

Title: Tensor Networks and TQFTs

Date: Oct 24, 2011 11:30 AM

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Abstract: In this talk, I will present a first principle construction of a holographic dual for gauged matrix models that include gauge theories. The dual theory is shown to be a closed string field theory coupled with an emergent two-form gauge field defined in one higher dimensional space. The bulk space with an extra dimension emerges as a well defined classical background only when the two-form gauge field is in the deconfinement phase. Based on this, it is shown that critical phases that admit holographic descriptions form a novel universality class with a non-trivial quantum order.

Tensor-Networks(TNs) & Topological Quantum Field Theories(TQFTs)

Zheng-cheng Gu (KITP)

Collaborators:

Part I(Duality):

Prof. M. Leivn(U. of Maryland)

Part II(TQFTs):

Prof. Xiao-gang Wen (MIT)

Dr. Xie Chen(MIT)

Prof. Zhenghan Wang (Station-q)

Dr. Zhen-Xin Liu(Tsinghua U.)

PI. Oct. 2011



Outline

- Intrinsic Topological(IT) order and TQFTs
- Symmetry Protected Topological(SPT) order



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- Symmetry Protected Topological(SPT) order
- Duality between IT and SPT orders
- TNs description of Intrinsic Topological order
- TQFTs of SPT order



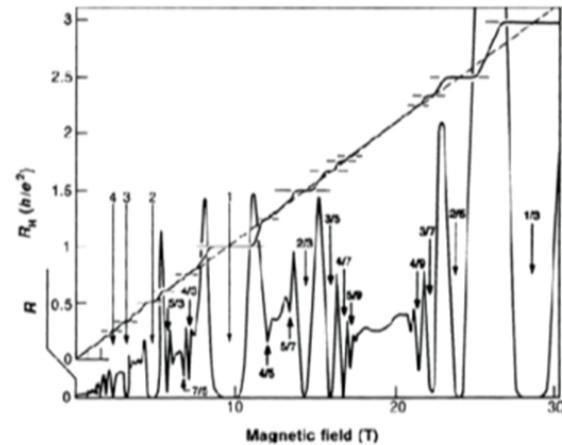
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- Symmetry Protected Topological(SPT) order
- Duality between IT and SPT orders
- TNs description of Intrinsic Topological order
- TQFTs of SPT order
- Fermionic TQFTs from Grassmann TNs



Topological Phenomena in strongly correlated systems

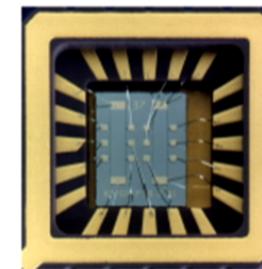
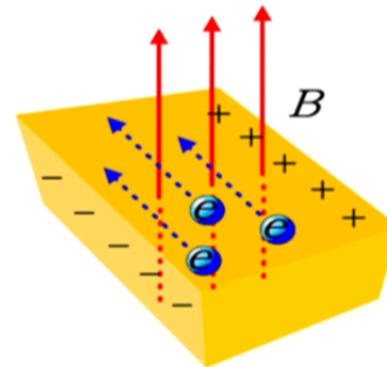
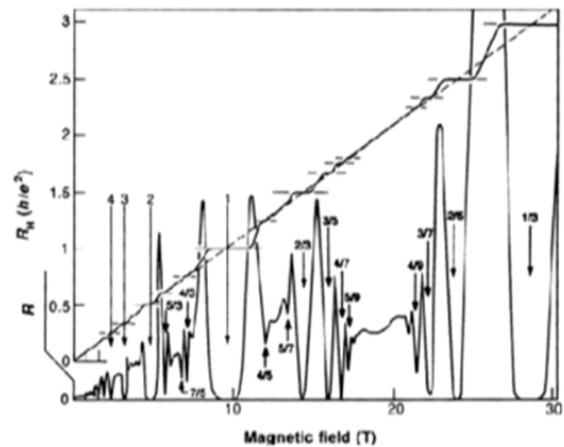
Fractional Quantum Hall Effect D.C.Tsui, *et al* 1982



Topological Phenomena in strongly correlated systems

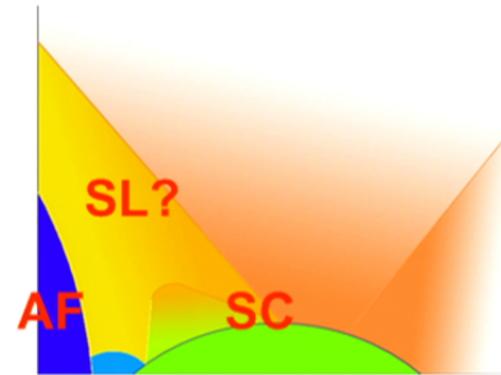
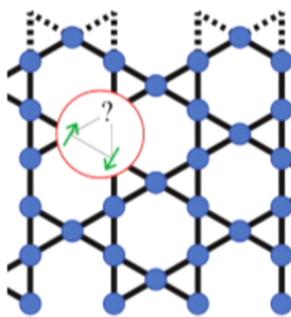
Fractional Quantum Hall Effect

D.C.Tsui, *etal* 1982



Gapped spin liquid

- Frustrated magnets
 - High-Tc cuprates



S. Yan, et al, Science, 2011

New phases of matter: (Intrinsic) Topological order

X.-G. Wen, *et al* 1989



New phases of matter: (Intrinsic) Topological order

X.-G. Wen, *et al* 1989

- Can have the same symmetry as disorder systems.
- No long range correlations.
- Ground state degeneracies depend on the topology of the space.
- Ground state degeneracies are robust against any local perturbations.



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- Can have the same symmetry as disorder systems.
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- Ground state degeneracies are robust against any local perturbations.
- Excitations carry fractional statistics.
- Protected chiral edge states(chiral topological order, e.g. FQHE).



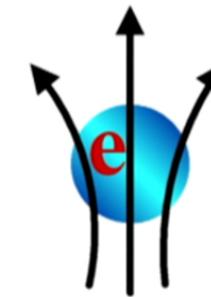
Topological quantum field theories (TQFTs)

- TQFTs are special classes of quantum field theories, whose actions are independent on space-time metric.
- TQFTs are stable fixed point actions. They are very powerful tools to help us understand various topological phenomena.

TQFTs for FQHE

$$\mathcal{L}_{\text{eff}} = \frac{2m+1}{4\pi} \epsilon^{\mu\nu\lambda} a_\mu \partial_\nu a_\lambda$$

S-C. Zhang, *et al* 1989
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Z_2 spin liquid and its TQFT

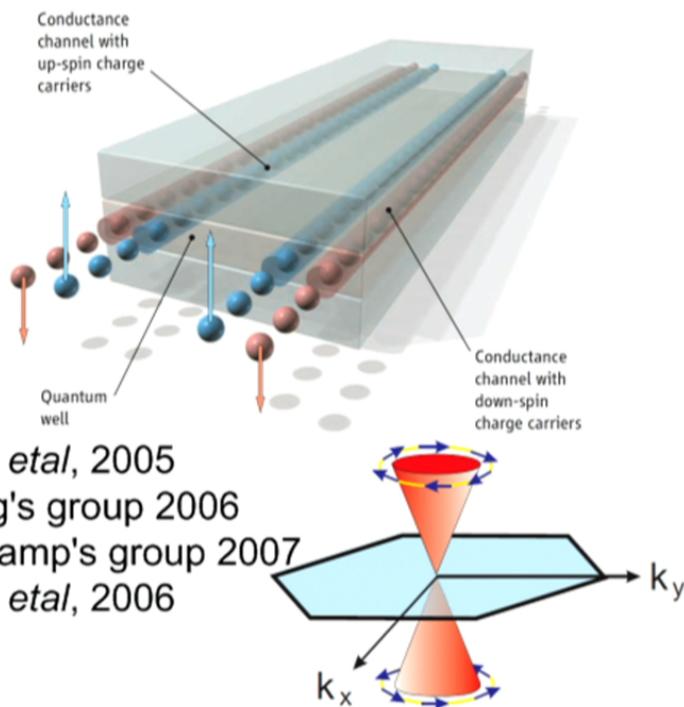
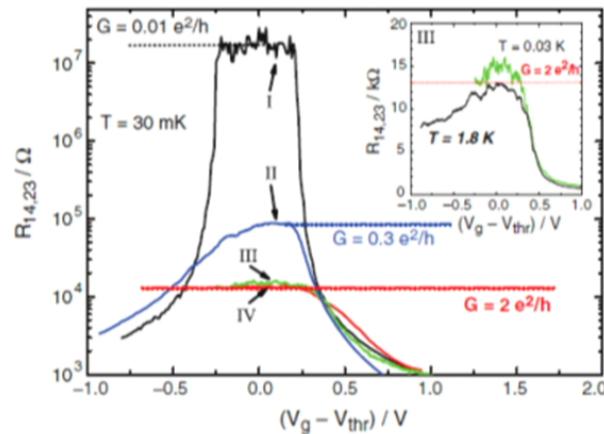
$$|\Psi_{RVB}\rangle = \left| \begin{array}{c} \text{Diagram of a triangular lattice with vertical cylinders representing spins} \end{array} \right\rangle + \left| \begin{array}{c} \text{Diagram of a triangular lattice with horizontal cylinders representing spins} \end{array} \right\rangle + \dots$$

$$L = \frac{1}{4\pi} K_{IJ} a_I \partial_\nu a_J \epsilon^{\mu\nu\lambda}, \quad I, J = 1, 2 \quad K = \begin{pmatrix} 0 & 2 \\ 2 & 0 \end{pmatrix}$$



Symmetry Protected Topological Phenomena

Topological insulator in 2D/3D

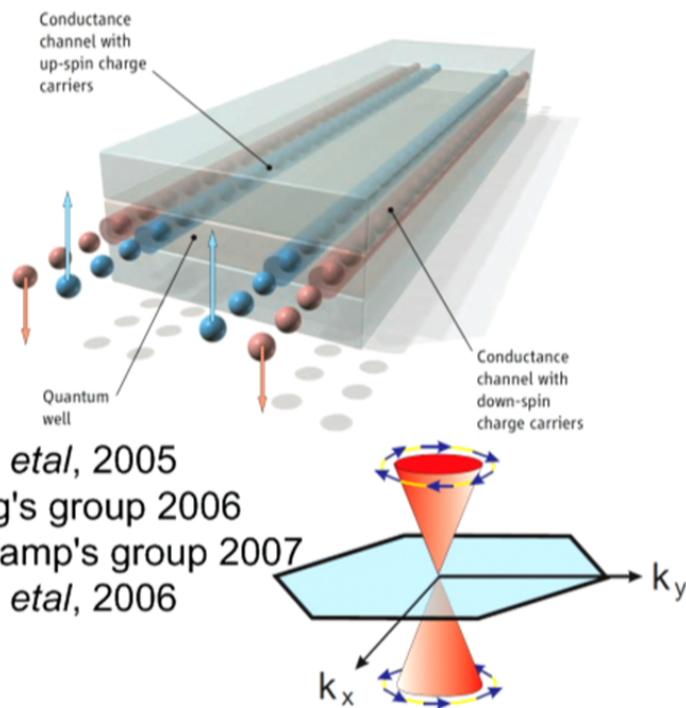
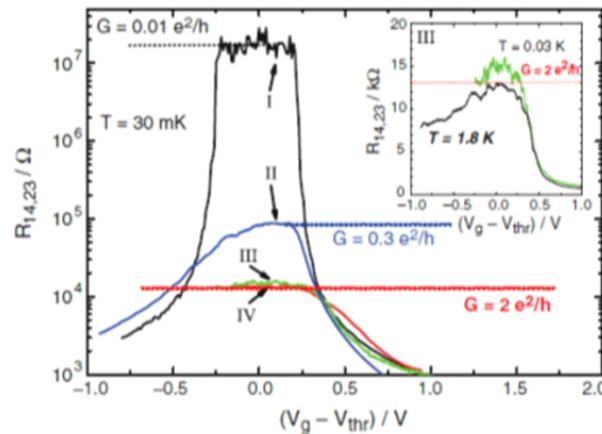


C L Kane, *et al*, 2005
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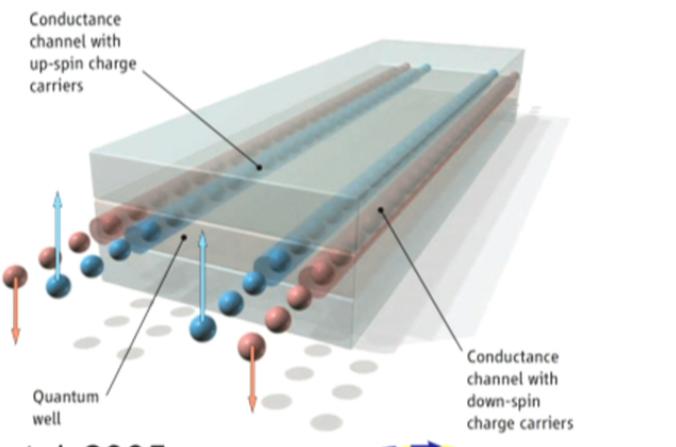
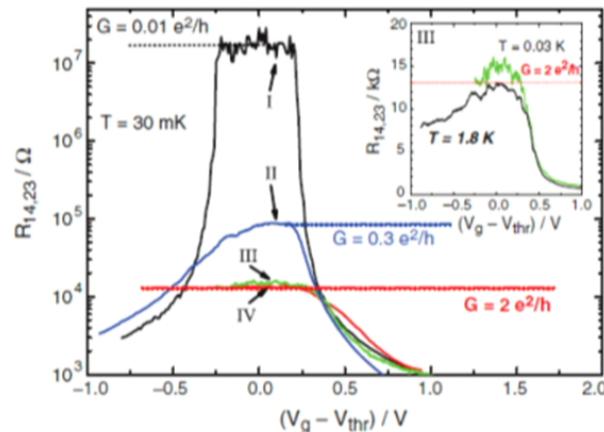


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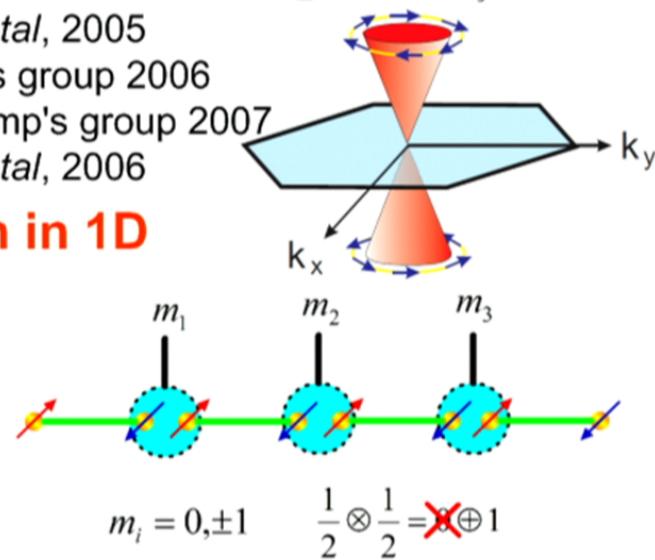


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Haldane phase of spin one chain in 1D

$$H = \sum_i P_2(\mathbf{S}_i + \mathbf{S}_{i+1}) \\ = \sum_i [\frac{1}{2} \mathbf{S}_i \cdot \mathbf{S}_{i+1} + \frac{1}{6} (\mathbf{S}_i \cdot \mathbf{S}_{i+1})^2 + \frac{1}{3}]$$

FDM Haldane, 1983, Ian Affleck, 1987
Z Gu, *et al*, 2009, F Pollmann, *et al*, 2010



New phases of matter: Symmetry Protected Topological Order

- Can have the same symmetry as trivial disorder systems.



New phases of matter: Symmetry Protected Topological Order

- Can have the same symmetry as trivial disorder systems.
- No long range correlations.
- Unique Ground state. Excitations do not carry fractional statistics.
- Indistinguishable with trivial disorder systems if symmetry is broken in bulk
- Protected gapless edge states if symmetry is not (spontaneously or explicitly)broken on the edge



How to understand the symmetry protected nature of SPT orders?



IT

Duality!



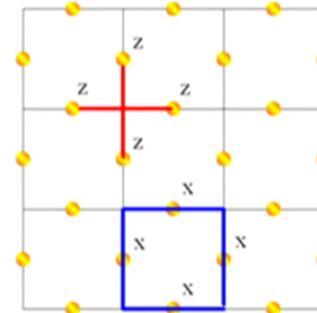
SPT



An exact solvable model of (intrinsic) topological order

toric code model: $H = U \sum_v \left(1 - \prod_{l \in v} \sigma_l^z \right) - t \sum_p \prod_{l \in p} \sigma_l^x$
(Z_2 spin liquid)

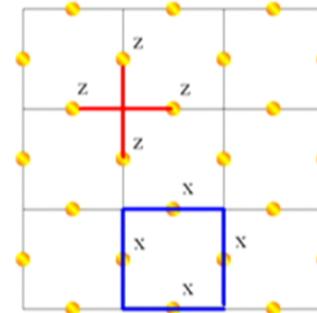
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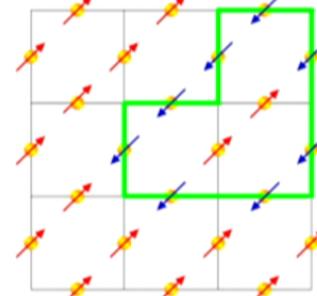
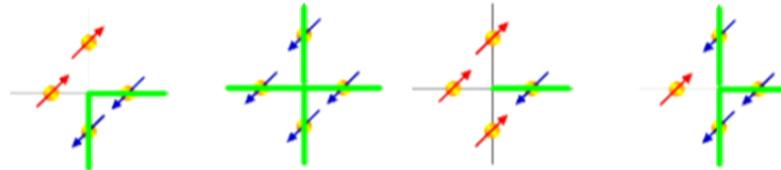


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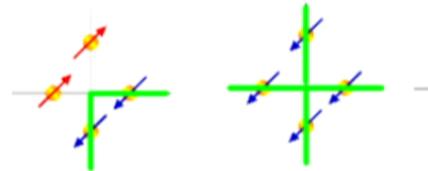
$| \uparrow \rangle \rightarrow$ no string; $| \downarrow \rangle \rightarrow$ one string



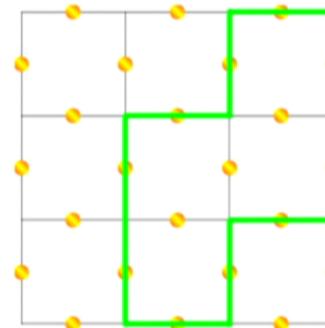
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Kitaev 2003, M. Levin and X.G. Wen 2005



$$|\Psi_{Z_2}\rangle = \sum |X_{\text{close}}\rangle$$



$| \uparrow \rangle$ → no string; $| \downarrow \rangle$ → one string

$$= \begin{array}{c} \text{Diagram 1} \\ \text{Diagram 2} \\ \text{Diagram 3} \end{array} + \cdots$$

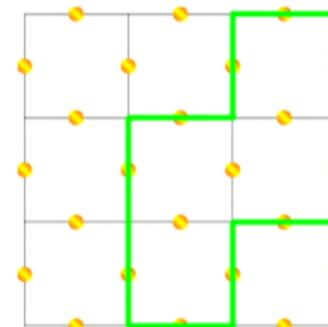
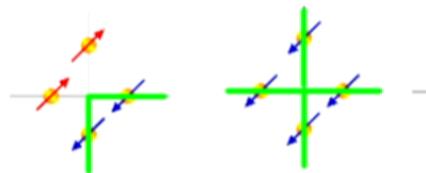


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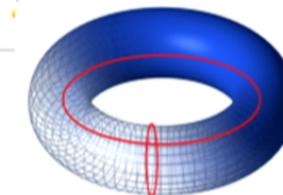
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$$= \begin{array}{c} \text{grid with one string} \\ + \end{array} \begin{array}{c} \text{grid with one string} \\ + \end{array} \begin{array}{c} \text{grid with one string} \\ + \cdots \end{array}$$



- Four fold ground state degeneracies
- Quasi particles carry fractional statistics

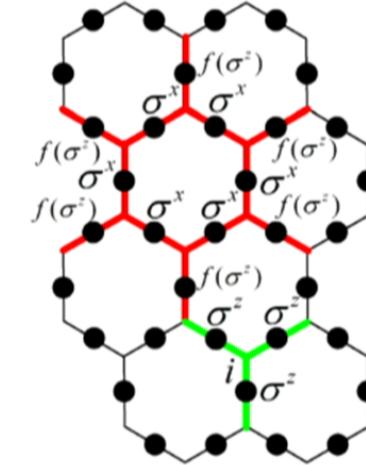


The twisted toric code: the double semion model

M. Levin and X.G. Wen 2005

$$H_{\text{dsemion}} = -U \sum_v \prod_{i \in v} \sigma_i^z - \sum_p \left(\prod_{i \in p} \sigma_i^x \prod_{\text{legs of } p} i^{\frac{1+\sigma_l^z}{2}} \right)$$

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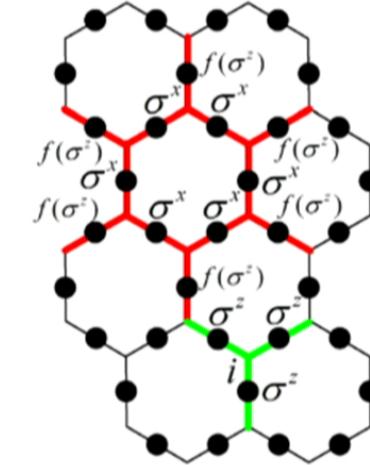


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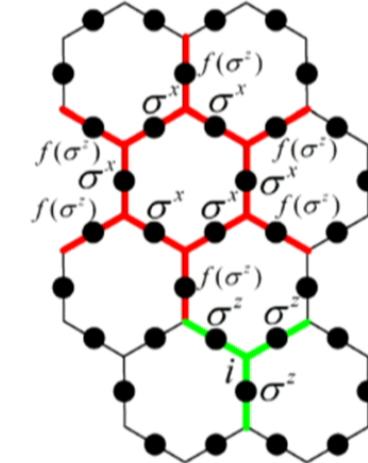
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Double semion model describes different intrinsic topological order

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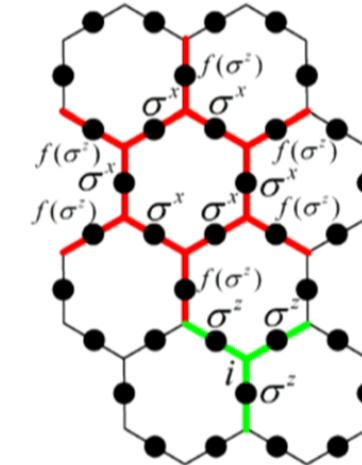
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Quasi-particle in toric code model:

1, e, m, f=em

Quasi-particle in double semion model:

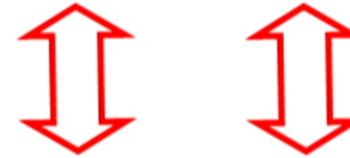
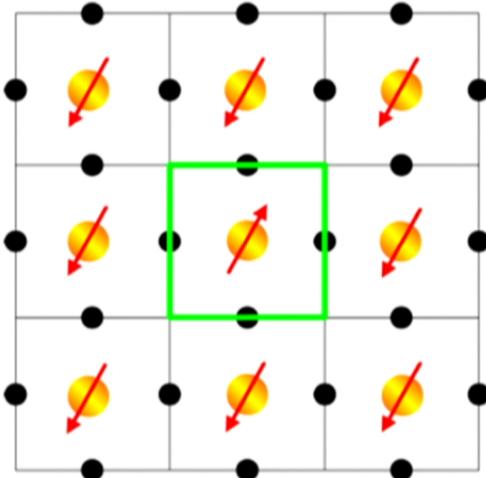
1, s, s, b=ss



Duality between \mathbb{Z}_2 gauge model and Ising model

$$H_{\mathbb{Z}_2} = -U \sum_v \prod_{i \in v} \sigma_i^z - \sum_p \prod_{i \in p} \sigma_i^x \quad \leftrightarrow \quad H_{\text{Ising}} = - \sum_p \sigma_p^x$$

$$H = U \sum_v \left(1 - \prod_{l \in v} \sigma_l^z \right) - t \sum_p \prod_{l \in p} \sigma_l^x - h \sum_l \sigma_z$$



$$H = -t \sum_p \sigma_p^x - h \sum_{\langle pq \rangle} \sigma_p^z \sigma_q^z$$

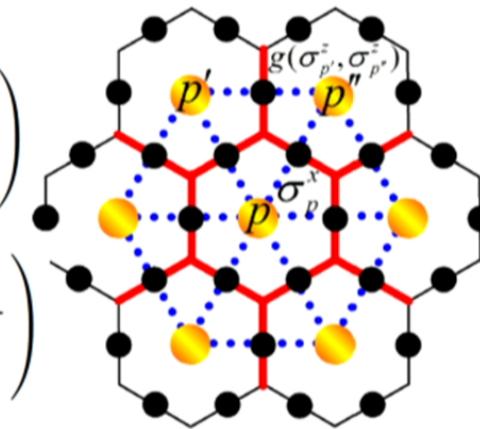
String-net condensation =
domain wall condensation



Dual theory of double semion model

$$H_{\text{dsemion}} = -U \sum_v \prod_{i \in v} \sigma_i^z - \sum_p \left(\prod_{i \in p} \sigma_i^x \prod_{\text{legs of } p} i^{\frac{1+\sigma_i^z}{2}} \right)$$

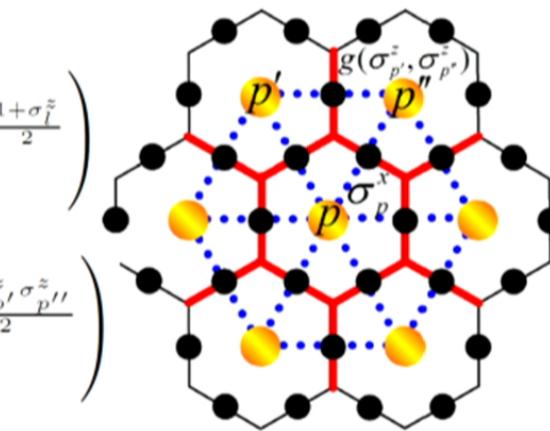
$$H_{\text{twistIsing}} = - \sum_p \tilde{\sigma}_p^x = - \sum_p \left(\sigma_p^x \prod_{\text{sites} \in p} i^{\frac{1+\sigma_{p'}^z \sigma_{p''}^z}{2}} \right)$$



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The dual theory of double semion model
is a SPT ordered phase!

$$H = -\alpha \sum_p \sigma_p^x - (1-\alpha) \sum_p \tilde{\sigma}_p^x$$

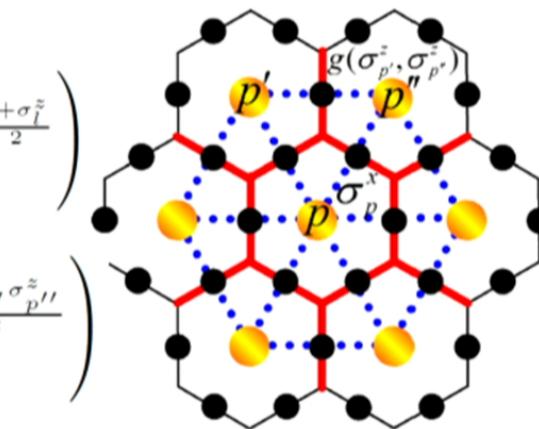
$$\begin{aligned} H' &= -U \sum_v \prod_{i \in v} \sigma_i^z - \alpha \sum_p \prod_{i \in p} \sigma_i^x \\ &\quad - (1-\alpha) \sum_p \left(\prod_{i \in p} \sigma_i^x \prod_{\text{legs of } p} i^{\frac{1+\sigma_i^z}{2}} \right) \end{aligned}$$



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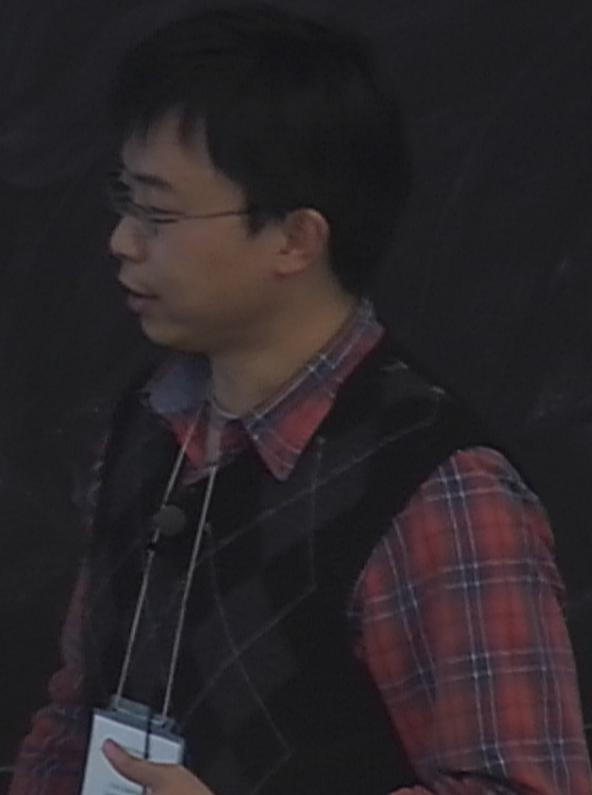
SPT order

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(intrinsic) topological order



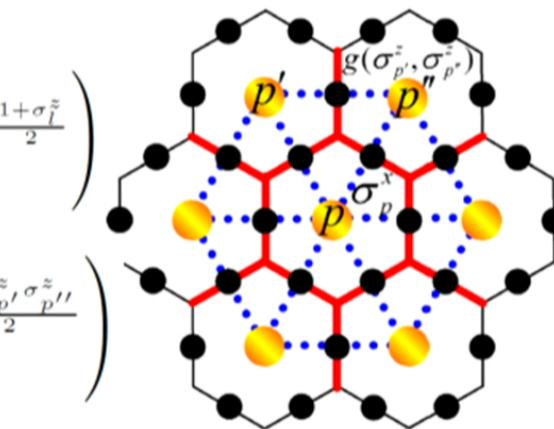
$$\text{I} \quad |\chi_{\text{close}}\rangle,$$
$$\text{II} \quad (-)^{N_G} |\chi_{\text{close}}\rangle$$



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(intrinsic) topological order

- Indeed, the above two models are self dual and the transition point is at 0.5



Edge theory of \mathbb{Z}_2 (bosonic) SPT order

Symmetry Breaking

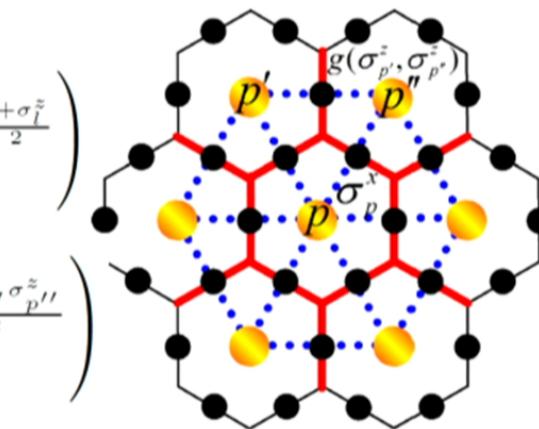
$$H_{\text{edge}} = \sum_p \sigma_p^z \sigma_{p+1}^z$$



Dual theory of double semion model

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(intrinsic) topological order

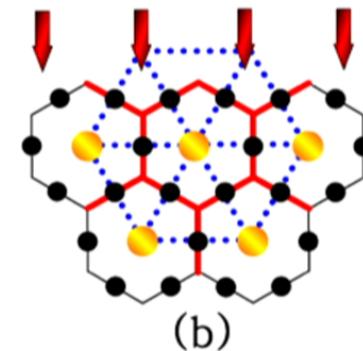
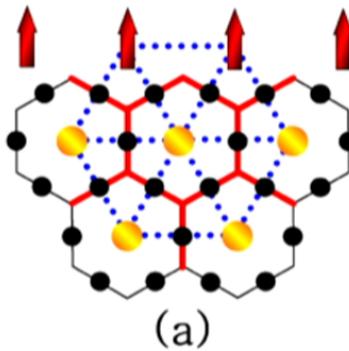
- Indeed, the above two models are self dual and the transition point is at 0.5



Edge theory of \mathbb{Z}_2 (bosonic) SPT order

Symmetry Breaking

$$H_{\text{edge}} = \sum_p \sigma_p^z \sigma_{p+1}^z$$



Gapless

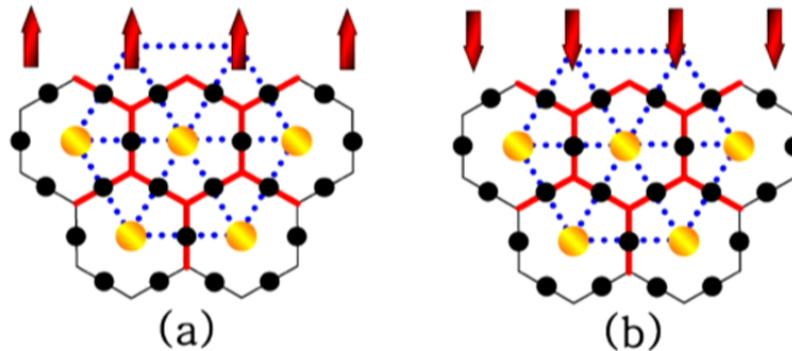
$$H_{\text{edge}} = \sum_p \tilde{\sigma}_p^x(\uparrow) + \tilde{\sigma}_p^x(\downarrow)$$



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Hamiltonian algebra
on edge

$$\begin{aligned} [\tilde{\sigma}_p^x(\uparrow), \tilde{\sigma}_{p'}^x(\uparrow)] &= 0; & [\tilde{\sigma}_p^x(\downarrow), \tilde{\sigma}_{p'}^x(\downarrow)] &= 0; \\ [\tilde{\sigma}_p^x(\uparrow), \tilde{\sigma}_{p'}^x(\downarrow)] &= 0 \quad \text{for } p' \neq p \pm 1; \\ \{\tilde{\sigma}_p^x(\uparrow), \tilde{\sigma}_{p'}^x(\downarrow)\} &= 0 \quad \text{for } p' = p \pm 1, \end{aligned}$$

$$\sigma_p^x(\uparrow) = i\gamma_{2p}\gamma_{2p+1}$$

$$H_{\text{edge}} = i \sum_p (\gamma_{2p}\gamma_{2p+1} + \gamma_{2p-1}\gamma_{2p+2})$$

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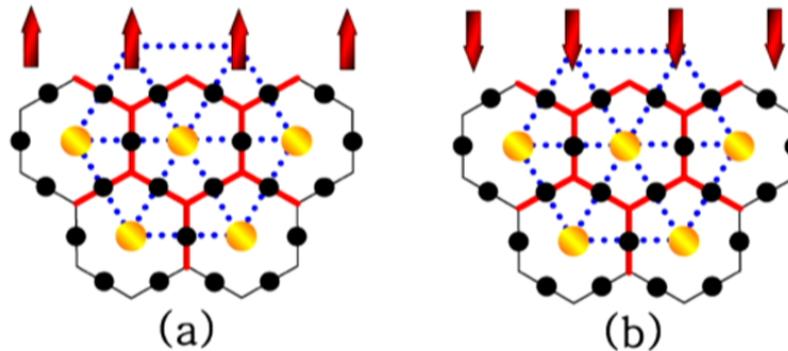
$$E_k = \pm \sqrt{2(1 + 2 \cos 2k)} = \pm 2 \sin k; \quad k > 0$$



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- After bosonization, we find all the relevant perturbations break symmetry.



- What are the TQFTs for SPT orders, especially for systems protected by discrete symmetry?



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- How to understand the topological phenomena in 3D. What are their TQFTs?
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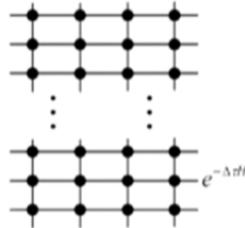
Tensor-Networks(TNs)



TNs representations for partition functions & their renormalization

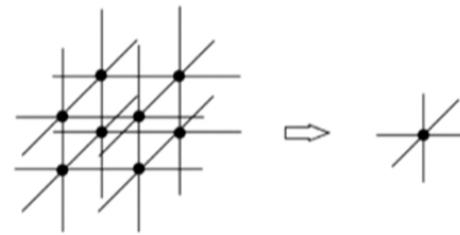
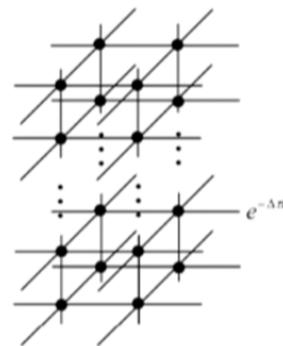
1D

$$Z = \sum_n \langle n | e^{-\beta H} | n \rangle =$$



2D

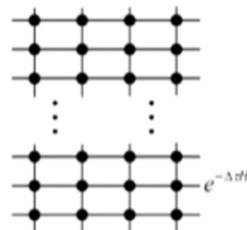
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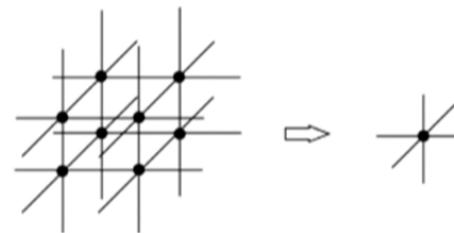
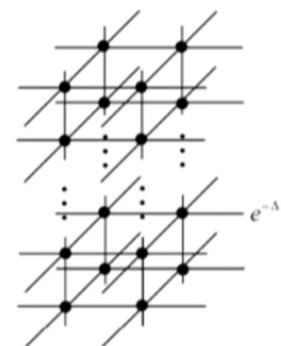
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Fixed point TQFTs
= fixed point tensors

$$\mathcal{L}^{(0)} \rightarrow \mathcal{L}^{(1)} \rightarrow \mathcal{L}^{(2)} \rightarrow \dots \rightarrow \mathcal{L}^{(*)}$$

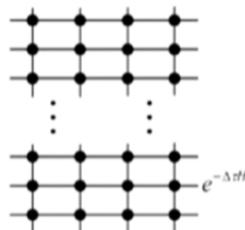
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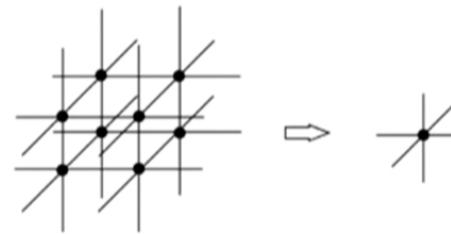
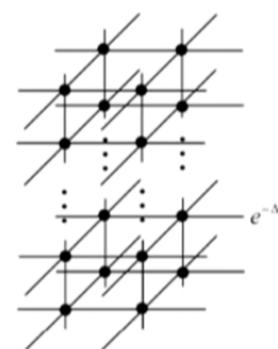
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Topological ordered phases are described by isolate fixed point tensors.

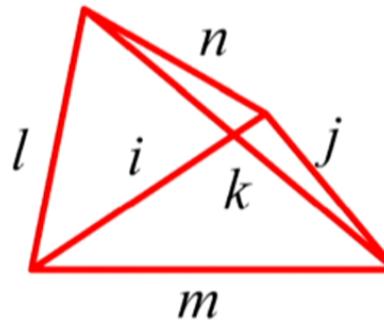


TNs representation for TQFTs of non-chrial (intrinsic) topological order

Partition function of string-net models: Turaev-Viro states sum invariants

$$Z = \frac{1}{D^{N_v}} \sum_{ijklmn \dots \text{ link}} \prod d_i \prod_{\text{tetrahedron}} G_{klm}^{ijm}$$

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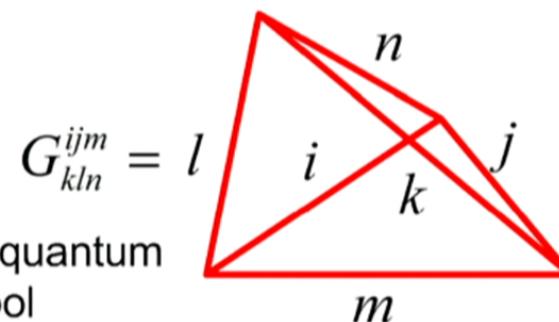
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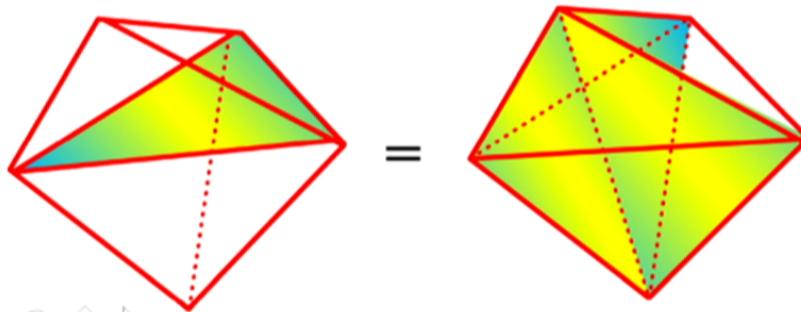
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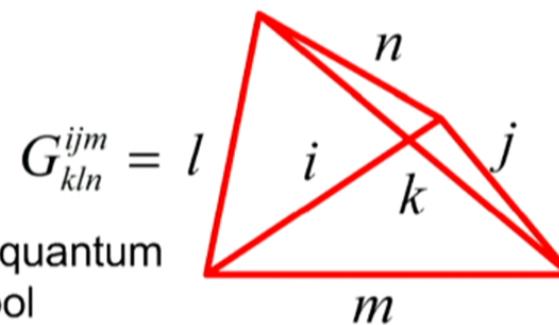


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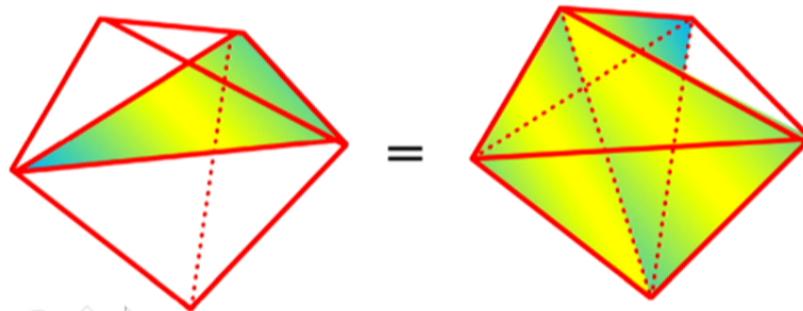
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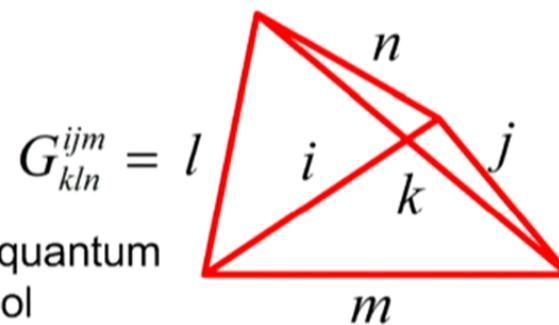


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A diagram illustrating the RG move from a tetrahedron to a pentagon. On the left, a tetrahedron is shown with dashed lines representing hidden edges. An equals sign follows, and on the right, a pentagon is shown with solid red lines.

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TQFTs for \mathbf{Z}_2 bosonic SPT order

$$\tilde{Z} = \frac{1}{D^{N_v}} \sum_{g_0 g_1 g_2 g_3 \dots} \prod_{\text{link}} d_i \prod_{\text{tetrahedron}} \tilde{G}(g_0, g_1, g_2, g_3),$$

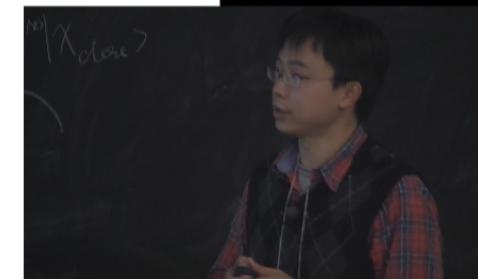
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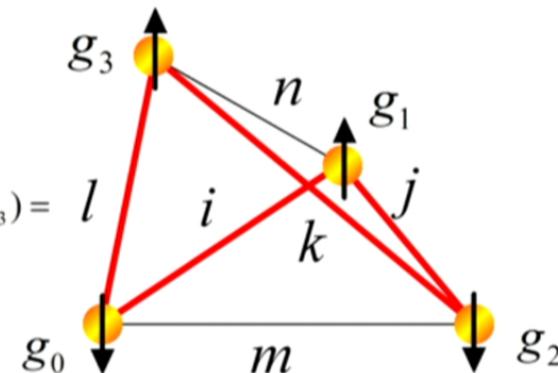
Topological term of (discrete) non-linear sigma model!

Pentagon equation



$$G_{klm}^{ijm} = \tilde{G}(g_0, g_1, g_2, g_3) =$$

Third group co-cycle condition



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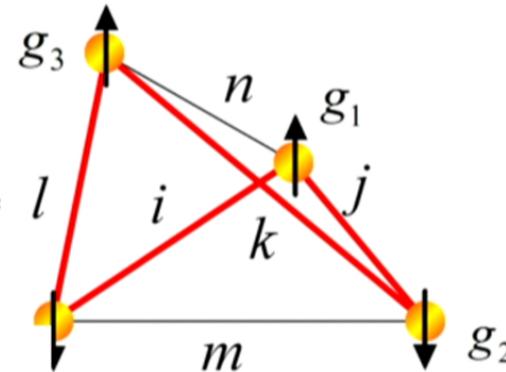
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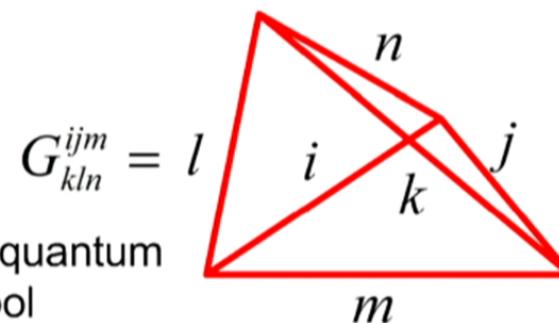
$$\begin{aligned} & \tilde{G}(g_0, g_2, g_3, g_4) \tilde{G}(g_0, g_1, g_2, g_4) \\ &= \tilde{G}(g_1, g_2, g_3, g_4) \tilde{G}(g_0, g_1, g_3, g_4) \tilde{G}(g_0, g_1, g_2, g_3) \quad \tilde{G}(g_0, g_2, g_3, g_4) \sim \nu_3(g_0, g_2, g_3, g_4) \end{aligned}$$



TNs representation for TQFTs of non-chrial (intrinsic) topological order

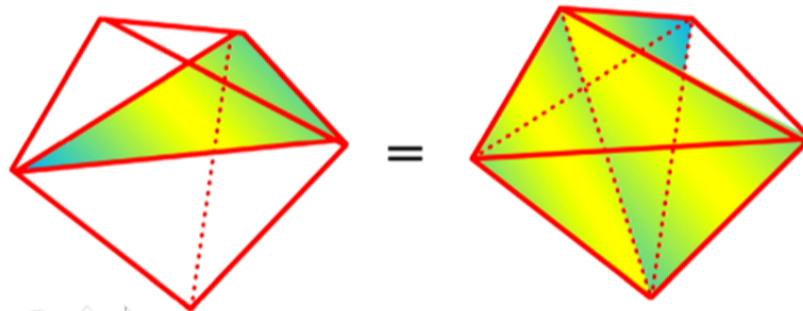
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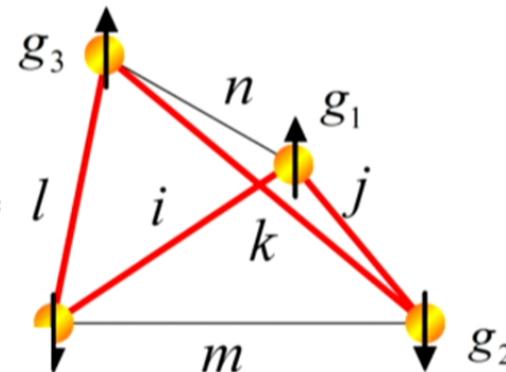
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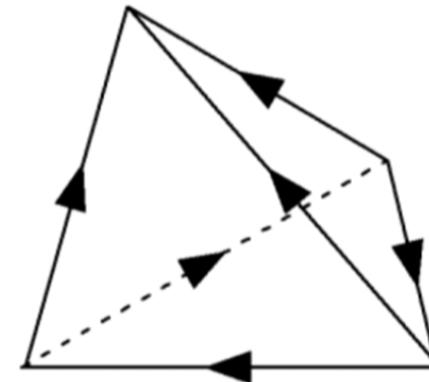
\mathbb{Z}_2 SPT order is classified by $H^3(\mathbb{Z}_2, U(1))$



TQFTs of (bosonic) SPT order in any dimensions with any symmetry

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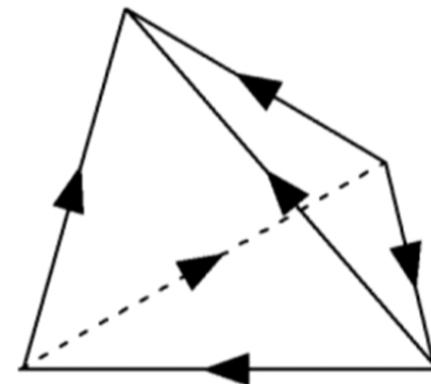
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Unique ground state which connects to a product state through local unitary transformation

$$\Psi_M(\{g_i\}_M) = \frac{\sum_{g_i \in \text{internal}}}{|G|^{N_v^{\text{internal}}}} \prod_{\{ij\dots k\}} \nu_{d+1}^{s_{ij\dots k}}(g_i, g_j, \dots, g_k)$$



**SPT order in any dimension with any group is classified
by $\mathcal{H}^{d+1}[G, U_T(1)]$**

Symmetry	$d = 0$	$d = 1$	$d = 2$	$d = 3$
$U(1) \rtimes Z_2^T$	\mathbb{Z}	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}_2^2
$U(1) \rtimes Z_2^T \times \text{trans}$	\mathbb{Z}	$\mathbb{Z} \times \mathbb{Z}_2$	$\mathbb{Z} \times \mathbb{Z}_2^3$	$\mathbb{Z} \times \mathbb{Z}_2^8$
$U(1) \times Z_2^T$	\mathbb{Z}_1	\mathbb{Z}_2^2	\mathbb{Z}_1	\mathbb{Z}_2^3
$U(1) \times Z_2^T \times \text{trans}$	\mathbb{Z}_1	\mathbb{Z}_2^2	\mathbb{Z}_2^4	\mathbb{Z}_2^9
Z_2^T	\mathbb{Z}_1	\mathbb{Z}_2	\mathbb{Z}_1	\mathbb{Z}_2
$Z_2^T \times \text{trans}$	\mathbb{Z}_1	\mathbb{Z}_2	\mathbb{Z}_2^2	\mathbb{Z}_2^4
$U(1)$	\mathbb{Z}	\mathbb{Z}_1	\mathbb{Z}	\mathbb{Z}_1
$U(1) \times \text{trans}$	\mathbb{Z}	\mathbb{Z}	\mathbb{Z}^2	\mathbb{Z}^4
Z_n	\mathbb{Z}_n	\mathbb{Z}_1	\mathbb{Z}_n	\mathbb{Z}_1
$Z_n \times \text{trans}$	\mathbb{Z}_n	\mathbb{Z}_n	\mathbb{Z}_n^2	\mathbb{Z}_n^4
$Z_2^T \times D_2 = D_{2h}$	\mathbb{Z}_2^2	\mathbb{Z}_2^4	\mathbb{Z}_2^6	\mathbb{Z}_2^9



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$Z_2^T \times D_2 = D_{2h}$	\mathbb{Z}_2^2	\mathbb{Z}_2^4	\mathbb{Z}_2^6	\mathbb{Z}_2^9



Continuum limit analogy

S=1 non-linear sigma model in 1+1D

$$S = \int dx dt \frac{1}{2g} (\partial \mathbf{n}(x, t))^2 + i\theta W, \quad \theta = 2\pi$$

$$W = (4\pi)^{-1} \int dt dx \mathbf{n}(t, x) \cdot [\partial_t \mathbf{n}(t, x) \times \partial_x \mathbf{n}(t, x)]$$

- The theta term has no contribution in the bulk but can induce a WZW term on the boundary describing the motion of a half spin



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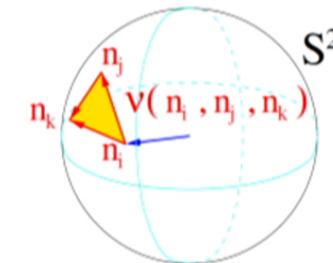
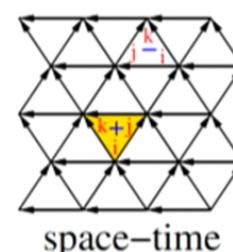
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Discrete version of theta term

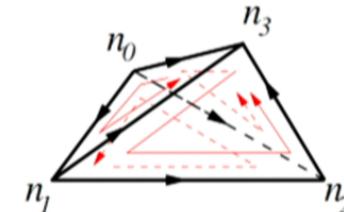
$$e^{-S} = \prod \nu^{s(i,j,k)}(\mathbf{n}_i, \mathbf{n}_j, \mathbf{n}_k),$$

$$\nu^{s(i,j,k)}(\mathbf{n}_i, \mathbf{n}_j, \mathbf{n}_k) = e^{-\int_{\Delta} dx dt}$$



$$\nu(g\mathbf{n}_i, g\mathbf{n}_j, g\mathbf{n}_k) = \nu(\mathbf{n}_i, \mathbf{n}_j, \mathbf{n}_k), \quad g \in SO(3)$$

$$\frac{\nu(\mathbf{n}_1, \mathbf{n}_2, \mathbf{n}_3)\nu(\mathbf{n}_0, \mathbf{n}_1, \mathbf{n}_3)}{\nu(\mathbf{n}_0, \mathbf{n}_2, \mathbf{n}_3)\nu(\mathbf{n}_0, \mathbf{n}_1, \mathbf{n}_2)} = 1.$$



Fermionic TQFTs constructed from Grassmann TNs



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- Fermionic TQFTs describe intrinsic/symmetry protected topological order in fermionic systems.

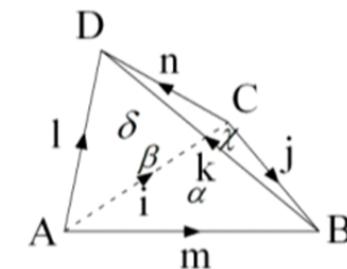


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- Fermionic TQFTs can not be realized in any bosonic systems. e.g., for FQHE, odd/even denote for fermion/boson systems.

Grassmann TNs construction for fermionic TQFTs

$$Z_{top} = \frac{1}{D^{N_v}} \sum_{\alpha\beta\gamma\delta\dots} \sum_{ijklmn\dots} \int \prod_{\text{face}} (d\theta_\alpha)^{P(\alpha)} (d\theta_{\alpha'})^{P(\alpha')} \prod_{\text{link}} d_i \\ \times \prod_{\text{tetrahedron}} G_{kln,\gamma\delta}^{ijm,\alpha\beta} (\theta_\alpha)^{P(\alpha)} (\theta_\beta)^{P(\beta)} (\theta_\gamma)^{P(\gamma)} (\theta_\delta)^{P(\delta)}$$



Fermionic TQFTs constructed from Grassmann TNs

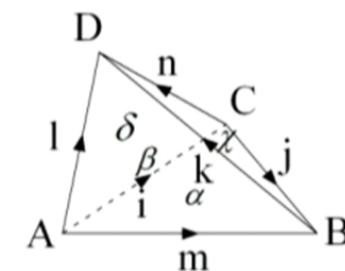
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A local fermionic TQFT requires fermion parity conservation

$$P(\alpha) + P(\beta) + P(\gamma) + P(\delta) = 0 \pmod{2}$$



Conclusion

- Intrinsic topological order and symmetry protected topological order can be related through duality map.
- We propose the TQFTs for SPT orders based on TNs representation. They are discrete analogy of topological terms of non-linear sigma model, whose boundary are described by generalized WZW terms.
- The Grassmann generalization of Turaev-Viro states sum invariants allow us to study fermionic TQFTs.
- The duality map can be generalized to any dimension. SPT orders in 3D will correspond to new intrinsic topological orders.



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- TQFTs constructed from TNs can be defined without space-time background. TNs provide an algebraic way to define TQFTs.
- TQFTs derived from TNs can be regarded as "background" of physical space-time. Space-time may emerge as collective modes of TQFTs.
- QG in 2+1D is a TQFT and can be regarded as a (double) quantum hall system.
- New TQFTs based on TNs construction could be potential candidates for quantum gravity(QG).



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$$\begin{array}{|c|c|} \hline & T \\ \hline g' & \\ \hline g & \\ \hline \end{array}$$

$$g^{\uparrow} = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}$$

$$g^{\downarrow} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$T_{\alpha\beta}(s) = \begin{cases} 0 & \text{if } \beta + s \text{ even} \\ 1 & \text{otherwise} \end{cases}$$

