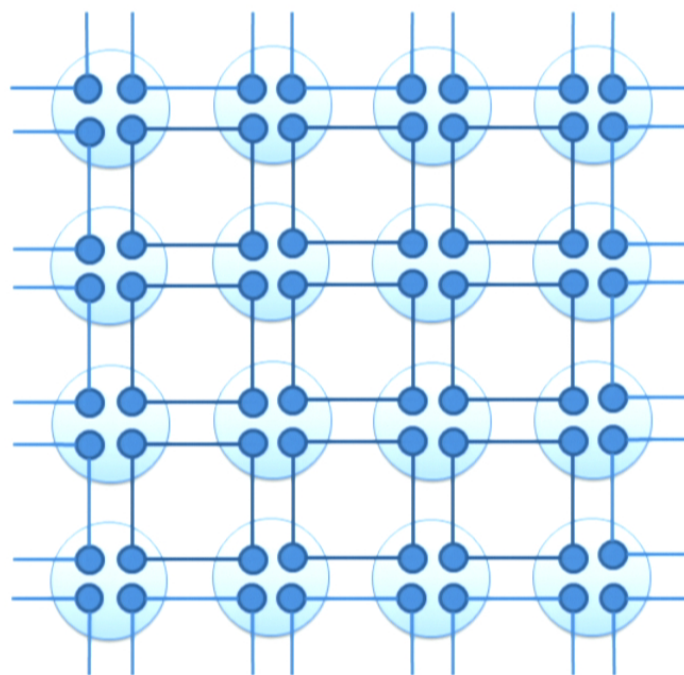
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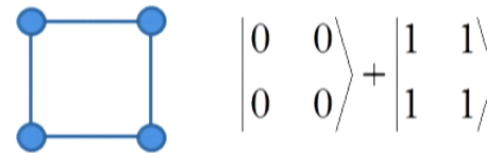


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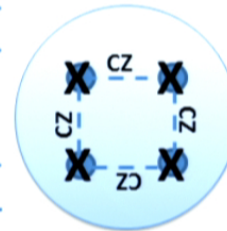
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● 2 level particle  $|0\rangle, |1\rangle$



Global Onsite  $Z_2$   
symmetry

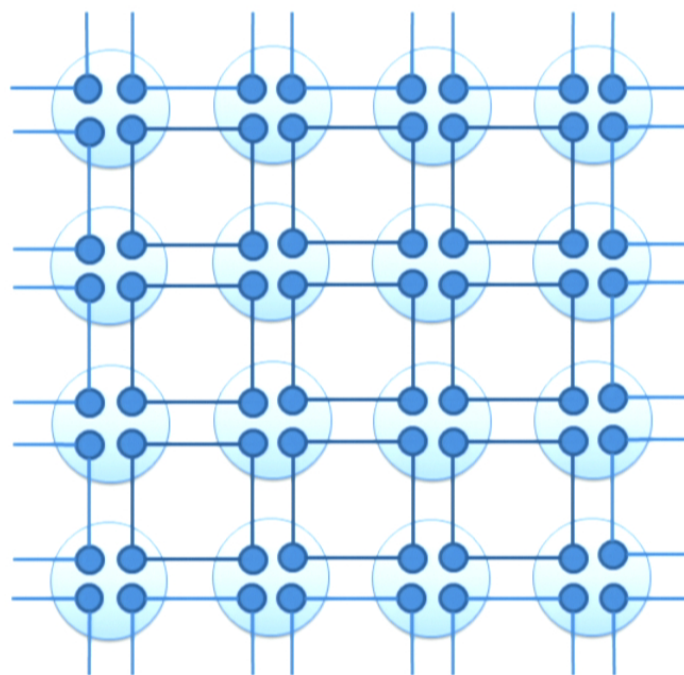


$$X = |0\rangle\langle 1| + |1\rangle\langle 0|$$

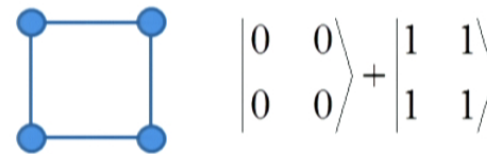
$$CZ = |00\rangle\langle 00| + |01\rangle\langle 01| + |10\rangle\langle 10| - |11\rangle\langle 11|$$

# Generalization to 2D: CZX model

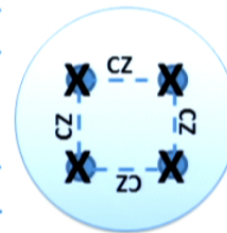
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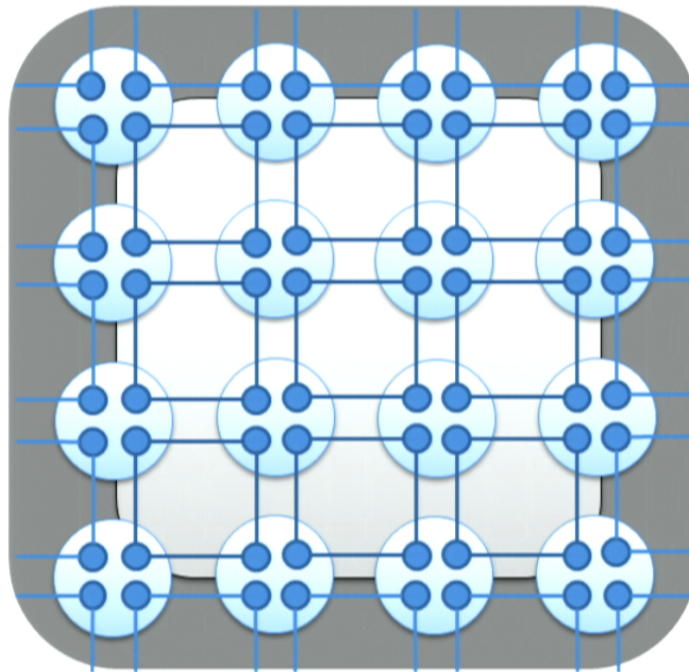


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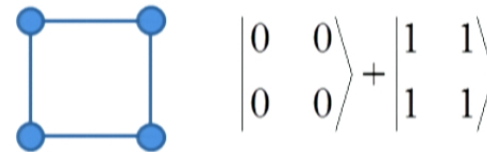
$$CZ = |00\rangle\langle 00| + |01\rangle\langle 01| + |10\rangle\langle 10| - |11\rangle\langle 11|$$

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Chen, Liu, Wen, arxiv:1106.4752, 2011



● 2 level particle  $|0\rangle, |1\rangle$

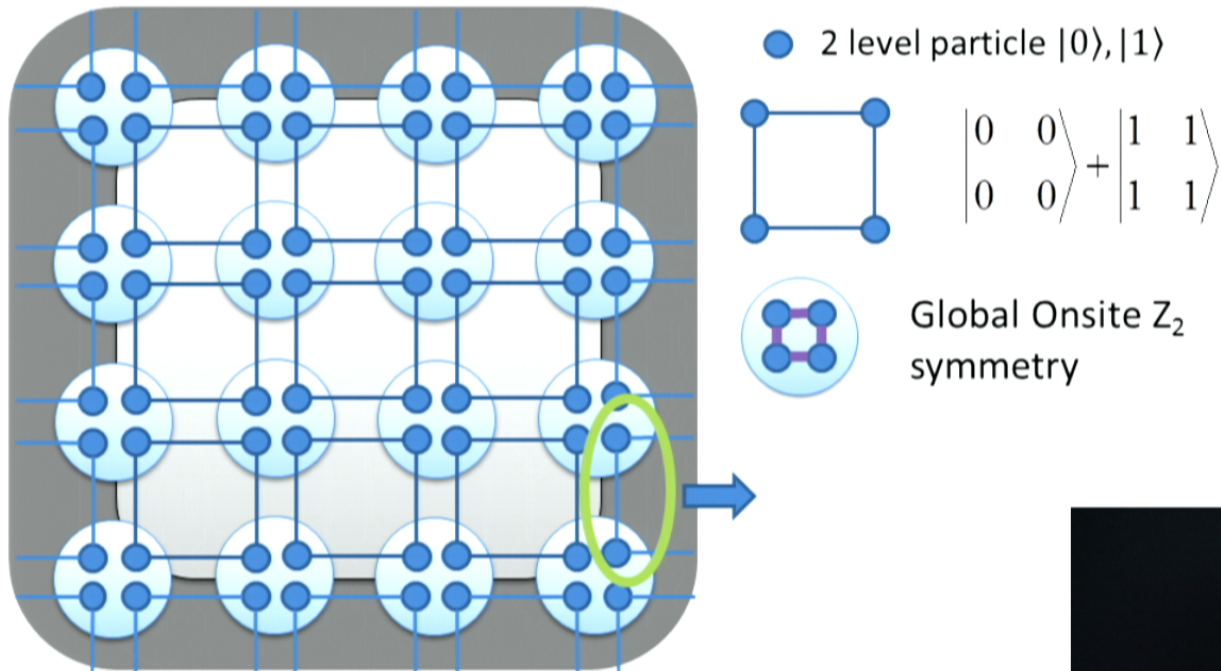


Global Onsite  $Z_2$  symmetry



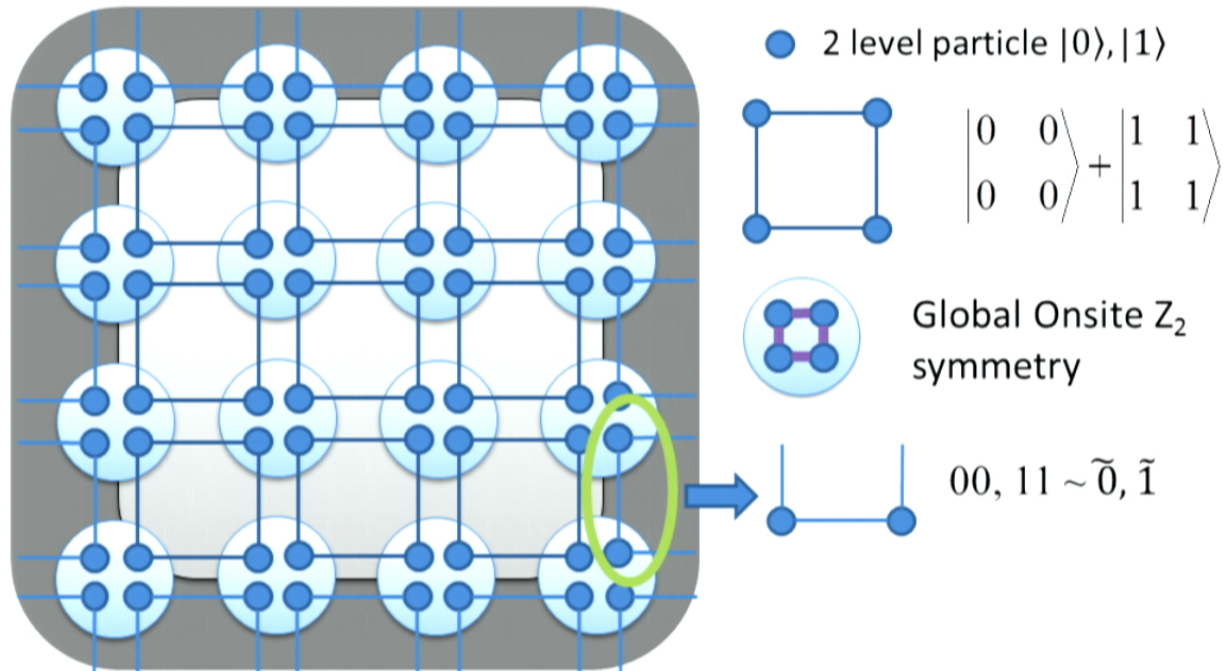
# Generalization to 2D: CZX model

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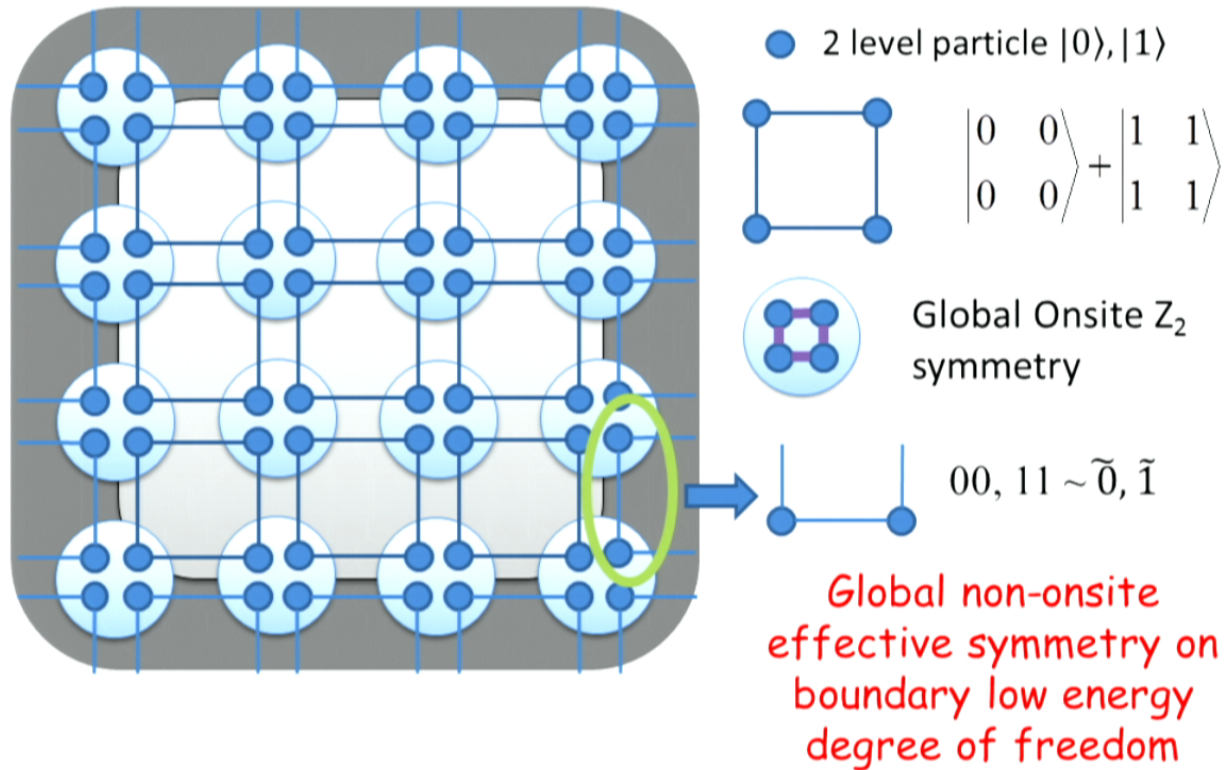
Chen, Liu, Wen, arxiv:1106.4752, 2011





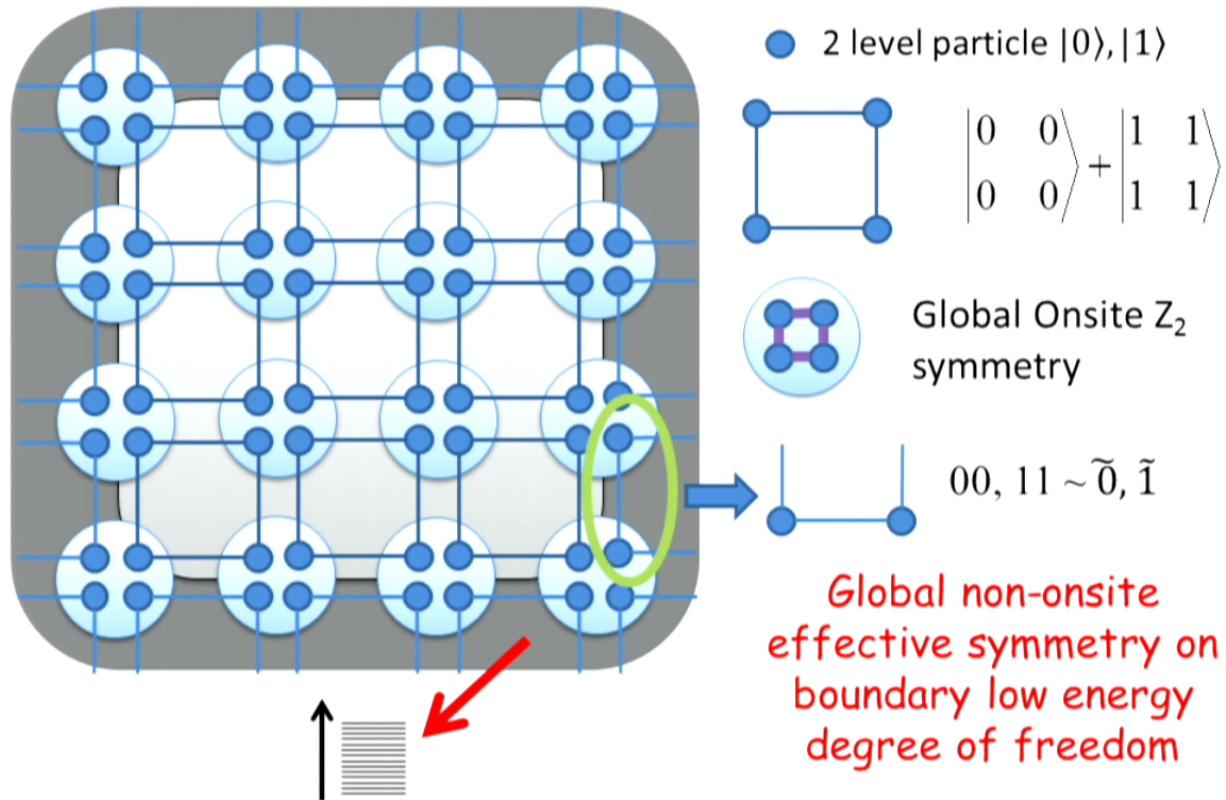
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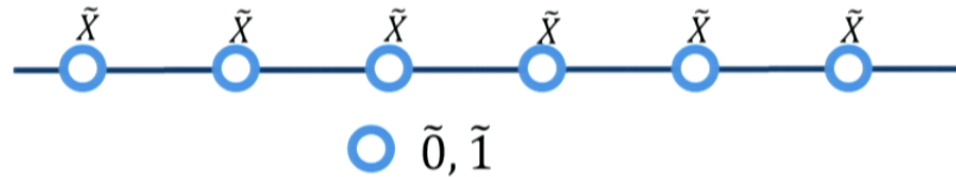
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# Generalization to 2D: CZX model

Boundary of CZX model

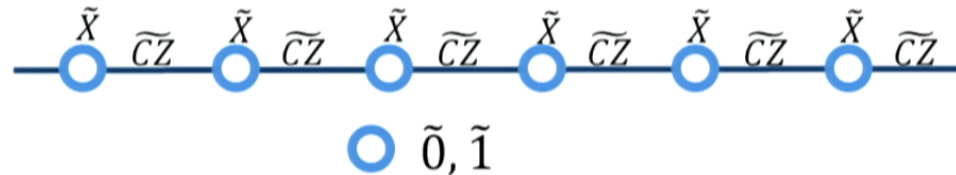


Global **Non-onsite** symmetry on boundary effective degree of freedom



# Generalization to 2D: CZX model

Boundary of CZX model



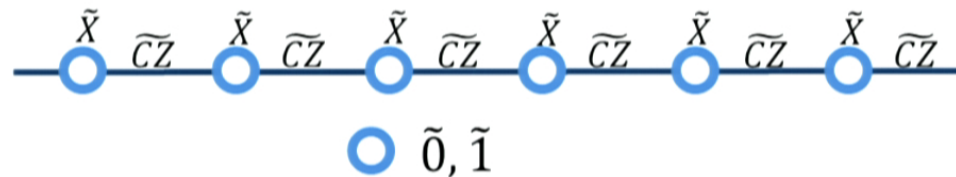
Global **Non-onsite** symmetry on boundary effective degree of freedom

$$H = \sum X_i + Z_{i-1} X_i Z_{i+1}$$



# Generalization to 2D: CZX model

Boundary of CZX model



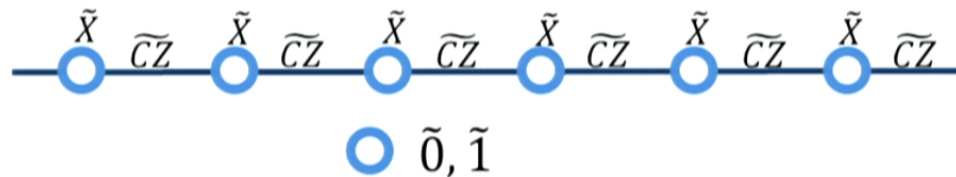
Global **Non-onsite** symmetry on boundary effective degree of freedom

$$H = \sum X_i + Z_{i-1} X_i Z_{i+1} \quad \text{Gapless}$$

$$H = - \sum Z_i Z_{i+1} \quad \text{Symmetry breaking}$$

# Generalization to 2D: CZX model

Boundary of CZX model

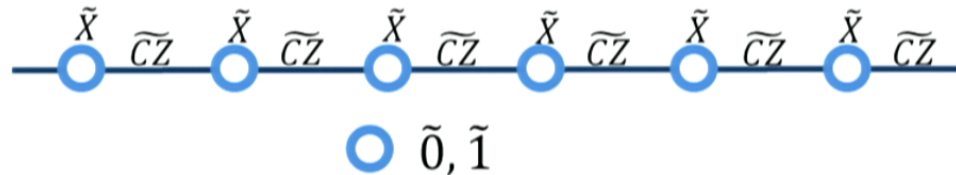


- Global **Non-onsite** symmetry on boundary effective degree of freedom
- Related to **nontrivial group cohomology** of the  $Z_2$  group

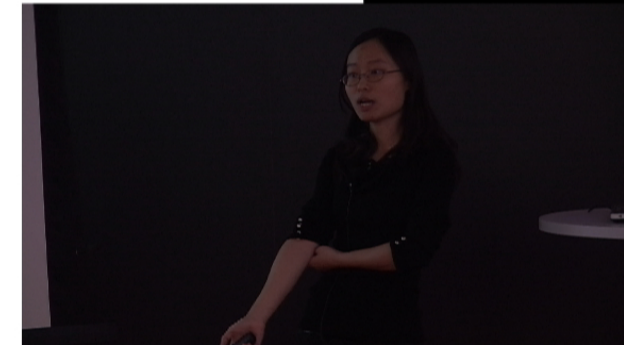


# Generalization to 2D: CZX model

Boundary of CZX model

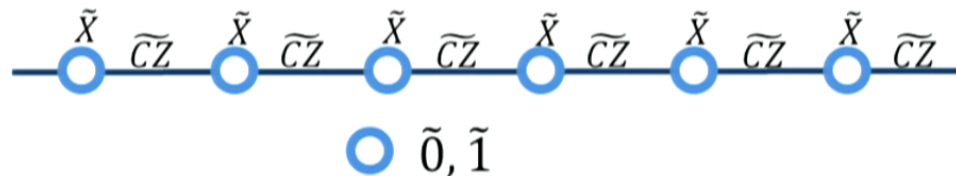


- Global **Non-onsite** symmetry on boundary effective degree of freedom
- Related to **nontrivial group cohomology** of the  $Z_2$  group
- No gapped symmetric phase exist under any interaction



# Generalization to 2D: CZX model

Boundary of CZX model



- Global **Non-onsite** symmetry on boundary effective degree of freedom
- Related to **nontrivial group cohomology** of the  $Z_2$  group
- No gapped symmetric phase exist under any interaction
- Either symmetry breaking or gapless
- Gapless mode stable under disorder and interaction, as long as  $Z_2$  symmetry is not broken



# Comparison with topological insulator

|  | Topological Insulator | CZX |
|--|-----------------------|-----|
| Gap in the bulk                                      | YES                   | YES |
| Gapless mode on boundary                             | YES                   | YES |
| Boundary modes can be destroyed by disorder          | NO                    | NO  |
| Boundary modes can be destroyed by breaking symmetry | YES                   | YES |

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New Symmetry Protected Topological Order in Interacting system

# 2D Symmetry Protected Topological Order

## Previous understanding

- Topological Insulator and topological superconductor
- Free fermion system
- Gapless boundary stable without interaction
- Classified for free fermion system

Kane, Mele, 2005; Xu, Moore, 2008; Kitaev, 2008; Schnyder, Ryu, Furusaki, Ludwig, 2009; Ryu, Schnyder, Furusaki, Ludwig 2009

## Our new model

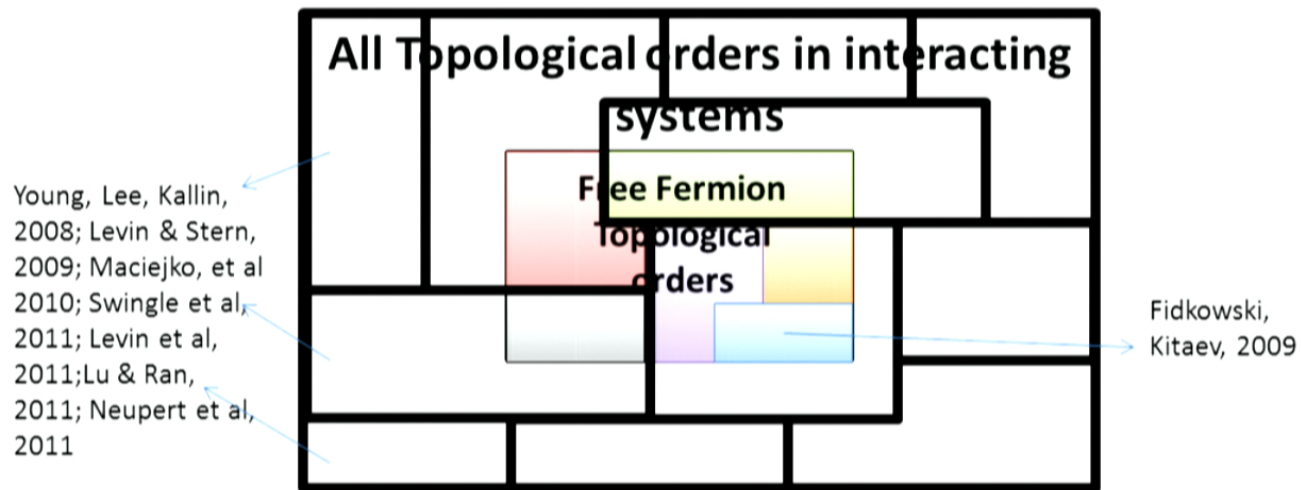
- Strongly interacting boson/fermion system
- Exactly solvable
- Gapless-ness of boundary proven, even with interaction
- Underlying mathematics—group cohomology
- Applies to systems with any symmetry and in any dimension
- Potential classification for all interacting systems

# Outline

- Background and question
  - Topological Order
  - What topological orders exist
- Results
  - Classification of 1D gapped topological order
  - Systematic construction of 2D topological order
- Approach
  - Many-body entanglement
  - Relation to quantum computation
- Outlook

# Hamiltonian Approach

Free fermion + interaction



# Ground State Approach

- Treat all interactions on equal footing
- Zero  $T$ , gap  $\rightarrow$  Ground state: universal structure
- Many-body Entanglement Structure Important!

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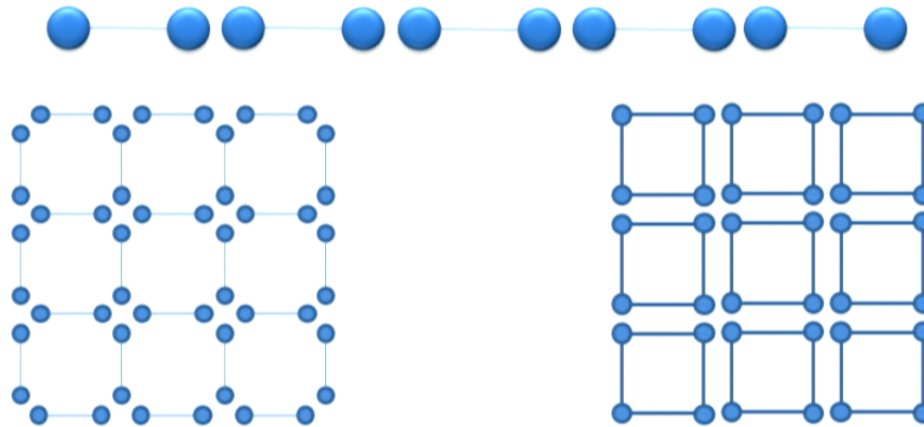
Un-entangled state  
No topological order

- Different many-body entanglement patterns correspond to different topological orders
- Matrix product state, tensor product state

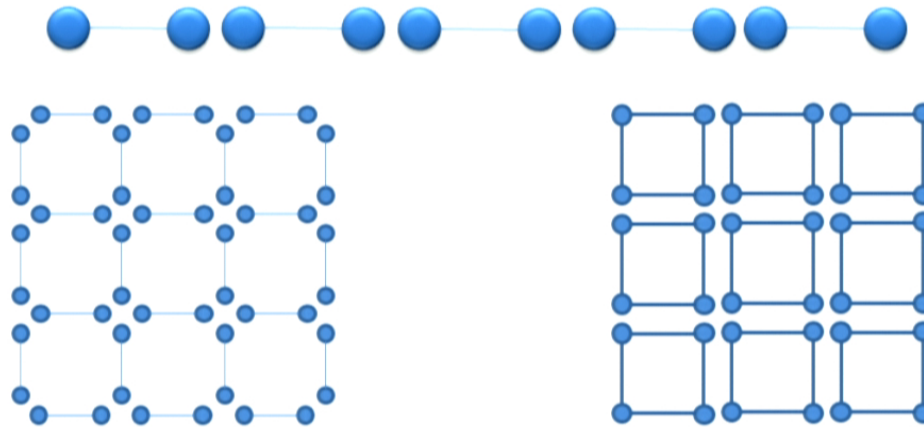
Fannes, Nachergaele, Werner, 1992; Vidal 2007; Perez-Garcia, Verstraete, Wolf, Cirac, 2007;



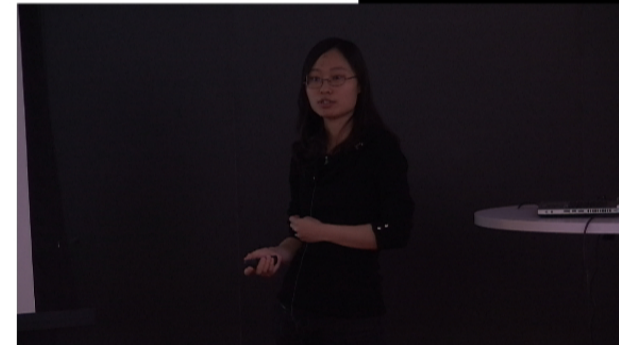
# Short Range Entanglement



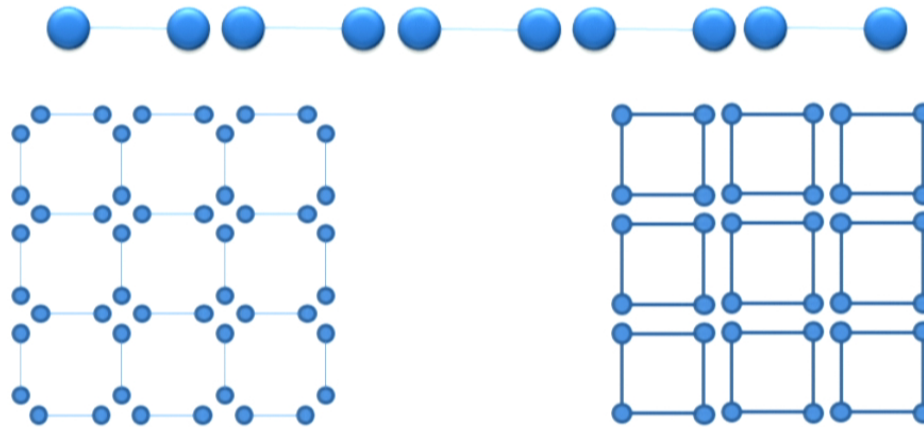
# Short Range Entanglement



- Can be smoothly deformed into a trivial product state without closing the gap of parent Hamiltonian

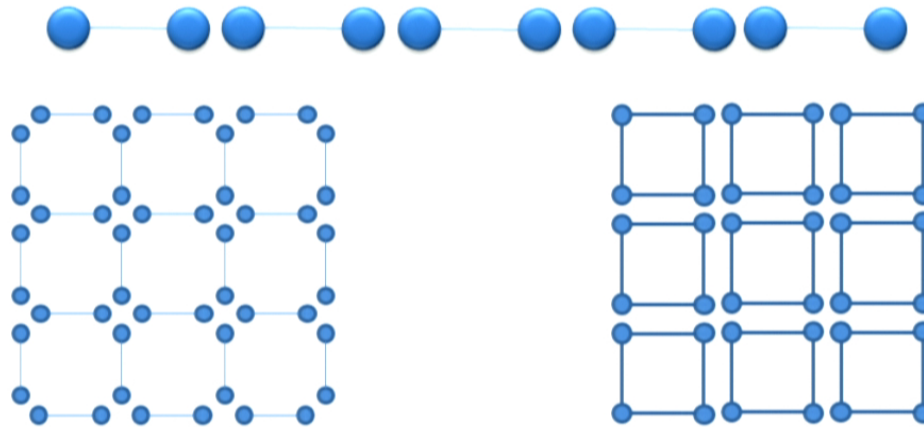


# Short Range Entanglement



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- CANNOT be smoothly deformed into product state under symmetry constraint

# Short Range Entanglement



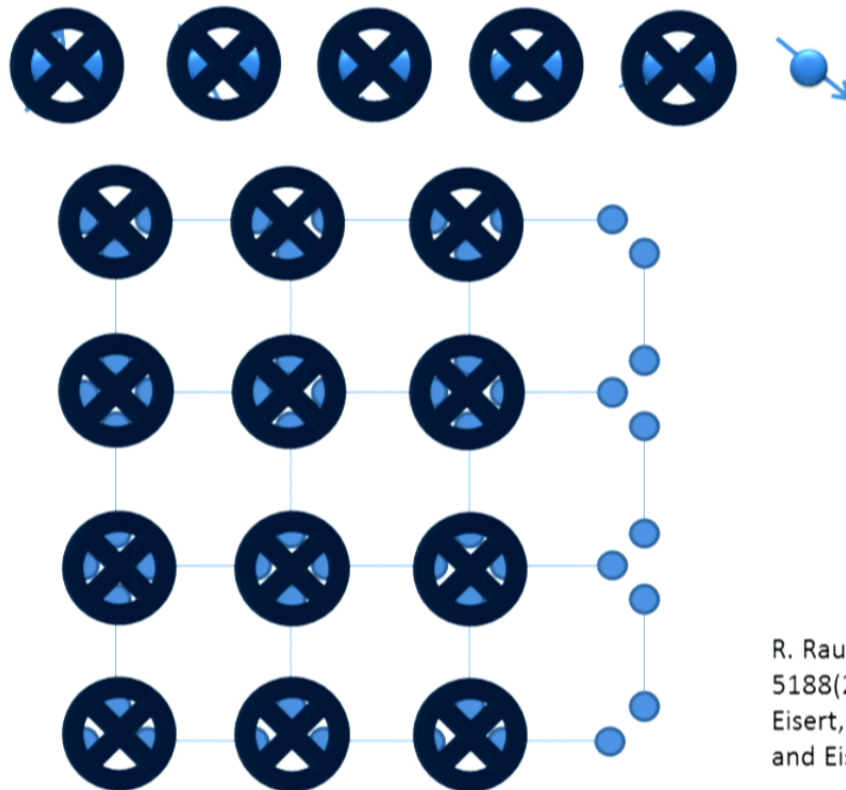
- Can be smoothly deformed into a trivial product state without closing the gap of parent Hamiltonian
- CANNOT be smoothly deformed into product state under symmetry constraint
- Symmetry protected topological order

# SRE and Quantum Computation



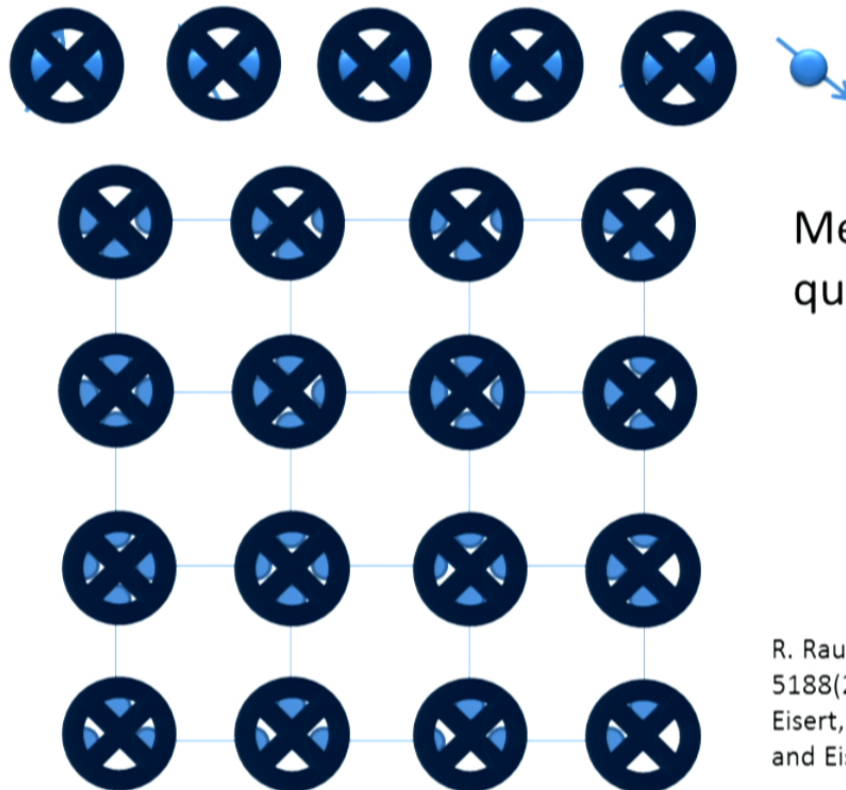
R. Raussendorf and H. Briegel, PRL 86, 5188(2001); Verstrate, Cirac, 2004; Gross, Eisert, Schuch, Perez-Garcia, 2007; Gross and Eisert, 2007

# SRE and Quantum Computation



R. Raussendorf and H. Briegel, PRL 86, 5188(2001); Verstrate, Cirac, 2004; Gross, Eisert, Schuch, Perez-Garcia, 2007; Gross and Eisert, 2007

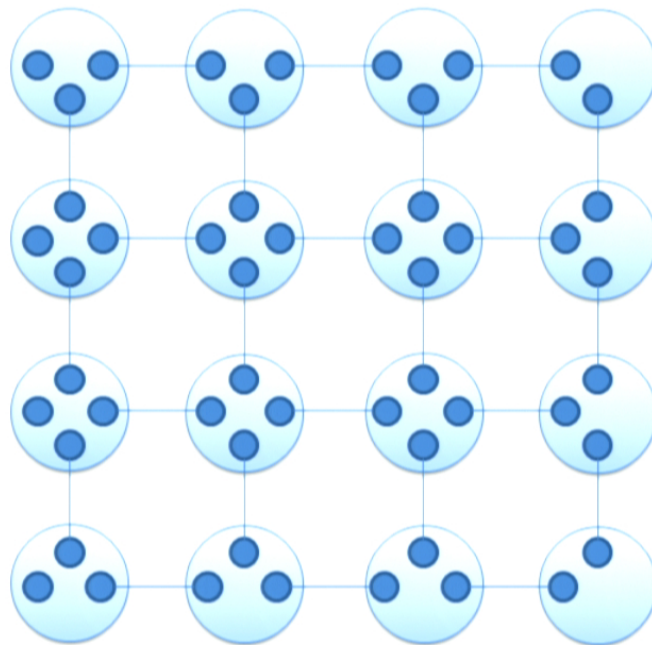
# SRE and Quantum Computation



Measurement-based  
quantum computation

R. Raussendorf and H. Briegel, PRL 86,  
5188(2001); Verstrate, Cirac, 2004; Gross,  
Eisert, Schuch, Perez-Garcia, 2007; Gross  
and Eisert, 2007

# SRE and measurement-based quantum computation





# LRE and quantum computation

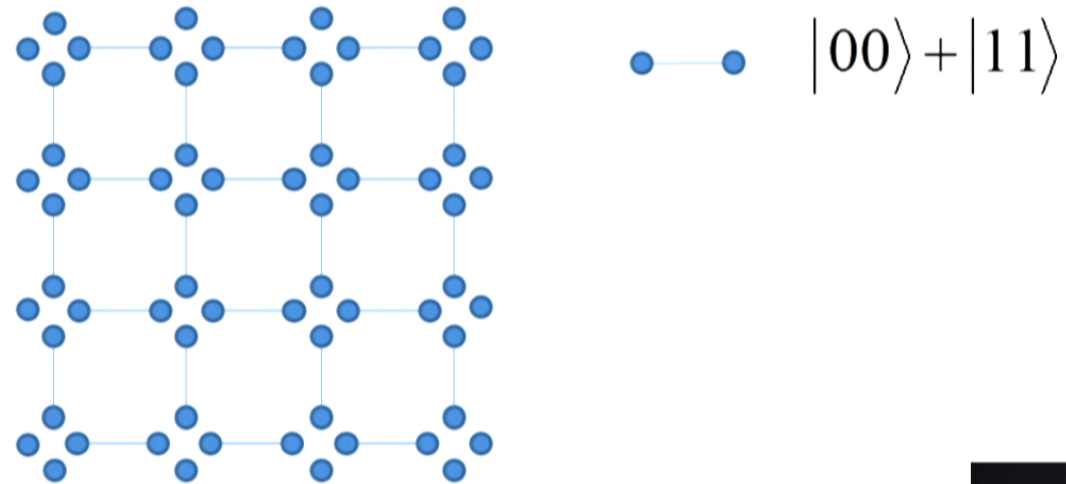
- Features:
  - Degeneracy depending on topology of system
  - fractional charge, fractional statistics

# LRE and quantum computation

- Features:
  - Degeneracy depending on topology of system
  - fractional charge, fractional statistics
- Quantum information application:
  - Topological quantum error correction code
  - Topological quantum memory
  - Topological quantum computation

# Long range entanglement

## Toric code



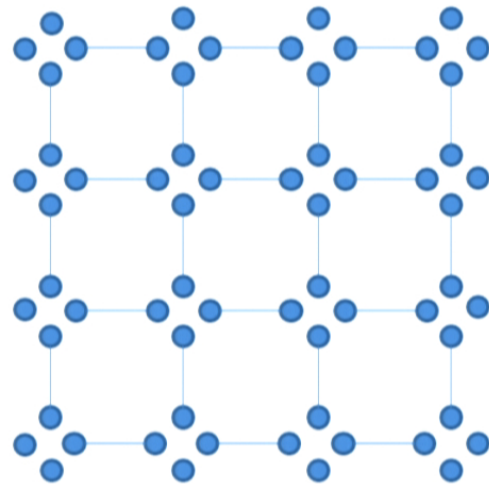
Bruerschaper, Aguado, Vidal, 2009; Gu, Levin, Swingle, Wen, 2009; Brian, Wen, 2010; Chen, Wen, 2010; Schuch, Cirac, Perez-Garcia, 2010; Verstraete, Wolf, Perez-Garcia, Cirac, 2006;




# Long range entanglement

Toric code

Entangled pair + gauge symmetry



  $|00\rangle + |11\rangle$

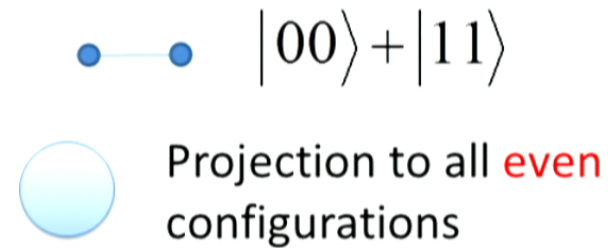
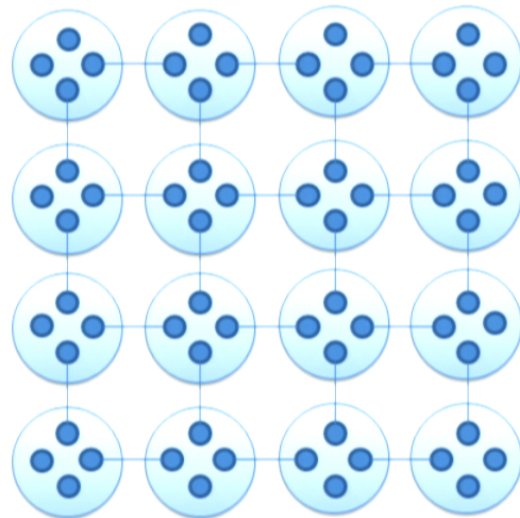
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# Long range entanglement

Toric code

Entangled pair + gauge symmetry



$$\begin{array}{cccc}
 \begin{array}{c} 0 \\ 0 \end{array} & \begin{array}{c} 0 \\ 1 \end{array} & \begin{array}{c} 1 \\ 1 \end{array} & \begin{array}{c} 1 \\ 0 \end{array} \\
 \begin{array}{c} 0 \\ 0 \end{array} & \begin{array}{c} 1 \\ 1 \end{array} & \begin{array}{c} 0 \\ 0 \end{array} & \begin{array}{c} 0 \\ 1 \end{array} \\
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 \end{array}$$

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Topological Order



Quantum Information

Many-body  
Entanglement

