

Title: Formulating Quantum Theory as a Causally Neutral Theory of Bayesian Inference

Date: Nov 08, 2011 03:30 PM

URL: <http://www.pirsa.org/11110114>

Abstract: Quantum theory can be thought of as a noncommutative generalization of Bayesian probability theory, but for the analogy to be convincing, it should be possible to describe inferences among quantum systems in a manner that is independent of the causal relationship between those systems. In particular, it should be possible to unify the treatment of two kinds of inferences: (i) from beliefs about one system to beliefs about another, for instance, in the Einstein-Podolsky-Rosen or “quantum steering” phenomenon, and (ii) from beliefs about a system at one time to beliefs about that same system at another time, for instance, in predictions or retrodictions about a system undergoing dynamical evolution or undergoing a measurement. I will present a formalism that achieves such a unification by making use of “conditional quantum states”, a noncommutative generalization of conditional probabilities. I argue for causal neutrality by drawing a comparison with a classical statistical theory with an epistemic restriction. (Joint work with Matthew Leifer).

Where I'm coming from...

$b_k^+$

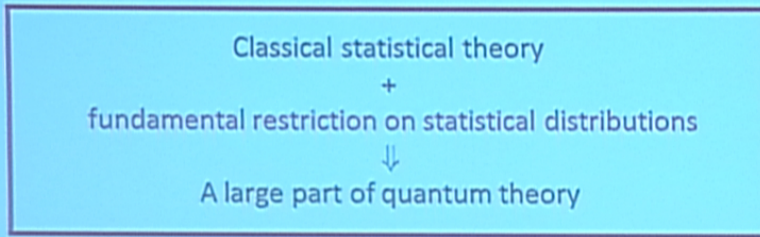
Canonical Transform  
 $[b_k, b_{k'}^+] = \delta_{kk'}$

choose  $u_k^2 - v_k^2 = 1 \quad \forall k \neq 0$   
But  $\frac{u_k}{v_k}$  is still arbitrary

$a_k^+$   $u_k b_k^+$   
trunk  
particle







In the sense of reproducing the operational predictions

$b_k^+$

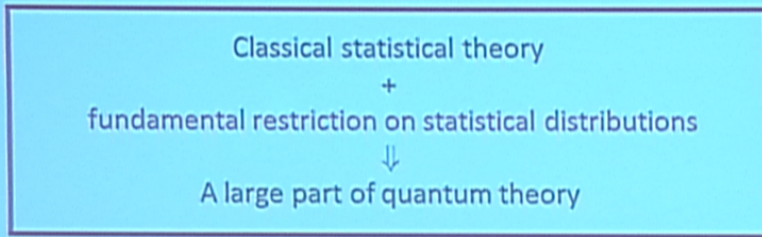
Canonical Transform  
 $[b_k, b_k^+] = \delta_{kk'}$

choose  $u_k = v_k = 1 \quad \forall k \neq 0$   
But  $\frac{u_0}{v_0}$  is still arbitrary

$a_k^+ = u_k b_k^+$   
truth  
particle  $\rightarrow$

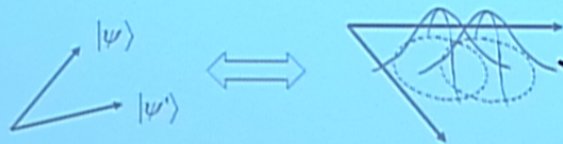






In the sense of reproducing the operational predictions

In the resulting model  
 quantum states are states of incomplete knowledge



Canonical Transform  
 $[b_k, b_k^\dagger] = \delta_{kk'}$   
 choose  $u_k^2 - v_k^2 = 1 \quad \forall k \neq 0$   
 But  $\frac{u_k}{v_k}$  is still arbitrary

$a_k^\dagger = u_k b_k^\dagger + v_k b_k$   
 track particles



The principle of classical complementarity:

An observer can only jointly know a set of variables if they commute relative to the Poisson bracket.

$b_k^+$

Canonical Transform  
 $[b_k^+, b_l^+] = \delta_{kl}$

choose  $u_k^2 - v_k^2 = 1 \quad \forall k \neq 0$   
But  $\frac{u_k}{v_k}$  is still arbitrary

$a_k^+ = u_k b_k^+ + v_k b_{-k}^+$   
frank  
particle



The principle of classical complementarity:

An observer can only jointly know a set of variables if they commute relative to the Poisson bracket.

A commuting pair  $\{F, G\} = 0$

e.g.  $\{Q_A, Q_B\}$ ,  $\{Q_A, P_B\}$ , and  $\{Q_A - Q_B, P_A + P_B\}$

$D_k^+$

Canonical Transform  
 $[b_k, b_k^\dagger] = \delta_{kk}$

choose  $u_k^2 - v_k^2 = 1 \quad \forall k \neq 0$   
But  $\frac{u_k}{v_k}$  is still arbitrary

$a_k^\dagger = u_k b_k^\dagger + v_k b_k$   
truth  
particle



The principle of classical complementarity:

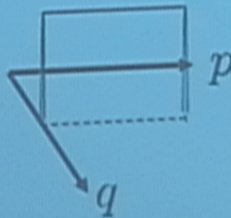
An observer can only jointly know a set of variables if they commute relative to the Poisson bracket.

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e.g.  $\{Q_A, Q_B\}$ ,  $\{Q_A, P_B\}$ , and  $\{Q_A - Q_B, P_A + P_B\}$

Valid preparations

$$P(q, p) \propto \delta(q - a)$$





The principle of classical complementarity:

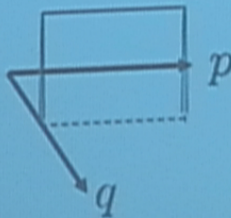
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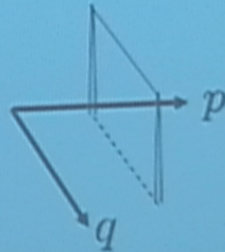
e.g.  $\{Q_A, Q_B\}$ ,  $\{Q_A, P_B\}$ , and  $\{Q_A - Q_B, P_A + P_B\}$

Valid preparations

$$P(q, p) \propto \delta(q - a)$$



$$P(q, p) \propto \delta(p - b)$$

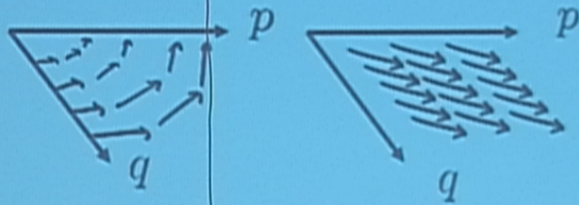


$$P(q_A, p_A, q_B, p_B) \propto \delta(q_A - q_B) \delta(p_A + p_B)$$

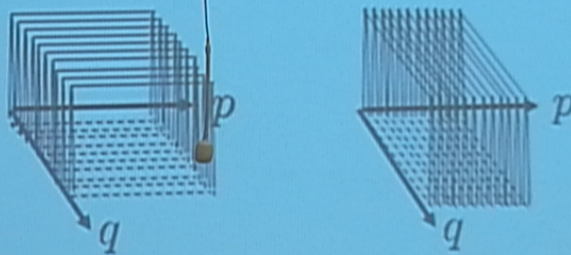
Analogue of EPR



## Valid transformations



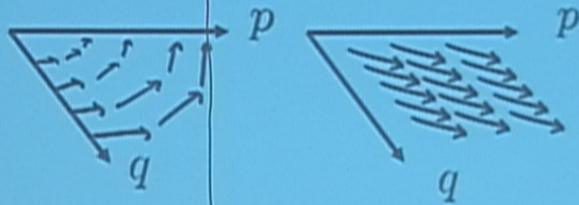
## Valid measurements



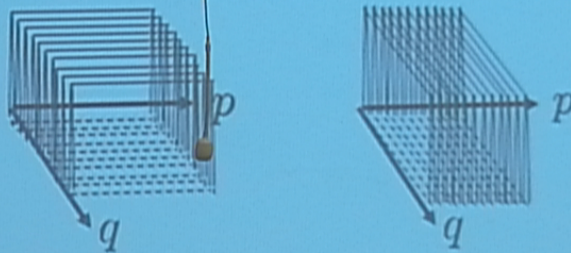
**Theorem:** Liouville mechanics with this statistical restriction is empirically equivalent to "quadrature quantum mechanics"



## Valid transformations



## Valid measurements



**Theorem:** Liouville mechanics with this statistical restriction is empirically equivalent to “quadrature quantum mechanics”



The principle of classical complementarity:

An observer can only jointly know a set of variables if they commute relative to the Poisson bracket.

A commuting pair  $\{F, G\} = 0$

e.g.  $\{Q_A, Q_B\}$ ,  $\{Q_A, P_B\}$ , and  $\{Q_A - Q_B, P_A + P_B\}$

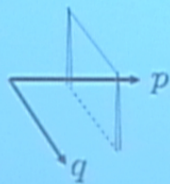
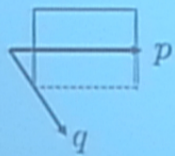
Valid preparations

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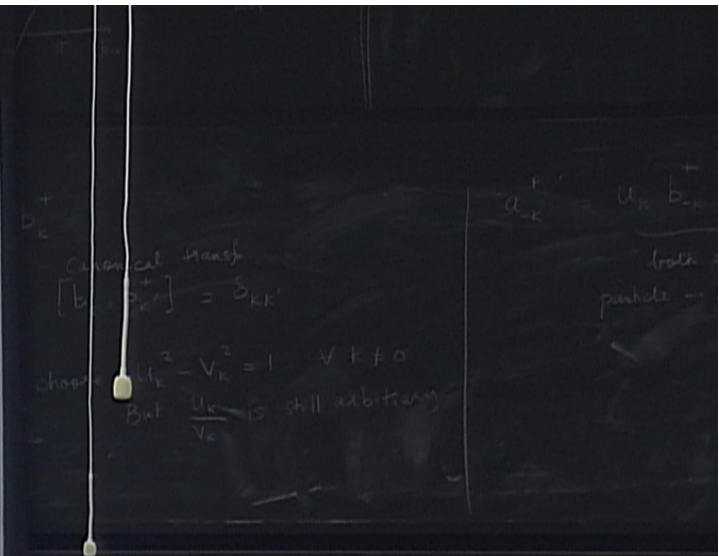
$$P(q, p) \propto \delta(p - b)$$

$$P(q_A, p_A, q_B, p_B)$$

$$\propto \delta(q_A - q_B) \delta(p_A + p_B)$$

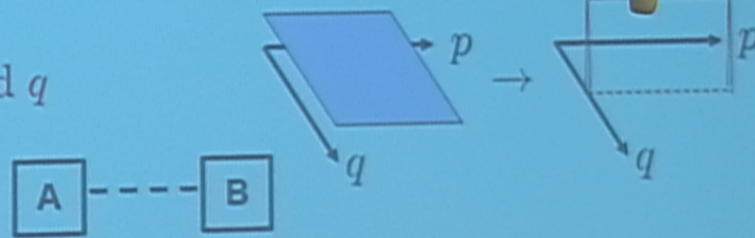


Analogue of EPR



# Model of EPR in epistemically restricted Liouville mechanics

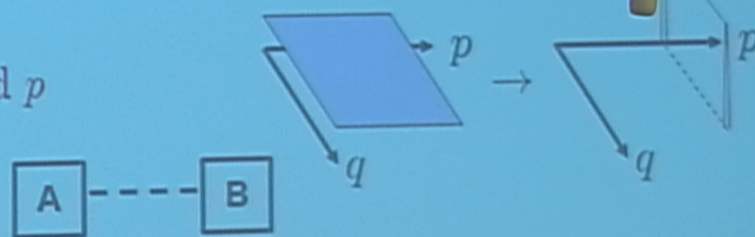
Measure  $Q_A$  find  $q$



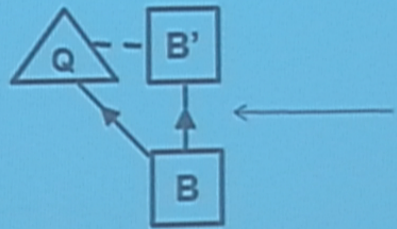


# Model of EPR in epistemically restricted Liouville mechanics

Measure  $P_A$  find  $p$



# Update of epistemic state as a result of measurement



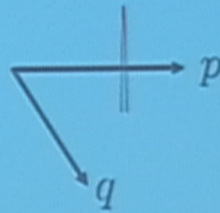
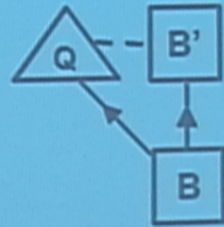
A causal connection  
e.g. B' and B denote the  
same system at two  
different times

Canonical  
 $[b_e, b_e^*]$   
...  
B...

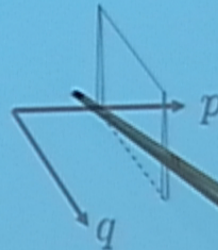


# Update of epistemic state as a result of measurement

Measure  $Q_B$  find  $q$

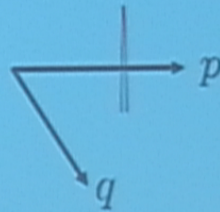
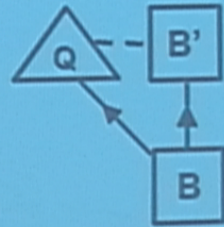


But this would violate the statistical restriction!

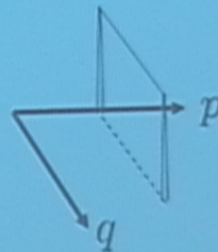


# Update of epistemic state as a result of measurement

Measure  $Q_B$  find  $q$



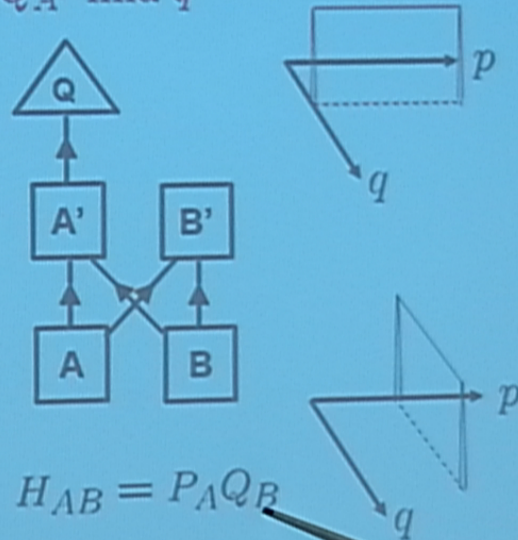
But this would violate the statistical restriction!





## Update of epistemic state as a result of measurement

Measure  $Q_{A'}$ , find  $q$



$$H_{AB} = P_A Q_B$$

$$Q_{A'} = Q_A + \gamma Q_B$$

$$P_{A'} = P_A$$

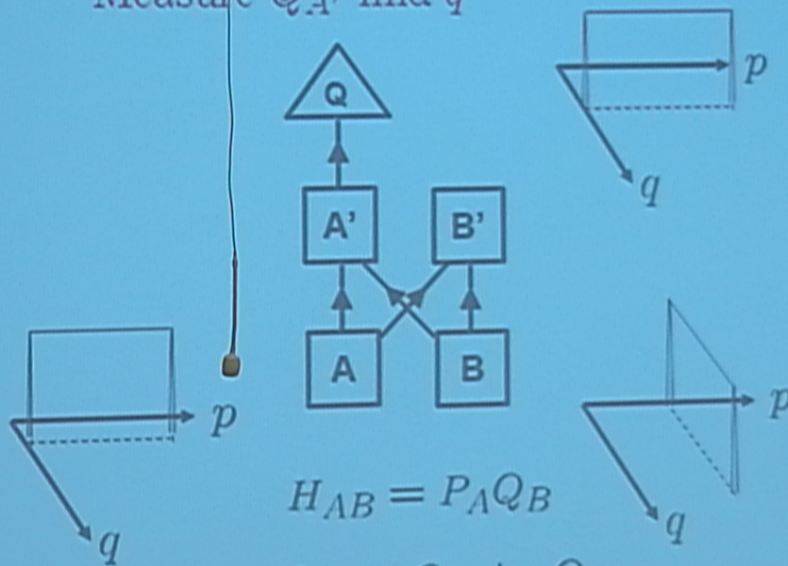
$$Q_{B'} = Q_B$$

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## Update of epistemic state as a result of measurement

Measure  $Q_{A'}$ , find  $q$



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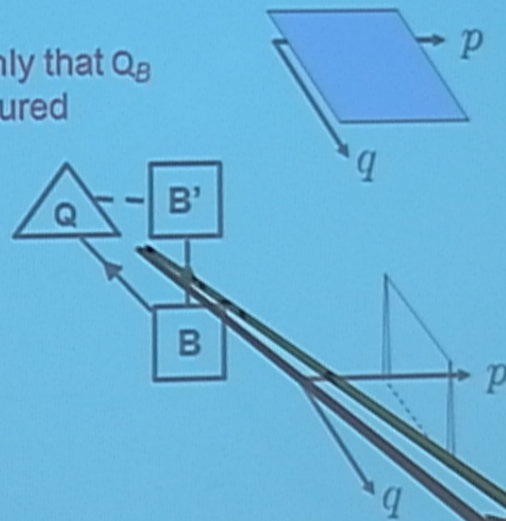
$$Q_{B'} = Q_B$$

$$P_{B'} = P_B + \gamma P_A$$



# Update of epistemic state as a result of measurement

Agent knows only that  $Q_B$  was measured



## Categorizing quantum phenomena

Those arising in a restricted  
statistical classical theory

Those not arising in a restricted  
statistical classical theory

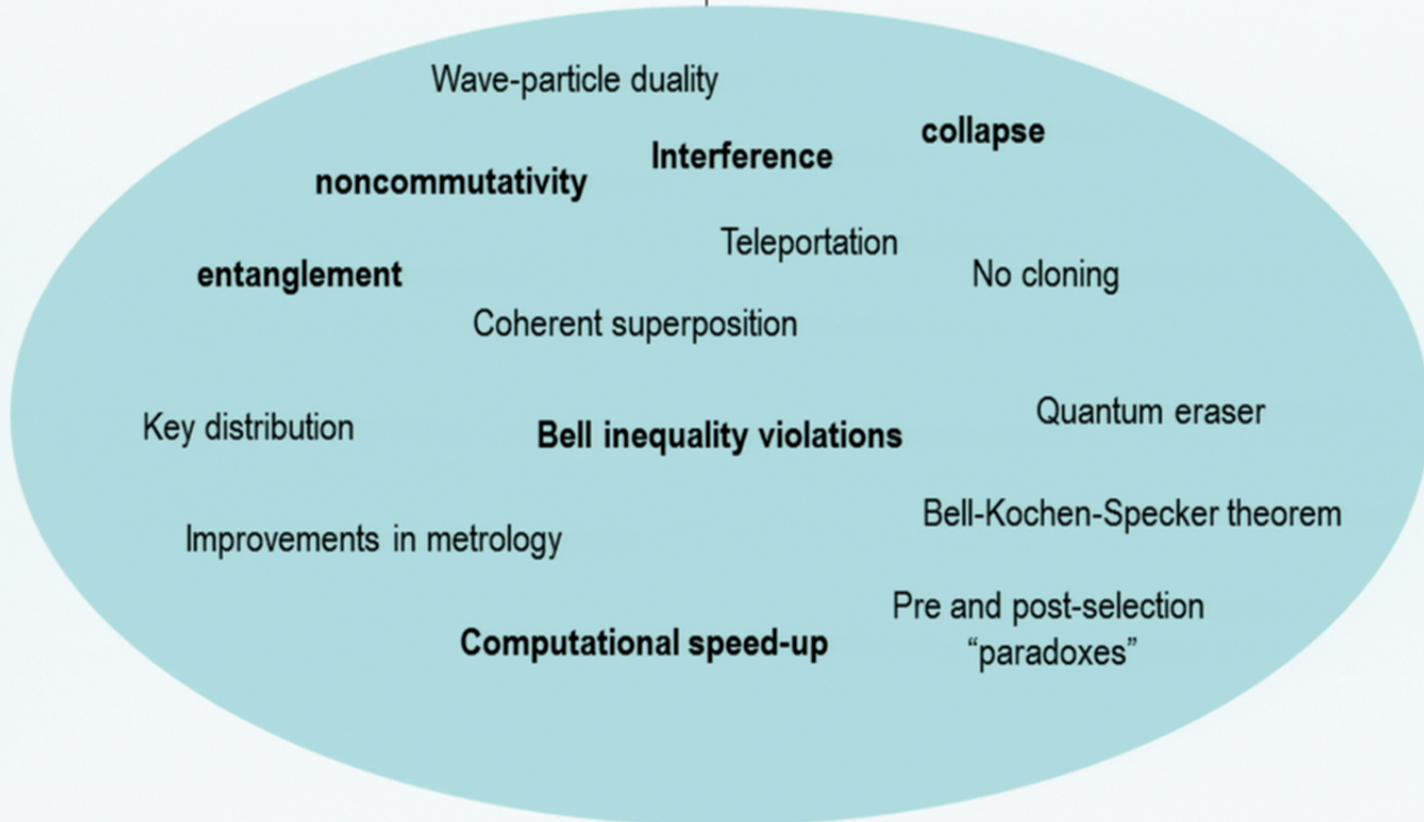
Canonical  
 $[b_c, b_c^\dagger]$   
...  
But



# Categorizing quantum phenomena

Those arising in a restricted statistical classical theory

Those not arising in a restricted statistical classical theory



## Categorizing quantum phenomena

### Those arising in a restricted statistical classical theory

**Interference**  
**Noncommutativity**  
**Entanglement**  
**Collapse**  
Wave-particle duality  
Teleportation  
No cloning  
Key distribution  
Improvements in metrology  
Quantum eraser  
Coherent superposition  
Pre and post-selection “paradoxes”  
Others...

### Those not arising in a restricted statistical classical theory

**Bell inequality violations**  
**Computational speed-up (if it exists)**  
Bell-Kochen-Specker theorem  
Certain aspects of items on the left  
Others...



## Classical

$R$  phase-space coordinates  
of a canonical system

$P(R)$  Probability  
distribution over  $R$

$P(R = r)$  probability that  $R=r$

$$\sum_R P(R) = 1$$

## Quantum

$A$  label for a quantum system

$\rho_A$  Operator on Hilbert  
space of  $A$

No obvious analogue (yet!)

$$\text{Tr}_A \rho_A = 1$$

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	Classical	Quantum
State of knowledge	$P(R)$	$\rho_A$
Normalization	$\sum_R P(R) = 1$	$\text{Tr}_A \rho_A = 1$
Joint state	$P(R, S)$	$\rho_{AB}$
Marginalization	$P(S) = \sum_R P(R, S)$	$\rho_B = \text{Tr}_A \rho_{AB}$

# The Three Pillars of Bayesian Inference

- Belief propagation

$$P(S) = \sum_R P(S|R)P(R)$$

- Bayes' theorem

$$P(R|S) = \frac{P(S|R)P(R)}{\sum_R P(S|R)P(R)}$$

- Bayesian conditioning

$$P(R) \rightarrow P(R|X = x)$$



Conditional probability

$$P(S|R)$$

Conditional state

$$\rho_{B|A}$$

Conditional probability

$$P(S|R)$$

Normalization condition

$$\sum_S P(S|R) = 1$$

Conditional state

$$\rho_{B|A}$$

Normalization condition

$$\text{Tr}_B(\rho_{B|A}) = I_A$$



Conditional probability

$$P(S|R)$$

Normalization condition

$$\sum_S P(S|R) = 1$$

Conditional state

$$\rho_{B|A}$$

Normalization condition

$$\text{Tr}_B(\rho_{B|A}) = I_A$$

See: Leifer, PRA 74, 042310 (2006)

Conditional probability

$$P(S|R)$$

Normalization condition

$$\sum_S P(S|R) = 1$$

Relation of conditional to joint

$$P(S|R) = \frac{P(R,S)}{P(R)}$$

Conditional state

$$\rho_{B|A}$$

Normalization condition

$$\text{Tr}_B(\rho_{B|A}) = I_A$$

Relation of conditional to joint

$$\rho_{B|A} = (\rho_A^{-1/2} \otimes I_B) \rho_{AB} (\rho_A^{-1/2} \otimes I_B)$$



Conditional probability

$$P(S|R)$$

Normalization condition

$$\sum_S P(S|R) = 1$$

Relation of conditional to joint

$$P(S|R) = \frac{P(R,S)}{P(R)}$$

$$P(R, S) = P(S|R)P(R)$$

Conditional state

$$\rho_{B|A}$$

Normalization condition

$$\text{Tr}_B(\rho_{B|A}) = I_A$$

Relation of conditional to joint

$$\rho_{B|A} = \rho_A^{-1/2} \rho_{AB} \rho_A^{-1/2}$$

$$\rho_{AB} = \rho_A^{1/2} \rho_{B|A} \rho_A^{1/2}$$

Conditional probability

$$P(S|R)$$

Normalization condition

$$\sum_S P(S|R) = 1$$

Relation of conditional to joint

$$P(S|R) = \frac{P(R,S)}{P(R)}$$

$$P(R,S) = P(S|R)P(R)$$

Classical belief propagation

$$P(S) = \sum_R P(S|R)P(R)$$

Conditional state

$$\rho_{B|A}$$

Normalization condition

$$\text{Tr}_B(\rho_{B|A}) = I_A$$

Relation of conditional to joint

$$\rho_{B|A} = \rho_{AB} * \rho_A^{-1}$$

$$\rho_{AB} = \rho_{B|A} * \rho_A$$

Quantum belief propagation

$$\rho_B = \text{Tr}_A(\rho_{B|A}\rho_A)$$



Two formulas for the  
joint probability

$$\begin{aligned}P(R, S) &= P(S|R)P(R) \\ &= P(R|S)P(S)\end{aligned}$$

Classical Bayes' theorem

$$P(S|R) = \frac{P(R|S)P(S)}{P(R)}$$

Two formulas for the  
joint state

$$\begin{aligned}\rho_{BA} &= \rho_A^{1/2} \rho_{B|A} \rho_A^{1/2} \\ &= \rho_B^{1/2} \rho_{A|B} \rho_B^{1/2}\end{aligned}$$

Quantum Bayes' theorem

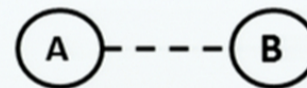
$$\rho_{B|A} = \rho_A^{-1/2} \rho_B^{1/2} \rho_{A|B} \rho_A^{-1/2} \rho_B^{1/2}$$



$$P(R, S)$$

$$P(S|R) = P(R, S)/P(R)$$

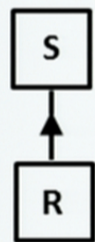
$$P(S) = \sum_R P(S|R)P(R)$$



$$\rho_{AB}$$

????

????



$$P(R, S)$$

$$P(S|R) = P(R, S)/P(R)$$

$$P(S) = \sum_R P(S|R)P(R)$$



????

????

????

$$\rho_B = \mathcal{E}_{A \rightarrow B}(\rho_A)$$



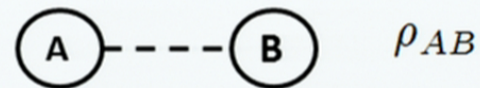


$$P(R, S)$$

$$P(S|R) = P(R, S)/P(R)$$

$$P(S) = \sum_R P(S|R)P(R)$$

$$P(S) = \Gamma_{R \rightarrow S}[P(R)]$$



$$\rho_{AB}$$

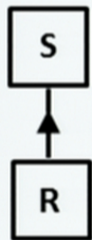
$$\rho_{B|A} = \rho_A^{-1/2} \rho_{AB} \rho_A^{-1/2}$$

$$\rho_B = \text{Tr}_A(\rho_{B|A} \rho_A)$$

$$\rho_B = \mathfrak{E}_{A \rightarrow B}(\rho_A)$$

$$\rho_{B|A} \geq 0$$

$$\mathfrak{E}_{A \rightarrow B} \circ T_A \text{ is CP}$$



$$P(R, S)$$

$$P(S|R) = P(R, S)/P(R)$$

$$P(S) = \sum_R P(S|R)P(R)$$

$$P(S) = \Gamma_{R \rightarrow S}[P(R)]$$



$$\varrho_{AB}$$

$$\varrho_{B|A} = \varrho_A^{-1/2} \varrho_{AB} \varrho_A^{-1/2}$$

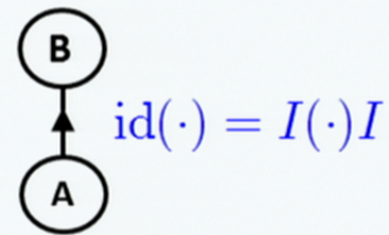
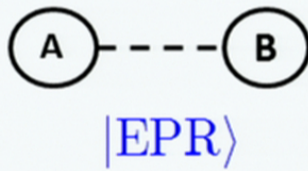
$$\rho_B = \text{Tr}_A(\varrho_{B|A} \rho_A)$$

$$\rho_B = \mathcal{E}_{A \rightarrow B}(\rho_A)$$

$$\varrho_{B|A}^{T_A} \geq 0$$

$$\mathcal{E}_{A \rightarrow B} \text{ is CP}$$

## Comparing causal and acausal correlations in Quantum Mechanics



	$\hat{Q}_A, \hat{Q}_B$	$\hat{P}_A, \hat{P}_B$
$ EPR\rangle$	C	A

	$\hat{Q}_A, \hat{Q}_B$	$\hat{P}_A, \hat{P}_B$
id	C	C



## Comparing causal and acausal correlations in Epistemically Restricted Liouville mechanics



$$P_{\text{EPR}}(q_A, p_A, q_B, p_B) \propto \delta(q_A - q_B) \delta(p_A + p_B)$$

$$Q_B - Q_A = 0$$

$$P_B + P_A = 0$$

$$P_{\text{id}}(q_B, p_B | q_A, p_A) \propto \delta(q_A - q_B) \delta(p_A - p_B)$$



$$Q_B = Q_A$$

$$P_B = P_A$$

	$Q_A, Q_B$	$P_A, P_B$
$P_{\text{EPR}}$	C	A

	$Q_A, Q_B$	$P_A, P_B$
$P_{\text{id}}$	C	C

## Jeffrey conditioning

Suppose  $P(S) = \sum_R P(S|R)P(R)$

If  $P(R) \rightarrow P^{\text{post}}(R)$

then  $P(S) \rightarrow P^{\text{post}}(S)$

where  $P^{\text{post}}(S) = \sum_R P(S|R)P^{\text{post}}(R)$

Suppose  $\rho_B = \text{Tr}_A(\rho_{B|A}\rho_A)$

If  $\rho_A \rightarrow \rho_A^{\text{post}}$

then  $\rho_B \rightarrow \rho_B^{\text{post}}$

where  $\rho_B^{\text{post}} = \text{Tr}_A(\rho_{B|A}\rho_A^{\text{post}})$

## Bayesian conditioning

Suppose  $P(S) = \sum_X P(S|X)P(X)$

If  $P(X) \rightarrow P^{\text{post}}(X) = \delta_{X,x}$

then  $P(S) \rightarrow P^{\text{post}}(S)$

where  $P^{\text{post}}(S) = \sum_X P(S|X)P^{\text{post}}(X)$   
 $= P(S|X=x)$

$P(S) \rightarrow P(S|X=x)$

Suppose  $\rho_B = \text{Tr}_X(\rho_{B|X}\rho_X)$

If  $\rho_X \rightarrow \rho_X^{\text{post}} = |x\rangle\langle x|_X$

then  $\rho_B \rightarrow \rho_B^{\text{post}}$

where  $\rho_B^{\text{post}} = \text{Tr}_X(\rho_{B|X}\rho_X^{\text{post}})$   
 $= \rho_{B|X=x}$

$\rho_B \rightarrow \rho_{B|X=x}$



## Other applications of the formalism

- Identify **analogies between multi-time and multi-system scenarios**  
e.g. no-broadcasting theorem  $\leftrightarrow$  monogamy of entanglement  
BB84 key distribution  $\leftrightarrow$  Ekert key distribution
- Accommodate **Aharonov *et al.* two-time and multi-time states** (pre and post selection is an instance of Bayesian inference)
- Obtain quantum analogues of key notions of Bayesian statistics  
e.g. **sufficient statistics, conditional independence, etc.**
- Multiple observers: **state compatibility criteria, state pooling rules, etc.**
- Quantum analogues of belief propagation algorithms are important for **quantum error correction**





