

Title: More is Different in the Quantum World, in its Own Way

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Abstract: In his article 'More is Different', P W Anderson (1972) sensitized the physics community about importance of emergence, by using concepts such as broken symmetry, emergent hierarchical structures, constructionists converse of reductionism etc. The manifestation of complexity and hierarchy in the quantum many particle systems go beyond broken symmetries. In certain quantum systems we have the opposite - emergence of new local gauge symmetries. This idea was introduced, for example, for Mott insulators, by Anderson and us in 1988. Properties such as entanglement, different possible statistics of identical particles, multi particle wave interference, vastness of Hilbert space etc., which are unique to the world of quantum, leads to a wonderful richness in physics and mathematics, including emergent gauge symmetries, gauge fields, quantum order, holographic correspondences and continuing surprises.

More is Different in Quantum World in its Own Way

Symmetry Enhancement
as opposed to
Symmetry Reduction

Emergence & Effective Field Theories

October 26-28, 2011



G Baskaran

PI





More is Different
Science (1972)

P W Anderson
July 2011

Introduction

More is different ...

symmetry breaking, emergence of hierarchy,
information bearing symmetry breaking, complexity, depth ...

Different in different ways in the World of Quantum Matter

Strong Correlation & Symmetry enhancement

Emergence of gauge symmetries, gauge particles, quantum order,
non Abelian anyons, topological protection, entanglement, holography ...

Mott Insulator & Quantum Spin Liquid as examples

More Is Different

Broken symmetry and the nature of the hierarchical structure of science.

P. W. Anderson

X	Y
solid state or many-body physics	elementary particle physics
chemistry	many-body physics
molecular biology	chemistry
cell biology	molecular biology
•	•
•	•
•	•
psychology	physiology
social sciences	psychology

But this hierarchy does not imply that science X is “just applied Y.”

SCIENCE

4 August 1972, Volume 177, Number 4047

at
each level of complexity entirely new
properties appear

we have begun to formulate a general theory of just how this shift from quantitative to qualitative differentiation takes place. This formulation, called the theory of “broken symmetry,” may be of help in making more generally clear the breakdown of the constructionist converse of reductionism.

World of Quantum Matter

Matter, including light

Natural

&

Man Made

Solids at Low T (phonons)

Magnets, Fe, Ni, Gd, ...

Superconductors, Hg, Pb, Nb ..

Insulators Diamond, Si, GaAs, ..

Oxides, sulfides, polymers. ...

Minerals

Biological Matter

**Photosynthesis, electron transfer
proton tunneling**

Bird navigation

**2 dimensional electron gas
(heterostructure) & large B**

**Doped La_2CuO_4 Superconductor
 Sr_2RuO_4 p + ip superconductor**

SrTiO_3 Quantum paraelectric

Heavy fermions, Kondo insulators

Superfluid & supersolid ^4He

**Quantum dots, Cold atom gas
Photonic crystals ...**

Quantum Mechanics

Schrodinger Equation

$$-i\hbar \frac{\partial \psi}{\partial t} = H\psi$$

ψ - complex
PROBABILITY
AMPLITUDE

NON CLASSICAL
CORRELATION

Hilbert Space
formed by physical states

Linear Vector Space

$|i\rangle, \dots, |j\rangle \dots$

$|\psi\rangle \equiv u|i\rangle + v|j\rangle$ is a
physically meaningful
state



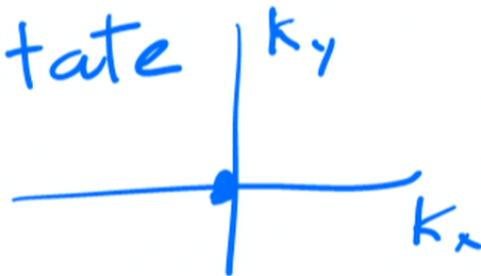
& ENTANGLEMENT

Identical particles

BOSONS

$$\psi(1,2) = \psi(2,1)$$

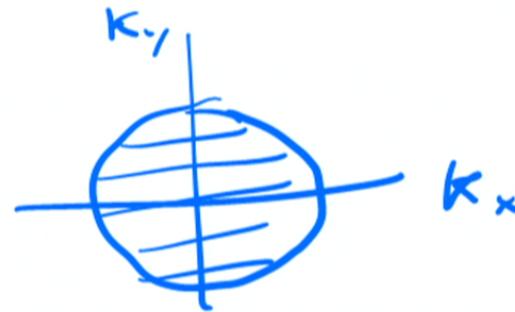
BEC
state



FERMIONS

$$\psi(1,2) = -\psi(2,1)$$

FERMI SEA



Classical Mechanics: $6N$ dimensional
phase space

VASTNESS OF HILBERT SPACE

$$\psi(\vec{r}_1 \sigma_1, \vec{r}_2 \sigma_2, \dots, \vec{r}_N \sigma_N)$$

In many problems,
Direct product of finite dimensional
Hilbert space:

N Qubits: 2^N dim.
Hilbert Space

Evolution in the World of Quantum Matter

**Effectively Non interacting
Non Entangled Ground States**

Photon (Einstein)

Crystals Phonon (Einstein, Debye)
Ferro, Antiferromagnets Magnon -spin waves
Superconductor, Superfluids
Charge density wave, spin density wave ...

Fermi sea

Band insulators, Topological insulators
Holes, electrons, excitons, plasmons, polaritons

Strongly Interacting Quantum Matter

**Mott Insulators, Quantum Spin liquids
Quantum Hall States**

**Luttinger liquids, Kondo Lattice system
Heavy Fermions**

**Superfluid ^4He , Supersolid ^4He
High T_c Superconductors (Cuprates, Fe pnictides)
P-wave superconductors (Sr_2RuO_4)**

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**Hilbert Space Restriction is a basic ingredient
for interaction effects to become important
and appearance of
unexpected properties of matter**

**Free electron gas → tight binding models
(periodic potential produces bands, narrow bands, gaps ...)
LaCuO₄ (narrow band), Quantum Spin Liquids
High T_c Superconductivity**

**Free electron gas → highly degenerate Landau Levels
A rich Quantum Hall Physics**

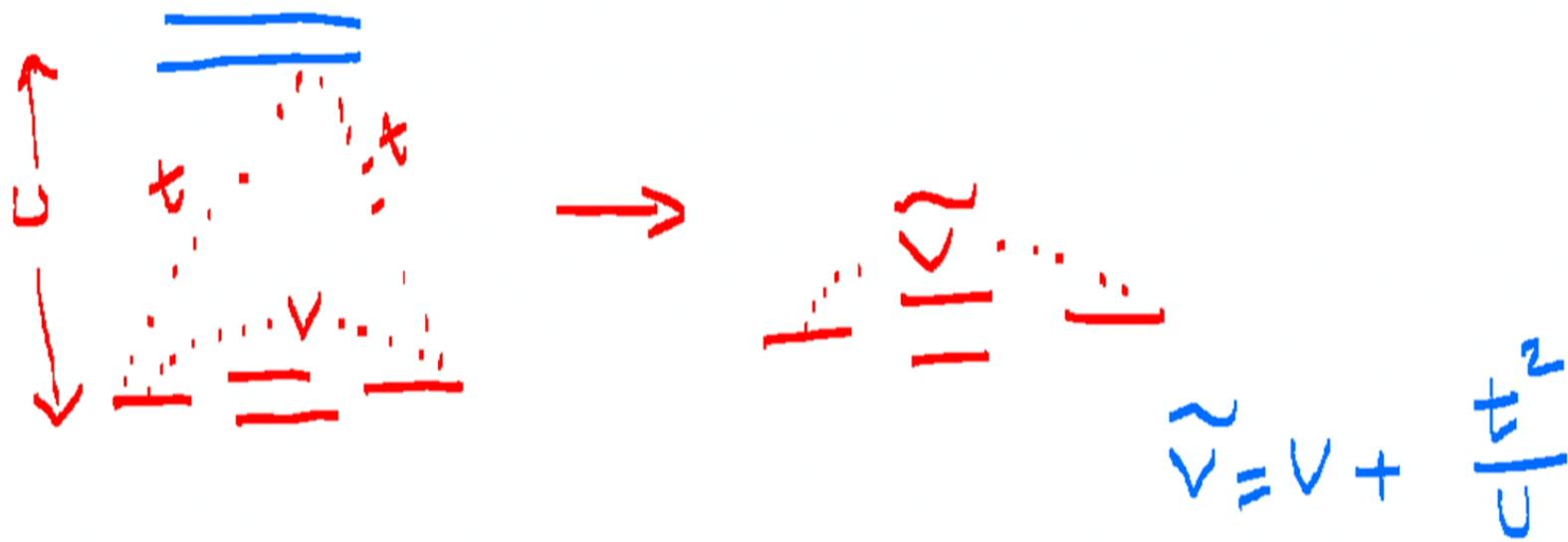
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**Finding the relevant Hilbert space and
 Focusing our attention there
 is a process of renormalization
 Often it is a finite number of discrete step**

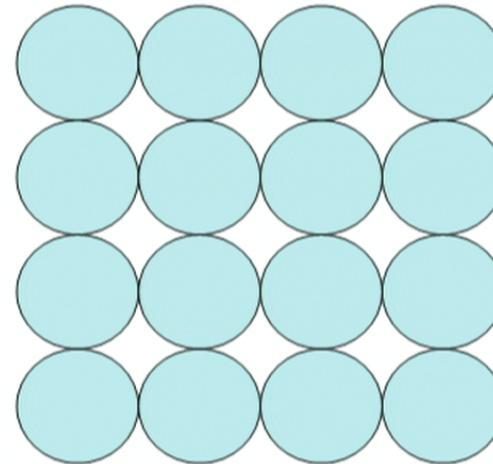
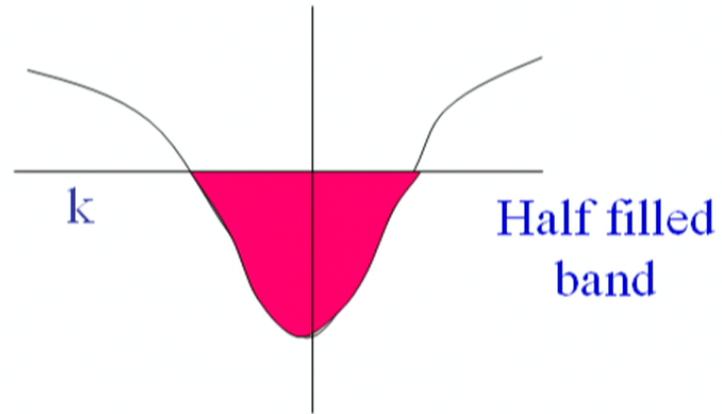
Elimination of high energy state above a gap



A collection of hydrogen atoms forming a hypothetical 3D lattice

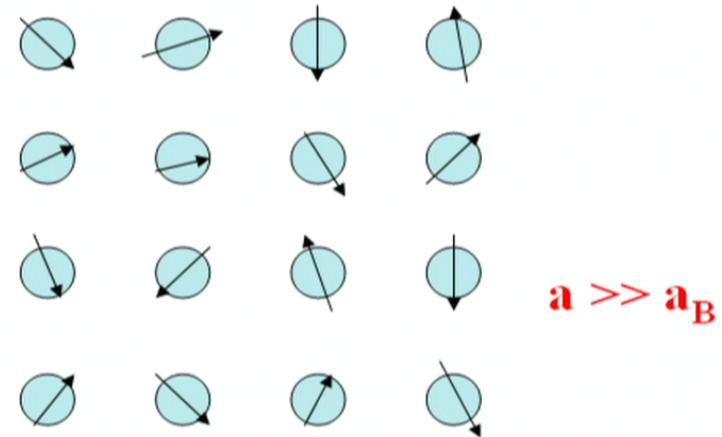
1s states of individual hydrogen atoms strongly overlap and form a tight binding half filled band

It is a metal



$$a \sim a_B$$

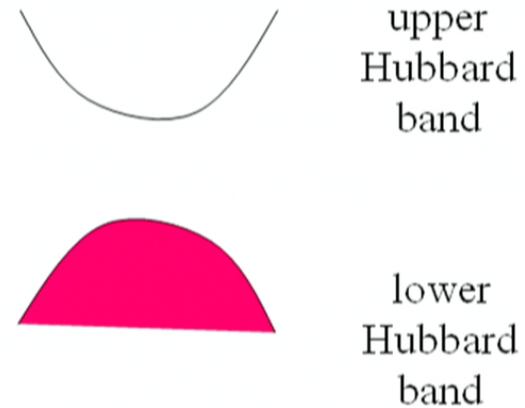
Let us expand the lattice

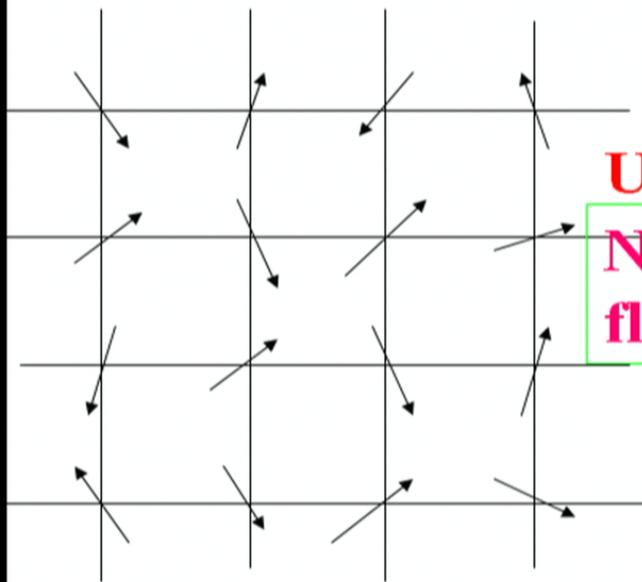


For $a \gg a_B$

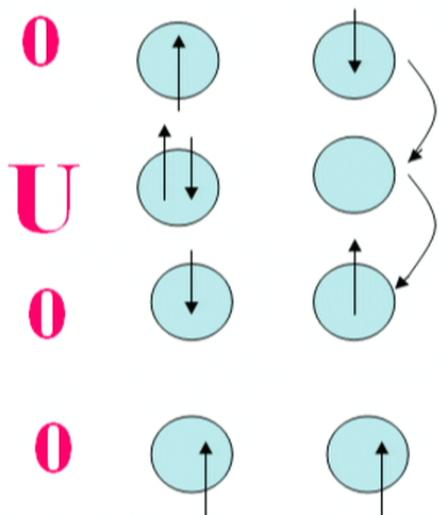
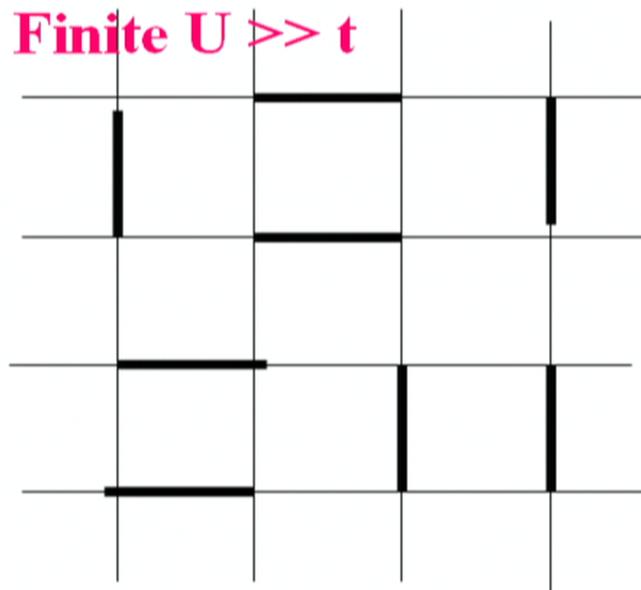
we get a Mott insulator

Spins are soft degrees of freedom
while charges are frozen





U = infinity
No quantum fluctuations



t

Energy gain = J
singlet

$$= \frac{-4t^2}{U}$$

$$= \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

triplet

Energy gain = 0 !

$$H = J \sum_{\langle ij \rangle} (S_i \cdot S_j - \frac{1}{4})$$

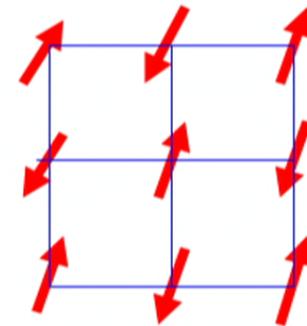
Resonating Valence Bond (RVB) States

Pauling
Anderson 1973
Fazekas

Quantum Heisenberg Antiferromagnets (spin half)

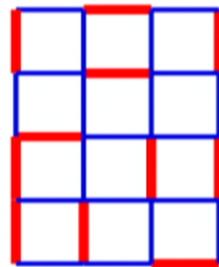
Quantum fluctuations

Encouraged by lattice frustration and lower dimensionality
may destroy long range AFM order (spin crystal)



Resulting in a Quantum Spin Liquid, a RVB state

$$|RVB\rangle = \sum_C |C\rangle$$



$$|RVB; \phi\rangle \equiv P_G \left(\sum_{ij} \phi_{ij} b_{ij}^\dagger \right)^{\frac{N}{2}} |0\rangle$$

Short range spin correlations

$$\langle S_{iz} S_{jz} \rangle \approx e^{\frac{-|i-j|}{\xi}} \quad \text{or} \quad \frac{1}{|i-j|^\alpha}$$

$$\text{---} = \frac{1}{\sqrt{2}} [|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle]$$

How do we get pseudo Fermi surface from theory ?

RVB Meanfield Theory (GB, Zou, Anderson 1987)

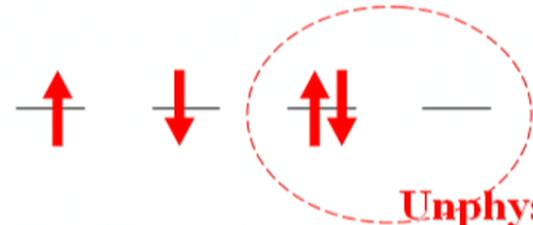
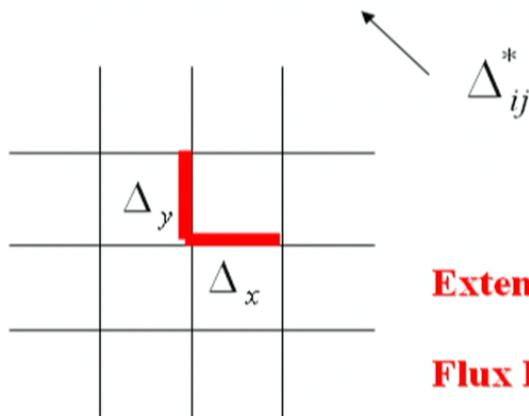
$$J(\mathbf{S}_i \cdot \mathbf{S}_j - \frac{1}{4} n_i n_j) \equiv -J b_{ij}^+ b_{ij}$$

$$\bar{S}_i \equiv C_{i\alpha}^+ \vec{\sigma}_{\alpha\beta} C_{i\beta} \quad n_{i\uparrow} + n_{i\downarrow} = 1$$

$$b_{ij}^+ = \frac{1}{\sqrt{2}} (C_{i\uparrow}^+ C_{j\downarrow}^+ - C_{i\downarrow}^+ C_{j\uparrow}^+)$$

Two complex fermion Hilbert space $2^2 = 4$

$$b_{ij}^+ b_{ij} \rightarrow \langle b_{ij}^+ \rangle b_{ij} + b_{ij}^+ \langle b_{ij} \rangle$$



Unphysical Hilbert space

Extended - S $\longrightarrow \Delta_x = \Delta_y$ BZA

Flux Phase $\longrightarrow \Delta_x = -\Delta_y$ Kotliar, Affleck, Marston

Extended - S $\longrightarrow \Delta_x = \Delta_y$

$$H_s = J \sum_{\langle ij \rangle} (\mathbf{S}_i \cdot \mathbf{S}_j - \frac{1}{4}) = -J \sum_{\langle ij \rangle} b_{ij}^\dagger b_{ij}$$

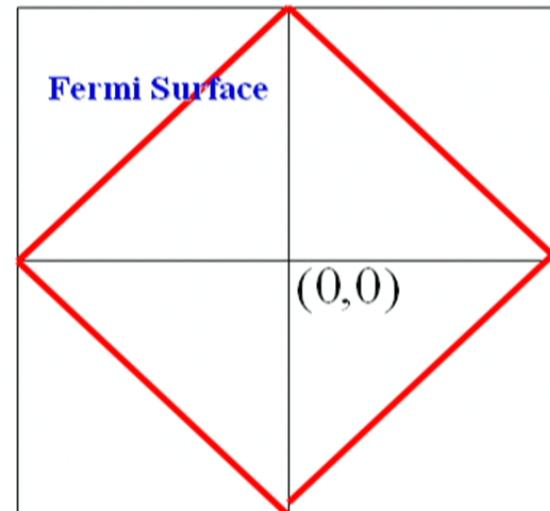
$$b_{ij}^\dagger b_{ij} \rightarrow \langle b_{ij}^\dagger \rangle b_{ij} + b_{ij}^\dagger \langle b_{ij} \rangle$$

$$H_{pair} = -J \sum_{k, k'} \gamma(\mathbf{k} - \mathbf{k}') c_{-k'\downarrow}^\dagger c_{k'\uparrow}^\dagger c_{k\uparrow} c_{-k\downarrow}$$

$$H_{mF} \sim J \sum_{k\alpha} |\cos k_x + \cos k_y| \alpha_{k\sigma}^\dagger \alpha_{k\sigma}$$

(π, π)

$$|2D RVB\rangle = P_G \prod_k (u_k + v_k c_{k\uparrow}^\dagger c_{-k\downarrow}^\dagger) |0\rangle$$



GB, Zou, Anderson 1987

Dynamically generated Gauge Fields

GB, Anderson 1987

$$2^N \rightarrow 4^N$$

**Spin waves (Goldstone modes) are elementary excitations
In magnetically ordered systems**

**Hilbert space enlargement helped us to see
presence of dynamically generated gauge field degree of freedom
in addition to topological excitations such as a spinon**

Local U(1) gauge symmetry

$$H_s = J \sum_{\langle ij \rangle} (\mathbf{S}_i \cdot \mathbf{S}_j - \frac{1}{4}) = -J \sum_{\langle ij \rangle} b_{ij}^\dagger b_{ij}$$

$$C_{i\alpha}^+ \rightarrow e^{i\theta_i} C_{i\alpha}^+$$

$$b_{ij}^+ \rightarrow e^{i\theta_i} b_{ij}^+ e^{i\theta_j}$$

$$\Delta_{ij}^* \rightarrow e^{i\theta_i} \Delta_{ij}^* e^{i\theta_j}$$

GB, Anderson

U(1) RVB magnetic field

$$\Re \int e^{i\oint \mathbf{A}(\mathbf{r}) \cdot d\mathbf{l}} \sim \mathbf{S}_i \times (\mathbf{S}_j \times \mathbf{S}_k)$$

Wen Wilczek Zee

Local SU(2) symmetry

$$C_{i\uparrow} \rightarrow u_i C_{i\uparrow} + v_i C_{i\downarrow}^\dagger \quad |u_i|^2 + |v_i|^2 = 1$$

Affleck Anderson Zou Hsu

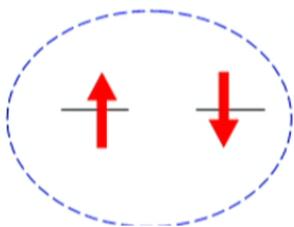
Two complex fermion

Hilbert space $2^2 = 4$

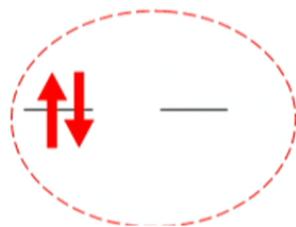
$$2^{2N} = 2^N \times \dots \times 2^N$$

$\leftarrow 2^N \text{ times} \rightarrow$

Physical spin



Pseudo spin



When a physical spin and a Pseudo spin get identified the above 2^N sectors become gauge copies of a Z_2 gauge theory

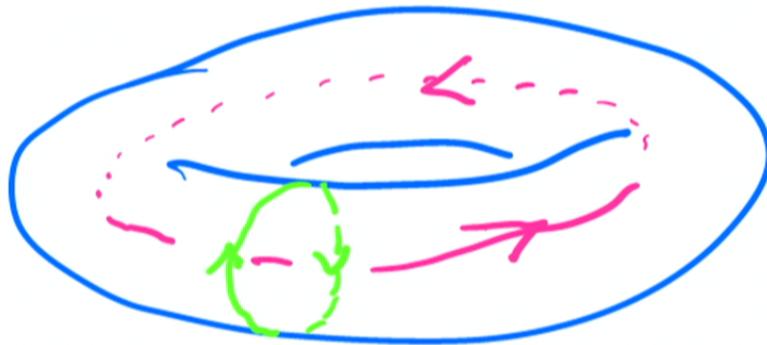
Quantum Order Topological Order

Kitaev, Rokhsar

Haldane
Rezayi

Wen

Gr, Anderson

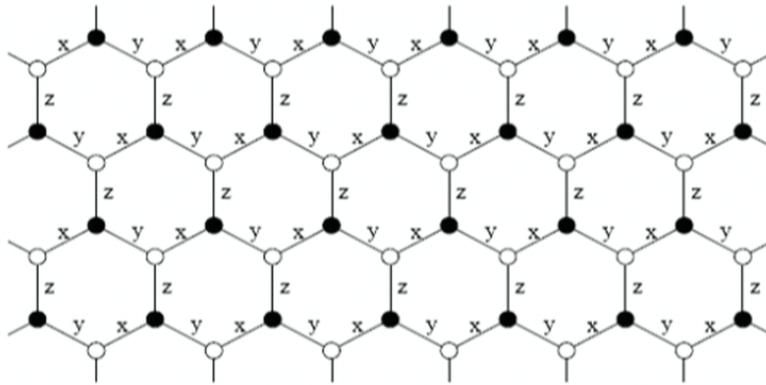


$$i \oint \vec{A} \cdot d\vec{\ell}$$

$$e = \pm 1$$

4 FOLD DEGENERACY

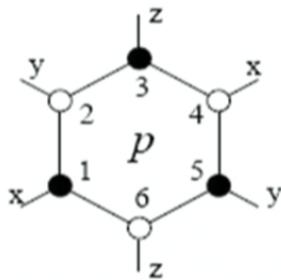
$\oint \vec{A} \cdot d\vec{\ell}$
dynamically generated RVB flux



Kitaev Model
it is a highly Frustrated
Quantum Spin System
 exactly solvable

$$H = -J_x \sum_{x\text{-links}} \sigma_j^x \sigma_k^x - J_y \sum_{y\text{-links}} \sigma_j^y \sigma_k^y - J_z \sum_{z\text{-links}} \sigma_j^z \sigma_k^z$$

Local Conserved Quantities



$$W_p = \sigma_1^x \sigma_2^y \sigma_3^z \sigma_4^x \sigma_5^y \sigma_6^z$$

$$W_p^2 = 1$$

Eigen values of $W_p = 1, -1$

Dirac or complex Fermion

$$\{\psi_i, \psi_j^\dagger\} = \delta_{ij}$$

$$\psi = \frac{1}{2}(\zeta + i\eta)$$

$$\{\zeta_i, \zeta_j\} = \delta_{ij} \quad \eta^2 = \zeta^2 = 1$$

$$D_\zeta = \sqrt{2}$$

$$D_\eta = \sqrt{2}$$

$$D_F = D_\zeta \times D_\eta = \sqrt{2} \times \sqrt{2} = 2$$

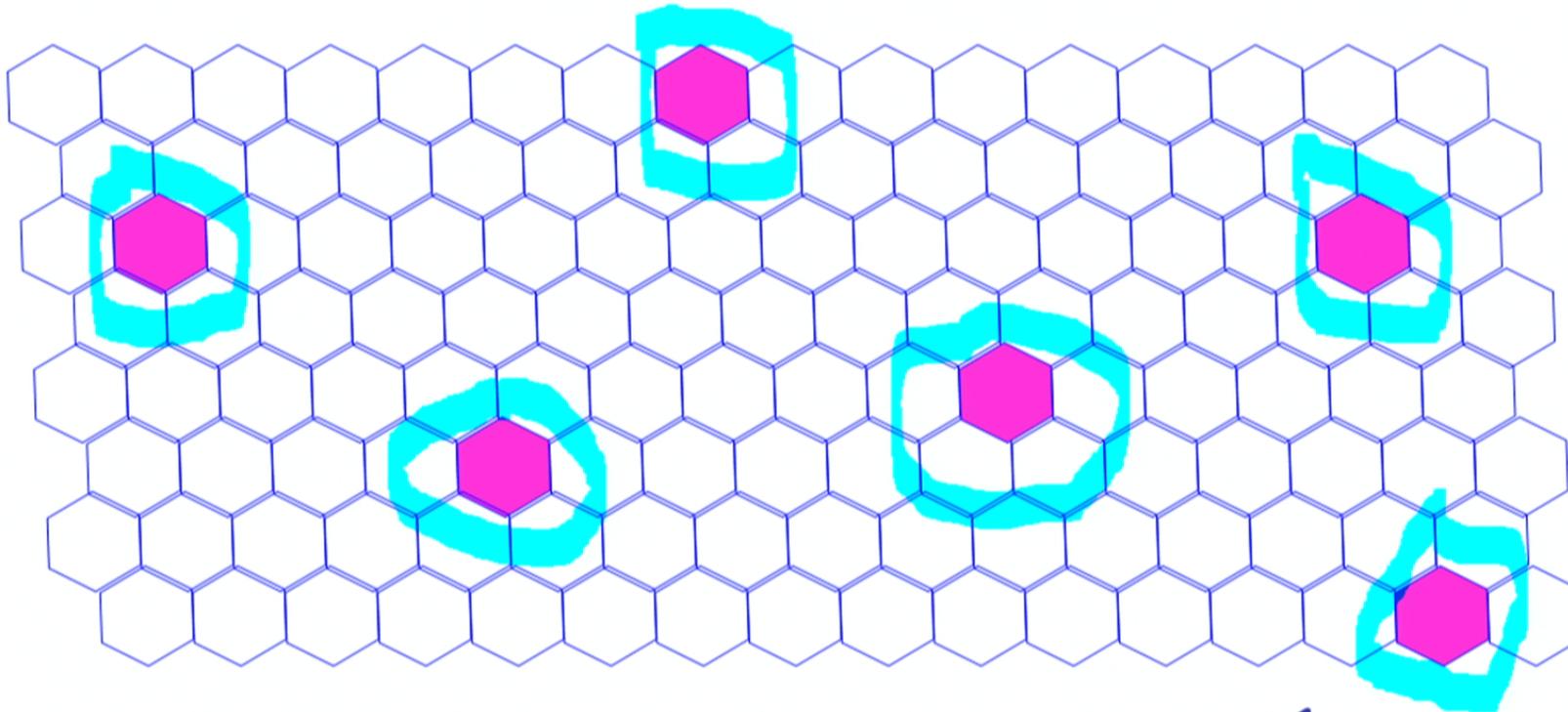
$$|0\rangle, |1\rangle$$

**Majorana or
Real fermions**

Hilbert space dimensions

M well separated vortices Ground state degeneracy is

$$(\sqrt{2})^M$$



6 well separated vortices Ground state degeneracy is

$$(\sqrt{2})^6 = 8$$