

Title: Emergence and Minimal Models in Condensed Matter Physics and Biology

Date: Oct 26, 2011 04:00 PM

URL: <http://pirsa.org/11100055>

Abstract: Our ability to understand the physical world has to a large extent depended on the existence of emergent properties, and the separation of scales that permits effective field theory descriptions to be useful. Exploiting this fact, we can construct minimal models that enable efficient calculation of desired quantities, as long as they are insensitive to microscopic details. This works in many instances in physics, and I give some examples drawn from the kinetics of phase transitions mediated by topological defects. In other fields, such as biology, it is not so clear that these concepts are useful, and I will discuss to what extent emergence and effective theories might be useful.

Emergence and minimal models in condensed matter physics and biology

Nigel Goldenfeld

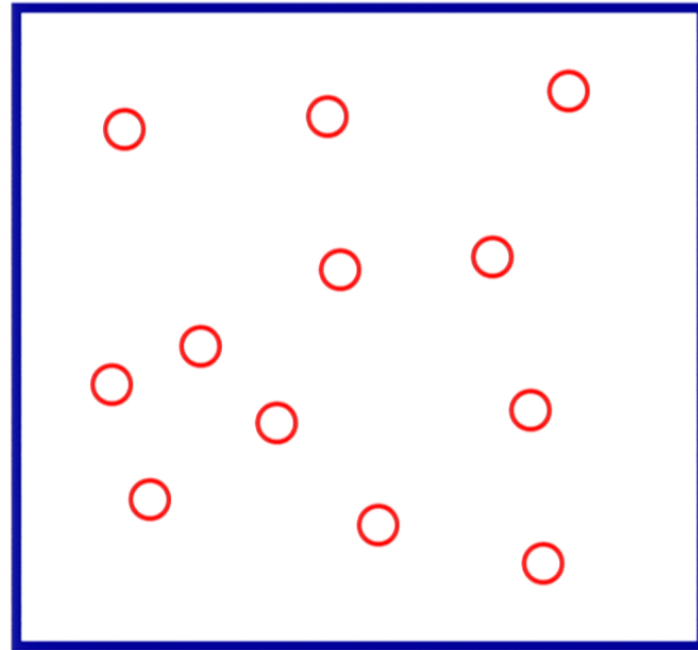
*Department of Physics
University of Illinois at Urbana-Champaign*

Collaborators: Yoshi Oono, M. Mondello (Cell dynamical systems); Kalin Vetsigian (evolution)

Work supported by the US National Science Foundation and NASA

Emergence of the solid state

- **Atoms in a box**
 - **Fluid state**
- **Low temperature**
 - **Crystalline solid forms**
- **Characteristics different from fluid state**
 - **Periodic crystal lattice**
 - **Static shear rigidity**



Emergence of the solid state

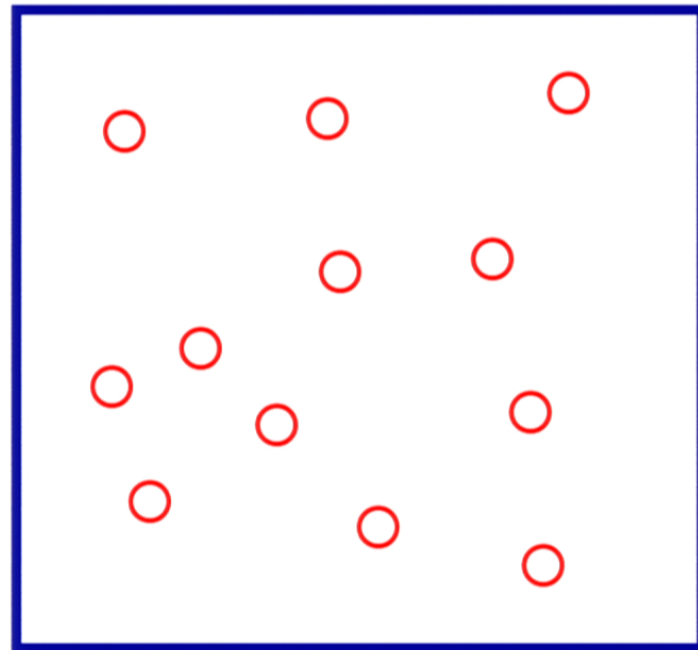
Q: *Why is this interesting?*

A: *Because the underlying physics of the crystal is exactly the same as the fluid!*

– **Total Energy =**

Kinetic energy+ Potential energy

– **Lowering temperature has not changed the interactions between the atoms**



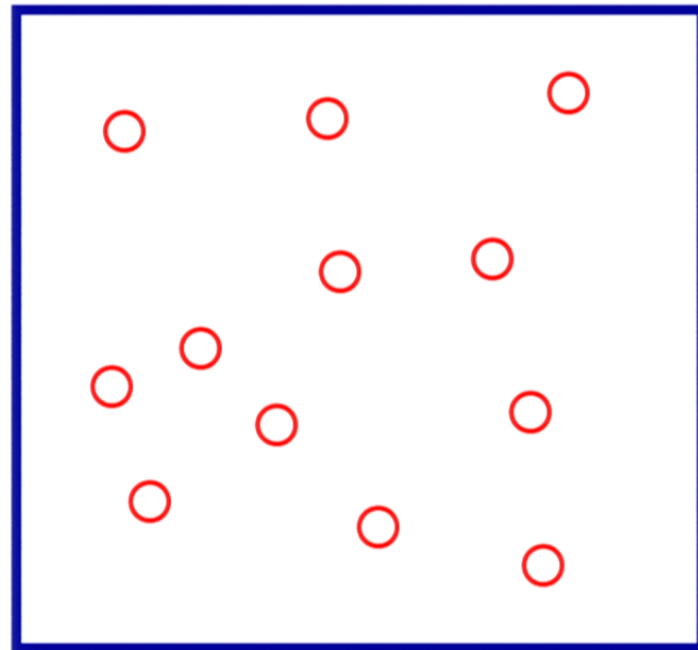
Emergence of the solid state

Conclusion:

Although the pairwise interactions between atoms did not change

the correlations between them changed when the temperature was lowered

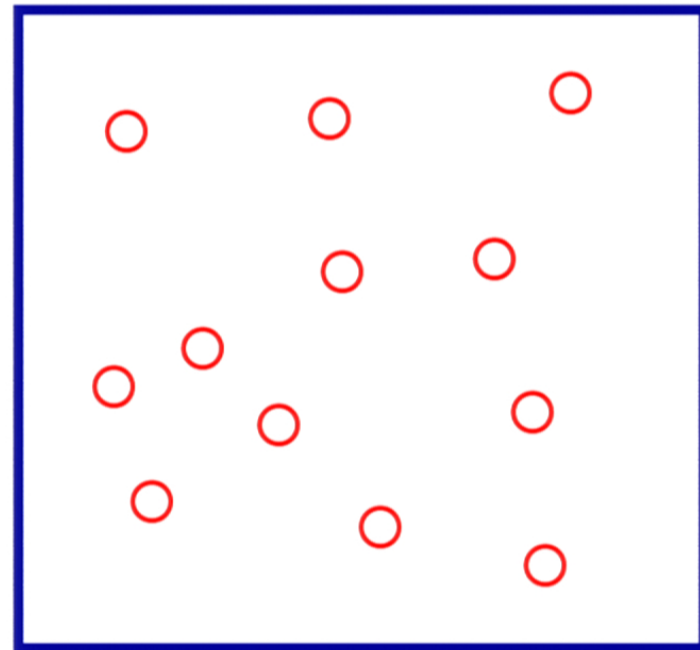
So the solid state emerged **as if by magic out of the fluid state!**



Emergence: the big deal

The solid state has its own new laws of physics!

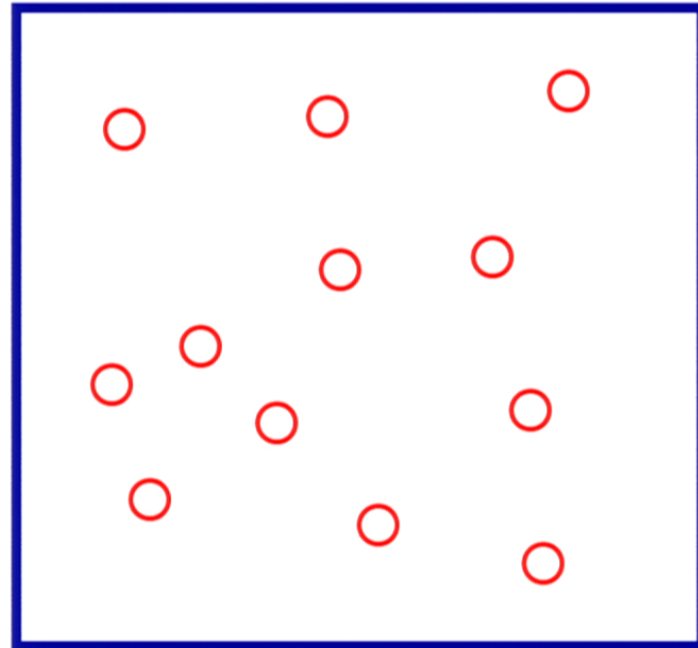
- New long-range forces that are completely distinct from the original forces between the atoms
 - **Elasticity!**
 - Response of solid to external disturbances completely different from that of a fluid
- Origin of these new laws: the **collective**, statistical behaviour of the atoms



Emergence: the big deal

The solid state has its own new laws of physics!

- New long-range forces that are completely distinct from the original forces between the atoms
 - **Elasticity!**
 - Response of solid to external disturbances completely different from that of a fluid
- Origin of these new laws: the **collective**, statistical behaviour of the atoms

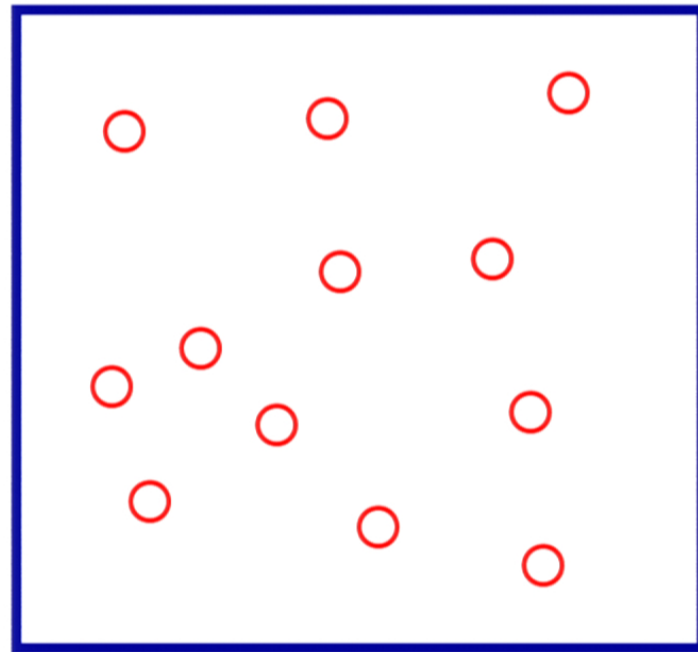


Emergence: new laws of physics?

The solid state has its own new "laws of physics"!

Really?

Well, yes! At the **level of description** of a solid, the actual interactions between the atoms cannot be detected --- only the continuum mechanics of the solid is observable!



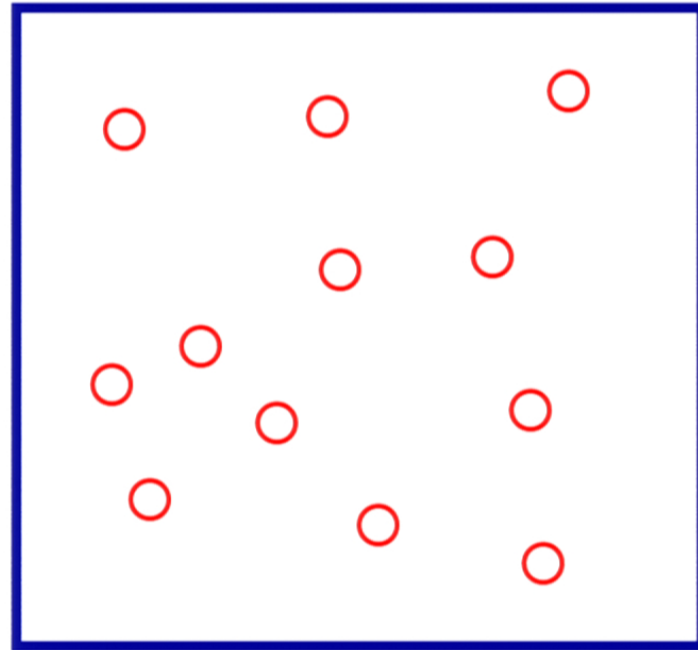
Emergence: new laws of physics?

This is why it took so long for the existence of atoms to be deduced!

Q: So where do atoms and their interactions matter?

A: The elastic forces are characterized by phenomenological parameters such as shear modulus that can only be computed from the atomic level of description.

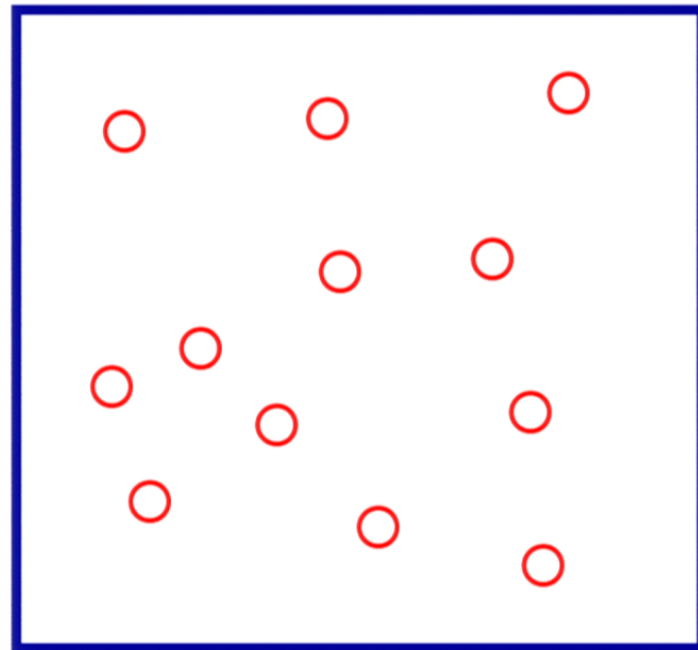
When we do atomic physics or chemistry, we do not need to take into account the mass of the top quark!



Emergence: new laws of physics?

Summary:

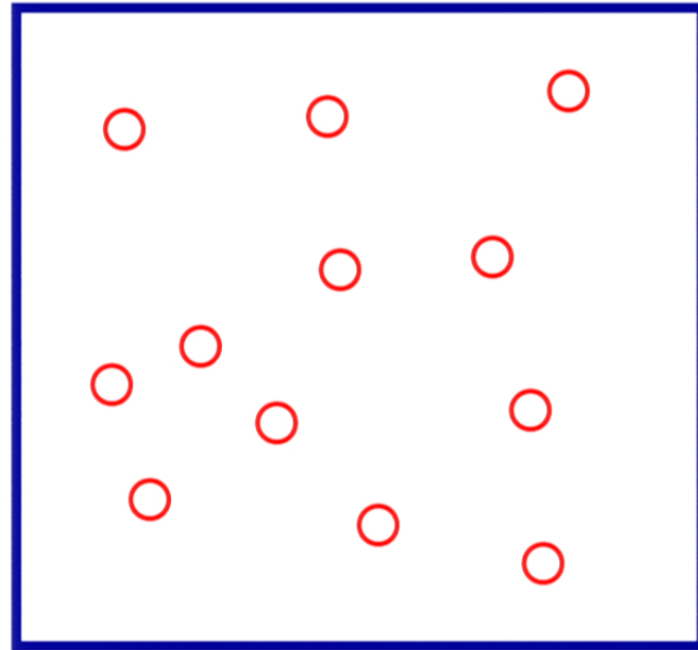
- Emergent levels of description absorb the properties of lower levels of description into phenomenological parameters
 - Typically a small number!



Emergence: new laws of physics?

Summary:

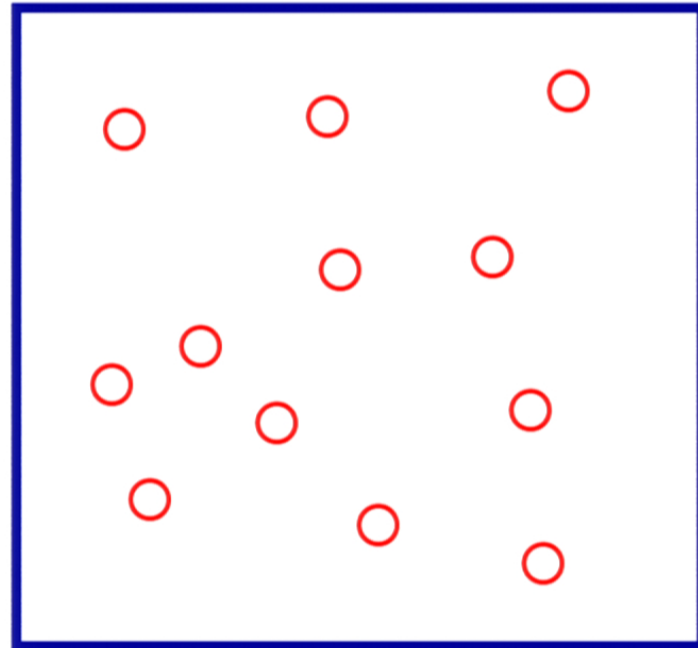
- Emergent levels of description absorb the properties of lower levels of description into phenomenological parameters
 - Typically a small number!
- This makes “fundamental physics” hard, because one “goes against the flow” of emergence



Emergence: new laws of physics?

Summary:

- Emergent levels of description absorb the properties of lower levels of description into phenomenological parameters
 - Typically a small number!
- This makes “fundamental physics” hard, because one “goes against the flow” of emergence
 - Doooms reductionism!
- But collective properties of matter can be understood, **and should be understood**, without requiring knowledge “all the way down”



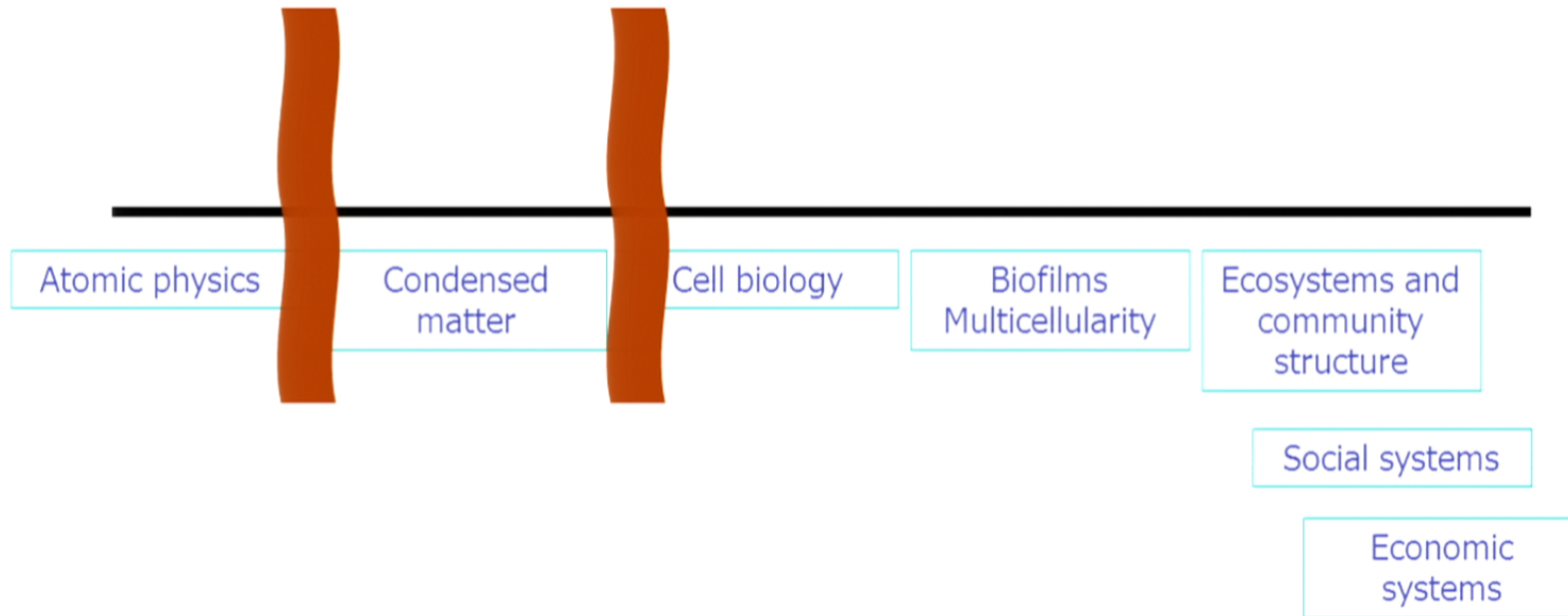
Characteristics of emergence

- **Emergence is associated with loss of uniqueness**
 - Fluid state still possible at low temperature but much less likely than solid state
 - Phase transitions separate states
 - Collective, only well-defined for thermodynamically large systems
- **New, collective properties arise in emergent states**
 - Crystals, magnets, superconductors, quantum hall effects ...
 - Turbulence? Gravity? Life?
 - Do emergent phenomena bring molecules to life?

Characteristics of emergence

- **Emergence is associated with loss of uniqueness**
 - Fluid state still possible at low temperature but much less likely than solid state
 - Phase transitions separate states
 - Collective, only well-defined for thermodynamically large systems
- **New, collective properties arise in emergent states**
 - Crystals, magnets, superconductors, quantum hall effects ...
 - Turbulence? Gravity? Life?
 - Do emergent phenomena bring molecules to life?

Emergence => complexity



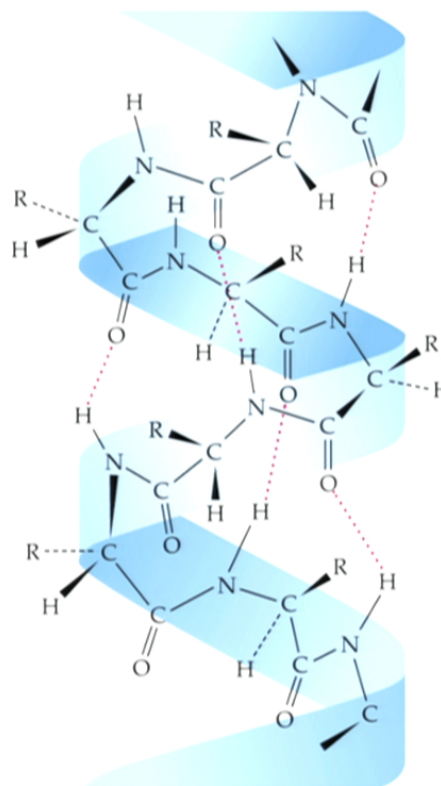
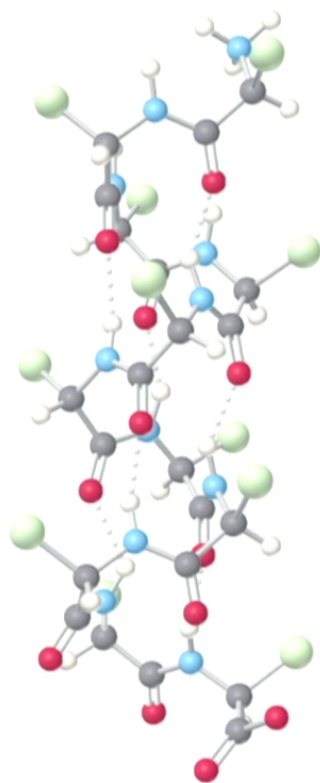
National Academies Keck Futures Initiative: Complexity

Nigel Goldenfeld on Emergence and Collective Phenomena

Overview

- **Minimal models and universality**
 - What aspects of physical phenomena should we try to capture?
 - What aspects of physical phenomena should we try to predict?
- **Modelling without calculus**
 - Cell dynamical systems
- **Space-time patterns during phase transitions**
 - Universal features of pattern formation
 - Comparison with experiment
- **Minimal models in biology – a waste of time?**

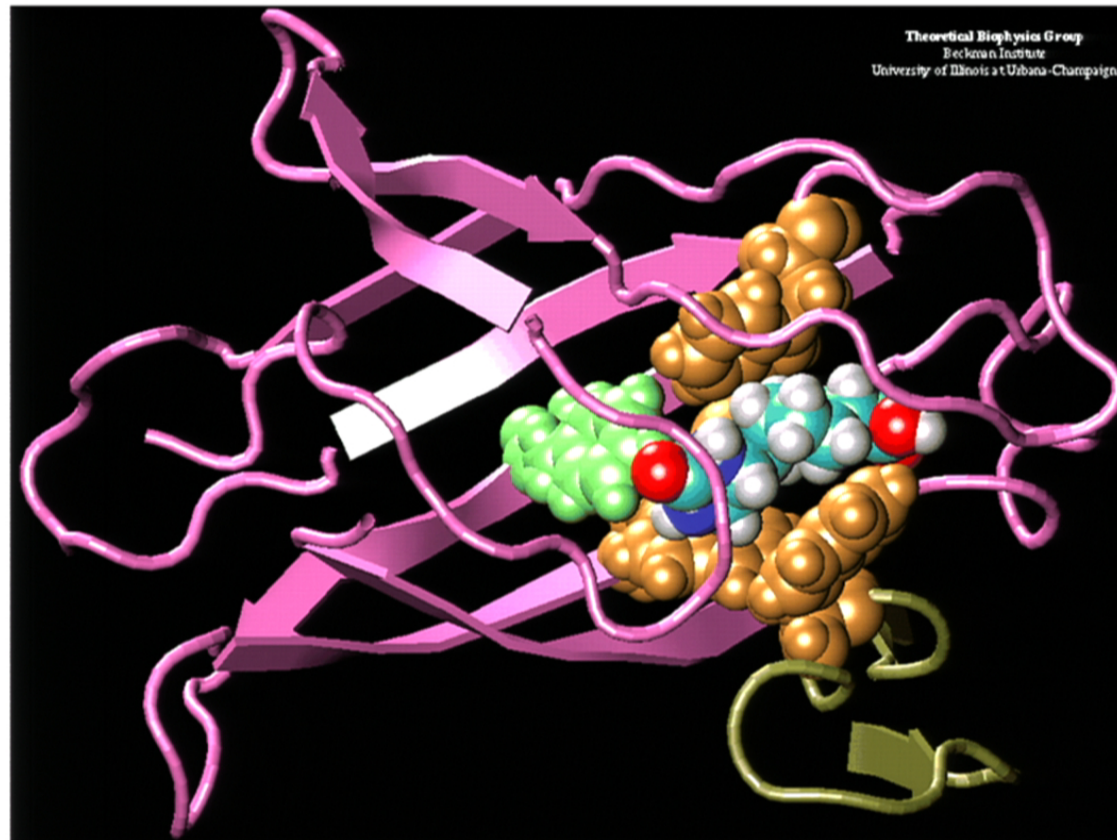
Perceptions of reality: polymers



From *Chemistry* by
McMurry & Fay

Chemist

Perceptions of reality: polymers



Computational biologist

Perceptions of reality: polymers

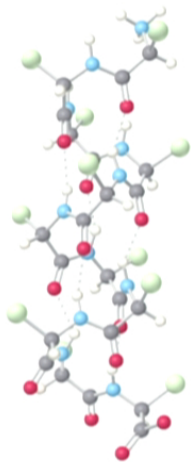


The Gutsy Gourmet

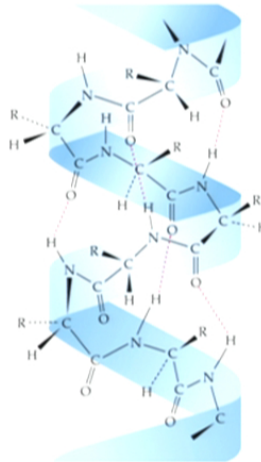
Theoretical physicist

Perceptions of reality: polymers

What questions can each of these representations answer?



Chemical bonding,
reactions, ...



Elasticity, large
scale motions,
folding



Thermodynamics,
light scattering,
rheology

Perceptions of reality: polymers

- How big is a polymer?



Total length of
polymer is $L = N \ell$

Perceptions of reality: polymers

- **How big is a polymer?**

- Chemistry

- model bond angles and energy barriers to rotational states, solvent molecules
 - Sample over long times, or many independent conformations
 - Predict R in Angstroms

- Physics

- **Random walk of N steps each of length ℓ**
 - **Einstein argument: $R^2 = N\ell^2 = \ell L$**
 - **Excluded volume: polymer cannot cross itself, because two atoms cannot occupy same position. Expect chain to swell:**

$$R = A(\ell) L^{0.588} \text{ as } L \rightarrow \infty$$

- **The scaling with L is:**

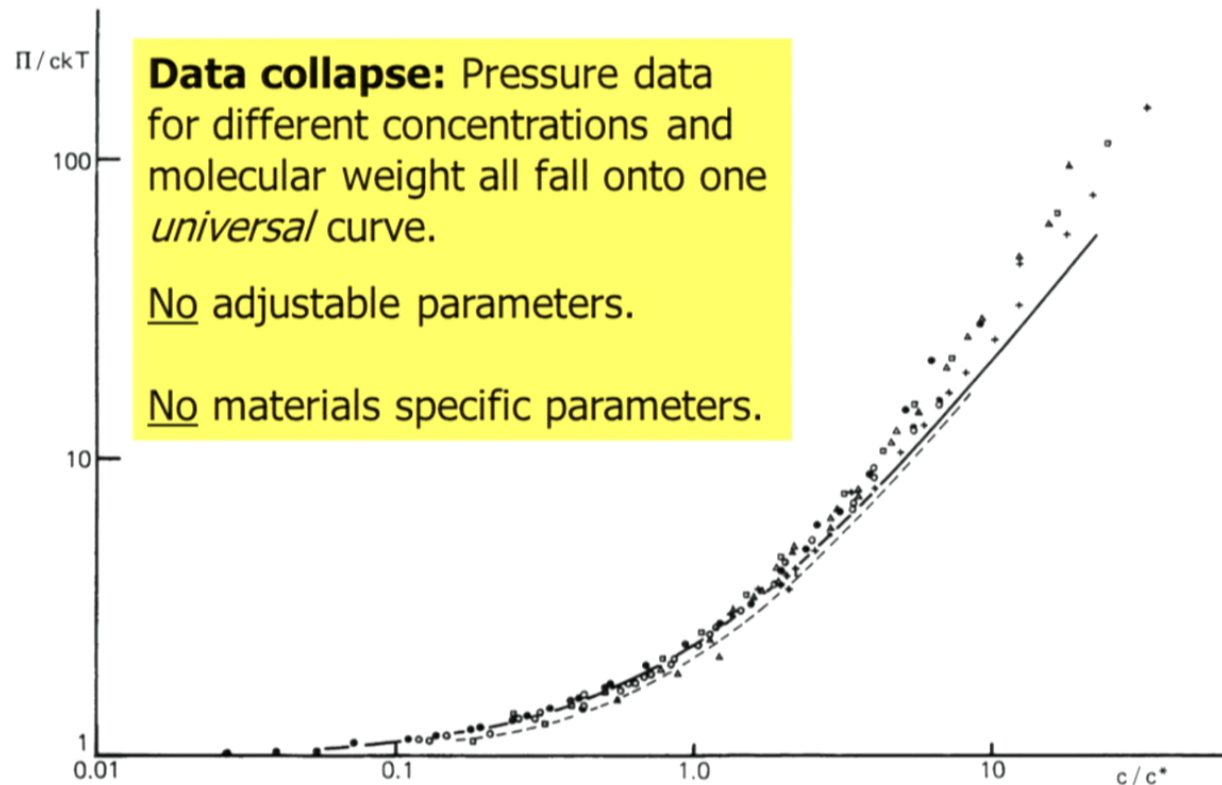
- Universal, i.e. independent of the chemical structure
 - Asymptotic for large chains

- **Cannot predict R in Angstroms**



Total length of polymer is $L = N \ell$

Perceptions of reality: polymers



Osmotic pressure of polystyrene in solution as a function of concentration. Data from Wiltzius *et al* (1983) for $10^4 < M_w < 10^7$ in two different solvents. Theory by Ohta and Oono (1982).

Perceptions of reality: polymers

- **The spaghetti model of a polymer is an example of a *minimal model***
 - Flexible, random walk
 - Self-avoidance: two atoms cannot occupy same point
 - Statistical thermodynamic equilibrium
- **Good for answering certain questions**
 - Universal quantities: pressure, scaling of size with number of monomers, ...
- **Lousy at answering other questions**
 - E.g. what is the size of the polymer coil (in Angstroms)?

Universality in physics

- **Successful scientific predictions are always the result of minimal models!**
 - Cannot include every detail in a model.
 - **Physics: usually only a few details important: some separation of energy scales. "Easy".**
 - Standard model of high energy physics is a minimal model (23 adjustable parameters!) Do you really think this is a "fundamental" theory?
 - **Chemistry: usually need a lot of details, but at the atomic level. Separation of energy scales. Don't worry about quarks. "Hard".**
 - **Economics: every detail matters, it seems. No obvious separation of scales. "Impossible".**

Part 1: Condensed matter

Universality in complex patterns

- **Complex patterns can arise from simple models, algorithms, or equations**
 - E.g. diffusion-limited interface dynamics
- **Can we model microstructure realistically using minimal models?**
- **Shape, scaling: yes**
Dimensions: no
 - Useful for predicting morphology phase diagram

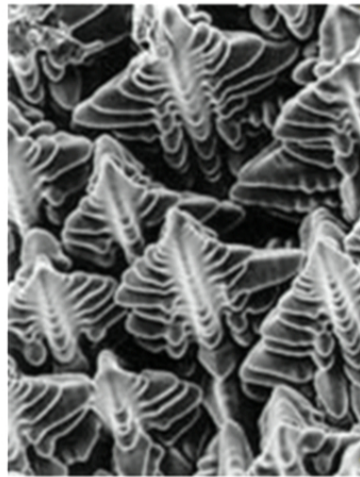
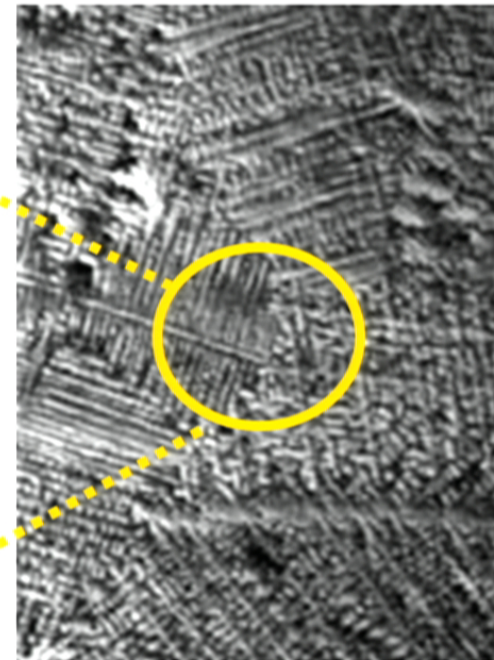
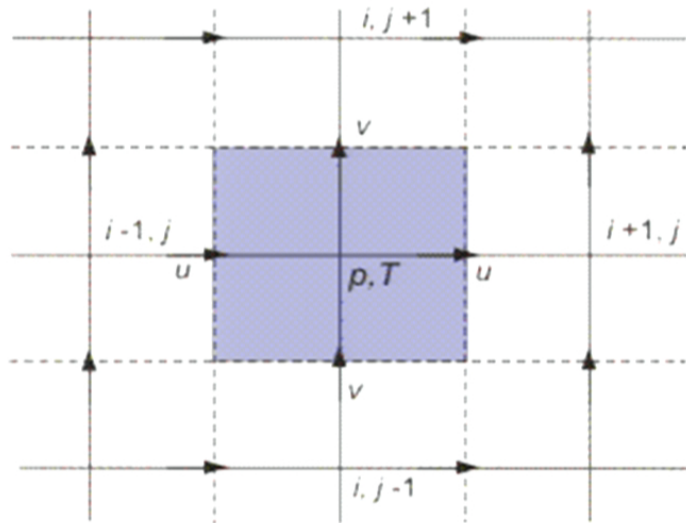


Image: David and Boatner (1997)



Numerical analysis: lattice vs. continuum



Can we study properties of materials without taking the continuum limit? Are the lattice model and the continuum model in the same universality class?

Conventional approach is to require grid convergence

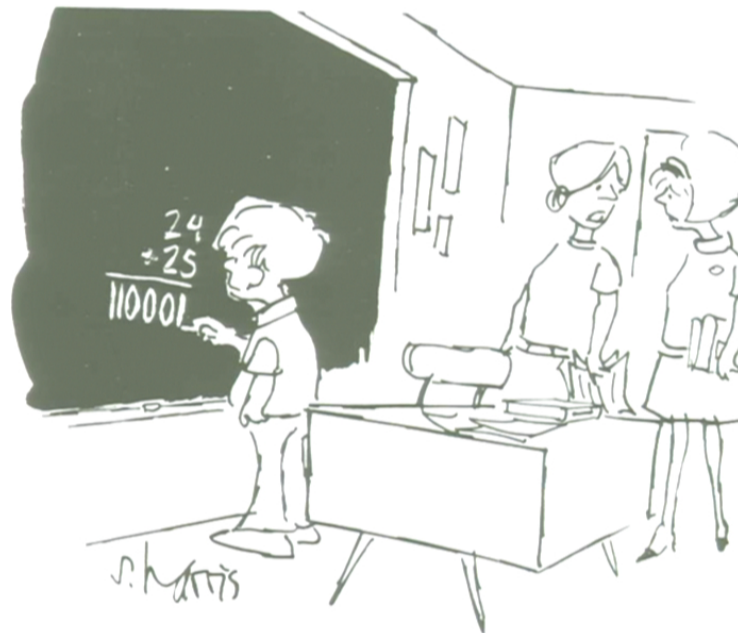
<http://www.thermopedia.com/content/5413/787NHTFig3.gif>

CONVENTIONAL

MODEL AS P.D.E.
DISCRETIZE P.D.E.
COUPLED MAPS
COMPUTE

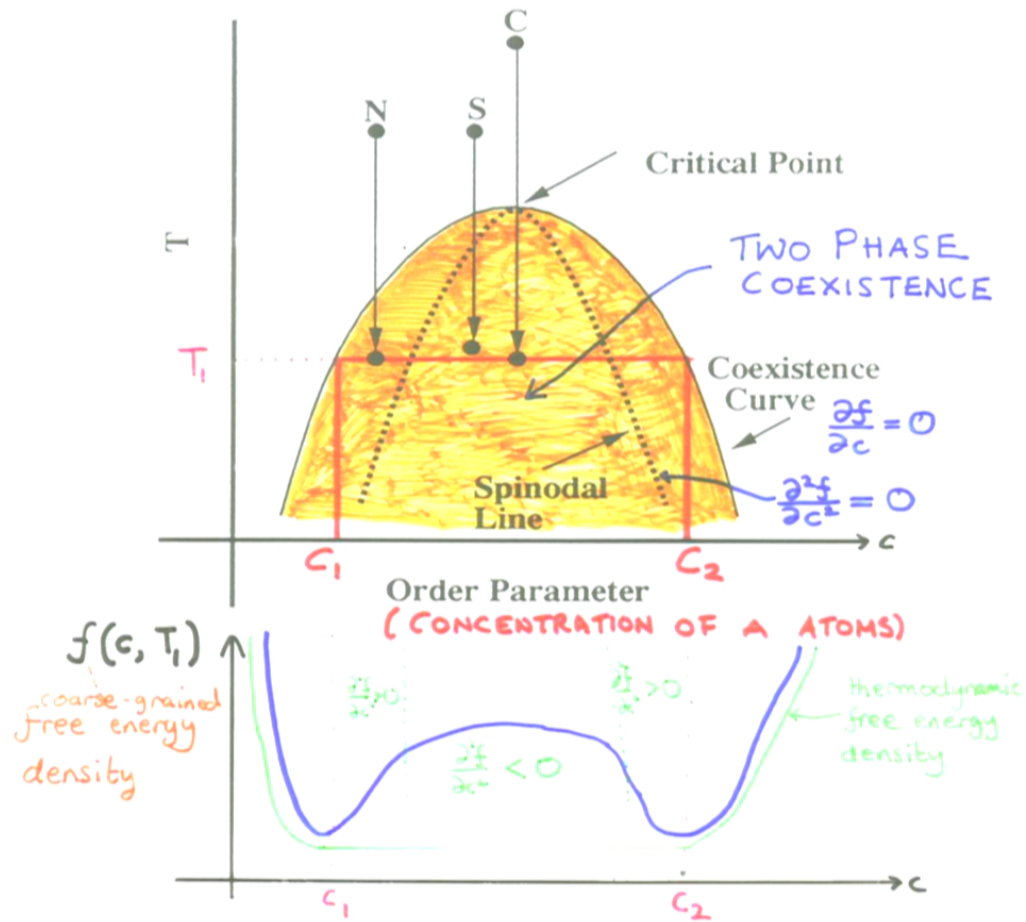
CDS

MODEL AS
COUPLED MAPS
COMPUTE



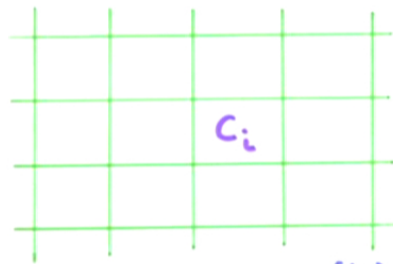
"It was bound to happen —
they're beginning to think like binary computers."

PHASE DIAGRAM OF A-B ALLOY

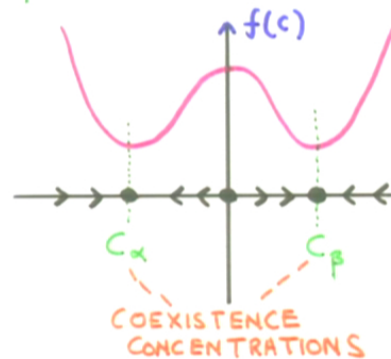


PROTOTYPE CDS: SPINODAL DECOMPOSITION

Y. OONO + S. PURI, Phys. Rev. Lett. 58, 836 (87)
Phys. Rev. A 38, 434 (88); *ibid.* 38, 1542 (88).



COARSE-GRAIN
SYSTEM ON
SCALE OF
CORRELATION LENGTH



IN EACH CELL,
CONCENTRATION
RELAXES IN "POTENTIAL"
FROM FREE ENERGY
DYNAMICS HAS THREE
FIXED POINTS

- MODEL CELL DYNAMICS BY ANY MAP WHICH HAS FIXED POINT STRUCTURE OF EVOLUTION EXPECTED ON PHYSICAL GROUNDS. **NOT DISCRETISATION OF PDE!**
- COUPLE CELLS: CELL CONCENTRATION ALSO EVOLVES ACCORDING TO THE AVERAGE CONCENTRATION OF NEIGHBOURING CELLS.
- IMPLEMENT CONSERVATION LAW

Minimal model for approach to equilibrium

- **Discrete coupled map lattice**
 - No continuum limit taken:
 - **$dt, dx = O(1)$**
 - But patterns and correlations correctly reproduced
- **Not naïve discretisation of Cahn-Hilliard**
 - Instead use onsite map that is injective
 - Algorithm is topologically stable
 - **No numerical instabilities as artifacts**
- **Use isotropic discretisation of Laplacian**
- **Impose conservation law after each time step**

CDS FOR SPINODAL DECOMPOSITION

$c_t(\eta) \equiv$ CONCENTRATION (i.e. OF SPECIES A) IN CELL AT η AT TIME STEP t .

EQUATION OF MOTION:

$$c_{t+1}(\eta) = c_t(\eta) + I_t(\eta) - \langle I_t(\eta) \rangle$$

EVOLUTION

IMPLEMENT CONSERVATION

EVOLUTION: $F(x) = A \tanh\left[\tanh^{-1}\left(\frac{1}{A}\right)x\right]$, $A=1.15$

$$I_t(\eta) = F(c_t(\eta)) - c_t(\eta) + D(\langle c_t(\eta) \rangle - c_t(\eta))$$

SINGLE CELL DYNAMICS

COUPLING BETWEEN CELLS = DIFFUSION
 $D = 0.7$

AVERAGING:

$$\langle U \rangle \equiv \frac{1}{6} \sum_{j \text{ n.n.}} u_j + \frac{1}{12} \sum_{j \text{ n.n.n.}} u_j (+ \dots)$$

NEAREST NEIGHBOURS

NEXT NEAREST NEIGHBOURS

CDS FOR SPINODAL DECOMPOSITION

Minimal model for approach to equilibrium

- Discrete coupled map lattice

- No continuum limit taken:

- $dt, dx = O(1)$

- But patterns and correctly reproduced

- Not naïve discrete Cahn-Hilliard

- Instead use one-to-one mapping

- Algorithm is topologically stable

- No numerical instabilities or artifacts

- Use isotropic discretisation of Laplacian

- Impose conservation law after each time step

$c_t(\underline{n}) \equiv$ CONCENTRATION (i.e. OF SPECIES A) IN CELL AT \underline{n} AT TIME STEP t .

EQUATION OF MOTION:

Algorithm models same thing as the PDE: both in same universality class as physical process.

Does it work?

$$c_t(\underline{n}) - \langle I_t(\underline{n}) \rangle$$

IMPLEMENT CONSERVATION

$$\tanh\left[\left(\tanh^{-1}\frac{1}{A}\right)x\right], A=1.15$$

$$+ D(\langle c_t(\underline{n}) \rangle - c_t(\underline{n}))$$

COUPLING BETWEEN CELLS = DIFFUSION
 $D = 0.7$

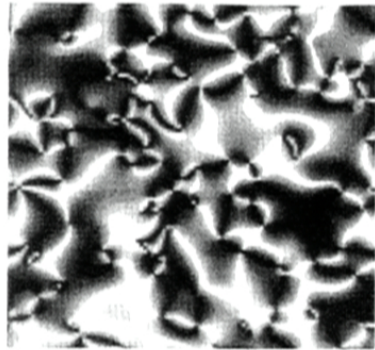
SINGLE CELL DYNAMICS

AVERAGING:

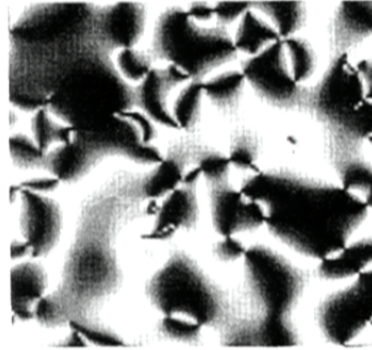
$$\langle U \rangle \equiv \frac{1}{6} \sum_{j \text{ n.n.}} u_j + \frac{1}{12} \sum_{j \text{ n.n.n.}} u_j (+ \dots)$$

NEAREST NEIGHBOURS

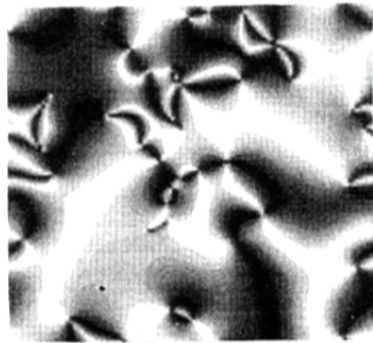
NEXT NEAREST NEIGHBOURS



a) $t = 500$



b) $t = 1000$



c) $t = 2000$

FIG. 17. Intensity of light transmitted through a uniaxial nematic film situated between crossed polarizers (the Schlieren pattern). We show a system of size 100×100 at times $t = 500$, $t = 1000$, and $t = 2000$ after the quench. The defects appear in the pictures as the intersections of two bright and two dark lines.

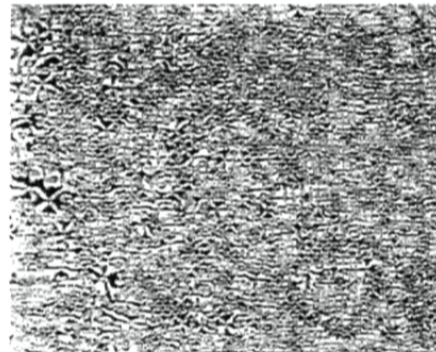
Zapotowcky et al. (1995).

Experimental data:
dynamics of
disclination
coarsening in liquid
crystals film.

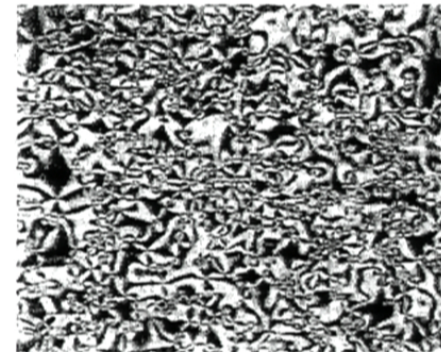
Cross-polarisers
shows the dynamics
of disclination lines
and other topological
defects

Coarsening reflected
in growth of mean
separation of defects

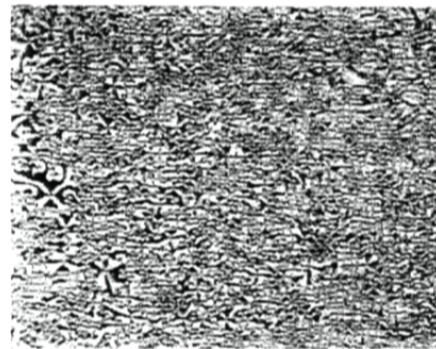
Nagaya *et al.* J.
Phys. Soc. Jpn.
64, 78 (1995).



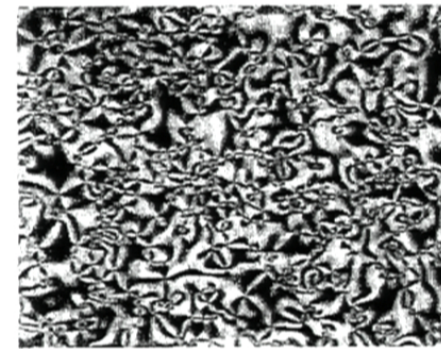
4 sec



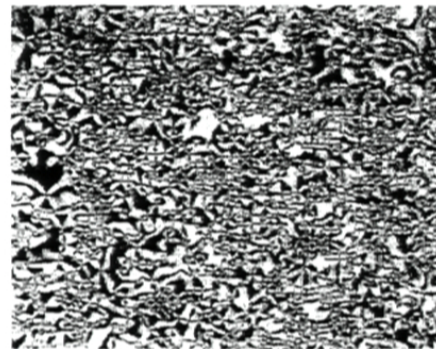
24 sec



8 sec



32 sec



16 sec



0.5 mm

Fig. 2. Time evolution of the disclination texture.¹⁹⁾

Scaled defect correlation function as a function of scaled distance.

Raw data are correlation function at different positions at different time.

Data collapse predicted by theory.

Experiment agrees with simulation with no adjustable parameters.

Mean field theory by Liu and Mazenko gives good description except near defect core.

Mondello and NG, *Phys. Rev. A* **45**, 657 (1992).

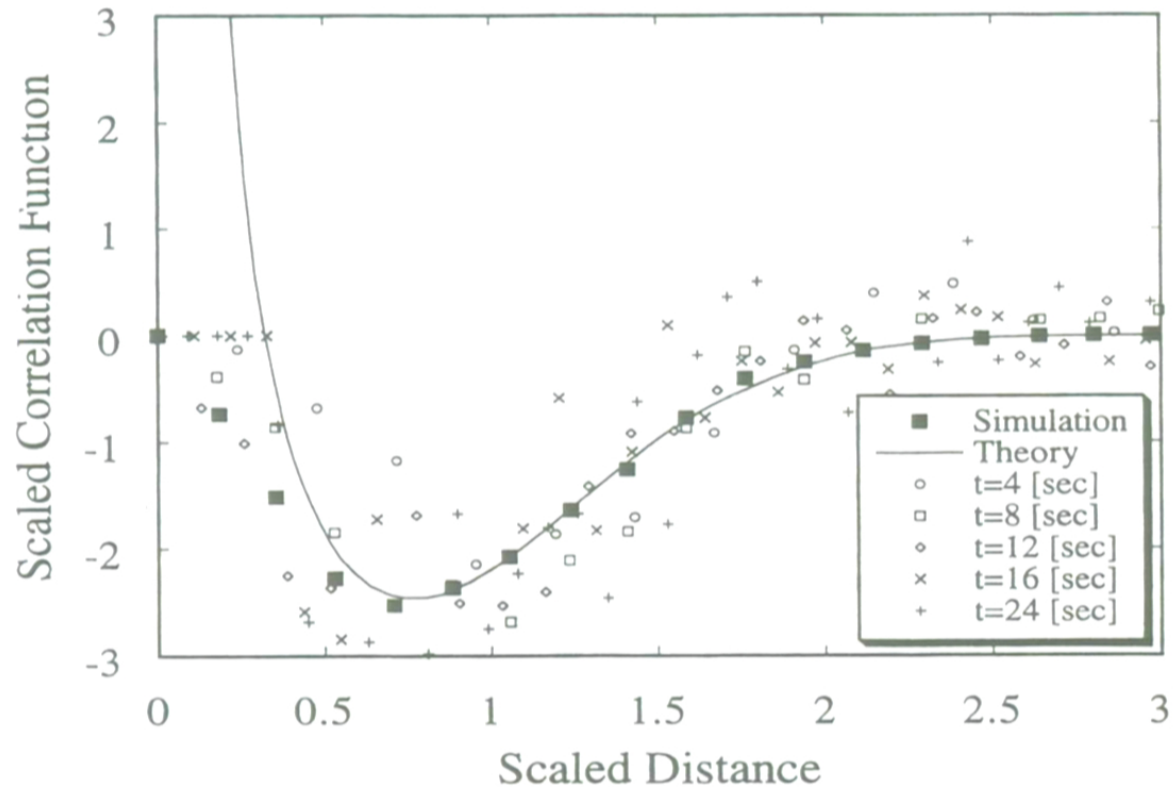


Fig. 6. The scaled correlation function of defect densities. The solid line and the large closed squares indicate the theoretical curve by Lue *et al.*¹¹⁾ and the data of simulation by Mondello *et al.*,⁸⁾ respectively.

Scaled defect correlation function as a function of scaled distance.

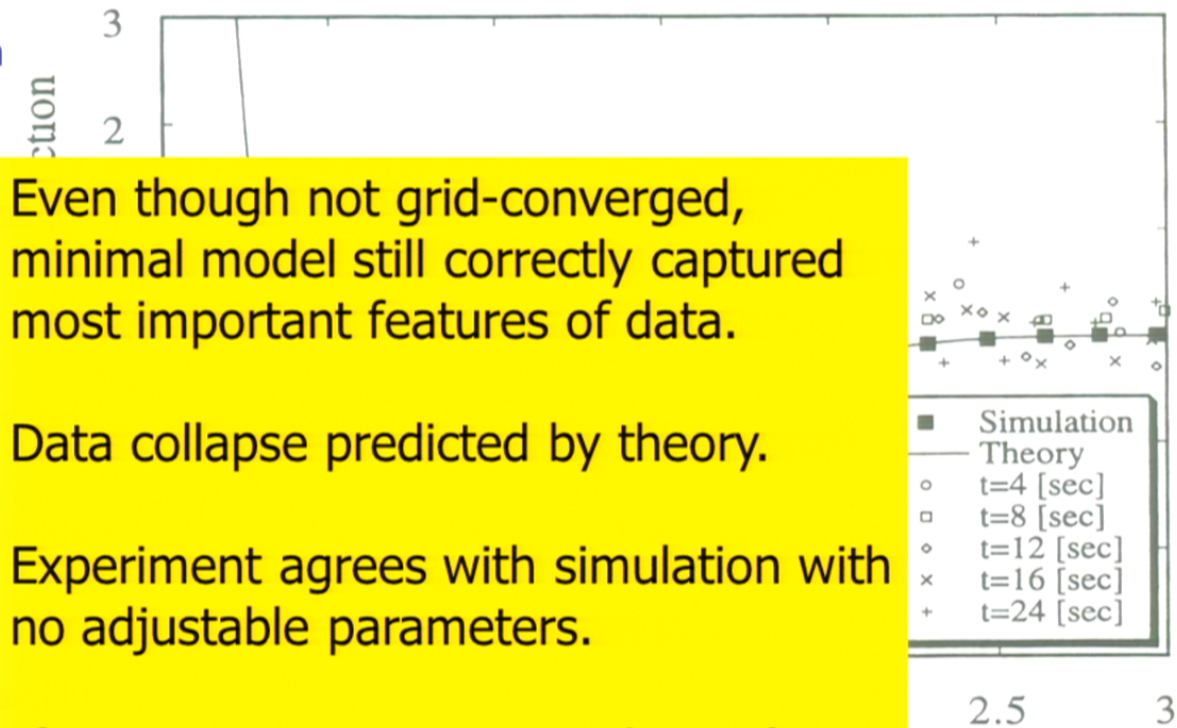
Raw data are correlation function at different positions at different time.

Data collapse predicted by theory.

Experiment agrees with simulation with no adjustable parameters.

Mean field theory by Liu and Mazenko gives good description except near defect core.

Mondello and NG, *Phys. Rev. A* **45**, 657 (1992).



Even though not grid-converged, minimal model still correctly captured most important features of data.

Data collapse predicted by theory.

Experiment agrees with simulation with no adjustable parameters.

If you want to compute scales in feet, you're out of luck: solution – measure one length.

defect densities. The solid line and the large closed squares indicate the theoretical curve by Lue *et al.*¹¹⁾ and the data of simulation by Mondello *et al.*,⁸⁾ respectively.

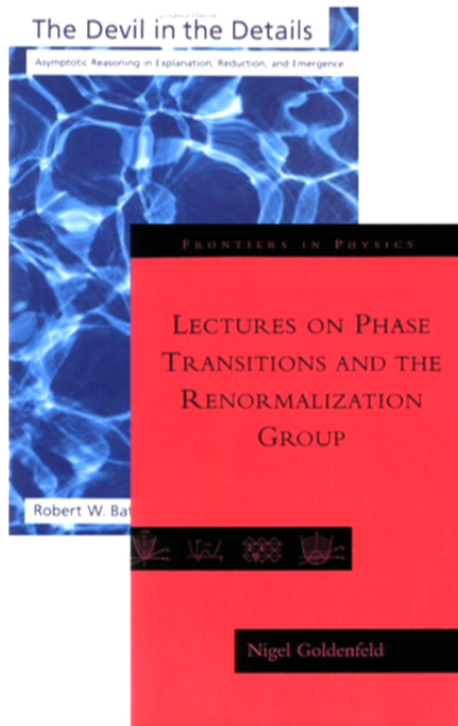
Summary

- **Predictions rely on existence of minimal models**
 - Parameter independent, universal quantities
- **Computational algorithms can model very simple processes that generate complicated space-time structures**
- **Quantitative prediction of universal quantities in excellent agreement with experiment**

Further reading

Tutorial on emergence and complexity

<http://media.nakfi.org/2008/complexity/goldenfeld/goldenfeld.htm>



Philosophy of science: R. Batterman has written about the use of asymptotic reasoning and RG as a scientific methodology:

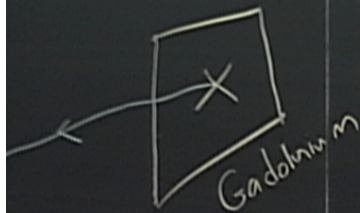
The Devil in the Details: asymptotic reasoning in explanation, reduction and emergence by Robert Batterman.

Renormalization group applied to modelling and the solution of partial differential equations:

Lectures on phase transitions and the RG by NG.

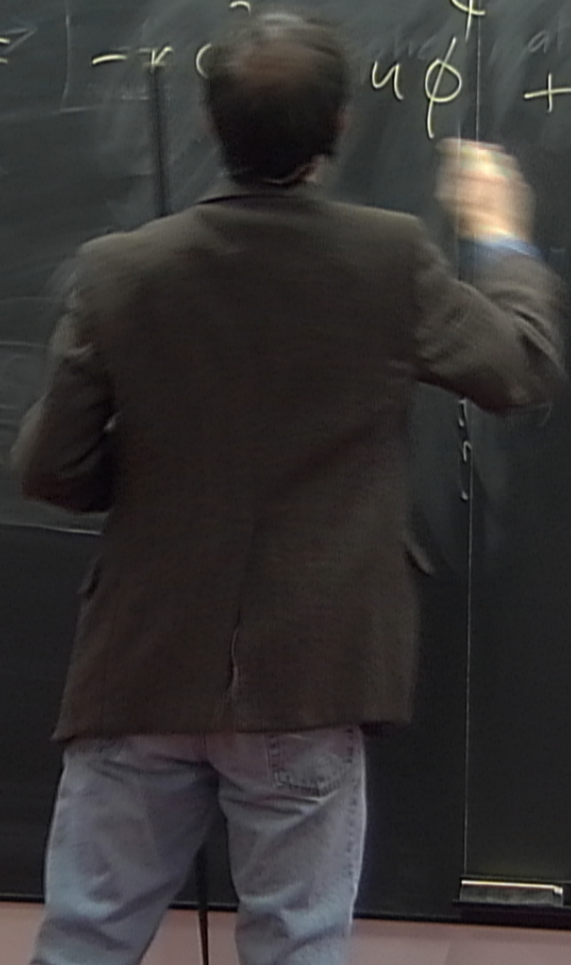
nickel
manifold of
phys Hams

manifold of
phys Hams



sher (1983)

$$F = \int \left[-\kappa(\nabla^2 \phi)^2 + u\phi + (\nabla \phi)^2 \right]$$



Manifold of phys Hams

Manifold of phys Hams

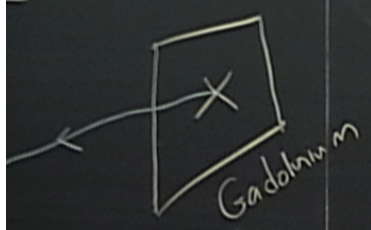


sher (198)

$$F = \int \left[-r \dot{\phi}^2 + u \phi^4 + (\nabla \phi)^2 \right] d^4x$$
$$= M \nabla \frac{\delta F}{\delta \phi}$$

nickel
manifold of
phys Hams

manifold of
phys Hams



sher (1983)

$$F = \int \left[-v \dot{\phi}^2 + u \phi^4 + (\nabla \phi)^2 \right] d^4x$$

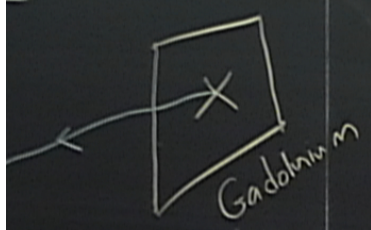
$$J = M \nabla \cdot \frac{\delta F}{\delta \phi}$$

$$\frac{\partial \phi}{\partial t} + \nabla \cdot \left(\frac{\delta F}{\delta \phi} \right) = 0$$

$$\frac{\partial \phi}{\partial t} = \nabla^2 \left(\frac{\delta F}{\delta \phi} \right)$$

nickel
manifold of
phys Hams

manifold of
phys Hams



sher (1983)

$$F = \int \left[-v \dot{\phi}^2 + u \phi^4 + (\nabla \phi)^2 \right] d^4x$$

$$J = M \nabla \cdot \frac{\delta F}{\delta \phi}$$

$$\frac{\partial \phi}{\partial t} + \nabla \cdot \left(\frac{\delta F}{\delta \phi} \right) = 0$$

$$\frac{\partial \phi}{\partial t} = \nabla^2 \left(\frac{\delta F}{\delta \phi} \right)$$

Manifold of phys Hams

Manifold of phys

she

$$F = \int \left[-v \dot{\phi}^2 + u \phi^4 + (\nabla \phi)^2 \right] d^4x$$

$$J = M \nabla \frac{\delta F}{\delta \phi}$$

$$\frac{\partial \phi}{\partial t} + \nabla \cdot J = 0$$

$$\frac{\partial \phi}{\partial t} = \nabla \cdot \left(\frac{\delta F}{\delta \phi} \right)$$