

Title: Recent Developments in Generalized ADM Formalism for the Dynamics of Compact Binaries with Spin

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Abstract: Based on tetrad-generalized canonical formalism by Arnowitt, Deser, and Misner most recent achievements in analytic calculations of higher order post-Newtonian Hamiltonians for spinning binary black holes and neutron stars are presented. The results of the generalized ADM formalism are put into mathematical relationship with those obtained within the Effective Field Theory approach.





Recent developments in generalized
ADM formalism for the dynamics of
compact binaries with spin

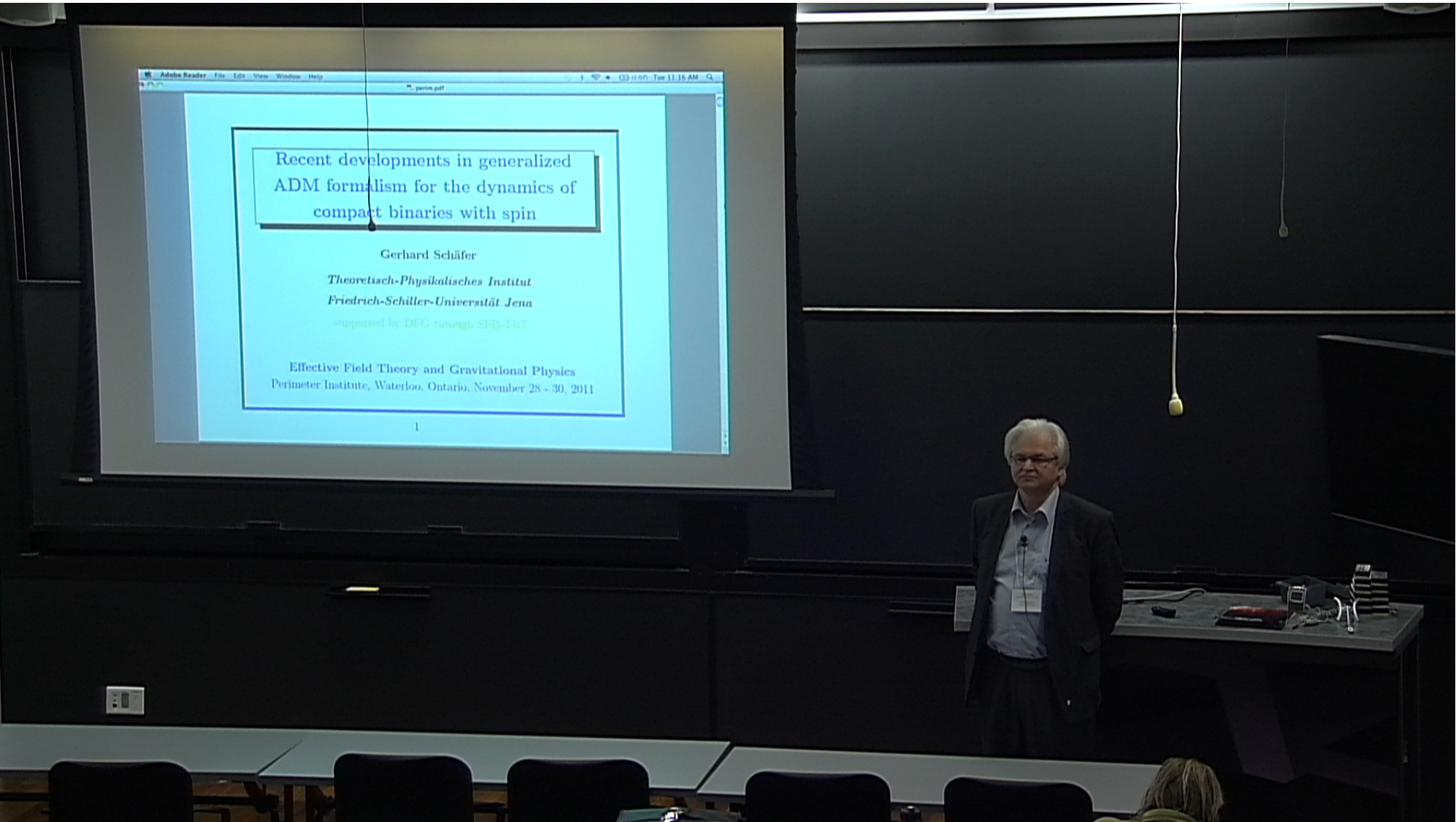
Gerhard Schäfer

*Theoretisch-Physikalisches Institut
Friedrich-Schiller-Universität Jena*

supported by DFG through SFB-TIT

Effective Field Theory and Gravitational Physics
Perimeter Institute, Waterloo, Ontario, November 28 - 30, 2011

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Outline

- Spin in Minkowski Space (introductory)
- Higher Order PN Hamiltonians (summary)
- Spin and Gravity (general)
- Hamiltonian Formulation for Spinning Objects
- Comparison with Results from EFT (recent preprint)
- Dirac Delta Functions and Black Holes

sometimes: $c = 1, G = 1$

Results based on the Hamiltonian approach

point-mass dynamics:

3PN: [Damour/Jaranowski/GS \(2001\)](#)

next-to-leading order rad. reaction (3.5PN): [Jaranowski/GS \(1997\)](#)

1PM: [Ledvinka/GS/Bicak \(2008\)](#)

linear-in-spin coupling: [Steinhoff/GS \(2009\)](#)

spin-orbit coupling (spin counted $1/c$):

next-to-leading order (2.5PN): [Damour/Jaranowski/GS \(2008\)](#)

next-to-next-to-leading order (3.5PN): [Hartung/Steinhoff \(2011\)](#)

leading order radiation reaction (4PN): [Steinhoff/Wang \(2010\)](#)

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spin(1)-spin(2) coupling (spin counted 1/c):

next-to-leading order (3PN): [Steinhoff/Hergt/GS \(2008\)](#)

next-to-next-to-leading order (4PN): [Hartung/Steinhoff \(2011\)](#)

lead. ord. rad. reaction (4.5PN): [Wang/Steinhoff/Zeng/GS \(2011\)](#)

spin(1)-spin(1) coupling (spin counted 1/c):

next-to-leading order for black holes and neutron stars (3PN):

[Steinhoff/Hergt/GS \(2008, 2010\)](#)

higher-order-in-spin coupling:

up to 3rd order in spin: $S(a)^n S(b)^{3-n} p(c)$: [Hergt/GS \(2008\)](#)

up to 4th order in spin: $S(a)^n S(b)^{4-n}$: [Hergt/GS \(2008\)](#)

spinning test-particle in Kerr: [Barausse/Racine/Buonanno \(2009\)](#)

el. charged particle with spin (m, e, S)

$$\frac{dp_\mu}{d\tau} = eF_{\mu\nu}u^\nu$$

$$\eta_{\mu\nu} \frac{dw^\nu}{d\tau} = g \frac{e}{2m} F_{\mu\nu} w^\nu + (g - 2) \frac{e}{2m^2} F_{\lambda\nu} u^\lambda w^\nu p_\mu$$

$$u^\mu = \frac{dx^\mu}{d\tau}, \quad -u^\mu u^\mu \eta_{\mu\nu} = 1, \quad p_\mu = m u_\mu, \quad -p_\mu p_\nu \eta^{\mu\nu} = m^2$$

$$w^\mu w^\nu \eta_{\mu\nu} = m^2 S^2, \quad p_\mu w^\mu = 0$$

el. charged particle with spin (m, e, \mathbf{S})

$$\frac{dp_\mu}{d\tau} = eF_{\mu\nu}u^\nu$$

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canonical variables

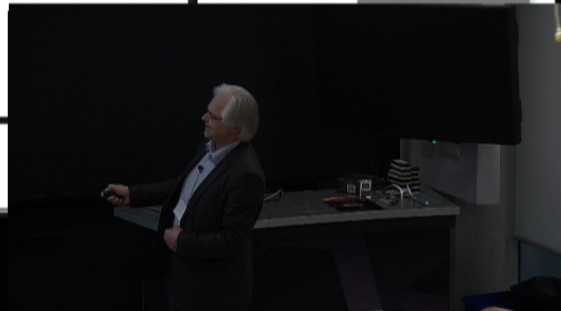
total angular momentum: $\mathbf{J} = \hat{\mathbf{X}} \times \mathbf{P} + \hat{\mathbf{S}}$

Lorentz boost: $\mathbf{K} = -t\mathbf{P} + H\hat{\mathbf{X}} - \frac{1}{H+m}\hat{\mathbf{S}} \times \mathbf{P}$

Hamiltonian: $H = \sqrt{m^2 + \mathbf{P}^2}$

center-of-energy: $\bar{\mathbf{X}} = \hat{\mathbf{X}} - \frac{1}{(H+m)H} \hat{\mathbf{S}} \times \mathbf{P}$

$$\mathbf{K} = -t\mathbf{P} + H\bar{\mathbf{X}}$$



$$w^0 = \mathbf{P} \cdot \mathbf{J} = \mathbf{P} \cdot \hat{\mathbf{S}}$$

$$\mathbf{w} = H\mathbf{J} + \mathbf{P} \times \mathbf{K} = H\hat{\mathbf{S}} - \frac{1}{H+m}\mathbf{P} \times (\hat{\mathbf{S}} \times \mathbf{P})$$

$$w^\mu = \frac{1}{2}\epsilon^{\mu\nu\alpha\beta}P_\nu J_{\alpha\beta} \quad J_{kl} = \epsilon_{klm}J_m \quad J_{0i} = K_i$$

$$w^\mu w_\mu \equiv -(w^0)^2 + \mathbf{w}^2 = m^2\hat{\mathbf{S}}^2$$

Poincaré algebra

$$\{P_i, H\} = \{J_i, H\} = 0$$

$$\{J_i, P_j\} = \varepsilon_{ijk} P_k, \quad \{J_i, J_j\} = \varepsilon_{ijk} J_k$$

$$\{J_i, G_j\} = \varepsilon_{ijk} G_k$$

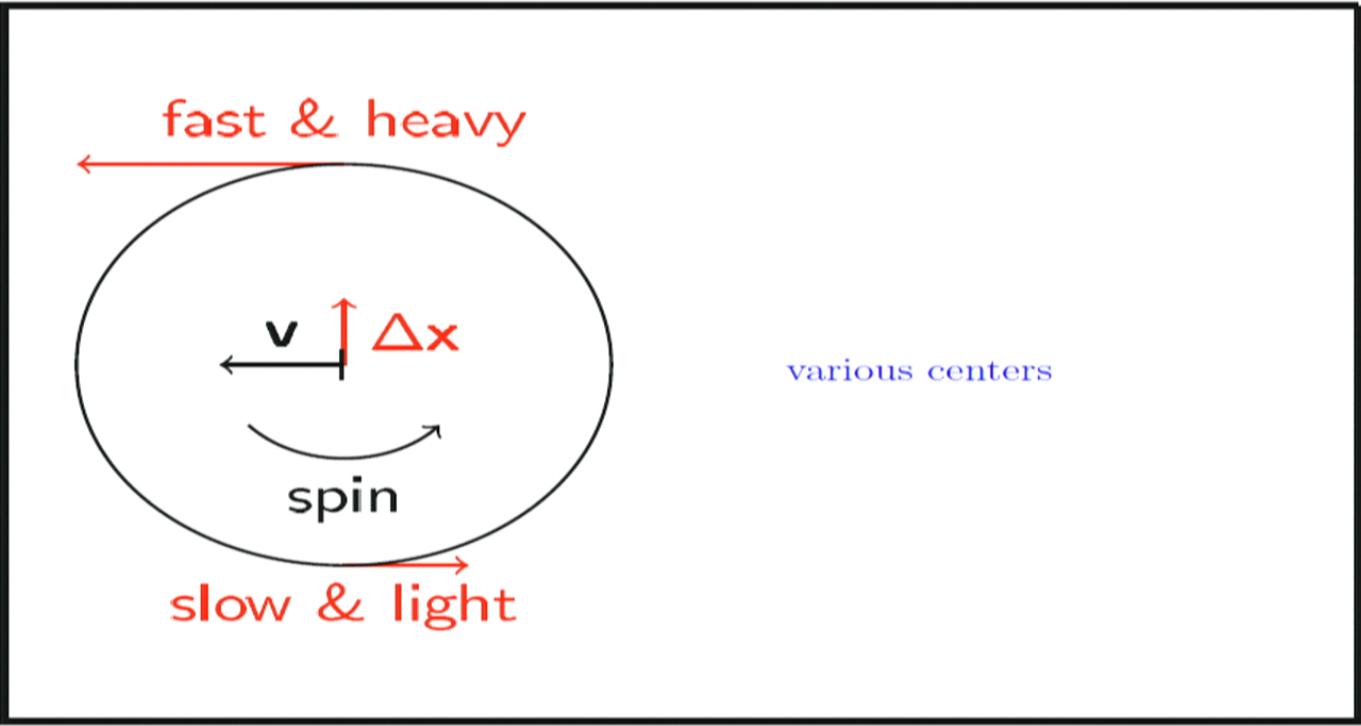
$$\{\mathbf{G}, H\} = \mathbf{P}$$

$$\{G_i, P_j\} = \frac{1}{c^2} H \delta_{ij}$$

$$\{G_i, G_j\} = -\frac{1}{c^2} \varepsilon_{ijk} J_k$$

$$\mathbf{K} = -t\mathbf{P} + \mathbf{G}$$

$$d\mathbf{K}/dt = \partial\mathbf{K}/\partial t + \{\mathbf{K}, H\} = -\mathbf{P} + \{\mathbf{G}, H\} = 0$$



various centers

center-of-spin: $\hat{\mathbf{X}}$; $\{\hat{X}^i, \hat{X}^j\} = 0$ (Newton-Wigner coordinates)

center-of-energy: $\bar{\mathbf{X}} = \hat{\mathbf{X}} - \frac{1}{(H+m)H} \hat{\mathbf{S}} \times \mathbf{P}$

center-of-inertia: $\mathbf{X} = \hat{\mathbf{X}} + \frac{1}{(H+m)m} \hat{\mathbf{S}} \times \mathbf{P}$

center-of-inertia: $S^{\mu\nu} P_\nu = 0$

center-of-energy: $\bar{S}^{\mu\nu} n_\nu = 0, \quad n_\mu = (-1, 0, 0, 0)$

center-of-spin: $m\hat{S}^{\mu\nu} n_\nu + \hat{S}^{\mu\nu} P_\nu = 0$

minimal radius of particle with spin: $\frac{\hat{S}}{mc}$

Compton wavelength: $2\frac{\hbar/2}{mc}$

radius of Kerr ring singularity (Boyer-Lindquist coord's): $\frac{\hat{S}}{mc}$

many-particle systems with interaction

$$\mathbf{P} = \sum_a \mathbf{p}_a$$

$$\mathbf{J} = \sum_a (\mathbf{r}_a \times \mathbf{p}_a + \mathbf{s}_a)$$

$$M = \sqrt{H^2 - \mathbf{P}^2}, \quad H = \sqrt{M^2 + \mathbf{P}^2}$$

$$\hat{\mathbf{X}} = \frac{\mathbf{G}}{H} + \frac{1}{M(H + M)} (\mathbf{J} - \frac{\mathbf{G}}{H} \times \mathbf{P}) \times \mathbf{P}$$

$$\{\hat{X}^i, \hat{X}^j\} = \{P^i, P^j\} = 0, \quad \{\hat{X}^i, P^j\} = \delta^{ij}$$

$$\{M, \hat{X}^j\} = \{M, P^j\} = \{M, H\} = 0$$

free particles with spin:

$$H = \sum_a h_a, \quad h_a = \sqrt{m_a^2 + \mathbf{p}_a^2}$$

$$\mathbf{G} = \sum_a \left(h_a \mathbf{r}_a - \frac{1}{h_a + m_a} \mathbf{s}_a \times \mathbf{p}_a \right)$$

Higher Order PN Hamiltonians (summary)

spin not counted $1/c$

binary black holes to 3.5 PN order

$$\begin{aligned} H(t) &= m_1 c^2 + m_2 c^2 + H_N + H_{1PN} \\ &+ H_{2PN} + H_{3PN} + \dots \\ &+ \frac{1}{c^5} H_{2.5PN}(t) + \frac{1}{c^7} H_{3.5PN}(t) + \dots \end{aligned}$$

$$\begin{aligned} \hat{H} &= (H - Mc^2)/\mu, \quad \mu = m_1 m_2 / M, \quad M = m_1 + m_2 \\ \nu &= \mu / M, \quad 0 \leq \nu \leq 1/4 \end{aligned}$$

test particles: $\nu = 0$, equal masses: $\nu = 1/4$

$$\begin{aligned} \text{CMF: } \mathbf{p}_1 + \mathbf{p}_2 &= 0, \quad \mathbf{p} = \mathbf{p}_1 / \mu, \quad r = r_{12} = |\mathbf{x}_1 - \mathbf{x}_2|, \\ p_r &= (\mathbf{n} \cdot \mathbf{p}), \quad \mathbf{q} = (\mathbf{x}_1 - \mathbf{x}_2) / GM, \quad \mathbf{n} = \mathbf{n}_{12} = \mathbf{q} / |\mathbf{q}| \end{aligned}$$

$$\hat{H}_N = \frac{p^2}{2} - \frac{1}{q}$$
$$c^2 \hat{H}_{1PN} = \frac{1}{8}(3\nu - 1)p^4 - \frac{1}{2}[(3 + \nu)p^2 + \nu p_r^2] \frac{1}{q} + \frac{1}{2q^2}$$
$$c^4 \hat{H}_{2PN} = \frac{1}{16}(1 - 5\nu + 5\nu^2)p^6$$
$$+ \frac{1}{8}[(5 - 20\nu - 3\nu^2)p^4 - 2\nu^2 p_r^2 p^2 - 3\nu^2 p_r^4] \frac{1}{q}$$
$$+ \frac{1}{2}[(5 + 8\nu)p^2 + 3\nu p_r^2] \frac{1}{q^2} - \frac{1}{4}(1 + 3\nu) \frac{1}{q^3}$$

$$\begin{aligned}
c^6 \hat{H}_{3PN} &= \frac{1}{128} (-5 + 35\nu - 70\nu^2 + 35\nu^3) p^8 \\
&+ \frac{1}{16} [(-7 + 42\nu - 53\nu^2 - 5\nu^3) p^6 + (2 - 3\nu) \nu^2 p_r^2 p^4 \\
&+ 3(1 - \nu) \nu^2 p_r^4 p^2 - 5\nu^3 p_r^6] \frac{1}{q} \\
&+ \left[\frac{1}{16} (-27 + 136\nu + 109\nu^2) p^4 + \frac{1}{16} (17 + 30\nu) \nu p_r^2 p^2 \right. \\
&+ \left. \frac{1}{12} (5 + 43\nu) \nu p_r^4 \right] \frac{1}{q^2} \\
&+ \left[\left(-\frac{25}{8} + \left(\frac{1}{64} \pi^2 - \frac{335}{48} \right) \nu - \frac{23}{8} \nu^2 \right) p^2 \right. \\
&+ \left. \left(-\frac{85}{16} - \frac{3}{64} \pi^2 - \frac{7}{4} \nu \right) \nu p_r^2 \right] \frac{1}{q^3} + \left[\frac{1}{8} + \left(\frac{109}{12} - \frac{21}{32} \pi^2 \right) \nu \right] \frac{1}{q^4}
\end{aligned}$$

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spin-gravity interaction ($\mathbf{S} \equiv \hat{\mathbf{S}}$)

leading order spin orbit

$$H_{SO}^{1PN} = \frac{G}{c^2} \sum_a \sum_{b \neq a} \frac{1}{r_{ab}^2} (\mathbf{S}_a \times \mathbf{n}_{ab}) \cdot \left[\frac{3m_b}{2m_a} \mathbf{p}_a - 2\mathbf{p}_b \right]$$

leading order spin(1)-spin(2)

$$H_{S_1 S_2}^{1PN} = \frac{G}{c^2} \sum_a \sum_{b \neq a} \frac{1}{2r_{ab}^3} [3(\mathbf{S}_a \cdot \mathbf{n}_{ab})(\mathbf{S}_b \cdot \mathbf{n}_{ab}) - (\mathbf{S}_a \cdot \mathbf{S}_b)]$$

leading order spin(1)-spin(1)

$$H_{S_1 S_1}^{1PN} = \frac{G}{c^2} \frac{1}{2r_{12}^3} [3(\mathbf{S}_1 \cdot \mathbf{n}_{12})(\mathbf{S}_1 \cdot \mathbf{n}_{12}) - (\mathbf{S}_1 \cdot \mathbf{S}_1)]$$

$$\begin{aligned}
 H_{SO}^{2PN} = & \frac{G}{c^4 r^2} \left[- ((\mathbf{p}_1 \times \mathbf{S}_1) \cdot \mathbf{n}_{12}) \left[\frac{5m_2 \mathbf{p}_1^2}{8m_1^3} + \frac{3(\mathbf{p}_1 \cdot \mathbf{p}_2)}{4m_1^2} \right. \right. \\
 & - \left. \frac{3\mathbf{p}_2^2}{4m_1 m_2} + \frac{3(\mathbf{p}_1 \cdot \mathbf{n}_{12})(\mathbf{p}_2 \cdot \mathbf{n}_{12})}{4m_1^2} + \frac{3(\mathbf{p}_2 \cdot \mathbf{n}_{12})^2}{2m_1 m_2} \right] \\
 & + ((\mathbf{p}_2 \times \mathbf{S}_1) \cdot \mathbf{n}_{12}) \left[\frac{(\mathbf{p}_1 \cdot \mathbf{p}_2)}{m_1 m_2} + \frac{3(\mathbf{p}_1 \cdot \mathbf{n}_{12})(\mathbf{p}_2 \cdot \mathbf{n}_{12})}{m_1 m_2} \right] \\
 & + ((\mathbf{p}_1 \times \mathbf{S}_1) \cdot \mathbf{p}_2) \left[\frac{2(\mathbf{p}_2 \cdot \mathbf{n}_{12})}{m_1 m_2} - \frac{3(\mathbf{p}_1 \cdot \mathbf{n}_{12})}{4m_1^2} \right] \\
 & + \frac{G^2}{c^4 r^3} \left[- ((\mathbf{p}_1 \times \mathbf{S}_1) \cdot \mathbf{n}_{12}) \left[\frac{11m_2}{2} + \frac{5m_2^2}{m_1} \right] \right. \\
 & \left. + ((\mathbf{p}_2 \times \mathbf{S}_1) \cdot \mathbf{n}_{12}) \left[6m_1 + \frac{15m_2}{2} \right] \right] + (1 \leftrightarrow 2)
 \end{aligned}$$

$$\begin{aligned}
 H_{S_1 S_2}^{2PN} = & (G/2m_1 m_2 c^4 r^3) [3((\mathbf{p}_1 \times \mathbf{S}_1) \cdot \mathbf{n}_{12})(\mathbf{p}_2 \times \mathbf{S}_2) \cdot \mathbf{n}_{12})/2 \\
 + & 6((\mathbf{p}_2 \times \mathbf{S}_1) \cdot \mathbf{n}_{12})(\mathbf{p}_1 \times \mathbf{S}_2) \cdot \mathbf{n}_{12}) \\
 - & 15(\mathbf{S}_1 \cdot \mathbf{n}_{12})(\mathbf{S}_2 \cdot \mathbf{n}_{12})(\mathbf{p}_1 \cdot \mathbf{n}_{12})(\mathbf{p}_2 \cdot \mathbf{n}_{12}) \\
 - & 3(\mathbf{S}_1 \cdot \mathbf{n}_{12})(\mathbf{S}_2 \cdot \mathbf{n}_{12})(\mathbf{p}_1 \cdot \mathbf{p}_2) + 3(\mathbf{S}_1 \cdot \mathbf{p}_2)(\mathbf{S}_2 \cdot \mathbf{n}_{12})(\mathbf{p}_1 \cdot \mathbf{n}_{12}) \\
 + & 3(\mathbf{S}_2 \cdot \mathbf{p}_1)(\mathbf{S}_1 \cdot \mathbf{n}_{12})(\mathbf{p}_2 \cdot \mathbf{n}_{12}) + 3(\mathbf{S}_1 \cdot \mathbf{p}_1)(\mathbf{S}_2 \cdot \mathbf{n}_{12})(\mathbf{p}_2 \cdot \mathbf{n}_{12}) \\
 + & 3(\mathbf{S}_2 \cdot \mathbf{p}_2)(\mathbf{S}_1 \cdot \mathbf{n}_{12})(\mathbf{p}_1 \cdot \mathbf{n}_{12}) - 3(\mathbf{S}_1 \cdot \mathbf{S}_2)(\mathbf{p}_1 \cdot \mathbf{n}_{12})(\mathbf{p}_2 \cdot \mathbf{n}_{12}) \\
 + & (\mathbf{S}_1 \cdot \mathbf{p}_1)(\mathbf{S}_2 \cdot \mathbf{p}_2) - (\mathbf{S}_1 \cdot \mathbf{p}_2)(\mathbf{S}_2 \cdot \mathbf{p}_1)/2 + (\mathbf{S}_1 \cdot \mathbf{S}_2)(\mathbf{p}_1 \cdot \mathbf{p}_2)/2] \\
 + & (3/2m_1^2 r^3) [-((\mathbf{p}_1 \times \mathbf{S}_1) \cdot \mathbf{n}_{12})(\mathbf{p}_1 \times \mathbf{S}_2) \cdot \mathbf{n}_{12}) \\
 + & (\mathbf{S}_1 \cdot \mathbf{S}_2)(\mathbf{p}_1 \cdot \mathbf{n}_{12})^2 - (\mathbf{S}_1 \cdot \mathbf{n}_{12})(\mathbf{S}_2 \cdot \mathbf{p}_1)(\mathbf{p}_1 \cdot \mathbf{n}_{12})] \\
 + & (3/2m_2^2 r^3) [-((\mathbf{p}_2 \times \mathbf{S}_2) \cdot \mathbf{n}_{12})(\mathbf{p}_2 \times \mathbf{S}_1) \cdot \mathbf{n}_{12}) \\
 + & (\mathbf{S}_1 \cdot \mathbf{S}_2)(\mathbf{p}_2 \cdot \mathbf{n}_{12})^2 - (\mathbf{S}_2 \cdot \mathbf{n}_{12})(\mathbf{S}_1 \cdot \mathbf{p}_2)(\mathbf{p}_2 \cdot \mathbf{n}_{12})] \\
 + & (6G^2(m_1 + m_2)/c^4 r^4) [(\mathbf{S}_1 \cdot \mathbf{S}_2) - 2(\mathbf{S}_1 \cdot \mathbf{n}_{12})(\mathbf{S}_2 \cdot \mathbf{n}_{12})]
 \end{aligned}$$

$$\begin{aligned}
H_{S_1 S_1}^{2PN} = \frac{G}{c^4 r^3} & \left[-\frac{5m_2}{4m_1^3} (\mathbf{p}_1 \cdot \mathbf{S}_1)^2 + \frac{m_2}{m_1^3} \mathbf{p}_1^2 \mathbf{S}_1^2 - \frac{21m_2}{8m_1^3} (\mathbf{p}_1 \cdot \mathbf{n})^2 \mathbf{S}_1^2 \right. \\
& - \frac{3m_2}{8m_1^3} \mathbf{p}_1^2 (\mathbf{S}_1 \cdot \mathbf{n})^2 + \frac{15m_2}{4m_1^3} (\mathbf{p}_1 \cdot \mathbf{n}) (\mathbf{S}_1 \cdot \mathbf{n}) (\mathbf{p}_1 \cdot \mathbf{S}_1) - \frac{3}{4m_1 m_2} \mathbf{p}_2^2 \mathbf{S}_1^2 \\
& + \frac{9}{4m_1 m_2} \mathbf{p}_2^2 (\mathbf{S}_1 \cdot \mathbf{n})^2 - \frac{1}{4m_1^2} (\mathbf{p}_1 \cdot \mathbf{p}_2) \mathbf{S}_1^2 - \frac{9}{4m_1^2} (\mathbf{p}_1 \cdot \mathbf{p}_2) (\mathbf{S}_1 \cdot \mathbf{n})^2 \\
& + \frac{3}{2m_1^2} (\mathbf{p}_1 \cdot \mathbf{S}_1) (\mathbf{p}_2 \cdot \mathbf{S}_1) - \frac{3}{2m_1^2} (\mathbf{p}_1 \cdot \mathbf{n}) (\mathbf{p}_2 \cdot \mathbf{S}_1) (\mathbf{S}_1 \cdot \mathbf{n}) \\
& - \frac{3}{2m_1^2} (\mathbf{p}_2 \cdot \mathbf{n}) (\mathbf{p}_1 \cdot \mathbf{S}_1) (\mathbf{S}_1 \cdot \mathbf{n}) + \frac{15}{4m_1^2} (\mathbf{p}_1 \cdot \mathbf{n}) (\mathbf{p}_2 \cdot \mathbf{n}) \mathbf{S}_1^2 \\
& \left. - \frac{15}{4m_1^2} (\mathbf{p}_1 \cdot \mathbf{n}) (\mathbf{p}_2 \cdot \mathbf{n}) (\mathbf{S}_1 \cdot \mathbf{n})^2 \right] \\
& - \frac{G^2 m_2}{2c^4 r^4} \left[5 \left(1 + \frac{m_2}{m_1} \right) ((\mathbf{S}_1 \cdot \mathbf{n})^2 - \mathbf{S}_1^2) + 4 \left(1 + \frac{2m_2}{m_1} \right) (\mathbf{S}_1 \cdot \mathbf{n})^2 \right]
\end{aligned}$$

NLO center-of-mass:

$$\begin{aligned}
 \mathbf{G}_{\text{SO}}^{2PN} = & - \sum_a \frac{\mathbf{P}_a^2}{8m_a^3} (\mathbf{P}_a \times \mathbf{S}_a) \\
 & + \sum_a \sum_{b \neq a} \frac{m_b}{4m_a r_{ab}} \left[((\mathbf{P}_a \times \mathbf{S}_a) \cdot \mathbf{n}_{ab}) \frac{5\mathbf{x}_a + \mathbf{x}_b}{r_{ab}} - 5(\mathbf{P}_a \times \mathbf{S}_a) \right] \\
 & + \sum_a \sum_{b \neq a} \frac{1}{r_{ab}} \left[\frac{3}{2} (\mathbf{P}_b \times \mathbf{S}_a) - \frac{1}{2} (\mathbf{n}_{ab} \times \mathbf{S}_a) (\mathbf{P}_b \cdot \mathbf{n}_{ab}) \right. \\
 & \quad \left. - ((\mathbf{P}_a \times \mathbf{S}_a) \cdot \mathbf{n}_{ab}) \frac{\mathbf{x}_a + \mathbf{x}_b}{r_{ab}} \right] \\
 \\
 \mathbf{G}_{\text{SS}}^{2PN} = & \frac{1}{2} \sum_a \sum_{b \neq a} \left\{ [3(\mathbf{S}_a \cdot \mathbf{n}_{ab})(\mathbf{S}_b \cdot \mathbf{n}_{ab}) - (\mathbf{S}_a \cdot \mathbf{S}_b)] \frac{\mathbf{x}_a}{r_{ab}^3} + (\mathbf{S}_b \cdot \mathbf{n}_{ab}) \frac{\mathbf{S}_a}{r_{ab}^2} \right\}
 \end{aligned}$$

$$\begin{aligned}
 H &= H_N + H_{1PN} + H_{2PN} + H_{2.5PN} + H_{3PN} + H_{3.5PN} \\
 &+ H_{SO}^{1PN} + H_{SO}^{2PN} + H_{SO}^{3PN} + H_{SO}^{3.5PN} \\
 &+ H_{S_1S_2}^{1PN} + H_{S_1S_2}^{2PN} + H_{S_1S_2}^{3PN} + H_{S_1S_2}^{3.5PN} \\
 &+ H_{S_1S_1}^{1PN} + H_{S_2S_2}^{1PN} + H_{S_1S_1}^{2PN} + H_{S_2S_2}^{2PN} \\
 &+ H_{p_1S_2^3} + H_{p_2S_1^3} + H_{p_1S_1S_2^2} + H_{p_2S_2S_1^2} \\
 &+ H_{S_1S_2^3} + H_{S_1^3S_2} + H_{S_1^2S_2^2}
 \end{aligned}$$

$$H = H_{1PM}$$

$$H = H(t) = H[p, q, h^{\text{TT}}(x; p', q'), \pi^{\text{TT}}(x; p', q')] = H(p, q, p', q')$$

tetrad field e_a^μ : $e_a^\mu e_{b\mu} = \eta_{ab}$ $e_{a\mu} e_{b\nu} \eta^{ab} = g_{\mu\nu} = g_{\nu\mu}$

local rotations: $e_a'^\mu = L^b{}_a e_b^\mu$ $L^a{}_c \eta_{ab} L^b{}_d = \eta_{cd}$

linear connection ω_μ^{ab} : $D_\mu \phi \equiv \partial_\mu \phi + \frac{1}{2} \omega_\mu^{ab} G_{[ab]} \phi$

local rotations: $\omega_\mu'^{ab} = L^a{}_c L^b{}_d \omega_\mu^{cd} + L^a{}_d \partial_\mu L^{bd}$ $\partial_\mu \equiv \frac{\partial}{\partial x^\mu}$

curvature tensor $R^ab{}_{\mu\nu}$: $D_\mu D_\nu \phi - D_\nu D_\mu \phi = R^ab{}_{\mu\nu} G_{[ab]} \phi$

$$R^ab{}_{\mu\nu} = \partial_\mu \omega_\nu^{ab} - \partial_\nu \omega_\mu^{ab} + \omega_\nu^{ac} \omega_\mu^{bd} \eta_{cd} - \omega_\mu^{ac} \omega_\nu^{bd} \eta_{cd}$$

Lagrangian for gravity

$$\mathcal{L}_G = \frac{1}{16\pi} \det(e_\gamma^c) e_a^\mu e_b^\nu R_{\mu\nu}^{ab}(\omega) + \partial_\mu C^\mu$$

vacuum Einstein equations:

$$0 = \frac{\delta \mathcal{L}_G}{\delta e_a^\mu} e_{a\nu} \equiv 2 \det(e_\gamma^c) (R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R)$$

$$0 = \frac{\delta \mathcal{L}_G}{\delta \omega_\mu^{ab}} \Rightarrow \omega_\mu^{ab} = \omega_\mu^{ab}(e, \partial_\nu e) \quad \text{no torsion !}$$

Lagrangian for spinning objects

$$\mathcal{L}_M = \int d\tau \left[\left(p_\mu - \frac{1}{2} S_{ab} \omega_\mu^{ab} \right) \frac{dz^\mu}{d\tau} + \frac{1}{2} S_{ab} \frac{d\theta^{ab}}{d\tau} \right] \delta_{(4)}$$

$$\mathcal{L}_C = \int d\tau \left[\lambda_1^a p^b S_{ab} + \lambda_2^{[i]a} \Lambda^{[i]a} p_a - \frac{\lambda_3}{2} (p^2 + m^2) \right] \delta_{(4)}$$

$$d\theta^{ab} = \Lambda_C^a d\Lambda^{Cb} = -d\theta^{ba}$$

matter action: $W_M = \int d^4x [\mathcal{L}_M + \mathcal{L}_C]$

equations of motion

$$\frac{DS_{ab}}{D\tau} = 0$$

$$\frac{Dp_{\mu}}{D\tau} = -\frac{1}{2}R_{\mu\rho ab}^{(4)}u^{\rho}S^{ab}$$

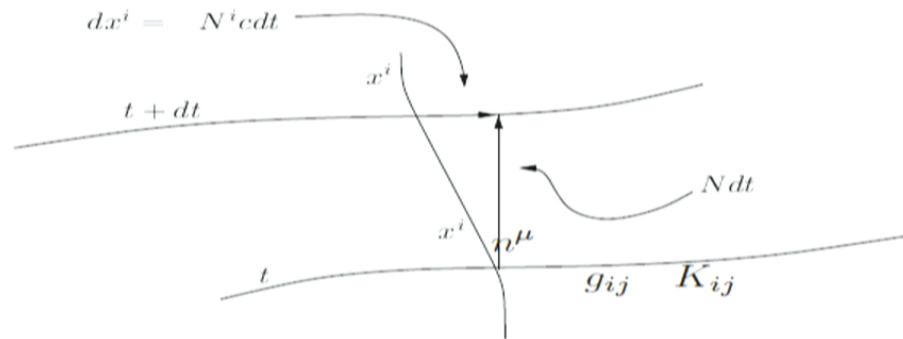
$$u^{\mu} \equiv \frac{dz^{\mu}}{d\tau} = \lambda_3 p^{\mu}$$

$$\sqrt{-g}T^{\mu\nu} = \int d\tau \left[\lambda_3 p^{\mu} p^{\nu} \delta_{(4)} + \left(u^{(\mu} S^{\nu)\alpha} \delta_{(4)} \right)_{||\alpha} \right]$$

3+1 splitting of spacetime

$$n^\mu = (1, -N^i)/N$$

$$n_\mu = (-N, 0, 0, 0)$$



$$K_{ij} = -N\Gamma_{ij}^0 = -Ng^{0\mu}(g_{i\mu,j} + g_{j\mu,i} - g_{ij,\mu})/2$$

$$ds^2 = -(Ncdt)^2 + g_{ij}(dx^i + N^i cdt)(dx^j + N^j cdt)$$

solution of the matter constraints

$$n^\mu = (1, -N^i)/N, \quad n_\mu = (-N, 0, 0, 0)$$

$$\lambda_3 : np \equiv n^\mu p_\mu = -\sqrt{m^2 + \gamma^{ij} p_i p_j}, \quad \gamma^{ik} g_{kj} = \delta_j^i$$

$$\lambda_1 : nS_i \equiv n^\mu S_{\mu i} = \frac{p_k \gamma^{kj} S_{ji}}{np} = g_{ij} nS^j$$

$$\lambda_2 : \Lambda^{[j](0)} = \Lambda^{[j](i)} \frac{P(i)}{p^{(0)}}, \quad \Lambda^{[0]a} = -\frac{p^a}{m}$$

time gauge for the tetrads

$$e_{(0)}^\mu = n^\mu, \quad e_{(0)}^0 = \frac{1}{N}, \quad e_{(0)}^i = -\frac{N^i}{N}$$

$$g_{ij} = e_i^{(m)} e_{(m)j}$$

$$\mathcal{L}_{MC} = -N\mathcal{H}^{\text{matter}} + N^i \mathcal{H}_i^{\text{matter}}$$

$$\mathcal{H}^{\text{matter}} = -np\delta - K^{ij} \frac{p_i n S_j}{np} \delta - (n S^k \delta)_{;k}$$

$$\mathcal{H}_i^{\text{matter}} = (p_i + K_{ij} n S^j) \delta + \left(\frac{1}{2} \gamma^{mk} S_{ik} \delta + \delta_i^{(k} \gamma^{l)m} \frac{p_k n S_l}{np} \delta \right)_{;m}$$

transformation to canonical matter variables

$$z^i = \hat{z}^i - \frac{nS^i}{m - np}, \quad nS_i = -\frac{p_k \gamma^{kj} \hat{S}_{ji}}{m}$$

$$S_{ij} = \hat{S}_{ij} - \frac{p_i n S_j}{m - np} + \frac{p_j n S_i}{m - np}$$

$$\Lambda^{[i](j)} = \hat{\Lambda}^{[i](k)} \left(\delta_{kj} + \frac{p^{(k)} p^{(j)}}{m(m - np)} \right)$$

$$P_i = p_i + K_{ij}nS^j + \hat{A}^{kl}e_{(j)k}e_{l,i}^{(j)} - \left(\frac{1}{2}S_{kj} + \frac{p(knS_j)}{np} \right) \Gamma^{kj}_i$$

$$g_{ik}g_{jl}\hat{A}^{kl} = \frac{1}{2}\hat{S}_{ij} + \frac{mp_{(i}nS_{j)}}{np(m-np)}$$

$$S^{ab}S_{ab} = \hat{S}_{(i)(j)}\hat{S}_{(i)(j)} = 2\hat{S}_{(i)}\hat{S}_{(i)} = 2s^2 = \text{const}$$

$$\hat{\Lambda}_{[k]}^{(i)}\hat{\Lambda}^{[k](j)} = \delta_{ij}$$

$$d\hat{\theta}^{(i)(j)} \equiv \hat{\Lambda}_{[k]}^{(i)}d\hat{\Lambda}^{[k](j)} = -d\hat{\theta}^{(j)(i)}$$

adding Lagrangian of gravity

$$\hat{\mathcal{L}}_{MK} = P_i \dot{z}^i \delta + \frac{1}{2} \hat{S}_{(i)(j)} \dot{\theta}^{(i)(j)} \delta$$

$$\hat{\mathcal{L}}_{GK} = \hat{A}^{ij} e_{(k)i} e_{j,0}^{(k)} \delta$$

$$\hat{\mathcal{L}}_{GK} + \mathcal{L}_G = \frac{1}{8\pi} [\pi^{ij} + 8\pi \hat{A}^{ij} \delta] e_{(k)i} e_{j,0}^{(k)} + \mathcal{L}_{GC} - \frac{1}{16\pi} \mathcal{E}_{i,i}$$

$$\mathcal{E}_i = g_{ij,j} - g_{jj,i}$$

total energy: $E = \frac{1}{16\pi} \oint d^2 s_i \mathcal{E}_i = \frac{1}{16\pi} \int d^3 x \mathcal{E}_{i,i}$

$$\mathcal{L}_{GC} = -N\mathcal{H}^{\text{field}} + N^i\mathcal{H}_i^{\text{field}}$$

$$\mathcal{H}^{\text{field}} = -\frac{1}{16\pi\sqrt{\gamma}} \left[\gamma R + \frac{1}{2} (g_{ij}\pi^{ij})^2 - g_{ij}g_{kl}\pi^{ik}\pi^{jl} \right]$$

$$\mathcal{H}_i^{\text{field}} = \frac{1}{8\pi} g_{ij}\pi^{jk}_{;k}$$

$$\pi^{ij} = \sqrt{\gamma}(\gamma^{ij}\gamma^{kl} - \gamma^{ik}\gamma^{jl})K_{kl} \quad \gamma \equiv \det(g_{ij})$$

$$W_G = \int d^4x \left[\frac{1}{8\pi} \pi^{ij} e_{(k)i} e_{j,0}^{(k)} - N\mathcal{H}^{\text{field}} + N^i\mathcal{H}_i^{\text{field}} - \frac{1}{16\pi} \mathcal{E}_{i,i} \right]$$

$$\mathcal{L}_{GC} = -N\mathcal{H}^{\text{field}} + N^i\mathcal{H}_i^{\text{field}}$$

$$\mathcal{H}^{\text{field}} = -\frac{1}{16\pi\sqrt{\gamma}} \left[\gamma R + \frac{1}{2} (g_{ij}\pi^{ij})^2 - g_{ij}g_{kl}\pi^{ik}\pi^{jl} \right]$$

$$\mathcal{H}_i^{\text{field}} = \frac{1}{8\pi} g_{ij}\pi^{jk}_{;k}$$

$$\pi^{ij} = \sqrt{\gamma}(\gamma^{ij}\gamma^{kl} - \gamma^{ik}\gamma^{jl})K_{kl} \quad \gamma \equiv \det(g_{ij})$$

$$W_G = \int d^4x \left[\frac{1}{8\pi} \pi^{ij} e_{(k)i} e_{j,0}^{(k)} - N\mathcal{H}^{\text{field}} + N^i\mathcal{H}_i^{\text{field}} - \frac{1}{16\pi} \mathcal{E}_{i,i} \right]$$

spatially symmetric time gauge for the tetrads

$$e_{(k)i} e_{j,\mu}^{(k)} = B_{ij}^{kl} g_{kl,\mu} + \frac{1}{2} g_{ij,\mu} \quad B_{ij}^{kl} = B_{[ij]}^{(kl)}$$

$$e_{(i)j} = e_{ij} = e_{ji}$$

$$e_{ij} e_{jk} = g_{ik} \quad e_{ij} = \sqrt{(g_{kl})}$$

$$\pi_{\text{can}}^{ij} = \pi^{ij} + 8\pi \hat{A}^{(ij)} \delta + 16\pi B_{kl}^{ij} \hat{A}^{[kl]} \delta$$

spacetime-coordinates conditions

$$3g_{ij,j} - g_{jj,i} = 0, \quad \pi_{\text{can}}^{ii} = 0$$

$$g_{ij} = \Psi^4 \delta_{ij} + h_{ij}^{\text{TT}}, \quad \pi_{\text{can}}^{ij} = \tilde{\pi}_{\text{can}}^{ij} + \pi_{\text{can}}^{ij\text{TT}}$$

transverse traceless: $h_{ii}^{\text{TT}} = \pi_{\text{can}}^{ii\text{TT}} = h_{ij,j}^{\text{TT}} = \pi_{\text{can},j}^{ij\text{TT}} = 0$

$$\tilde{\pi}_{\text{can}}^{ij} = V_{\text{can},j}^i + V_{\text{can},i}^j - \frac{2}{3} \delta_{ij} V_{\text{can},k}^k$$

constraints: $\mathcal{H}^{\text{field}} + \mathcal{H}^{\text{matter}} = 0, \quad \mathcal{H}_i^{\text{field}} + \mathcal{H}_i^{\text{matter}} = 0$

total action in canonical form

$$W = \frac{1}{16\pi} \int d^4x \pi_{\text{can}}^{ij\text{TT}} h_{ij,0}^{\text{TT}} + \int dt \left[P_i \dot{z}^i + \frac{1}{2} \hat{S}_{(i)(j)} \dot{\theta}^{(i)(j)} - E \right]$$

Hamiltonian: $E \equiv H_{\text{ADM}} = -\frac{1}{2\pi} \int d^3x \Delta\Psi \left[\dot{z}^i, P_i, \hat{S}_{(i)(j)}, h_{ij}^{\text{TT}}, \pi_{\text{can}}^{ij\text{TT}} \right]$

$$\{\dot{z}^i, P_j\} = \delta_{ij}, \quad \{\hat{S}_{(i)}, \hat{S}_{(j)}\} = \epsilon_{ijk} \hat{S}_{(k)}$$

$$\{h_{ij}^{\text{TT}}(\mathbf{x}, t), \pi_{\text{can}}^{kl\text{TT}}(\mathbf{x}', t)\} = 16\pi \delta_{ij}^{\text{TT}kl} \delta(\mathbf{x} - \mathbf{x}')$$

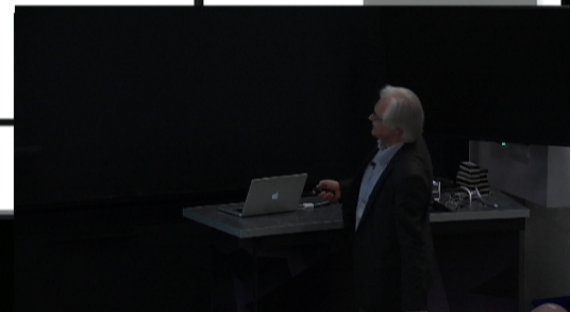
total action in canonical form

$$W = \frac{1}{16\pi} \int d^4x \pi_{\text{can}}^{ij\text{TT}} h_{ij,0}^{\text{TT}} + \int dt \left[P_i \dot{z}^i + \frac{1}{2} \hat{S}_{(i)(j)} \dot{\theta}^{(i)(j)} - E \right]$$

Hamiltonian: $E \equiv H_{\text{ADM}} = -\frac{1}{2\pi} \int d^3x \Delta\Psi \left[\dot{z}^i, P_i, \hat{S}_{(i)(j)}, h_{ij}^{\text{TT}}, \pi_{\text{can}}^{ij\text{TT}} \right]$

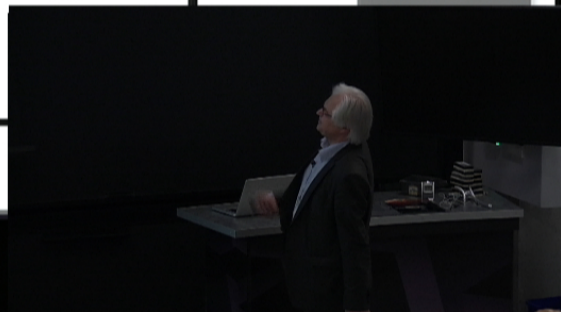
$$\{\dot{z}^i, P_j\} = \delta_{ij}, \quad \{\hat{S}_{(i)}, \hat{S}_{(j)}\} = \epsilon_{ijk} \hat{S}_{(k)}$$

$$\{h_{ij}^{\text{TT}}(\mathbf{x}, t), \pi_{\text{can}}^{kl\text{TT}}(\mathbf{x}', t)\} = 16\pi \delta_{ij}^{\text{TT}kl} \delta(\mathbf{x} - \mathbf{x}')$$

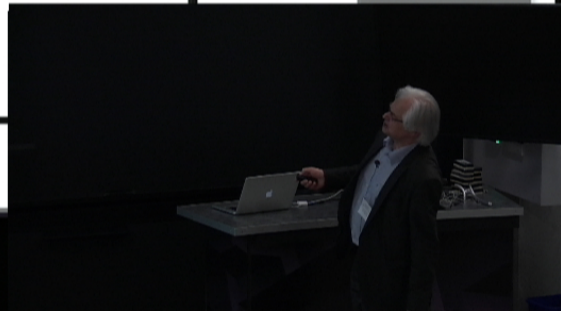


Comparison with Results from EFT (selective)

in this section: -2 signature



$$\begin{aligned}
 L_{\text{SO(L)}}^{\text{NLO}} = & \frac{Gm_2}{r_{12}^2} \mathbf{S}_1 \cdot \left[\mathbf{v}_1 \times \mathbf{n}_{12} \left(\frac{1}{2} \mathbf{v} \cdot \mathbf{v}_2 - \frac{1}{2} v_2^2 - \frac{3}{2} (\mathbf{v}_1 \cdot \mathbf{n}_{12})(\mathbf{v}_2 \cdot \mathbf{n}_{12}) \right) \right. \\
 & + \mathbf{v}_2 \times \mathbf{n}_{12} (\mathbf{v}_1 \cdot \mathbf{v}_2 - v_2^2 + 3(\mathbf{v}_1 \cdot \mathbf{n}_{12})(\mathbf{v}_2 \cdot \mathbf{n}_{12})) \\
 & \left. + \mathbf{v}_1 \times \mathbf{v}_2 \left(\frac{1}{2} \mathbf{v}_1 \cdot \mathbf{n}_{12} + \mathbf{v}_2 \cdot \mathbf{n}_{12} \right) \right] \\
 & + \frac{G^2 m_2}{r_{12}^3} \mathbf{S}_1 \cdot \left[\mathbf{v}_1 \times \mathbf{n}_{12} \left(2m_1 - \frac{1}{2} m_2 \right) + \mathbf{v}_2 \times \mathbf{n}_{12} (2m_2) \right] \\
 & + \frac{Gm_2}{r_{12}^2} \left[S_1^{(0)(i)} n_{12}^i \left(1 - \frac{3}{2} \mathbf{v}_1 \cdot \mathbf{v}_2 + \frac{3}{2} v_2^2 - \frac{3}{2} (\mathbf{v}_1 \cdot \mathbf{n}_{12})(\mathbf{v}_2 \cdot \mathbf{n}_{12}) \right) \right. \\
 & \left. + S_1^{(0)(i)} v_2^i \left(-\frac{3}{2} \mathbf{v}_1 \cdot \mathbf{n}_{12} \right) \right] \\
 & + \frac{Gm_2}{r_{12}} \left[\frac{3}{2} \dot{S}_1^{(0)(i)} v_2^i \right] - \frac{G^2 m_2}{r_{12}^3} S_1^{(0)(i)} n_{12}^i [m_1 + 2m_2]
 \end{aligned}$$



$$\begin{aligned}
 L_{\text{SO(PR)}}^{\text{NLO}} = & -\frac{Gm_2}{r_{12}^2} \left[\left\{ S_1^{(i)(0)} \left(1 + 2\mathbf{v}_2^2 - 2\mathbf{v}_1 \cdot \mathbf{v}_2 - \frac{3}{2}(\mathbf{v}_2 \cdot \mathbf{n}_{12})^2 \right. \right. \right. \\
 & - \left. \frac{G}{r} (3m_1 + 2m_2) \right) + \left(1 - \frac{3}{2}(\mathbf{v}_2 \cdot \mathbf{n}_{12})^2 + \frac{G}{2r} (4m_1 - m_2) \right) S_1^{(i)(j)} \mathbf{v}_1^j \\
 & - \left(2 - 2\mathbf{v}_1 \cdot \mathbf{v}_2 - 3(\mathbf{v}_2 \cdot \mathbf{n}_{12})^2 + 2\mathbf{v}_2^2 - \frac{G}{2r} (2m_1 + 5m_2) \right) S_1^{(i)(j)} \mathbf{v}_2^j \left. \right\} \mathbf{n}_{12}^i \\
 & + S_1^{(i)(0)} (\mathbf{v}_1 - \mathbf{v}_2)^i \mathbf{v}_2 \cdot \mathbf{n}_{12} + S_1^{(i)(j)} \mathbf{v}_1^j \mathbf{v}_2^i \mathbf{v}_2 \cdot \mathbf{n}_{12} \left. \right] + (1 \leftrightarrow 2)
 \end{aligned}$$

$$\begin{aligned}
 V_{S_1 S_2}^{\text{NLO}(\text{PR})} = & -\frac{G}{r_{12}^3} \left[(\delta^{ij} - 3n_{12}^i n_{12}^j) \left(S_1^{(i)(0)} S_2^{(j)(0)} + \frac{1}{2} \mathbf{v}_1 \cdot \mathbf{v}_2 S_1^{(i)(k)} S_2^{(j)(k)} \right. \right. \\
 & + v_1^m v_2^k S_1^{(i)(k)} S_2^{(j)(m)} - v_1^k v_2^m S_1^{(i)(k)} S_2^{(j)(m)} + S_1^{(i)(0)} S_2^{(j)(k)} (v_2^k - v_1^k) \\
 & \left. \left. + S_1^{(i)(k)} S_2^{(j)(0)} (v_1^k - v_2^k) \right) \right. \\
 & + \frac{1}{2} S_1^{(k)(i)} S_2^{(k)(j)} \left(3\mathbf{v}_1 \cdot \mathbf{n}_{12} \mathbf{v}_2 \cdot \mathbf{n}_{12} (\delta^{ij} - 5n_{12}^i n_{12}^j) \right. \\
 & \left. + 3\mathbf{v}_1 \cdot \mathbf{n}_{12} (v_2^j n_{12}^i + v_2^i n_{12}^j) + 3\mathbf{v}_2 \cdot \mathbf{n}_{12} (v_1^j n_{12}^i + v_1^i n_{12}^j) - v_1^i v_2^j - v_2^i v_1^j \right) \\
 & \left. + (3n_{12}^l \mathbf{v}_2 \cdot \mathbf{n}_{12} - v_2^l) S_1^{(0)(k)} S_2^{(k)(l)} + (3n_{12}^l \mathbf{v}_1 \cdot \mathbf{n}_{12} - v_1^l) S_2^{(0)(k)} S_1^{(k)(l)} \right] \\
 & + \left(\frac{G}{r_{12}^3} - \frac{3(m_1 + m_2)G^2}{r_{12}^4} \right) S_1^{(j)(k)} S_2^{(j)(i)} (\delta^{ki} - 3n_{12}^k n_{12}^i) \\
 & + \frac{Gm_2}{r_{12}^2} n_{12}^j S_1^{(j)(0)} - \frac{Gm_1}{r_{12}^2} n_{12}^j S_2^{(j)(0)}
 \end{aligned}$$

$$p_\mu = mg_{\mu\nu} \frac{u^\nu}{\sqrt{u_\lambda u^\lambda}}, \quad g^{\mu\nu} p_\mu p_\nu = m^2, \quad u^0 = 1$$

$$\begin{aligned} \{z_I^i, p_{Jj}\} &= \delta_{ij} \delta_{IJ} \\ \{S_{Iab}, S_{Icd}\} &= S_{Ica} \eta_{bd} - S_{Ida} \eta_{bc} + S_{Idb} \eta_{ac} - S_{Icb} \eta_{ad} \\ \{S_{Iab}, \Lambda_{Ic}^A\} &= \eta_{bc} \Lambda_{Ia}^A - \eta_{ac} \Lambda_{Ib}^A \end{aligned}$$

remaining supplementary conditions:

$$\begin{aligned} S_{ab} u^b &= 0 \Leftrightarrow S_{(i)b} u^b = 0 \Leftrightarrow S_{(i)(0)} + S_{(i)(j)} \frac{u^{(j)}}{u^{(0)}} = 0 \\ \Lambda^{[i]a} u_a &= 0 \Leftrightarrow \Lambda^{[i](0)} + \Lambda^{[i](j)} \frac{u^{(j)}}{u^{(0)}} = 0 \\ \Lambda_{[0]a} &= \frac{u_a}{\sqrt{u_b u^b}} \end{aligned}$$

$$\{\Phi_A, \Phi_B\} \equiv C_{AB}$$

$$\{Q_1, Q_2\}_D = \{Q_1, Q_2\} - \{Q_1, \Phi_A\} C^{AB} \{\Phi_B, Q_2\}$$

$$\dot{z}_I^i = \{z_I^i, H_{\text{eff}}\}_D, \quad \dot{p}_{Ii} = \{p_{Ii}, H_{\text{eff}}\}_D, \quad \dot{S}_{I(i)(j)} = \{S_{I(i)(j)}, H_{\text{eff}}\}_D$$

aimed at variables:

$$\{\hat{z}_I^i, \hat{p}_{Jj}\}_D = \delta_j^i \delta_{IJ}$$

$$\{\hat{\Lambda}_I^{[i](j)}, \hat{S}_I^{(m)(n)}\}_D = -\delta_{jm} \hat{\Lambda}_I^{[i](n)} + \delta_{jn} \hat{\Lambda}_I^{[i](m)}$$

$$\{\hat{S}_I^{(i)(j)}, \hat{S}_I^{(k)(l)}\}_D = \delta_{jl} \hat{S}_I^{(i)(k)} - \delta_{jk} \hat{S}_I^{(i)(l)} + \delta_{ik} \hat{S}_I^{(j)(l)} - \delta_{il} \hat{S}_I^{(j)(k)}$$

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$$
$$u_I^a = e^a_{\mu}(z_I) u_I^\mu = \eta^{a\mu} \left(\eta_{\mu\nu} + \frac{1}{2} h_{\mu\nu} - \frac{1}{8} h_{\mu\rho} \eta^{\rho\sigma} h_{\sigma\nu} + \dots \right) u_I^\nu$$
$$e_{(0)0} = 1 + \frac{1}{2} h_{00} = 1 - G \frac{m_1}{r_1} - G \frac{m_2}{r_2} + \dots$$
$$e_{(0)i} = e_{(i)0} = \frac{1}{2} h_{0i} = 2Gv_1^i \frac{m_1}{r_1} - GS_1^{(i)(j)} \frac{n_1^j}{r_1^2} + (1 \leftrightarrow 2) + \dots$$
$$e_{(i)j} = -\delta_{ij} + \frac{1}{2} h_{ij} = -\delta_{ij} \left(1 + G \frac{m_1}{r_1} + G \frac{m_2}{r_2} \right) + \dots$$

$$0 = \Phi_A$$

$$\begin{aligned}
 0 &= \Lambda_1^{[0](i)} - \frac{u_1^{(i)}}{\sqrt{(u_1^{(0)})^2 - (u_1^{(i)})^2}} \\
 &= \Lambda_1^{[0](i)} - (1 - v_1^2)^{-\frac{1}{2}} v_1^i - 2G(v_1^i - v_2^i) \frac{m_2}{r_{12}} - GS_2^{ij} \frac{n_{12}^j}{r_{12}^2} + \dots \\
 &= \Lambda_1^{[0](i)} - \frac{p_{1i}}{m_1} - \frac{G}{2r_{12}} \left(3p_{2i} - 2\frac{m_2}{m_1} p_{1i} + (\mathbf{n}_{12} \cdot \mathbf{p}_2) n_{12}^i \right) \\
 &+ \frac{Gn_{12}^j}{r_{12}^2} \left(S_2^{(i)(j)} + \frac{m_2}{m_1} S_1^{(i)(j)} \right) + \dots
 \end{aligned}$$

$$0 = \Phi_A$$

$$\begin{aligned}
 0 &= \Lambda_1^{[0](i)} - \frac{u_1^{(i)}}{\sqrt{(u_1^{(0)})^2 - (u_1^{(i)})^2}} \\
 &= \Lambda_1^{[0](i)} - (1 - v_1^2)^{-\frac{1}{2}} v_1^i - 2G(v_1^i - v_2^i) \frac{m_2}{r_{12}} - GS_2^{ij} \frac{n_{12}^j}{r_{12}^2} + \dots \\
 &= \Lambda_1^{[0](i)} - \frac{p_{1i}}{m_1} - \frac{G}{2r_{12}} \left(3p_{2i} - 2\frac{m_2}{m_1} p_{1i} + (\mathbf{n}_{12} \cdot \mathbf{p}_2) n_{12}^i \right) \\
 &+ \frac{Gn_{12}^j}{r_{12}^2} \left(S_2^{(i)(j)} + \frac{m_2}{m_1} S_1^{(i)(j)} \right) + \dots
 \end{aligned}$$

$$\mathbf{p}_I = \hat{\mathbf{p}}_I$$

$$\begin{aligned}
 z_1^i = \hat{z}_1^i &- \left[\frac{1}{2m_1^2} p_{1k} \hat{S}_{1(i)(k)} \left(1 - \frac{p_1^2}{4m_1^2} \right) - G \frac{m_2}{m_1^2} \frac{p_{1k} \hat{S}_{1(i)(k)}}{\hat{r}_{12}} \right. \\
 &+ \frac{3}{2} G \frac{p_{2k} \hat{S}_{1(i)(k)}}{m_1 \hat{r}_{12}} \\
 &+ \frac{G}{2} \frac{\hat{n}_{12}^k (\hat{\mathbf{n}}_{12} \cdot \mathbf{p}_2) \hat{S}_{1(i)(k)}}{m_1 \hat{r}_{12}} + G \frac{m_2}{m_1^2} \frac{\hat{S}_{1(k)(l)} \hat{S}_{1(i)(l)} \hat{n}_{12}^k}{\hat{r}_{12}^2} \\
 &\left. + G \frac{\hat{n}_{12}^k \hat{S}_{1(i)(l)} \hat{S}_{2(k)(l)}}{m_1 \hat{r}_{12}^2} \right] + \dots
 \end{aligned}$$

$$\begin{aligned}
 S_{1(i)(j)} &= \hat{S}_{1(i)(j)} - \left[\frac{p_{1[i\hat{S}_1(j)](k)} p_{1k}}{m_1^2} \left(1 - \frac{p_1^2}{4m_1^2} \right) \right. \\
 &- \frac{2Gm_2}{m_1^2 \hat{r}_{12}} p_{1[i\hat{S}_1(j)](k)} p_{1k} \\
 &+ \frac{3G}{m_1 \hat{r}_{12}} p_{1[i\hat{S}_1(j)](k)} p_{2k} + \frac{G}{m_1 \hat{r}_{12}} p_{1[i\hat{S}_1(j)](k)} \hat{n}_{12}^k (\hat{n}_{12} \cdot p_2) \\
 &+ \frac{2Gm_2}{m_1^2 \hat{r}_{12}^2} p_{1[i\hat{S}_1(j)](l)} \hat{S}_{1(k)(l)} \hat{n}_{12}^k \\
 &\left. + \frac{2G}{m_1 \hat{r}_{12}^2} p_{1[i\hat{S}_1(j)](l)} \hat{S}_{2(k)(l)} \hat{n}_{12}^k \right] + \dots
 \end{aligned}$$

Hergt/Steinhoff/GS: arXiv:1110.2094

Dirac Delta Functions and Black Holes



independent field variables

$$g_{ij} = \left(1 + \frac{1}{8}\phi\right)^4 \delta_{ij} + h_{ij}^{\text{TT}}$$

$$\pi^{ii} = 0, \quad \pi^{ij} = -\gamma^{1/2}(K^{ij} - \gamma^{ij}K), \quad \pi_i^i = \pi^{ij}h_{ij}^{\text{TT}}$$

unique decomposition: $\pi^{ij} = \tilde{\pi}^{ij} + \pi^{ij\text{TT}}$

$$\tilde{\pi}^{ij} = \partial_i \pi^j + \partial_j \pi^i - \frac{2}{3}\delta_{ij}\partial_k \pi^k$$

$\pi^{ij\text{TT}}c^3/16\pi G$: canonical conjugate to h_{ij}^{TT}



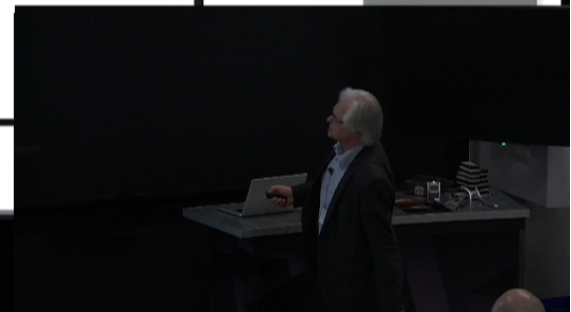
Hamilton and momentum constraints

$$g^{1/2}R = \frac{1}{g^{1/2}} \left(\pi_j^i \pi_i^j - \frac{1}{2} \pi_i^i \pi_j^j \right) + \frac{16\pi G}{c^3} \sum_a (m_a^2 c^2 + \gamma^{ij} p_{ai} p_{aj})^{1/2} \delta_a$$

$$\sqrt{-g}G^{00} = \frac{8\pi G}{c^4} \sqrt{-g}T^{00}$$

$$-2\partial_j \pi_i^j + \pi^{kl} \partial_i g_{kl} = \frac{16\pi G}{c^3} \sum_a p_{ai} \delta_a$$

$$\sqrt{-g}G_i^0 = \frac{8\pi G}{c^4} \sqrt{-g}T_i^0$$

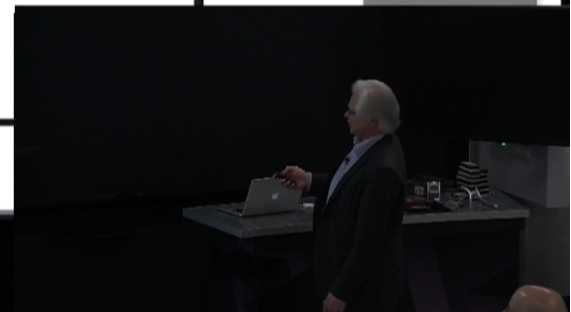


isolated BH

$$\begin{aligned} ds^2 &= - \left(\frac{1 - \frac{Gm}{2rc^2}}{1 + \frac{Gm}{2rc^2}} \right)^2 c^2 dt^2 + \left(1 + \frac{Gm}{2rc^2} \right)^4 \delta_{ij} dx^i dx^j \\ &= - \left(\frac{1 - \frac{Gm}{2Rc^2}}{1 + \frac{Gm}{2Rc^2}} \right)^2 c^2 dt^2 + \left(1 + \frac{Gm}{2Rc^2} \right)^4 \delta_{ij} dX^i dX^j \end{aligned}$$

symmetry transformation (inversion): $Rr = \left(\frac{Gm}{2c^2}\right)^2$

$$R^2 = X^i X^i, \quad r^2 = x^i x^i$$



Brill-Lindquist BHs

maximally sliced

$$ds^2 = - \left(\frac{1 - \frac{\beta_1 G}{2r_1 c^2} - \frac{\beta_2 G}{2r_2 c^2}}{1 + \frac{\alpha_1 G}{2r_1 c^2} + \frac{\alpha_2 G}{2r_2 c^2}} \right)^2 c^2 dt^2 + \left(1 + \frac{\alpha_1 G}{2r_1 c^2} + \frac{\alpha_2 G}{2r_2 c^2} \right)^4 d\mathbf{x}^2$$

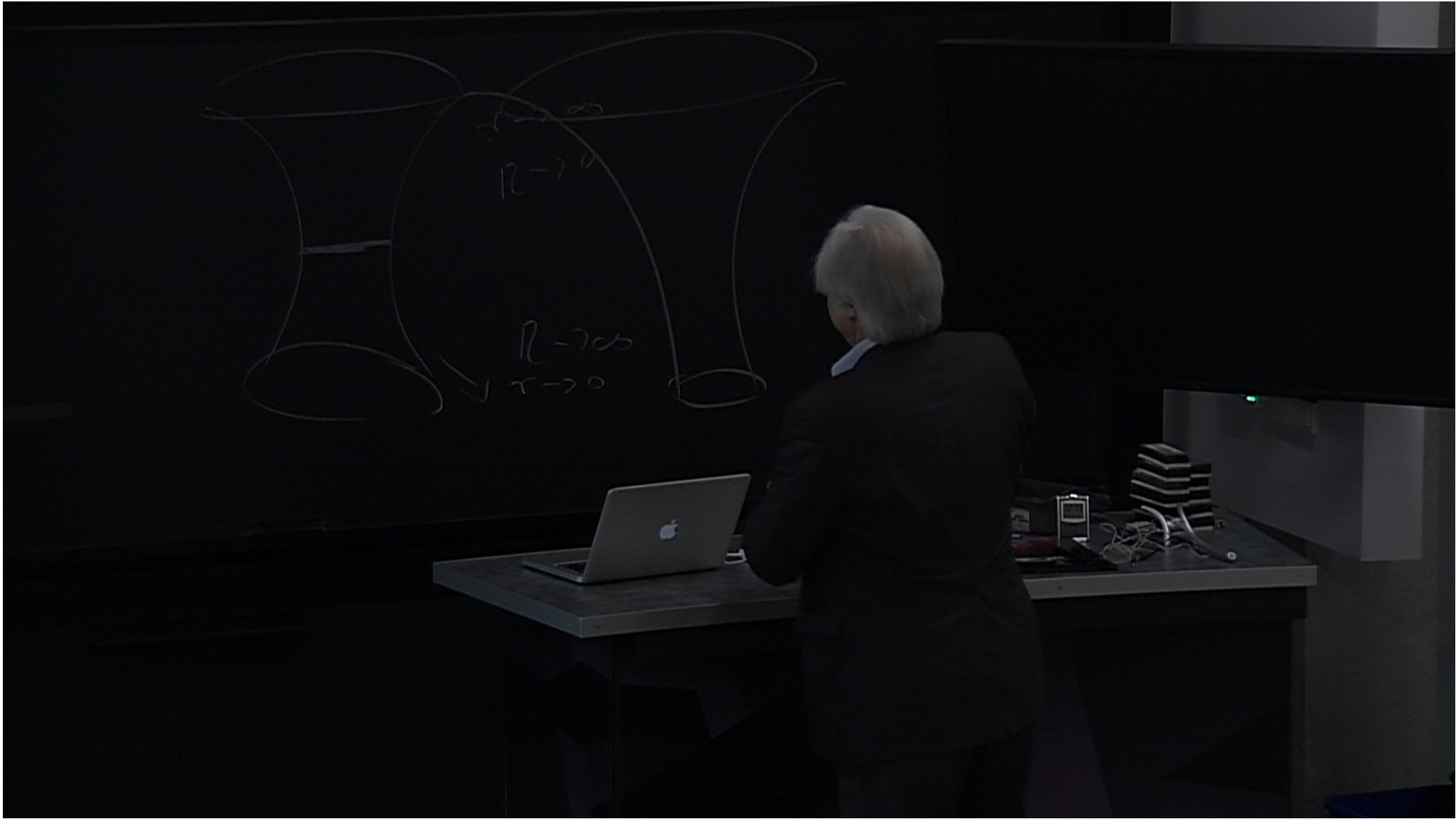
total energy: $E_{ADM} = -\frac{c^4}{2\pi G} \oint_{i_0} ds_i \partial_i \Psi = (\alpha_1 + \alpha_2) c^2$

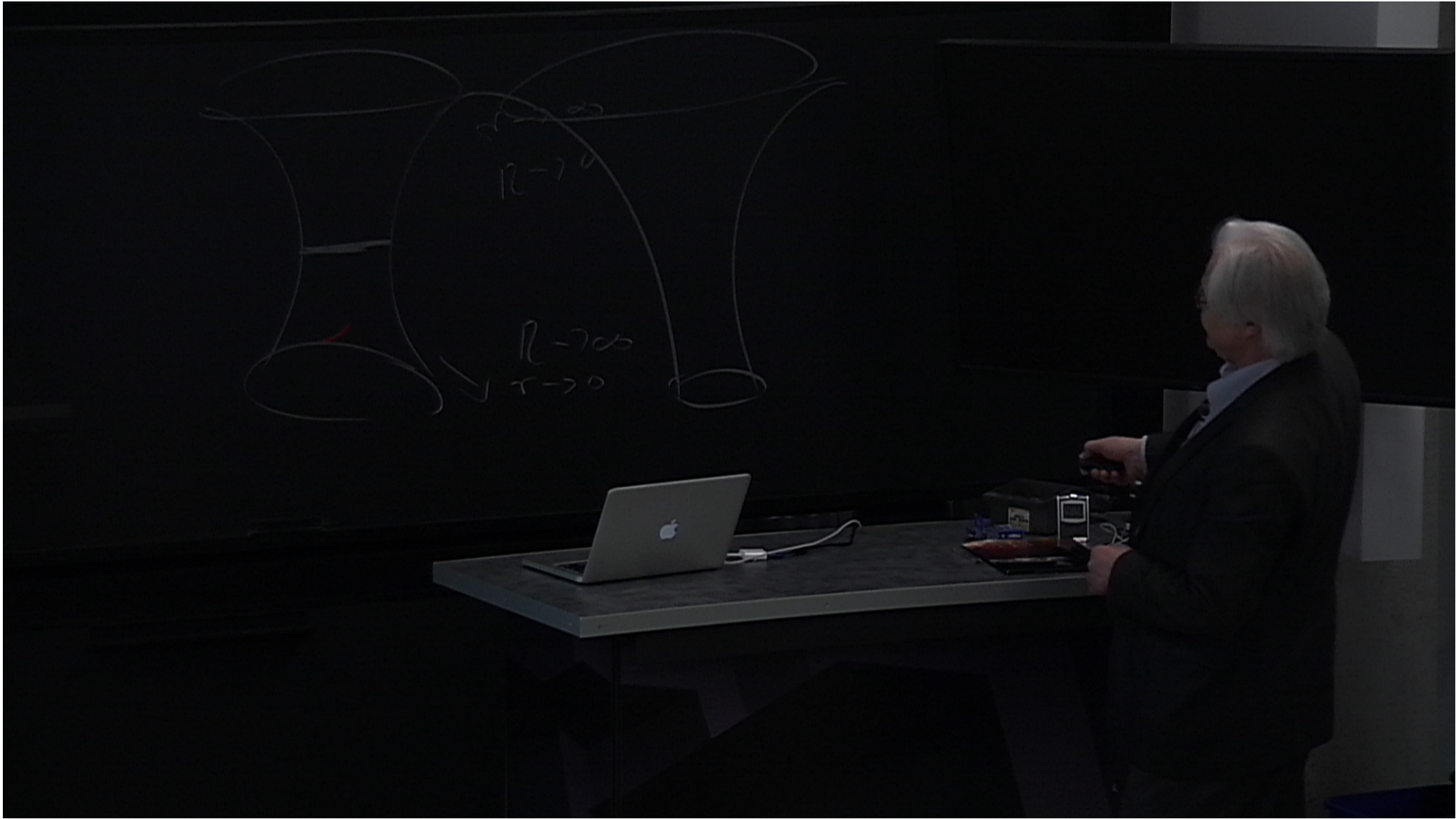
$$\Psi = 1 + \frac{\alpha_1 G}{2r_1 c^2} + \frac{\alpha_2 G}{2r_2 c^2}$$

inversion map of the three-metric at the throat of black hole 1

$$\mathbf{r}'_1 = \mathbf{r}_1 \alpha_1^2 G^2 / 4c^4 r_1^2$$

$$\mathbf{r}'_1 = \mathbf{x}' - \mathbf{x}_1, \quad \mathbf{r}_1 = \mathbf{x} - \mathbf{x}_1, \quad r_1 = |\mathbf{x} - \mathbf{x}_1|$$





$$dl^2 = \Psi^4 d\mathbf{X}^2 = \left(1 + \frac{\alpha_1 G}{2r_1 c^2} + \frac{\alpha_2 G}{2r_2 c^2}\right)^4 d\mathbf{X}^2$$

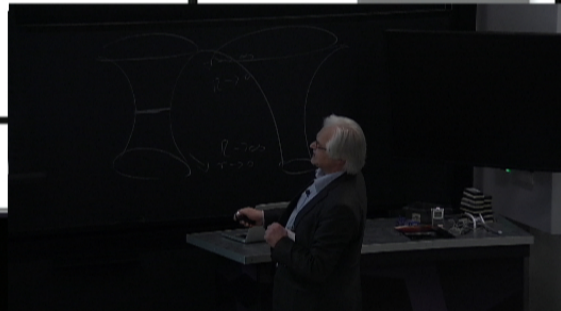
$$dl^2 = \Psi'^4 d\mathbf{X}'^2 = \left(1 + \frac{\alpha_1 G}{2r_1' c^2} + \frac{\alpha_1 \alpha_2 G^2}{4r_2 r_1' c^4}\right)^4 d\mathbf{X}'^2$$

$$\mathbf{r}_2 = \frac{\alpha_1^2 G^2}{4c^4} \frac{\mathbf{r}_1'}{r_1'^2} + \mathbf{r}_{12}, \quad \mathbf{r}_{12} = \mathbf{r}_1 - \mathbf{r}_2$$

$$m_1 = -\frac{c^2}{2\pi G} \oint_{i_0} ds'_i \partial'_i \Psi' = \alpha_1 + \frac{\alpha_1 \alpha_2 G}{2r_{12} c^2}$$

$$\Psi' = 1 + \frac{\alpha_1 G}{2r_1' c^2} + \frac{\alpha_1 \alpha_2 G^2}{4r_2 r_1' c^4}$$

$$-\left(1 + \frac{1}{8} \phi\right) \Delta\phi = \frac{16\pi G}{c^2} \sum_a m_a \delta_a \quad (h_{ij}^{\text{TT}} = \pi^{ij\text{TT}} = p_{ai} = 0)$$
$$\phi = \frac{4G}{c^2} \left(\frac{\alpha_1}{r_1} + \frac{\alpha_2}{r_2} \right)$$
$$\alpha_a = m_a - \frac{m_a + m_b}{2} + \frac{c^2 r_{ab}}{G} \left(\sqrt{1 + \frac{m_a + m_b}{c^2 r_{ab}/G} + \left(\frac{m_a - m_b}{2c^2 r_{ab}/G} \right)^2} - 1 \right)$$
$$H_{\text{BL}} = (\alpha_1 + \alpha_2) c^2 = (m_1 + m_2) c^2 - G \frac{\alpha_1 \alpha_2}{r_{12}}$$



Metric in d-dimensional conformally flat space:

$$g_{ij} = \left(1 + \frac{1}{4} \frac{d-2}{d-1} \phi \right)^{\frac{4}{d-2}} \delta_{ij}$$

$$\phi = \frac{4G}{c^2} \frac{\Gamma(\frac{d-2}{2})}{\pi^{\frac{d-2}{2}}} \left(\frac{\alpha_1}{r_1^{d-2}} + \frac{\alpha_2}{r_2^{d-2}} \right)$$

$$\Psi = 1 + \frac{1}{4} \frac{d-2}{d-1} \phi$$

$$-\Delta^{-1}\delta = \frac{\Gamma((d-2)/2)}{4\pi^{d/2}} r^{2-d}$$
$$\Psi = 1 + \frac{G(d-2)\Gamma((d-2)/2)}{c^2(d-1)\pi^{(d-2)/2}} \left(\frac{\alpha_1}{r_1^{d-2}} + \frac{\alpha_2}{r_2^{d-2}} \right)$$
$$\left(1 + \frac{G(d-2)\Gamma((d-2)/2)}{c^2(d-1)\pi^{(d-2)/2}} \left(\frac{\alpha_1}{r_1^{d-2}} + \frac{\alpha_2}{r_2^{d-2}} \right) \right) \alpha_1 \delta_1 = m_1 \delta_1$$

$1 < d < 2$

$$\left(1 + \frac{G(d-2)\Gamma((d-2)/2)}{c^2(d-1)\pi^{(d-2)/2}} \frac{\alpha_2}{r_{12}^{d-2}} \right) \alpha_1 \delta_1 = m_1 \delta_1$$

